



INTERNATIONAL ATOMIC ENERGY AGENCY  
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**INTERNATIONAL CENTRE FOR THEORETICAL PHYSICS**  
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H4-SMR 1012 - 33

## AUTUMN COLLEGE ON PLASMA PHYSICS

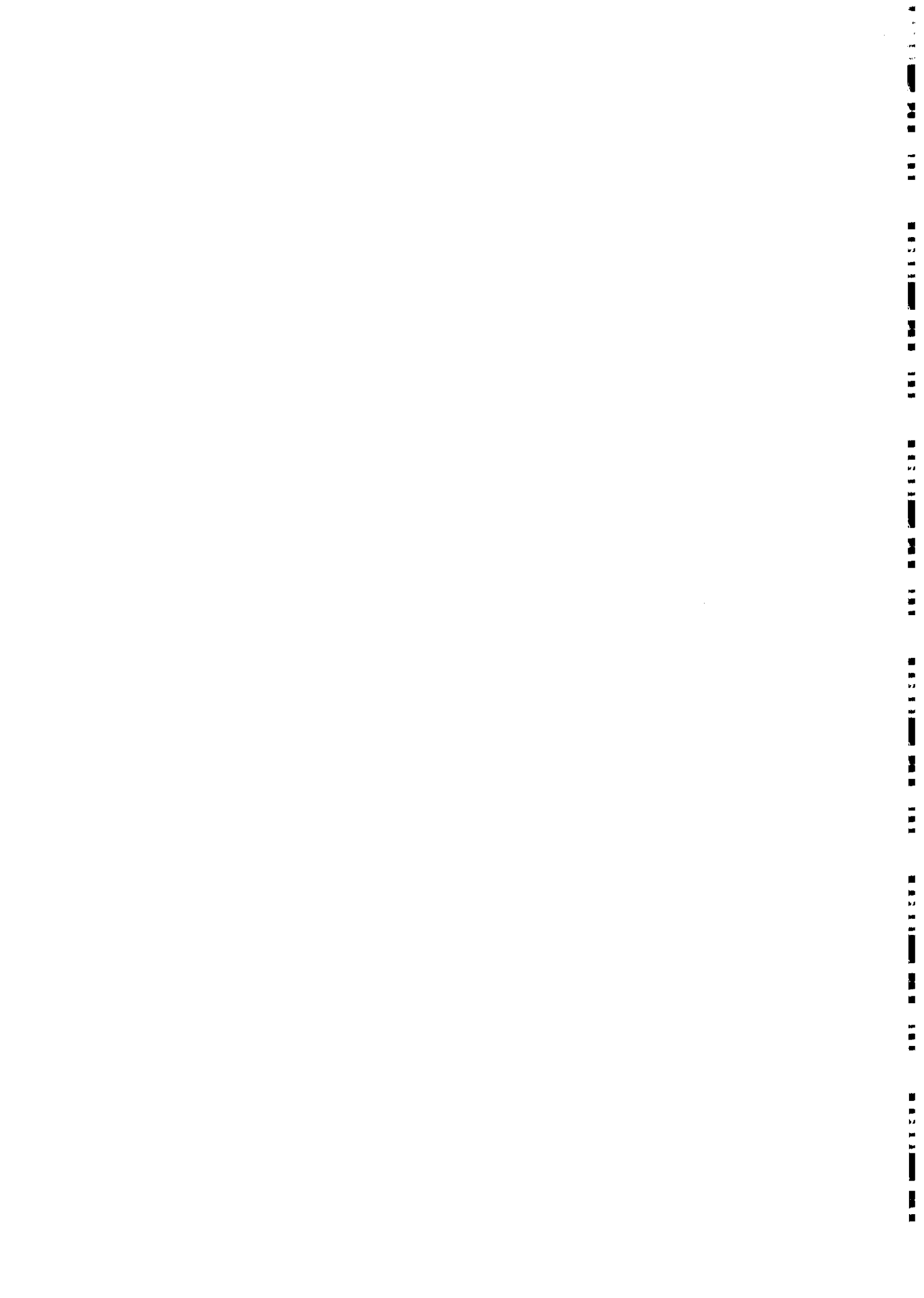
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### ELECTRON SCATTERING AND LOSS FROM THE INNER MAGNETOSPHERE

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These are lecture notes, intended for distribution to participants.



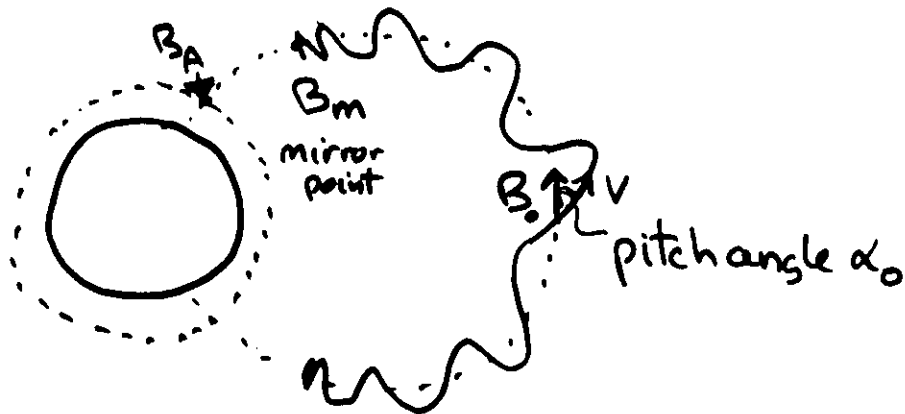
# Electron Scattering and Loss from the Inner Magnetosphere.

R. M. Thorne      UCLA

B. Abel      Olympi College, WA.

- ① Review of Adiabatic Motion
- ② Non Adiabatic Processes : Injection and Loss from the Radiation Belts
- ③ Resonant Wave-Particle Scattering
- ④ Equilibrium Pitch-Angle Distributions
- ⑤ Electron Lifetimes : Dominant Waves

## Geomagnetic Trapping :

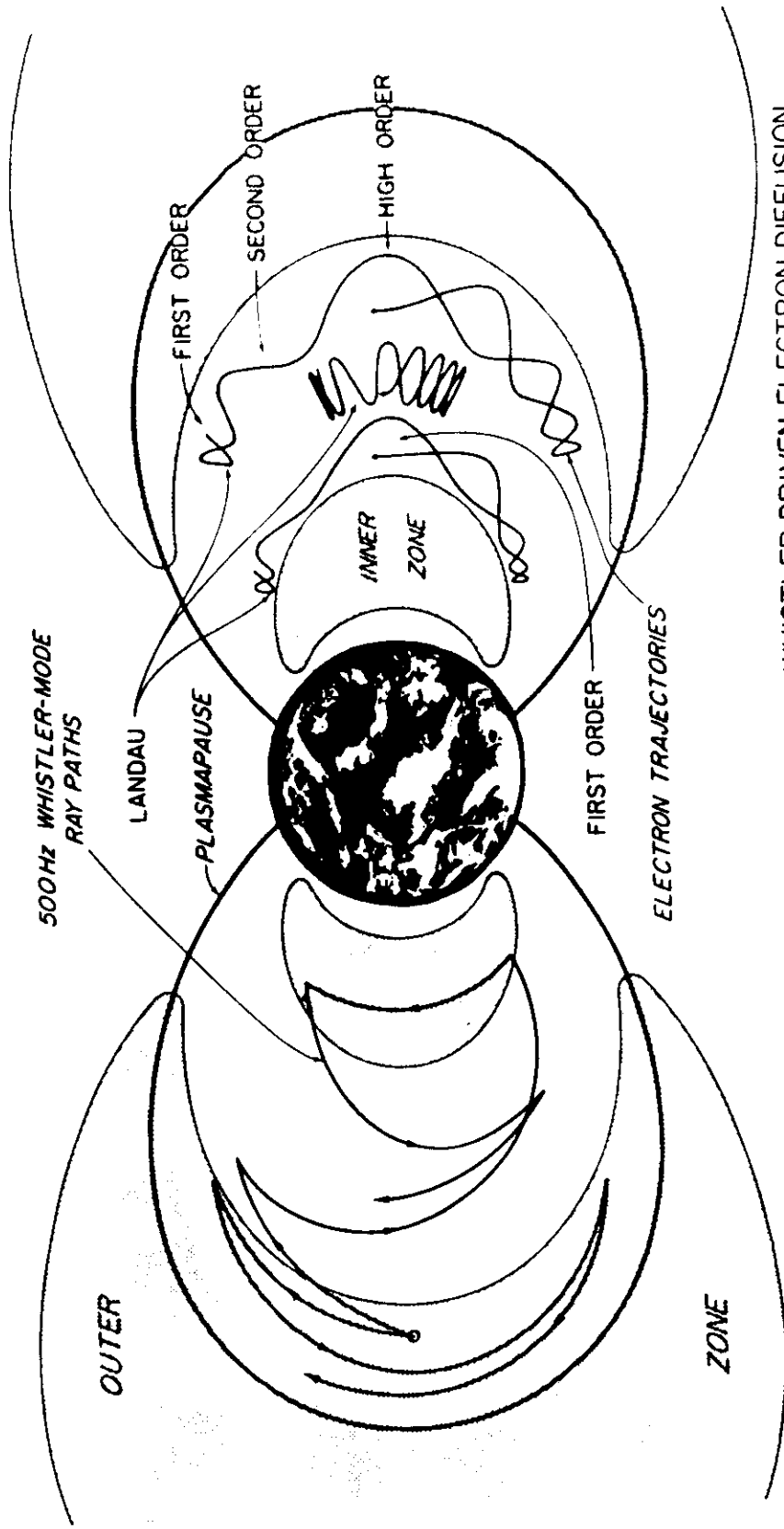


Conservation of 1<sup>st</sup> invariant  $\mu = E_v/B \Rightarrow \boxed{\frac{\sin^2 \alpha_0}{B_0} = \text{const.} = \frac{1}{B_m}}$

Loss Cone : due to collisions in atmosphere below a critical altitude ( $\leq 100$  km)

For  $\alpha < \alpha_c$  such that  $\boxed{\frac{\sin^2 \alpha_c}{B_0} = \frac{1}{B_a}}$

the particle mirror point will lie below the top of the atmosphere! Such particles will be removed by collisions. Generally, the loss cone will be empty.  $\boxed{f(\alpha < \alpha_c) \Rightarrow 0}$



### WAVE PROPAGATION

### WHISTLER DRIVEN ELECTRON DIFFUSION

Illustration of plasmaspheric whistler-mode wave propagation and electron pitch angle diffusion driven by observed plasmaspheric whistler-mode hiss (from Lyons *et al.*, 1972). The left-hand side of the figure shows examples of 500 Hz whistler-mode ray paths (obtained using the Stanford ray tracing program) which illustrate how waves generated near the plasmapause can propagate so as to fill the plasmasphere. The right-hand side of the figure illustrates the parasitic resonant interactions that occur between the plasmaspheric wave band and radiation belt electrons along three typical magnetospheric electron trajectories. Electrons near the inner edge of the slot (inner trajectory) are diffused primarily by first-order cyclotron resonance encountered near the equator. The diffusion of electrons in the outer regions of the slot (outer trajectory) is controlled by the combination of first order cyclotron resonance at high latitudes and higher harmonic resonances at lower latitudes. All radiation belt electrons encounter the Landau resonance near the mirror points, however such diffusion is most important for electrons of high equatorial pitch angles (middle trajectory).

**Lyons, Thorne, Kennel : SAR, '72**

## The rate of Pitch-Angle Scattering by Waves.

Under pure pitch-angle scattering (no energy change) the particle distribution function is controlled by:

$$\left[ \frac{\partial f}{\partial t} = \frac{1}{\sin \alpha} \frac{\partial}{\partial \alpha} \left( D_{\alpha\alpha} \frac{\partial f}{\partial \alpha} \right) + S - L \right] \text{loss}$$

Rate of Diffusion

Source

$$D_{\alpha\alpha} = \left\langle \frac{(\Delta \alpha)^2}{2 \Delta t} \right\rangle$$

$$D_{\alpha\alpha} \sim \Omega B_{\text{wave}}^2$$

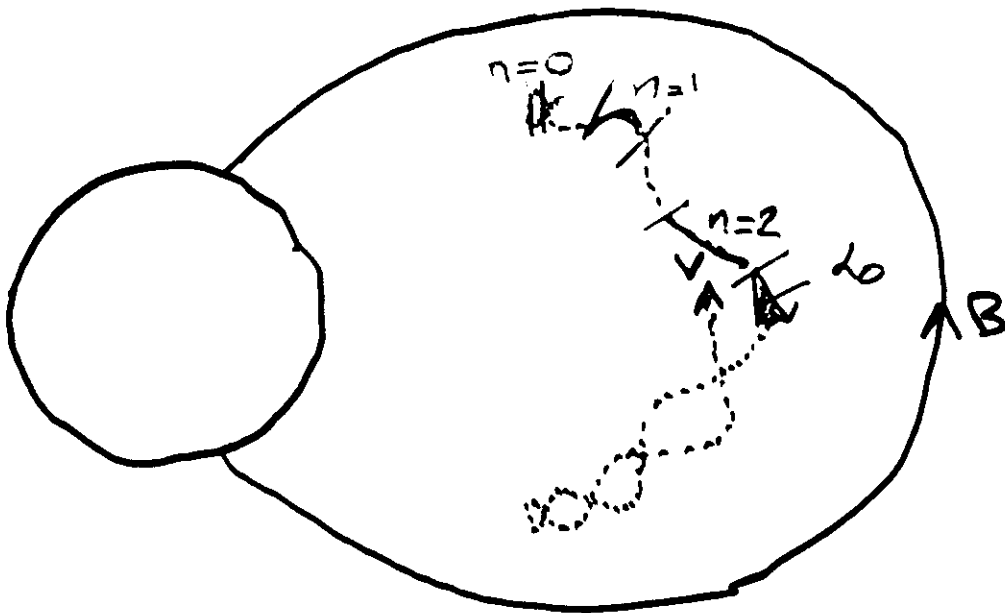
← This is controlled by the wave Power Spectral Density  $P(\omega, \alpha)$  and the order of resonance.

The theoretical formalism to compute  $D_{\alpha\alpha}$  has been described by Lyons, Thorne, Kennel; JGR, 77, 3455, 1972.

To utilize this for geomagnetically trapped electrons the rate of scattering must first be averaged over the particle's bounce orbit. Numerical calculations of  $D(\alpha_0)$  can then be used to calculate the particle lifetime  $\tau_c$  and the equilibrium distribution function  $f(\alpha_0)$  of electrons.

Since Coulomb lifetimes are very long for high energy electrons; wave scattering can be a dominant loss process.

# Bounced Averaged Rate of Scattering.



- ① Classify each electron in terms of  $\alpha_0$  at  $\lambda=0$ .
- ② Determine latitude range along bounce orbit for each harmonic resonance ( $n=0, 1, 2, \dots$ ) and evaluate  $D_{\alpha\alpha}(\lambda)$  for all resonant waves.
- ③ Obtain bounce orbit averaged diffusion rate

$$D(\alpha_0) = \frac{1}{T_B} \int_0^{T_B} D(t) \left( \frac{\partial \alpha_0}{\partial \alpha} \right)^2 dt$$

Bounce orbit.

$$D(\alpha_0) = \frac{1}{T(\alpha_0)} \int_0^{T_m} D(\lambda) \frac{\cos \alpha}{\cos^2 \alpha_0} \cos^2 \lambda d\lambda.$$

$$T(\alpha_0) \approx 1.3 - 0.56 \sin \alpha_0$$

... with  $\alpha$

$$\frac{\sin^2 \alpha(\lambda)}{B(\lambda)} = \frac{\sin^2 \alpha_0}{B_0}$$

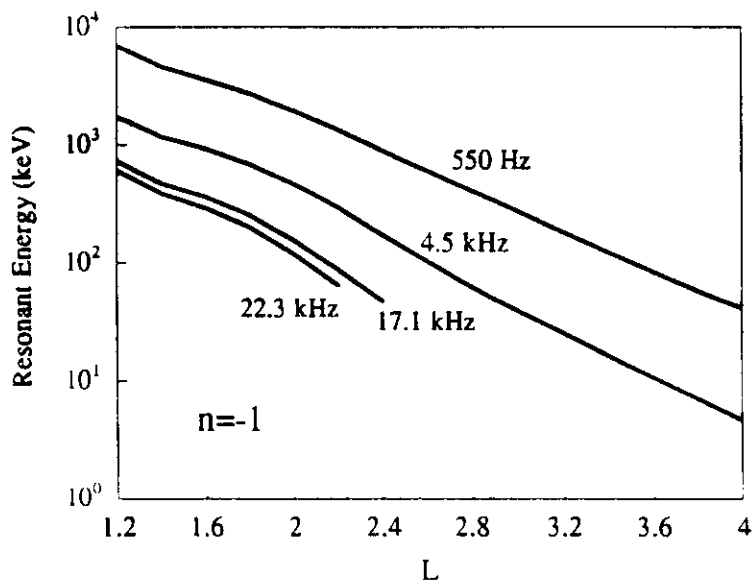
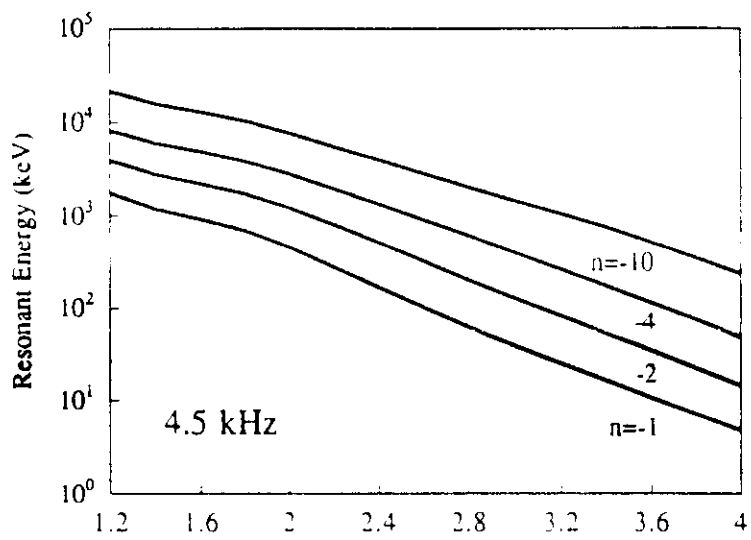
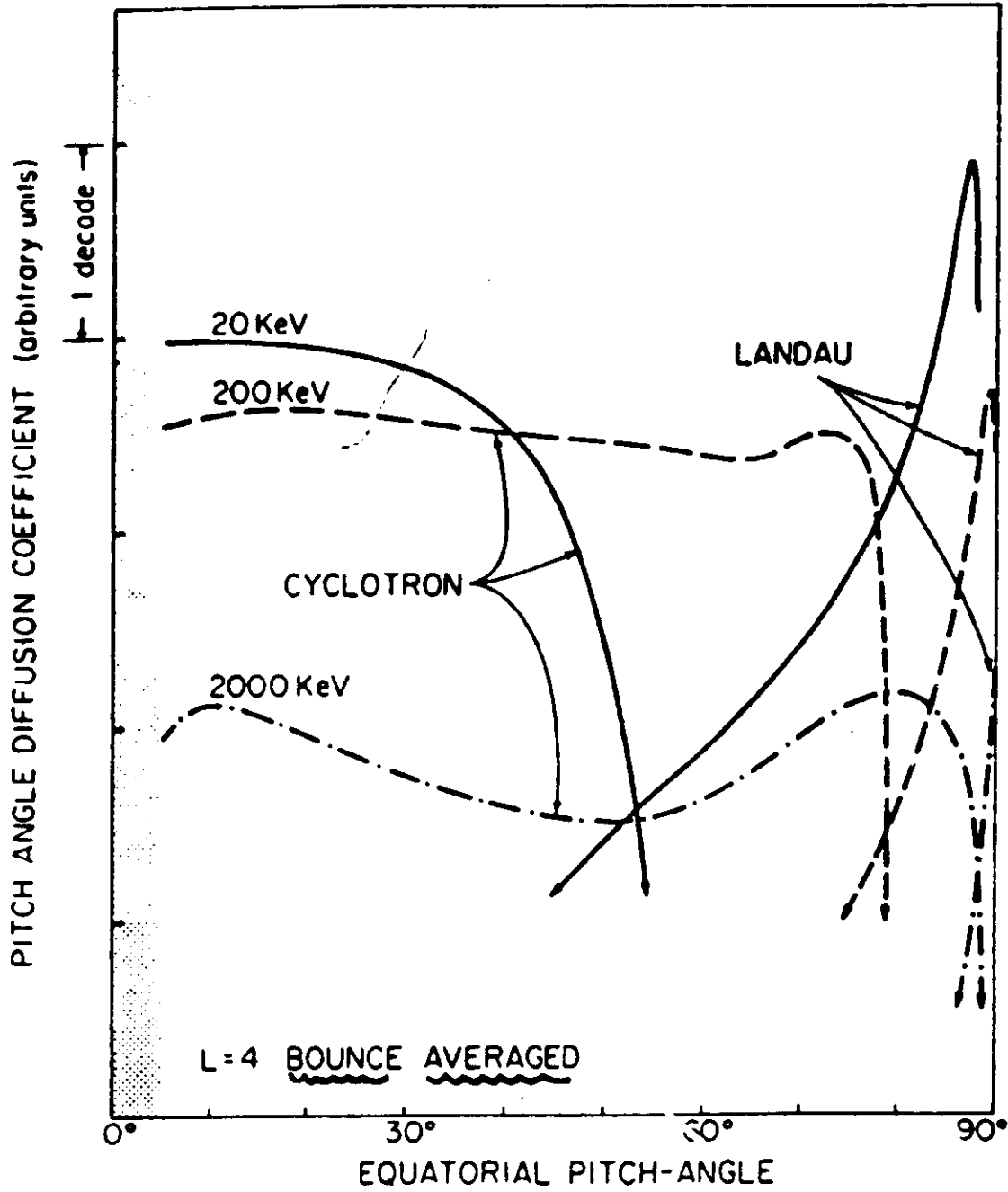


Fig. 2



Bounce Averaged  $D_{\parallel}(v)$  for plasmaspheric hiss  
in the outer plasmasphere.



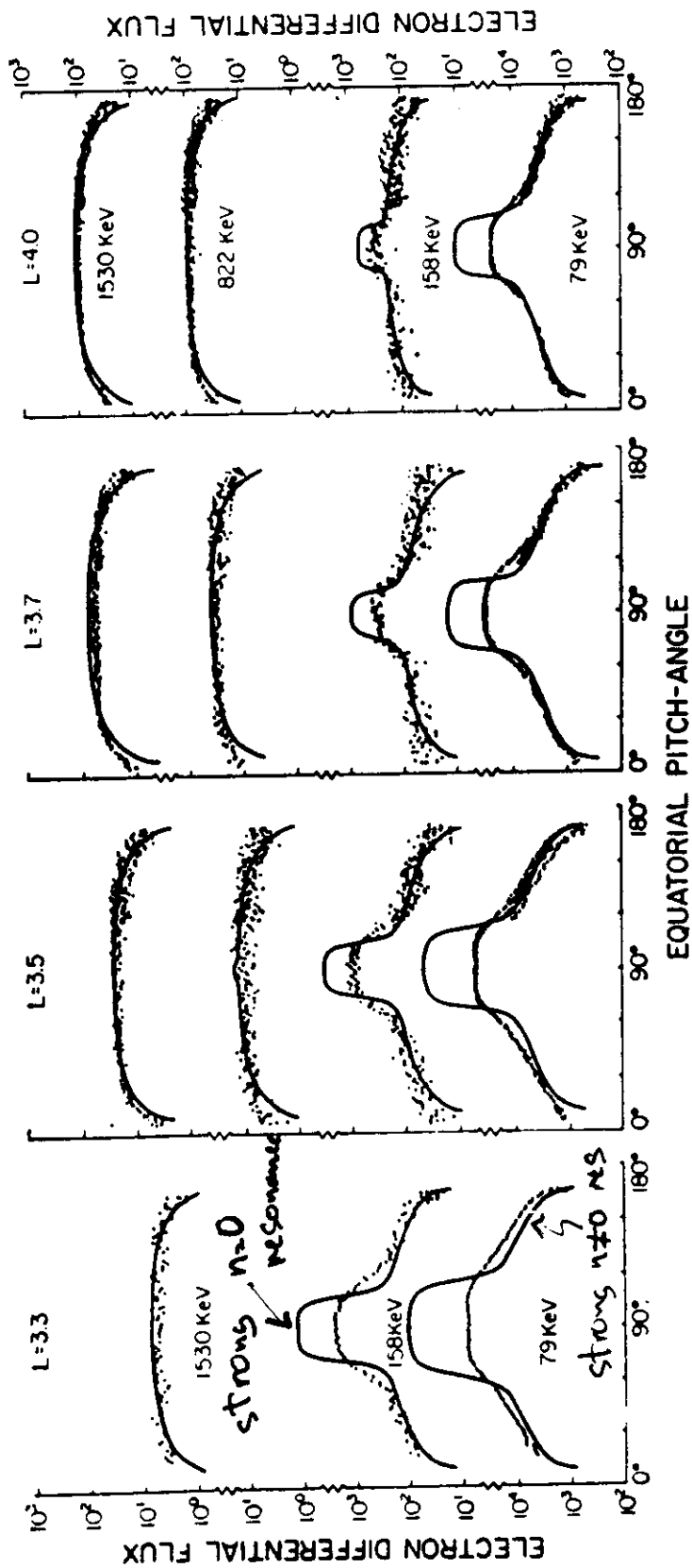


Fig. 6. Comparison of calculated equatorial pitch-angle distributions with equatorial distributions observed within the electron slot during the decay phase following an injection event (West, personal communication, 1971). The data points (electrons/cm<sup>2</sup>-sec-ster-kev) are plotted as a function of pitch angle at the  $L$  values and energies indicated on the figure. The vertical positionings of the corresponding calculated distributions, shown by solid lines, are arbitrary on a logarithmic scale and have therefore been adjusted so as to best fit the data.

Lyons et al 77, 3455, 1972.

## Wave Parameters

<u>Type of Wave</u>	<u>Plasmaspheric Hiss</u>	<u>Lightning-Generated Whistlers</u>	<u>VLF Transmitters</u>
Frequency	<u>550±300 Hz</u>	<u>4.5±2 kHz</u>	17.1 kHz ± 50 Hz and 22.3 kHz ± 50 Hz
Propagation Angle	45° ± 22.5°	45° ± 22.5°	45° ± 22.5°
Wave Intensity	10 my	10 my	Variable
Occurrence Rate	10% for L<2 ramps linearly to 50% at L=3, 50% for L≥3	( <sup>↑</sup> <i>Serutani et al.</i> ) 3%	2.4% each

Electron Energy Range: 100-1,500 keV

### Multi-Species Density Calculations:

Neutrals: MSIS90 [*Hedin*, 1991]

Electrons: [*Chiu et al.*, 1979]

Ions: IRI90 [*Rawer*, 1982] Anchored to Electrons

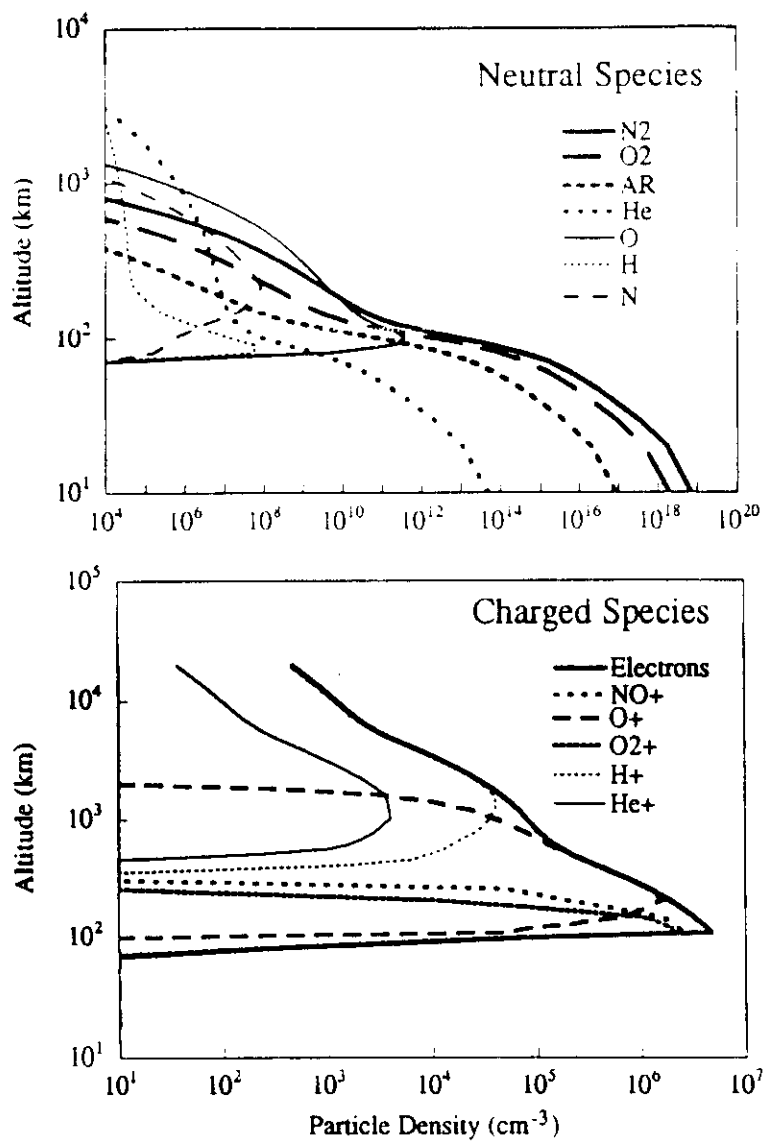


Fig. 1

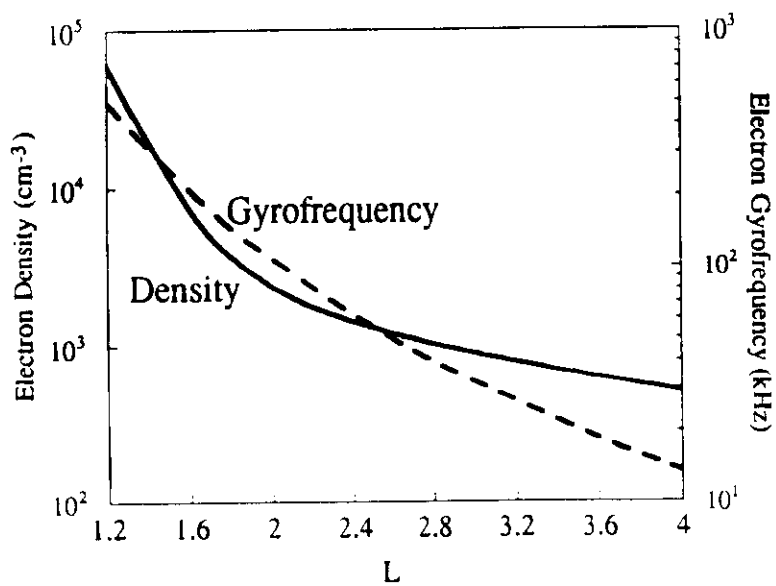
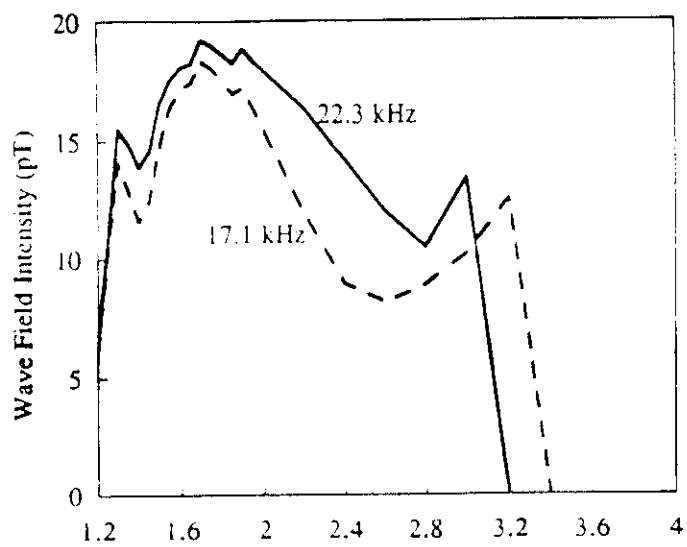


Fig.

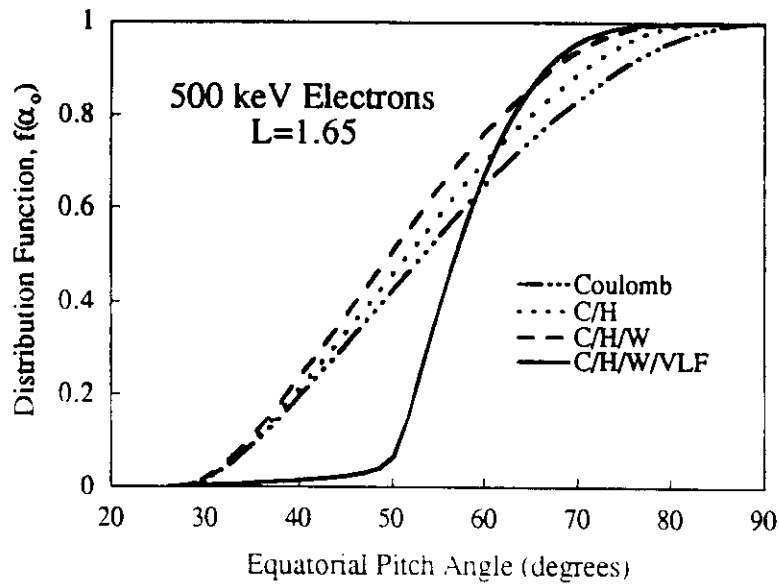
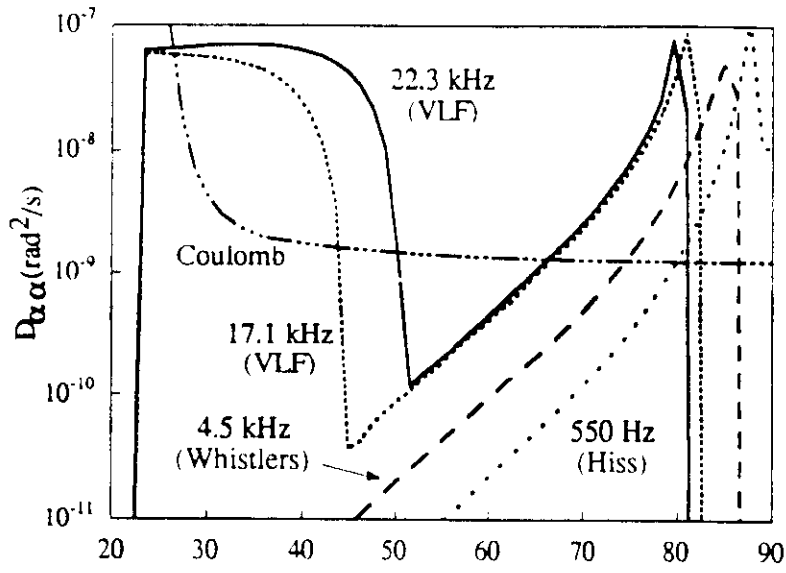
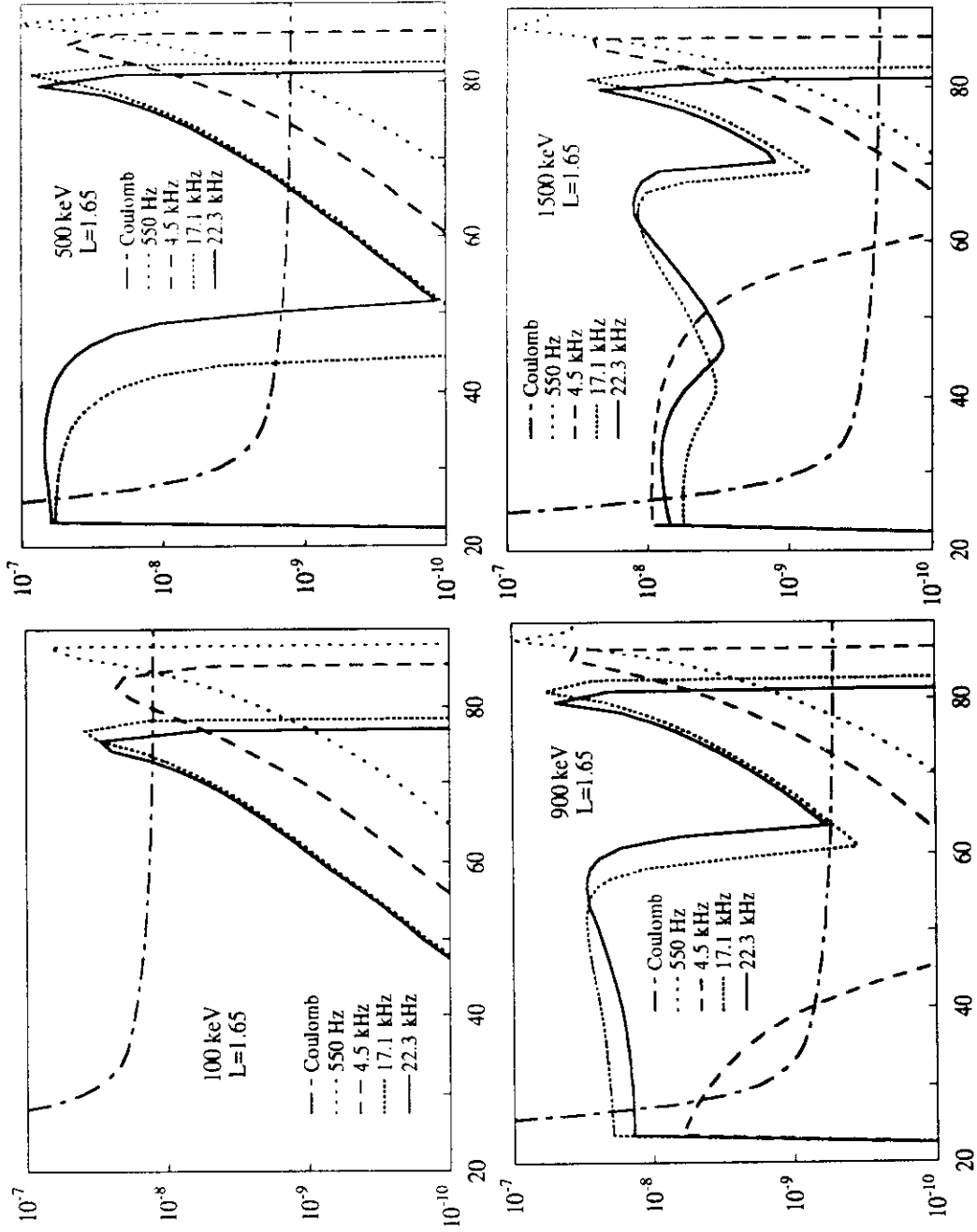
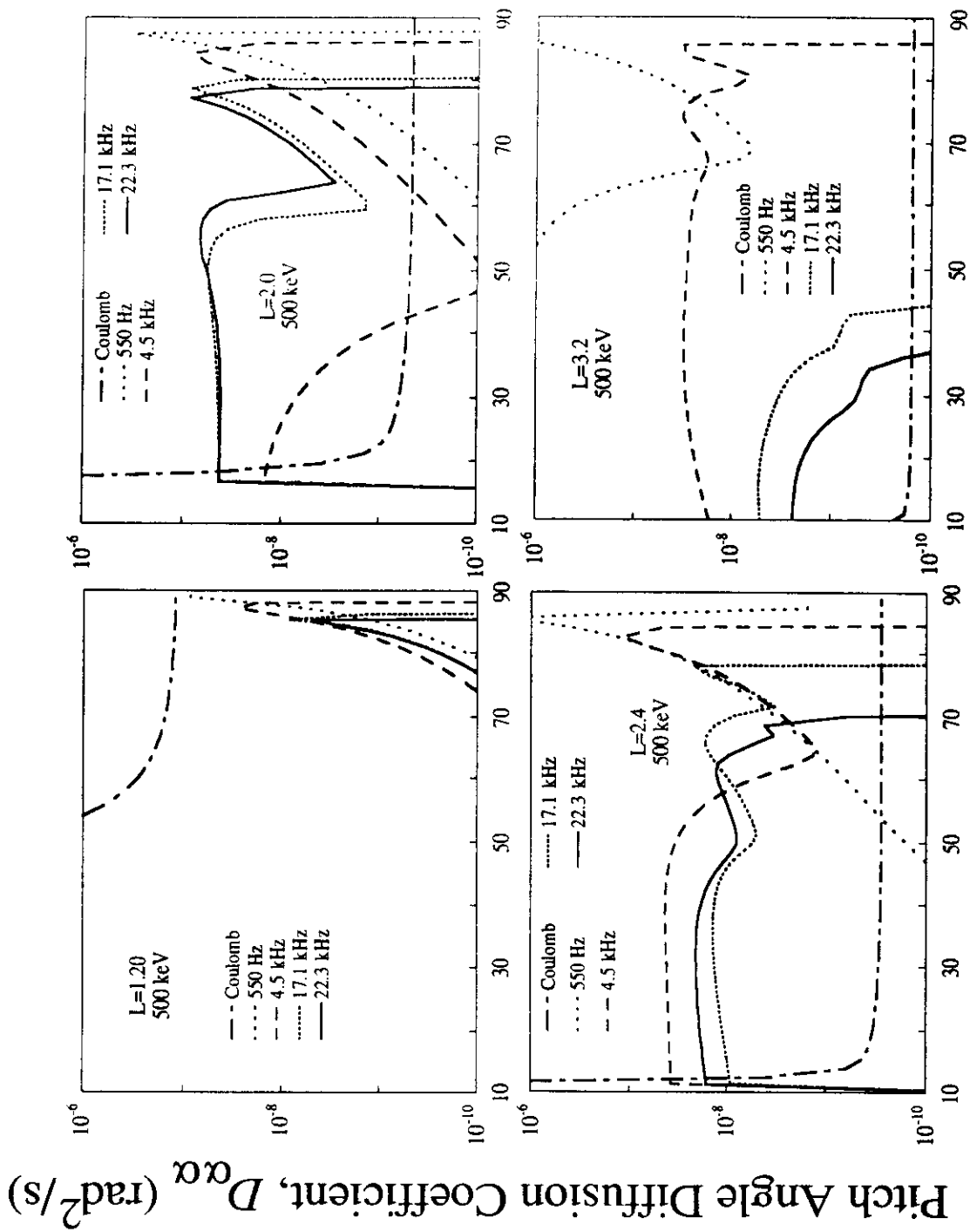


Fig. 7

Pitch Angle Diffusion Coefficient,  $D_{\alpha\alpha}$  ( $\text{rad}^2/\text{s}$ )



Equatorial Pitch Angle,  $\alpha_0$  (degrees)



Equatorial Pitch Angle,  $\alpha_0$  (degrees)



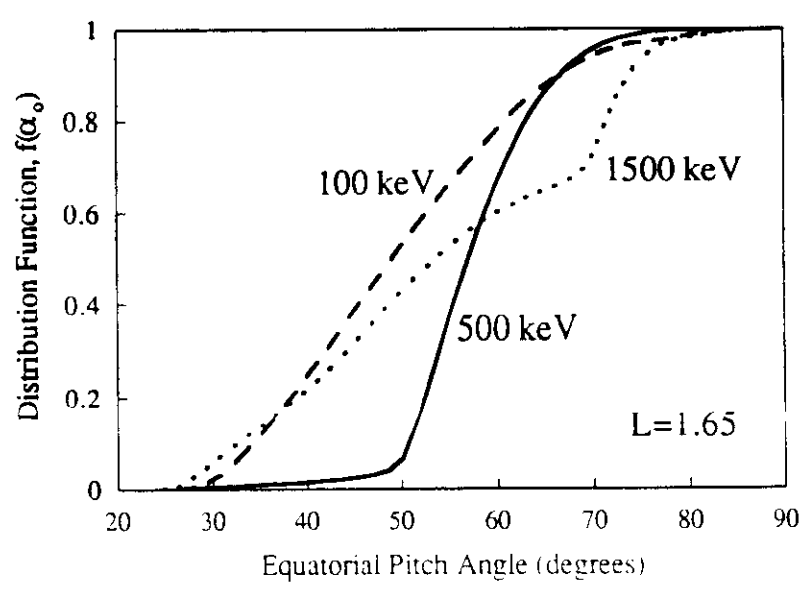
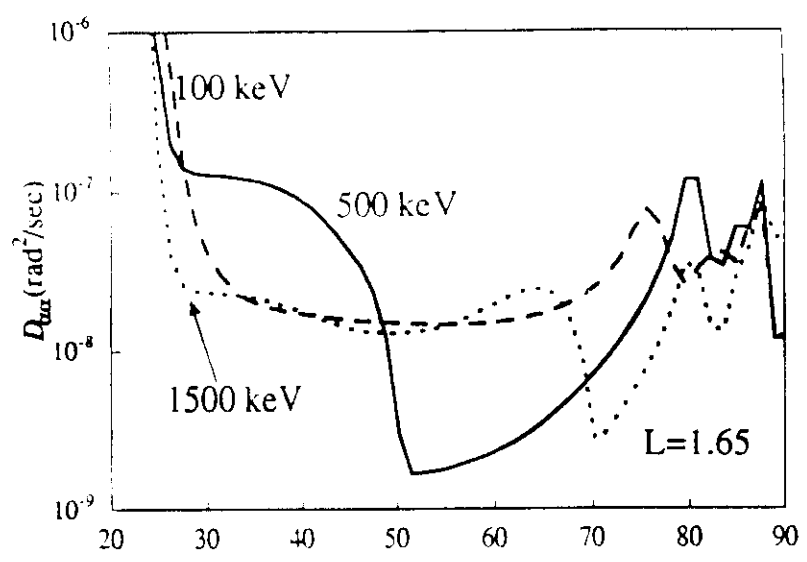


Fig. 7

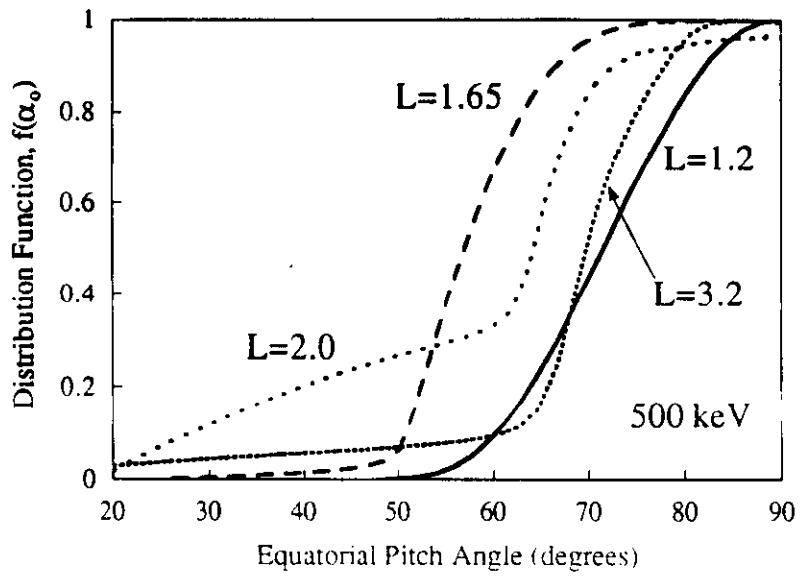
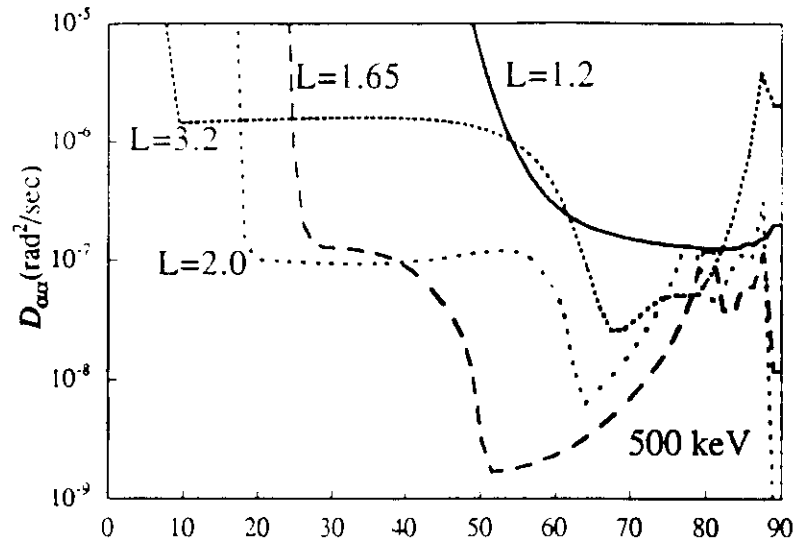


Fig. 6

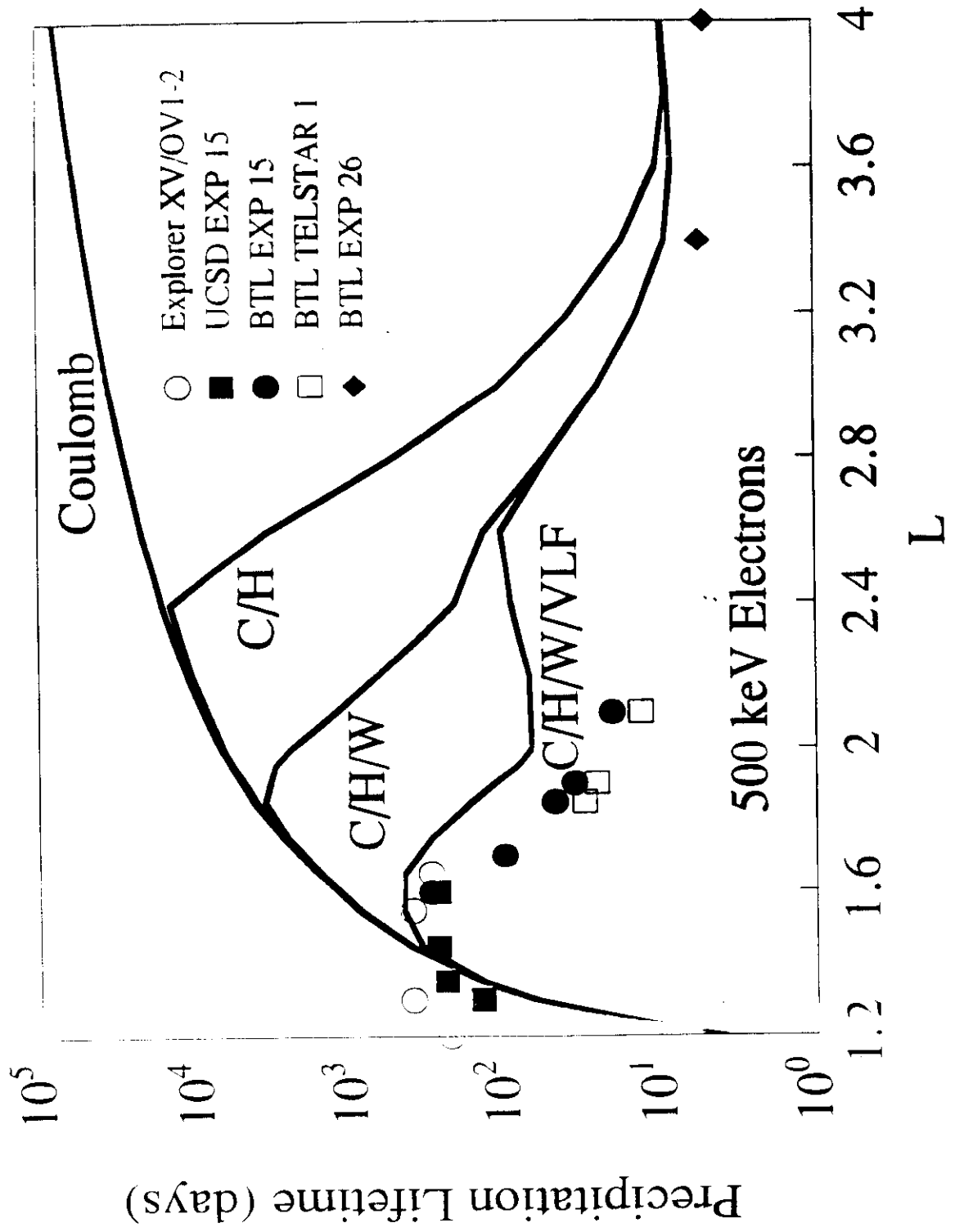
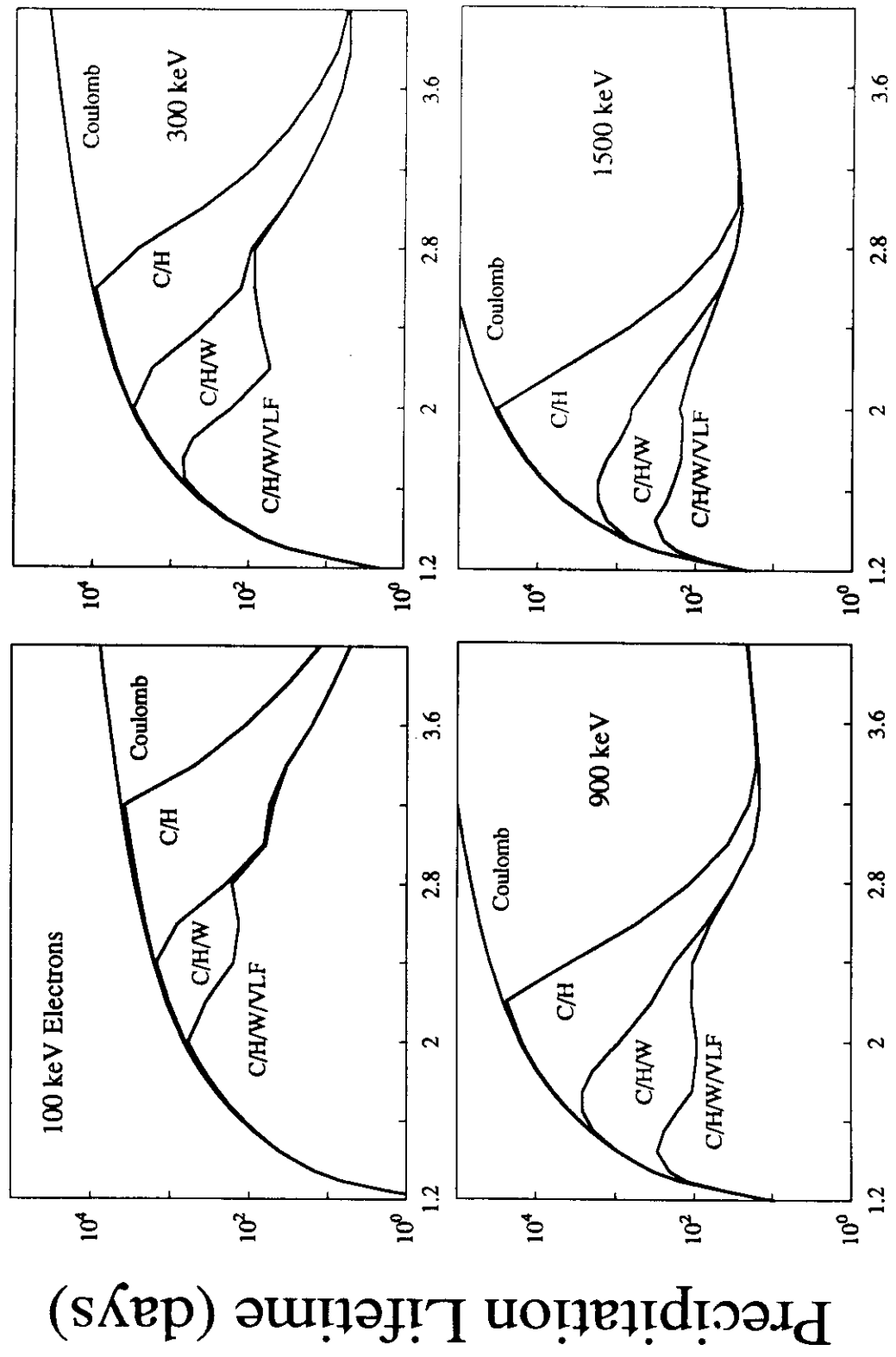


Fig. 1



L Value

Precipitation Lifetime (days)

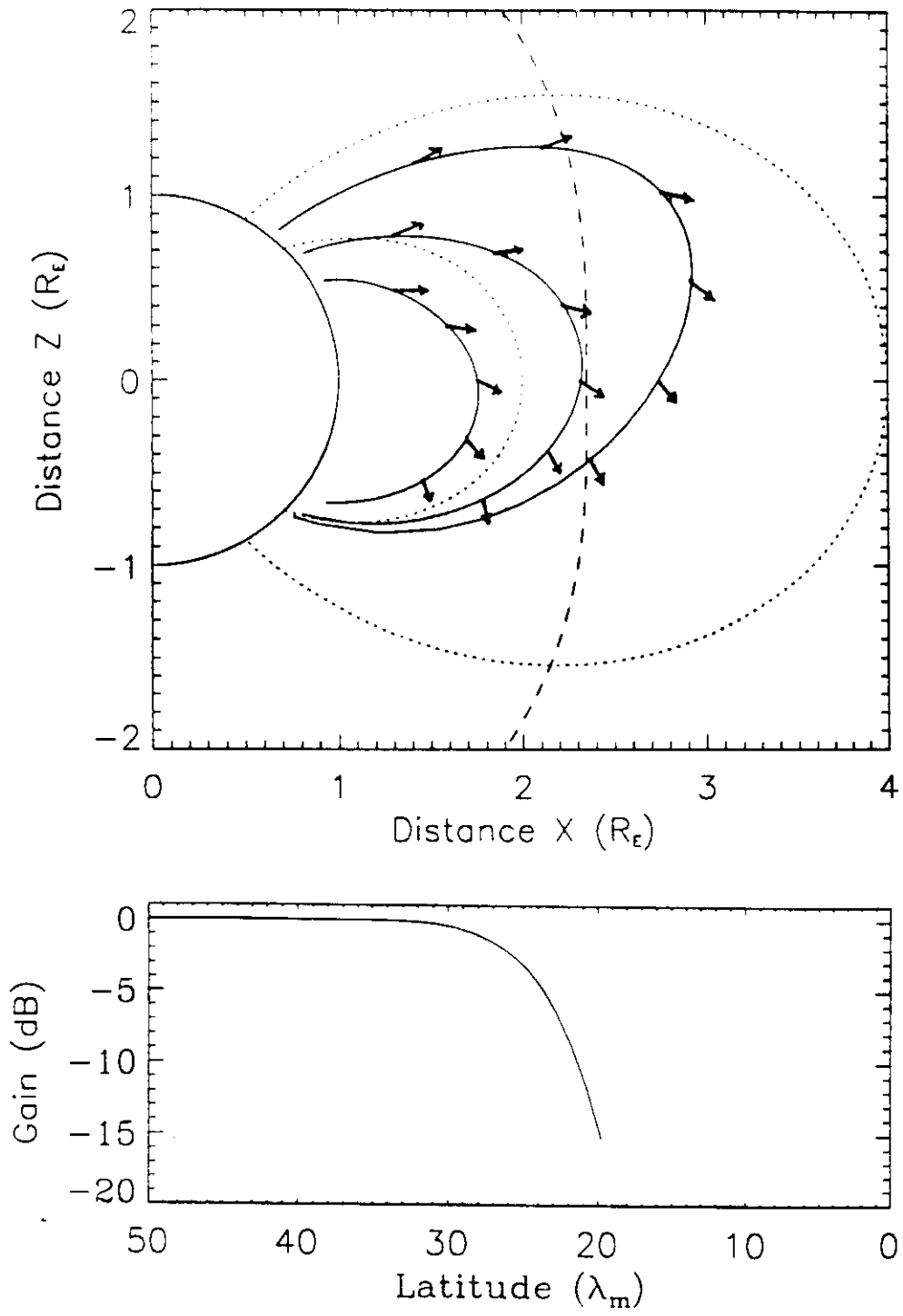


Fig. 1

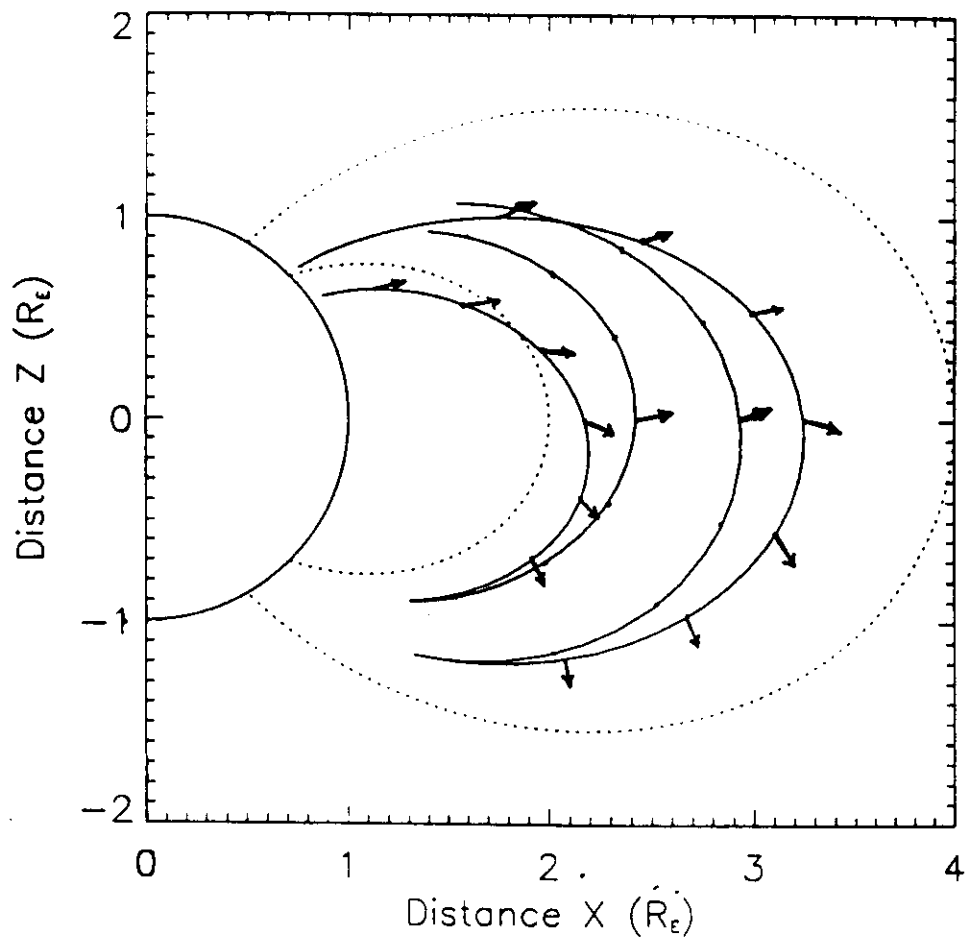


Fig. 2

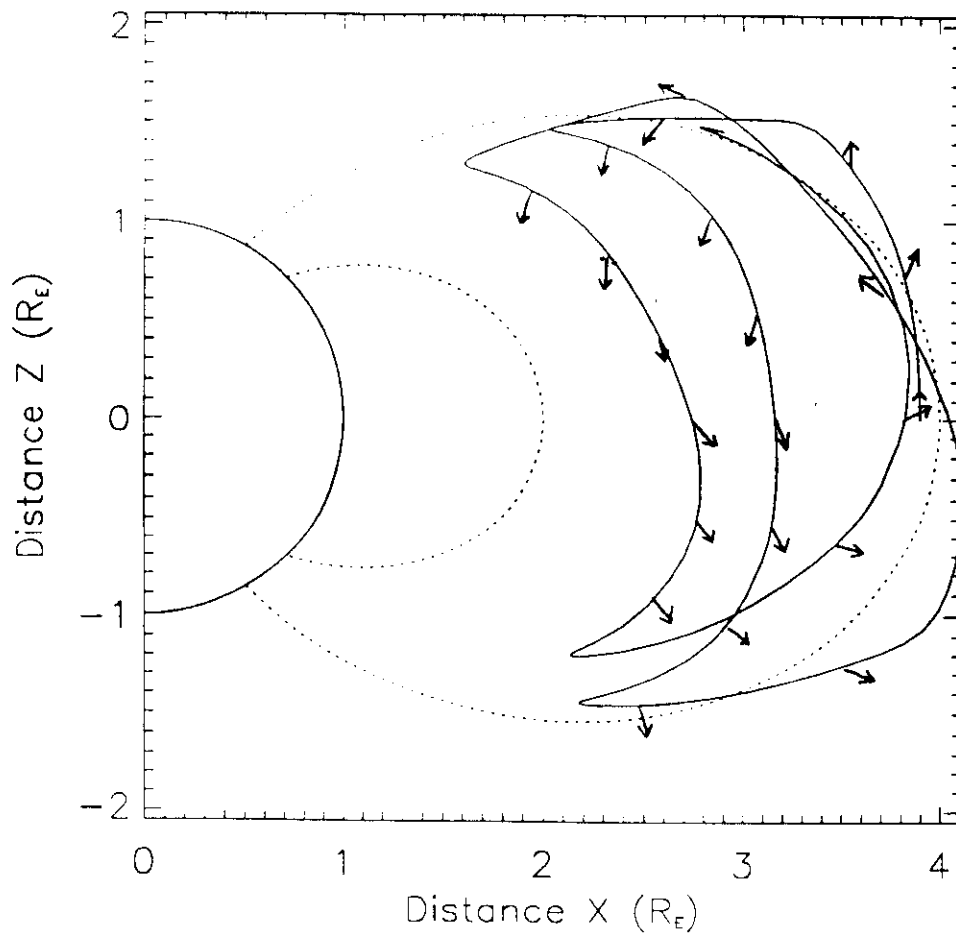


Fig. 3

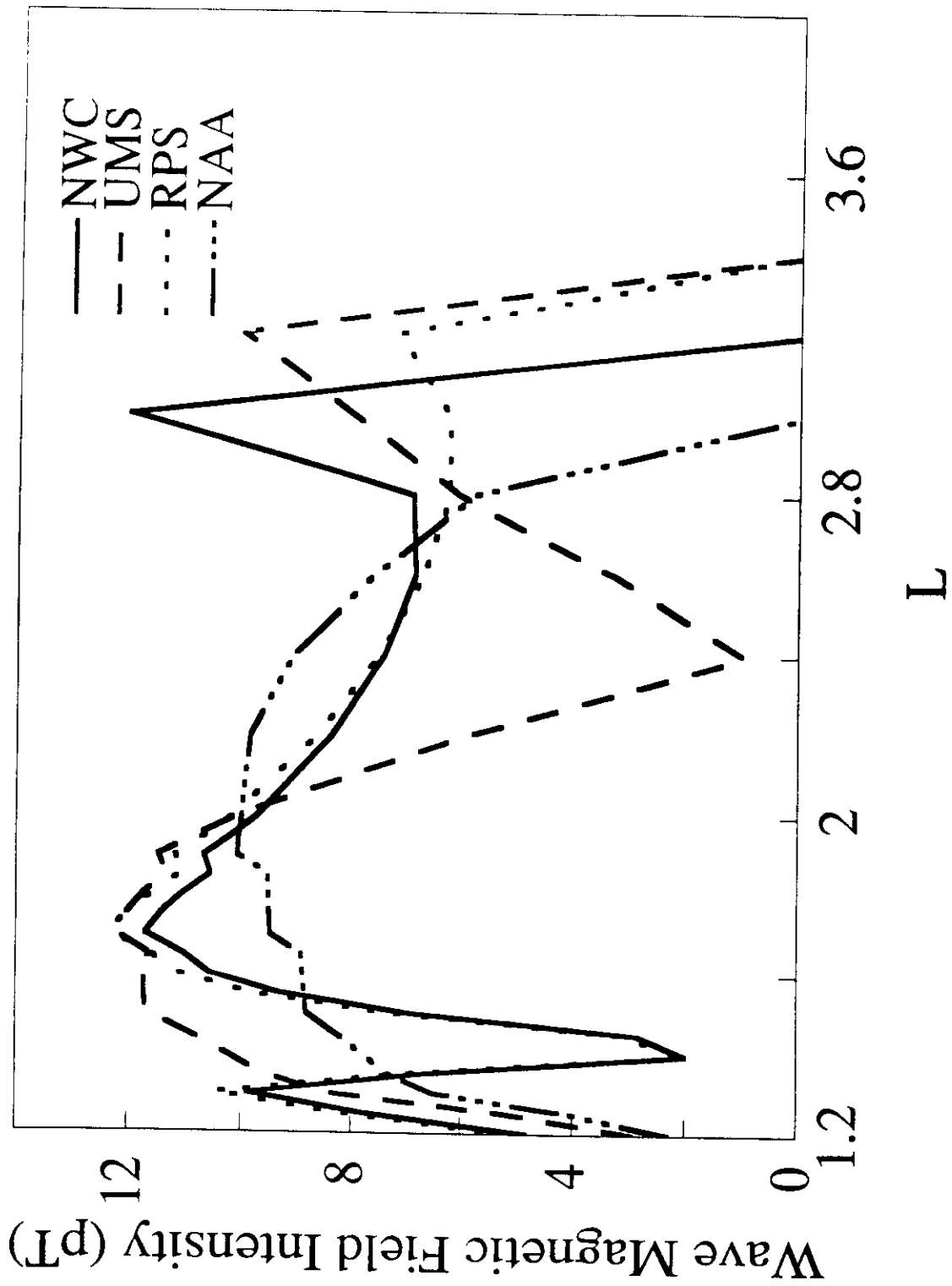
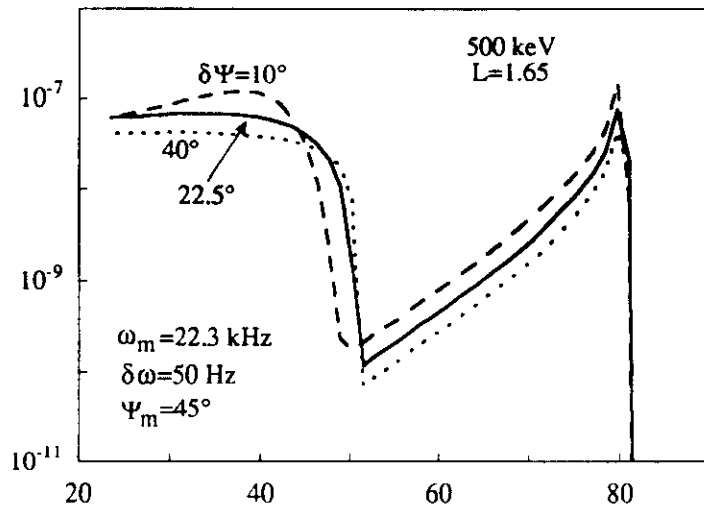
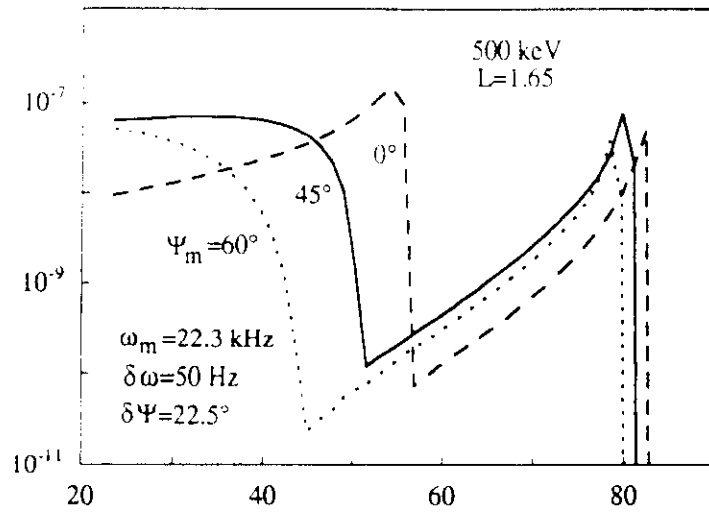


Fig. 4

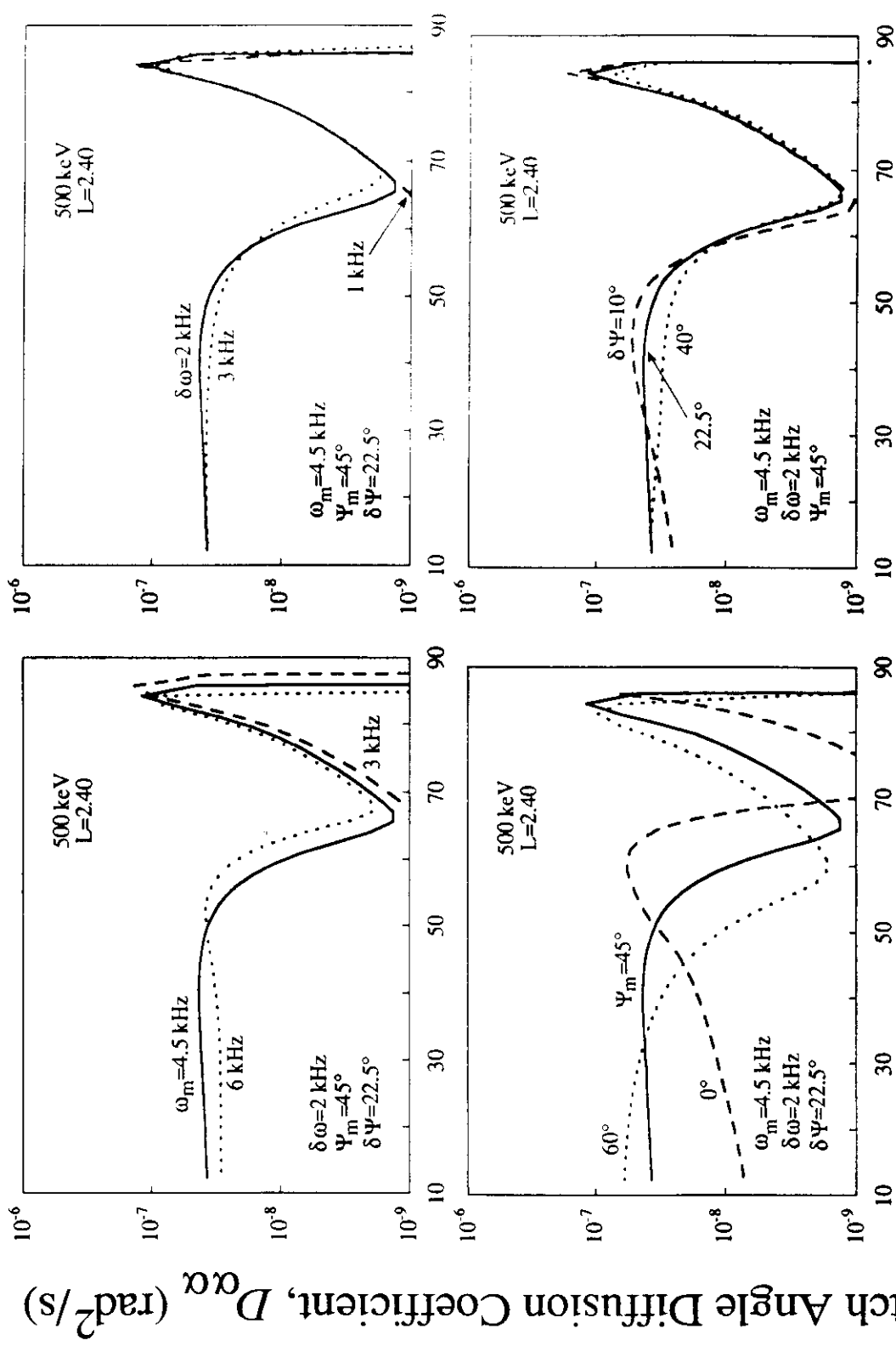


Pitch Angle Diffusion Coefficient,  $D_{\alpha\alpha}$  ( $\text{rad}^2/\text{s}$ )



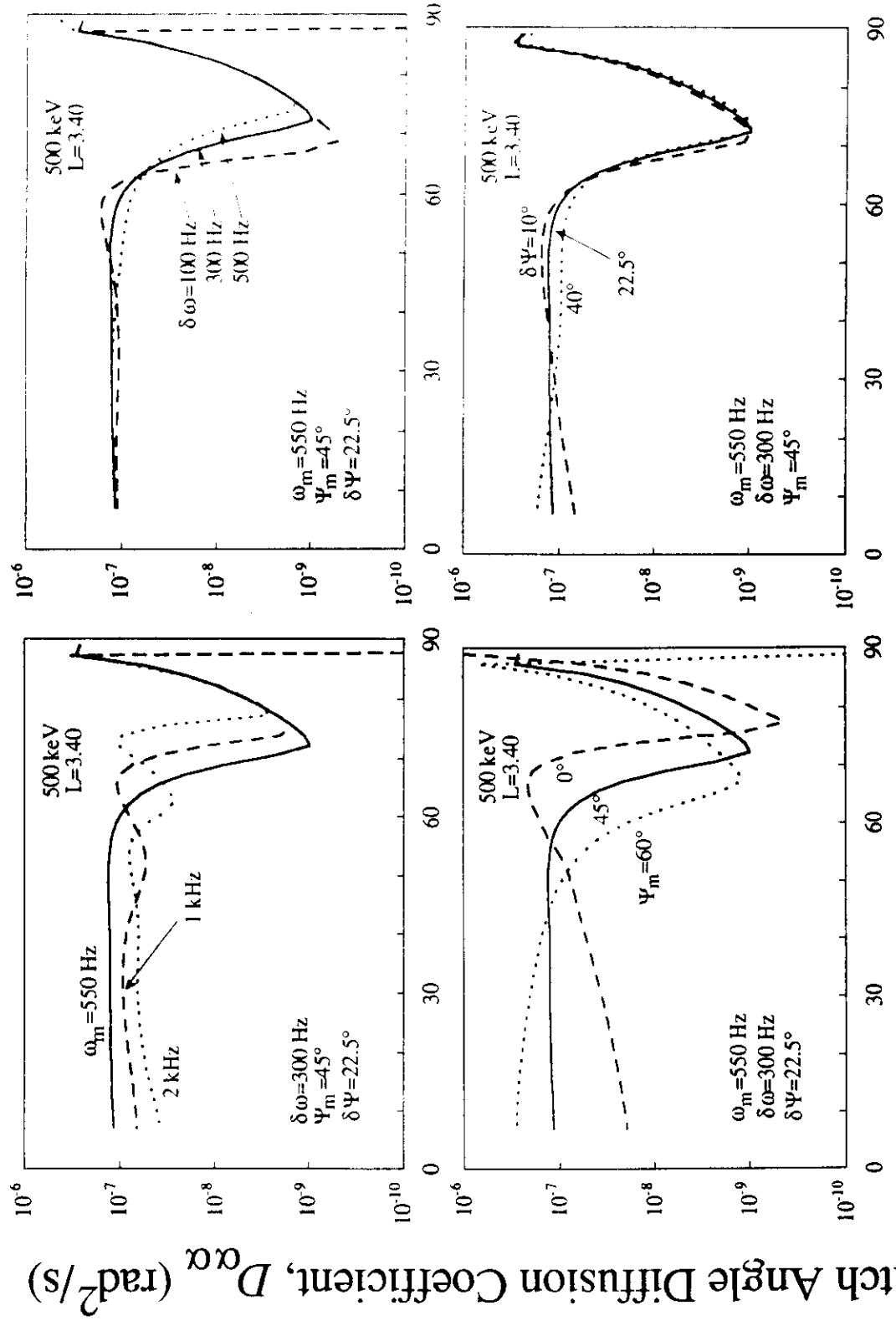
Equatorial Pitch Angle,  $\alpha_0$  (degrees)

Fig. 5



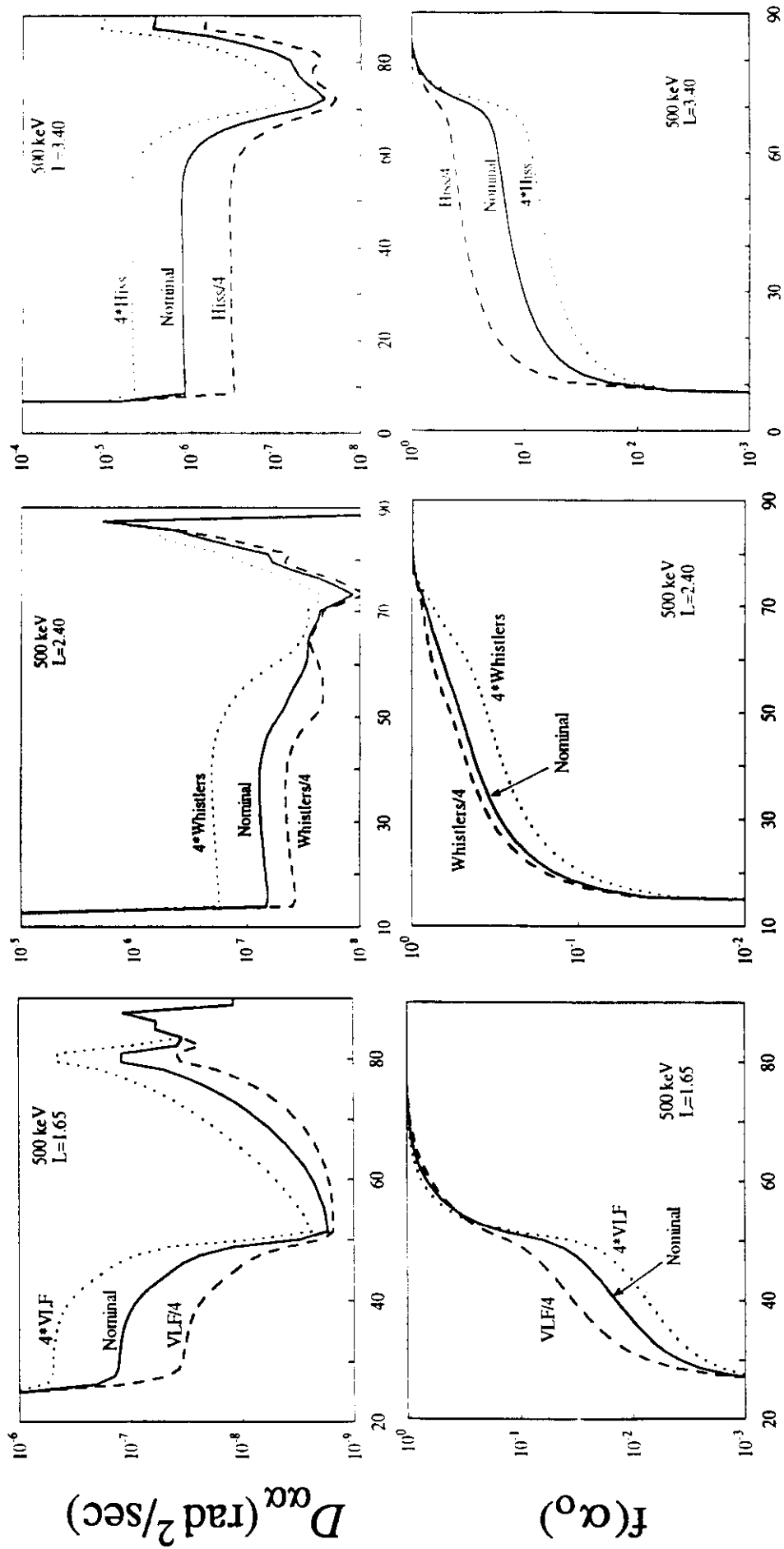
Pitch Angle Diffusion Coefficient,  $D_{\alpha\alpha}$  (rad<sup>2</sup>/s)

Equatorial Pitch Angle,  $\alpha_0$  (degrees)



Equatorial Pitch Angle,  $\alpha_0$  (degrees)

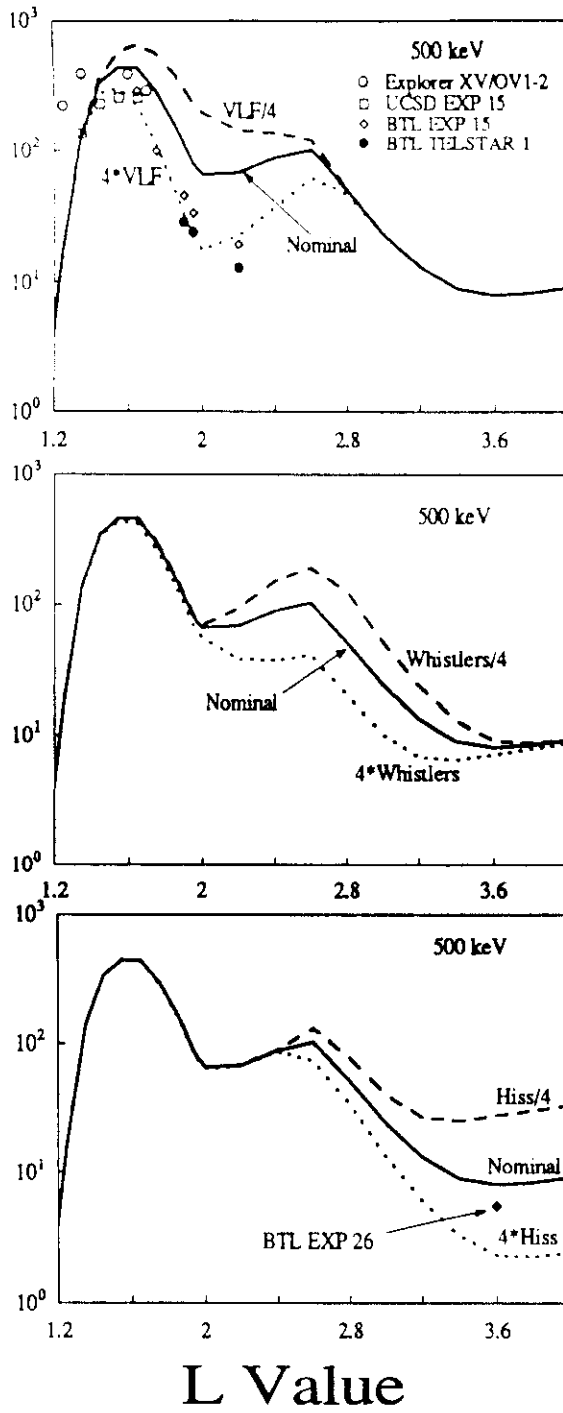
Fig. 7



Equatorial Pitch Angle,  $\alpha_0$  (degrees)

Fig. 8

Precipitation Lifetime (days)



Precipitation Lifetime (days)

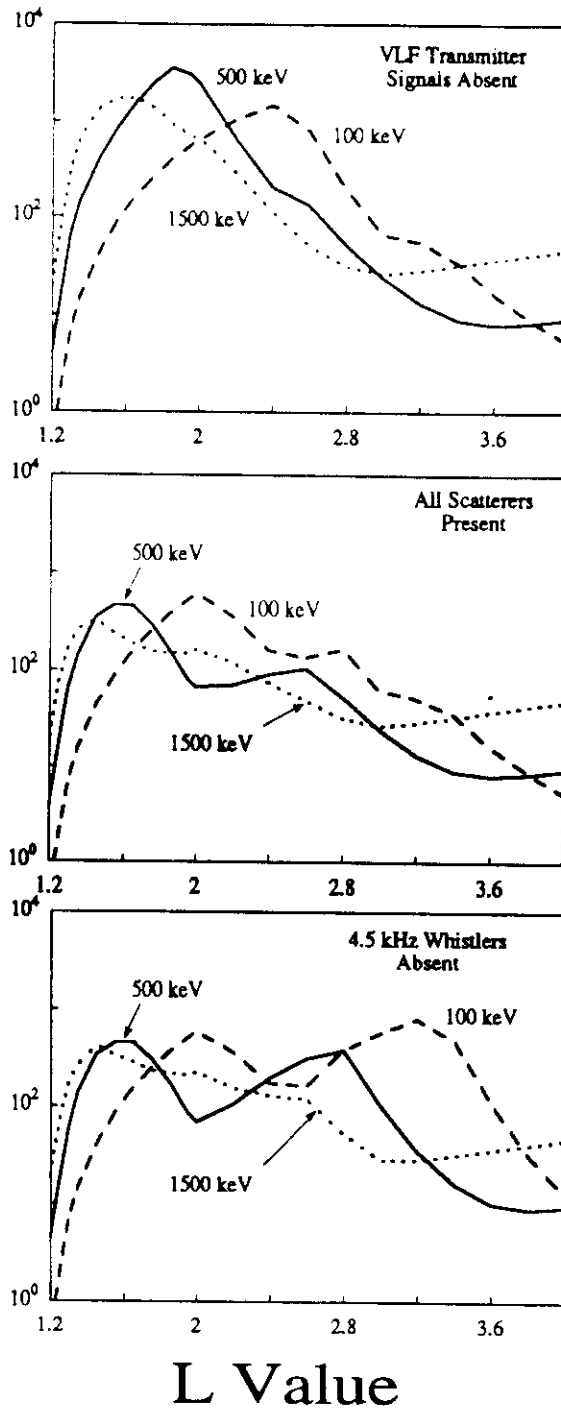
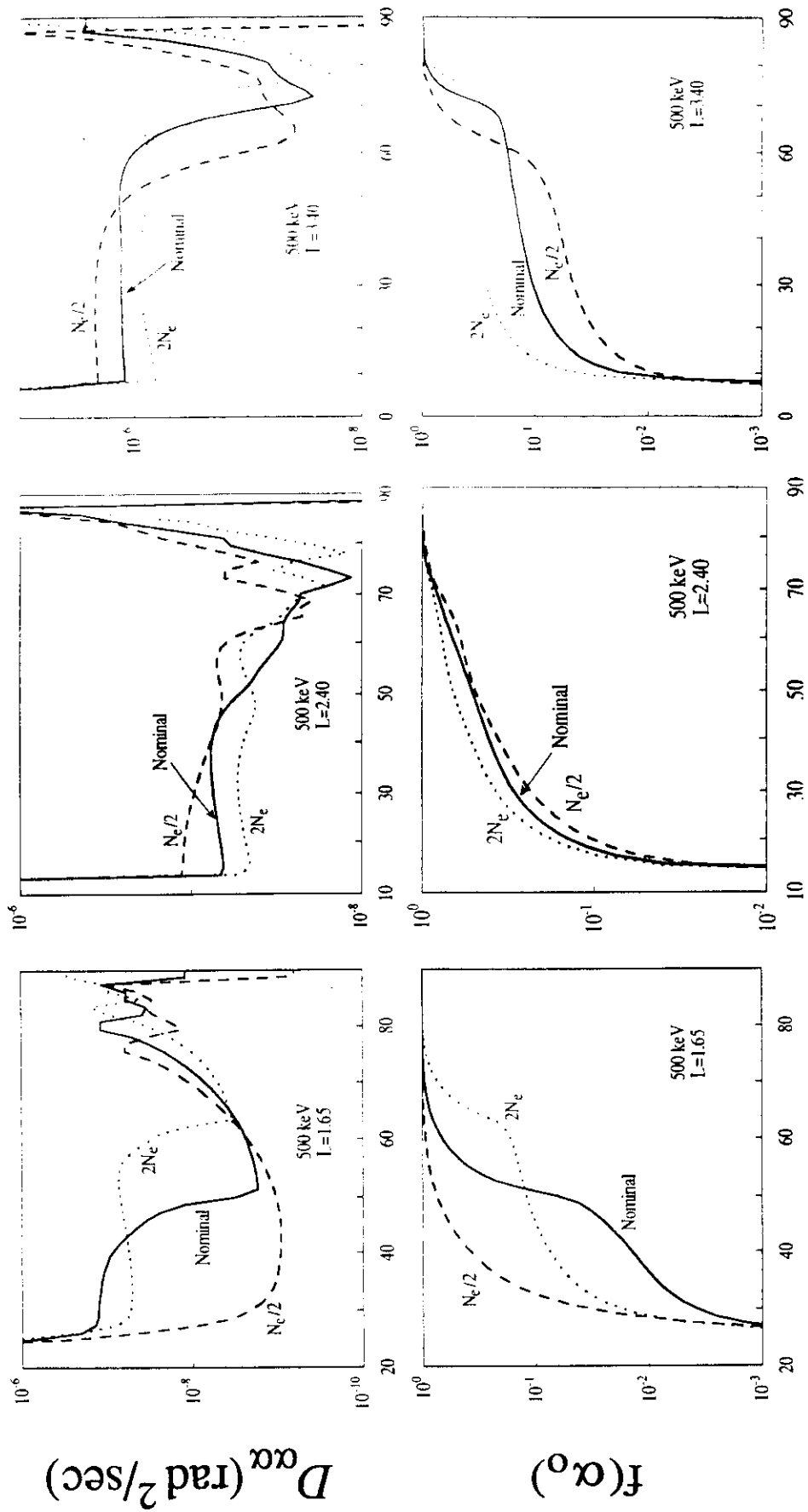


Fig. 10



Equatorial Pitch Angle,  $\alpha_0$  (degrees)

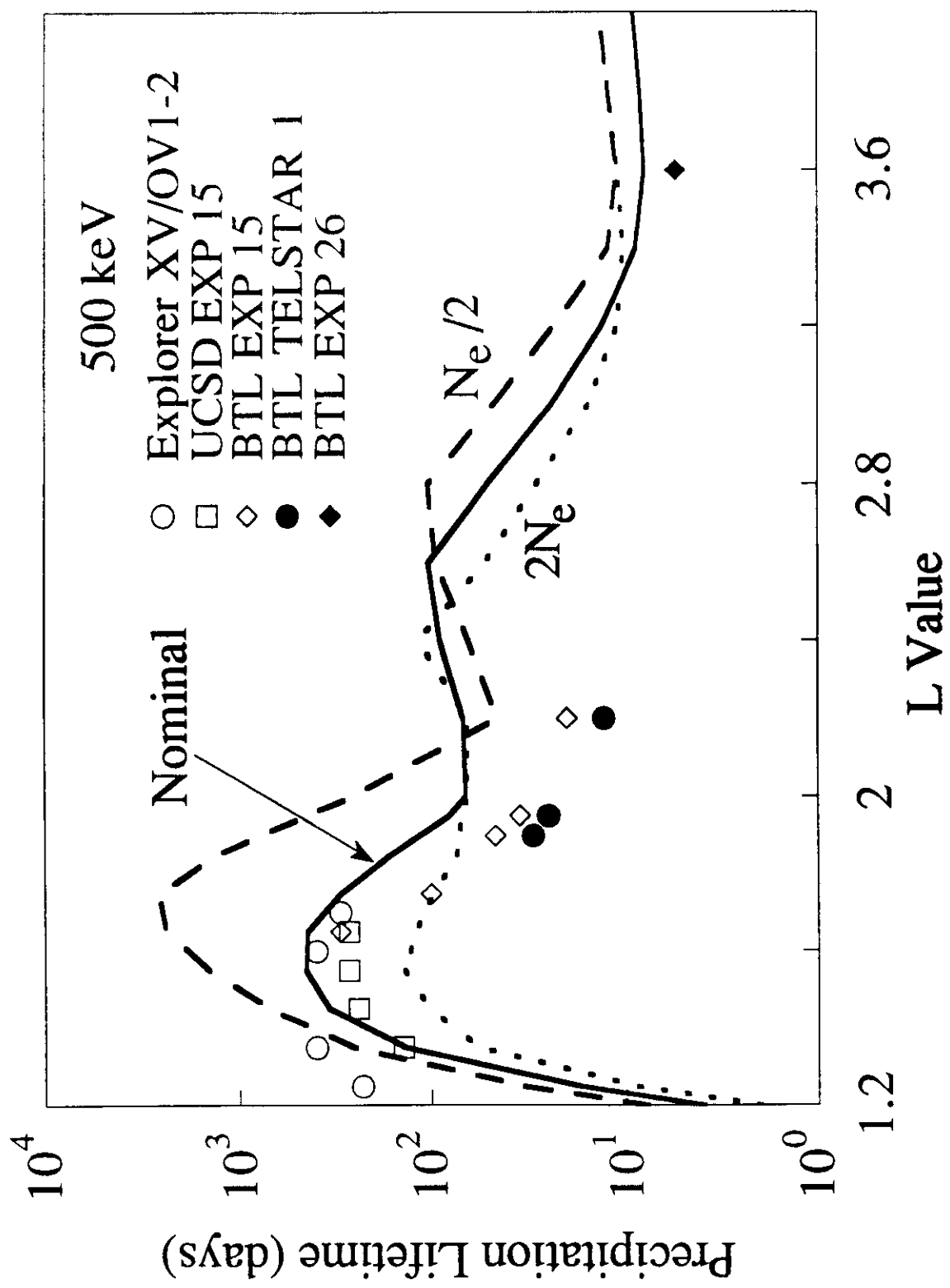


Fig. 12



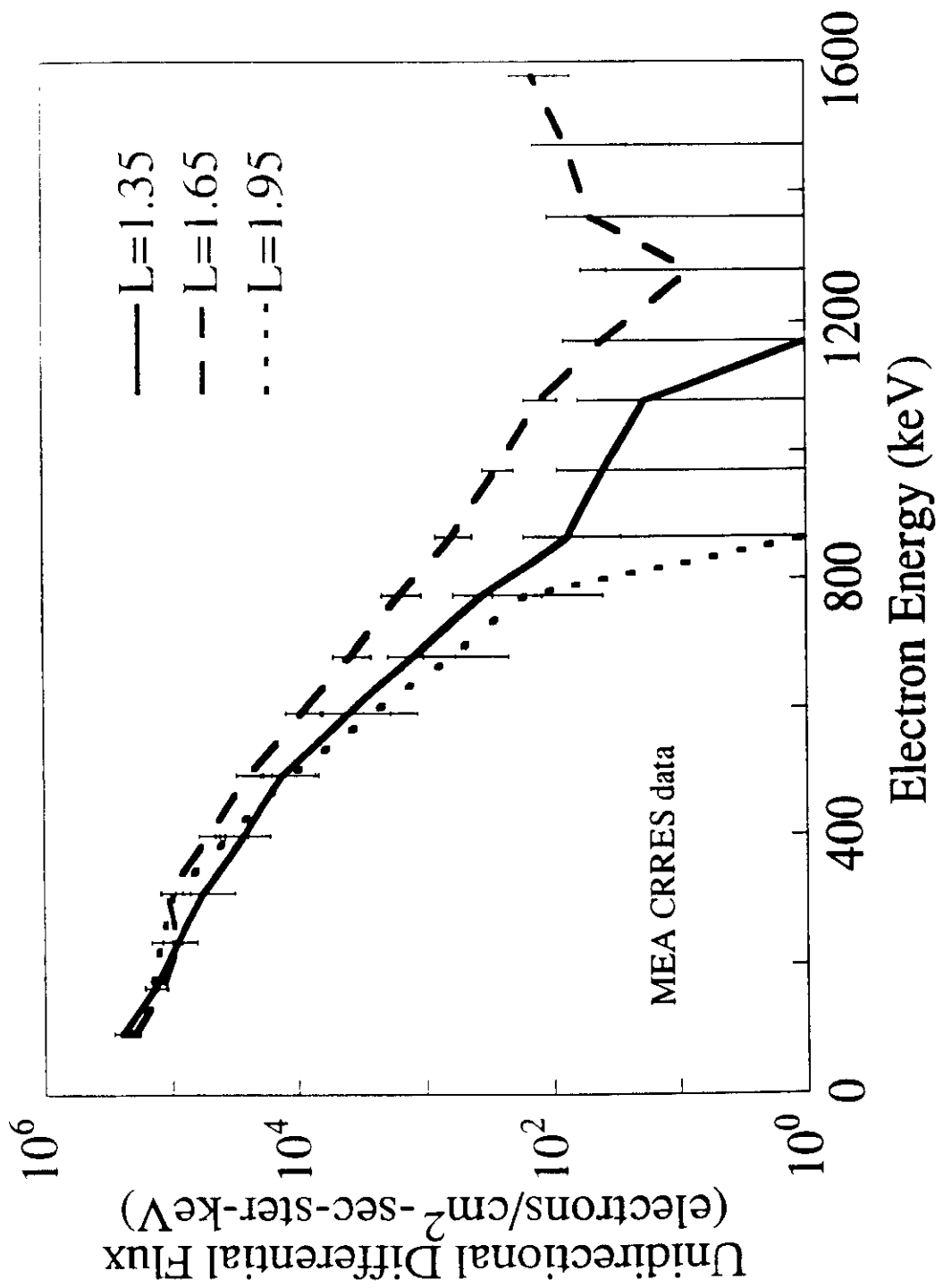


Fig. 1

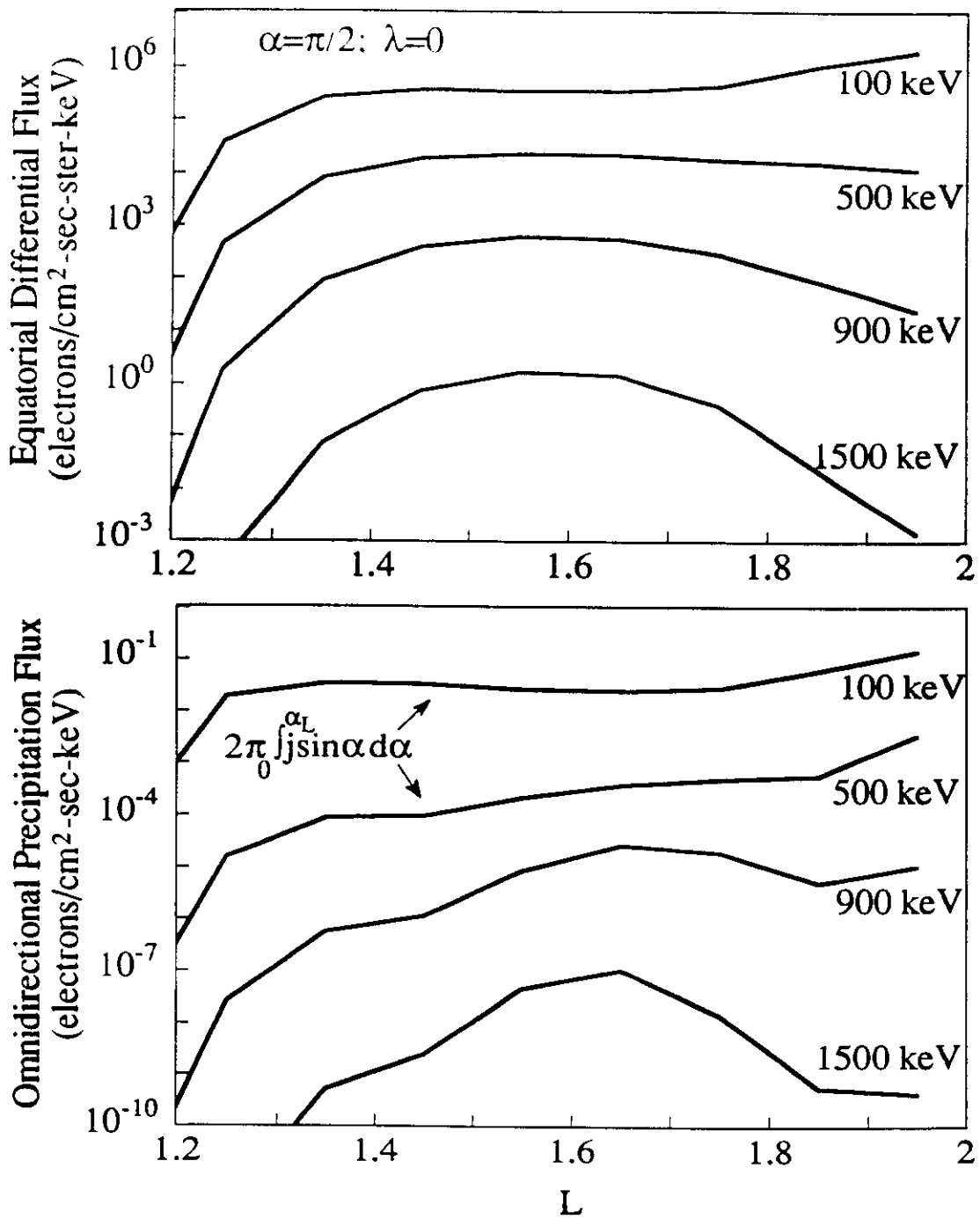


Fig. 2

