



the
abdus salam
international centre for theoretical physics

ICTP/UCSB/TWAS
MINIWORKSHOP ON "*FRONTIERS IN MATERIALS SCIENCE*"
15 - 18 May 2001

301/1311-12

*"Pushing the Limits of Electronic Structure Theory.
Can We Design New Spintronic Materials?"*

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USA

Please note: These are preliminary notes intended for internal distribution only.

Pushing the limits of
electronic structure theory
Can we design new spin-
tronic materials?

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Stefano Sanvito

Acknowledgments

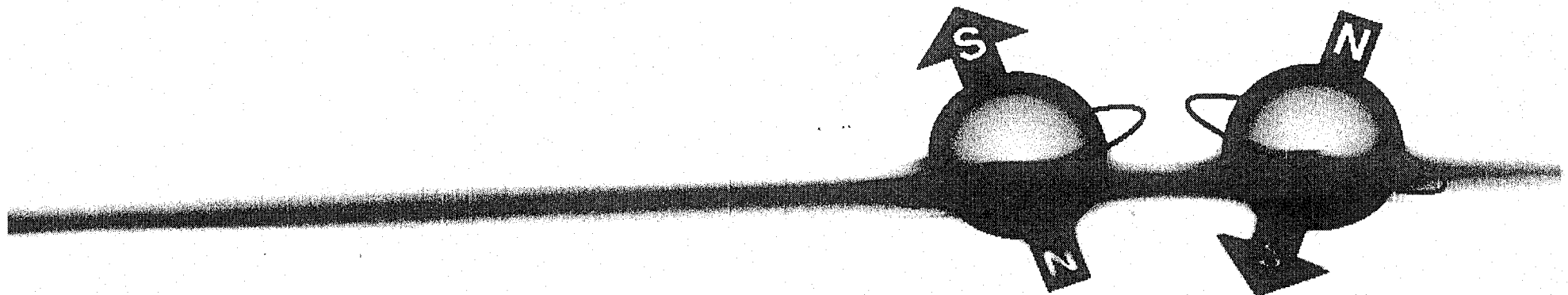
ONR
ACS / PRF
JNC + IISc Bangalore

Outline

- What is spintronics
- Theoretical techniques -
successes + shortcomings
- Solutions! Spinor + Siesta
- Results for materials

What is Spintronics?

- **Spin polarized electronics**
- exploits the fact that electrons have spin as well as charge
- tunable/controllable by magnetic fields
- extra degree of freedom

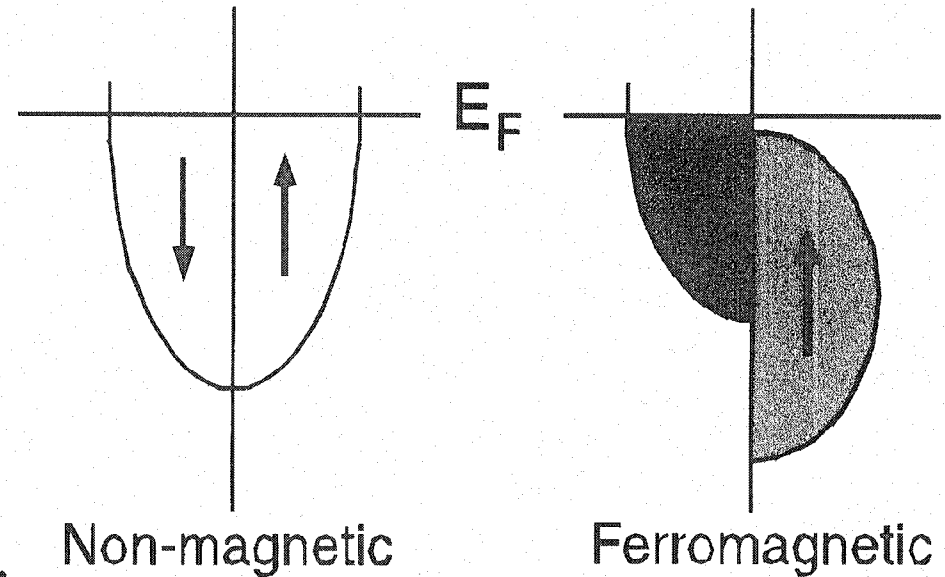


Application: Quantum computing!

- Fundamentally new mode of information processing
- Uses quantum interference and entanglement
- Requires stable, long-lived, coherent quantum states
- Good for cryptography, searching, factorization

Spin-polarized transport

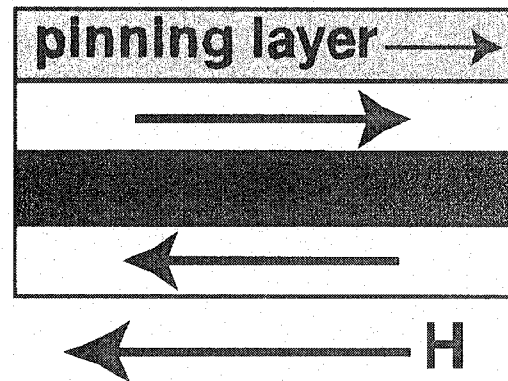
- imbalance of spin populations at Fermi level
- e.g. ferromagnetic metals
- Half-metallic ferromagnets!



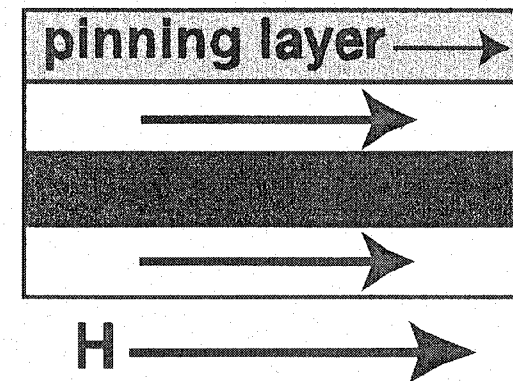
Applications

Spin valves

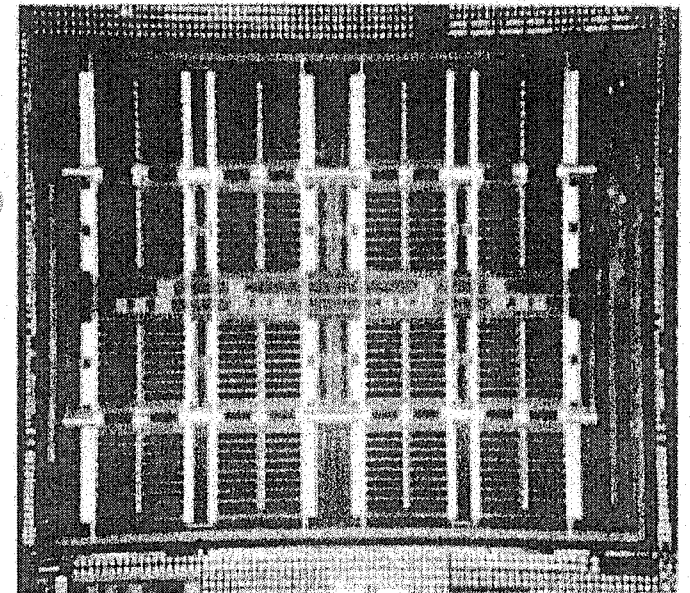
high resistance



low resistance



used in: magnetic field sensors
recording heads
MRAM



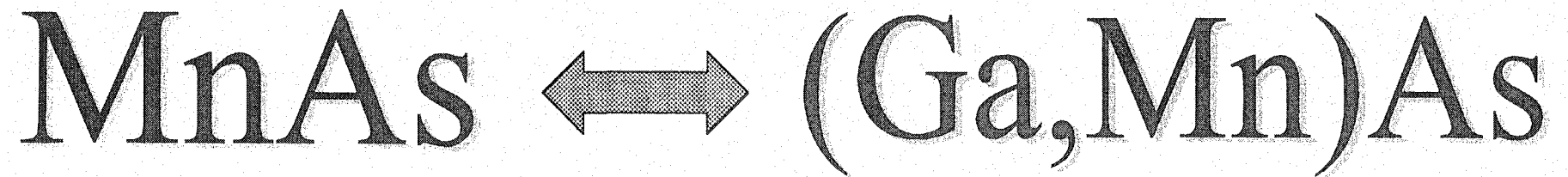
Technology-enabling material

Need:

- Room temperature ferromagnet
- Compatible with existing processing techniques
- Compatible interfaces with semiconductors

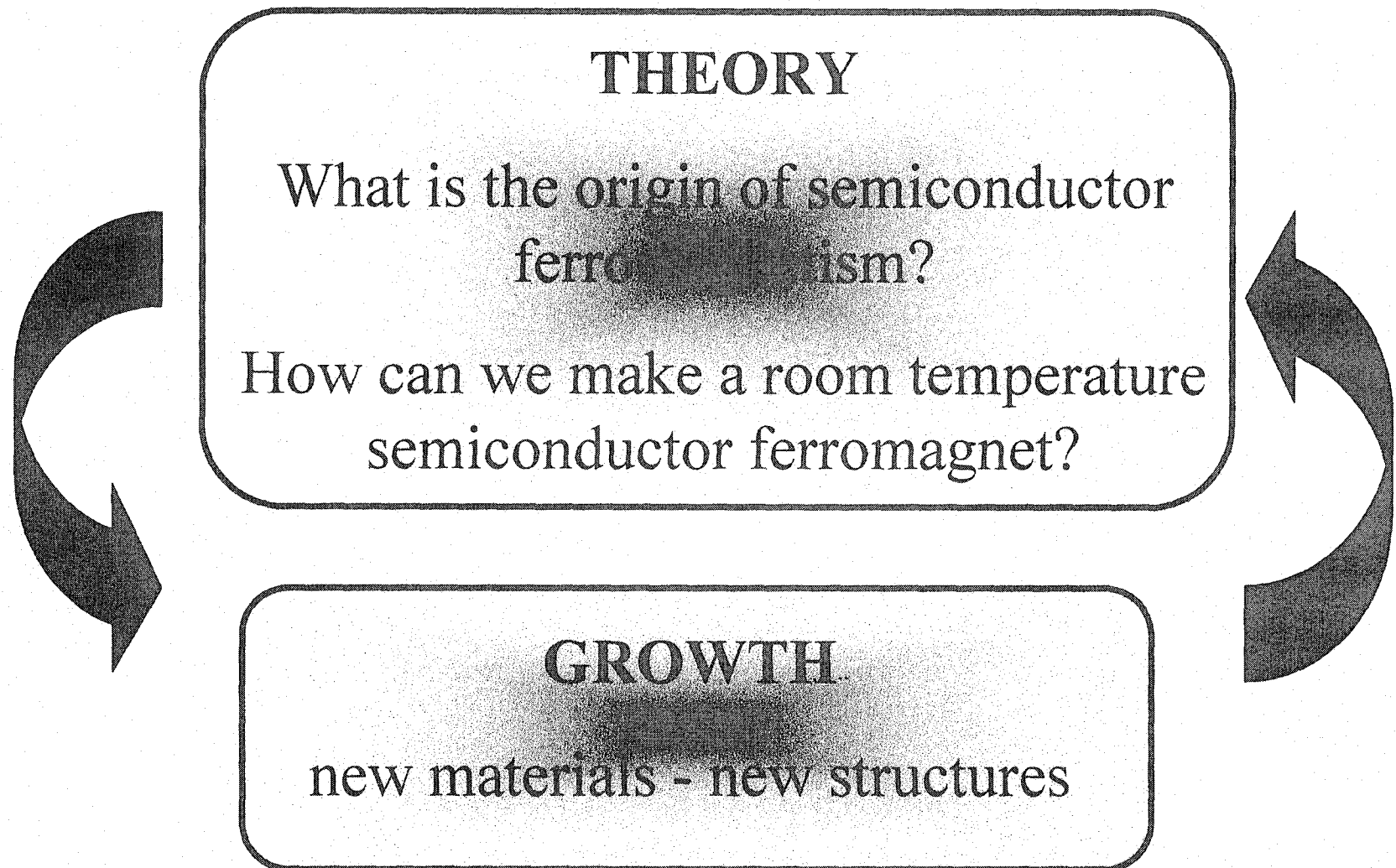
**FERROMAGNETIC
SEMICONDUCTOR!**

Promising systems

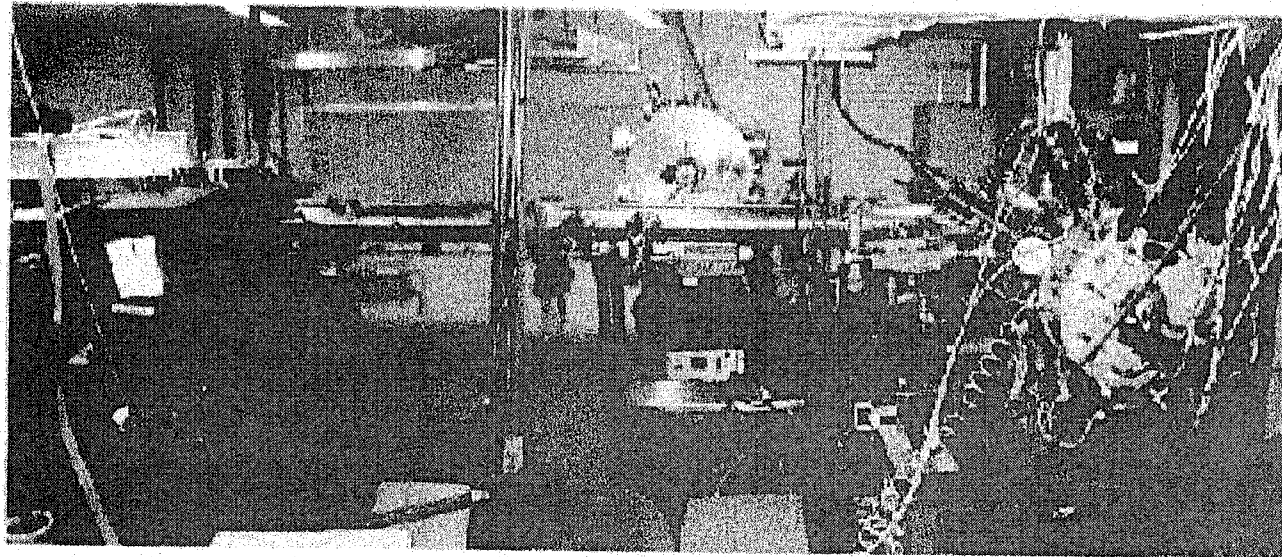


- 1) they can be grown on commonly used GaAs,
- 2) the interfaces with GaAs are thermodynamically stable,
- 3) the growth process is compatible with existing III-V MBE technology.
- 4) They can be ferromagnetic **BUT** T_c is only around 100K!

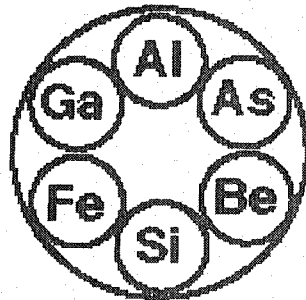
UCSB Research Effort



Spintronics MBE Laboratory

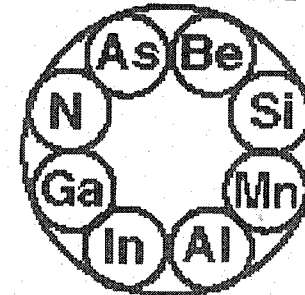


EPI-620
Ferromagnetic Metals
on Semiconductors



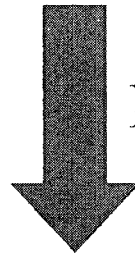
Analysis Chamber

Gen-II
Ferromagnetic
Semiconductors



Density Functional Theory

interacting many-electron system



Kohn-Sham Equations

system of non-interacting electrons

Hohenberg and Kohn (1964)

Kohn and Sham (1965)

The Kohn-Sham Equations

$$\{T + V_{ei}(\mathbf{r}) + V_H(\mathbf{r}) + V_{xc}(\mathbf{r})\} \phi_i(\mathbf{r}) = \varepsilon_i \phi_i(\mathbf{r})$$

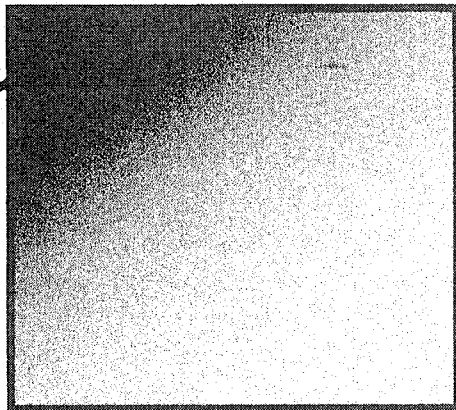
where both the Hartree potential, V_H , and the exchange correlation, V_{xc} , depend on the density

$$\rho(\mathbf{r}) = \sum \phi_i^*(\mathbf{r}) \phi_i(\mathbf{r})$$

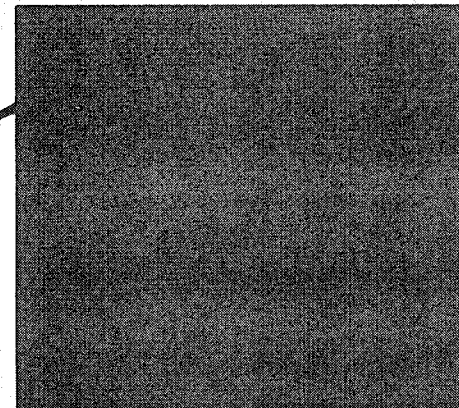
Local Density Approximation (LDA)

assumes that the exchange-correlation energy per electron at a point \mathbf{r} in the electron gas is equal to the exchange-correlation energy per electron in a homogeneous electron gas that has the same density as the electron gas at point \mathbf{r} .

$E_{xc}(\mathbf{r})$
(True)



$E_{xc}(\mathbf{r})$
(LDA)

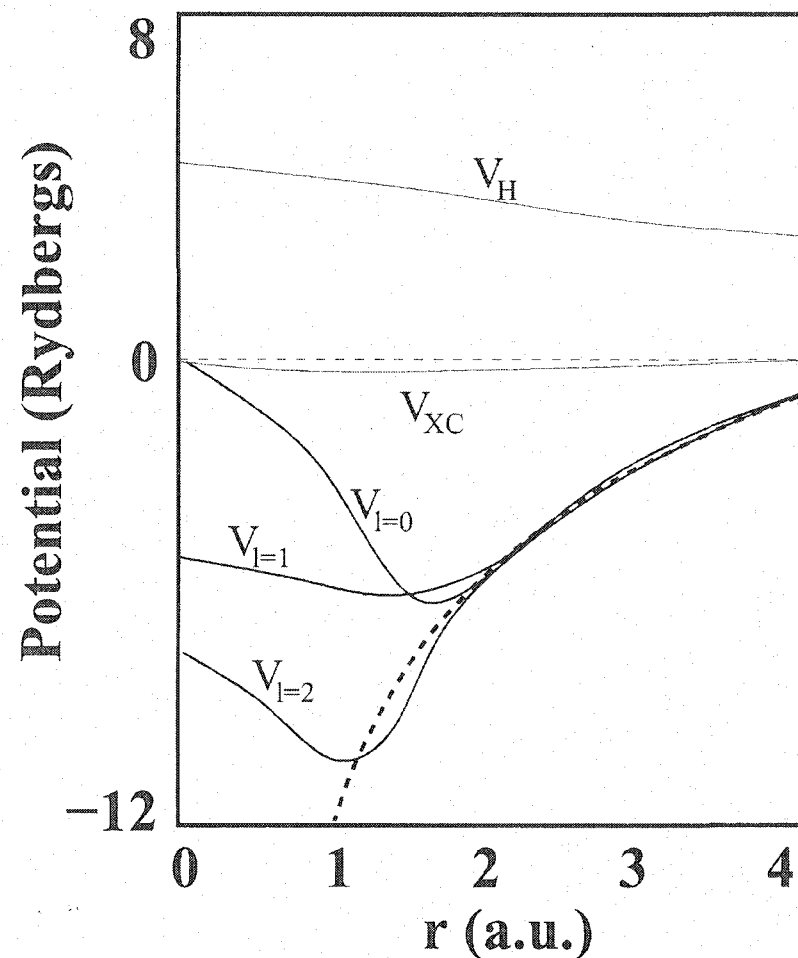


Pseudopotential Approximation

- Replacement:

$$V_{nuc} \rightarrow V_{ion}$$

- Treat only valence electrons
- Nodeless wave functions
- Plane waves are o.k.



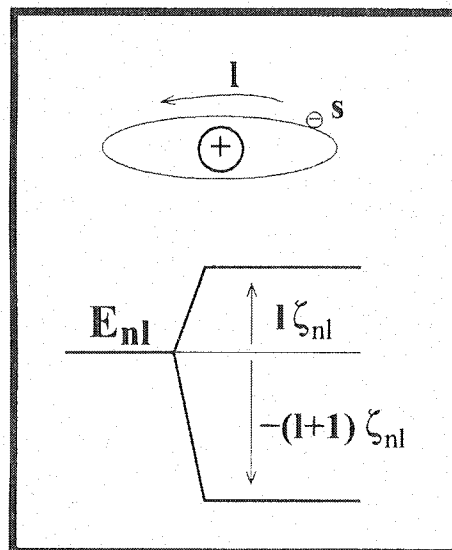
Pseudopotential Success Story

The *ab initio* pseudopotential method has been successfully employed in the prediction of structural and electronic properties for a wide range of materials, including:

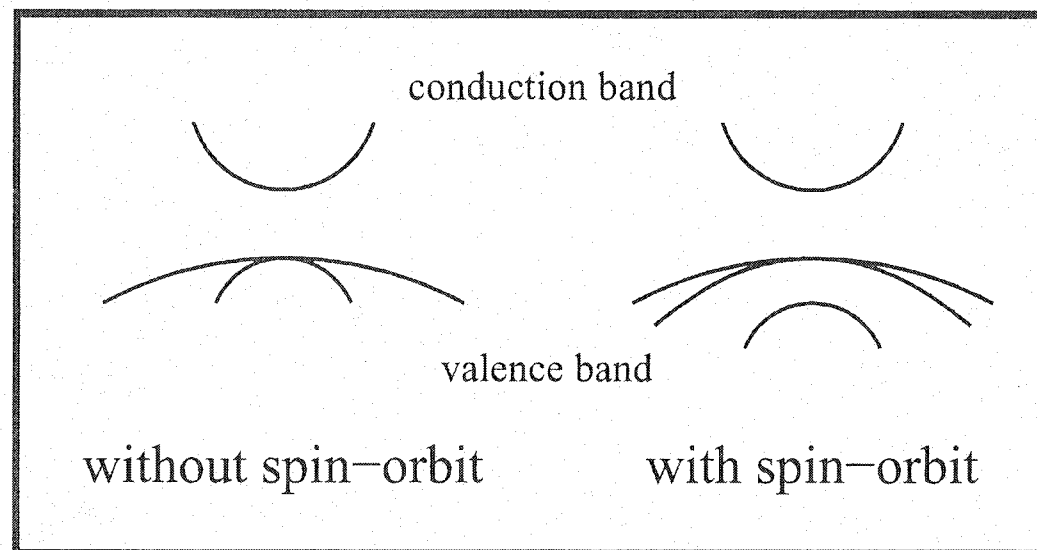
- Silica Frameworks
- Semiconductors
- Ferromagnetic Ferroelectrics !
- Metals
- Transition metals

What's missing? Spin-Orbit Coupling

Atom:



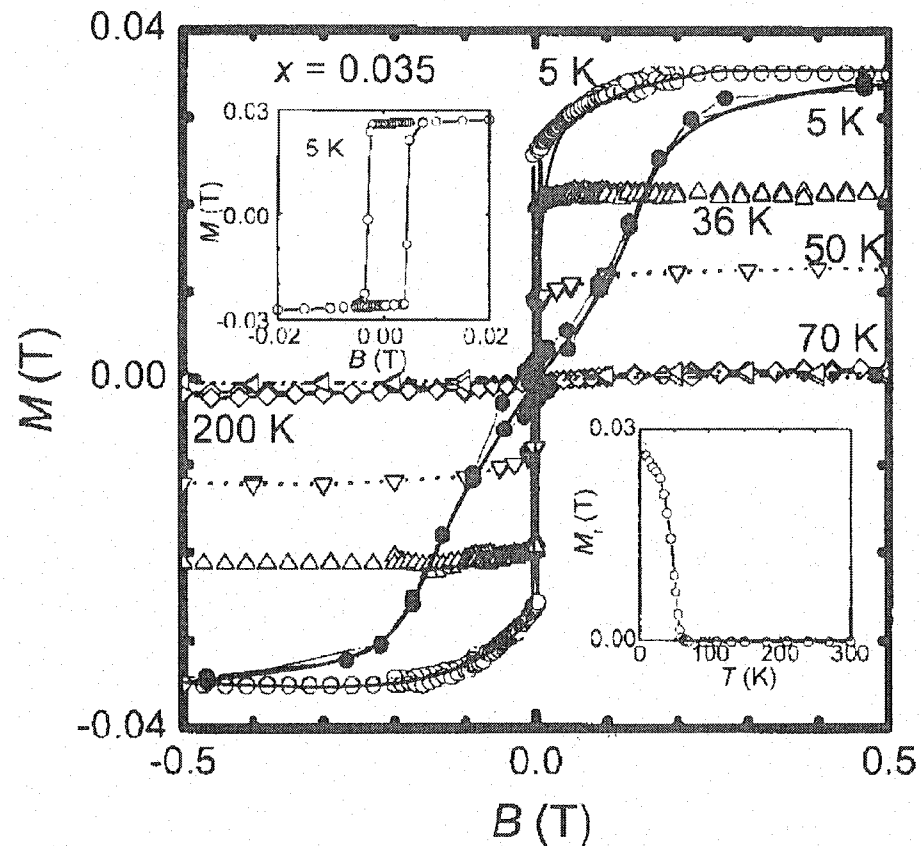
Semiconductor:



What's missing?

Magnetocrystalline Anisotropy

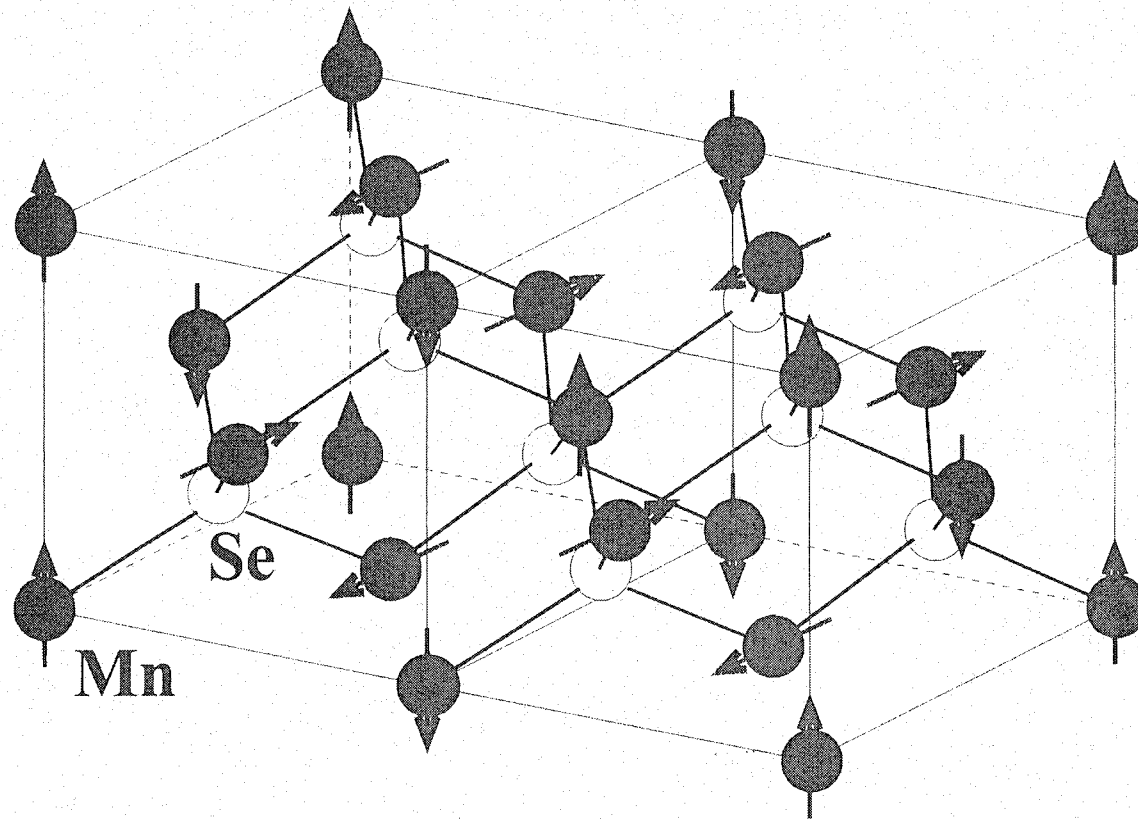
Easy & hard axis: (Ga,Mn)As



H. Ohno and F. Matsukura, Solid State Commun. **117**, 179 (2001)

What's missing? Non-collinear Magnetism

Anisotropic superexchange (Dzyaloshinskii–Moriya): MnSe



How can we fix it?

- Wave functions become general spinors

$$\psi = \phi_{\uparrow}(r)|\uparrow\rangle + \phi_{\downarrow}(r)|\downarrow\rangle .$$

- Local density approximation needs to be generalized

$$\vec{b}_{xc}(\vec{r}) = \frac{\delta E_{xc}[n, \vec{m}]}{\delta \vec{m}(\vec{r})} = \frac{\delta E_{xc}[n, m]}{\delta m(\vec{r})} \hat{m}(\vec{r}) .$$

- Pseudopotential needs to include spin–orbit coupling.

Fully Separable j -dependent Pseudopotentials

$$V_{pp}^{KB} = V_{loc} + \sum_{l,j,m_j} \frac{|\delta V_{l,j} \phi_{l,j,m_j}\rangle \langle \phi_{l,j,m_j} \delta V_{l,j}|}{\langle \phi_{l,j,m_j} | \delta V_{l,j} | \phi_{l,j,m_j} \rangle}$$

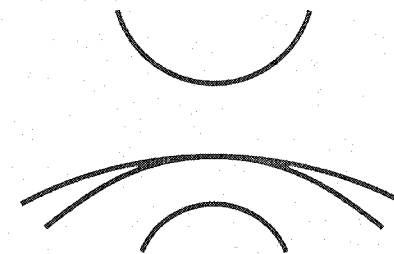
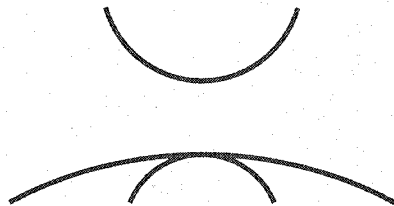
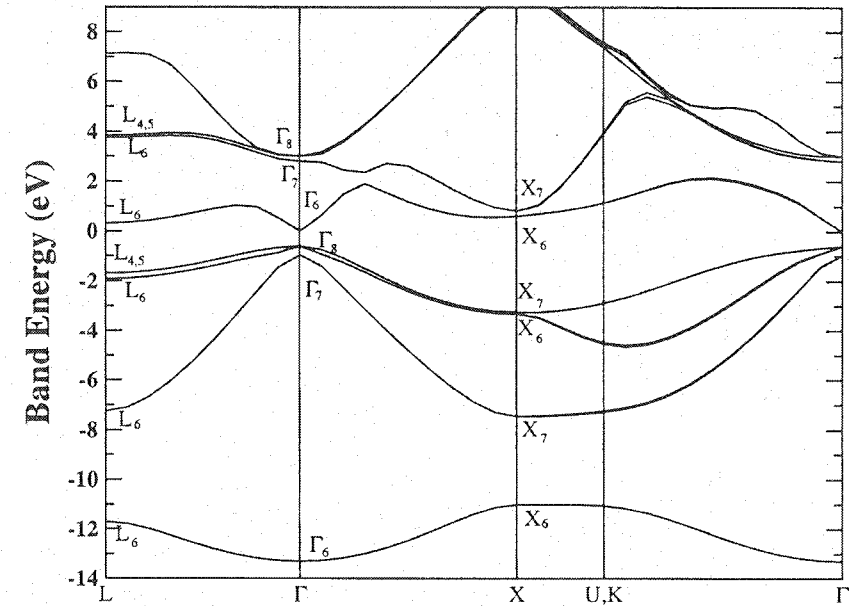
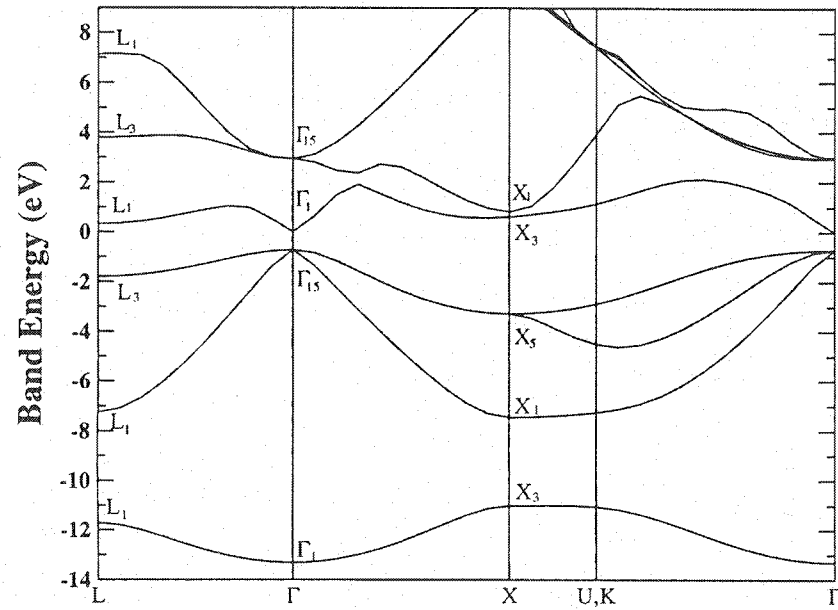
Pseudo wave function: $|\phi_{l,j,m_j}^{i_s,i_a}\rangle = |R_{l,j}^{i_s,i_a}\rangle |\Phi_{m_j}^{l,j}\rangle$

$$\text{for } j = l + \frac{1}{2}: \quad |\Phi_{m_j}^{l,j}\rangle = \left(\frac{l+m+1}{2l+1}\right)^{\frac{1}{2}} |Y_l^m\rangle |\uparrow\rangle + \left(\frac{l-m}{2l+1}\right)^{\frac{1}{2}} |Y_l^{m+1}\rangle |\downarrow\rangle$$

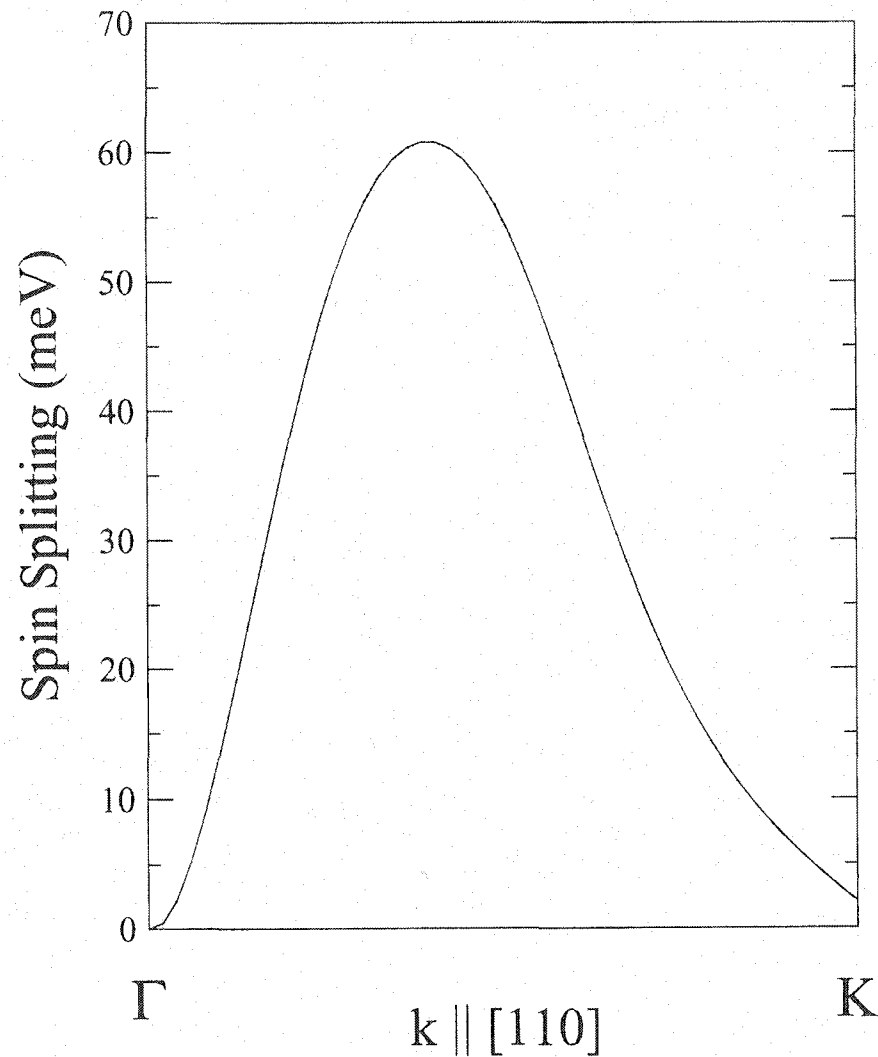
$$\text{for } j = l - \frac{1}{2}: \quad |\Phi_{m_j}^{l,j}\rangle = \left(\frac{l-m+1}{2l+1}\right)^{\frac{1}{2}} |Y_l^{m-1}\rangle |\uparrow\rangle - \left(\frac{l+m}{2l+1}\right)^{\frac{1}{2}} |Y_l^m\rangle |\downarrow\rangle$$

<http://www.mrl.ucsb.edu/~theurich/Spinor>

GaAs: scalar vs. fully relativistic



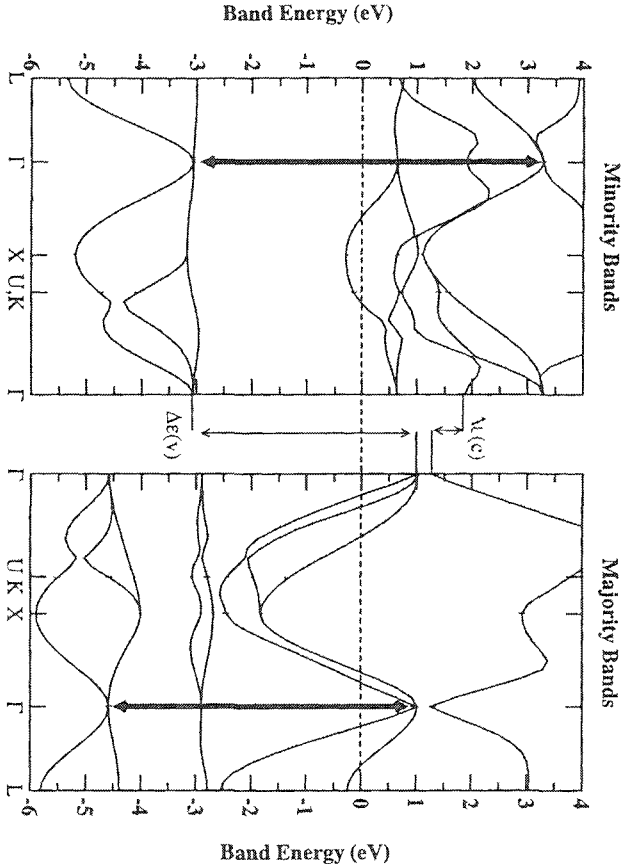
Spin-splitting in GaAs



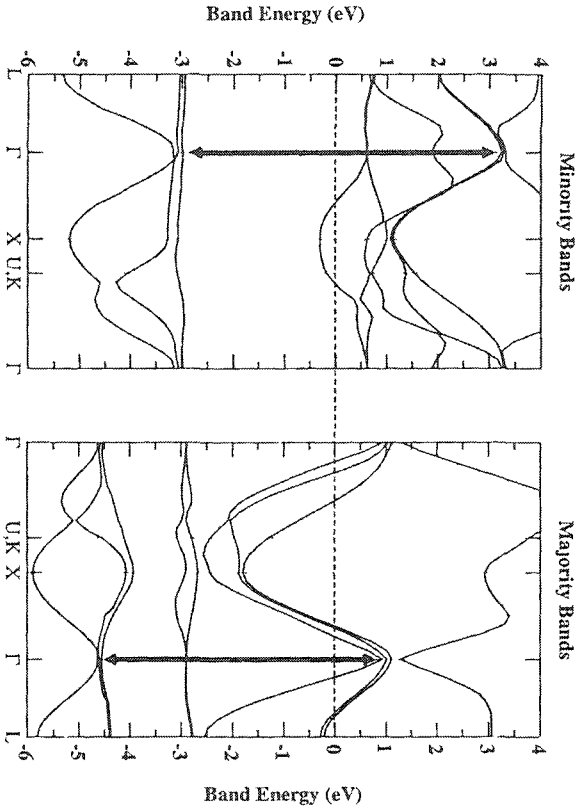
- Zincblende inversion asymmetry allows spin splitting
- Spin splitting is k dependent
- Spin splitting causes spin relaxation

[001] FM zincblende MnSe

Scalar-Relativistic



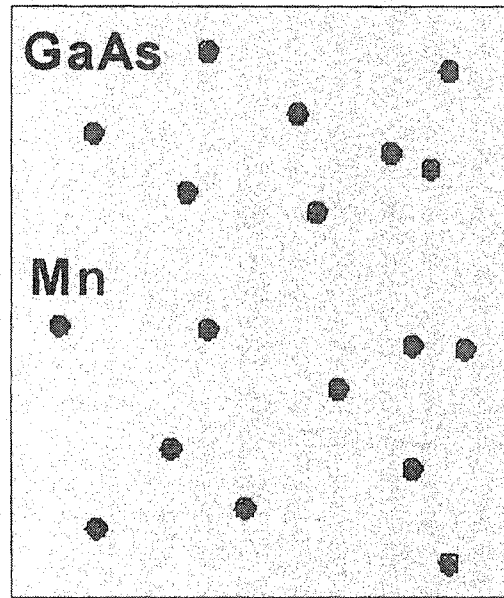
Fully Relativistic



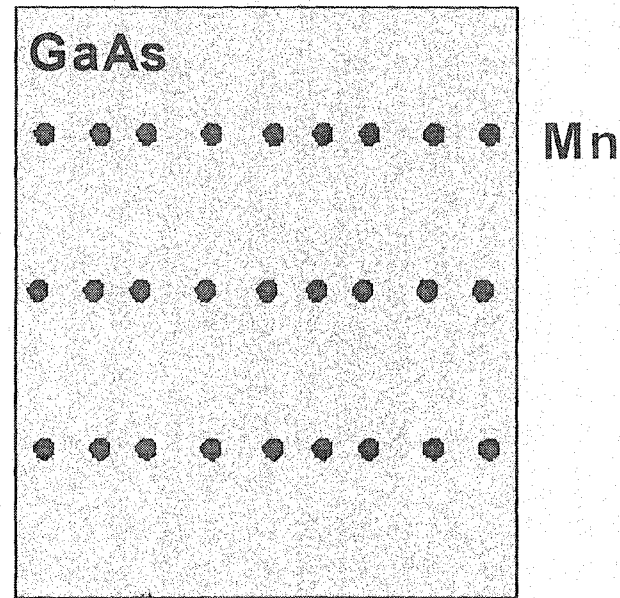
	Fully Relativistic Scalar-Relativistic	M (μ_B)	$N\alpha(\text{eV})$	$N\beta(\text{eV})$
theory		4.41	0.28	-1.86
theory		4.42	0.28	-1.85
experiment	Zn _{1-x} Mn _x Se (Twardowski et al., 1984)	-	0.26	-1.31

New materials and structures

Random Alloy



“DFH”



Theoretical Techniques

Density functional theory within the LSDA. Details:

MnAs

“SPINOR”

<http://www.mrl.ucsb.edu/~theurich/Spinor>

- Plane-wave basis (cutoff energy 70 Ryd) [2]
- Optimized Pseudopotentials
- Small unit cells
- non-collinear magnetism and spin-orbit coupling

(Ga,Mn)As

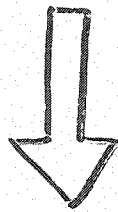
“SIESTA”

D. Sanchez-Portal et al., Int. J. Quant. Chem. **65**, 453 (1997).

- Localized multiple- ζ Pseudo-atomic orbital basis
- Troullier-Martin Pseudopotentials
- Super-cells with up to 100 atoms

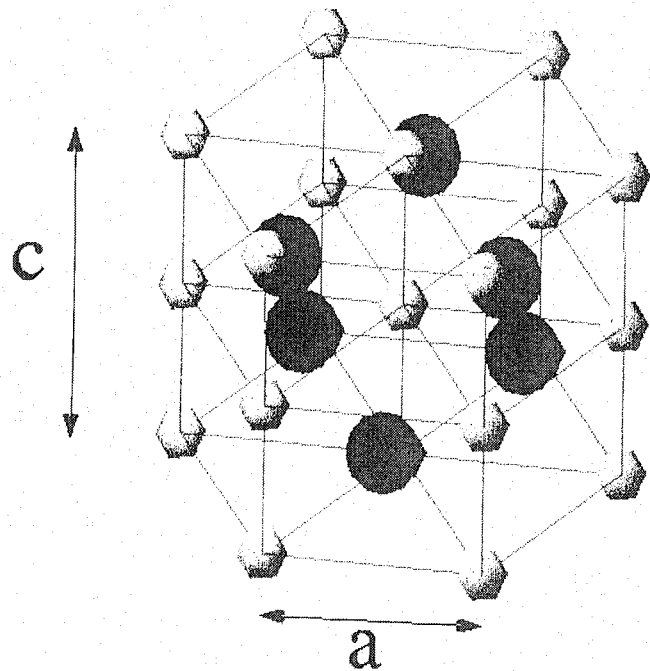
QUESTIONS WE'D LIKE TO ANSWER:-

- Can we grow thicker MnAs layers
(and do we want to?)
- Where does the ferromagnetism
come from?
- Can we increase T_c ?
- Smaller is Better. how will quantum
confinement affect exchange interactions:

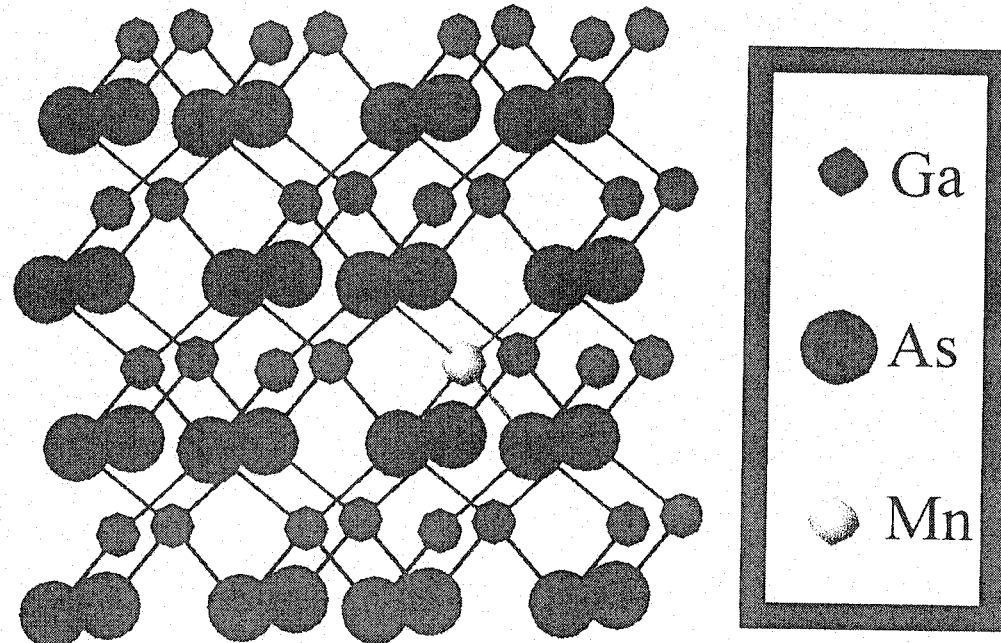


Design of new (and better!)
spintronic materials.

MnAs: two crystalline phases



NiAs-type

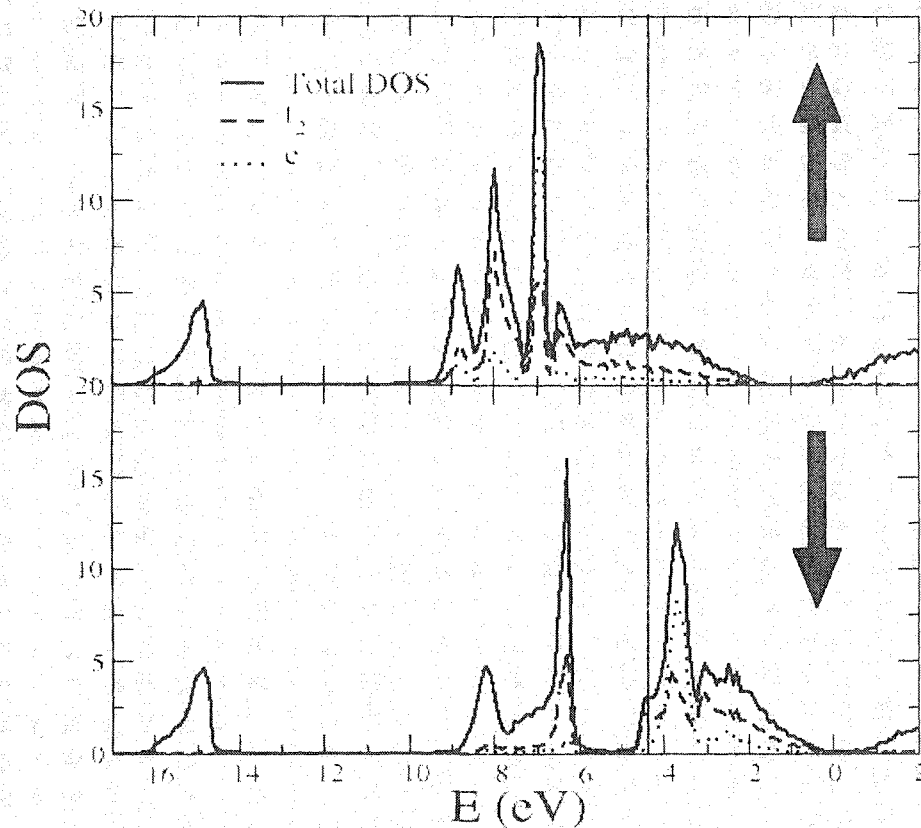


Zincblende (Ga,Mn)As

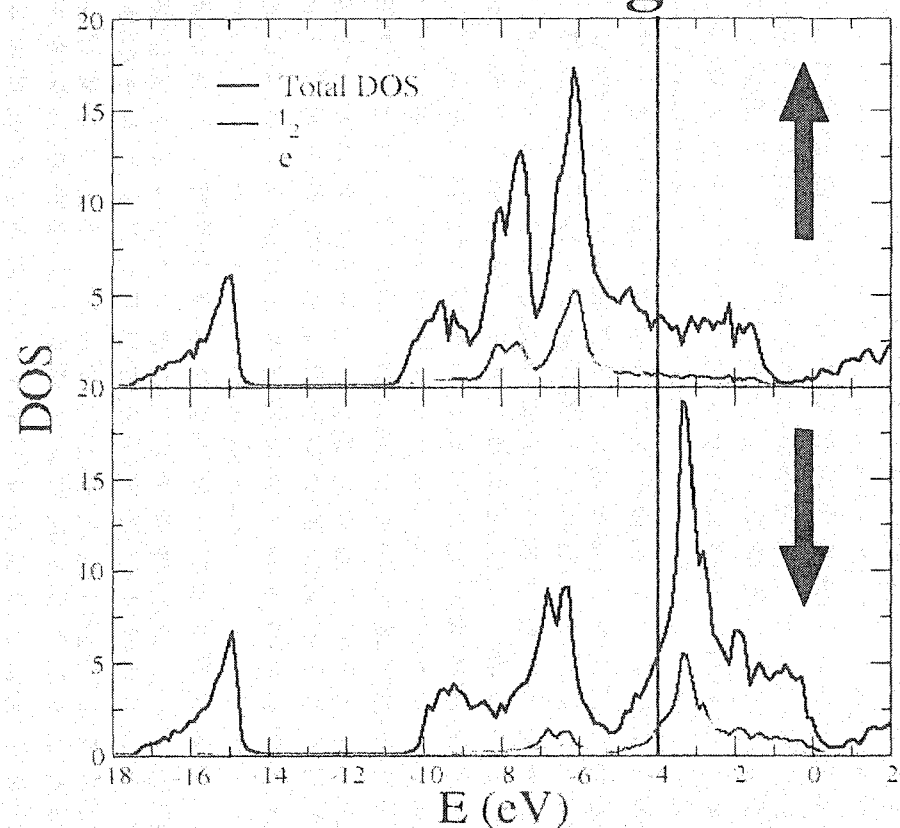
- 1) Bulk MnAs has NiAs-type structure
- 2) ZB MnAs can be MBE-grown up to 1 monolayer
- 3) ZB MnAs important to understand the Mn-Mn exchange coupling in (Ga,Mn)As
- 4) Experimentally (Ga,Mn)As $\text{Mn} < 7\%$

MnAs: ZB vs Hexagonal

LDOS Zincblende

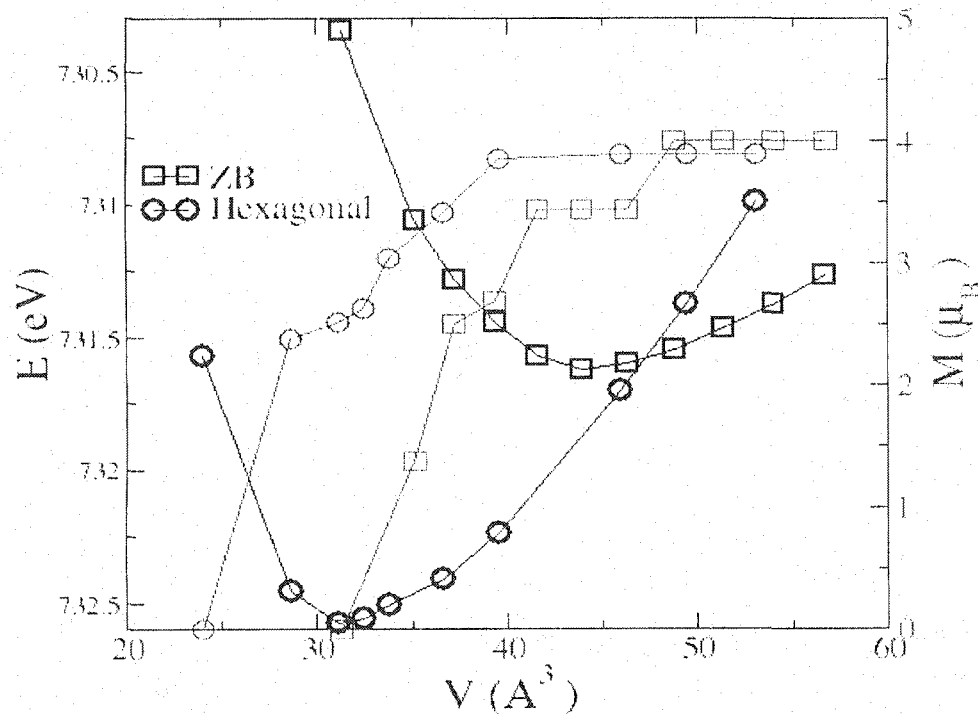


LDOS Hexagonal



- In the minority band of ZB MnAs the Fermi Energy cuts through a dispersionless d-band. For large lattice spacing it is **Half Metallic**

MnAs: ZB vs Hexagonal

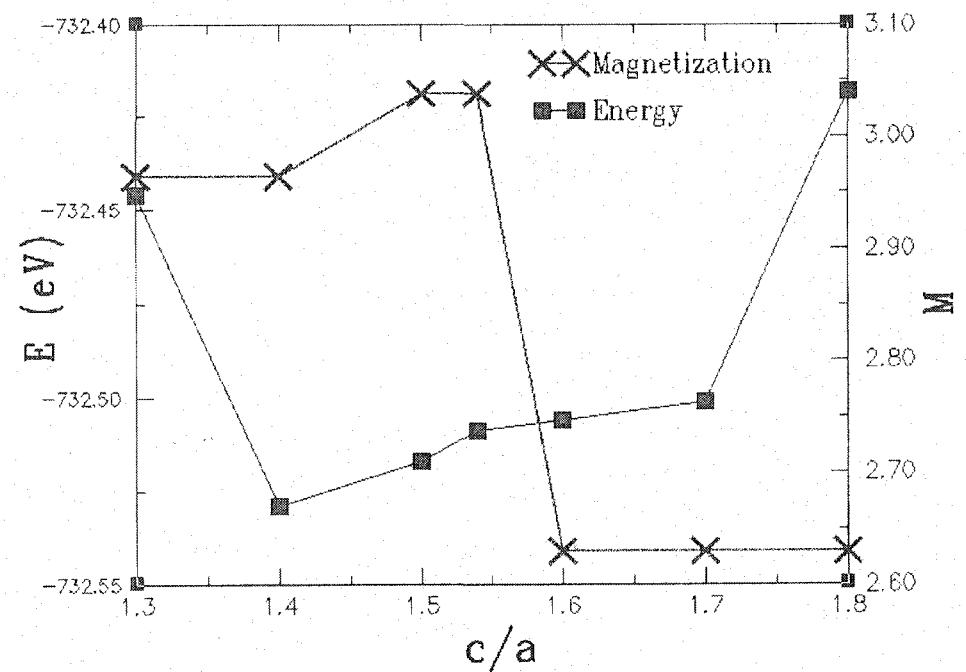


- NiAs-type is more stable with smaller unit cell volume

- Crossover volume for large stress

- NiAs-type MnAs can accommodate large distortions

- ZB MnAs difficult to stabilize



Mulliken Analysis for ZB MnAs

Mulliken Population analysis confirms the **antiferromagnetic coupling** between Mn-d and As-p states in MnAs:

$$n(\text{Mn} - d \uparrow) - n(\text{Mn} - d \downarrow) = 3.787$$

$$n(\text{As} - p \uparrow) - n(\text{As} - p \downarrow) = -0.285$$

Conclude - FM is mediated
by As-p based holes

MnAs: Conclusion

• ZB MnAs presents half-metal behavior and large spin polarization: ideal conditions for spintronics applications

but

• NiAs-type is more stable with smaller unit cell volume.
It can also accommodate large distortions

• band spin split and Mulliken population analysis

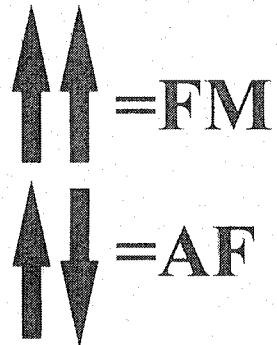
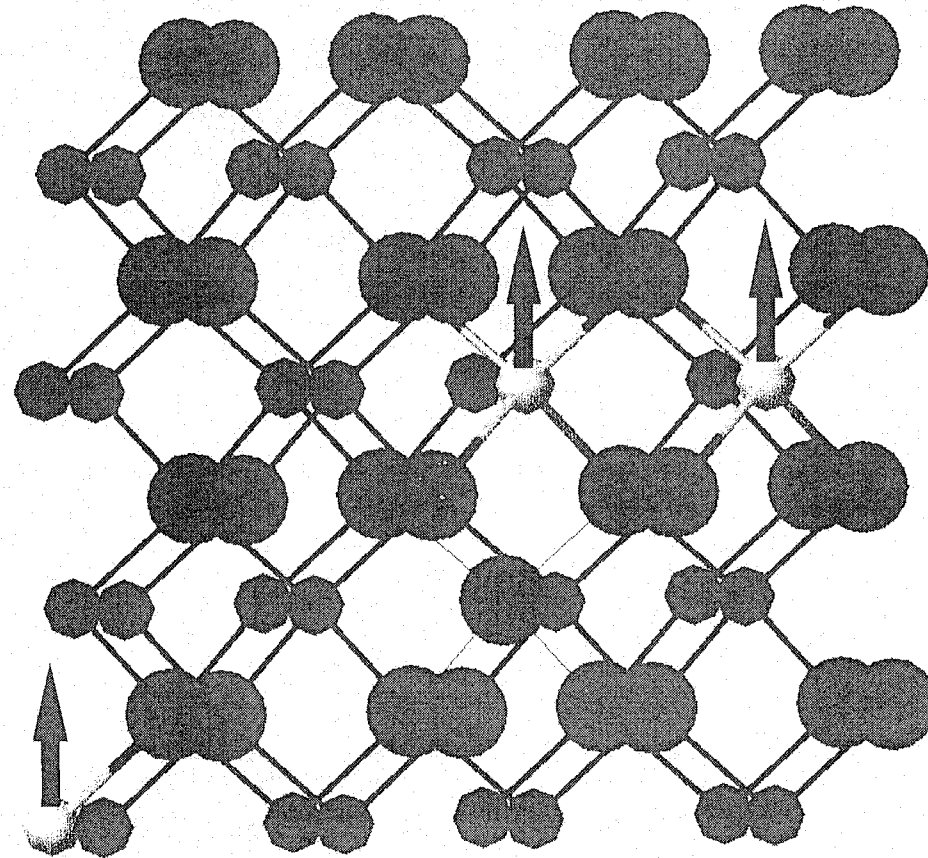
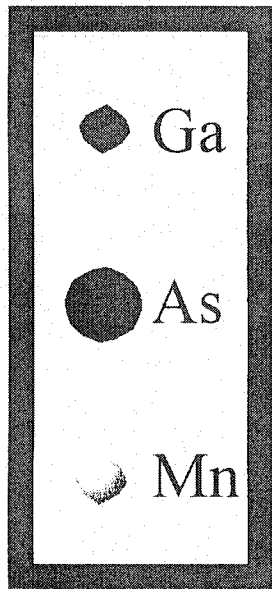
p-d coupling is

**ANTIFERROMAGNETIC
and drives the ferromagnetism**

(Ga,Mn)As

- Grown by low temperature MBE, with random or digital arrangement of Mn ions.
- Inclusion of As antisites is unavoidable
- Study effect of As antisites on magnetic properties. Each As antisite is a double donor, so for a Mn density twice as large as the As antisite density compensation is expected
- Study effect of Mn arrangement (random versus DFHs).

(Ga,Mn)As

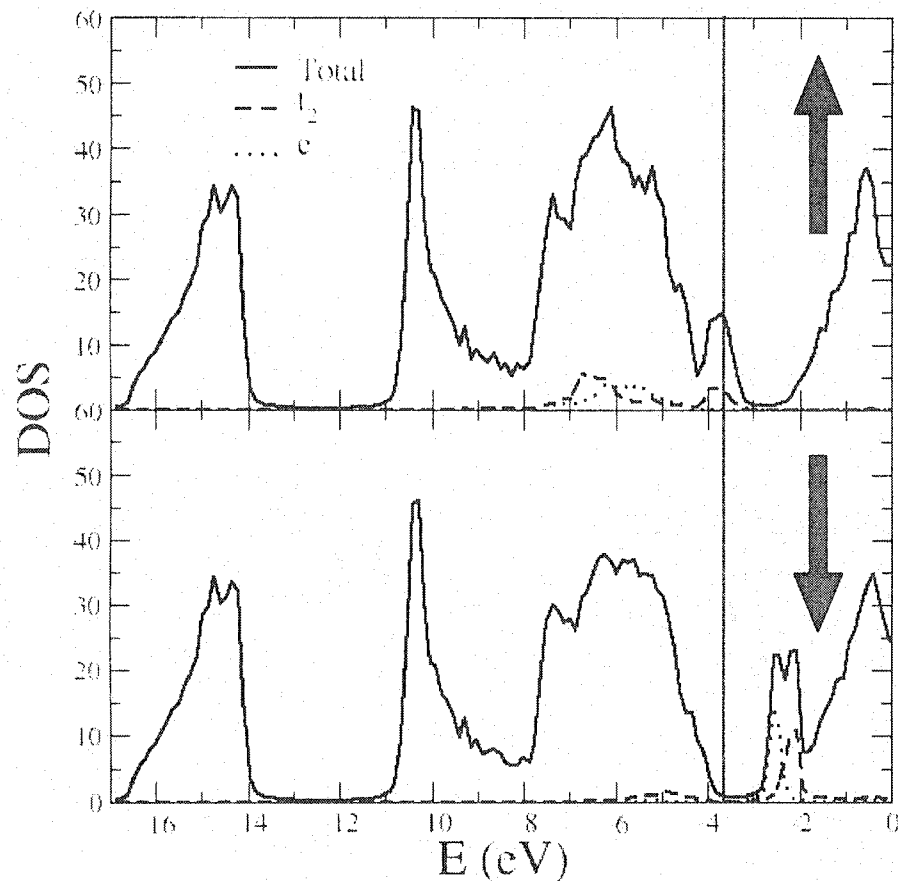


1) Effects of As antisites
on the ferromagnetism

$$\text{exp: } [\text{Mn}] > p$$

2) Effects of the Mn
distribution on the
ferromagnetism

GaAs:Mn Local DOS



32 atom unit cell with 1 Mn impurity

- Large spin-splitting of Mn-d orbitals

- The Fermi energy cuts through the Mn-d impurity band in the majority spin band

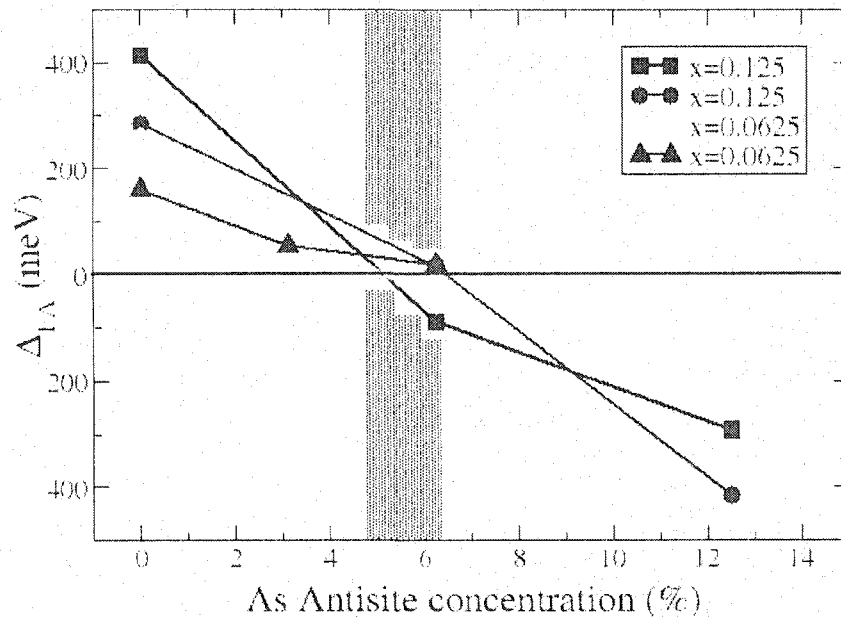
- Almost no-occupation of the Mn-d orbitals in the minority band

- Mn impurity band strongly hybridized with the As-p orbitals of the nearest neighbors

(Ga,Mn)As and As antisites

Do As antisites weaken the ferromagnetism in (Ga,Mn)As?
We consider the energy split Δ between the ferromagnetic and antiferromagnetic alignment of Mn ions in a large GaAs cell

$$\Delta_{FA} = E_{FM} - E_{AF}$$



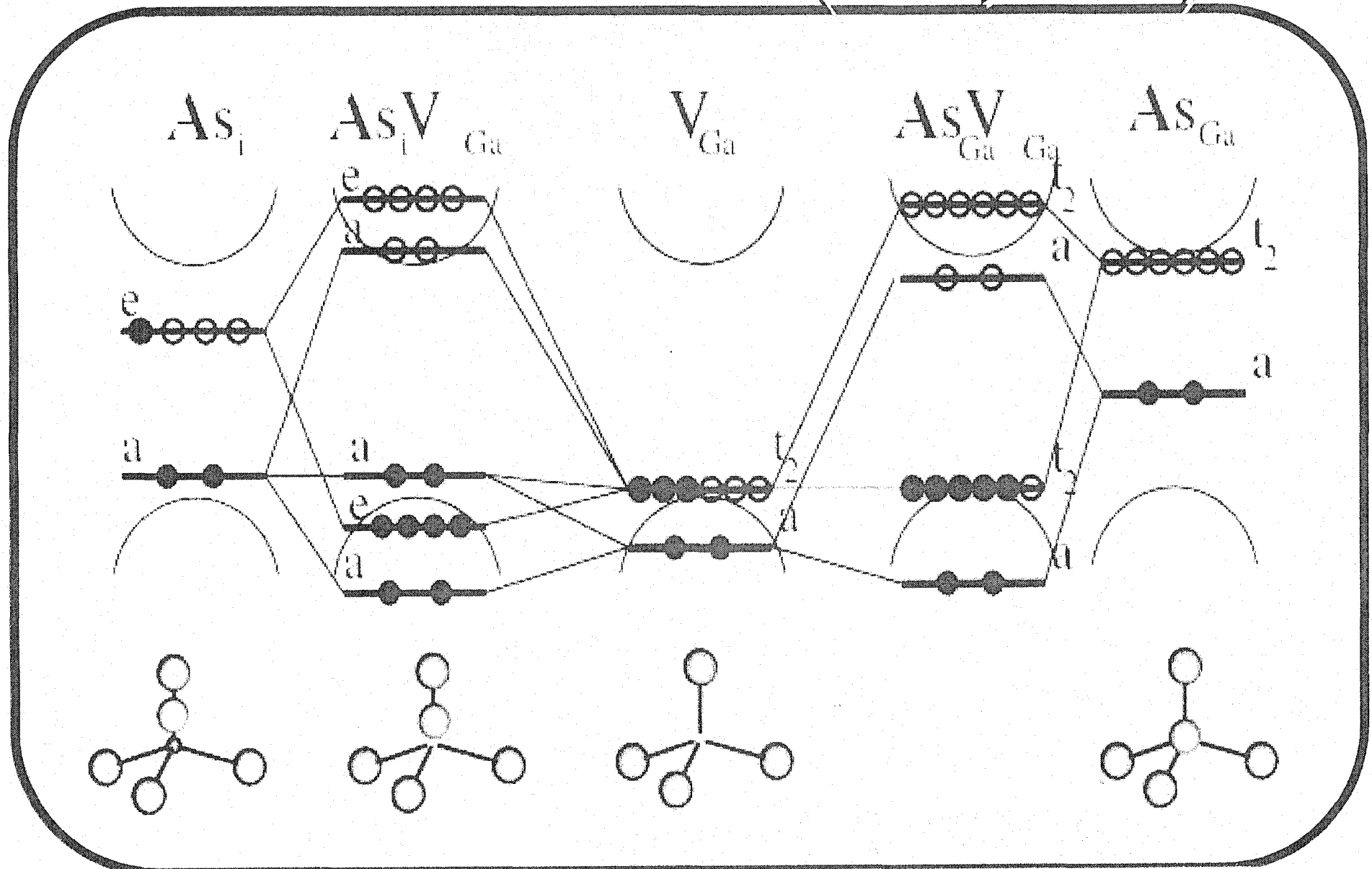
- No As antisites = strong FM order

- Presence of As antisites weakens FM alignment

- Picture of Zener exchange model not strictly valid since ferromagnetic order at compensation

- Also short range AF contribution to ferromagnetism

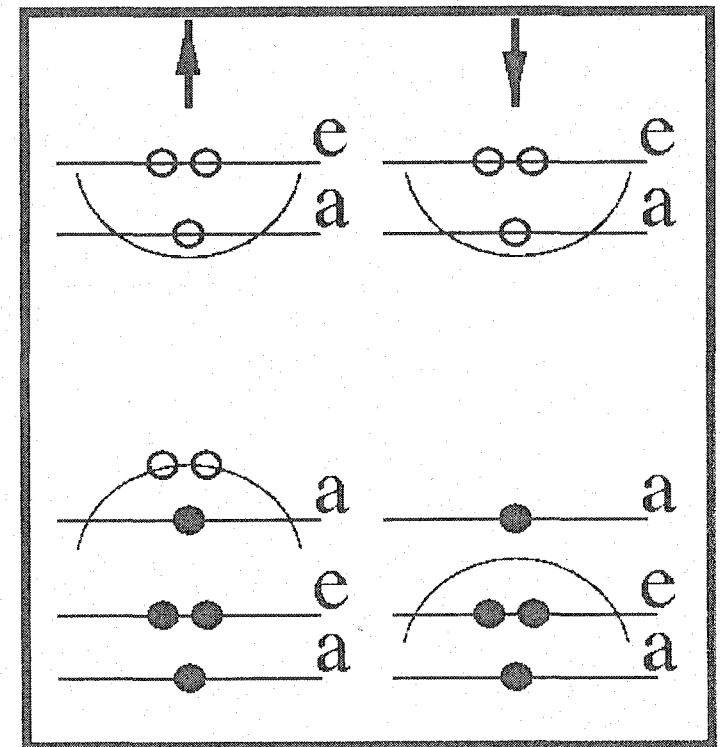
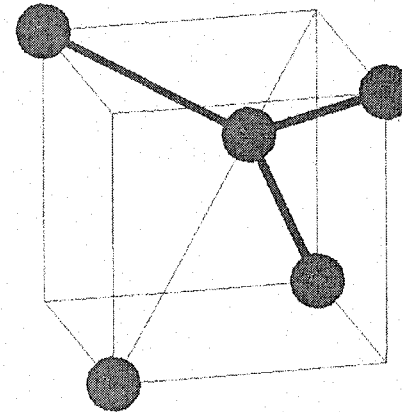
Other defects in (Ga,Mn)As



As interstitial/Ga-vacancy pair

- No suppression of the FM coupling $\Delta_{AF} \cong 120\text{meV}$
- Polarization similar to the defect-free case $P_{\text{Mn}} = 0.723$
- No extra-carrier
- Metastable-state from As antisite

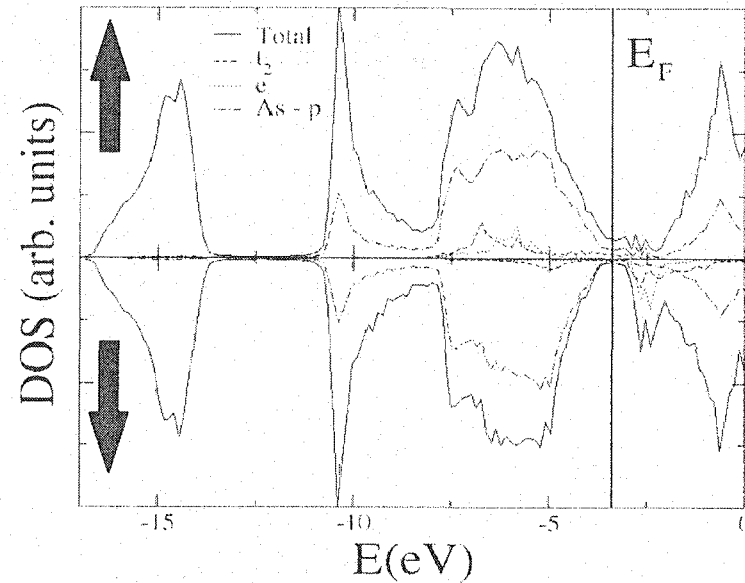
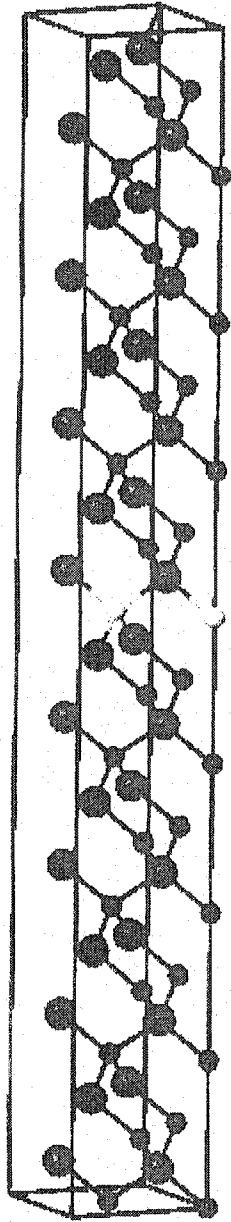
The formation of this complex from As antisite strongly enhances the ferromagnetic order of (Ga,Mn)As. This metastable configuration can be induced optically



(Ga,Mn)As Predictions

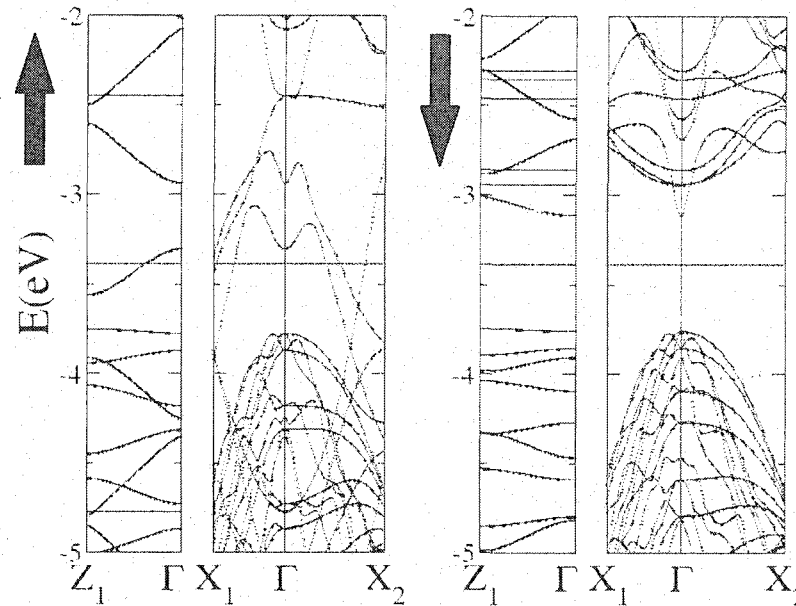
- Conversion of As antisites into As interstitial - Ga vacancy pairs will strongly enhance the ferromagnetism.
- Such a conversion can be achieved by photoexcitation [Scheffler et al.]
- PHOTOEXCITATION will increase the ferromagnetic Curie temperature in (Ga,Mn)As.

Digital Magnetic Structures



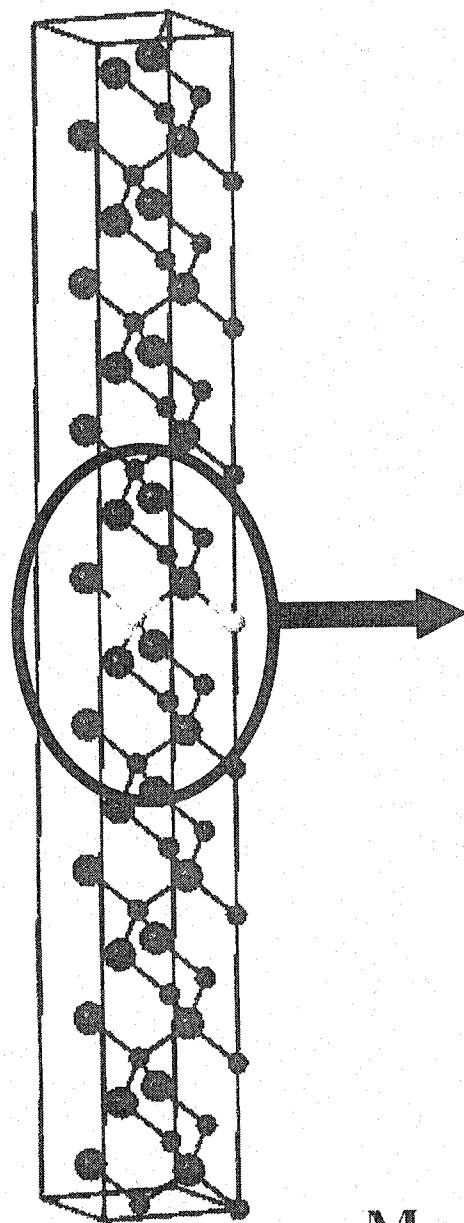
Half-metal

The DOS is similar to the random alloy but small at the Fermi energy

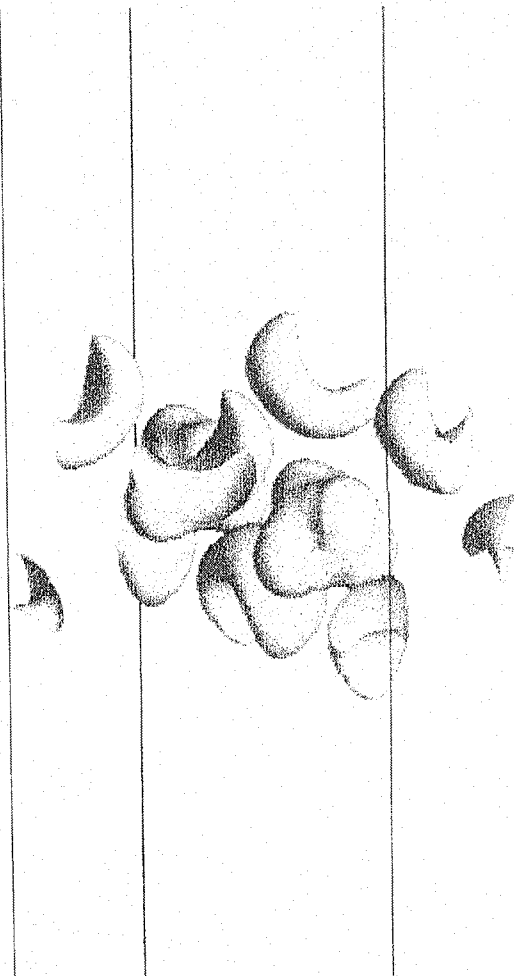


Large dispersion parallel to MnAs plane

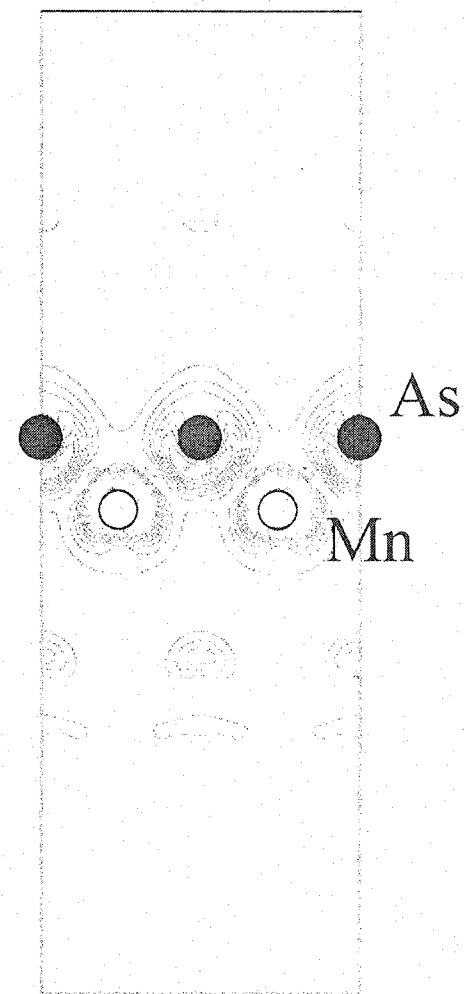
Small dispersion perpendicular to MnAs plane



Isosurface



Contour plot

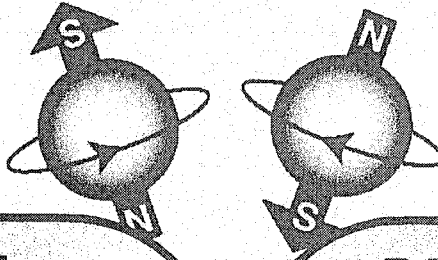


**MnAs monolayers embedded into GaAs behave like
two-dimensional half-metals**

(Ga,Mn)As Conclusion

- The atomic configuration of Mn in GaAs is compatible with both d^5 and d^6 electron densities
- The Mn-d band is **ANTIFERROMAGNETICALLY** coupled with the As-p band
- In the absence of As antisites (Ga,Mn)As is found to be ferromagnetic
- The arrangement of the Mn ions influences the strength of the ferromagnetic ordering
- As antisites weaken the ferromagnetism and a transition to an antiferromagnetic alignment is possible
- The positions of the antisites also affect the strength of the exchange interactions.

Future work



New Materials

**Mn-doped GaN - stronger
p-d hybridization**

**Mn-doped ZnO -
ferromagnetic and
piezoelectric!?**

New Techniques

**Spin-orbit coupling
versus p-d hybridization**

Non-collinear magnetism