

SMR 1232 - 21

**XII WORKSHOP ON
STRONGLY CORRELATED ELECTRON SYSTEMS**

17 - 28 July 2000

RVB REDUX

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These are preliminary lecture notes, intended only for distribution to participants.

NAMES

RVB: The Originators

Z Zou, G Baskaran, T Hsu, PWA; I Affleck, G Kotliar, X-G Wen
R B Laughlin et al
KRS (Kivelson et al)

Inhabitants of the Big Tent--(spin gap = RVB)

PWA, Baskaran, Muthukumar
P A Lee, Wen, Nagaosa et al (Gauge Th.)
Fisher, Nayak, Senthil et al (Nodons)
Ting, Weng, Sheng et al (Schwinger Bosons)
Fukuyama et al
John, Berciu (instantons)

Fellow-travelers: (NFL, but.....)

Varma, Abrahams (MFL)
Rice, Laughlin, Lederer, etc (Flux phases)
Chakravarty, Sudbo, et al (ILT)

Many Experimentalists --too many to mention but especially:

N-P Ong, T Timusk, H Drew and many other persistent, careful transport measurers
J-C Campuzano, Mike Norman and the Argonne PES group
P Johnson, Z-X Shen and many other PES workers

RVB Redux

The Theory that got lost
Outline of Talk

I Historical (1957-88) RVB

II Successes, realized and unrealized

III Reasons for Abandoned
and (III a) why they were
wrong

IV New RVB

V What about SC?

I. Historical RVB of '87-
(Baskaran, Zou, Tsun., P)

Cuprates based on Mott insu-
lator with ^{Hubbard} gap: t - J model

from Hubbard Model

$$\sum_{i,j} t c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

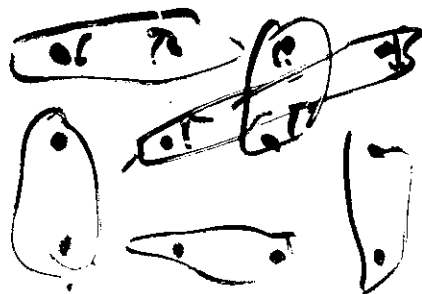


$P T P$

$$+ J \sum_{i,j} S_{i\uparrow} S_{j\uparrow}$$

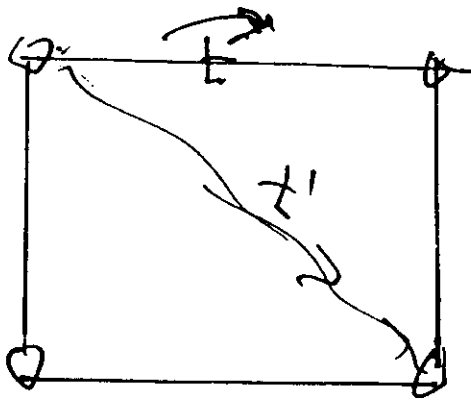
P : ~~X double occupancy X~~
perturbation in $1/U$.

RVB: a hypothetical spin
liquid of singlet pairs



+ etc.

THE HUBBARD MODEL—ONE BAND!



(TL) : U

$$H = t \sum_{\langle i,j \rangle} c_{i\sigma}^\dagger c_{j\sigma} + t' \sum_{\langle\langle i,j \rangle\rangle} c_{i\sigma}^\dagger c_{j\sigma} + U n_{i\uparrow} n_{i\downarrow} \quad U \gg 0 > t$$

+ Hopping between planes, sometimes.

Justifiable from chemical reasoning—all but the essentials abstracted out. But we have NO reason to think more is needed!

THE T-J TRANSFORM:

Antiferromagnetic exchange from a PROJECTIVE transformation. Take U as H_0 , t as a perturbation, and eliminate double occupancy:

$$n_{i\uparrow} n_{i\downarrow} \rightarrow 0$$

you get the "t-J" Hamiltonian

$$H = P \sum_{ij} t_{ij} c_{i\sigma}^\dagger c_{j\sigma} P + J \vec{S}_i \cdot \vec{S}_j$$

for large U , small x , $J \approx 4t^2/U$; it is a DERIVED quantity, always $< t$, U , somewhat x -dependent.

$$P = \text{projector} ; \vec{S}_i = \sum_{\sigma\sigma'} a_{i\sigma}^\dagger (\vec{\sigma})_{\sigma\sigma'} a_{i\sigma'}$$

You can't have AF exchange without the projection: FLT and AF J are logically incompatible. In FLT exchange is ferromagnetic. (SDW's are caused by nesting, not by superexchange)

A possible ground state of
Heisenberg Model

To construct : Bose Condensate
of singlet pairs $\equiv$ BCS
wave function, project on

$(\hat{n}_3) \equiv 0$ so $\hat{n}_3 (= 1 - \hat{n}_{\uparrow} - \hat{n}_{\downarrow})$
drops out.

All self-energy anom
alous

$$\begin{pmatrix} \psi \\ \psi^\dagger \end{pmatrix} \begin{pmatrix} 0 & \Delta_r + i\Delta_i \\ \Delta_r - i\Delta_i & 0 \end{pmatrix} \begin{pmatrix} \psi \\ \psi^\dagger \end{pmatrix}$$

Eigen fns

$\psi \pm \psi^\dagger$

"spinors" =

Real Fermions

eigen excitations = Majorana Fs

(projection means only created in
pairs!)

ENORMOUS SYMMETRY : Local
(Kohn; BASKARAN) $U(1)$

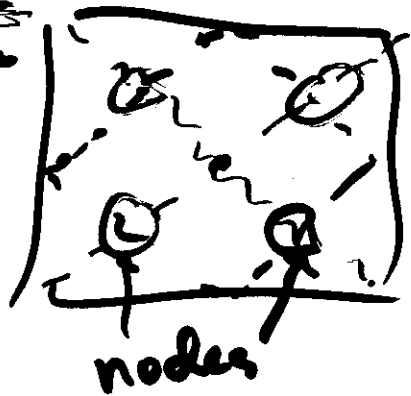
(or $Z_2: \pm 1$ - G.B. Yule, G. Rossi?)

(doesn't care about sign of w.f.)

eg $\cosh k_1 + \cosh k_2 \equiv \cosh k_1 - \cosh k_2$

can use both: "s + id" = $\Delta_1 + i\Delta_2$

is best solution ~~the~~
variationally.



Barkeran: define an $F(\Delta)$ with $U(1)$ symmetry

$$F = F_1(1/\Delta^2)$$

$$F_1 \rightarrow "s"$$

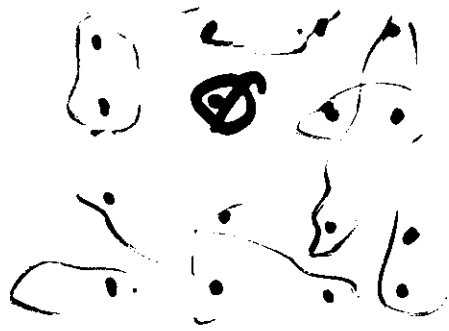
+ $F_2 \sum$ plaquet terms

$$\Delta_{12} \Delta_{23} \Delta_{34} \Delta_{41}$$

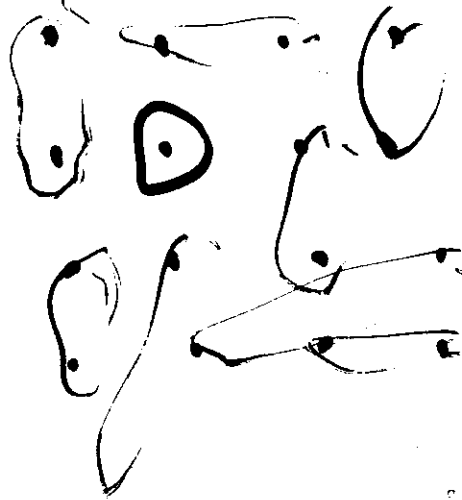
Plaquet term $\rightarrow "d" = \text{resonance}$

(Diff of energy $\rightarrow$ of symmetric vs anti: $||\pm\rangle$)

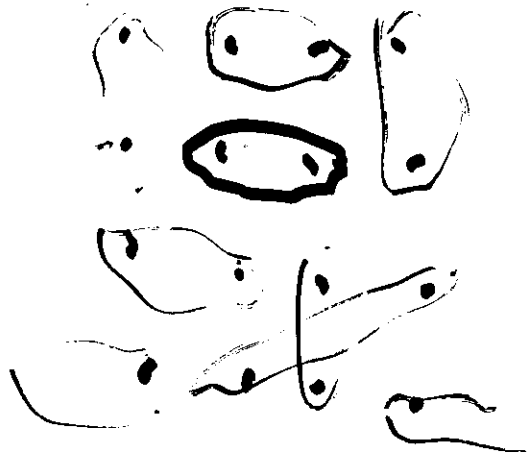
spinon



holon



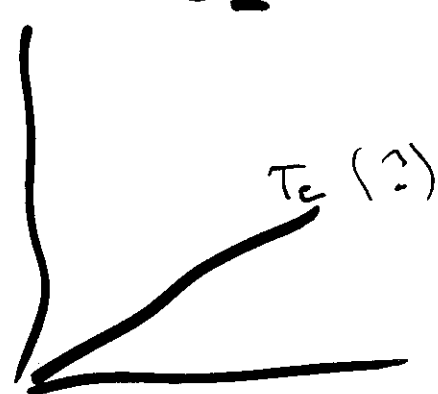
sc



For $1/2$ full, symmetry
says $F_1 \approx F_2 = J$

Now DOPE T_c No T_c at $x=0$
phase stiffness increases
with x

but Δ forms
at $T \sim J$



NO PHASE TRANS. x

II QUALITATIVE SUCCESSES OF RVB.

(1) ABOVE ALL: SOLUTION
OF THE DEEP, QUANTITATIVE PROBLEM:
PAIRING WITHOUT PHONONS,
PURELY FROM REPULSIVE
ELECTRON-ELECTRON INTERACTION

- (2) PREDICTED : PSEUDOGAP !
 - (3) " : NO PHASE TRANSITION TO PSEUDOGAP
 - (4) " : d-wave gap ! ^{Kotliar} _{Liu}
 - (5) "protectorate": insensitive to phonons, impurities (magnetic)
 - (6) Dep of T_c on χ
-

III So why abandoned ?

- (1) Heisenberg Model not RVB (though close to Hsu)
- (2) Insensitivity of T_c suggested s-wave
- (3) Optimally doped ^(?!) has Fermi surface $\therefore$ NFL theory
- (4) $T_c = f(\text{layer } \#)$: ILT
- (5) $T^*, E_g^* = f(x)$
- (6) Main : no formal th of T_c .

Radun : solve these problems

- (1) Irrelevant: AF $\rightarrow$ RVB with doping (critical pt.). Soft mode = neutron resonance!
- (2) d-wave! insensitivity $\rightarrow$ quantum protectorate
- (3) combine RVB, NFL!
see below
- (4) Experimental red Herring?
partly true
- (5) see (3)
- (6) Still a problem: but energy scales are clear.

NEW RVB

(1) Charge - spin separation
is generic in 2D.

Expt - ARPES EDC's
Transport theory; impurities

Theory - chiral anomaly,
infrared catastrophe

Symmetry: $U(2) \rightarrow SU(2) \times U(1) \times Z_2^{\text{hole-particle}}$

(2) Therefore: spin sector

~~still~~ Lagrangian retains
 $U(1)$ invariance: Elitzur's

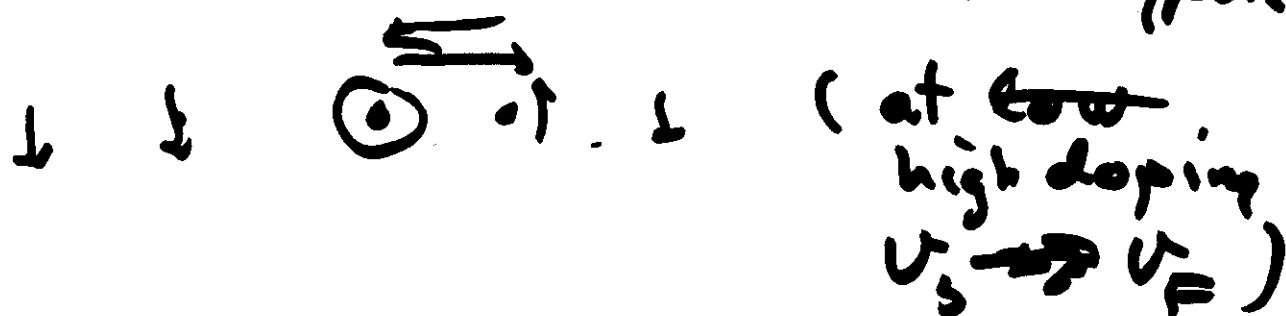
theorem allows no phase
transition at T^*, E_g^*

(3) Effective energy in spin
sector = $F(\Delta_b^2) + K \square$

F increases with ^{d-wave} term
For $J + xt$

K decreases with doping

+ sign. Weakening of Umklapp
 v_s increases because of ~~non~~
 repeating paths - Brattner-Rice effect.

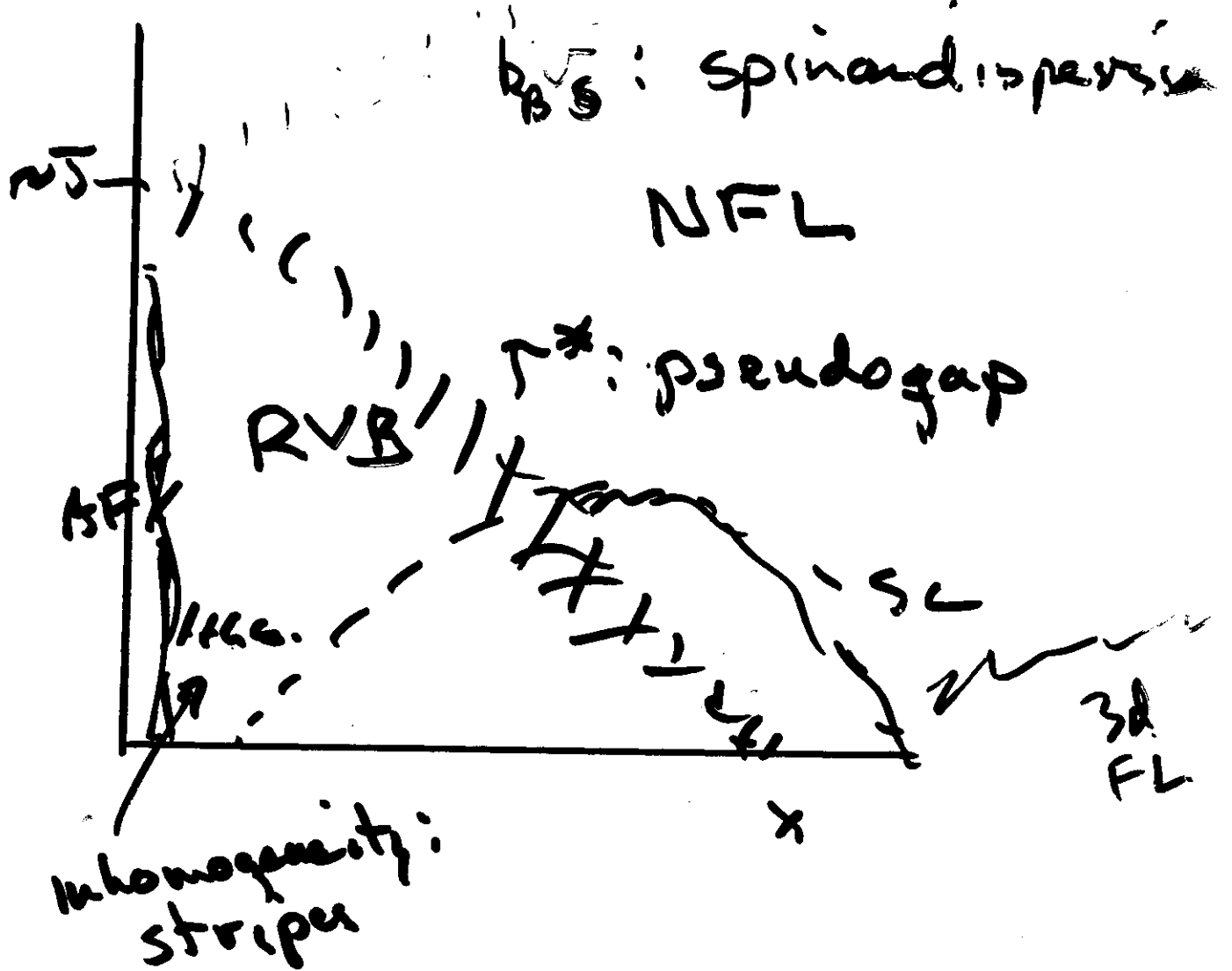


(How do we define v_s ?
 make an electron excitation
 with charge in ground state,
 this has defined momentum,
 satisfies Luttinger thm.)

So we have

- (1) pseudo gap, no transition
- (2) neutron res in RVB state only

Phase Diagram



Why does T_{eff} decrease?

Resonance energy decreases because loop paths of holes have ferromagnetic sign. Hole makes $\square$ with

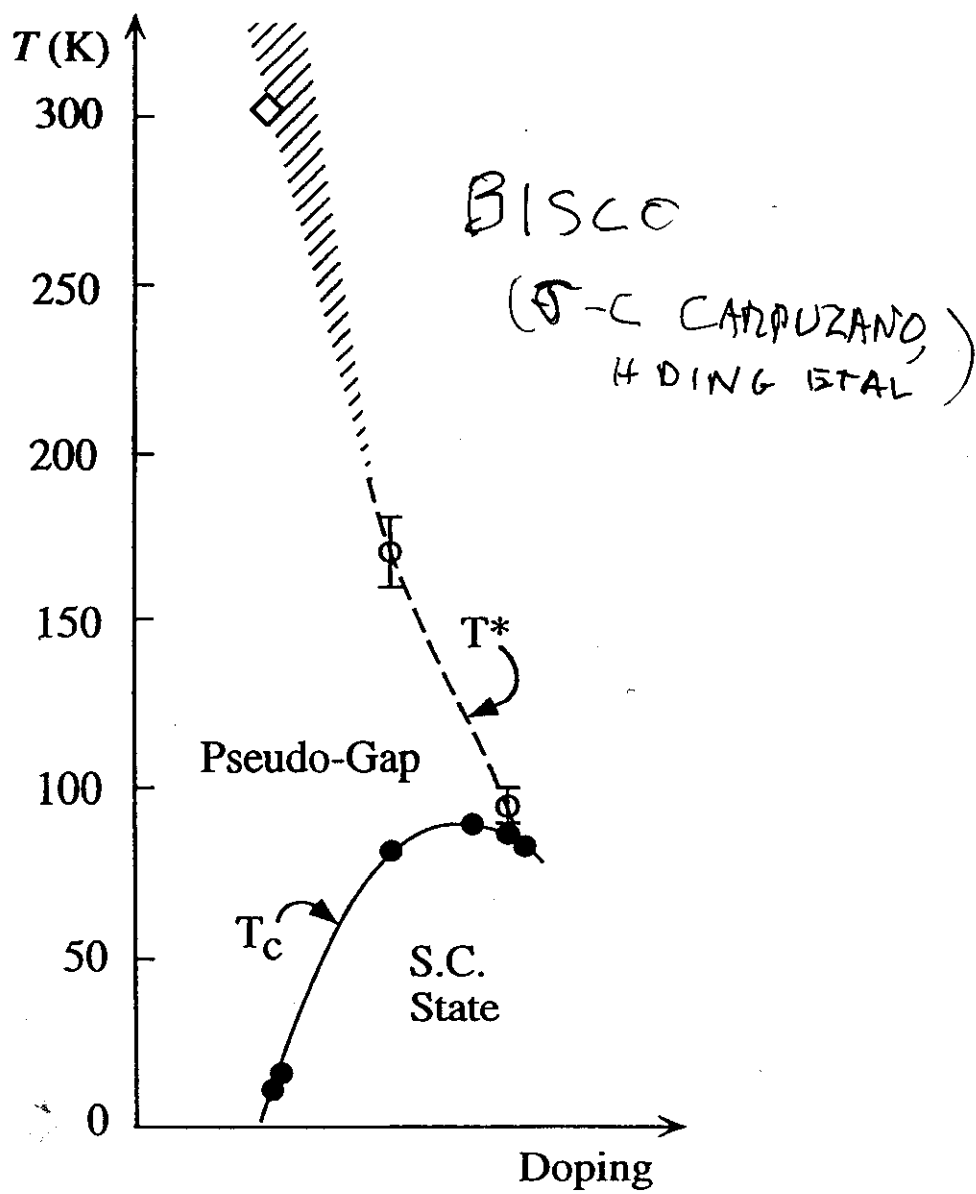


Fig. 4

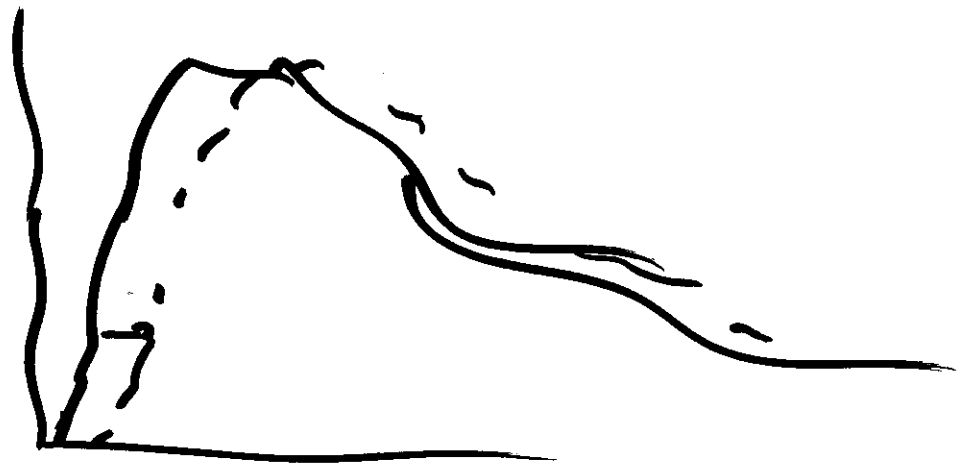
2. $T^*(x)$

Why + How Superconductivity?

Why is easy! ~~AE~~

Spin gap causes big increase in kinetic energy.

ARPES



IR: BIG INCR ω IN M^* (see Fig)

THOMAS SECRET: Kohn sum

RULE SAYS K.E.

KE RE GAINED BY COHERENT

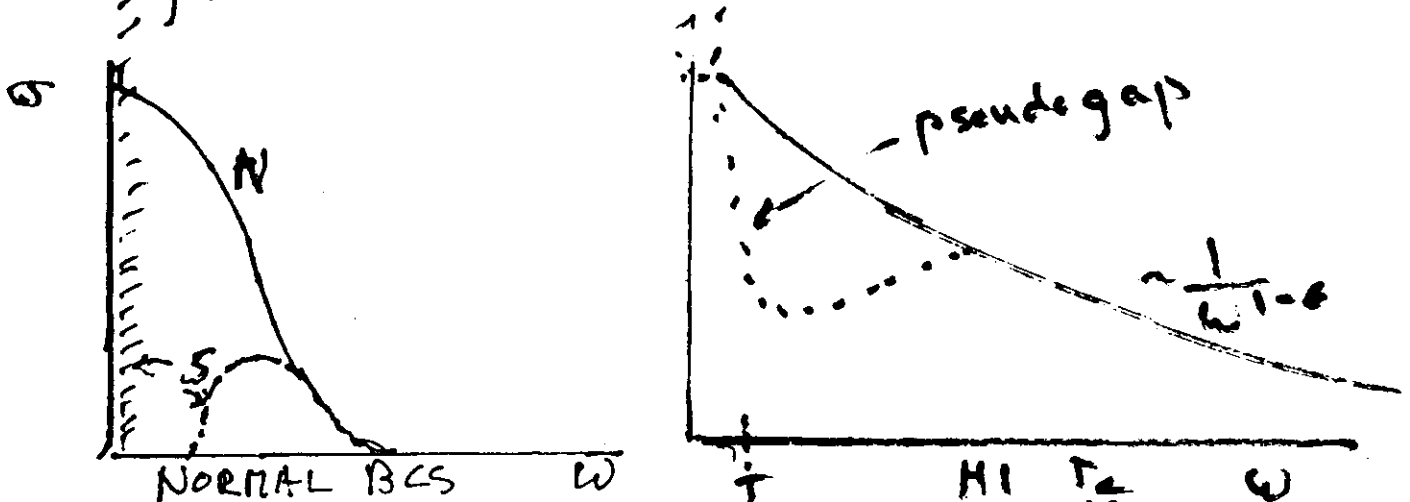
CHARGE FLUCTUATIONS AS IN

17 FOMELL - FINKHAM IDENTITY

SUPERCONDUCTIVITY AT LAST

The role of the (Kohn)-Ferrell-Tinkham sum rule

$$\int \sigma d\omega = e^2 \times \text{Total K.E.}$$



In BCS, $K.E. \gg$ condensation energy; therefore no gap without SC--compensation of coherent pair hopping vs coherent single-particle motion is nearly exact.

But in the cuprate the conductivity is $oc x$, and much more widely distributed in frequency; so the gap in the conductivity is less severe and roughly $oc x$. This is why the spin gap can open for small x , ie why phase II exists.

But eventually, the kinetic energy wins and charge must condense into hole pairs: hence T_c . The K E can be both interlayer and intralayer in multilayers.

On the overdoped side, T^* and T_c coincide-- T_c is TWO transitions. Hole condensation is s-wave, hence not much affected by scattering? But may have isotope shift.

$$\sigma \approx \frac{ne^2}{i m^*(\omega)\omega + \frac{1}{\tau(\omega)}}$$

m^* Big in gap

The pseudogap in high-temperature superconductors

Timush's little secret. 93

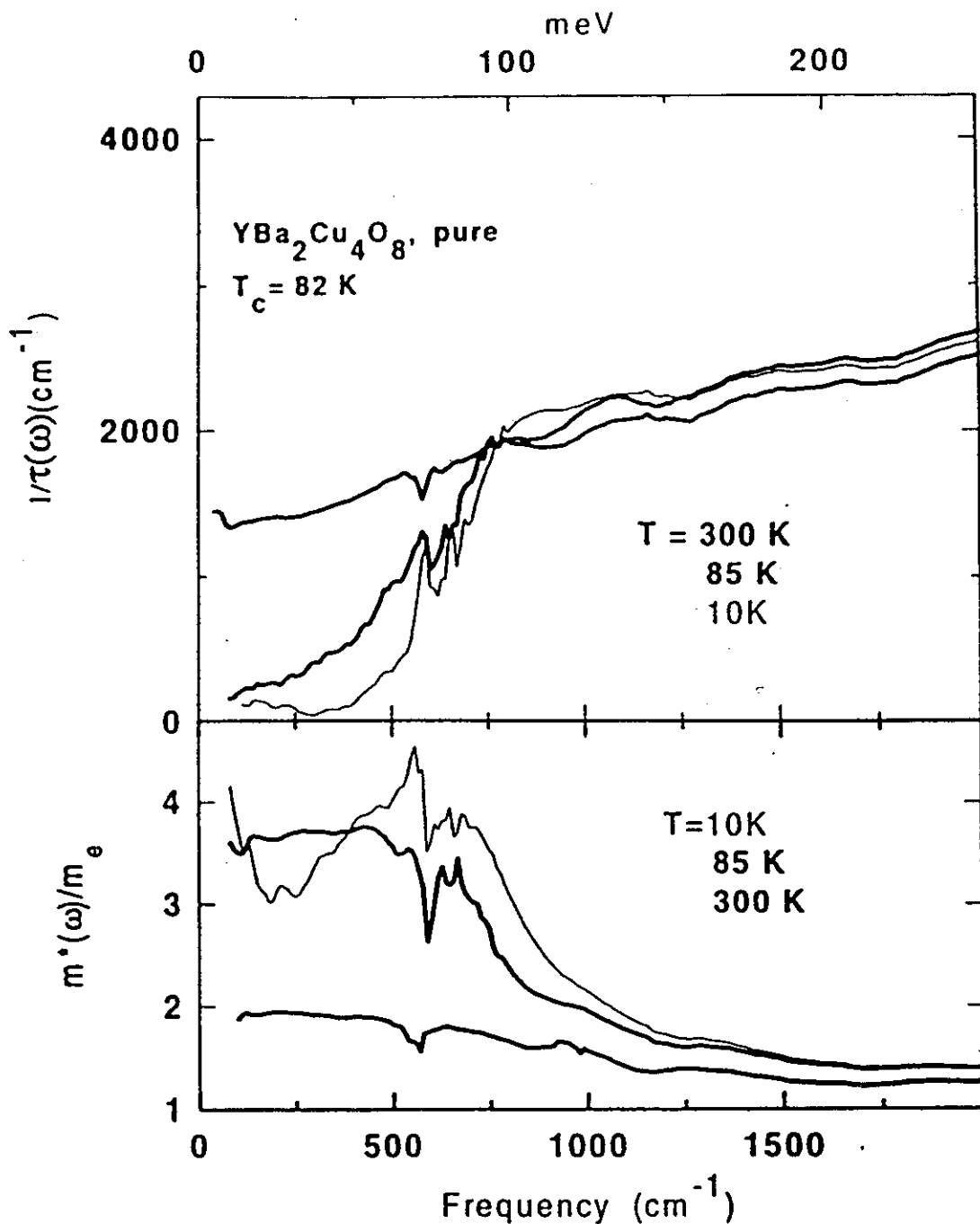
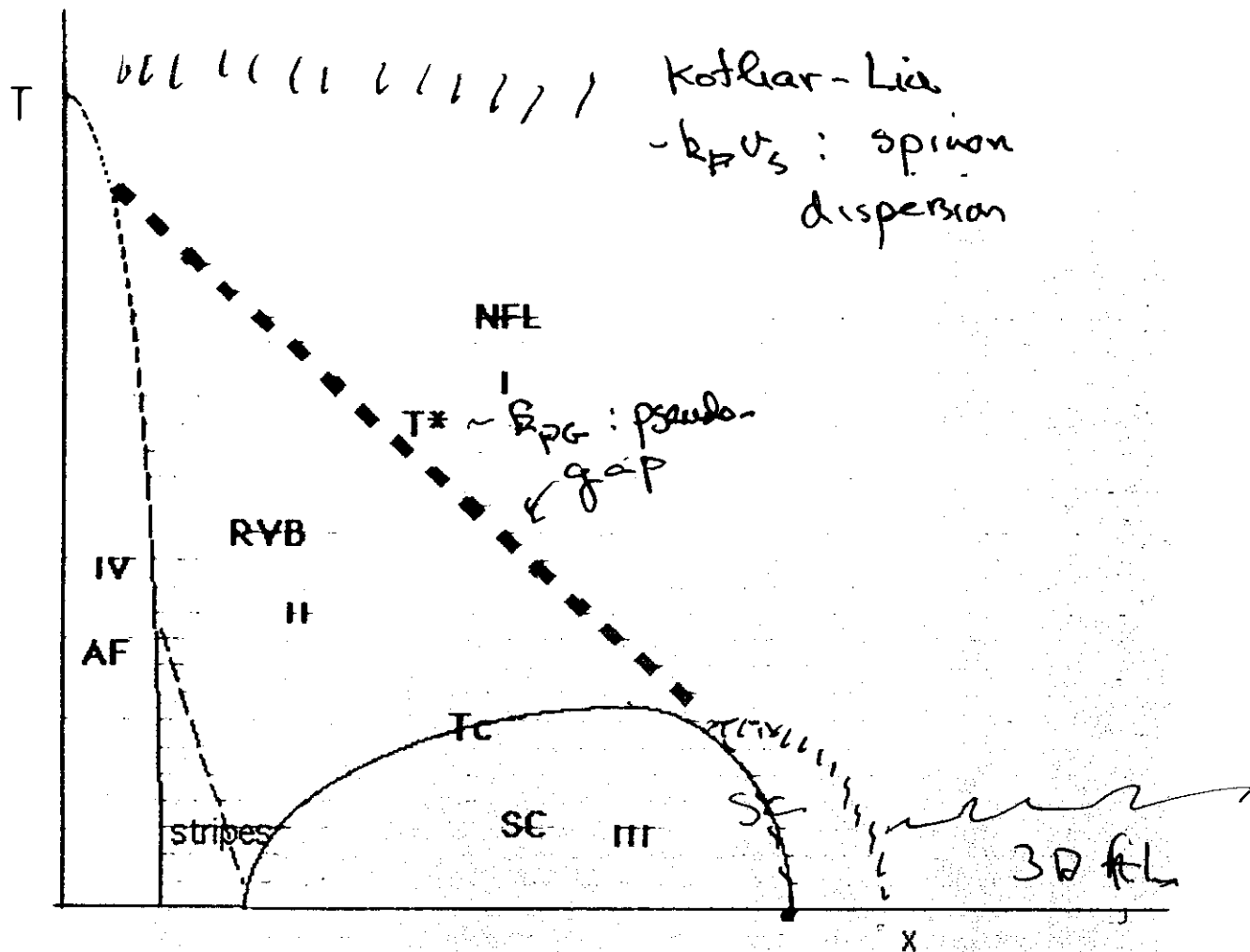


Figure 28. The frequency-dependent scattering rate and the effective mass of YBa₂Cu₄O₈. The scattering rate varies linearly at room temperature but develops a gap like depression in the normal state. At the same time the effective mass of the carriers develops a resonance peak at $\approx 600 \text{ cm}^{-1}$ (75 meV).



GENERALISED PHASE DIAGRAM

spin prefers to bind to
pairs of holes

CONCLUSIONS

Cuprates have unique properties:

- single, narrow band, without self-trapping
- proximity of an $S=1/2$ Mott Antiferromagnet
- 2D= non-Fermi liquid
- allows AF- $\rightarrow$ RVB with doping

THE CRUCIAL TRANSITION IS THE SPINON PAIRING AT T^*

Superconductivity is then inevitable.

We were so close in 1987!!