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INTERNATIONAL ATOMIC ENERGY AGENCY
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H4.SMR/1001-6

**IX TRIESTE WORKSHOP ON
OPEN PROBLEMS IN
STRONGLY CORRELATED SYSTEMS**

14 - 25 July 1997

**OUT-OF-PLANE CHARGE TRANSPORT IN
HIGH- T_c CUPRATES**

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These are preliminary lecture notes, intended only for distribution to participants.

Out-of-plane charge transport in high- T_c cuprates

- *distinct ρ_c behavior in underdoped Y124*

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collaborators

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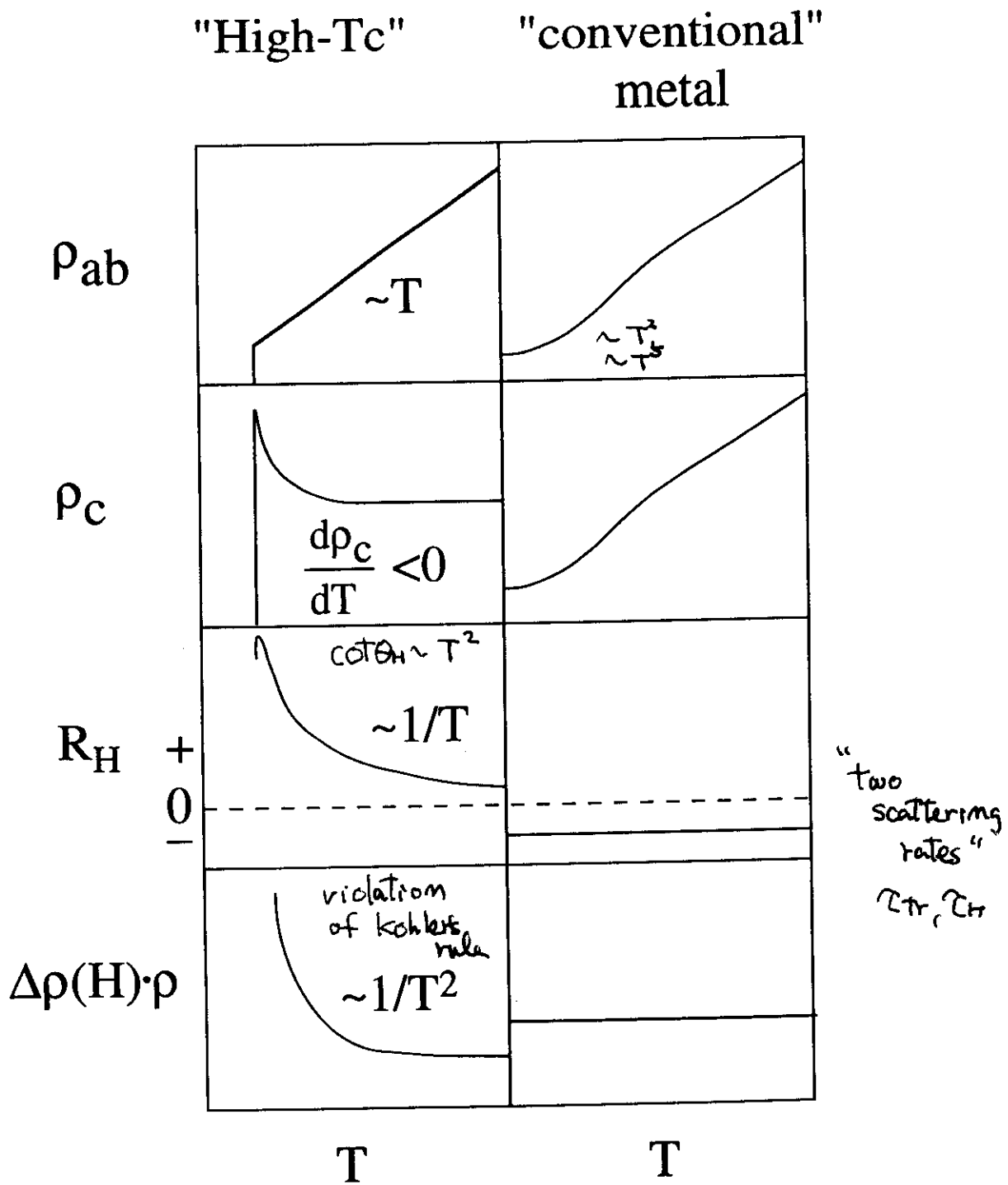
ISSP

S.Adachi
K.Tanabe

SRL, ISTECS

1. overview on out-of-plane transport
2. Y124, underdoped cuprate with spin gap
3. anisotropic 3D transport in Y124
4. field-induced dimensionality control

Hallmarks of Anomalous Charge Transport in High-Tc Cuprates.



Completely different conduction mechanism along c-axis

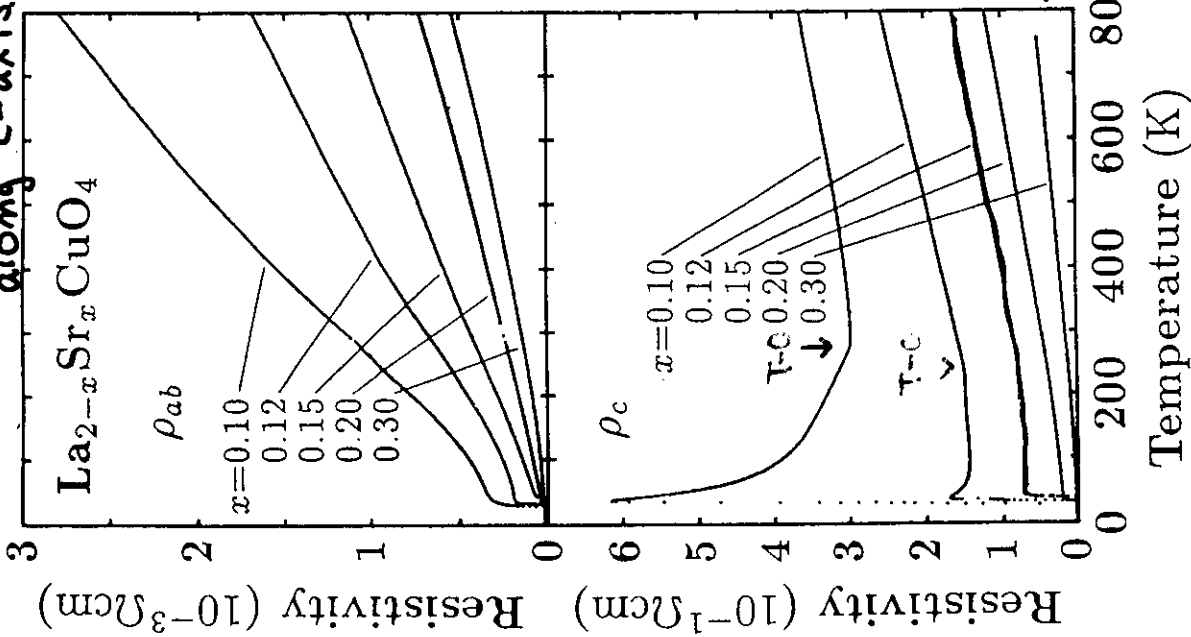


Fig. 2

2D-3D

"crossover"

disappearance of SC

1 Temp dep.

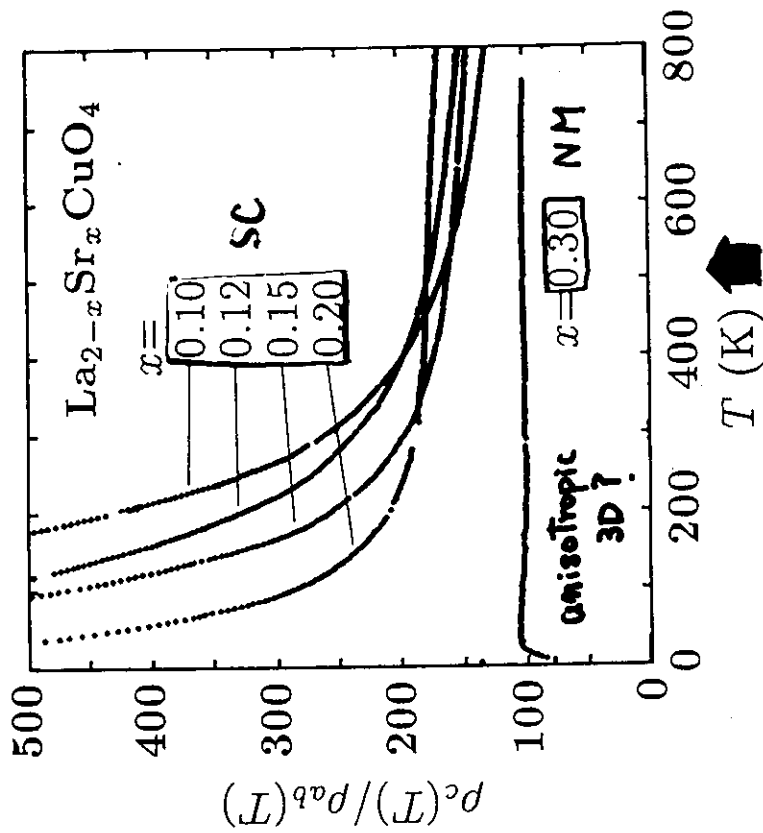
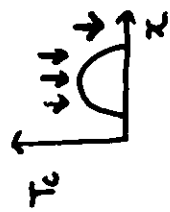
completely different

from ρ_{ab}

2 undeveloped

$\rho_c \rightarrow \infty$

$T \rightarrow 0$



Set, RH anomaly

"anomaly" initial

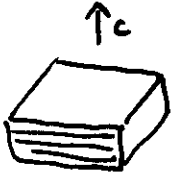
Fig. 3

No correlation between ρ_c and T_c

ρ_c strongly

materials dependent

no correlation with T_c



$$\frac{\rho_c}{\rho_{ab}} = 10^2 \sim 10^6$$

$\downarrow$
 $\downarrow$

band calc.

charge are
strongly
confined
within the plane

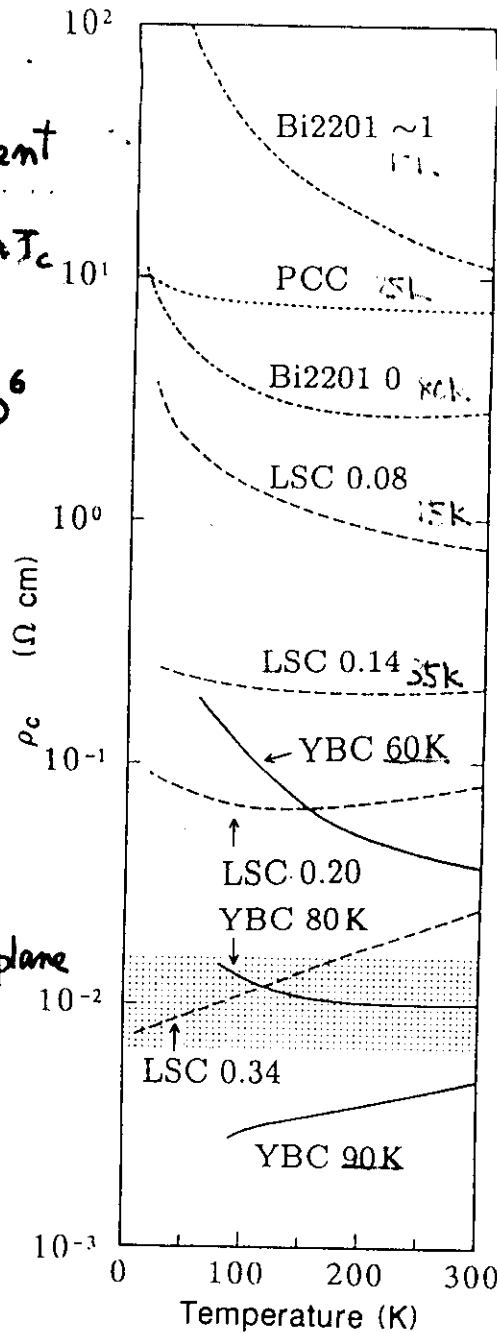


FIG. 2 Temperature dependences of ρ_c for $\text{YBa}_2\text{Cu}_3\text{O}_{7-y}$, $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$, $\text{Pr}_{2-x}\text{Ce}_x\text{CuO}_{4-y}$ ($x \approx 0.15$) and $\text{Bi}_2\text{Sr}_{2-x}\text{La}_x\text{CuO}_6$, abbreviated to YBC, LSC, PCC and Bi2201, respectively. The numbers following LSC and Bi2201 indicate the Sr and La composition x , respectively (the value of $x = 0.14$ for LSC is from ref. 6). The shaded region indicates the critical resistivity along the c axis, which is the extended Mott-Ioffe-Regel limit estimated from the values given in the text.

Recovery of coherent transport;
in "overdoped" normal metal

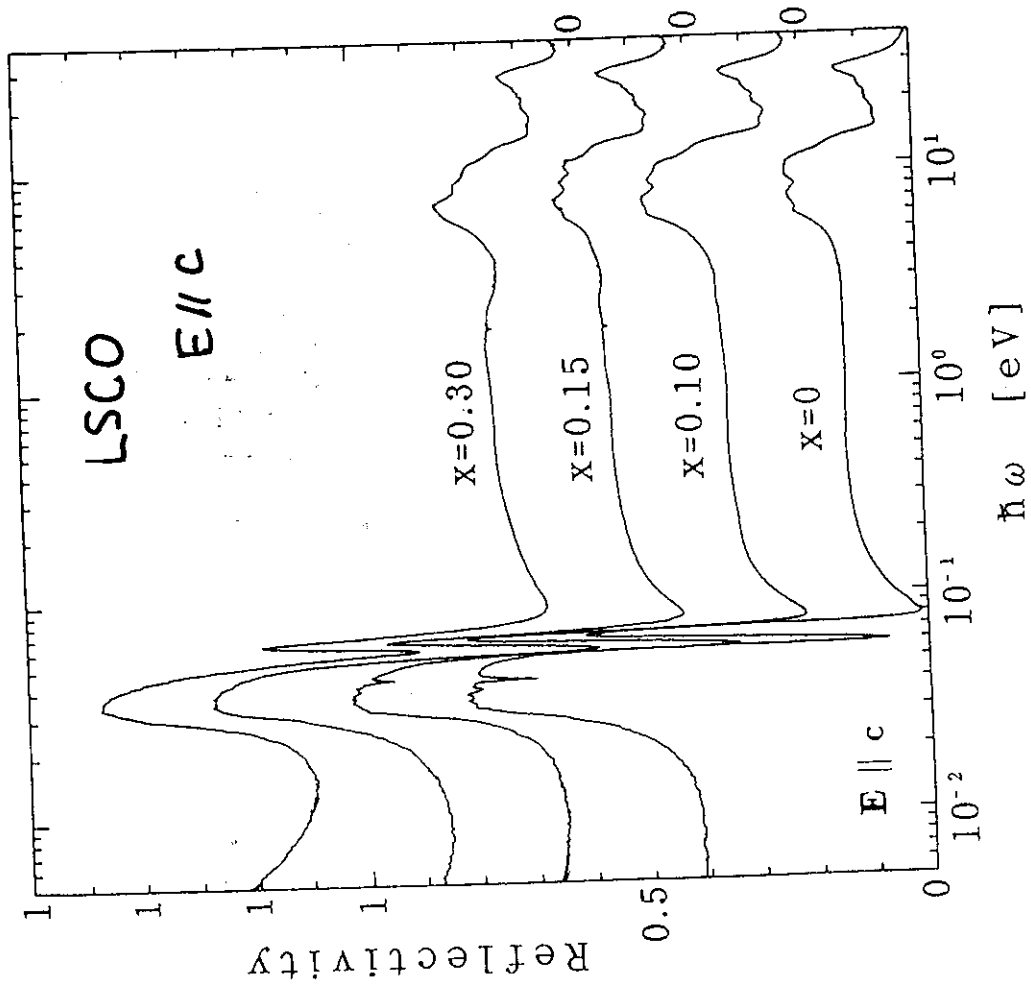
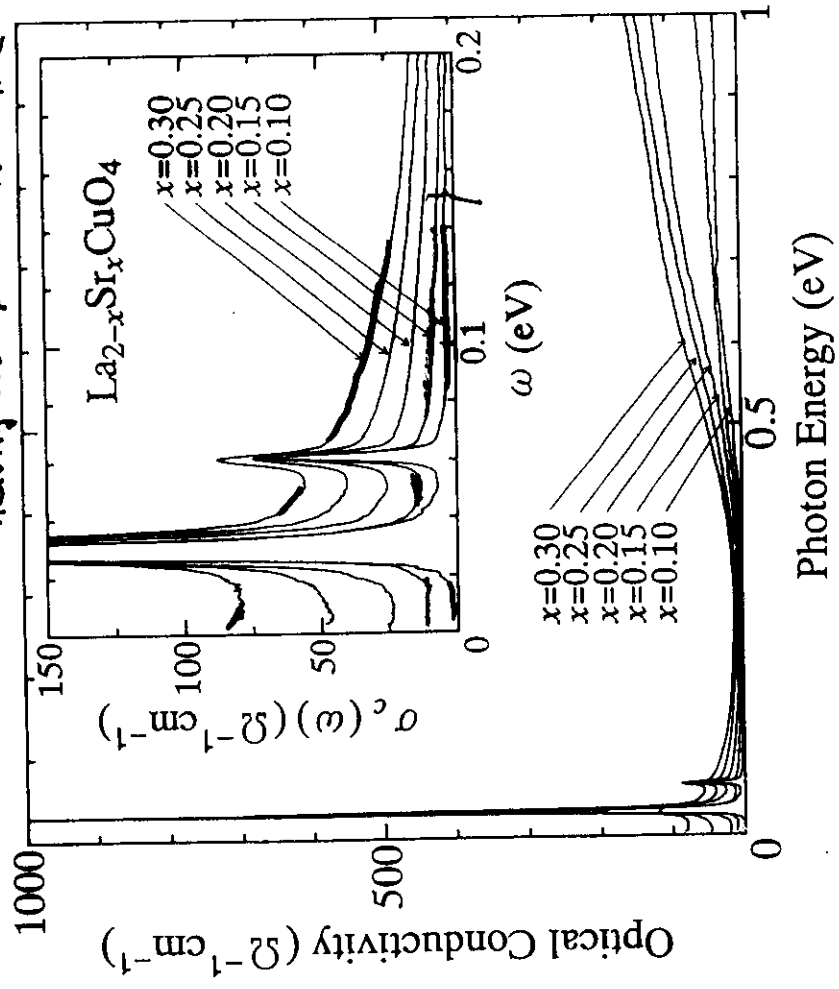


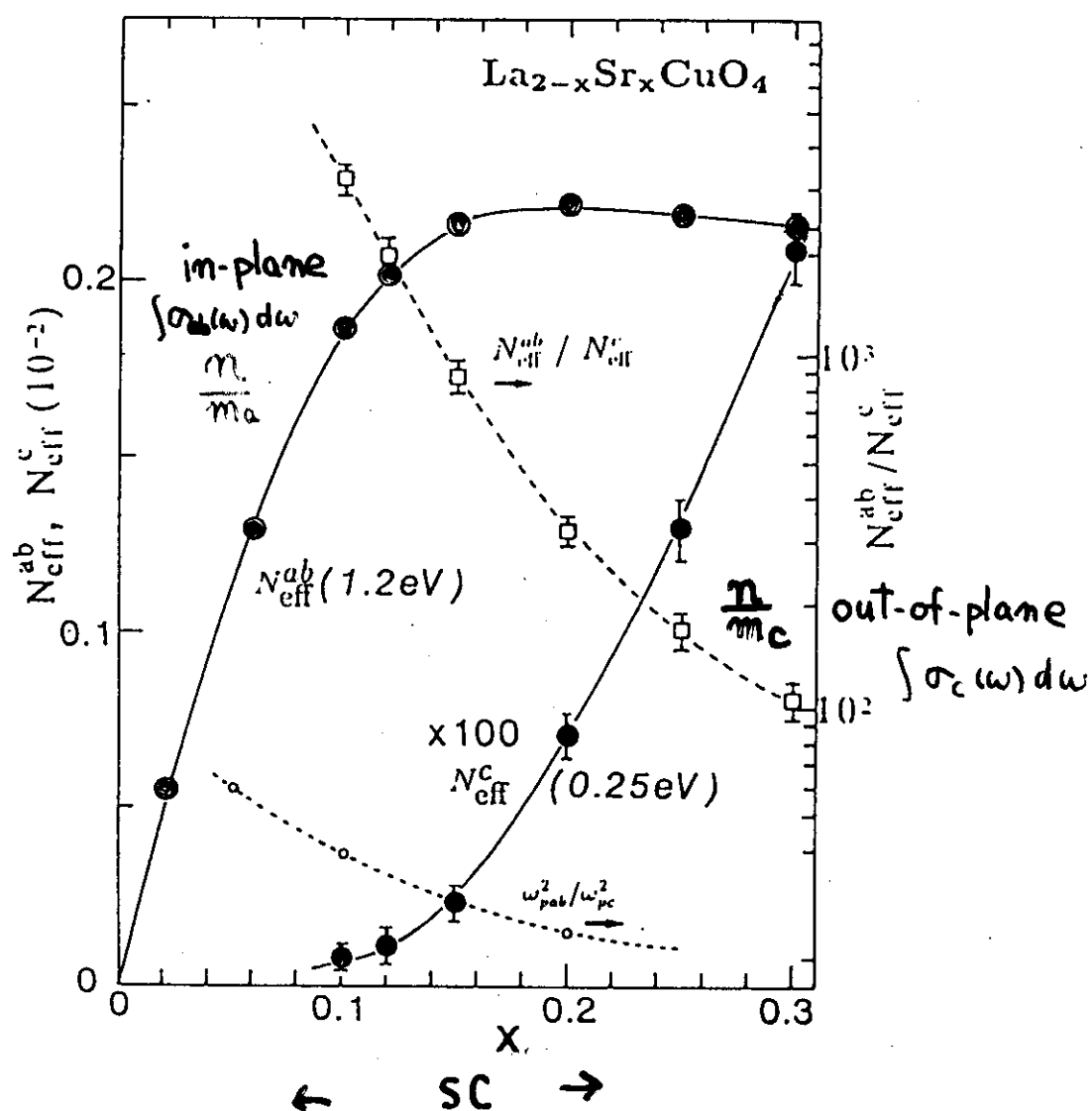
Fig. 1 K. Tanasaku et al.

under - optimum $\sigma(\omega)$, w-indip, incoherent, heavily overdoped Drude-like



energy scale ~ 10⁻²

Strong suppression of out of plane spectral weight

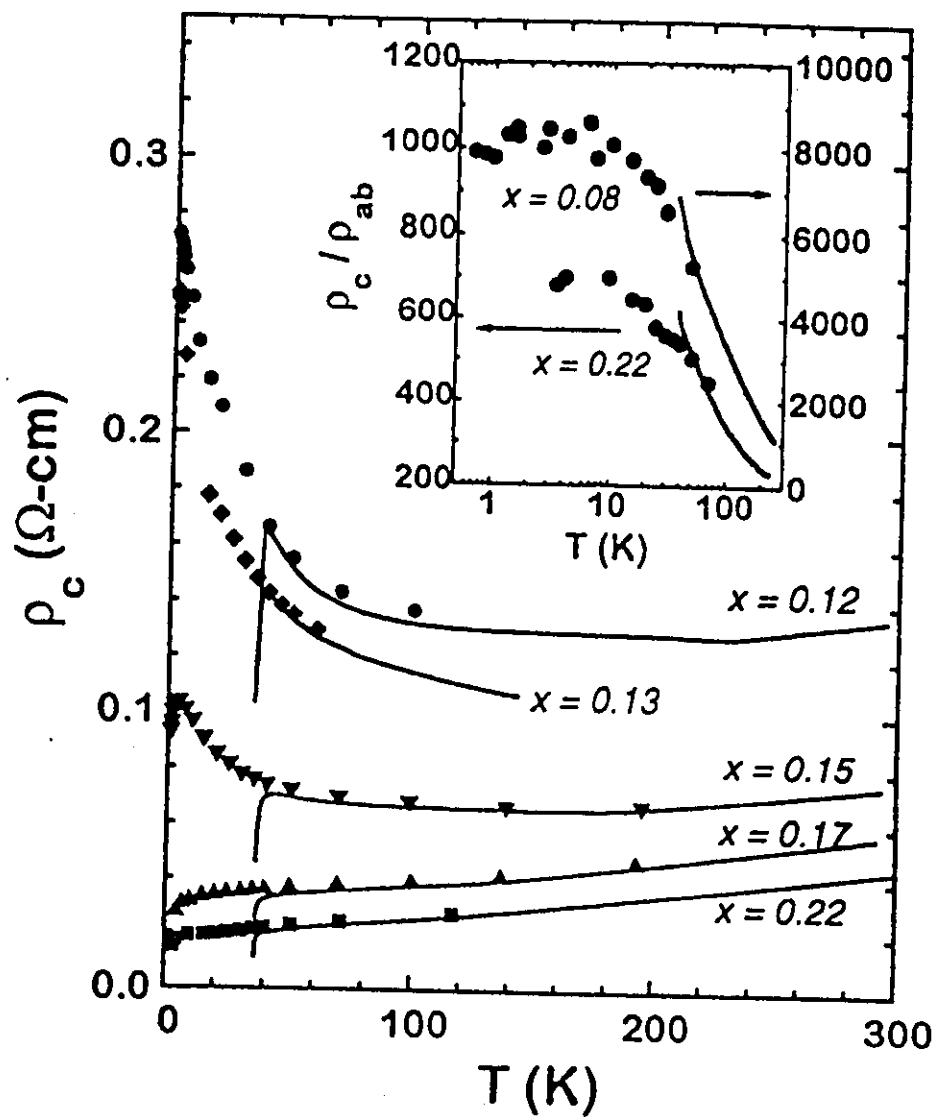


$T \rightarrow 0$ limit

Boebinger

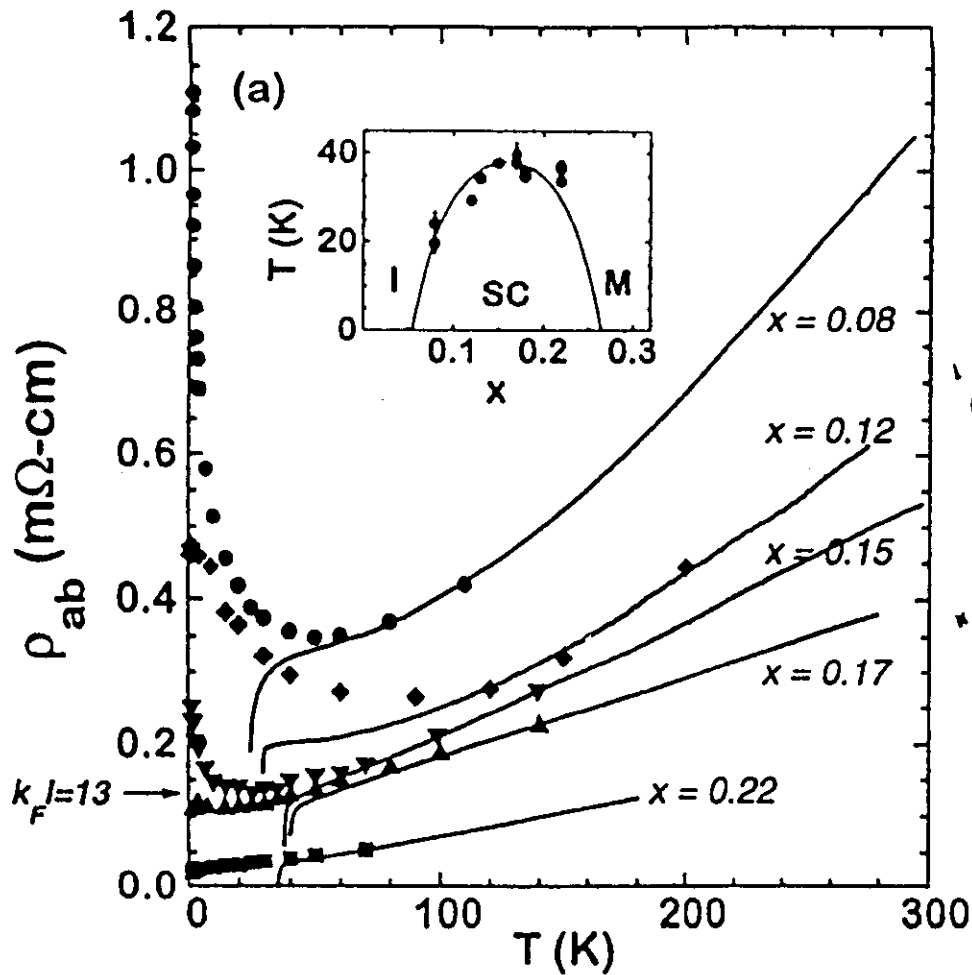
underdoped $\rho_c \rightarrow \infty$ ($\propto -\ln T$)

overdoped $\rho_c \rightarrow \text{const}$

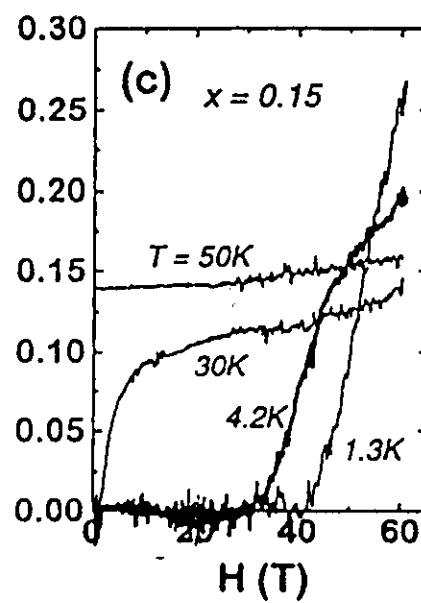
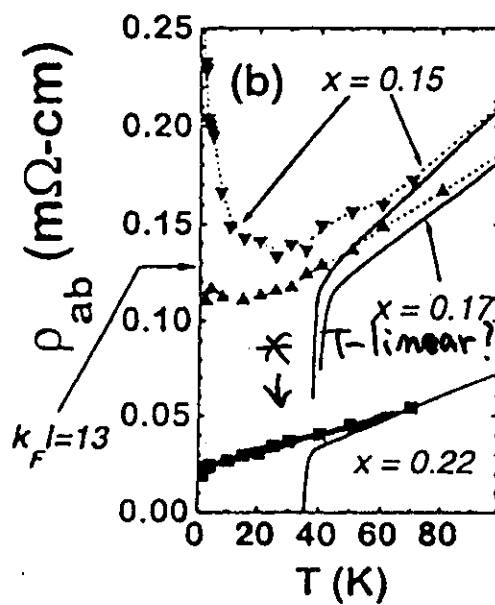


ρ_c divergence

accompanied with $\log T$ divergence of ρ_a



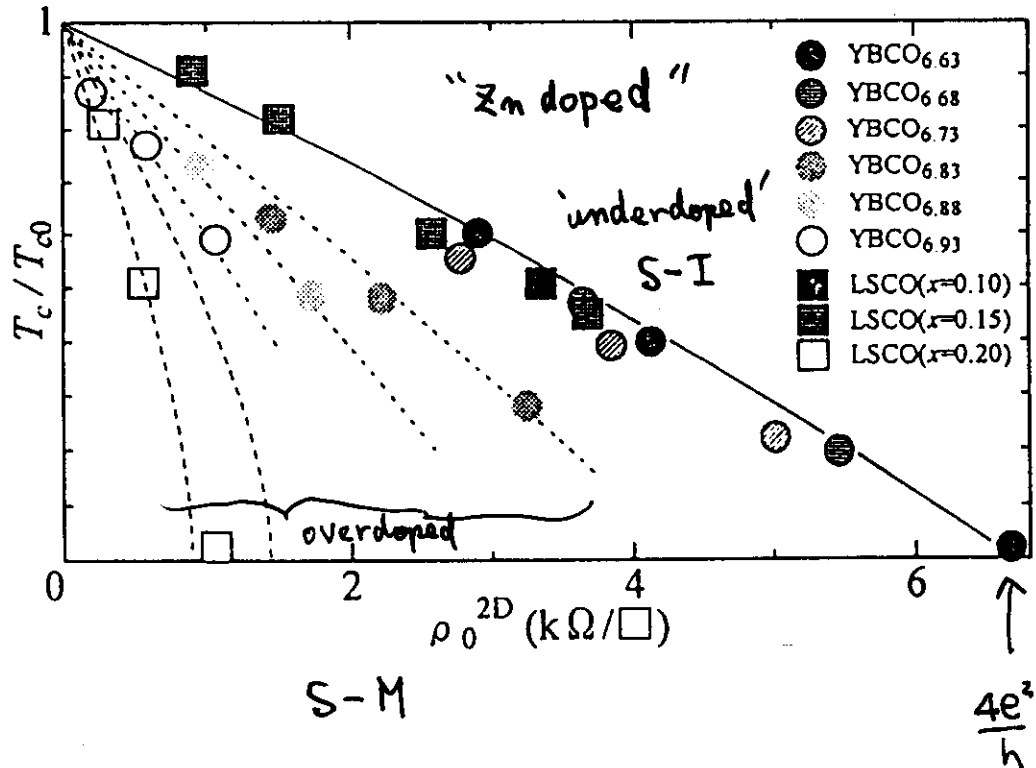
'underdoped'
 $\rho \propto \log T$
 at low T
 * $1/8$ anomaly

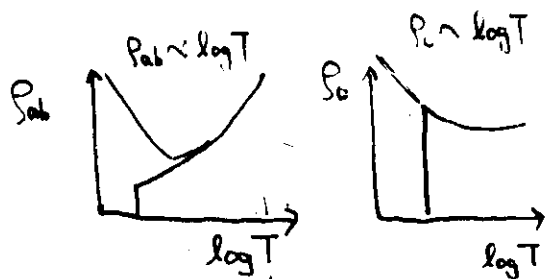


2D S-I transition (disorder induced)

$$R_0 = \frac{4e^2}{h} \sim 6.7 \text{ k}\Omega$$

$$\rightarrow \rho_{3d} = R_0 \cdot d \sim 1.6 \text{ m}\Omega \text{ cm LSCO}$$

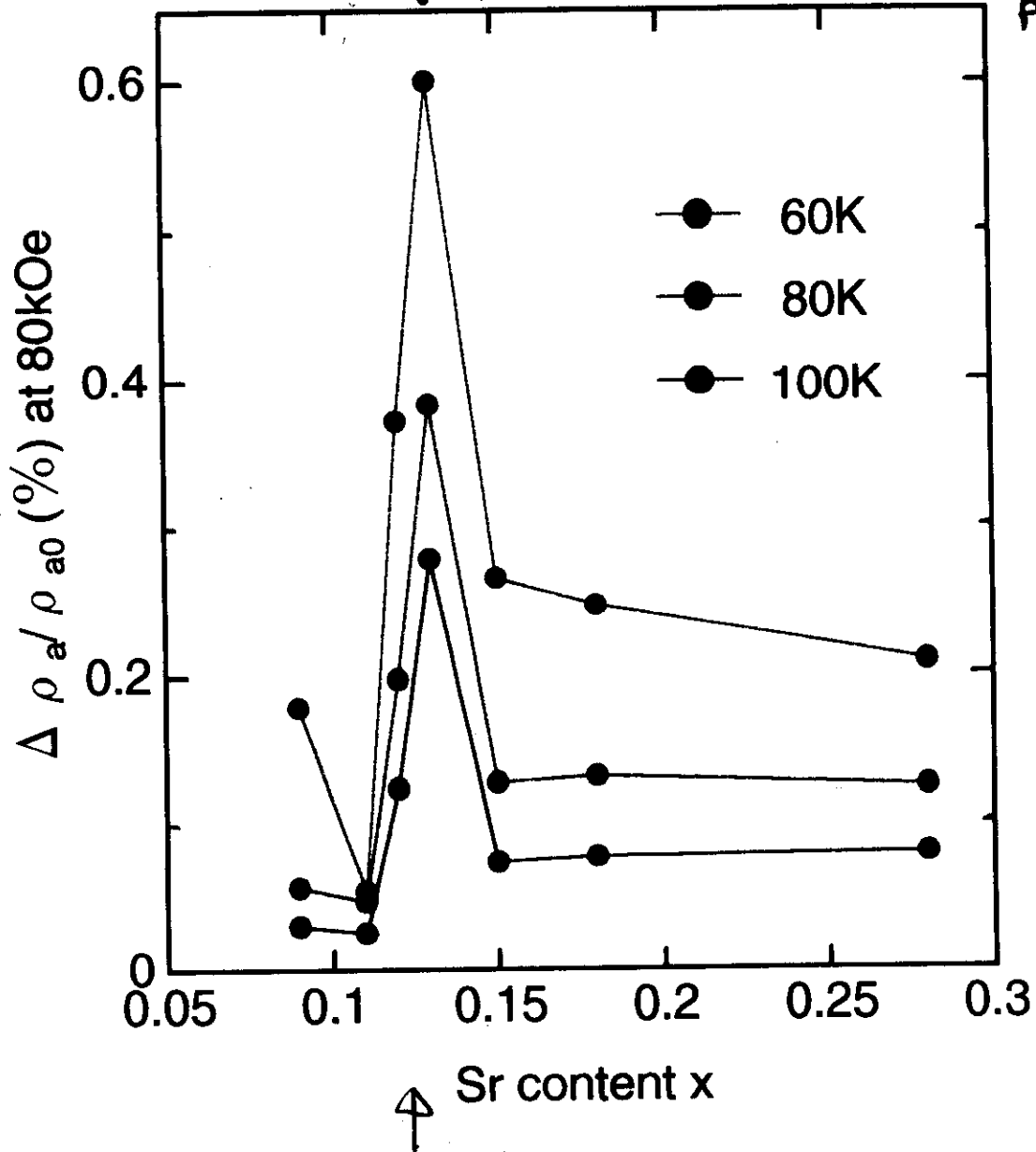




Im-plane MR
pronounced anomaly
around " $1/8$ "

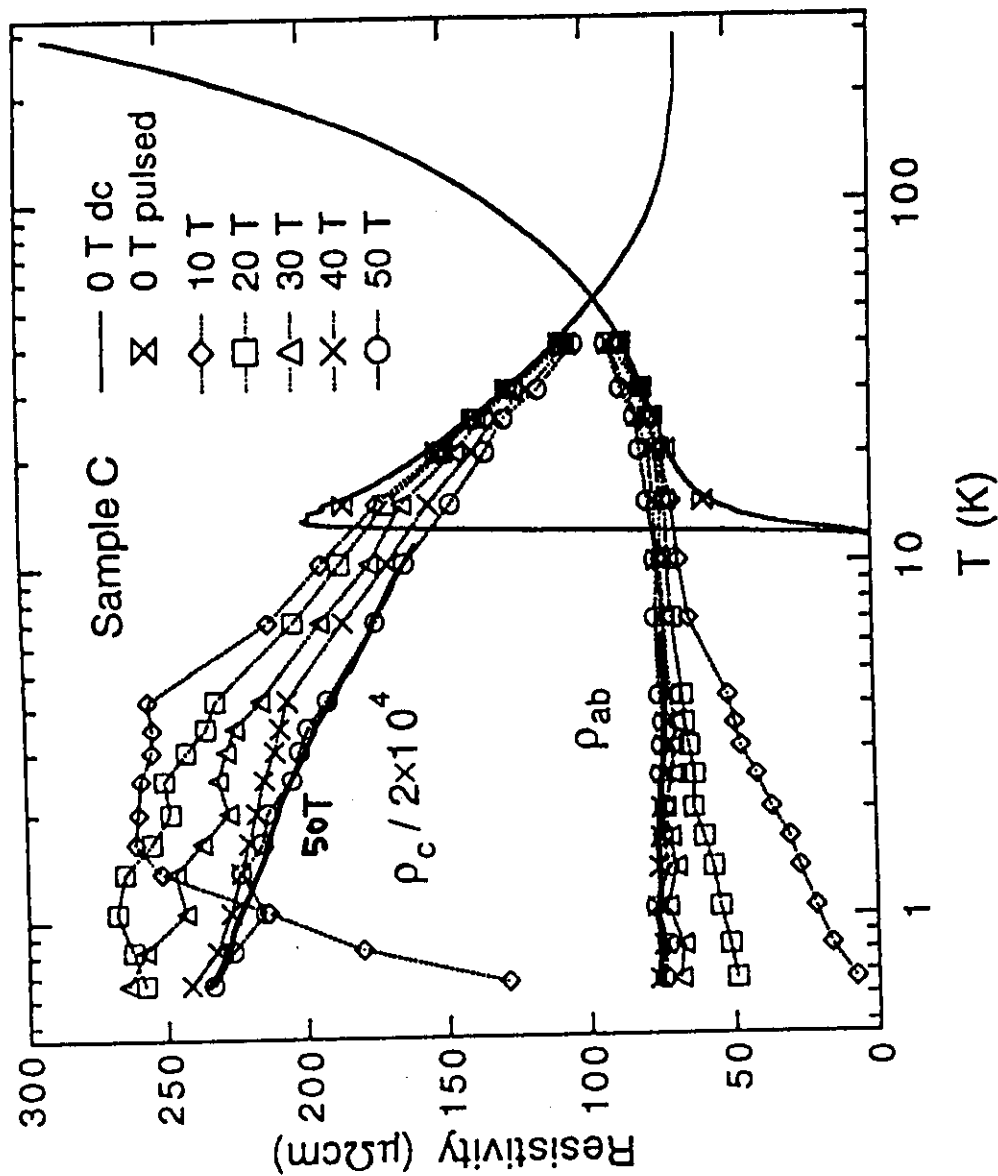
T. Kimura, H.T

PRB 96



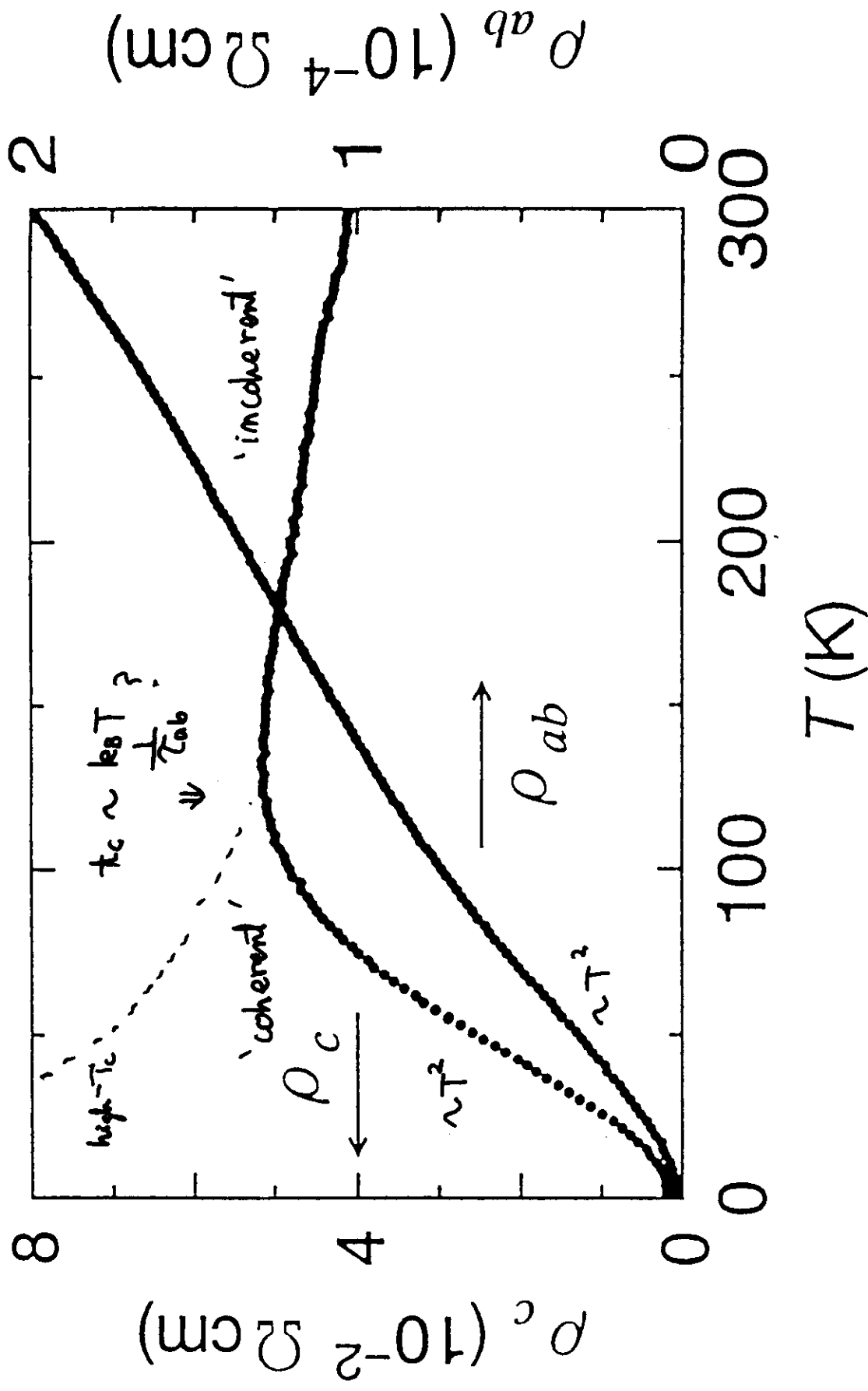
" $1/8$ "
field induced
charge ordering?

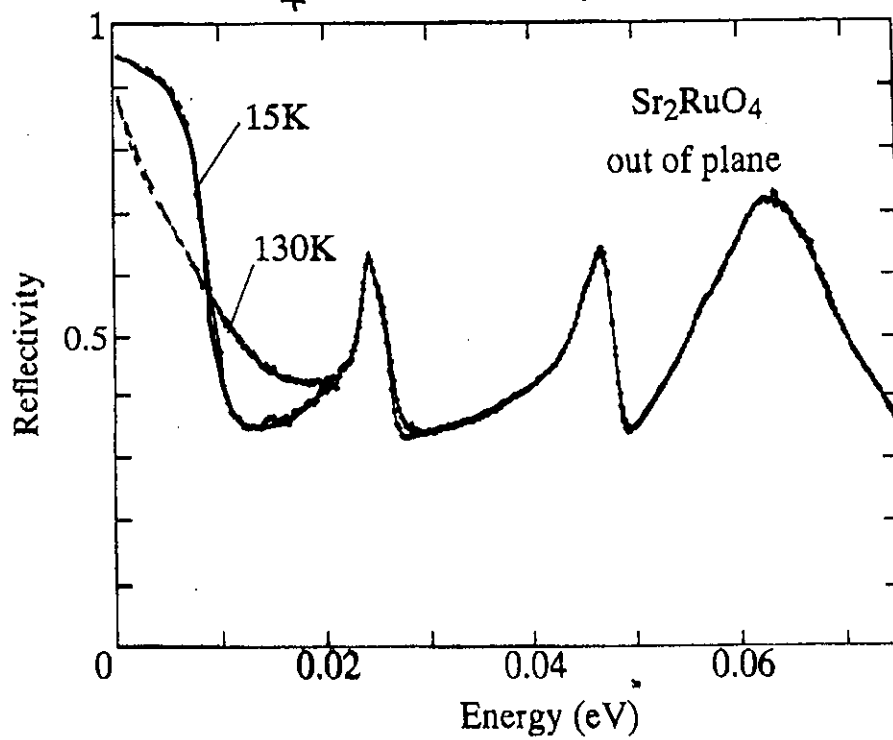
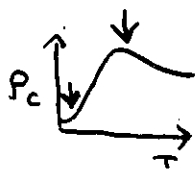
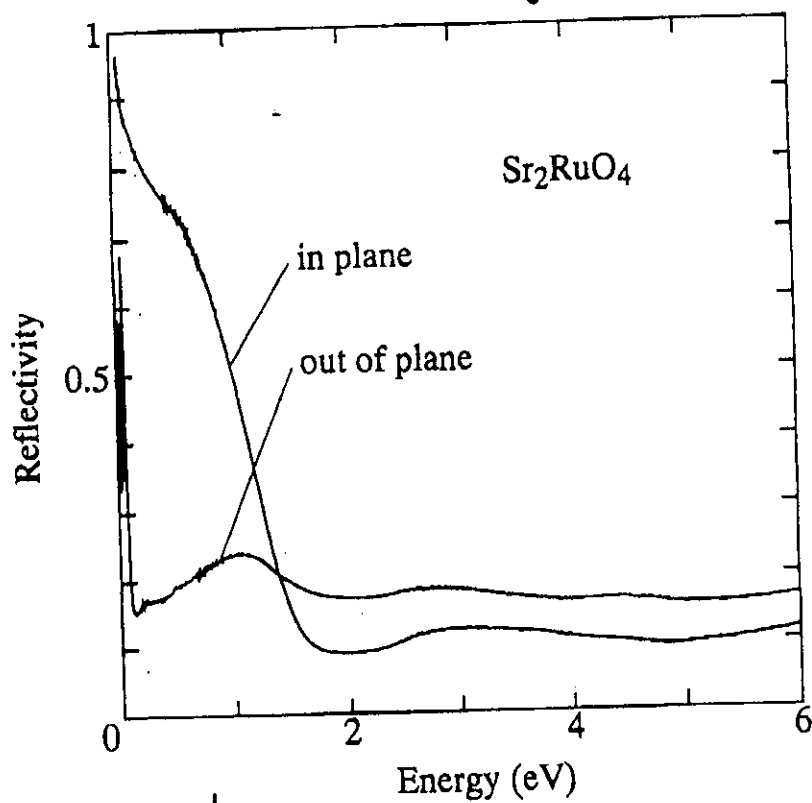
$\text{Bi}_2\text{Sr}_2\text{CuO}_6$ $\rho_c \rightarrow \infty$ but $\rho_a \rightarrow \text{const}$ FL? Ando - Boeinger
 $\rho_a \rightarrow \infty$ LSCO



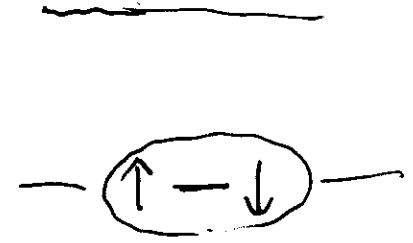
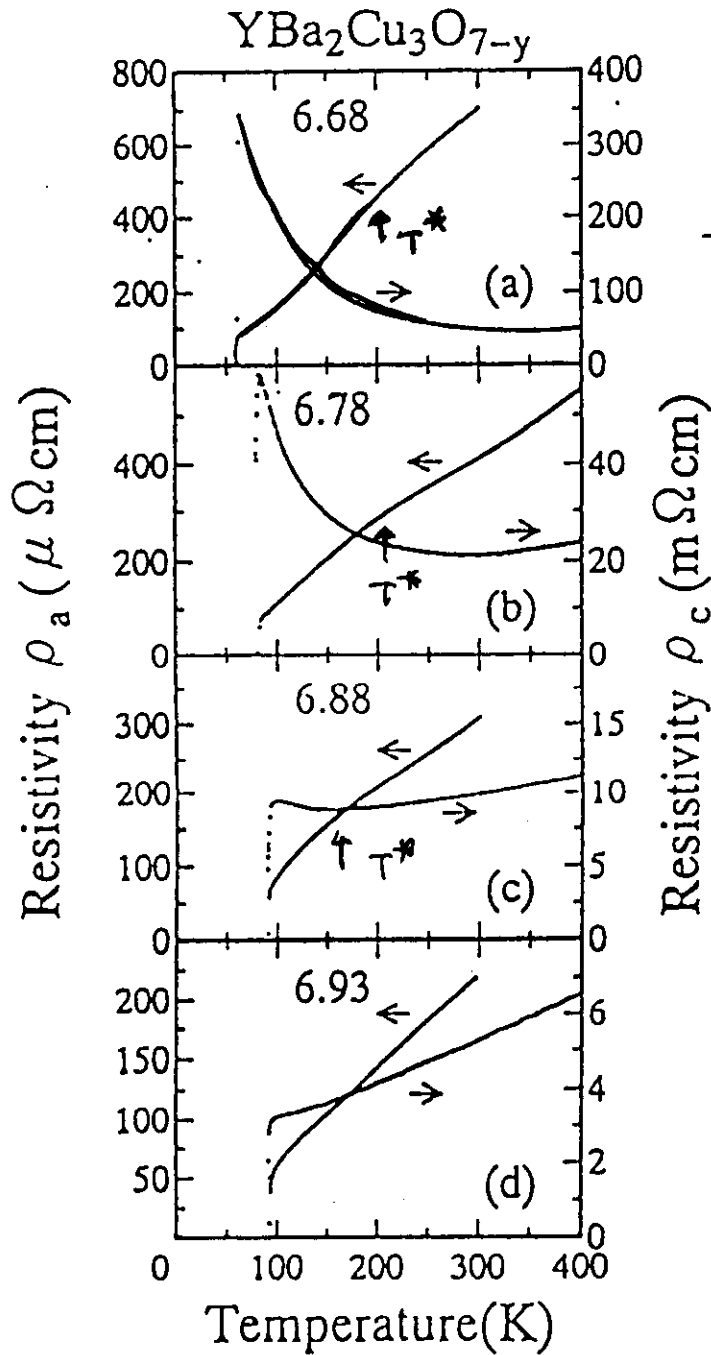
Sr₂RuO₄ (K₂NiF₄ structure)

Maeno et al. Nature





Charge Confinement associated with "spin gap"



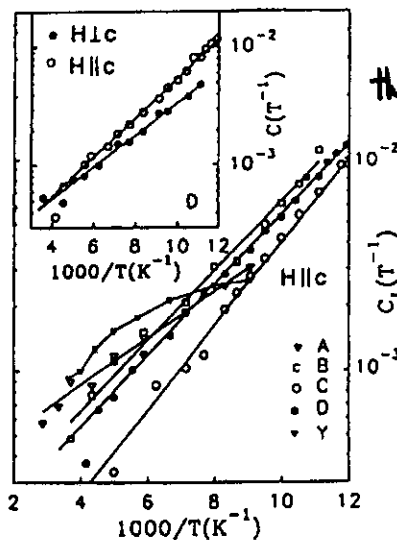
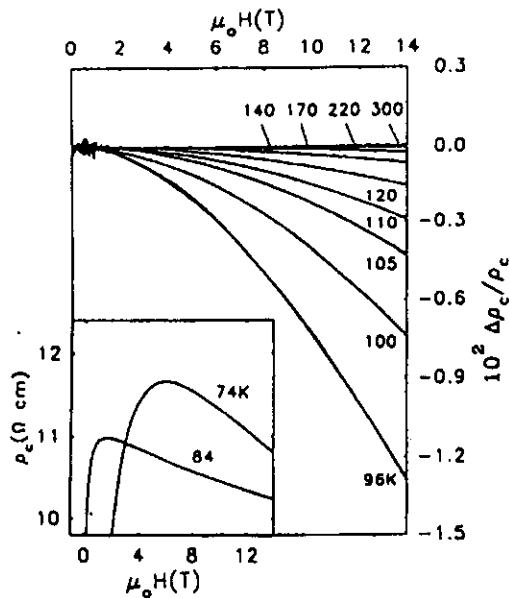
insulating ρ_c
 $\updownarrow$
 'spin gap'

c-axis Magnetoresistance

UNDERDOPED

Negative isotropic MR proportional to B^2

Effect of pseudogap on $\rho_c(T)$ ^{spin}



thermally activated

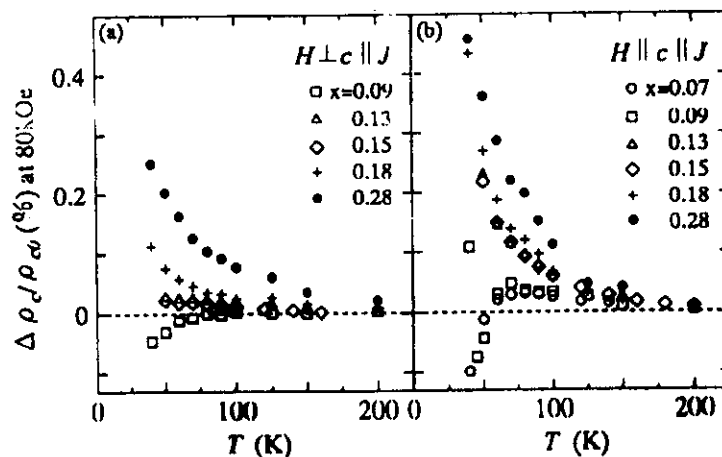
Yan *et al.*
PRB
52 R751
(1995)

OVERDOPED

Crossover to positive orbital B^2 MR

Gradual metalizing of c-axis with increased doping

Unusual longitudinal MR



Kimura *et al.*
PRB
53 8733
(1996)

pseudo gap in $\sigma_c(\omega)$ above T_c in underdoped Y123

Homes
Tajima

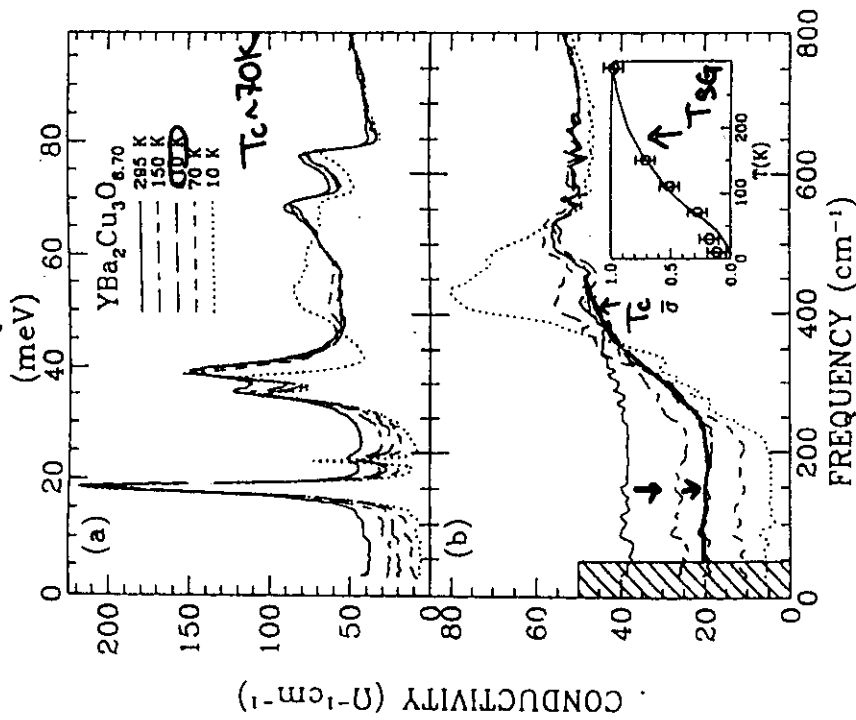


FIG. 2. The optical conductivity of $\text{YBa}_2\text{Cu}_3\text{O}_{6.70}$ along the c axis from ≈ 25 to 800 cm^{-1} obtained by a Kramers-Kronig analysis of the reflectance with (a) the phonons at 150, 152, 286, 317, 557, and 630 cm^{-1} present and (b) subtracted to yield the electronic background. Note that the formation of a pseudogap is clearly visible well above T_c (63 K). The shaded area represents the spectral weight of the condensate (305 cm^{-1}) for $T \ll T_c$. Inset: The conductivity at 50 cm^{-1} normalized with respect to the room temperature conductivity (open circles), compared to the normalized Knight shift for $\text{Cu}(2)$ (solid line) in $\text{YBa}_2\text{Cu}_3\text{O}_{6.63}$ (after Ref. [2]).

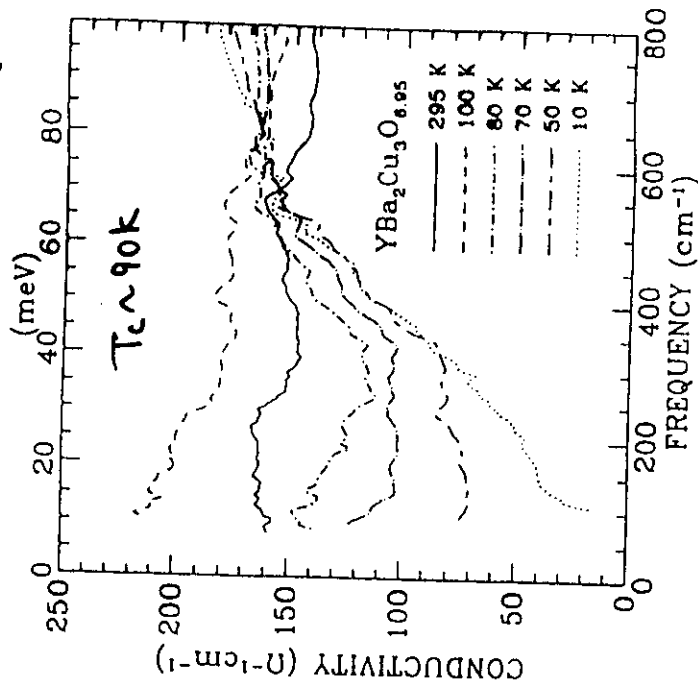


FIG. 3. The optical conductivity of $\text{YBa}_2\text{Cu}_3\text{O}_{6.95}$, along the c axis from ≈ 90 to 800 cm^{-1} , above and below T_c (93 K), with the phonons normally present at 153, 195, 287, 313, and 567 cm^{-1} removed to show the electronic background. The normal-state behavior is metallic. Below T_c a gaplike depression appears in the conductivity, but shifts to zero frequency with decreasing temperature, leaving residual conductivity at $\approx 10 \text{ K}$ (dotted line) down to the lowest measured frequency point $\approx 100 \text{ cm}^{-1}$.

insulating ρ_c
↕
pseudo gap in $\sigma_c(\omega)$
↓
appears to
be connected
with SC gap
smoothly

What is generally believed for the out-of-plane transport in high-T_c cuprates

ρ_c/ρ_a by far larger than those predicted by band calculations

temperature dependence distinctly different from ρ_a

no apparent correlation between ρ_c/ρ_a and T_c

insulating behavior pronounced in the underdoped region, in particular when spin gap is developed ($T < T^*$)

- closely related with pseudo gap in $\sigma_c(\omega)$

- isotropic negative MR

- charge confinement associated with the spin gap?

- T=0 limit

logarithmic divergence with $T \rightarrow 0$,
accompanied by in-plane insulating behavior
underdoped LSCO, but universal?

ρ_c diverge with $T \rightarrow 0$ (Bi2201) but not for ρ_a

tend to recover coherent behavior when overdoped

- ρ_c is metallic and appears to stay finite with $T \rightarrow 0$ (LSCO, Y123 90K, Tl)

Y124
 ρ_c
distinctly
different

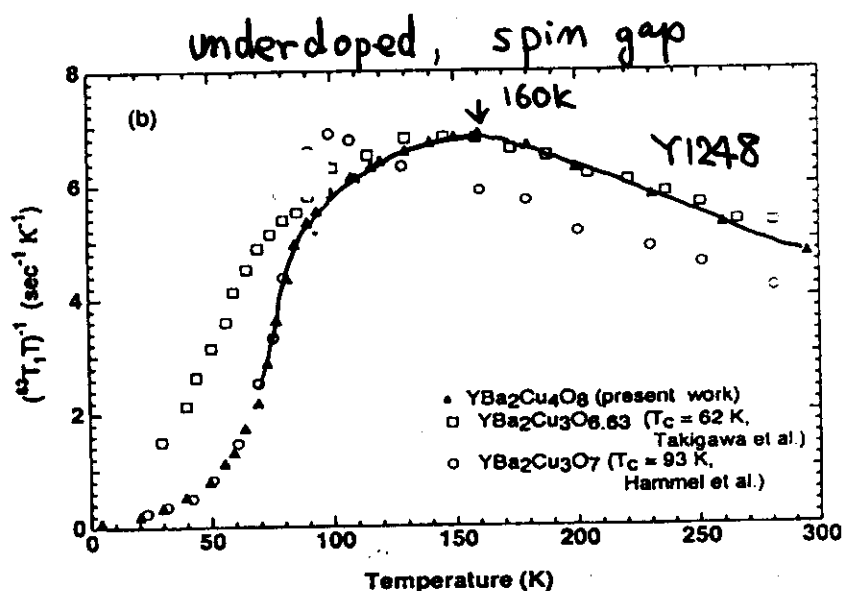
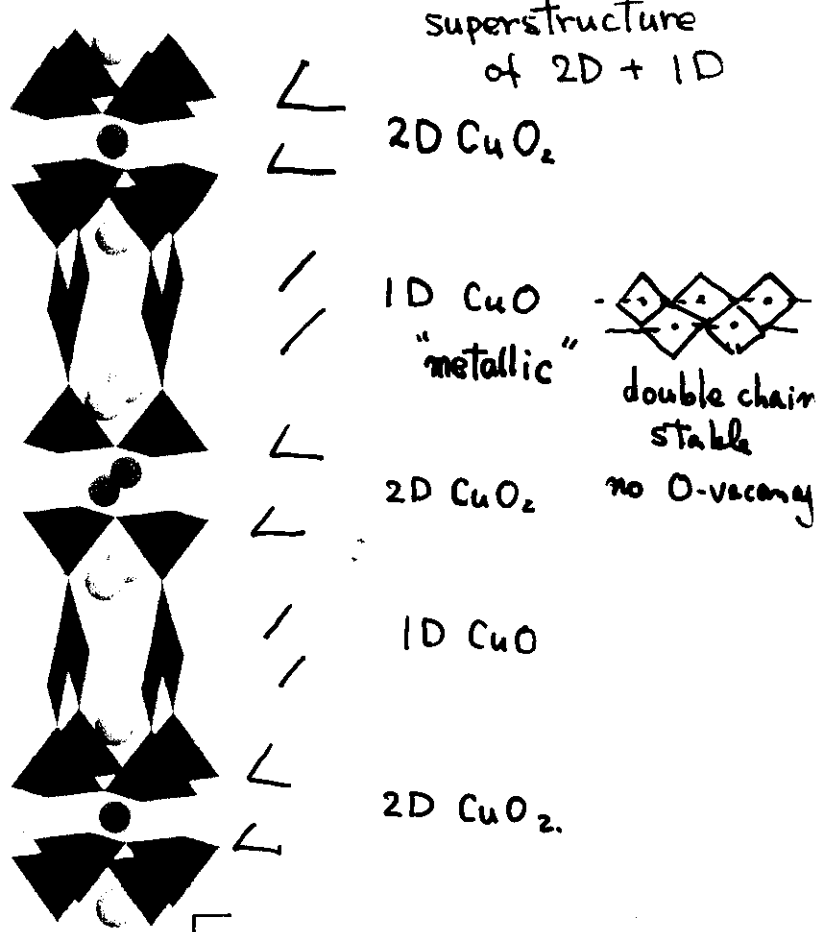
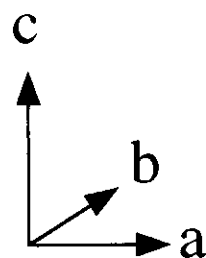
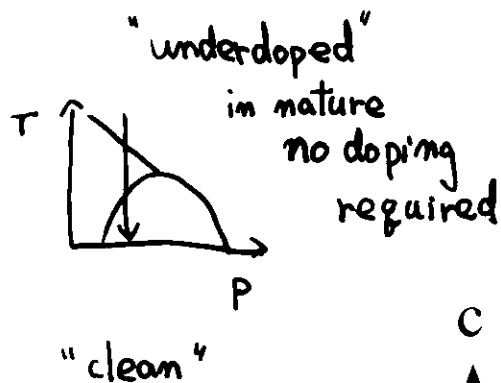
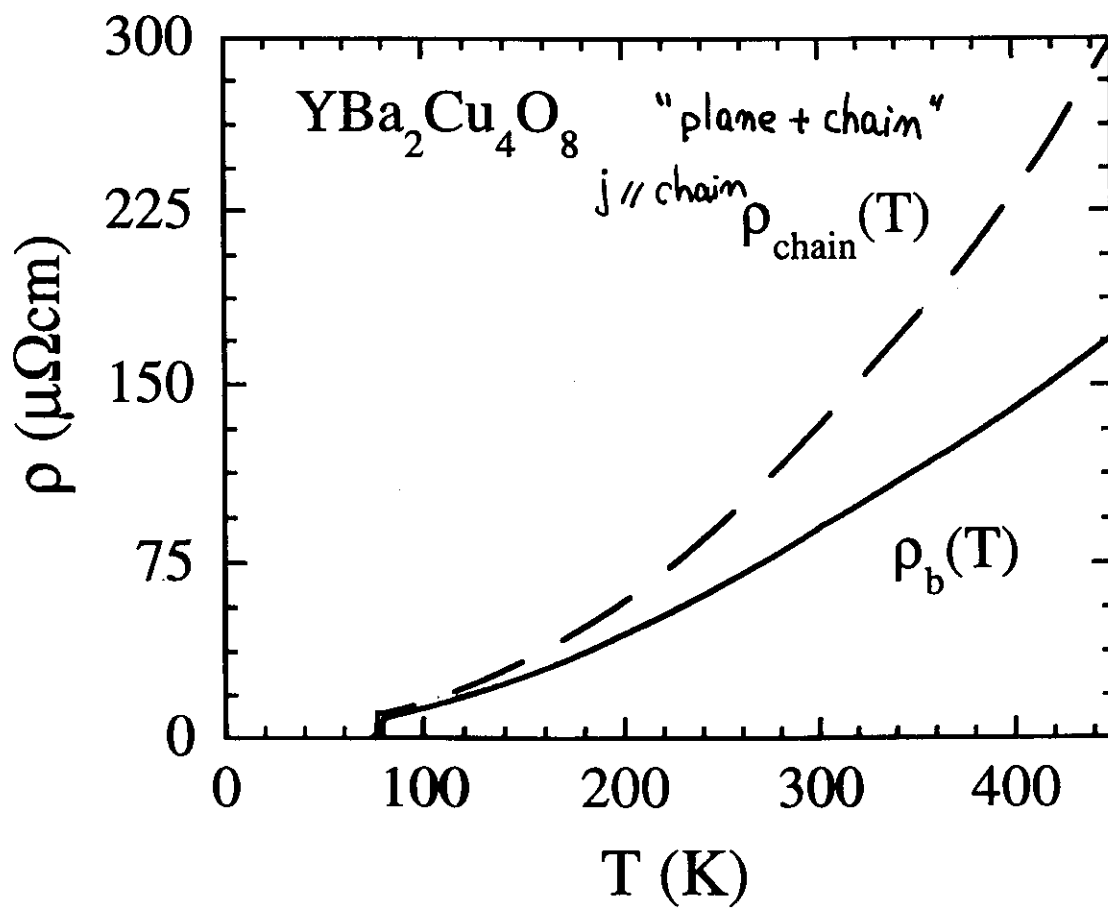
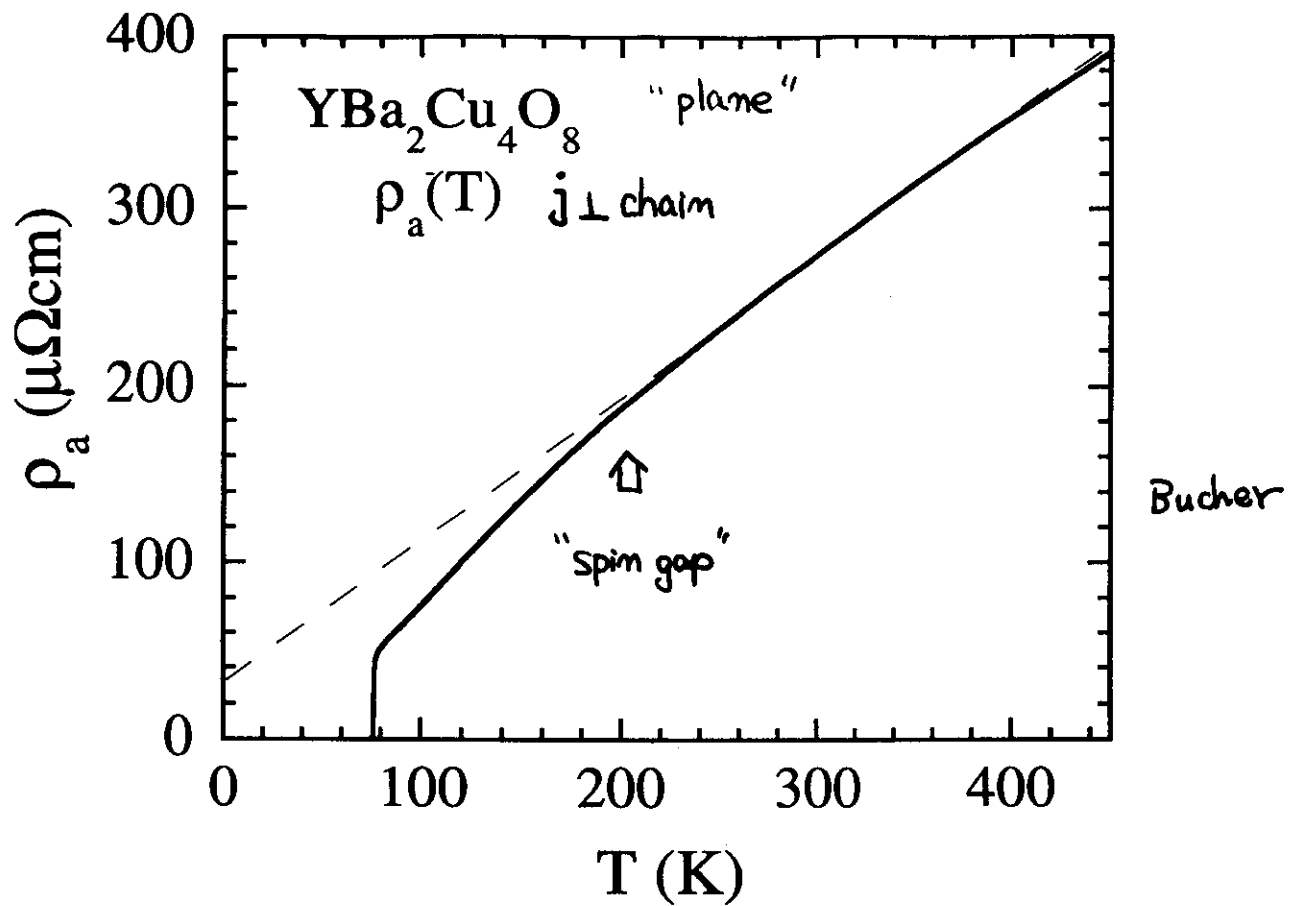


FIG. 8. Comparison of present $(^aT_1T)^{-1}$ data for YBa₂Cu₄O₈ (triangles) with those for YBa₂Cu₃O_{6+x}. The data for the 90-K phase (circles) and the 60-K phase (squares) of YBa₂Cu₃O_{6+x} were taken from Hammel *et al.* (Ref. 2) and Takigawa *et al.* (Ref. 3), respectively. (a) $(^{17}T_1T)^{-1}$ at the O(2,3) site. (b) $(^{63}T_1T)^{-1}$ at the Cu(2) site.

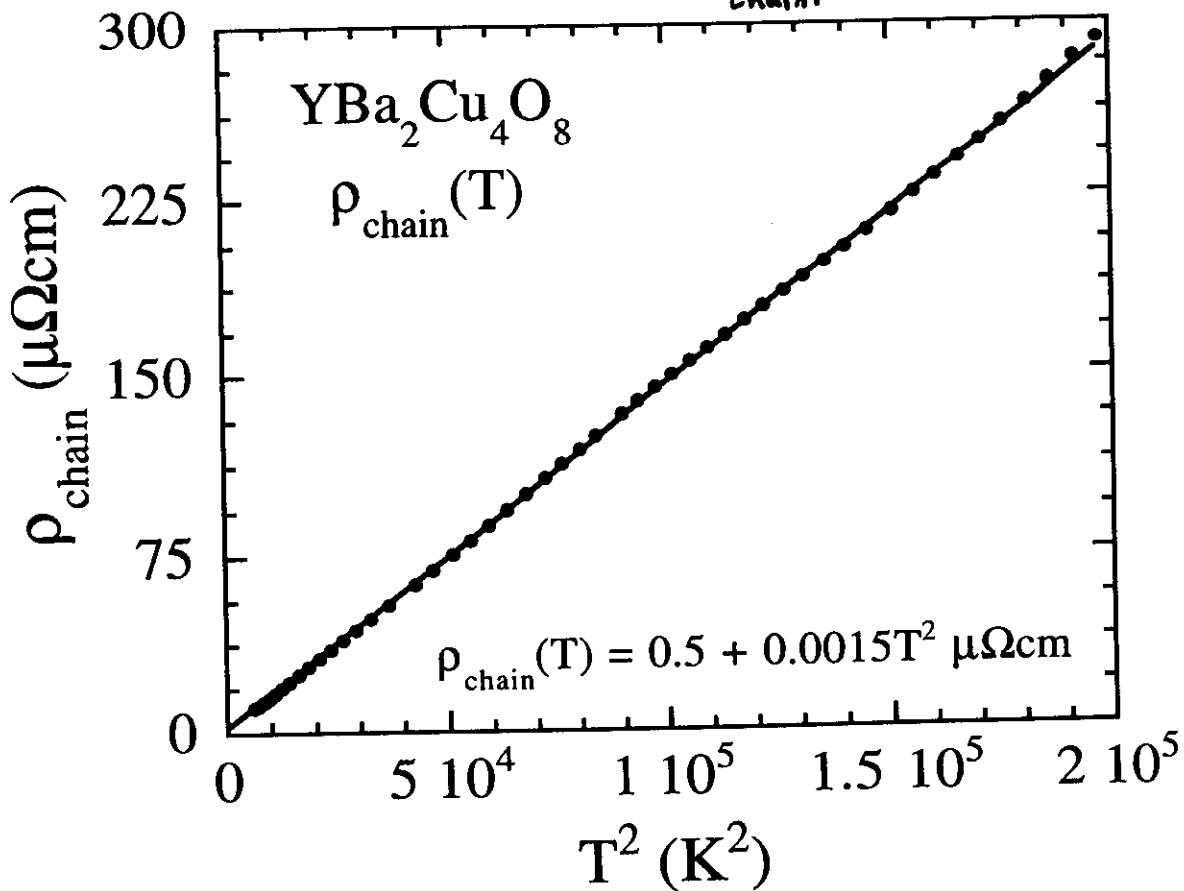
I. Tomeno, et al., *PRB.*, **49**, 15327(1994).

In-Plane resistivity



$$\frac{1}{\rho_{\text{chain}}} = \frac{1}{\rho_b} - \frac{1}{\rho_a}$$

plane + chain ↑ plane



spin gap effect
not clear
in ρ_{chain}

1 $\rho_{\text{chain}} \sim 5 \mu\Omega\cdot\text{cm}$

$\rightarrow 700 \text{ \AA} \text{ at } 80 \text{ K} !$

2 $\rho = \rho_0 + AT^2$

$\rho_0 \sim 0.5 \mu\Omega\cdot\text{cm}$

3. T^2 behavior like conventional FL

3D coupling of chains
through 2D CuO₂ sheets.

Out-of-plane ρ_c

metallic at low-T
though underdoped
with spin gap.

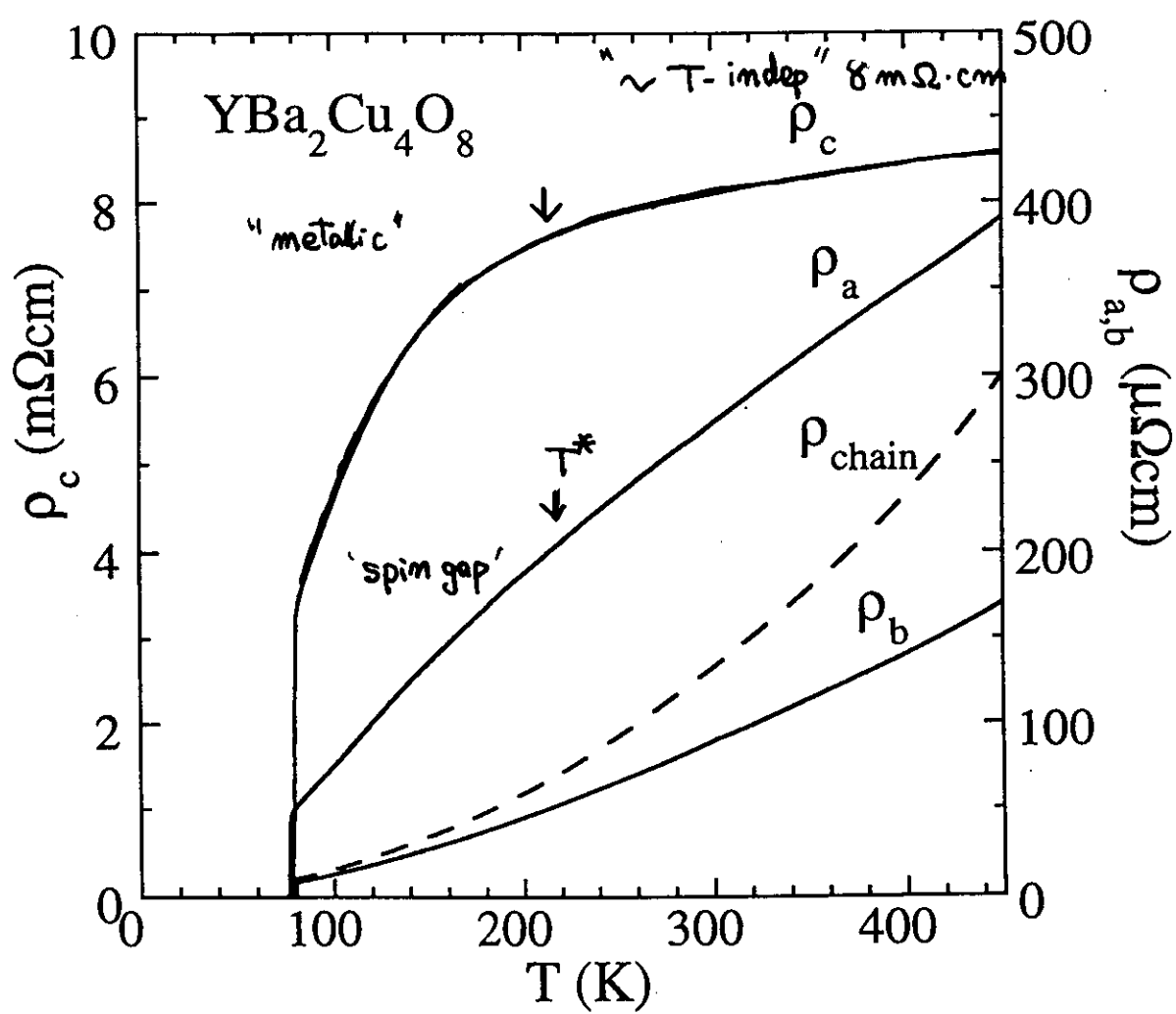
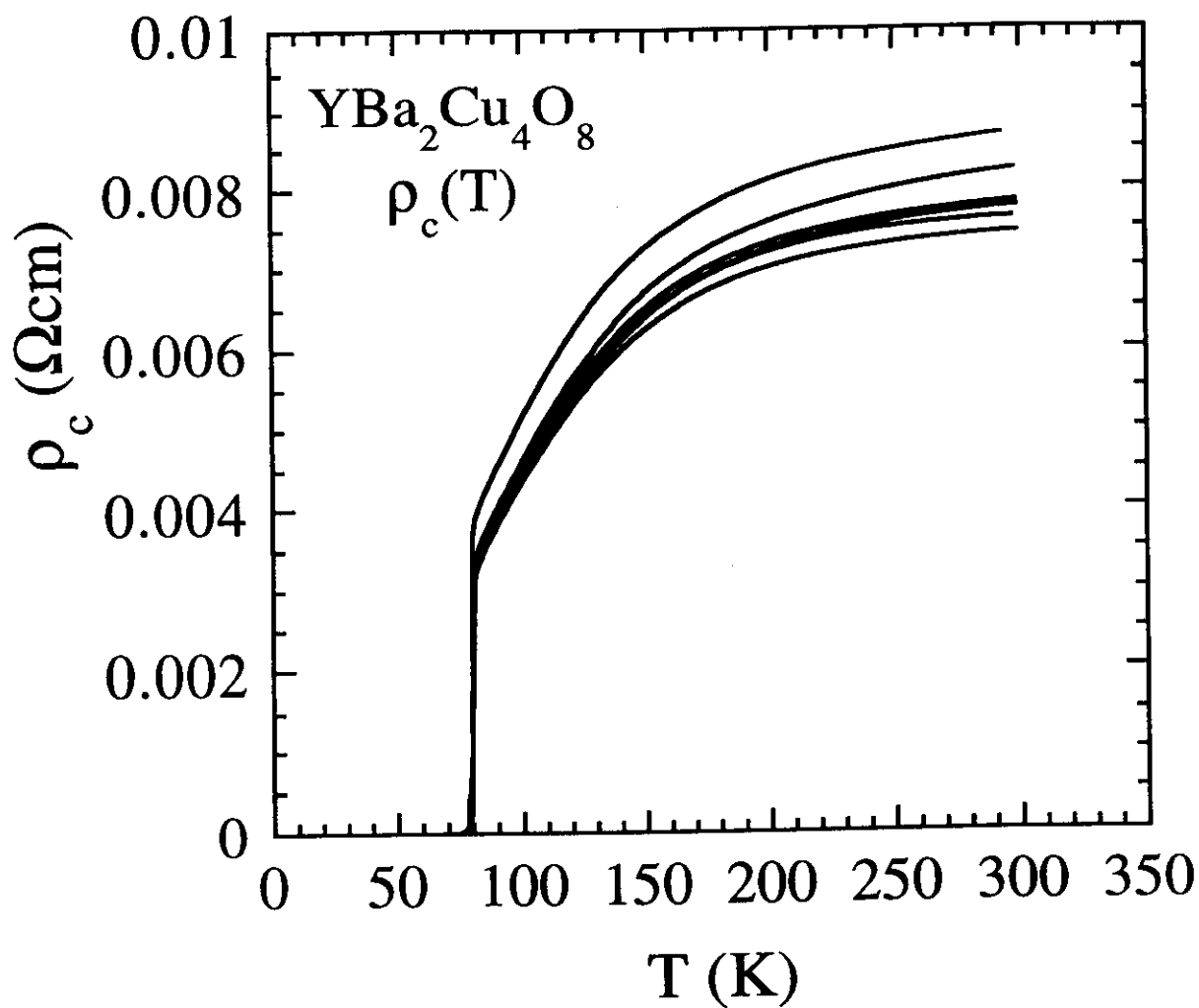
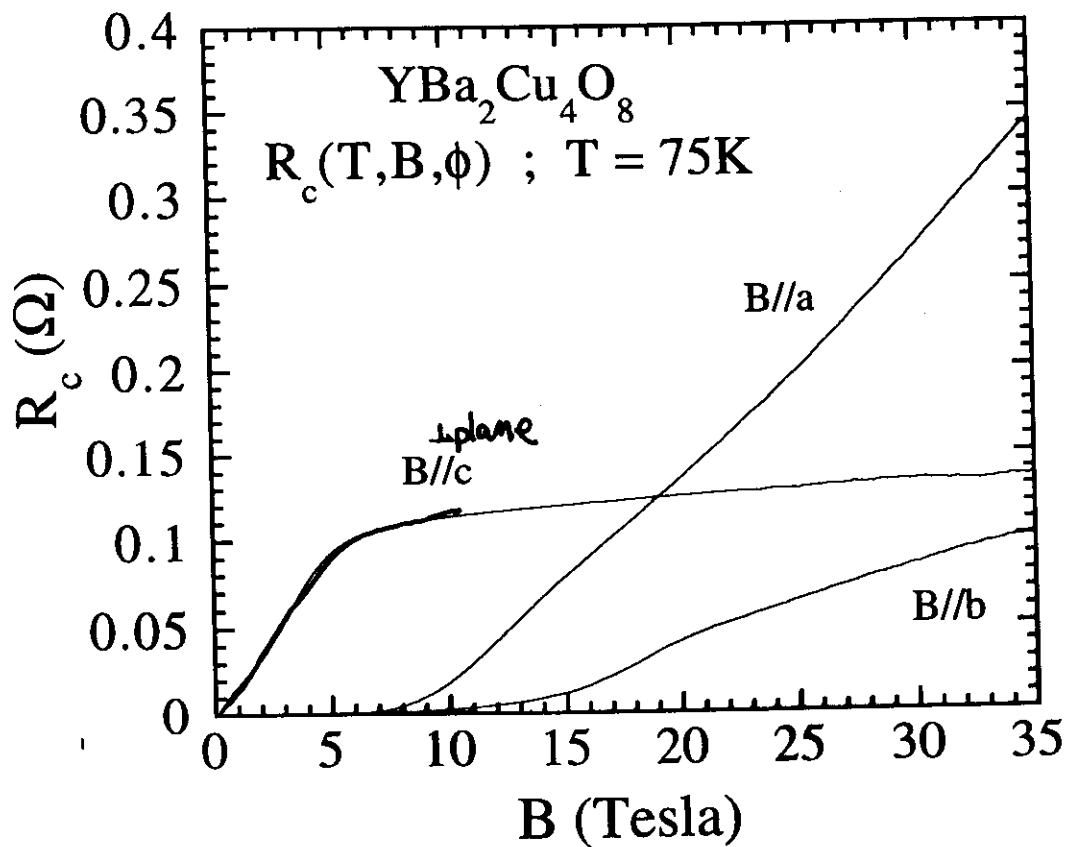
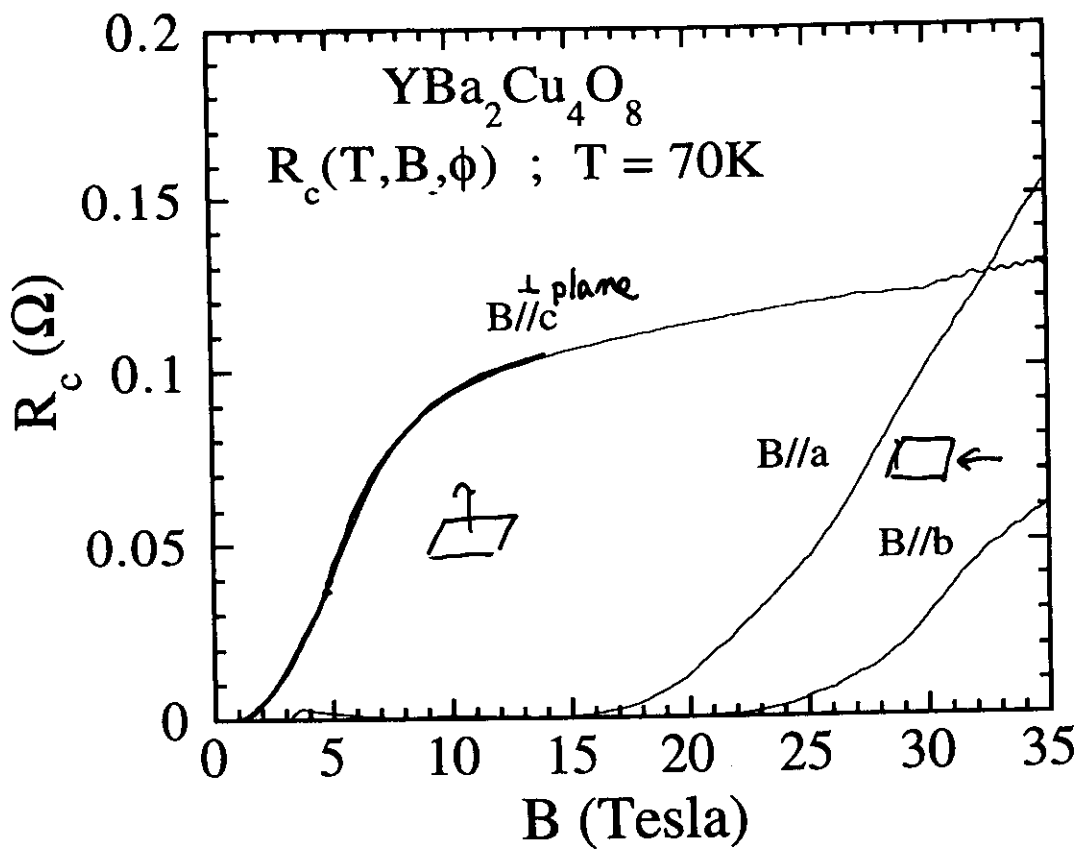


Figure 2
N.E. Hussey et al

Shorted out due to misoriented domains!

"reproducibility"





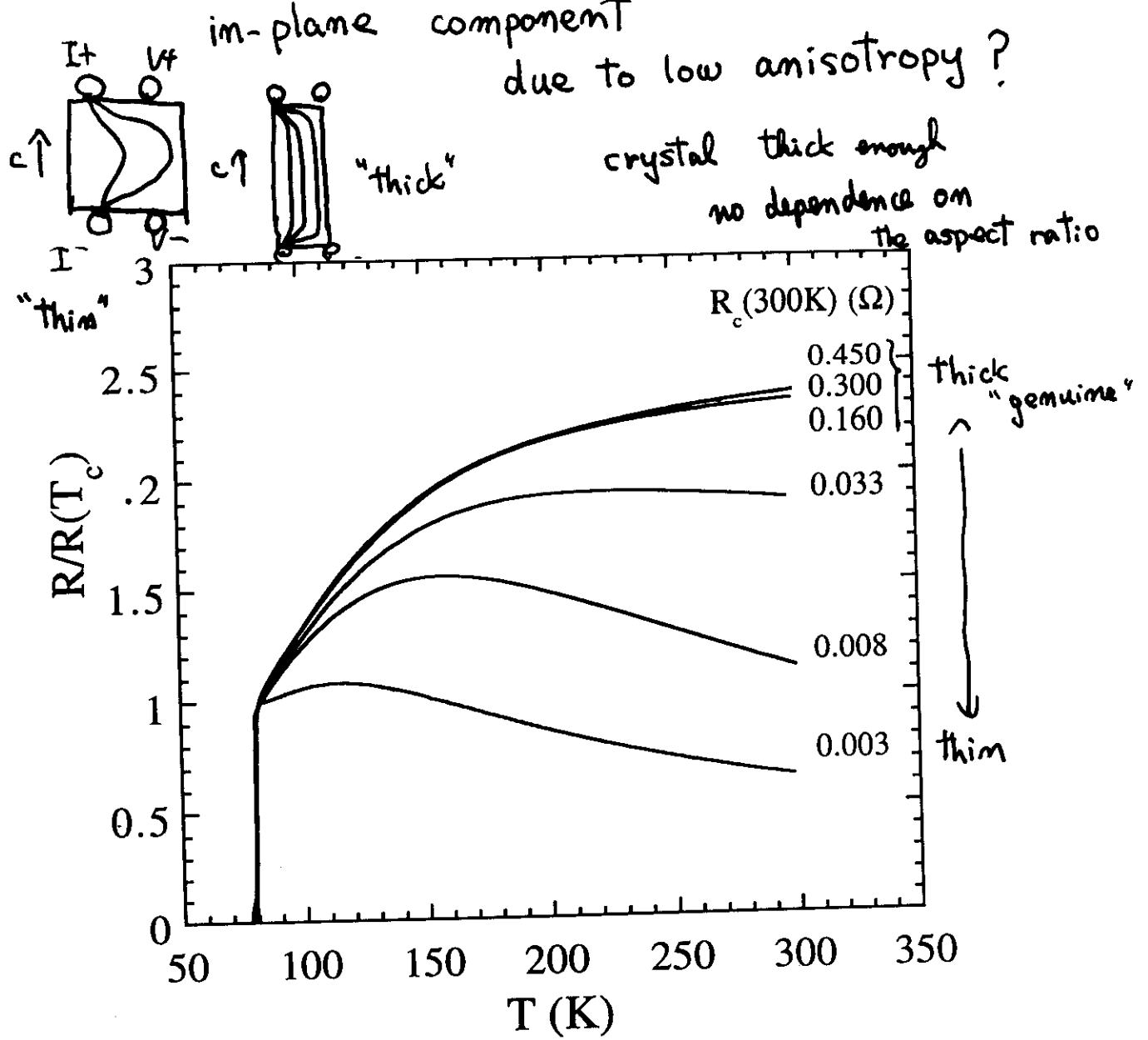


Figure 1
N.E. Hussey et al.

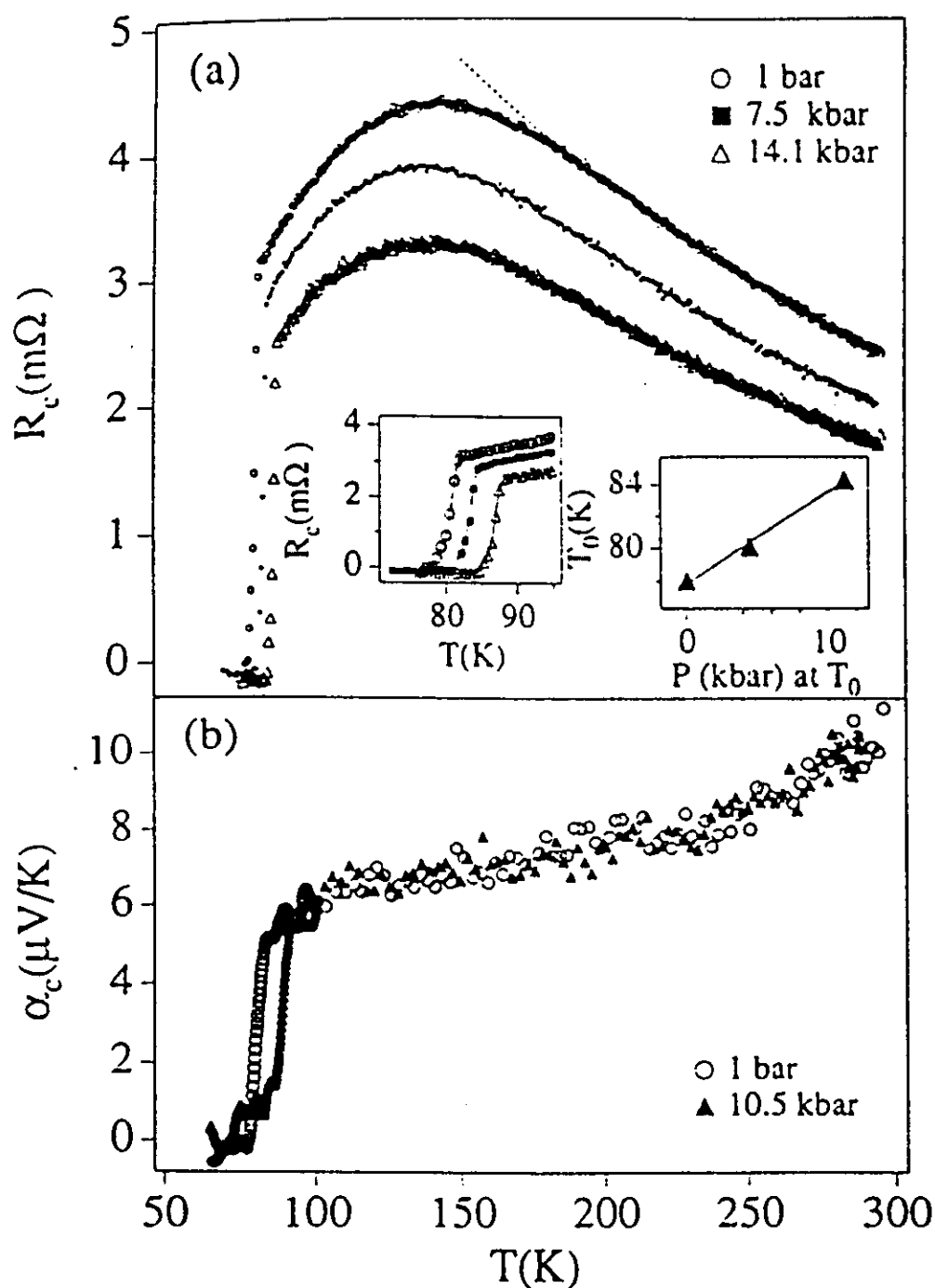
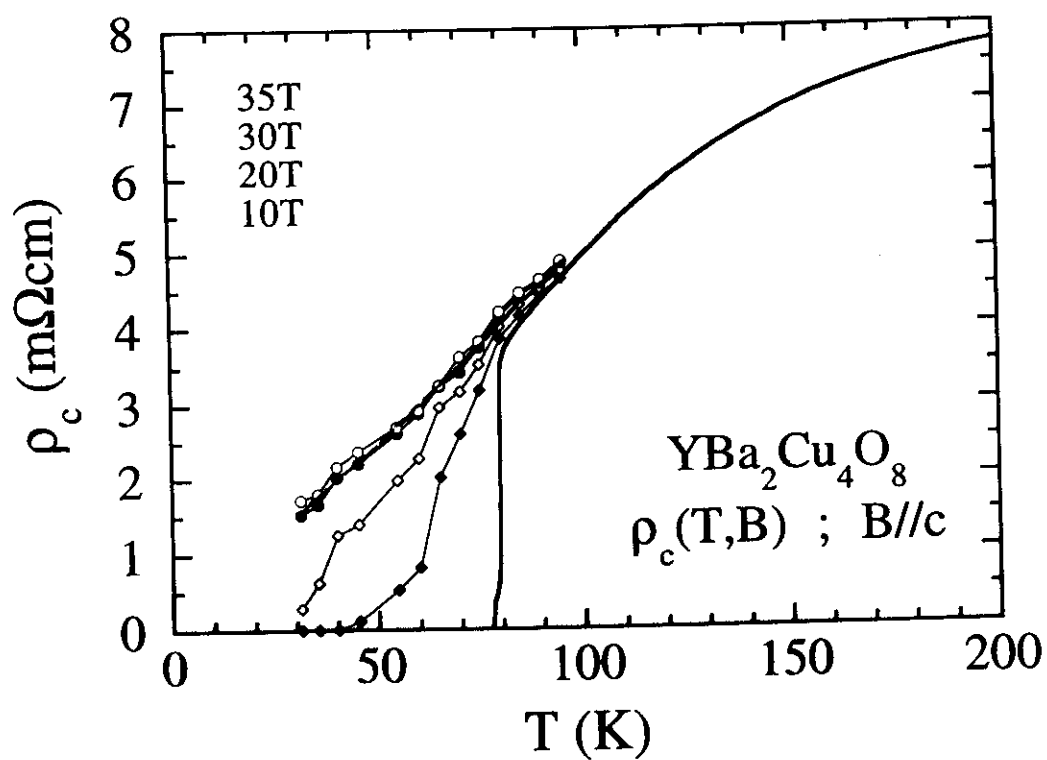
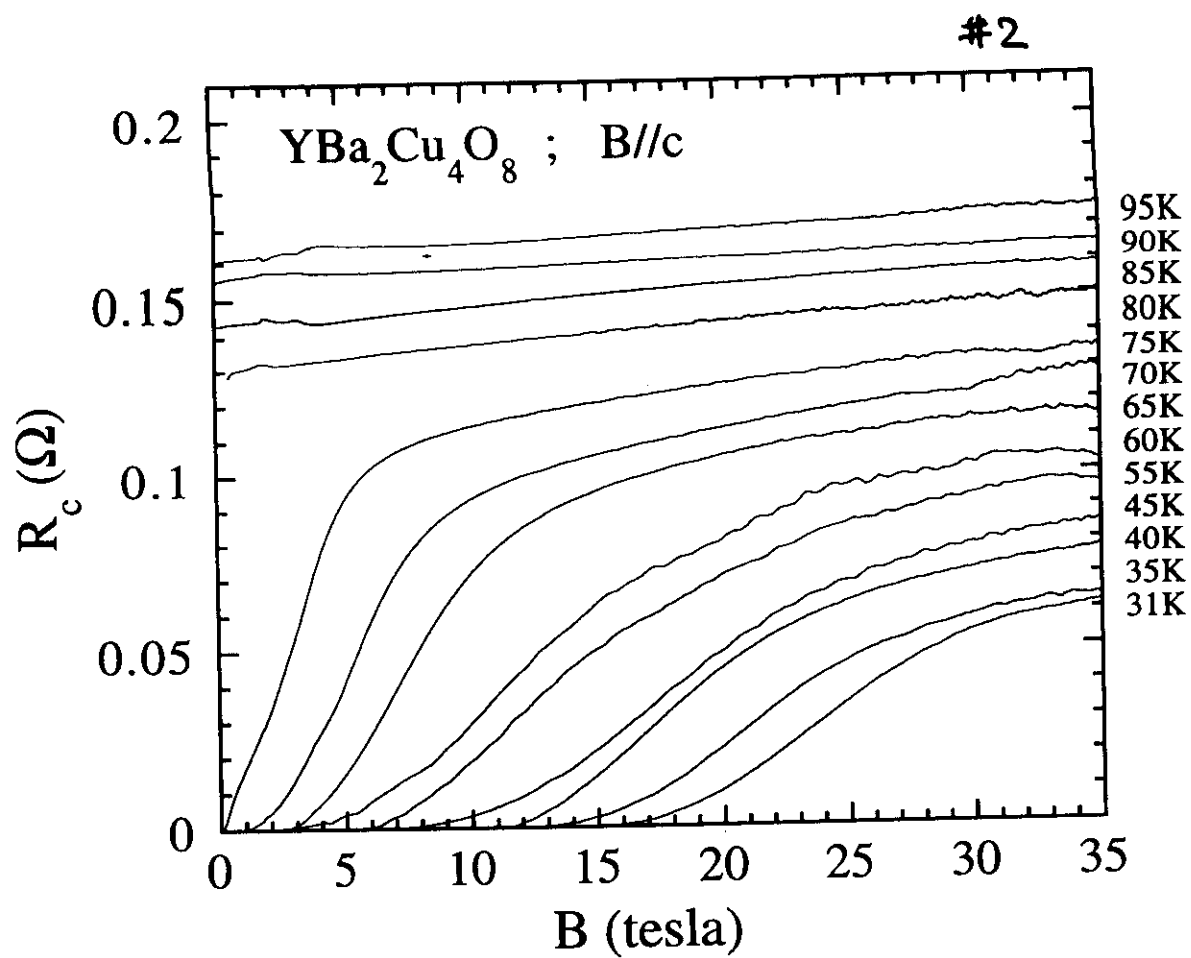
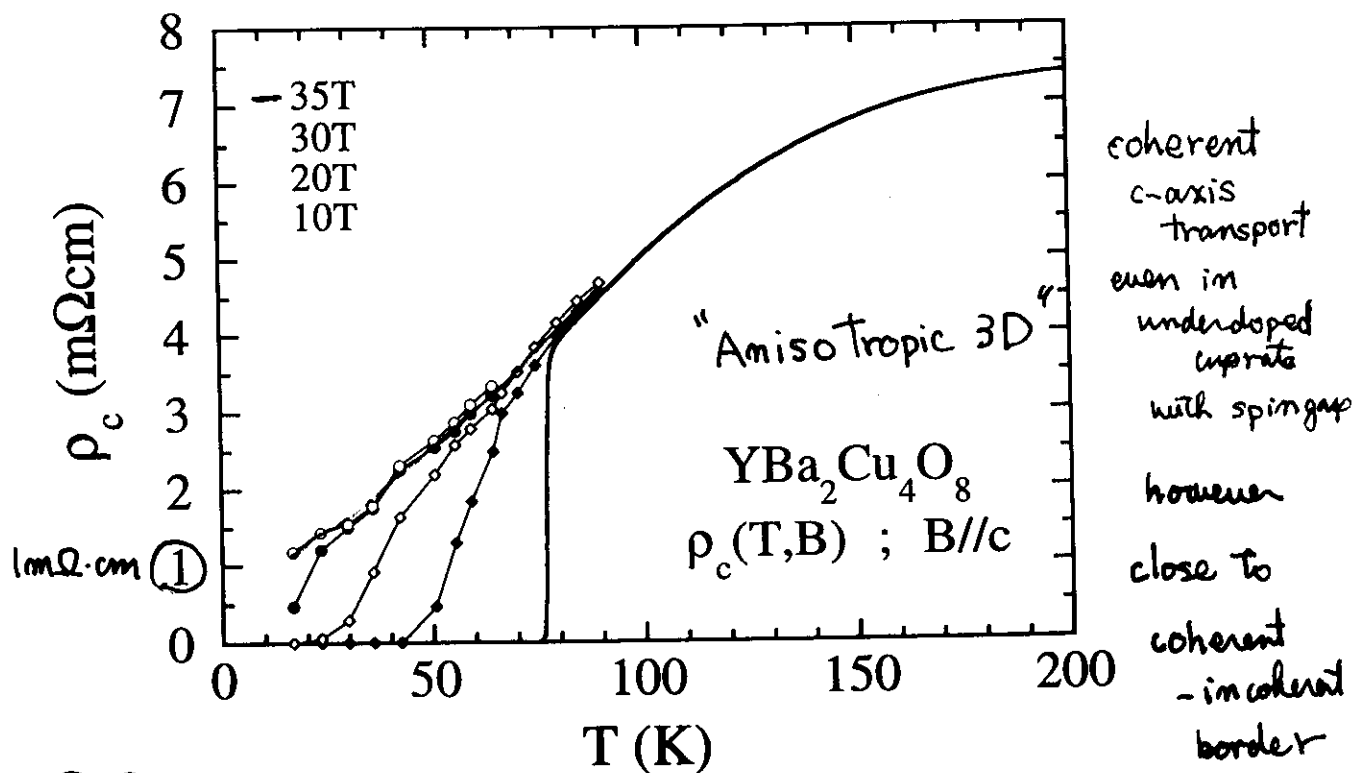
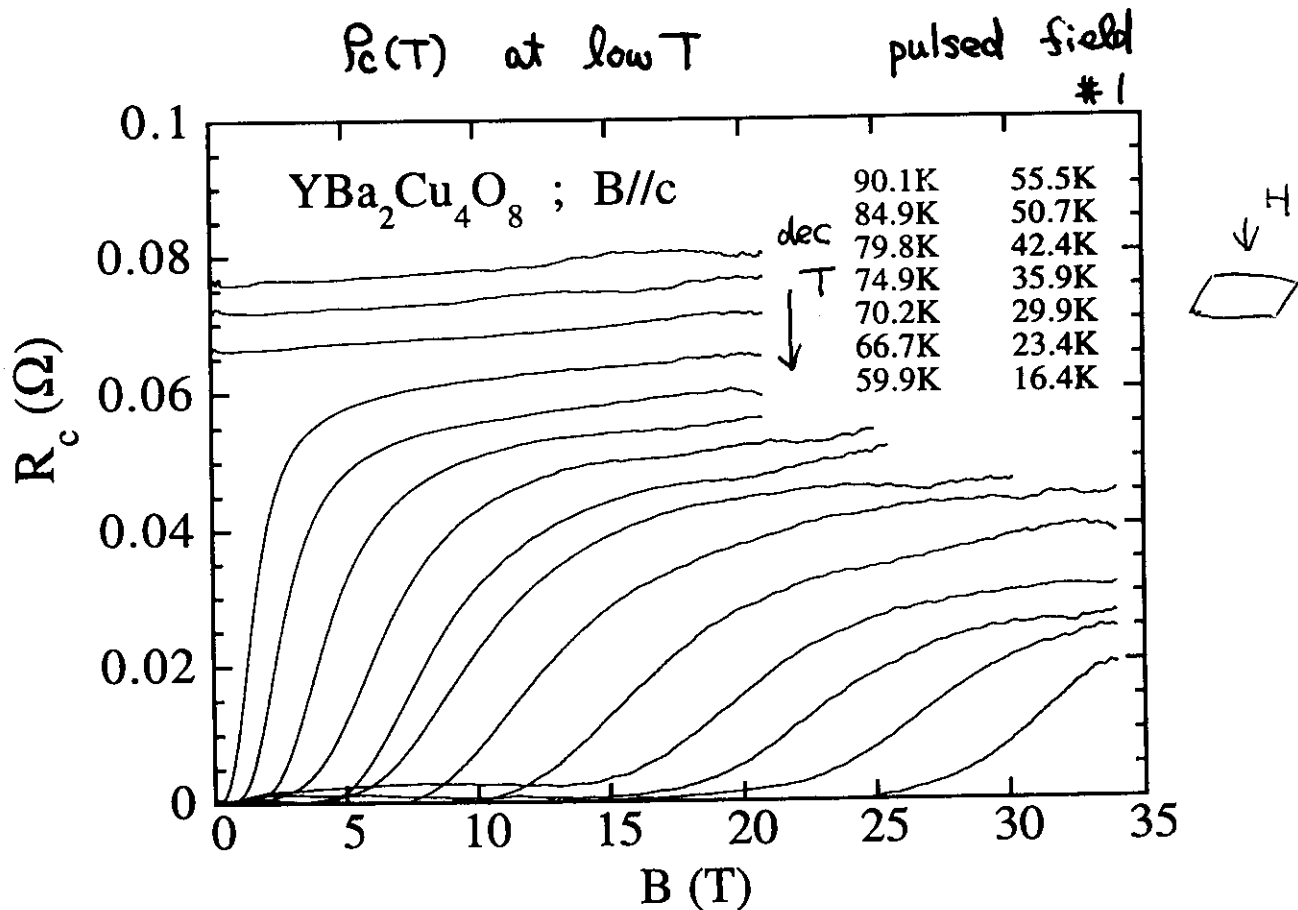


FIG. 4. The resistance (a) and TEP (b) along the c axis versus temperature under several pressures. The dashed line is an exponential fitting to the resistance at 1 bar. Inset: detail of the resistivity near transition temperature and superconductive transition temperature T_0 versus pressure.



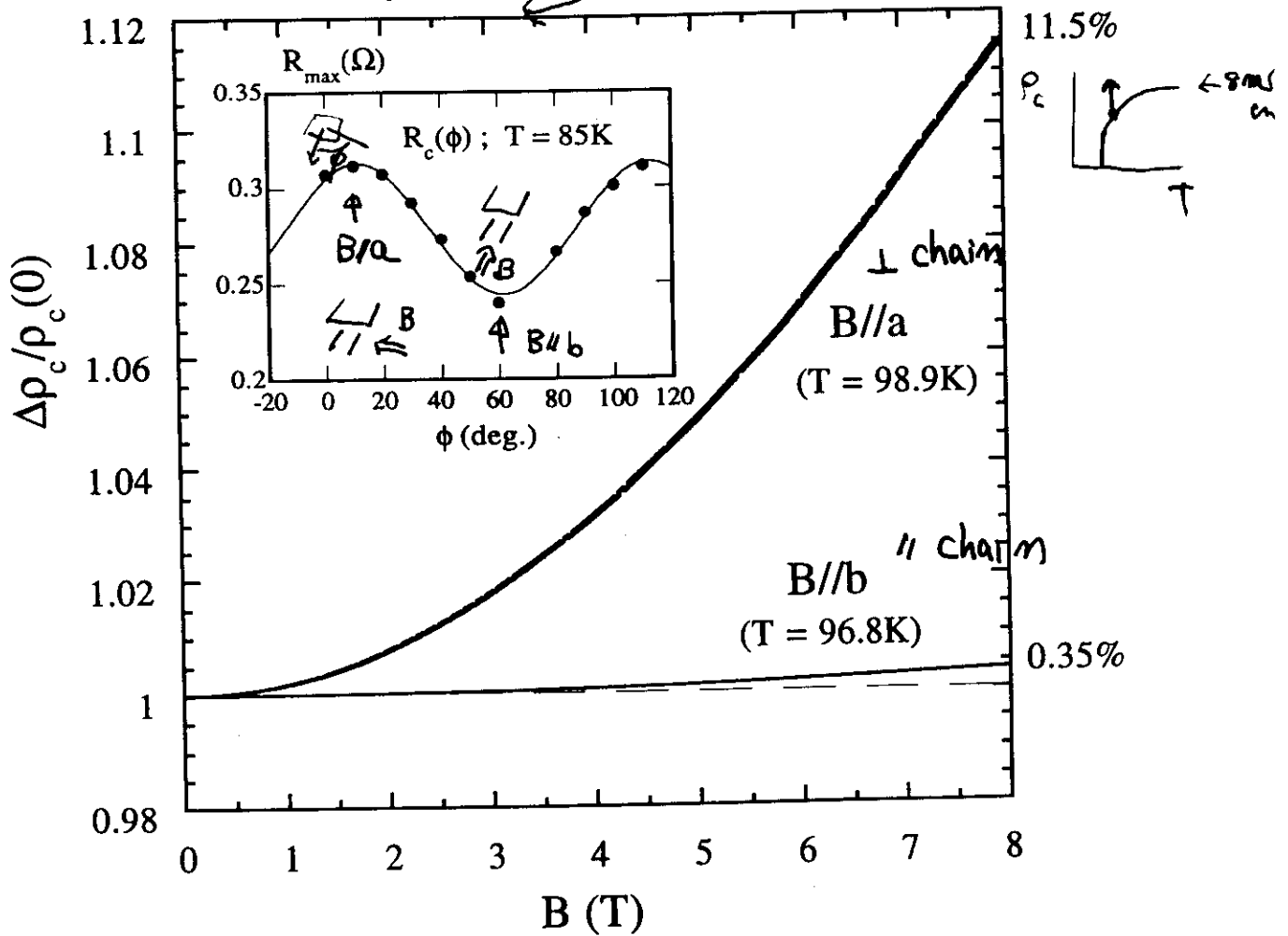


Sr₂RuO₄
 $\rho_c(0) \approx 1 \text{ m}\Omega\text{cm}$
 below 3D Mott limit.

C-axis MR reflects the symmetry of chains

- Importance of clean and metallic

chain for the coherent c-axis transport



metallic $\rho_c(T)$

$\longleftrightarrow$ metallic chain

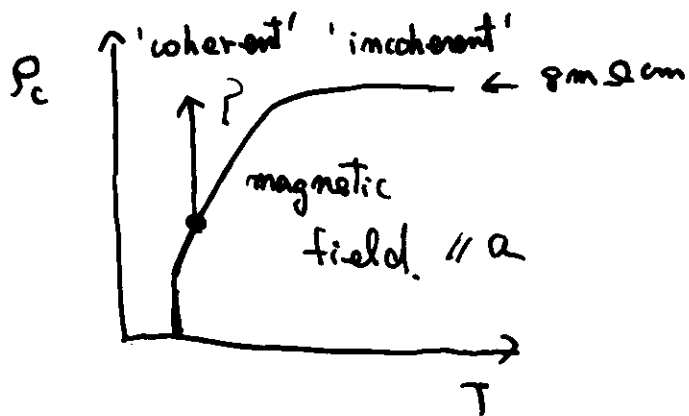
chain carries alone

can lead the coherent c-axis transport??

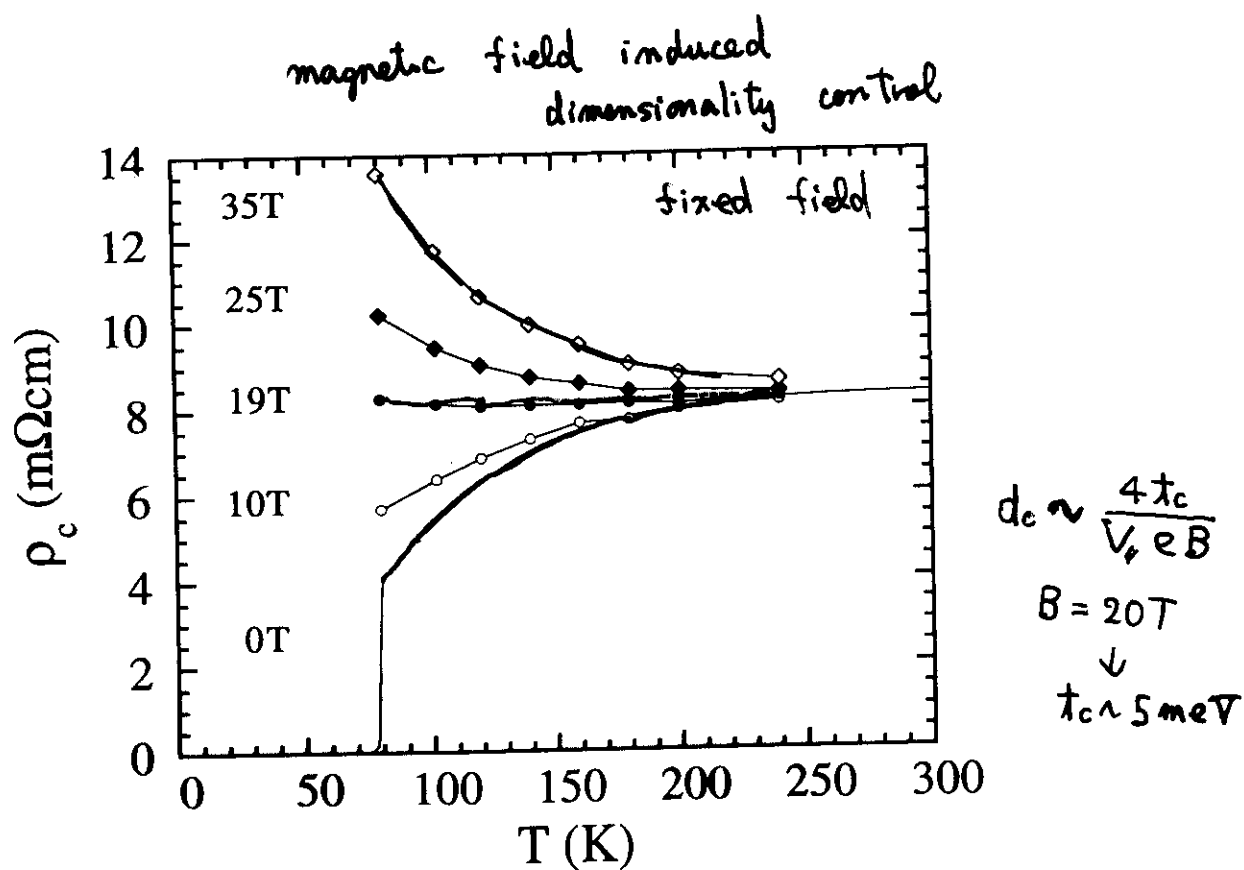
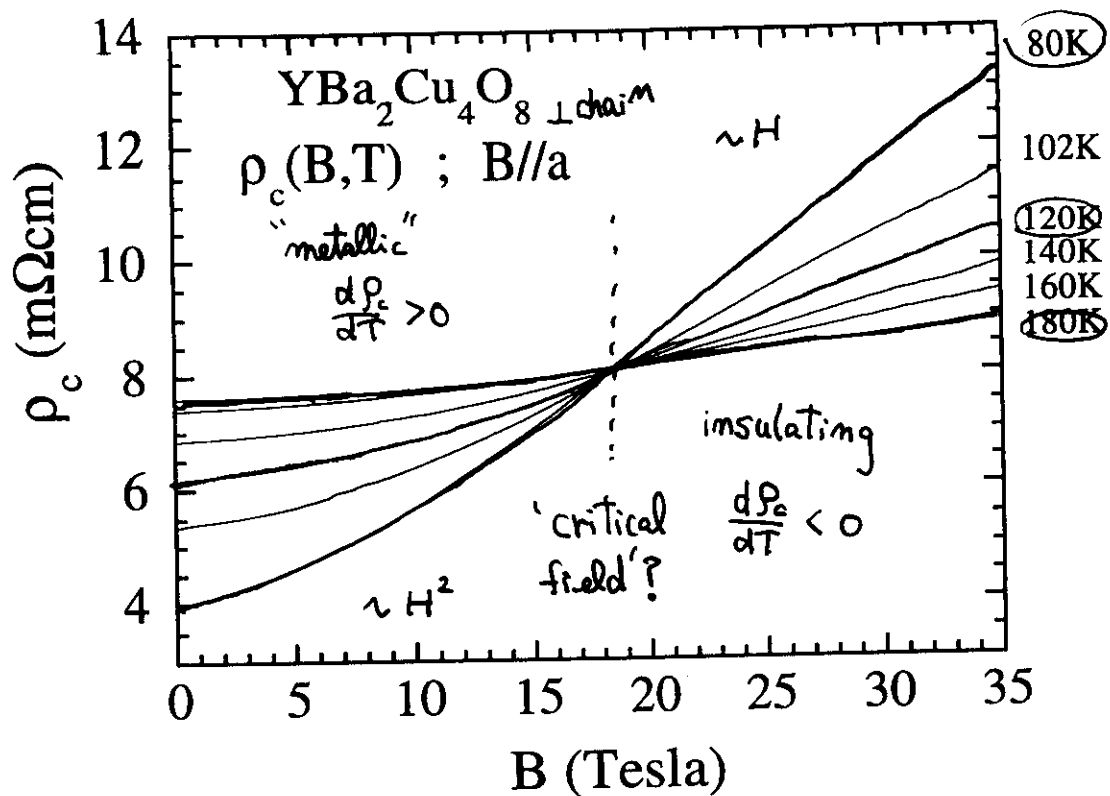
$\rightarrow$ to achieve coherent c-axis transport,

chain carriers must be coupled with plane carriers

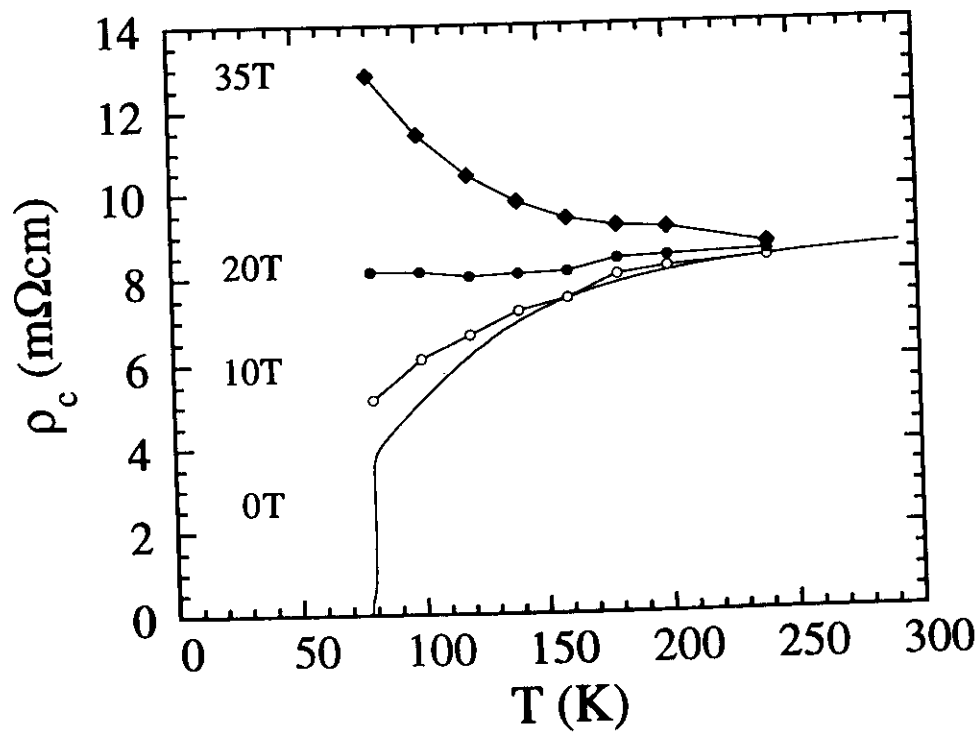
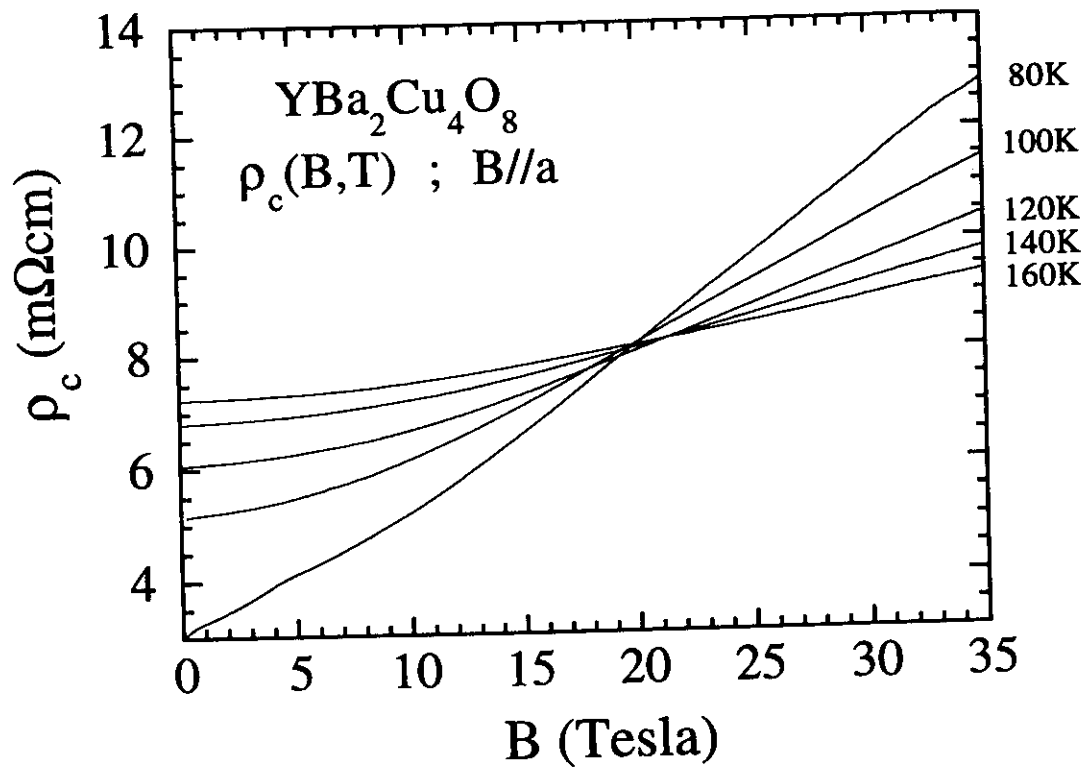
What will happen if $\rho_c \rightarrow 8m\Omega_{cm}$
with high magnetic field.

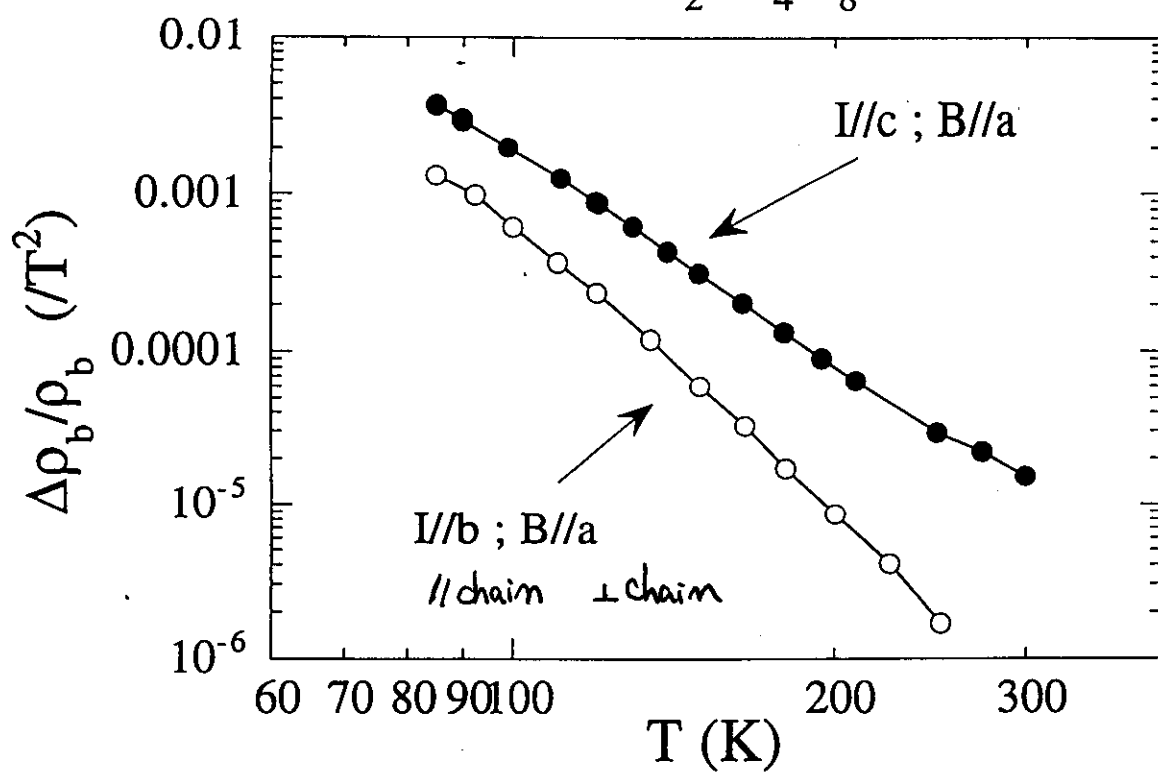
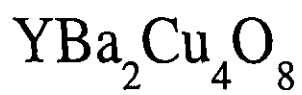


High-field normal state c-axis MR



sample #2.





B $\perp$ chain

$$\frac{\Delta\rho_b}{\rho_b} \sim \frac{1}{3} \frac{\Delta\rho_c}{\rho_c}$$

Y124 *summary*

Although Y124 is underdoped with a spin gap, ρ_c shows a incoherent to coherent crossover with decreasing temperature

- ρ_c at low temperatures falls below Mott resistivity
- anisotropic 3D transport

Can charge still be confined within the plane?

clean and surprisingly metallic double CuO chains is very likely responsible for the coherent out-of-plane transport

magnetic field can decouple planes and chains

- possible field induced dimensionality change

coherent but close to coherent - incoherent crossover

What is generally believed for the out-of-plane transport in high- T_c cuprates

ρ_c/ρ_a by far larger than those predicted by band calculations

temperature dependence distinctly different from ρ_a

no apparent correlation between ρ_c/ρ_a and T_c

insulating behavior pronounced in the underdoped region, in particular when spin gap is developed ($T < T^*$) ?

- closely related with pseudo gap in $\sigma_c(\omega)$
- isotropic negative MR
- charge confinement associated with the spin gap?
- $T=0$ limit

logarithmic divergence with $T \rightarrow 0$,
accompanied by in-plane insulating behavior
underdoped LSCO, but universal?

ρ_c diverge with $T \rightarrow 0$ (Bi2201) but not for ρ_a

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