

SUMMER SCHOOL
on
LOW-DIMENSIONAL QUANTUM SYSTEMS:
Theory and Experiment
(16 - 27 JULY 2001)

PLUS

PRE-TUTORIAL SESSIONS
(11 - 13 JULY 2001)

TRANSPORT AND DIMENSIONAL CROSSOVER
IN QUASI-ONE DIMENSIONAL CONDUCTORS

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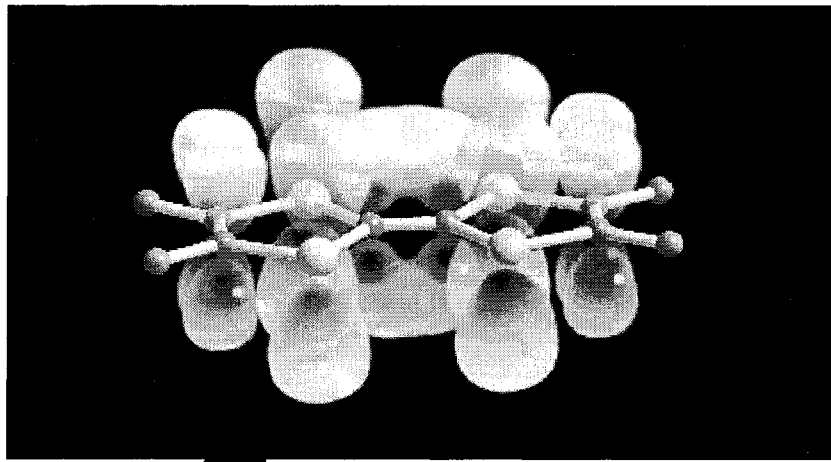
These are preliminary lecture notes, intended only for distribution to participants

Transport and dimensional crossover in quasi-one dimensional conductors

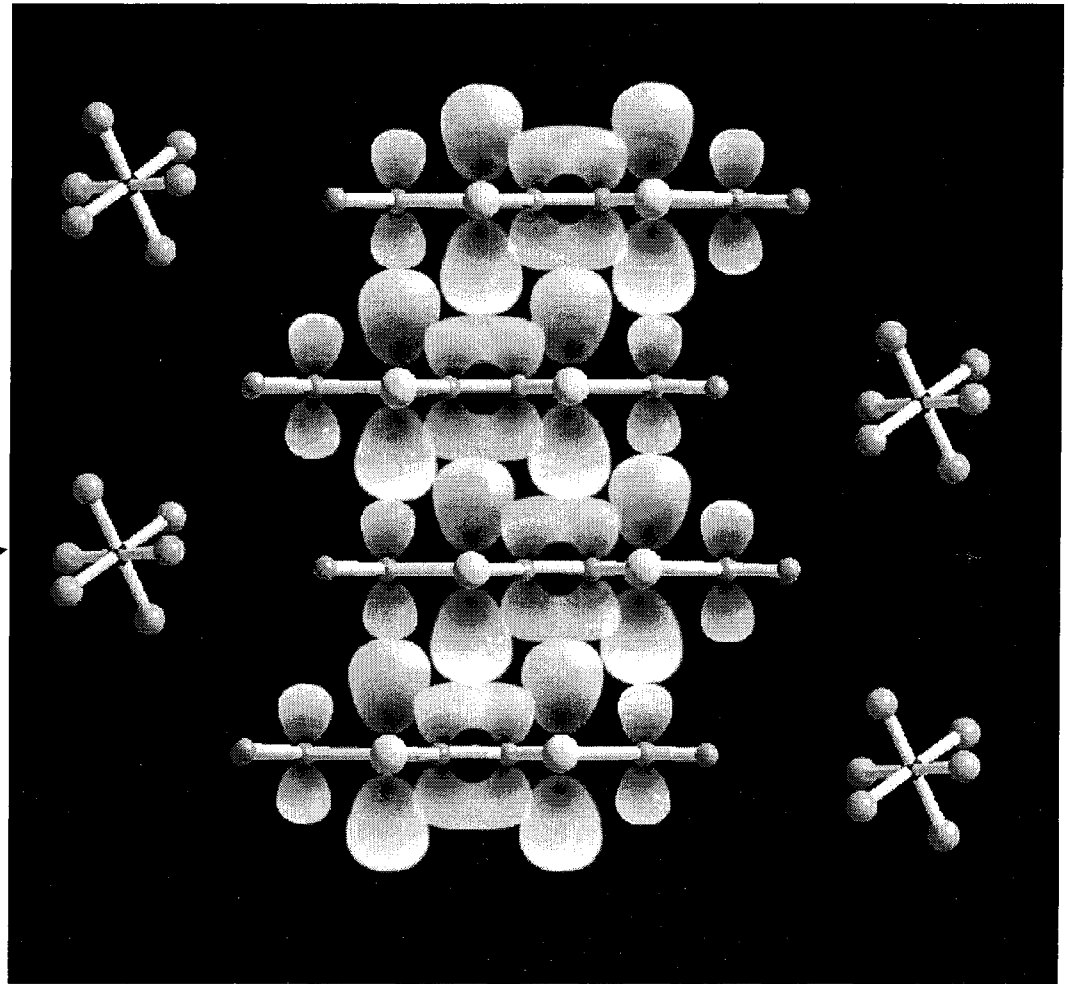
G. Gruner, L. DeGiorgi
A. Schwarz, V. Vescoli, M. Dressel
A. Georges, N. Sandler
S. Biermann, A. Lichtenstein



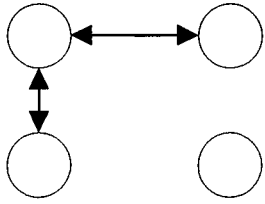
- Organic Superconductors: quasi-one dimensional systems (Bechgaard salts)
- TMTTF and TMTSF molecules
- Remarkable properties :
 - Non Fermi liquid behavior
 - Quantum Hall effect
 - Superconductivity and Frohlich conductivity



TMTSF₂(X)

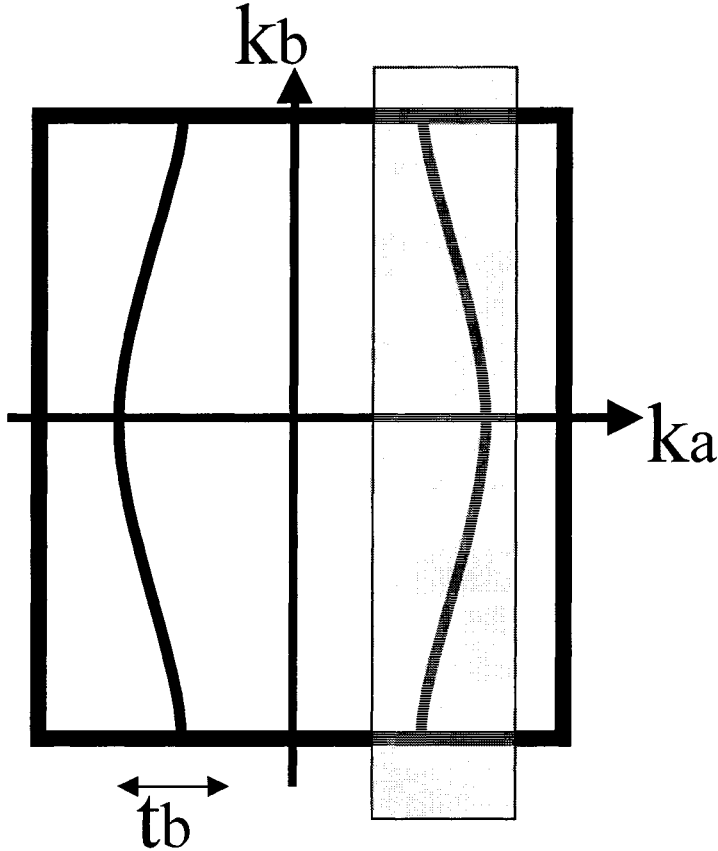


- Propagation of electrons along the chains
- Commensurate filling (1/4 filling)



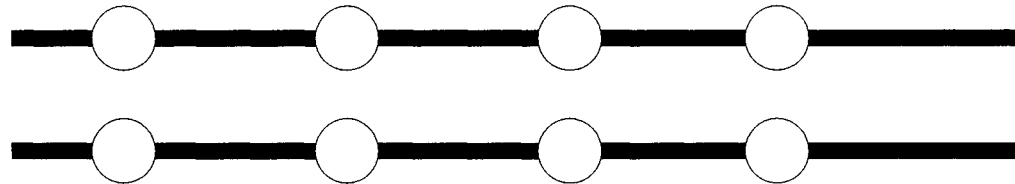
$$t_a > t_b > t_c$$

3000K, 300K, 20K



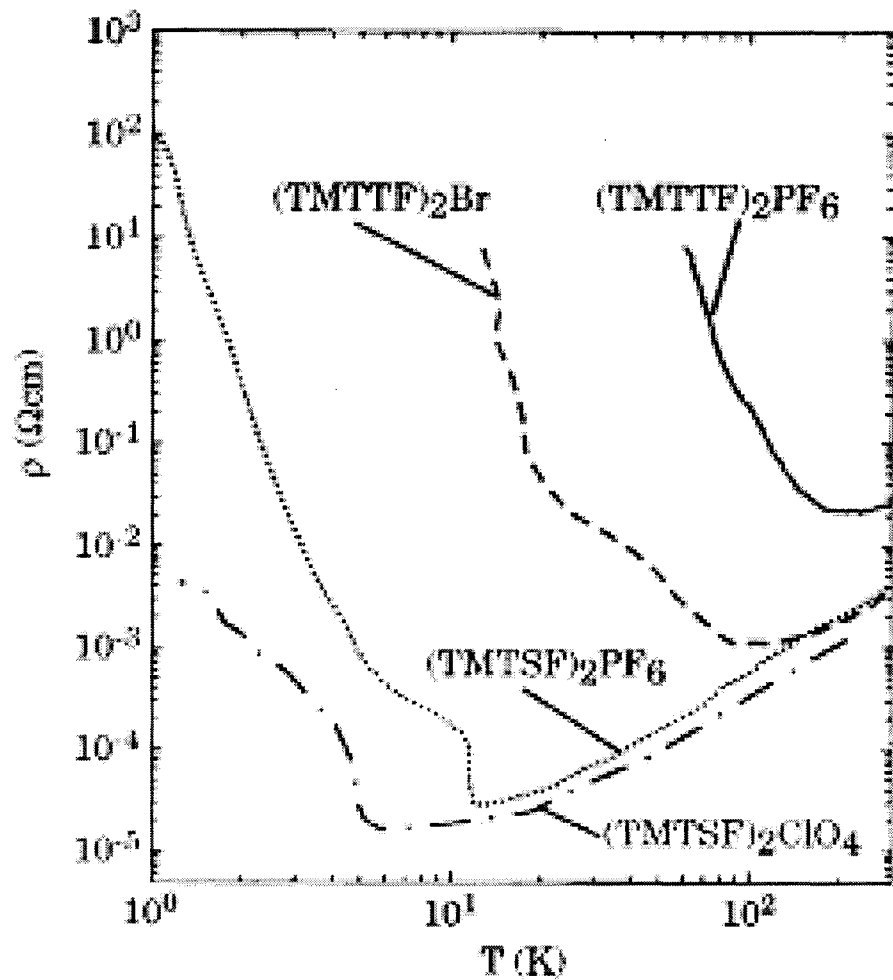
- High Energy (T, ω): 1D
- Low Energy (T, ω): 2D, 3D

Dimensional crossover

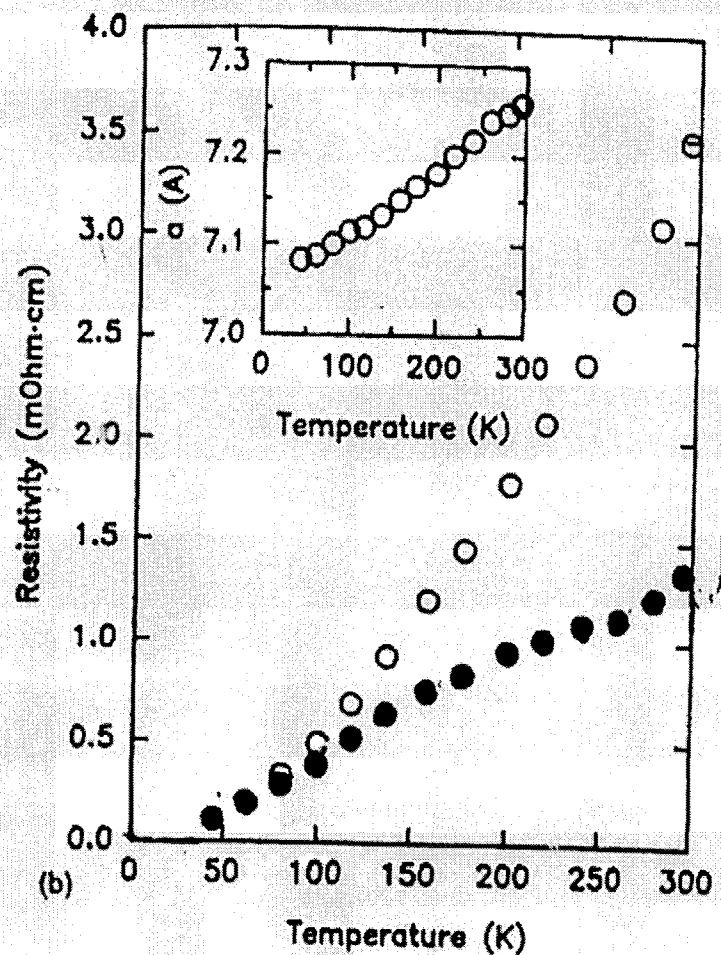


Competition Mott insulator/Interchain hopping

-9-



V. Vescoli et al.
Euro Phys J B 13
503 (2000)



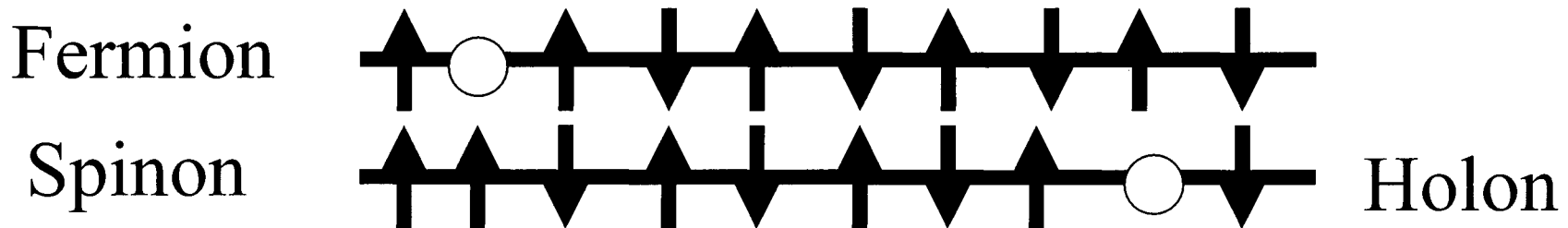
D. Jérôme in "Organic
superconductors", ed. J.P.
Farges

Questions

- Are TMTSF strongly correlated systems ?
- What is the strength of interactions ?
- What causes the difference between TF and SF (weak vs strong dimerization ?) ?
- At what scale does the dimensional crossover takes place ? t_b renormalized by interactions (C. Bourbonnais). 300K vs 20 K !

Luttinger liquid

- Spin charge separation



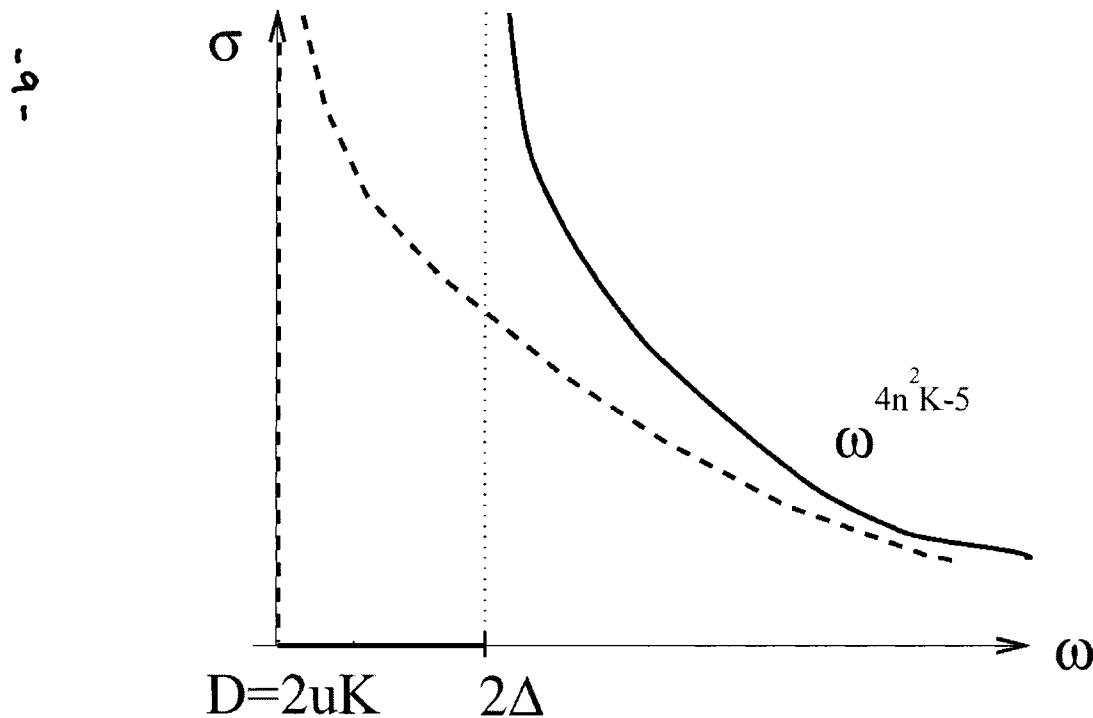
- No fermionic quasiparticles
- Power law decay of correlation functions

$$\langle S(x)S(0) \rangle = \frac{1}{x^2} + \cos(2k_F x) \left(\frac{1}{x}\right)^{1+K_\rho} + \dots$$

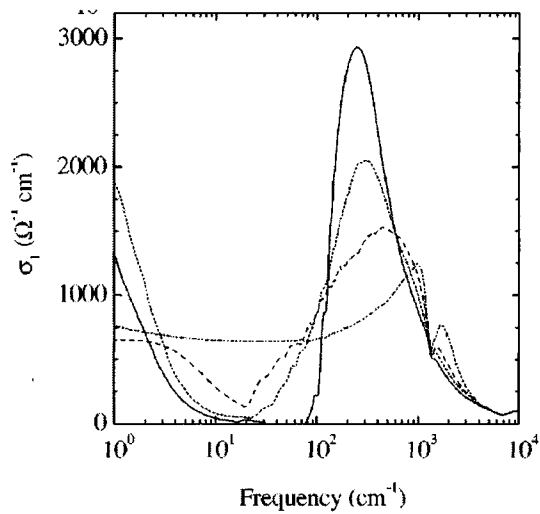
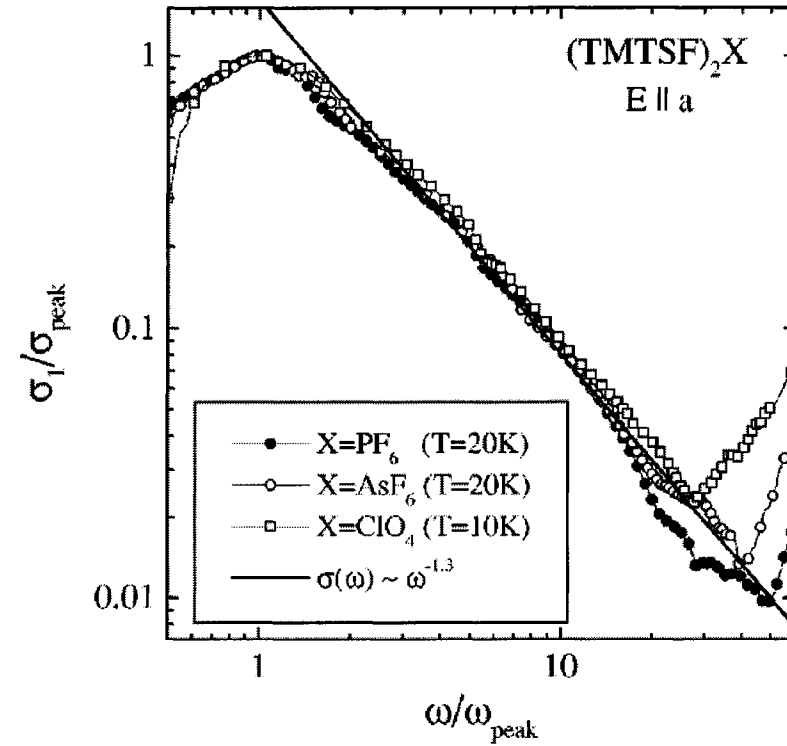
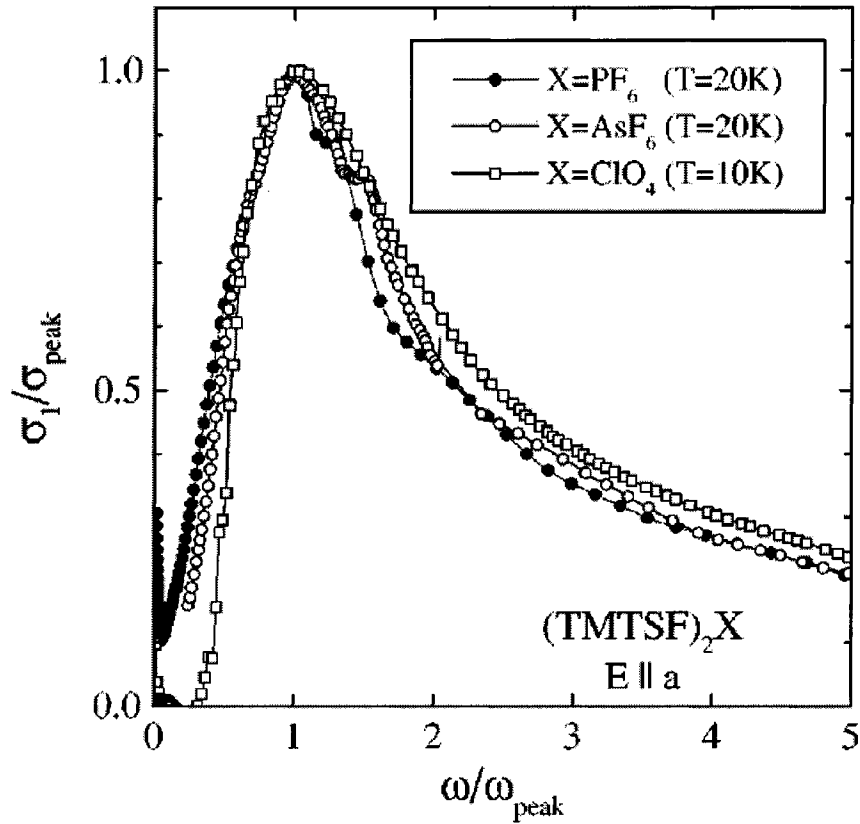
K_ρ contains all information about interactions

Transport in a LL

TG PRB 44 2905 (91); Physica B 230 975 (97)



- Mott insulator for $\nu = 1$ filling also !
- Power law in $\sigma(\omega)$ determines $K\rho$
- Deviations from 1D law gives E_c



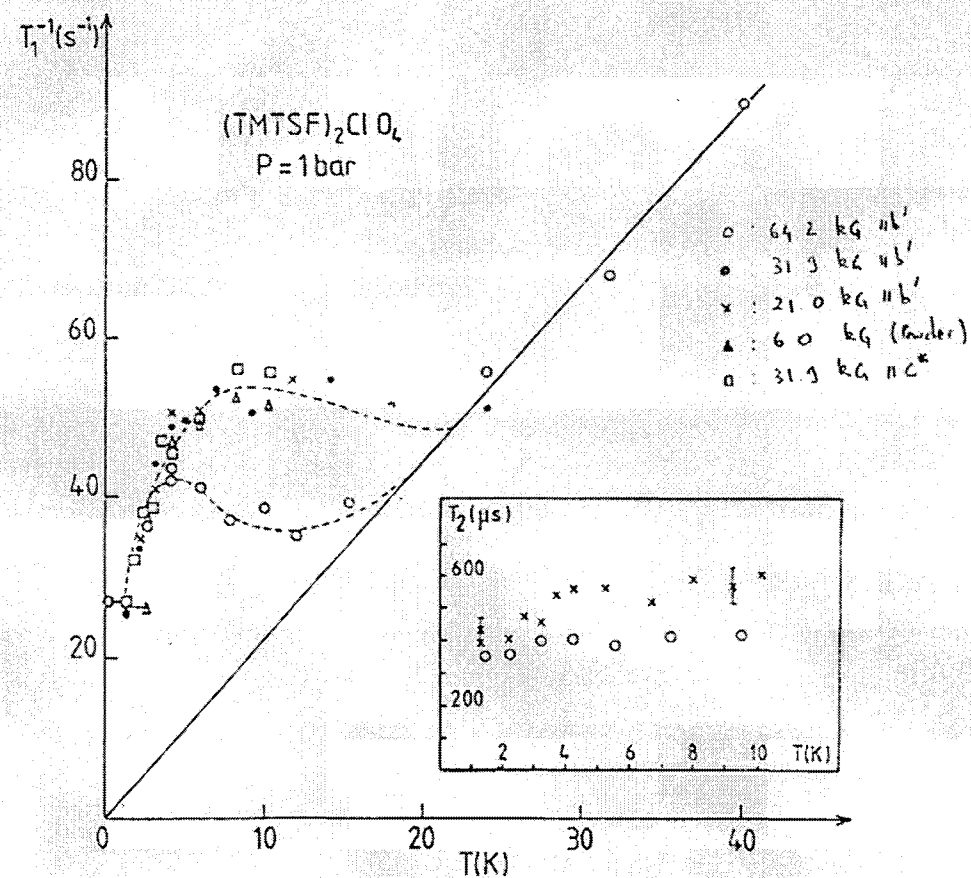
A. Schwartz et al. PRB 58 1261 (1998)

Consequences

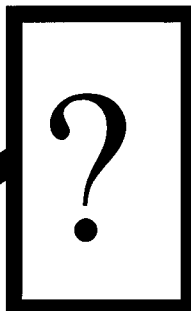
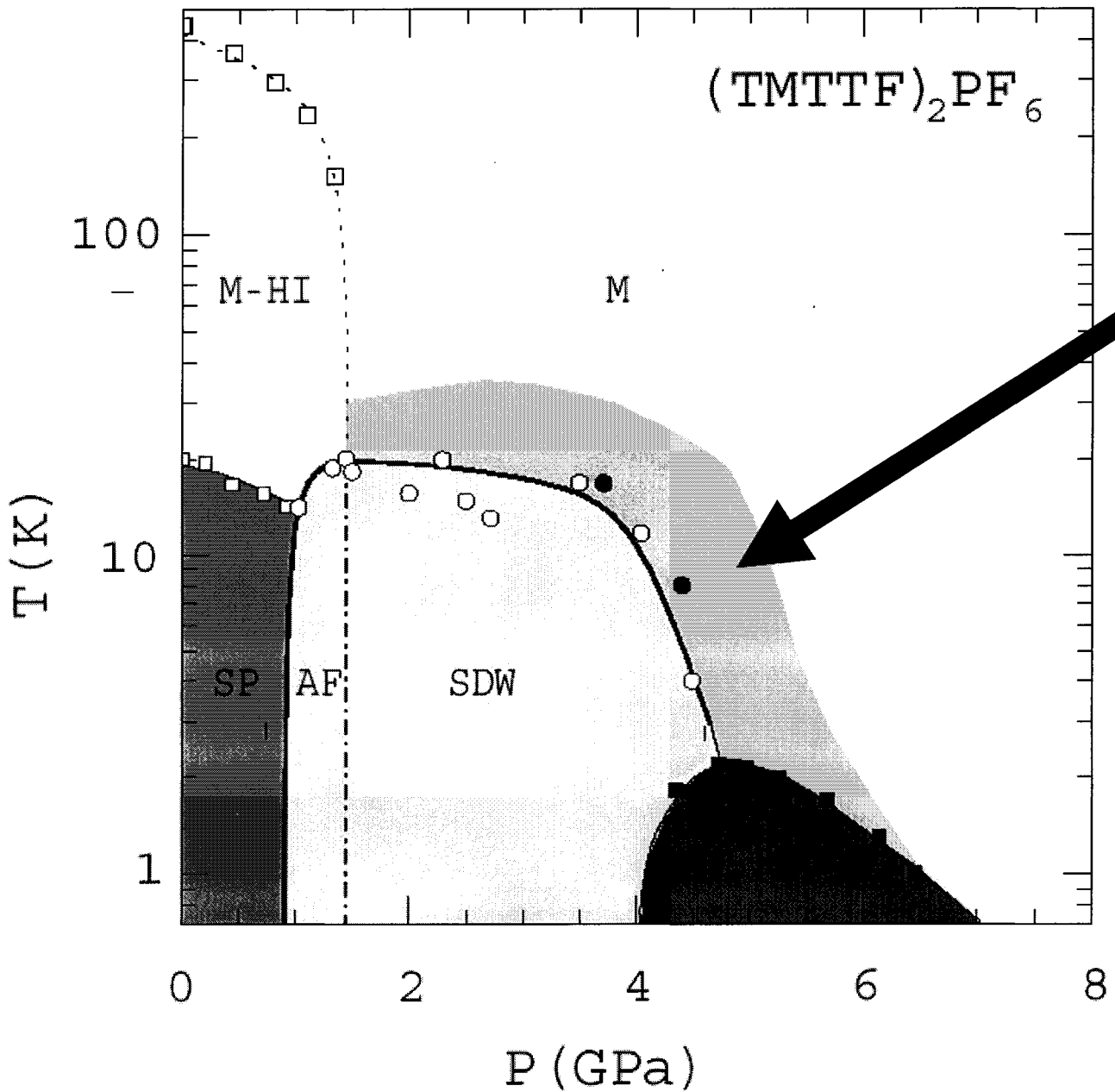
- Luttinger Liquid !
- $K\rho = 0.23$ very strong interactions
- Dimerization not important (1/4 filled Mott insulator) [confirmed in new non-dimerized compound]
- TF vs SF likely to be due to change of interactions (quantum critical point)
- Dimensional crossover at $E=100\text{K}$!

Puzzle from NMR

F. Creuzet et al., J. Phys. Lett 45 L755 (84)



Non fermi liquid
below E_c ?

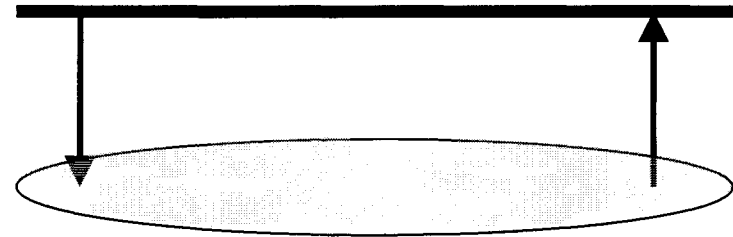


After D. Jérôme

- 3/ -

Dimensional Crossover

A. Georges, TG, N. Sandler PRB 61 16393 (00)



- $d = \infty + 1$ (cDMFT)
- Effective one dimensional theory

$$\begin{aligned}
S_{\text{eff}} = & - \int \int_0^\beta d\tau d\tau' \sum_{ij,\sigma} c_{i\sigma}^+(\tau) \mathcal{G}_0^{-1}(i-j, \tau - \tau') c_{j\sigma}(\tau') \\
& + \int_0^\beta d\tau H_{\text{ID}}^{\text{int}}[\{c_{i\sigma}, c_{i\sigma}^+\}], \tag{4}
\end{aligned}$$

$$G(k, i\omega_n) = \int d\epsilon_\perp \frac{D(\epsilon_\perp)}{i\omega_n + \mu - \epsilon_k - \Sigma(i\omega_n, k) - \epsilon_\perp}.$$

$$\begin{aligned}
\text{Re} \sigma_\perp(\omega, T) \propto t_\perp^2 \int d\epsilon_\perp D(\epsilon_\perp) \int \frac{dk}{2\pi} \int d\omega' A(\epsilon_\perp, k, \omega') \\
\times A(\epsilon_\perp, k, \omega + \omega') \frac{f(\omega') - f(\omega' + \omega)}{\omega}.
\end{aligned}$$

- Self consistent theory for Σ
- Feedback of t_{\perp} in Σ (a priori important for deconfinement)
- Different from RPA

$$G(k, k_{\perp}, i\omega_n) = \frac{1}{i\omega_n - \varepsilon_k - \varepsilon_{\perp}(k_{\perp}) - \Sigma_{1D}(k, i\omega_n)}$$

- Difficult to solve the equations analytically

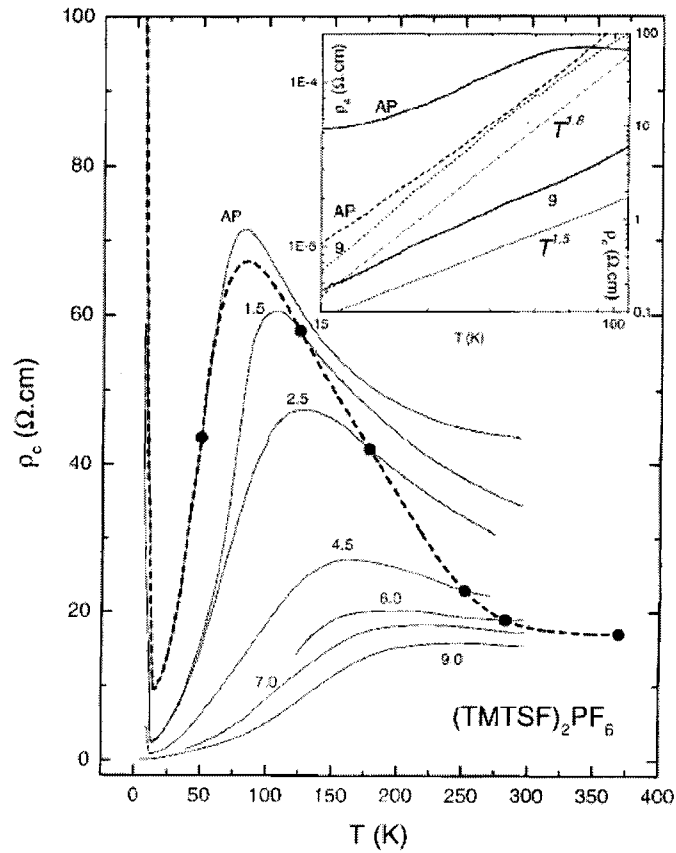
Transverse transport

$$\sigma(\omega, T) \propto (\omega, T)^{2\alpha-1}$$

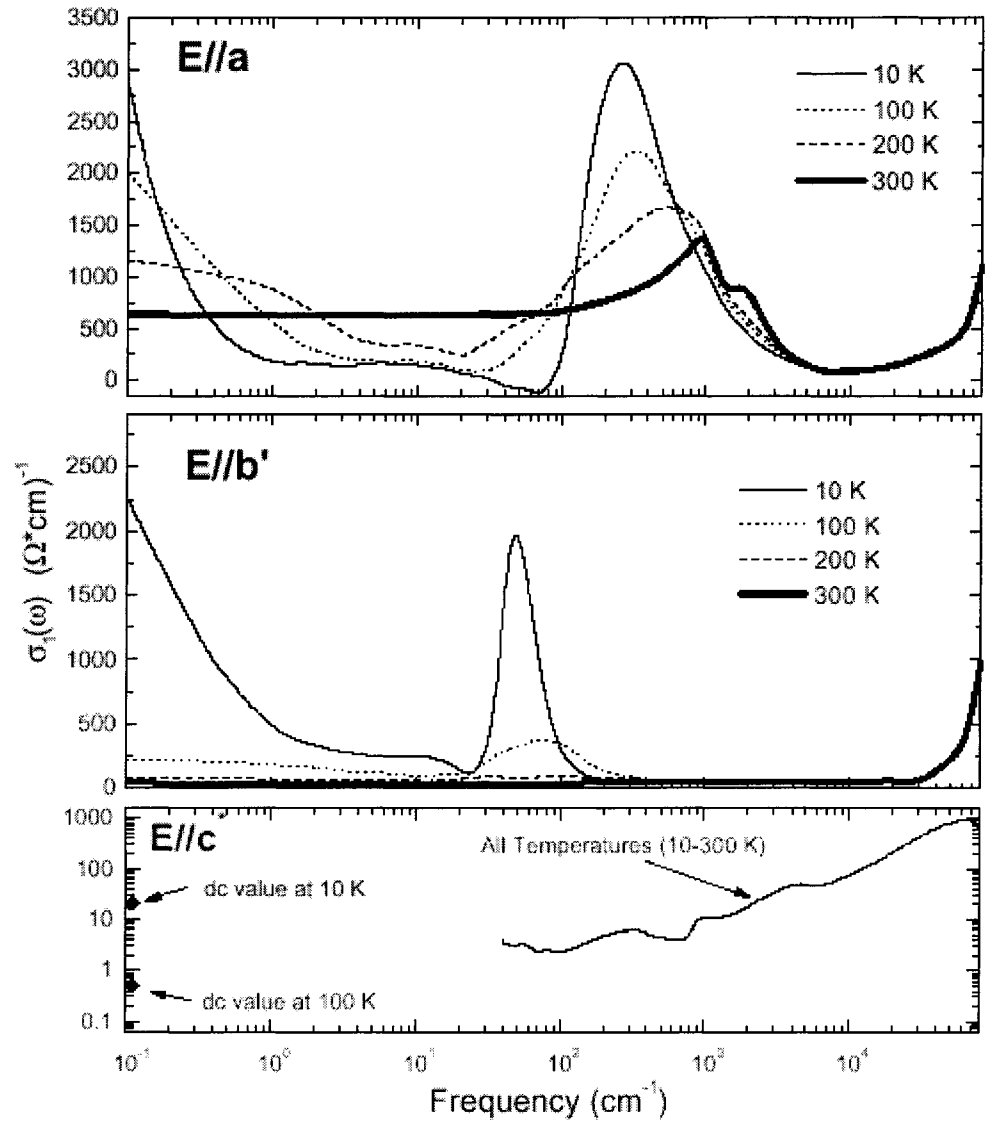
$$\alpha = \frac{1}{4}(K + K^{-1}) - \frac{1}{2}$$

$$\alpha \simeq 0.6$$

$$\omega_{D\perp}^2 / \omega_{P\perp}^2 \propto (t_{\perp} / t)^{2\alpha/(1-\alpha)} = Z^2$$



J. Moser et al. Euro Phys. J. B 1 39 (1998)



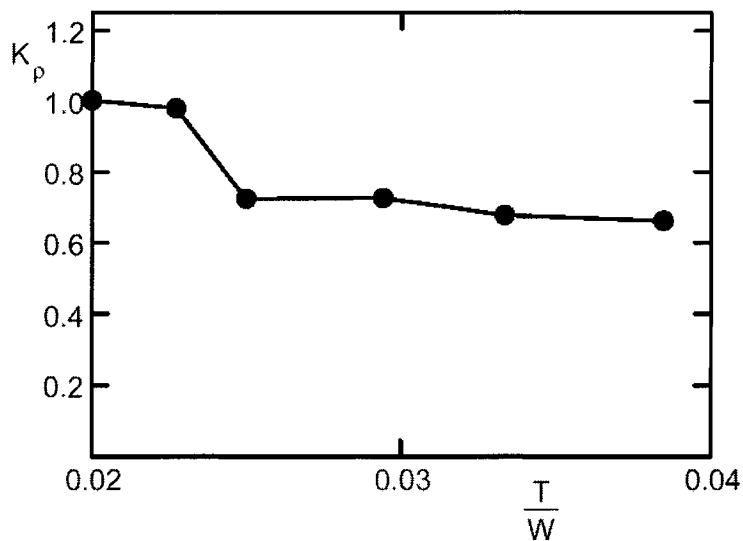
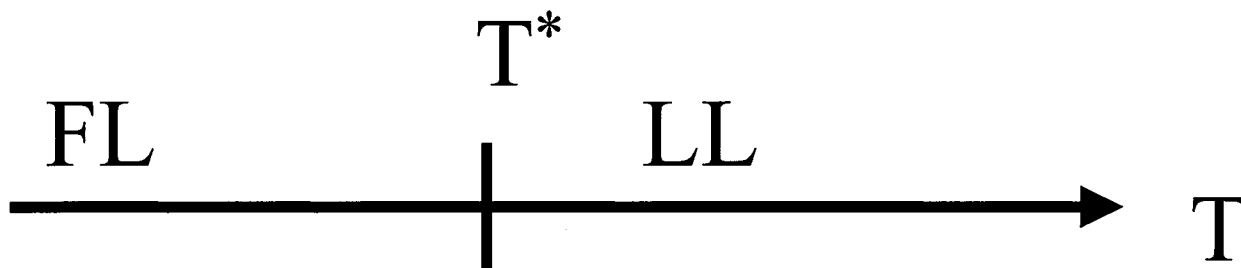
V. Vescoli et al. Euro Phys J B 11 365 (1999)

Numerical Solution

S. Biermann, A. Georges, A. Lichtenstein, TG

- Hubbard Model
- QMC (16 – 32 sites)
- 32 time slices ($T/W = 1/50$)

Incommensurate case



$$T^* \sim \frac{t_\perp}{\pi} \frac{t_\perp}{t} \sqrt{\frac{\theta}{1-\theta}}$$

$$T^* \sim 0.5 \frac{t_\perp}{\pi}$$

FIG. 2: Effective K_ρ vs. temperature in the doped case (filling $n \simeq 0.8$) for $U/W = 1.0$, $t_\perp/W = 0.14$.

~20~

Commensurate case

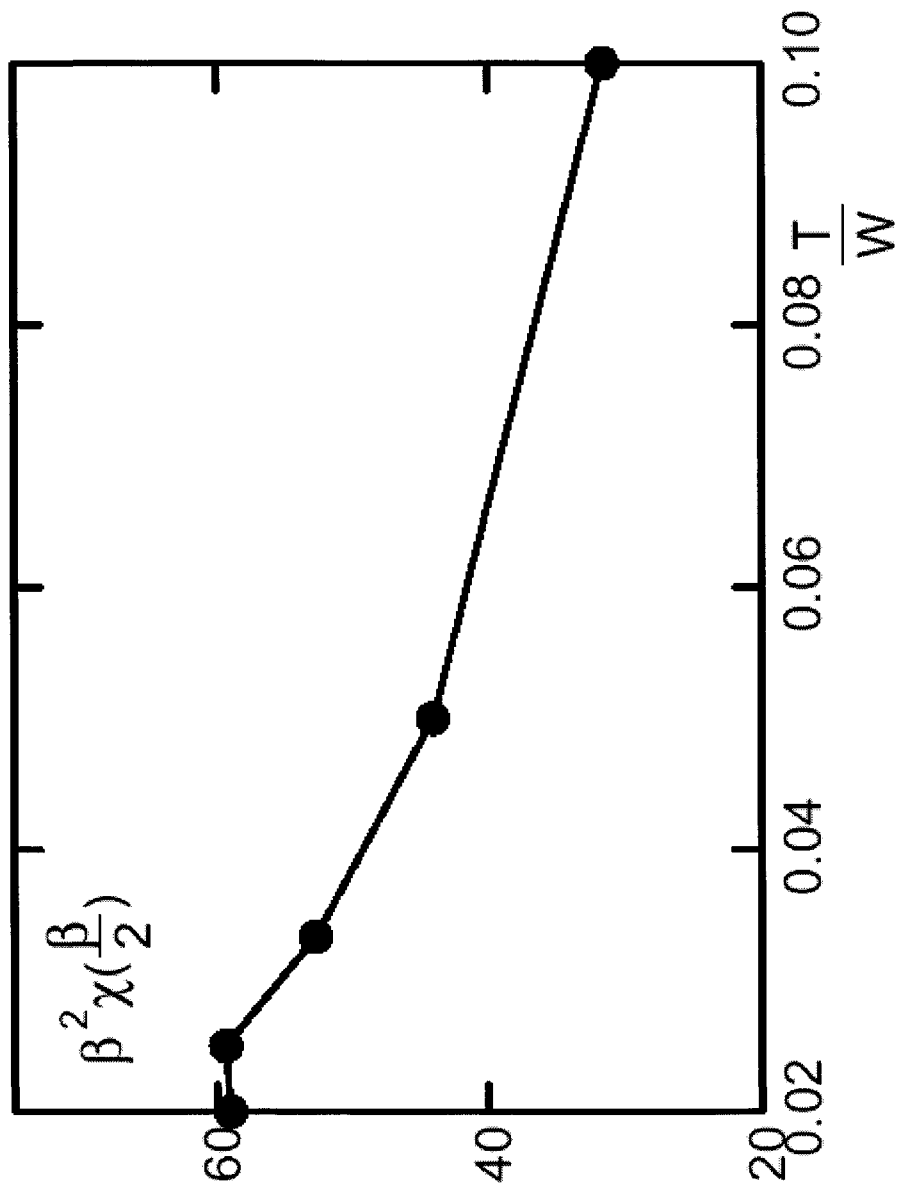


TABLE II: Effective K_{ρ} at half-filling, as a function of t_{\perp}/W for $U/W = 0.65$ and $T/W = 1/40$.

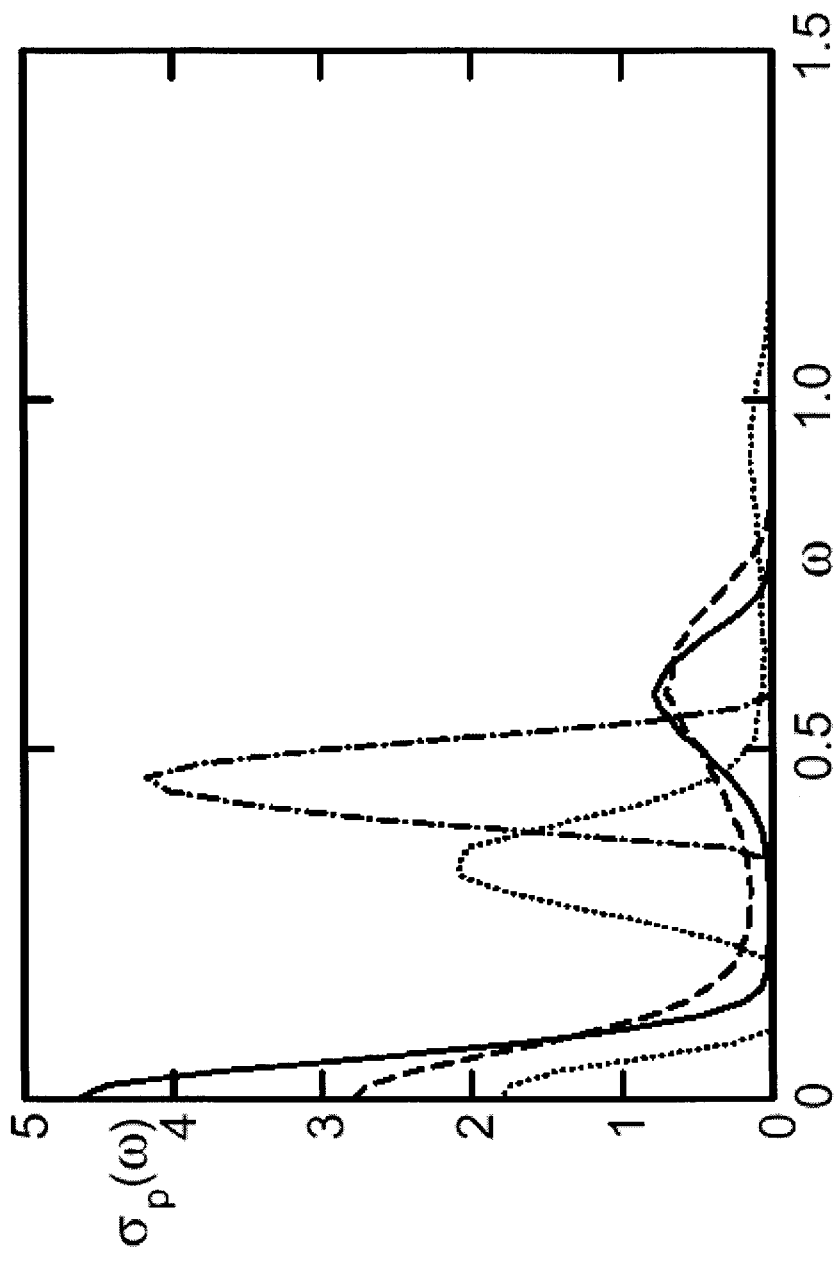
t_{\perp}/W	0.00	0.04	0.07	0.11	0.14	0.16	0.18
K_{ρ}	0.00	0.02	1.01	1.09	1.07	1.06	1.04

$$t_{\perp}^* \cup \Delta_{1D}$$

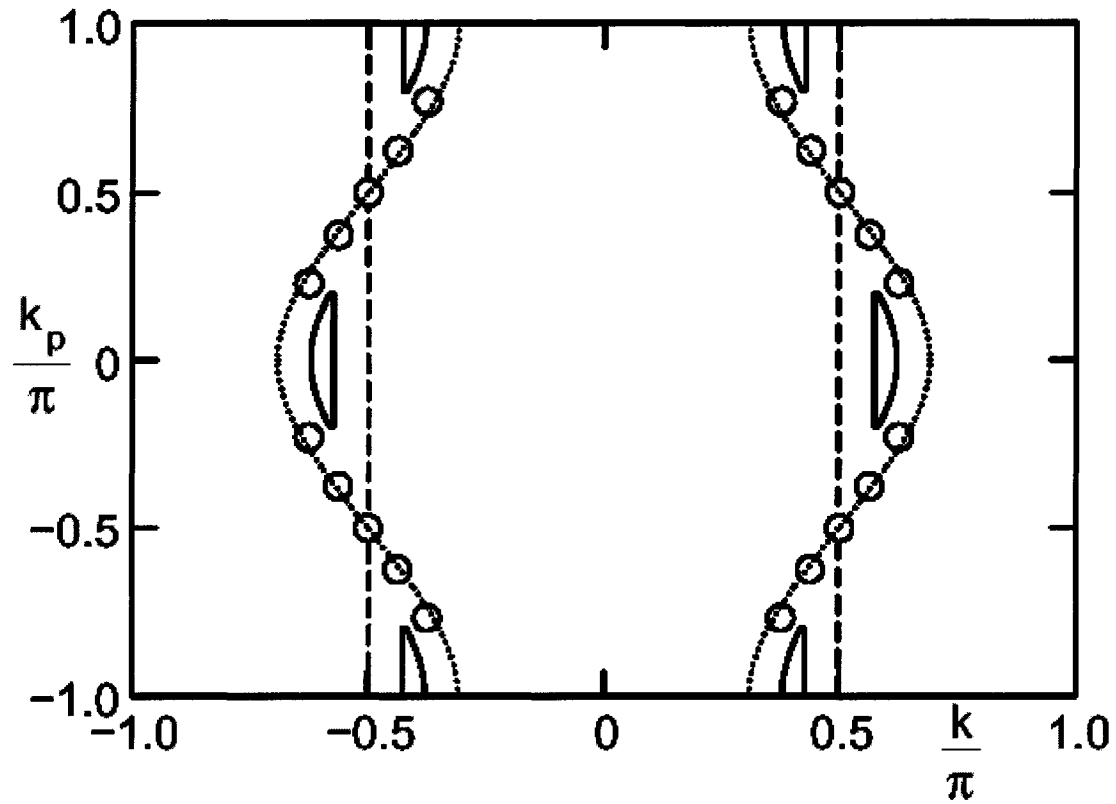
NMR



Transverse conductivity



Fermi Surface



-410-

Z

k_{\perp}/π	0.23	0.38	0.50	0.62	0.77
$Z(k_{\perp})$	0.78	0.77	0.76	0.77	0.79

Conclusions

- Transport proves LL nature of high energy phase
- $K=0.23$, ν filled Mott insulators
- $E_c = 100K$ much higher than expected
- Nature of the 2D phase ? 2D non Fermi liquid ?
- cDMFT Good method to tackle the dimensional crossover



- Filling with cDMFT
- Other physical quantities : Hall effect
- Ordered phases (superconductivity)