

SMR 1232 - 36

**XII WORKSHOP ON
STRONGLY CORRELATED ELECTRON SYSTEMS**

17 - 28 July 2000

GLASSES, OPTIMIZATION AND COMPUTATION

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These are preliminary lecture notes, intended only for distribution to participants.

GLASSES, OPTIMIZATION AND COMPUTATION

L. B. IOFFE

PC

I. MOTIVATION

II. A MODEL

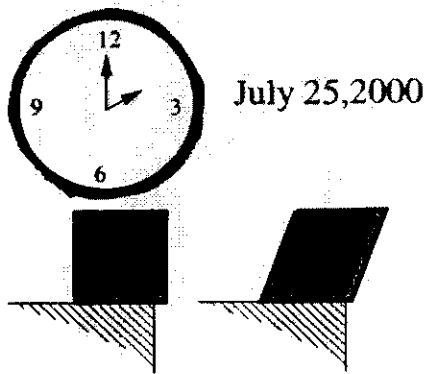
III. NEW APPROACHES TO OPTIMIZATION

IV. COMPUTATION?

V. SUMMARY / ???s

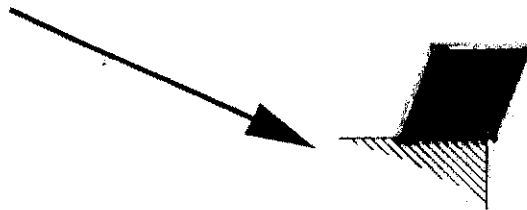
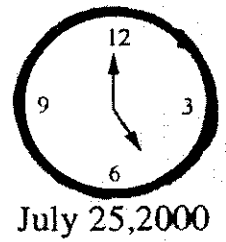


RESPONSE TO SHEAR

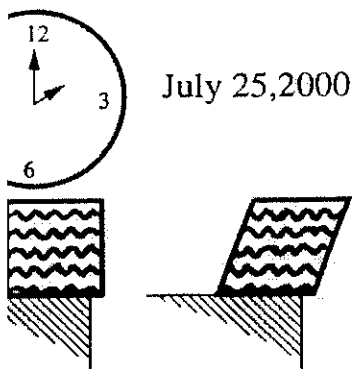


IDEAL

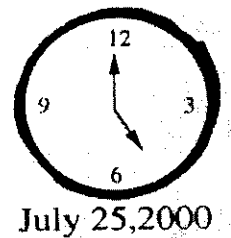
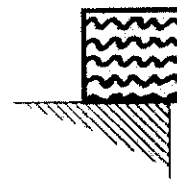
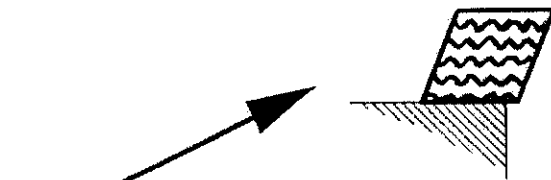
ELASTIC SOLID



July 25, 3000



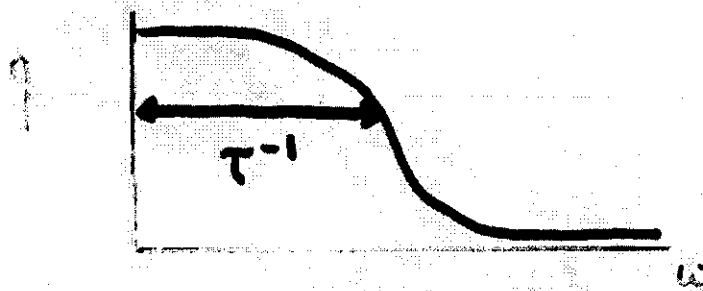
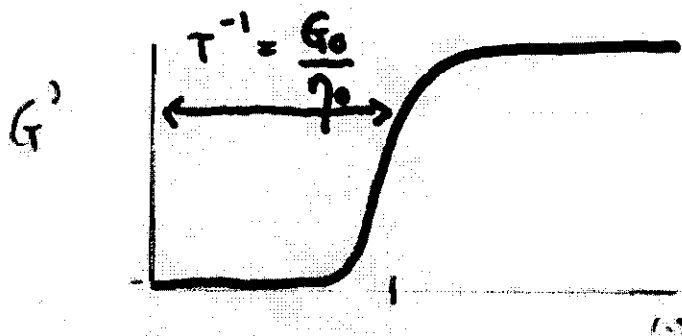
GLASS



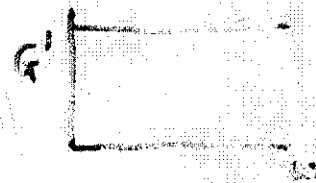
July 25, 3000

	LIQUID $\rightarrow$ SOLID	METAL $\rightarrow$ SUPERCONDUCTOR
CONDUCTIVITY EQUATION	$\Pi = \eta \dot{u}$	$J = \sigma E = -QA$
"DRUDE"	$\eta = \eta_0 \frac{1}{1-i\omega\tau}$	$\sigma = \frac{ne^2\tau}{m} \frac{1}{1-i\omega\tau}$
LONDON KERNEL	$\Pi = G u$	$J = -QA$
KUBO RELATION	$G = -i\omega\eta$	$Q = -i\omega\sigma$

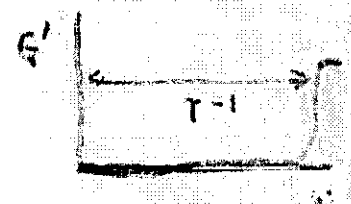
$$G = \frac{G_0}{\eta_0} \frac{-i\omega\tau}{1-i\omega\tau}$$



IDEAL
SOLID



LIQUID



WANTED : A REDUCTIONIST MODEL

(e.g. ISING MODEL)

FOR GLASSES

- HOW IS A SYSTEM SIMULTANEOUSLY
OUT-OF-EQUILIBRIUM AND
STABLE ON VERY LONG
(e.g. MILLENNIA)
TIME-SCALES ?

- HOW DOES ONE KEEP AN INTRINSICALLY
NON-RANDOM SYSTEM AWAY FROM
ITS THERMODYNAMIC GROUND-STATE

WEISS MODELS OF DISORDER-FREE GLASSINESS

MAIN DIFFICULTY : INEQUIVALENCE OF SITE
(IN A GIVEN STATE M_i VARIES FROM SITE TO SITE)

N.B. GLASSES WITH
QUENCHED DISORDER

⇒

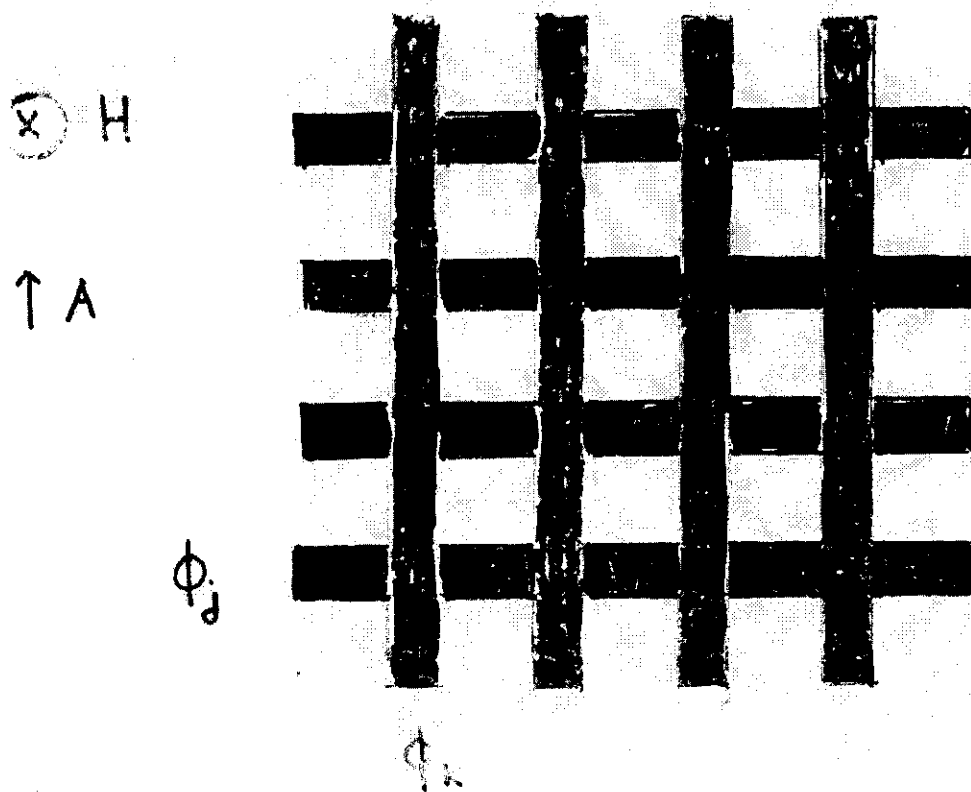
EQUIVALENCE OF SITE
AFTER DISORDER
AVERAGING

CURRENT APPROACHES

- FIDUCIARY HAMILTONIAN
- REPLICA / CLONING
- DYNAMICS

REVIEW: BOUCHAUD, CUGLIANDOLO, KURCHAN + MEZARD, IN YOUNG
SPIN GLASSES AND RANDOM BELDS (WORLD SCIENTIFIC,
SINGAPORE, 1999), pp. 101 - 225.

A MODEL

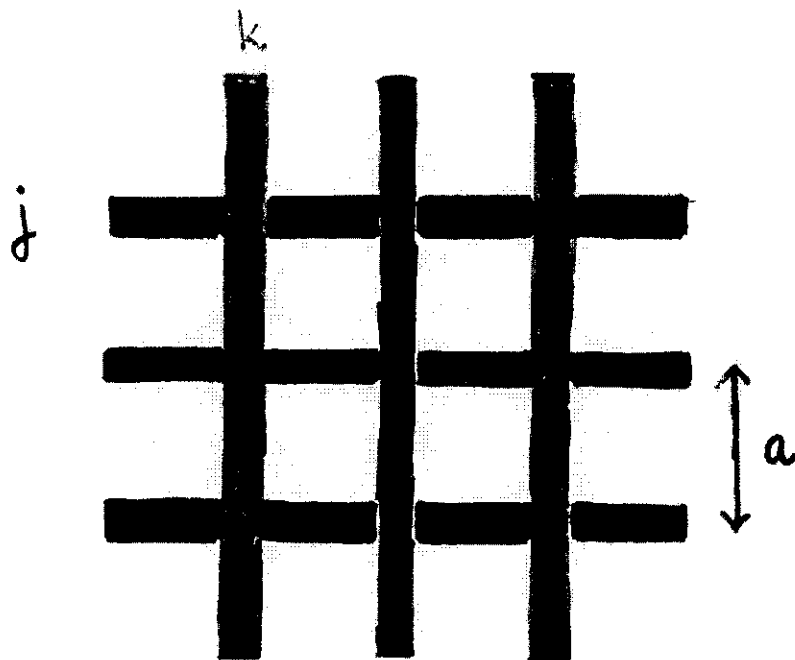


$$Z = N$$

$$E = \int_V (j_s)^2 dV - J_0 \sum_{jk} \cos \Delta \phi_{jk}$$

$$j_s \approx \nabla \phi - \frac{2\pi}{\phi_0} \vec{A}$$

(X) H



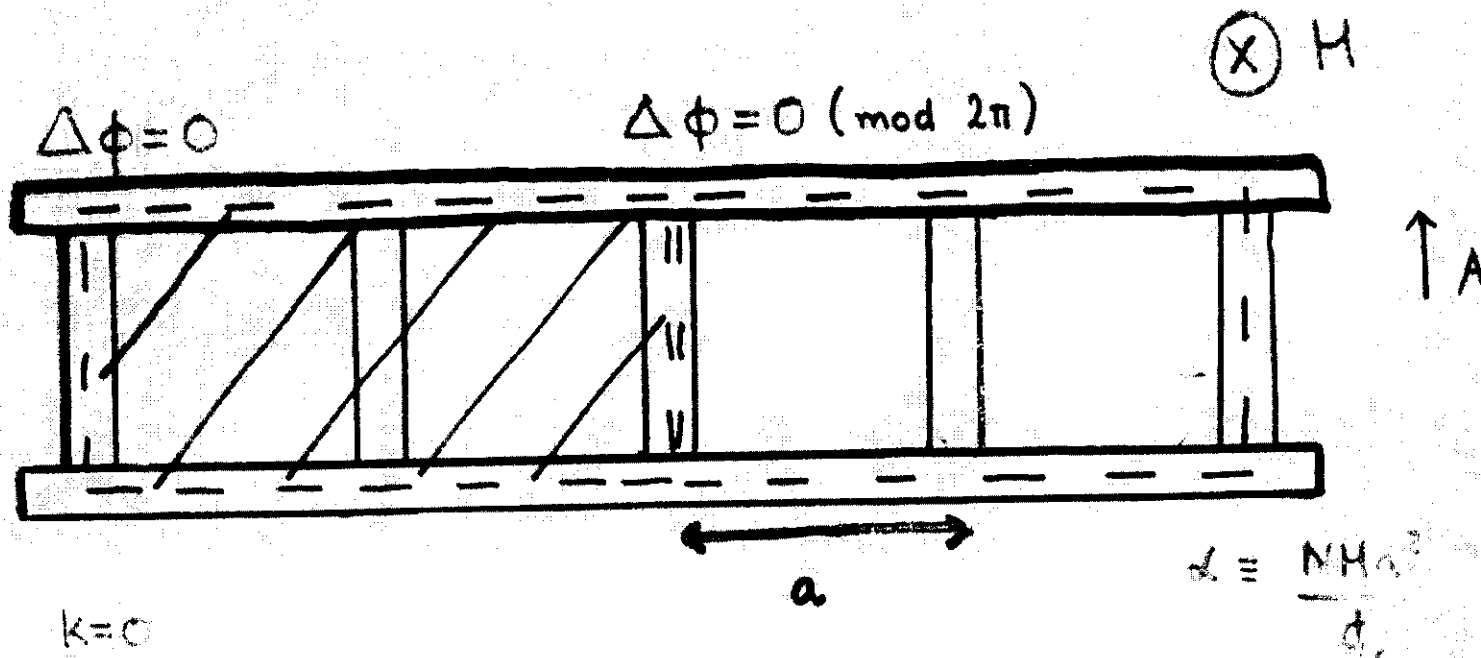
$$d = \frac{NHa^2}{\Phi_0}$$

$$\mathcal{H} = -J \sum_{jk} \cos \left(\phi_j - \phi_k - \frac{2\pi \Phi_{jk}}{\Phi_0} \right)$$

$$= -\text{Re} \sum_{jk} s_j^* J_{jk} s_k$$

$$s_j = e^{i\phi_j}$$

$$J_{jk} = J \exp \frac{2\pi i \Phi_{jk}}{\Phi_0} = J \exp \frac{2\pi i d_{jk}}{N}$$

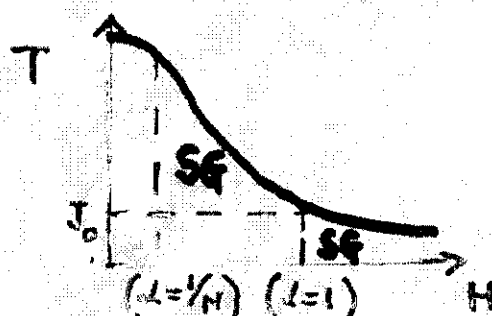


$$\sum_k^N J_{jk} J_{k\ell}^+ = \sum_k \exp \frac{2\pi i \lambda k (j-\ell)}{N}$$

COMMENSURATE IFF
INTEGER

DISORDERED CASE

(VINOKUR, IOFFE, LARKIN,
FEIGELMAN 1987)



PERIODIC CASE

$$\frac{1}{N} < \lambda < 1$$

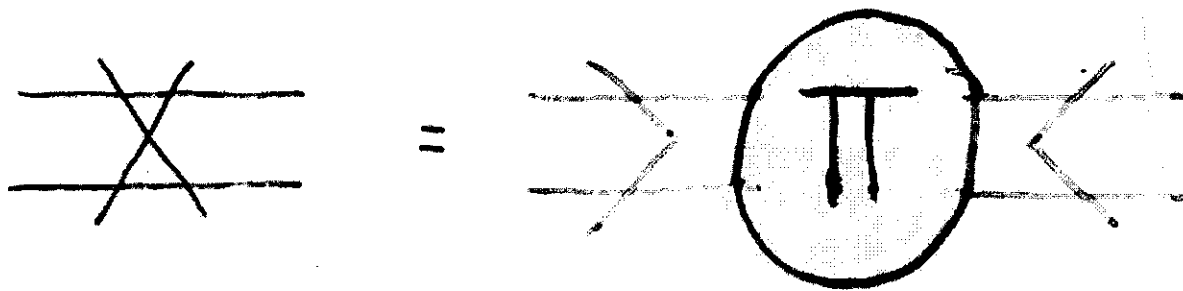
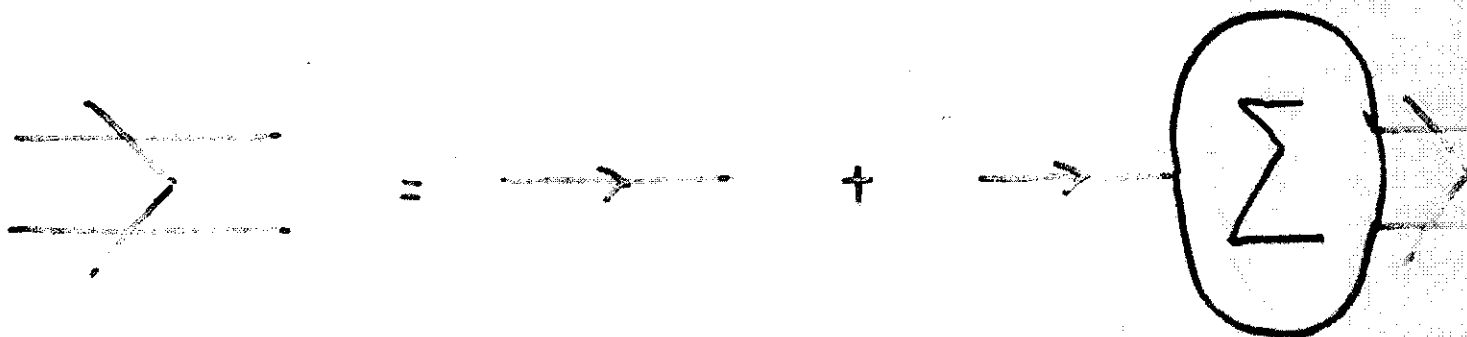
ABSENCE OF COMMENSURABILITY !!

DYNAMICS

(SOFT SPINS W/ LANGEVIN)

$$D(t, t') = \cancel{\text{diagram}} = \frac{1}{2N} \sum_{j=1}^{2N} \langle s_j(t) s_j(t') \rangle$$

$$G(t, t') = \Rightarrow \text{diagram} = \frac{1}{2N} \sum_{j=1}^{2N} \frac{\partial \langle s_j(t) \rangle}{\partial h_j(t')} \bigg|_{h_j=0, t > t'}$$



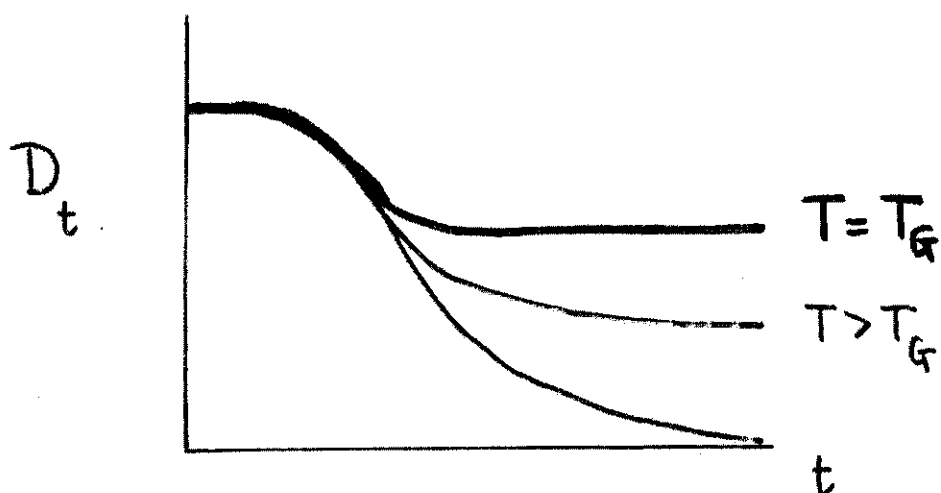
$$\Sigma \propto \text{diagram} ; \quad \Pi \propto \text{diagram}$$

$T > T_g$

$$C(t-t') = - \frac{\partial D}{\partial t}(t-t') \phi(t-t')$$

$\Downarrow$

$$\ddot{D}_t + \gamma \dot{D}_t + \Omega^2 D_t + \int_0^t \Pi(t-t') \dot{D}_{t'} dt' = 0 \quad (*)$$



(*) SIMILAR TO MODE - COUPLING EQ.

(DENSITY - DENSITY CORRELATIONS)

(*) IDENTICAL TO DYNAMICAL EQUATION

FOR THE $P=4$ (DISORDERED) SPHERICAL
MODEL ("SIMPLEST SPIN GLASS")

Transition Temperature of Josephson Junction Arrays with Long-Range Interaction

H. R. Shea and M. Tinkham

Department of Physics and Division of Engineering and Applied Sciences, Harvard University, Cambridge, Massachusetts 02
(Received 17 June 1997)

We report measurements of the dependence on magnetic field and array size of the resistive transition of Josephson junction arrays with long-range interaction. Because every wire in these arrays has a large number of nearest neighbors (9 or 18 in our case), a mean-field theory should provide an excellent description of this system. Our data agree well with this mean-field calculation, which predicts that T_c (the temperature below which the array exhibits macroscopic phase coherence) shows very strong commensurability effects and scales with array size. [S0031-9007(97)04071-4]

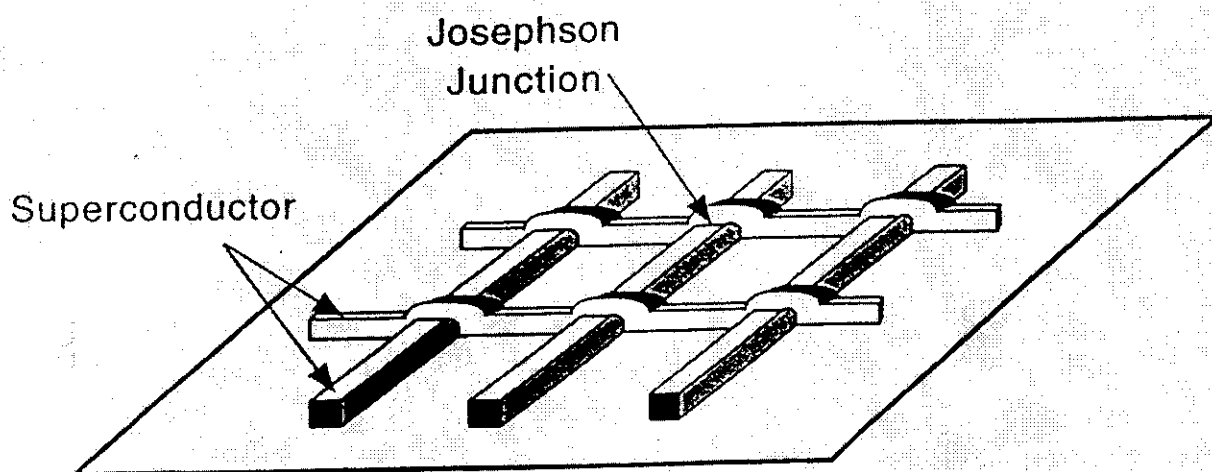
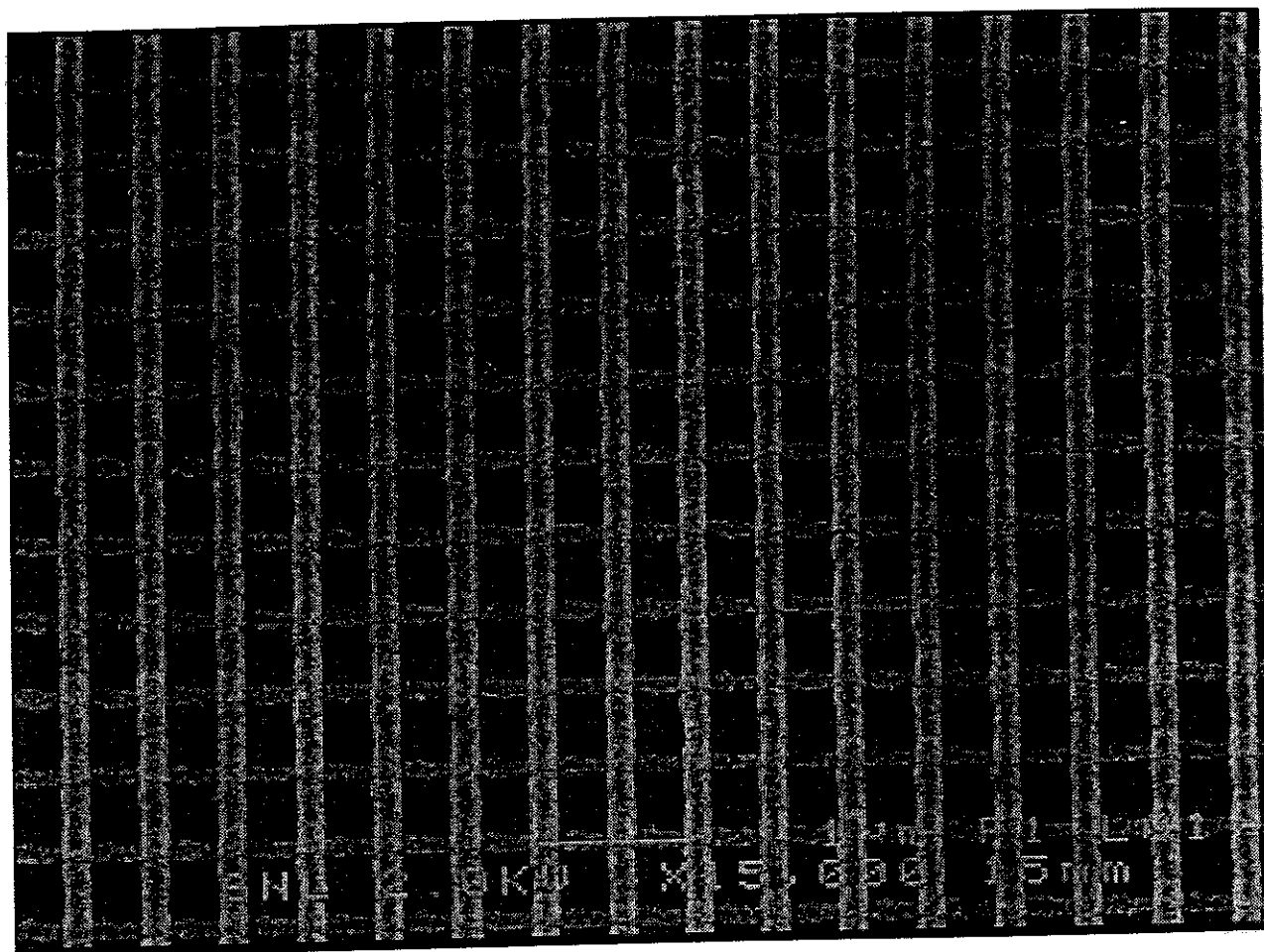


FIGURE 1

H.R. Shea and M. Tinkham
"Transition Temperature of Josephson Junction
Arrays with Long-Range Interaction"



(M. GERSHENSON ET AL.)

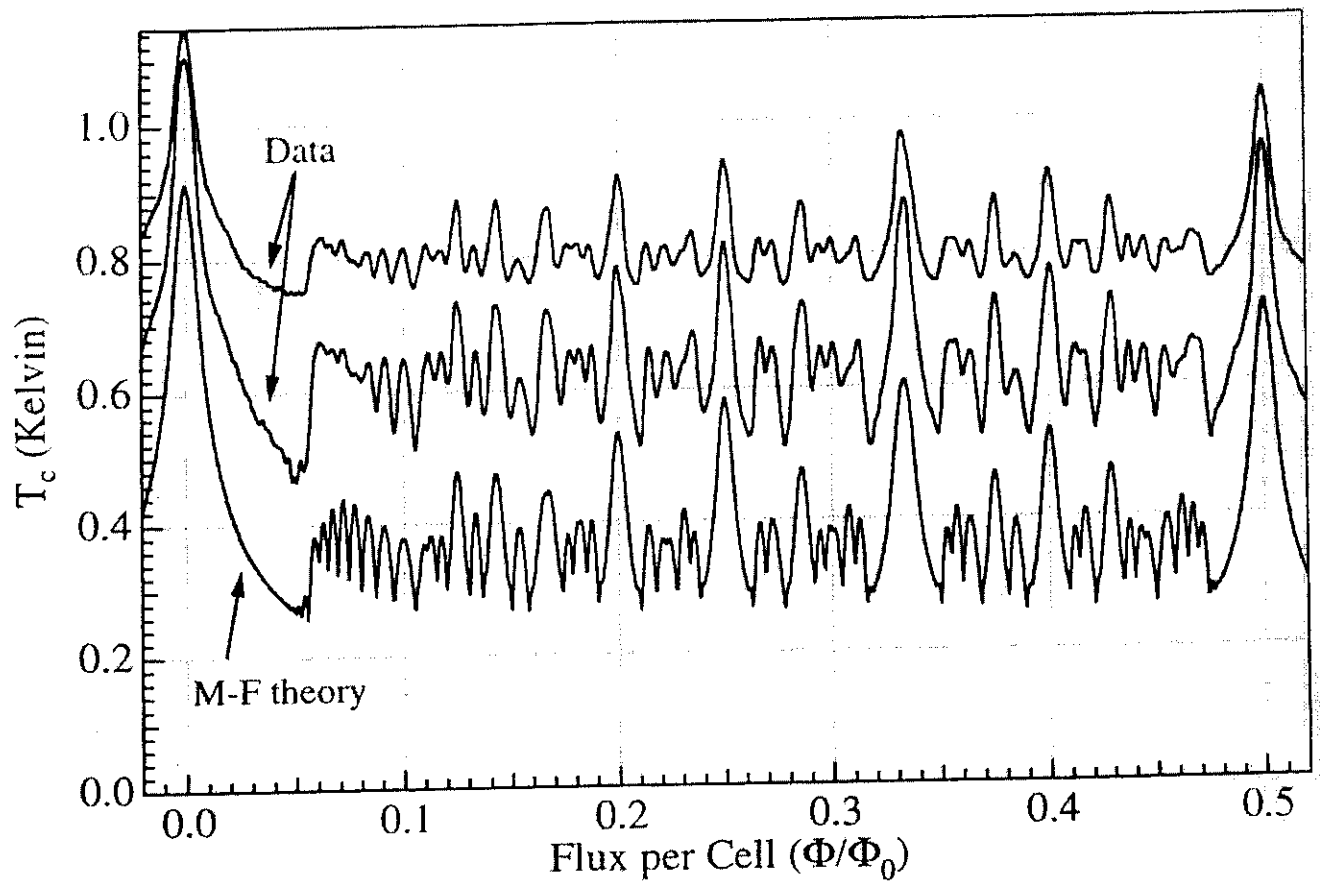


Figure 4

H.R. Shea and M. Tinkham
 "Transition Temperature of Josephson Junction Arrays with Long-Range Interaction"

UNITARY CASE ($\alpha=1$)

$$H = \sum_{i,j} h_{ij} a_i^\dagger a_j$$

$a_i^\dagger, a_i$ are fermionic creation and annihilation operators

$$[a_i, a_j] = 0, [a_i, a_j^\dagger] = \delta_{ij}$$
$$a_i^2 = 0, (a_i^\dagger)^2 = 0$$

$\psi_i =$ SPINOR FIELD

GROUND STATE

$$E_0 = - \sum_j |h_j|$$

$$\sum_k h_k^2 = N \Rightarrow \text{MINIMAL } E_0 \Leftrightarrow |h_k| = 1$$

GROUND
STATES

$\Leftrightarrow$

UNITARY SEQUENCE
WITH FLAT FOURIER
TRANSFORM

$$B_r \equiv \sum_l s_{l+r}^* s_l = 0 \quad \text{ZERO AUTO CORRELATION}$$

$$\text{SINCE } \sum_k e^{2\pi i \frac{kr}{N}} |h_k|^2 = B_r = \delta_{0r}$$

FAMOUS PSEUDORANDOM SEQUENCES

GAUSS

$$s_{jk} = e^{2\pi i \frac{k^2}{N}}$$

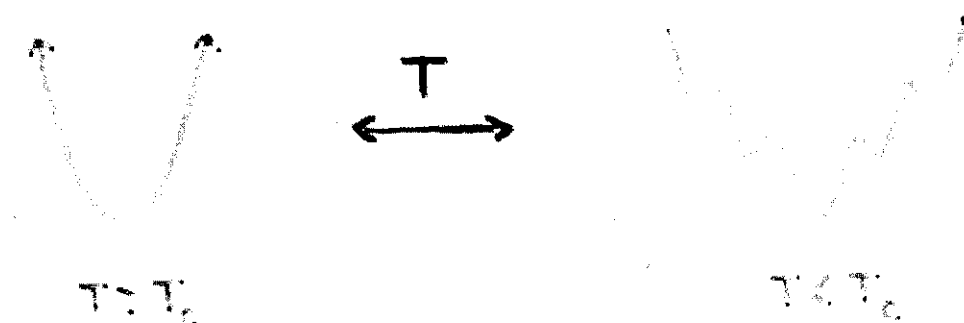
GALOIS

$$s_{jk} = (-1)^{a_k} \quad k=1 \dots N$$

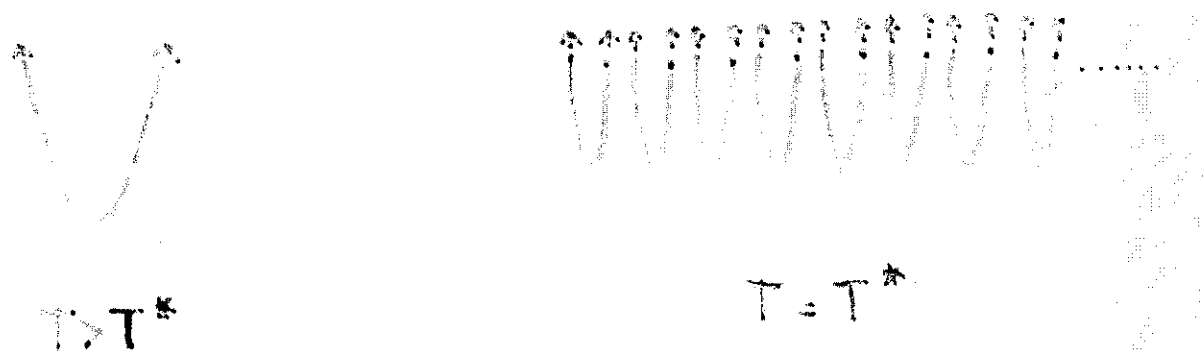
OPTIMIZATION

OPTIMAL STATE $\longleftrightarrow$ STATE w/
LARGE BASIN
OF ATTRACTION

E.G. SIMULATED ANNEALING



HOWEVER....



TUNING PARAMETER ?

FAMILY OF MODELS

$$\sqrt{x} \sum_{i_1 i_2} J_{i_1 i_2} s_{i_1} s_{i_2} + \sqrt{1-x} \sum_{i_1 i_2 i_3} J_{i_1 i_2 i_3} s_{i_1} s_{i_2} s_{i_3}$$

$$\left(\sum_{i=1}^N s_i^2 = 1 \right)$$

$p > 2$

- ∞ RANGE

- ALL METASTABLE STATES APPEAR AT T_c AND ARE ORTHOGONAL

(CRISANTI + SOMME
KURCHAN, PARISI
+ VIRASORO)

- HISTORY-DEPENDANCE AND AGING

(CUGLIANDOLO +
KURCHAN)

$x = 1$

($p=2$) ONE SOLUTION (KOSTERLITZ)

$x = 0$

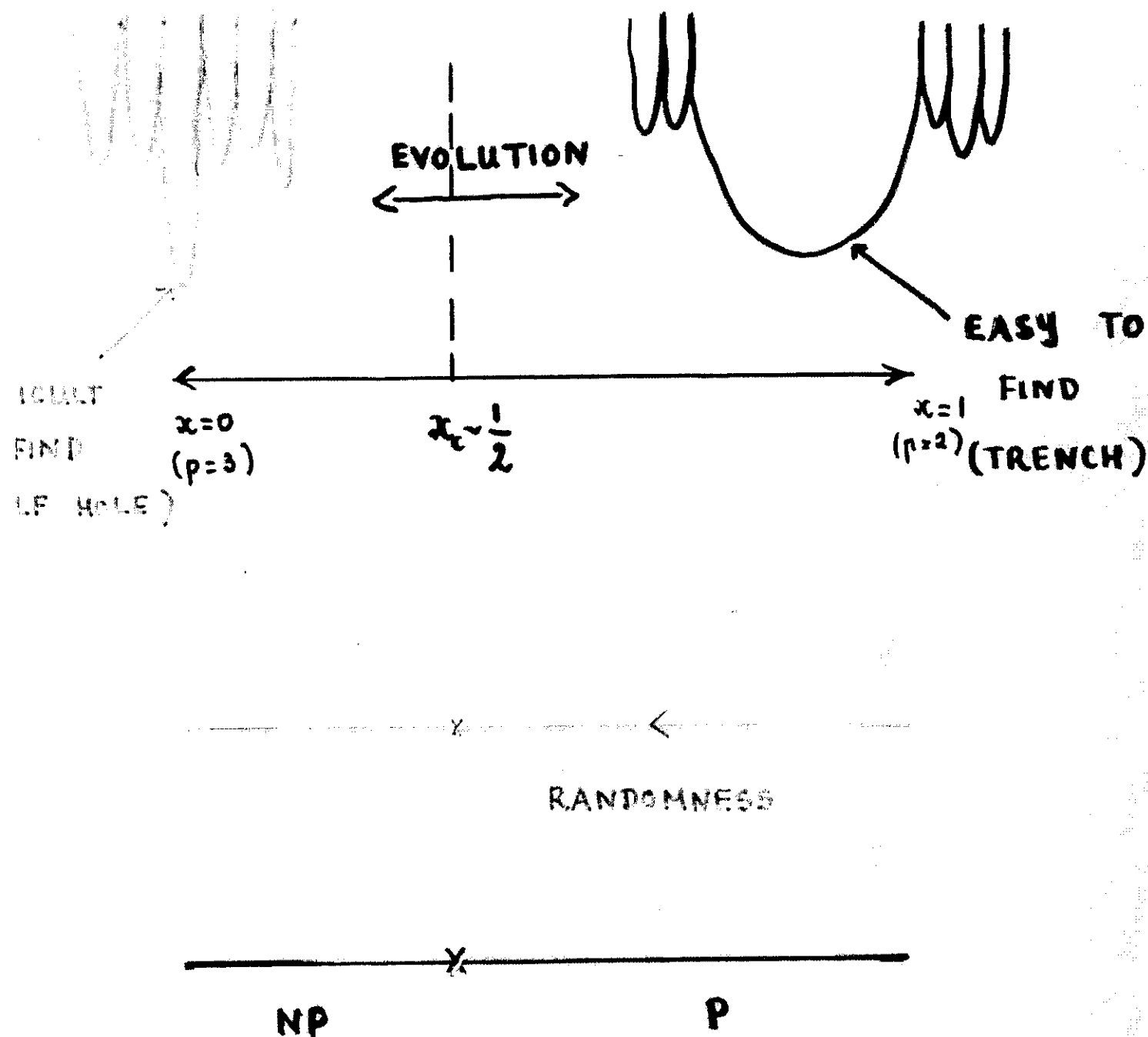
($p=3$) ATTRACTION BASINS
EXPONENTIALLY SMALL (BARRAT
+ FRANK)

$0 \leq x < 1$

- 1 STEP RSB

- $S_{\text{CONF}} \equiv \ln n_s \propto N$

QUALITATIVE RESULTS

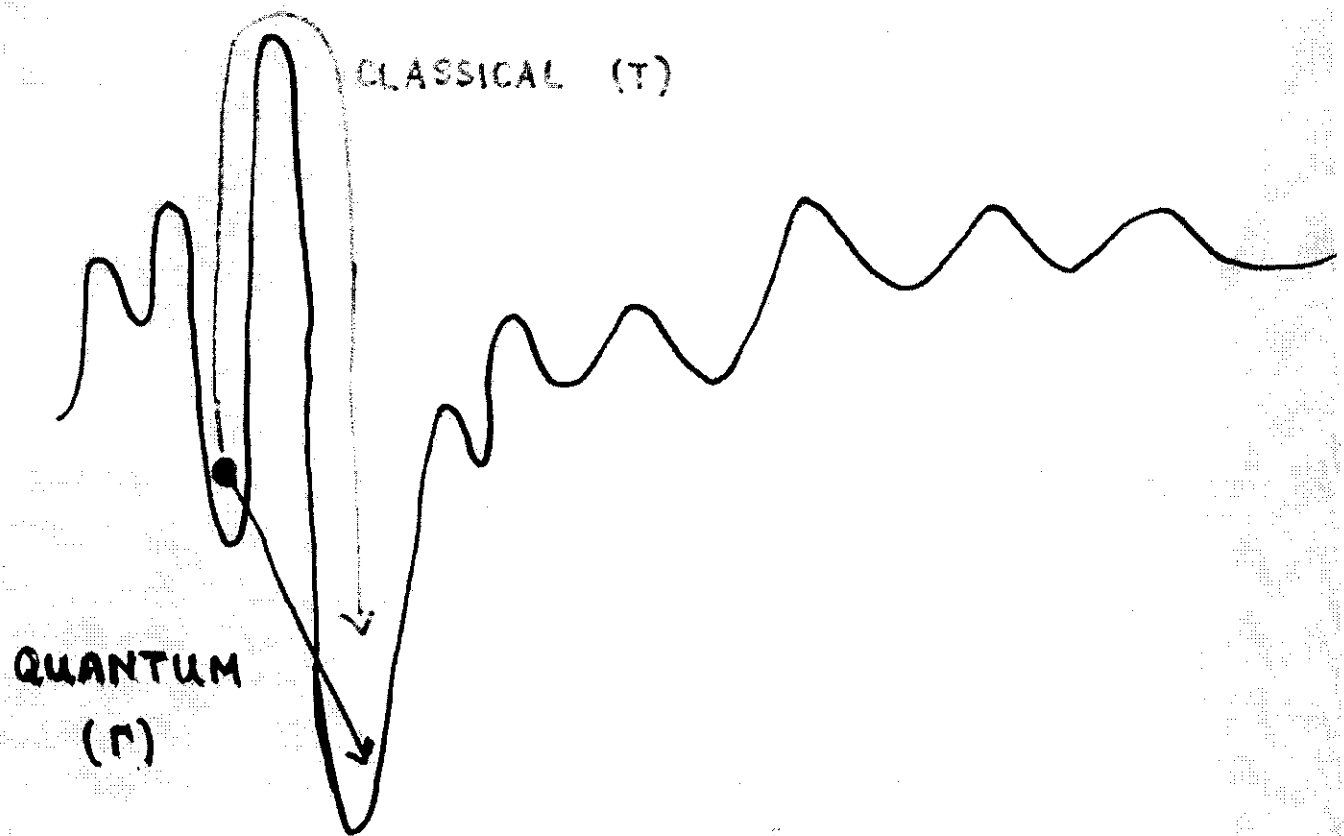


K-SAT PROBLEMS

MONASSON, ZECCINA, KIRKPATRICK

ET AL, NATURE 400, 133 (1999)

QUANTUM ANNEALING



$$H = \sum_{ij}^N J_{ij} \sigma_i^z \sigma_j^z - \Gamma \sum_i \sigma_i^x$$

• EXPERIMENT

BROOKE ET AL. SCIENCE 284, 779 (99)

• THEORY

KADOYAKI + NISHIMORI COND-MAT / 9804280

CUGLIANDOLE ET AL COND-MAT / 003208

COMPUTATION ?

MANIPULATION

ROBUST TO

OF STATES

ENVIRONMENT



WANTED :

QUANTUM GLASS W/

"HALL LIQUID"

FEATURES

- SPIN LIQUID G. S.

- GAP

SUMMARY

- DISORDER-FREE GLASS MODELS



CODING SEQUENCES

- ALTERNATIVE OPTIMIZATION METHODS

QUESTIONS

- UNIFYING PRINCIPLES UNDERLYING
SYSTEMS FAR FROM EQUILIBRIUM

• DRIVEN VS. SPONTANEOUS

Q. DOTS

GLASSES

TURBULENCE



IMPORTANCE
OF DISORDER

- COMBINED COMPLEXITIES OF CLASSICAL
+ QUANTUM MECHANICS ?
(NONLINEARITY + SUPERPOSITION) ?

