

*COLLEGE ON MEDICAL PHYSICS
AND
WORKSHOP ON
NUCLEAR DATA FOR SCIENCE AND TECHNOLOGY:
MEDICAL APPLICATIONS
(20 SEPTEMBER - 15 OCTOBER 1999)*

**"Physics and Applications of ETRAN and ITS
Electron-Photon Transport
Monte Carlo Codes"**

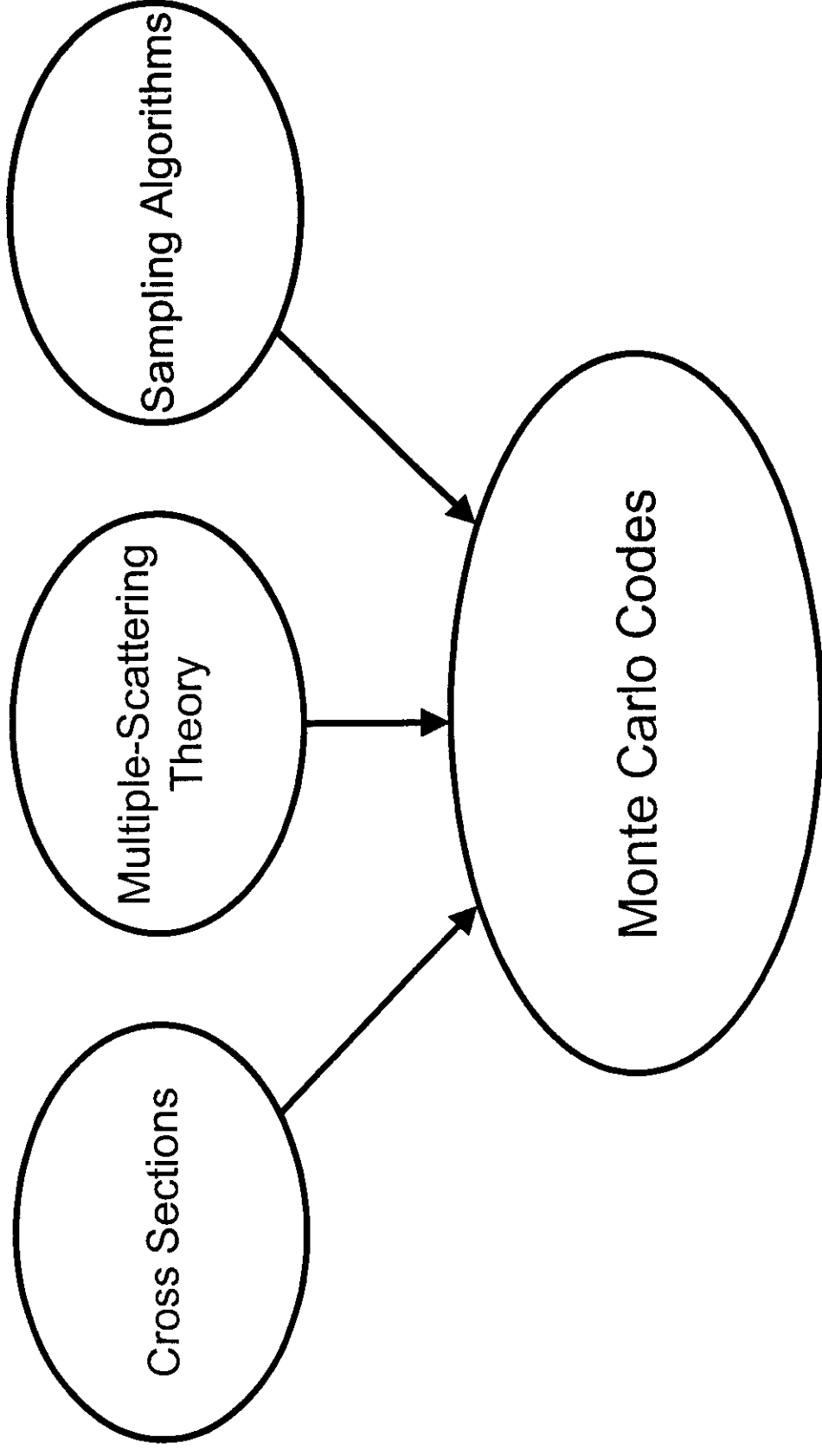
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Physics and Applications of ETRAN and ITS Electron-Photon Transport Monte Carlo Codes

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**Workshop on Nuclear Data for Science and Technology
Medical Applications
International Center for Theoretical Physics
Trieste, October 4-15, 1999**

Development of Radiation Interaction Data and Transport Calculations



Coupled Electron-Photon Monte Carlo Calculations

PHOTON TRANSPORT (Analog)

Photon Interactions

- Coherent (Rayleigh) scattering**
- Incoherent (Compton) scattering**
 - Binding corrections**
- Photoelectric absorption**
- Pair production (screened nuclear field)**
- Triplet production (field of atomic electrons)**
- Photonuclear absorption**

ELECTRON TRANSPORT (Condensed History)

Energy-loss

- Multiple-scattering treatment**
 - Class I (complete grouping)**
 - CSDA**
 - Energy-loss straggling**
 - Class II (catastrophic+CSDA)**
- Electronic stopping power (mean energy loss)**
 - Density effect**
- Radiative stopping power (mean energy loss)**
- Bremsstrahlung production (single scattering)**
 - Energy distribution**
 - Analytical theory**
 - “Exact”**

Elastic scattering

- Multiple-scattering treatment**
 - Goudsmit-Saunderson (exact)**
 - Molière (small-angle approximation)**
- Single-scattering cross sections**
 - Screened Rutherford**
 - “Exact”**
- Spatial displacements (“wiggleness” corrections)**
 - e.g.*, PRESTA (transverse)
 - e.g.*, TLC (transverse and longitudinal)
 - e.g.*, Random hinging

Coupled Electron-Photon Monte Carlo Calculations

SECONDARY RADIATIONS

Bremsstrahlung photons

Energy distribution

Analytical theory

"Exact"

Intrinsic angular distribution

Analytical theory

"Exact"

Knock-on electrons (Møller or Bhabha cross section)

Compton electrons

Binding corrections

Photoelectrons

Angular distribution

Pair "electrons"

Energy distribution

Angular distribution

Annihilation quanta

Positron at rest

Positron in flight

Atomic relaxation

Characteristic x rays

Auger electrons

ALGORITHM CONSIDERATIONS

Boundary crossing

Track-end disposition

Variance reduction

Geometry

Generality

Speed

Ease of use

Sampling

Numerical distributions

Analytical formulae

Ersatz functions

PHOTON TRANSPORT

PHOTON TRANSPORT

Single-scattering treatment

- pair (and triplet) production
- photo-electric absorption
- incoherent scattering
- coherent scattering

Production and transport of secondary radiation

- electron-positron pairs
- photoelectrons, characteristic x rays, Auger electrons
- Compton electrons

ELECTRON AND POSITRON TRANSPORT

ELECTRON AND POSITRON TRANSPORT

Basically "condensed"-history or path-segment model

Multiple-scattering treatment of all collision for:

- **ionization and excitation energy loss (including fluctuations)**
- **elastic scattering**

Sample individual interactions for production of secondary radiations:

- **bremsstrahlung photons (contributes also to energy-loss straggling)**
- **knock-on electrons (Møller [or Bhabha] cross section)**
- **characteristic x rays and Auger electrons from ionization events**
- **photons from annihilation of positrons at rest [and in flight]**

STEP SIZE

CONDENSED-HISTORY STEP SIZE

- Should be *short* enough: energy loss and angular deflection is small
- Should be *long* enough: enough collisions for multiple scattering treatment; not too many steps per history

CONDENSED-HISTORY STEP SIZE

- **Compromise:**

(a) Major step length for mean fractional energy loss of

$$2^{-1/k}, \quad \text{i.e.,} \quad T_{n+1} = 2^{-1/k} T_n$$

$$s_n = \int_{T_{n+1}}^{T_n} \left[\frac{1}{\rho} S(T') \right]^{-1} dT' = r_o(T_n) - r_o(T_{n+1})$$

Typically $k = 8$; mean fractional energy loss = 8.3%

(b) Divide major step into m sub-steps for small angular deflections. Typically $m = 2$ to 20

Mean deflection cosines for typical electron steps in the ETRAN random walk. Pathlengths for energy reduction by factor $2^{1/8}$ ($= 0.917$), further subdivided into m equal sub-steps.

	C	A	Cu	Ag	Pb
	$m = 2$	3	7	11	18
T(MeV)					
64	0.9989	0.9990	0.9994	0.9995	0.9996
32	0.9975	0.9973	0.9982	0.9985	0.9988
16	0.9948	0.9939	0.9952	0.9959	0.9963
8	0.9901	0.9876	0.9893	0.9901	0.9904
4	0.9825	0.9776	0.9793	0.9796	0.9788
2	0.9720	0.9634	0.9648	0.9640	0.9604
1	0.9597	0.9463	0.9477	0.9451	0.9375
0.5	0.9478	0.9297	0.9313	0.9271	0.9158
0.25	0.9382	0.9164	0.9185	0.9134	0.9014
0.125	0.9316	0.9072	0.9101	0.9056	0.8969
0.0625	0.9273	0.9012	0.9054	0.9033	0.9010
0.03125	0.9246	0.8975	0.9039	0.9057	0.9092
0.015625	0.9226	0.8951	0.9054	0.9117	0.9183
0.0078125	0.9209	0.8939	0.9096	0.9190	0.9280

ANGULAR DEFLECTIONS

GOUDSMIT-SAUNDERSON DISTRIBUTION

Disadvantages

- **Sampling more complicated than from Molière**
- **Cannot evaluate on the fly. Need to pre-store distributions for pre-selected set of pathlengths**

Advantages

- **Not in small-angle approximation, good for all scattering angles**
- **Formally exact for any cross section (e.g., e^- and e^+ differences)**
- **Can apply to smaller pathlengths**

MULTIPLE ELASTIC-SCATTERING ANGULAR DEFLECTIONS

Goudsmit-Saunderson Distribution

$$A_{GS}(\vartheta, s) = \sum_{\ell=0}^{\infty} \frac{2\ell+1}{2} e^{-sG_{\ell}} P_{\ell}(\cos \theta)$$

$$G_{\ell} = 2\pi N \int_{-1}^{+1} \frac{d\sigma(\theta, T)}{d\Omega} [1 - P_{\ell}(\cos \theta)] d(\cos \theta)$$

Energy loss included by $sG_{\ell} \rightarrow \int_0^s G_{\ell}(s') ds'$ evaluated in CSDA

ELASTIC-SCATTERING CROSS SECTION

Factored cross section

$$\frac{d\sigma(\theta)}{d\Omega} = \frac{d\sigma_{Ruth}(\theta)}{d\Omega} K_{scr}(\theta) K_{rel}(\theta)$$

Spin-relativistic correction

$$K_{rel}(\theta) = \frac{d\sigma_{Mott}}{d\Omega} / \frac{d\sigma_{Ruth}}{d\Omega}$$

Screening correction

$$K_{scr}(\theta) = \left(\frac{1 - \cos \theta}{1 - \cos \theta + 2\eta} \right)^2$$

Molière screening angle

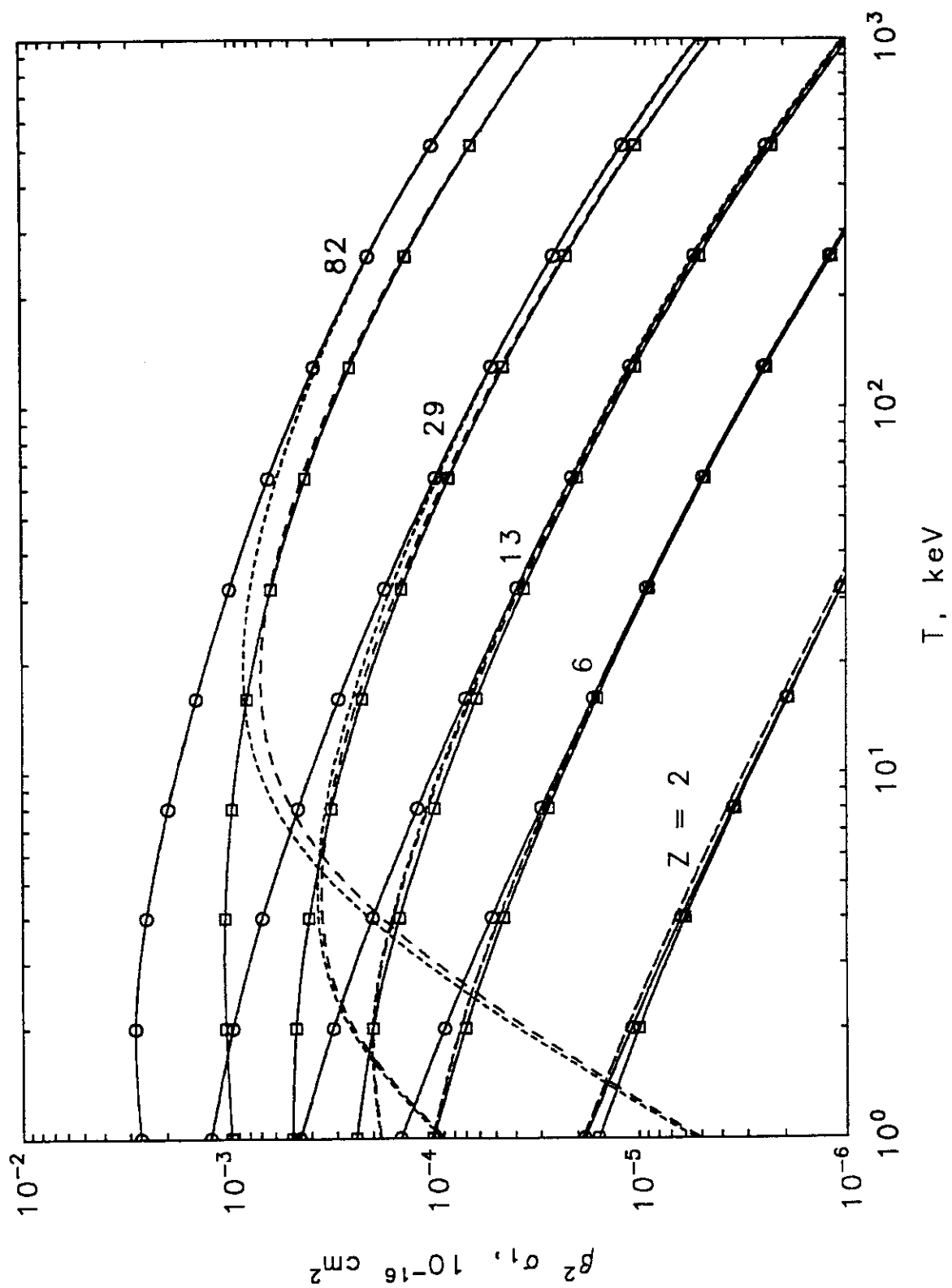
$$\eta = \frac{1}{4} \left(\frac{\alpha mc}{0.885 p} \right)^2 Z^{2/3} [1.13 + 3.76 (\alpha Z / \beta)^2]$$

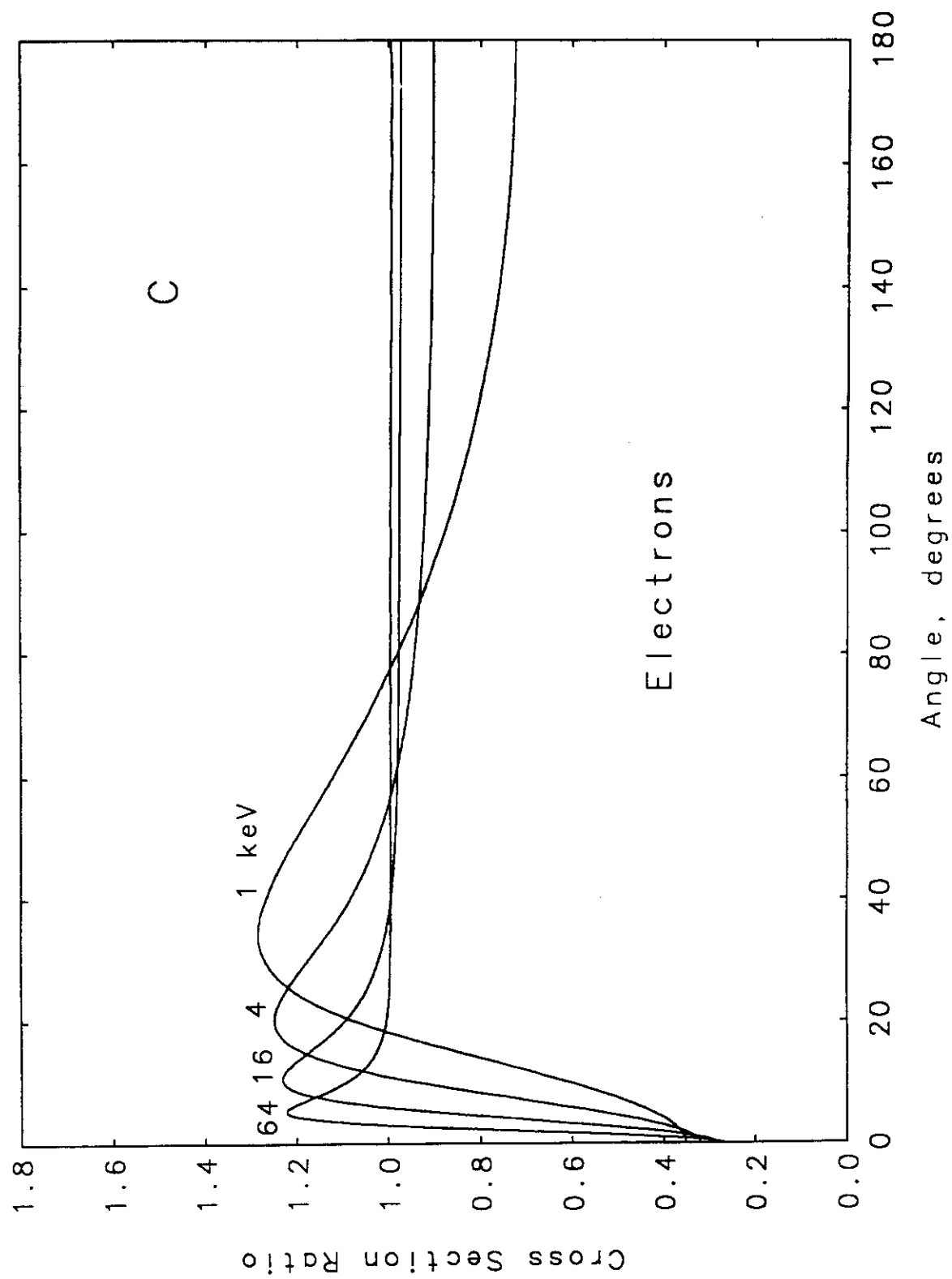
ELASTIC-SCATTERING CROSS SECTION

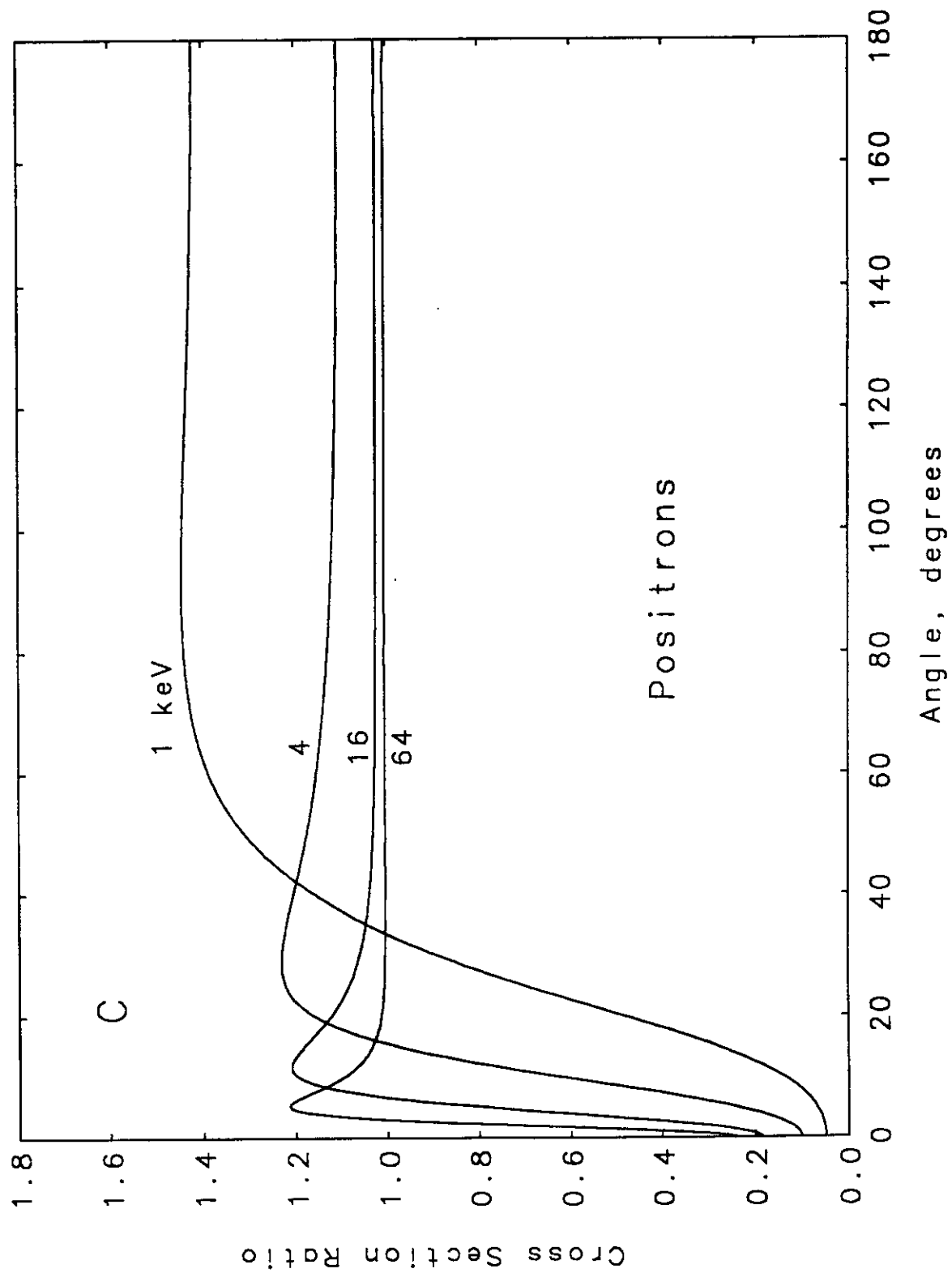
Transport (or momentum-transfer) cross section

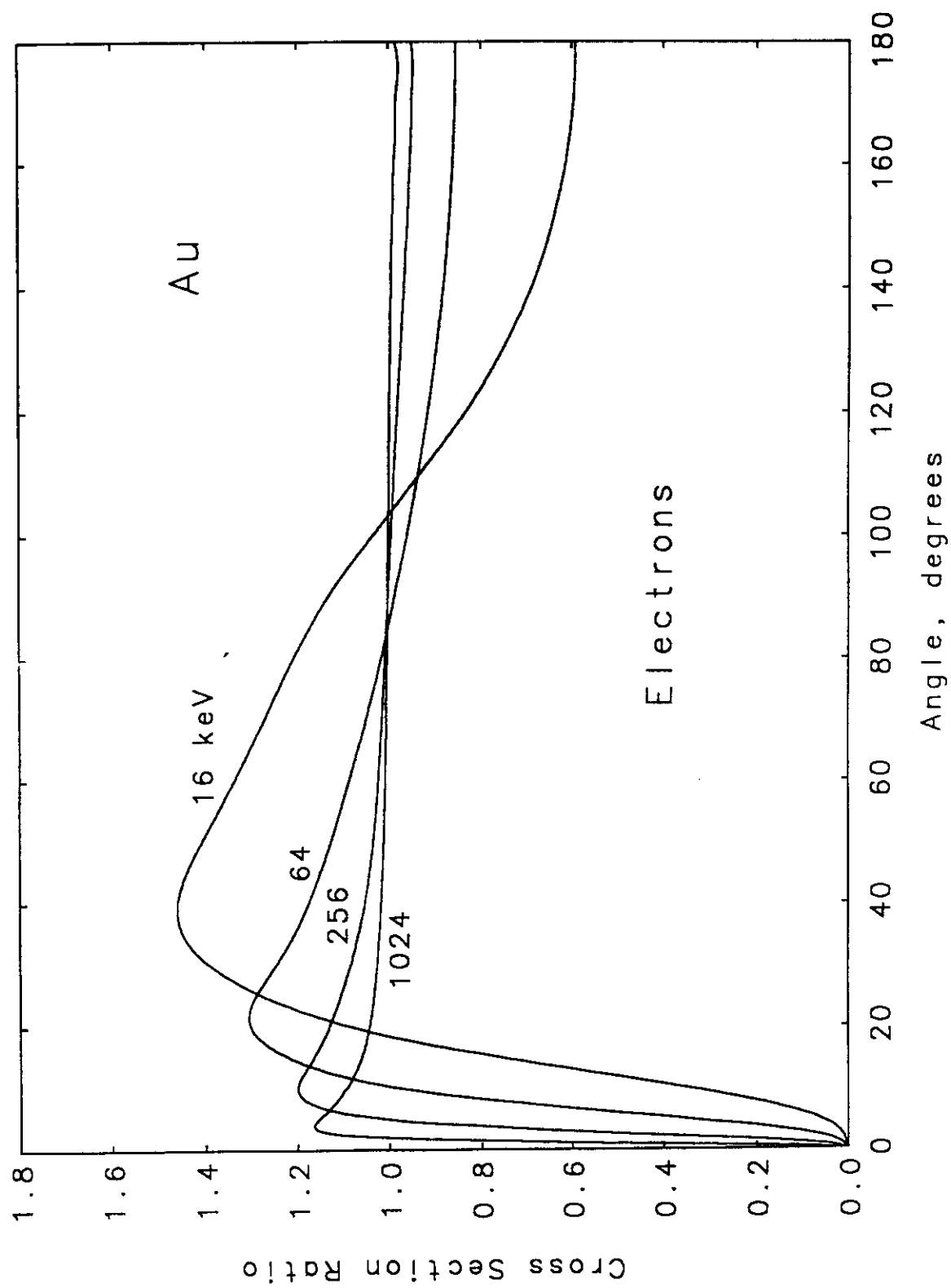
$$\sigma_t = 2\pi \int_{-1}^{+1} \frac{d\sigma(\theta)}{d\Omega} (1 - \cos\theta) d(\cos\theta)$$

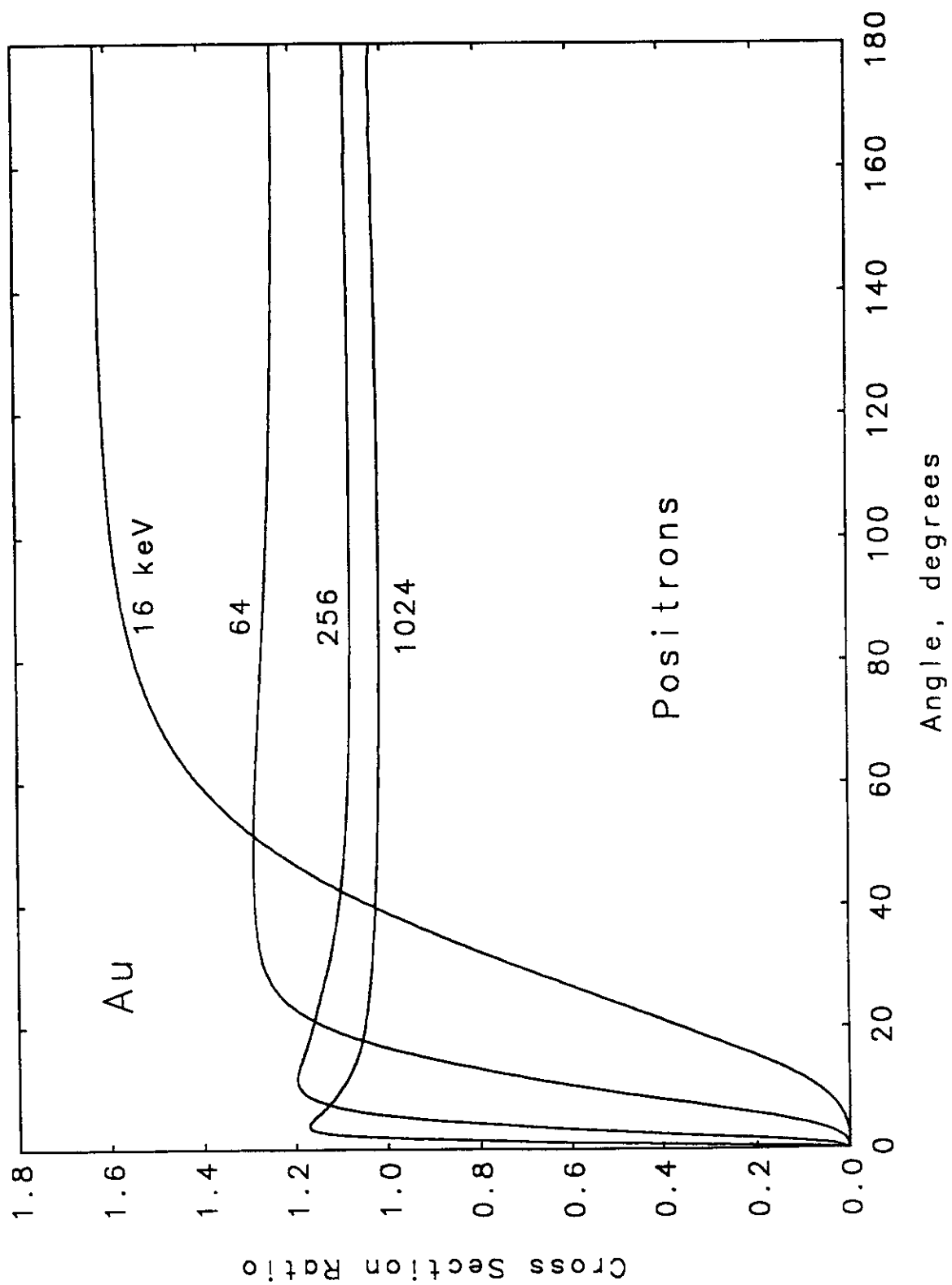
has dominant influence on multiple-scattering process

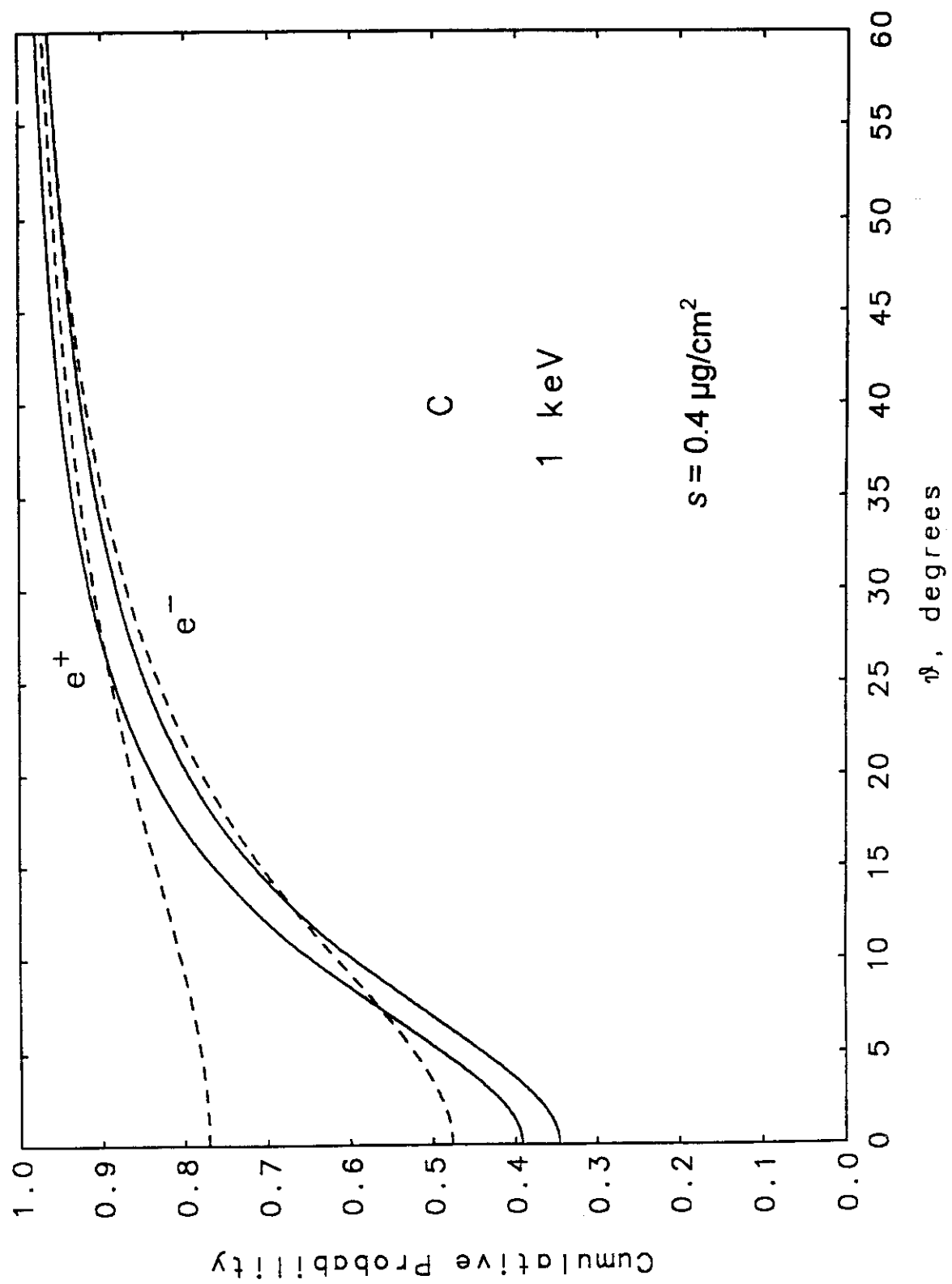


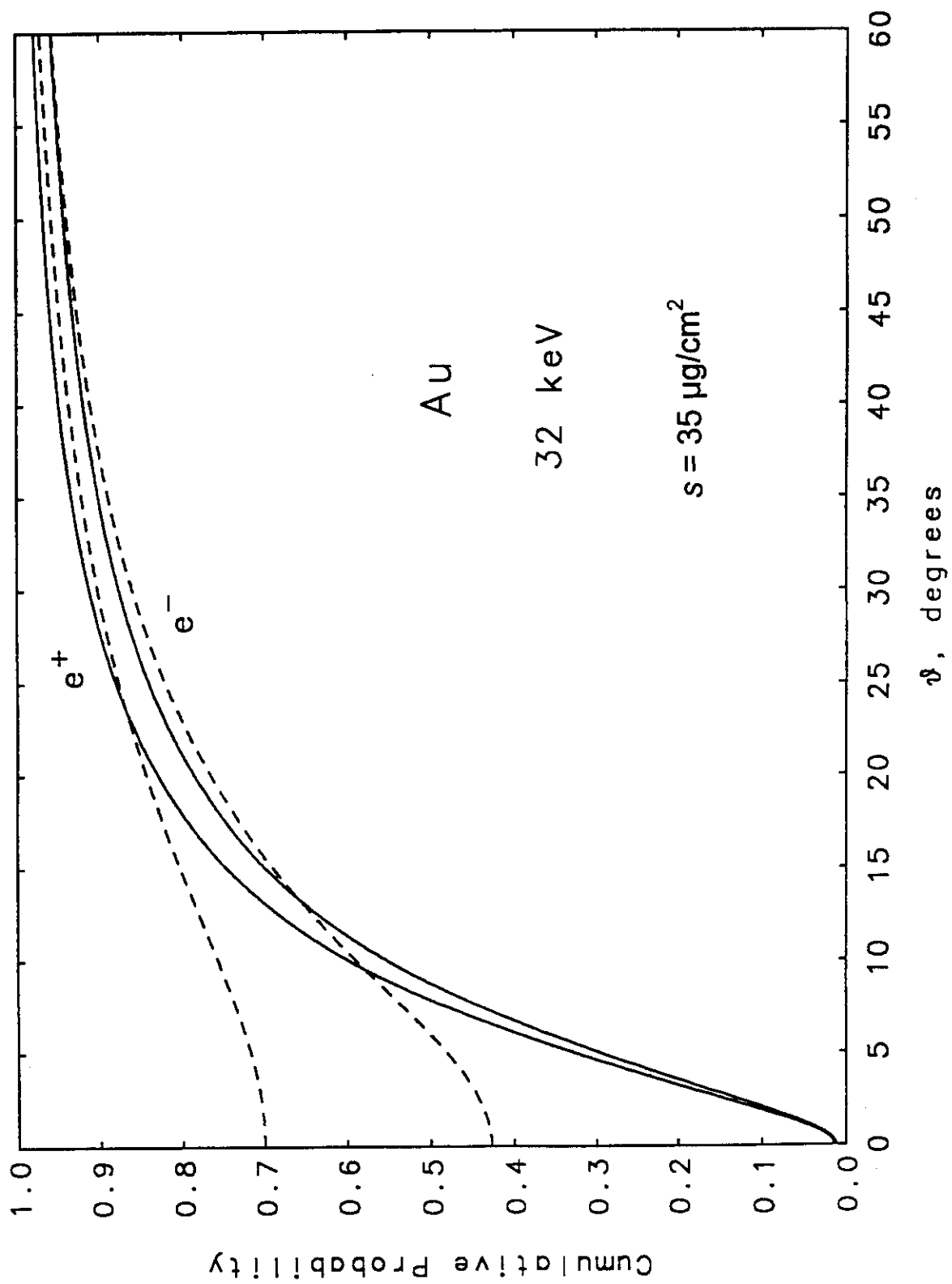












Dependence of foil transmission and reflection coefficients on the elastic-scattering cross sections used in ETRAN calculations. The results pertain to 32-keV electrons normally incident on tantalum foils of thickness z , expressed in units of the csda range $r_0 = 4.61 \text{ mg/cm}^2$.

Fraction Transmitted						
z/r_0	Number			Energy		
	Riley ^a	Factored ^b	b/a	Riley ^a	Factored ^b	b/a
0.05	0.823	0.826	1.004	0.776	0.778	1.003
0.10	0.647	0.649	1.003	0.560	0.562	1.004
0.15	0.492	0.483	1.002	0.385	0.385	1.000
0.20	0.346	0.350	1.012	0.246	0.249	1.012
0.25	0.226	0.227	1.004	0.147	0.148	1.007
0.30	0.135	0.135	1.000	0.0813	0.0815	1.002
0.35	0.0745	0.0744	0.999	0.0421	0.0418	0.993

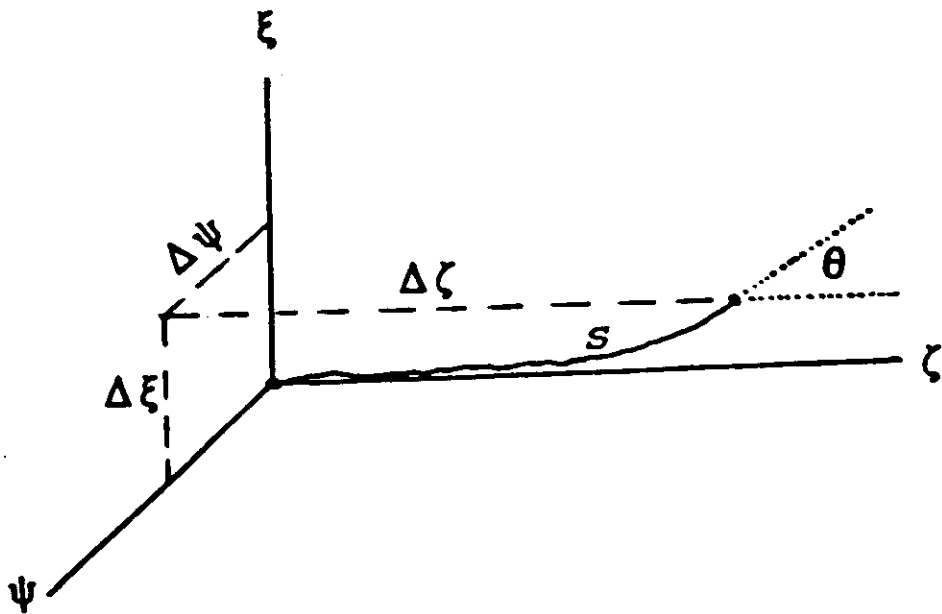
Fraction Reflected						
z/r_0	Number			Energy		
	Riley ^a	Factored ^b	b/a	Riley ^a	Factored ^b	b/a
0.35	0.490	0.490	1.000	0.378	0.378	1.000

^a Using elastic-scattering cross sections from partial-wave calculations with the Riley code.

^b Using elastic-scattering cross sections from the factored cross section, with the screening angle chosen to give the same transport cross section as from the partial-wave calculations.

SPATIAL DISPLACEMENT CORRECTIONS

Transverse and Longitudinal Corrections (TLC)



$$\Delta\zeta = s \frac{1 + \cos\theta}{2}$$

$$\Delta\xi = \frac{s}{2} \left[\sin\theta \cos\Delta\varphi + \nu_x \left(\frac{\langle \theta^2 \rangle}{6} \right)^{1/2} \right]$$

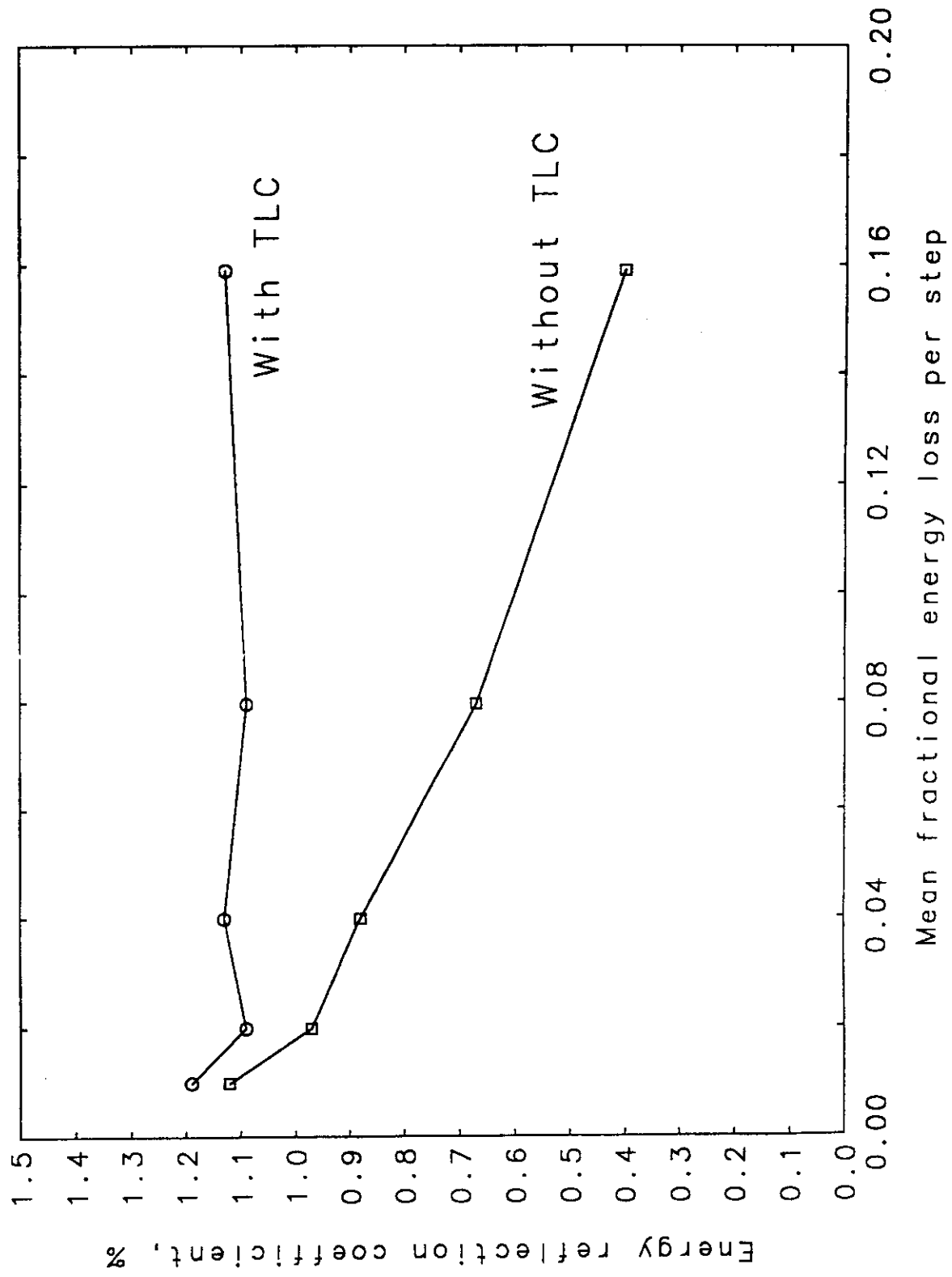
$$\Delta\psi = \frac{s}{2} \left[\sin\theta \sin\Delta\varphi + \nu_y \left(\frac{\langle \theta^2 \rangle}{6} \right)^{1/2} \right]$$

ν_x and ν_y are normally distributed random numbers

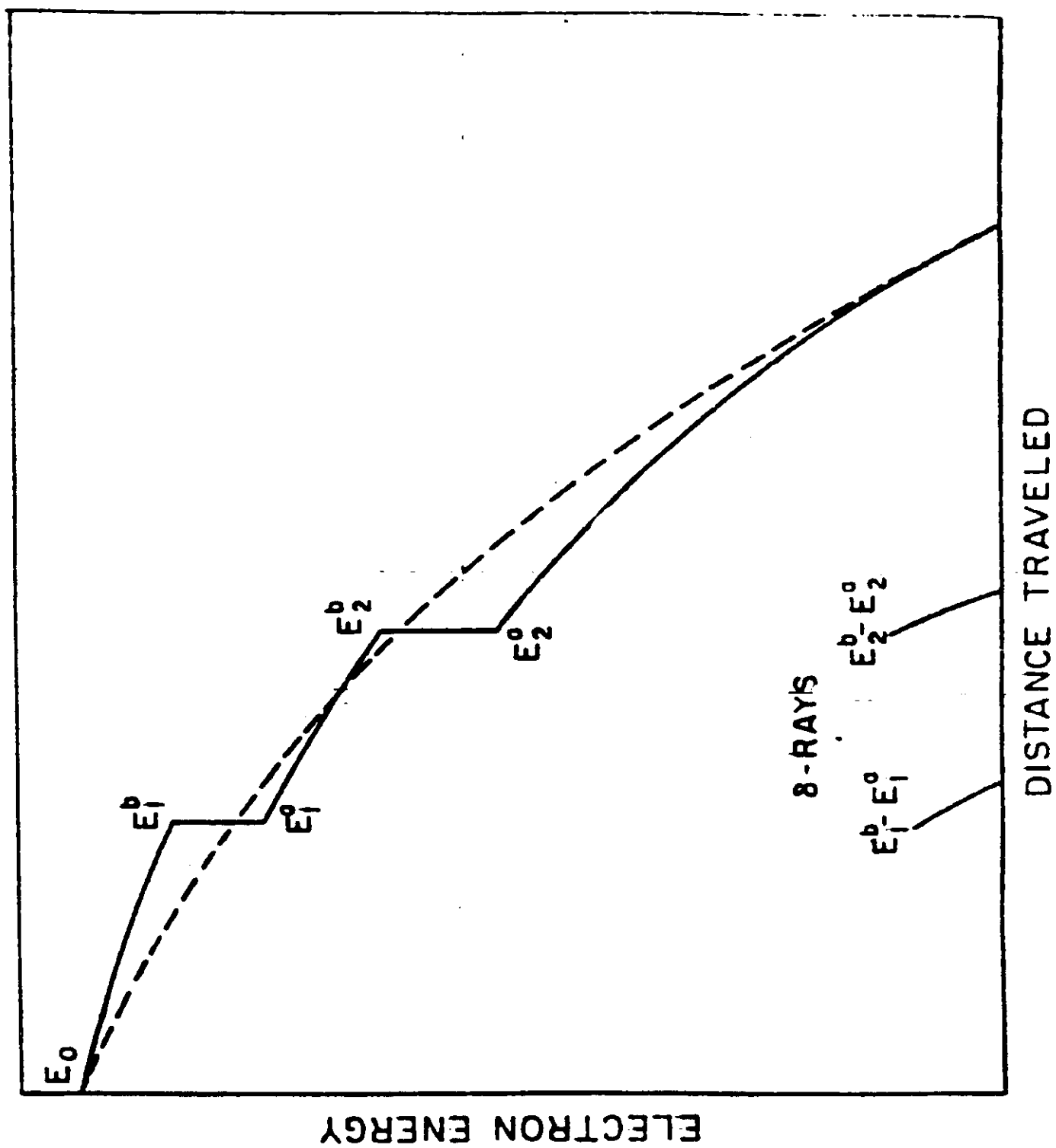
Transverse and Longitudinal Corrections (TLC)

$$\langle \theta^2 \rangle = 4\pi N_s \int_{\cos\theta_c}^1 \left[\frac{d\sigma_{Ruth}}{d\Omega} K_{scr}(\theta) \right] (1 - \cos\theta) d(\cos\theta)$$

$$2\pi N_s \int_{-1}^{\cos\theta_c} \left[\frac{d\sigma_{Ruth}}{d\Omega} K_{scr}(\theta) \right] d(\cos\theta) = n \sim 1$$



ENERGY LOSS



IONIZATION ENERGY-LOSS FLUCTUATIONS

Landau theory

- Many collisions
- One velocity (small mean energy loss)
- Rutherford energy-transfer cross section, $\sim \epsilon^{-2}$ (for $\epsilon \gg I$)
- First moment of cross section for small energy transfers from

Bethe theory

- Maximum energy transfer $\epsilon_{max} = \infty$

Distribution of energy loss Δ in terms of universal function

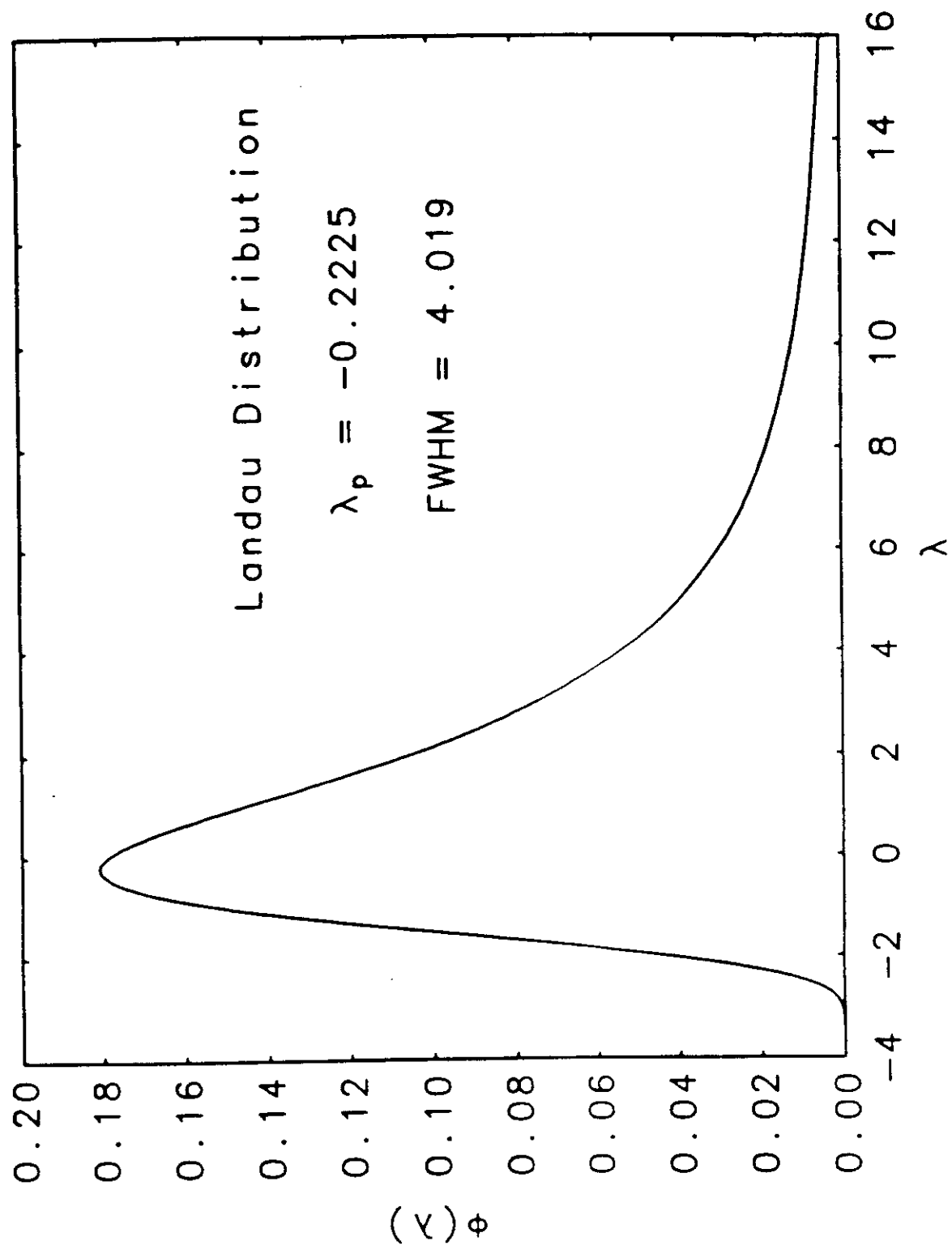
$$f(\Delta, s) d\Delta = \phi(\lambda) d\lambda$$

$$\lambda = \frac{\Delta}{\xi} - \ln \left[\frac{2\xi mc^2 \beta^2}{(1 - \beta^2)I^2} \right] + \beta^2 + \delta - 0.42278$$

$$\xi = \frac{2\pi r_e^2 mc^2 N Z}{\beta^2} s$$

$$\phi(\lambda) = \frac{1}{2\pi i} \int_{a-i\infty}^{a+i\infty} \exp(u \ln u + \lambda u) du$$

Valid for $T \gg \xi \gg I$



IONIZATION ENERGY-LOSS FLUCTUATIONS

Because $\epsilon_{max} = \infty$, $\langle \Delta \rangle = \int_0^\infty \Delta f(\Delta, s) d\Delta = \infty$

Need to impose correct mean energy loss.

Re-write Landau variable in terms of $\bar{\Delta}$ from stopping-power,

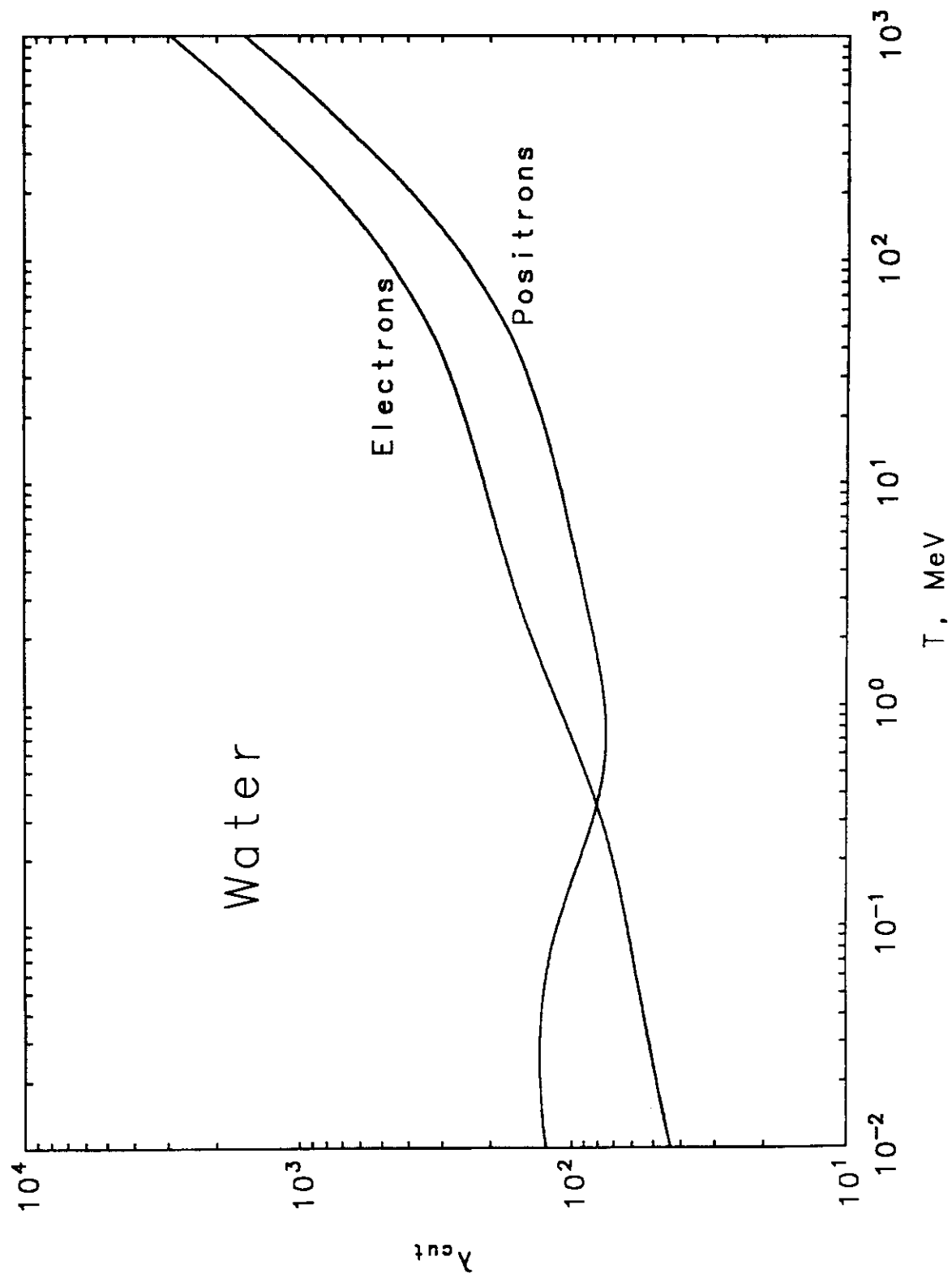
$$\lambda = \frac{\Delta - \bar{\Delta}}{\xi} + \nu^\pm$$

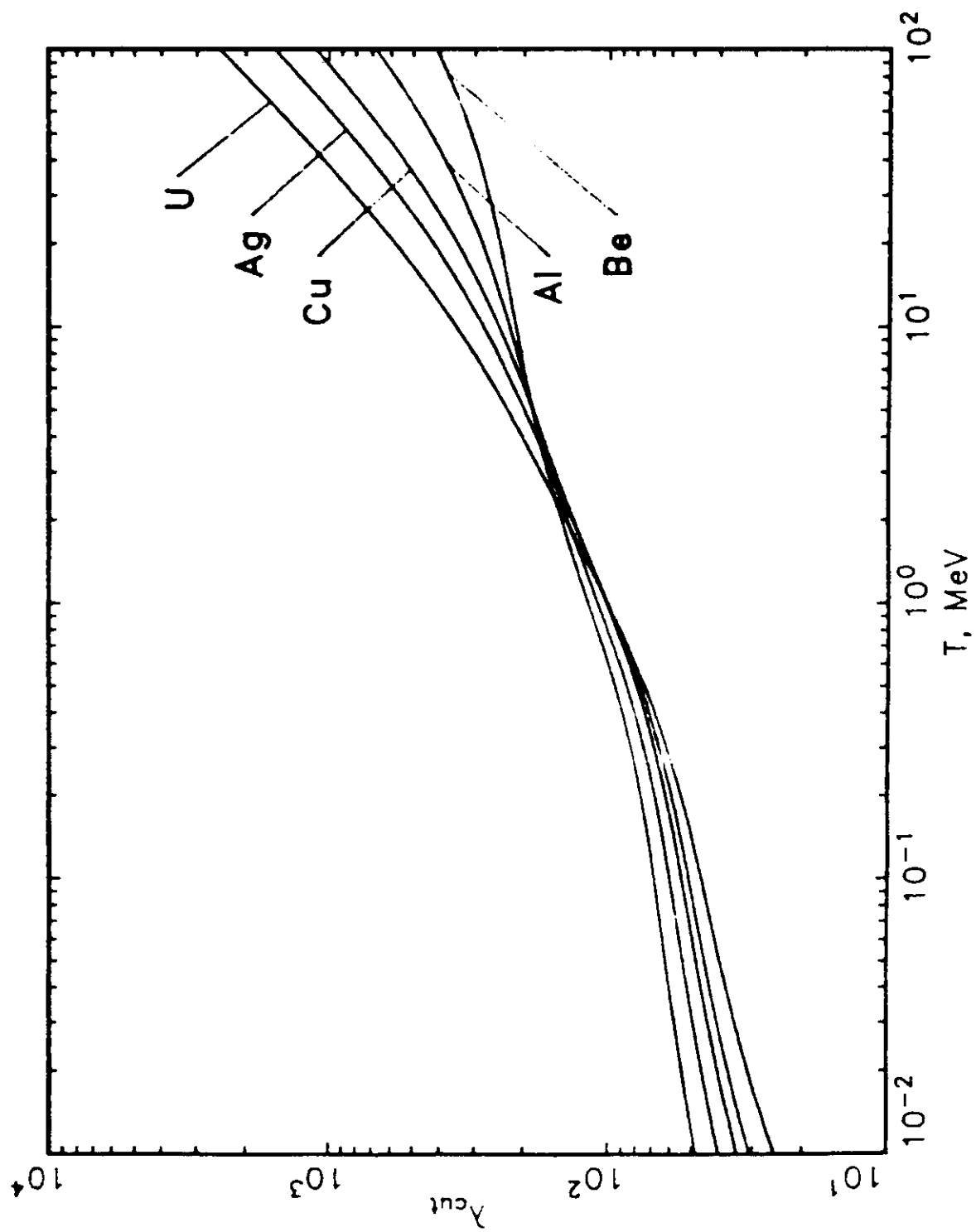
$$\nu^- = \ln\left(\frac{T}{\xi}\right) - 0.80907 + \left[\frac{\tau^2}{8} - (2\tau + 1)\ln 2\right]/(\tau + 1)^2$$

$$\nu^+ = \ln\left(\frac{T}{\xi}\right) - 0.422784 - \frac{\beta^2}{12} \left[11 + \frac{14}{\tau + 2} + \frac{10}{(\tau + 2)^2} + \frac{4}{(\tau + 2)^3}\right]$$

Define λ_{cut} such that

$$\bar{\lambda} = \nu^\pm$$





IONIZATION ENERGY-LOSS FLUCTUATIONS

Blunck & Leisegang included second moment of the cross section for small energy transfers

$$f_{BL}(\Delta, s) = \frac{1}{\sqrt{2\pi}\sigma} \int_{-\infty}^{+\infty} f(\Delta', s) \exp\left[-\frac{(\Delta - \Delta')^2}{2\sigma^2}\right] d\Delta'$$

Blunck & Westphal:

$$\sigma_{BW}^2 = 10eV \cdot Z^{4/3} \bar{\Delta}$$

Chechin & Ermilova estimate relative error of Landau-Blunck-Leisegang distribution (rather, its Laplace transform)

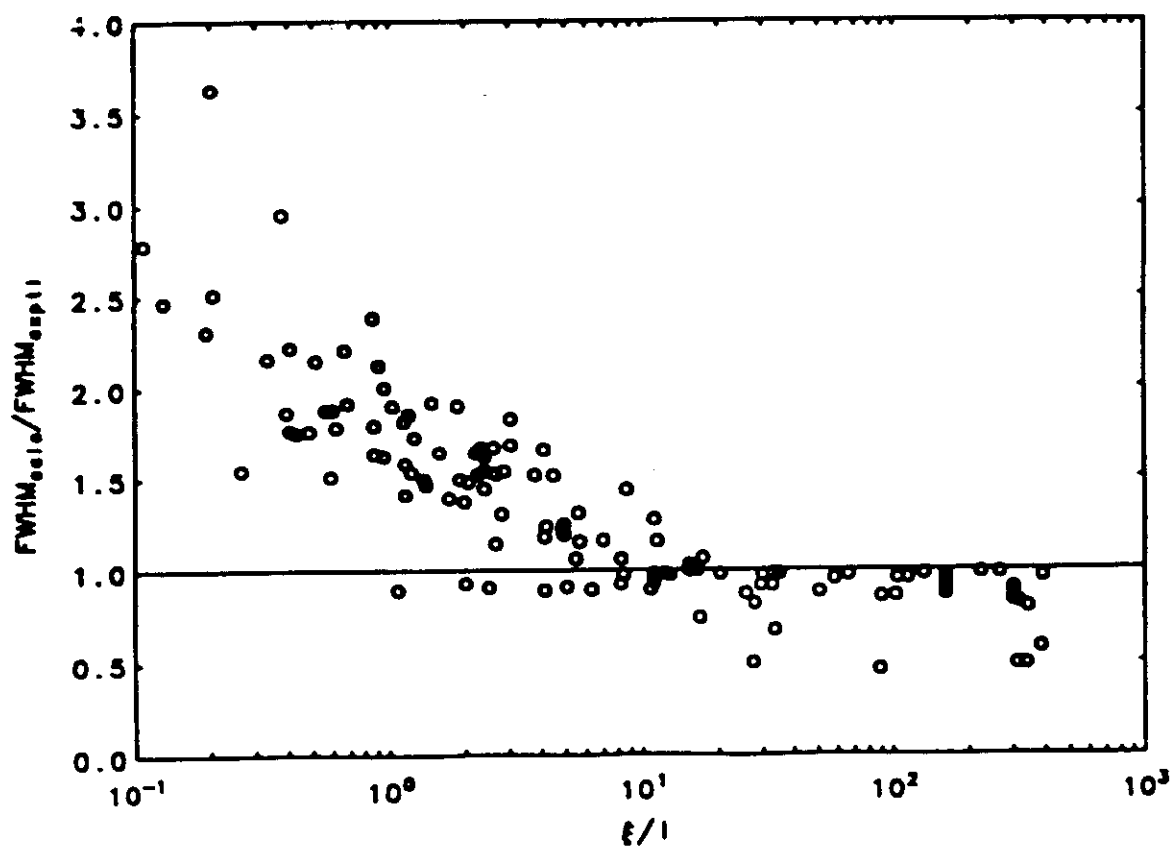
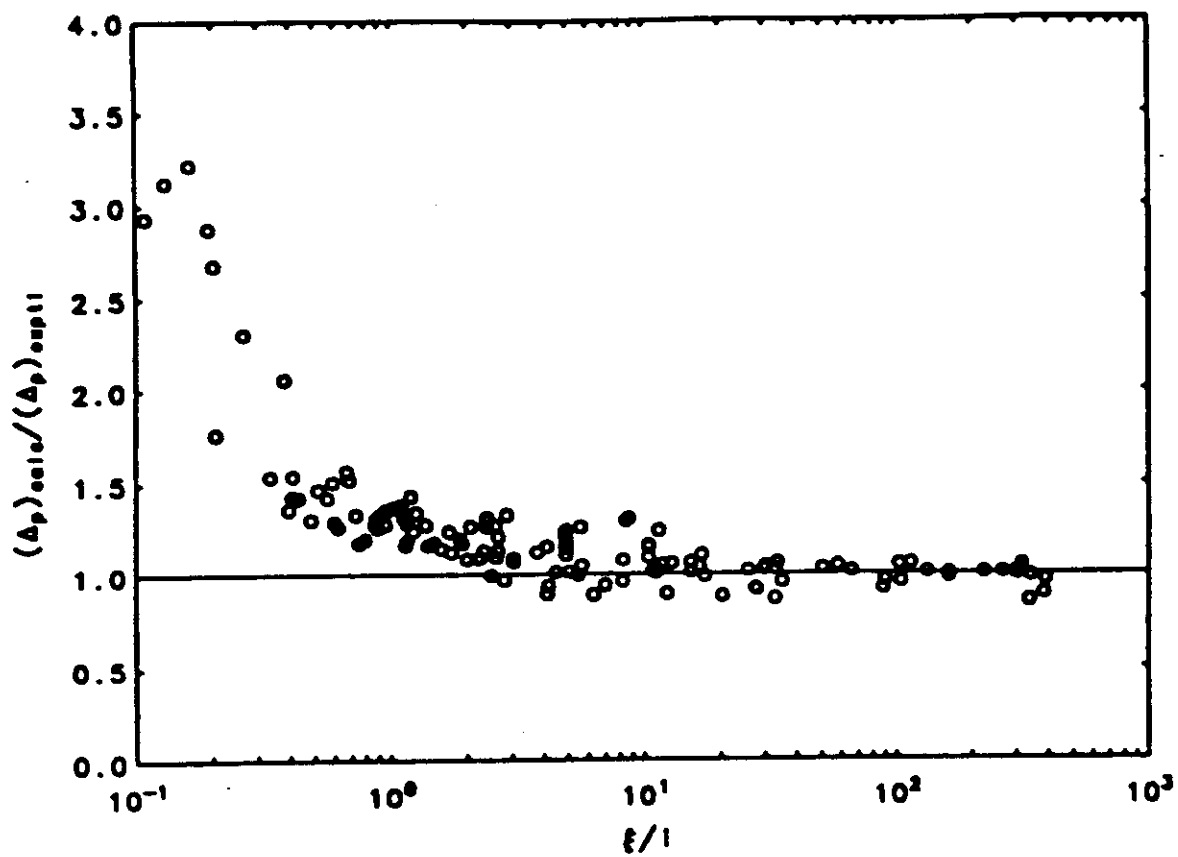
$$\epsilon \sim \left[10 \frac{\xi}{I} \left(1 + \frac{1}{10} \frac{\xi}{I} \right)^3 \right]^{-1/2}$$

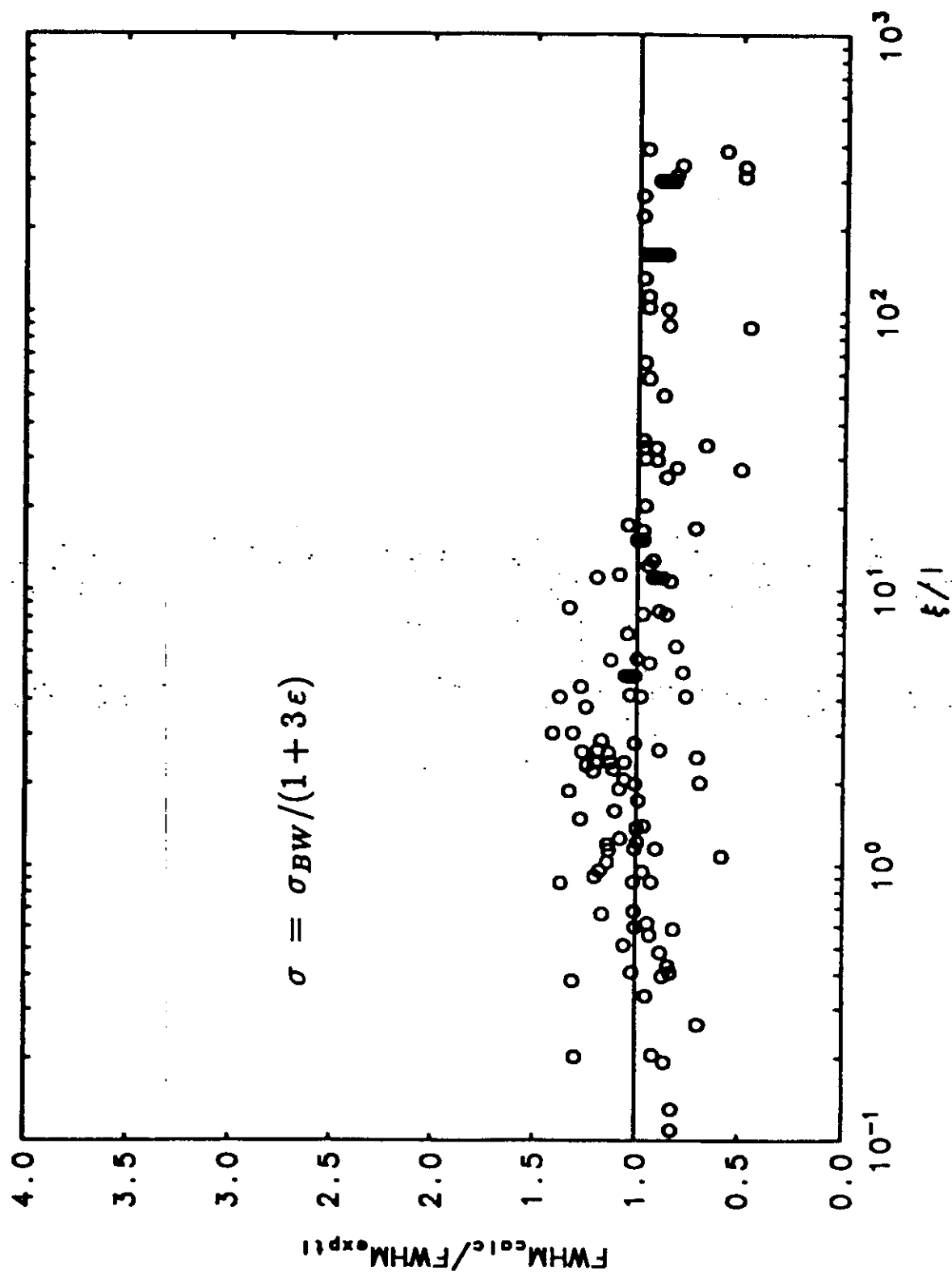
e.g.,

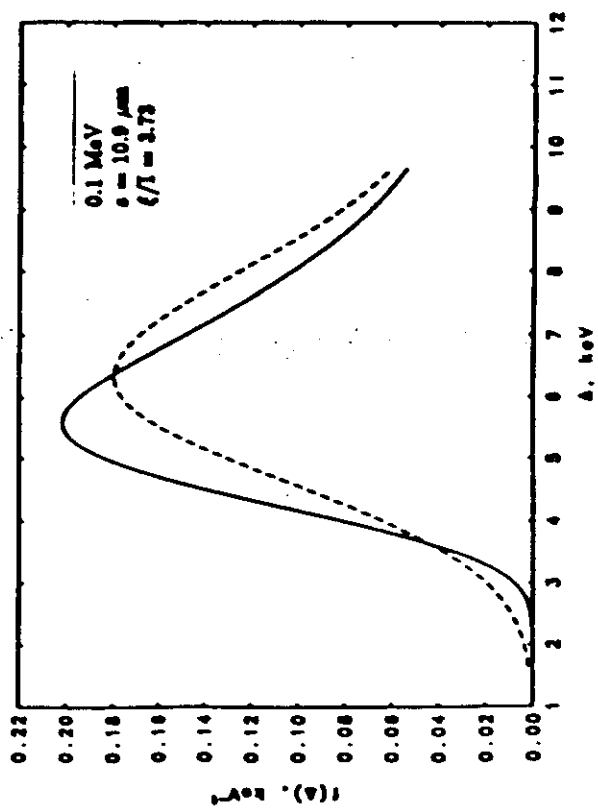
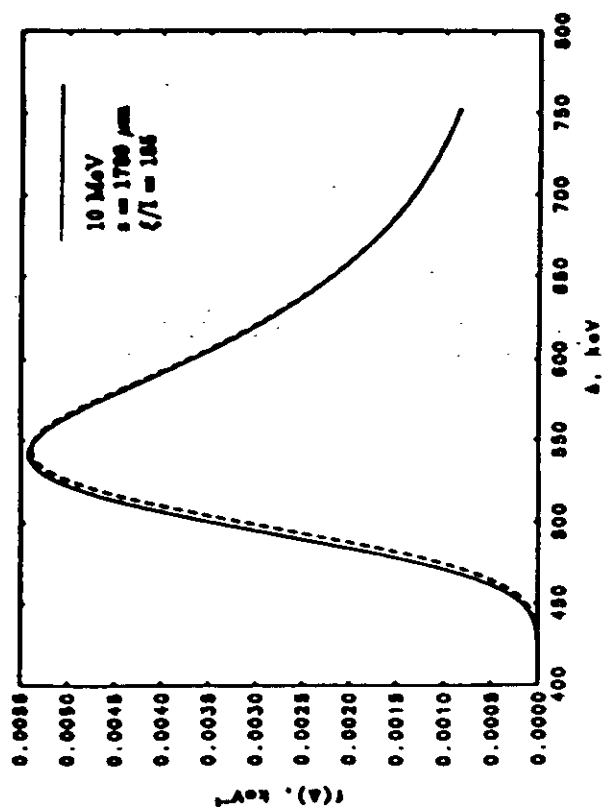
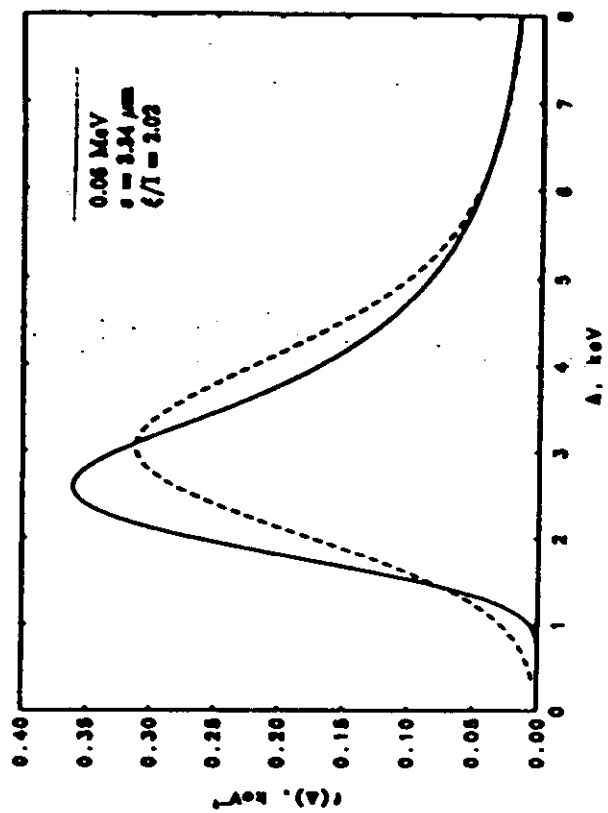
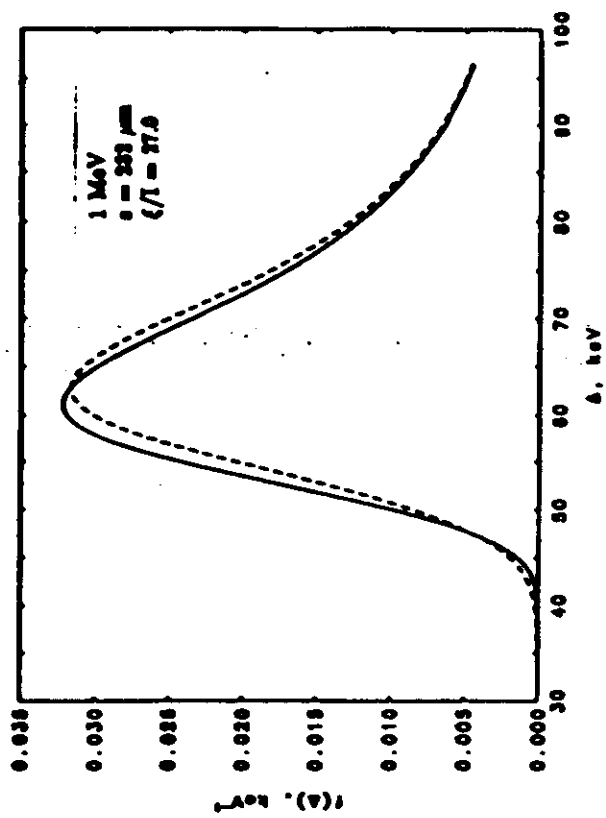
$$\epsilon \lesssim 10\% \quad \text{for} \quad \frac{\xi}{I} \gtrsim 4$$

Values of ξ/I for typical ETRAN pathlengths in which mean total energy loss is 8.3%

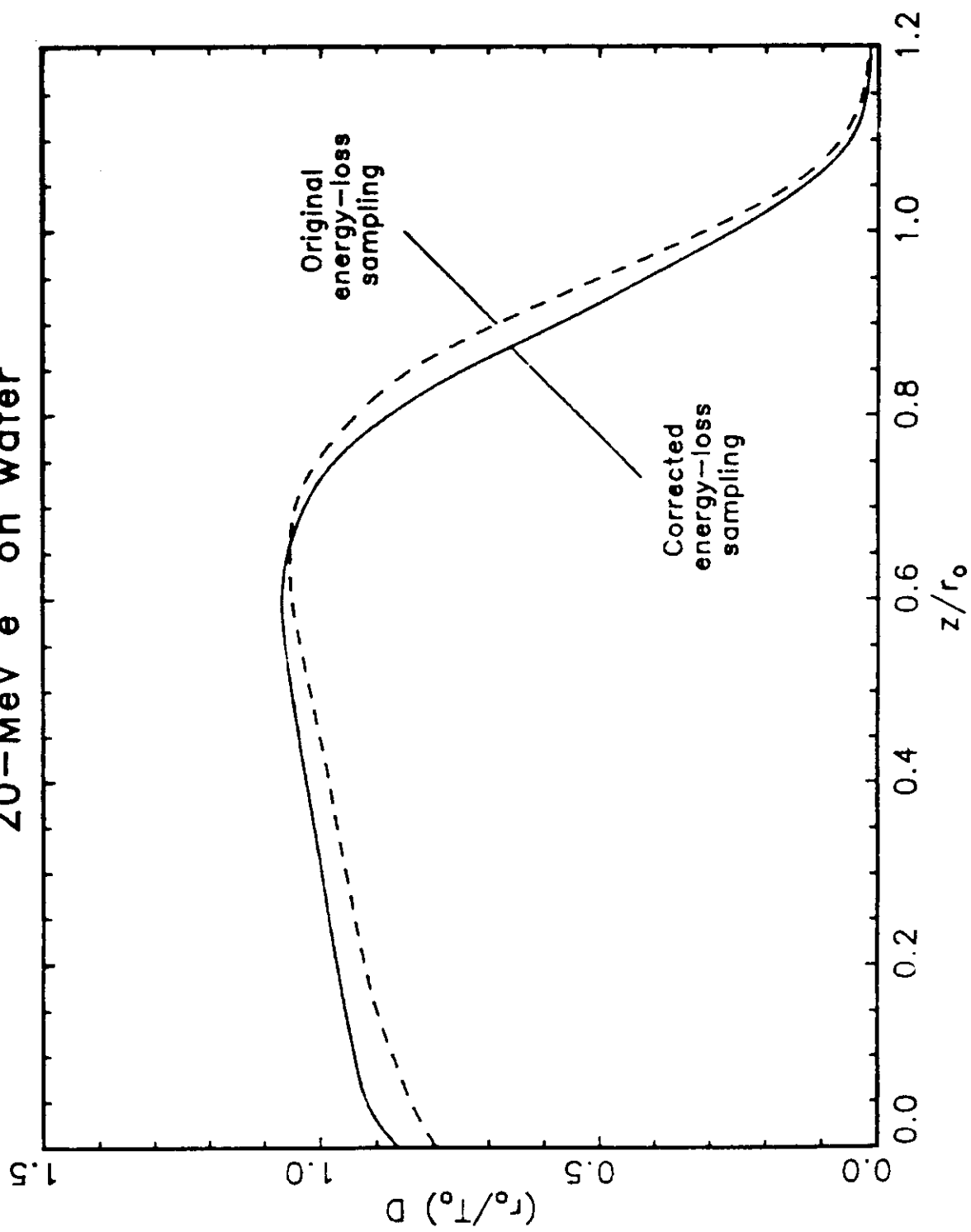
T(MeV)	H	C	Al	Cu	Ag	Gd	Au	U
1000		3421	907	241	111	68.6	44.0	35.0
100		2091	687	215	104	66.4	43.2	34.7
10	1481	434	195	92.1	55.4	39.7	27.9	23.4
5	806	238	111	57.2	36.5	27.3	19.8	16.8
2	360	106	51.2	27.9	18.9	14.7	10.9	9.44
1	196	56.9	28.3	15.7	10.9	8.68	6.54	5.75
0.5	106	30.7	15.6	8.73	6.20	5.00	3.82	3.40
0.2	46.8	13.6	7.05	4.00	2.89	2.37	1.84	1.66
0.1	25.2	7.42	3.89	2.23	1.63	1.35	1.06	0.96
0.05	13.6	4.08	2.17	1.26	0.93	0.78	0.62	0.56
0.02	6.15	1.89	1.02	0.61	0.46	0.39	0.32	0.29
0.01	3.40	1.07	0.60	0.36	0.28	0.24	0.20	0.18







20-MeV e^- on water



BREMSSTRAHLUNG

BREMSSTRAHLUNG ENERGY SPECTRA

$$\frac{d\sigma}{dk} = \frac{d\sigma_n}{dk} + \frac{Z d\sigma_e}{dk} = \left(1 + \frac{\eta}{Z}\right) \frac{d\sigma_n}{dk}$$

- Electron-nucleus cross section

$T > 50$ MeV, high-energy analytical theory with screening
and Coulomb corrections

$T \leq 2$ MeV, results of phase-shift calculations

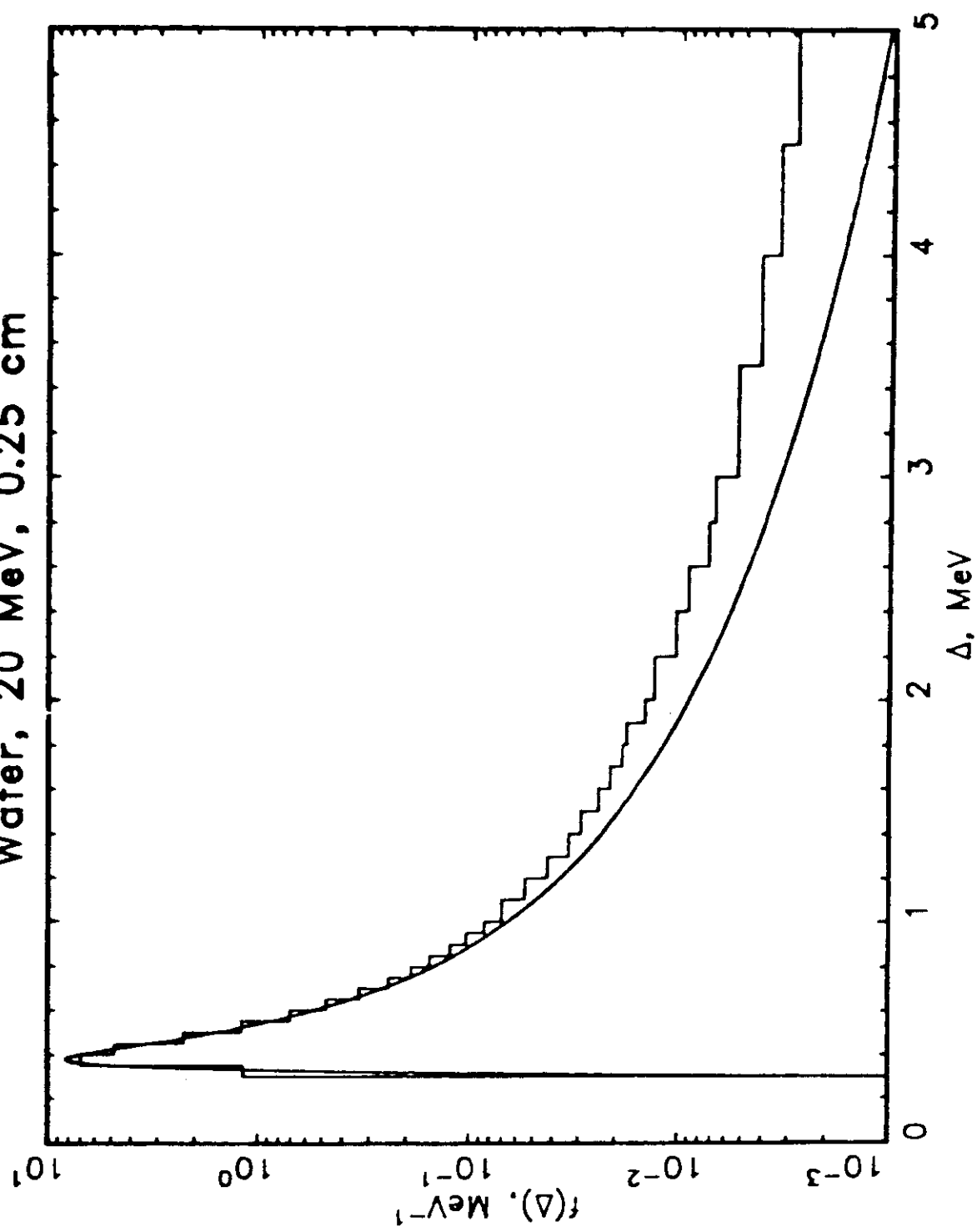
- Electron-electron cross section

Haug's theory with screening corrections

- Positron bremsstrahlung

Spectral shape of electron cross section, but normalized to
positron radiative stopping power

Water, 20 MeV, 0.25 cm



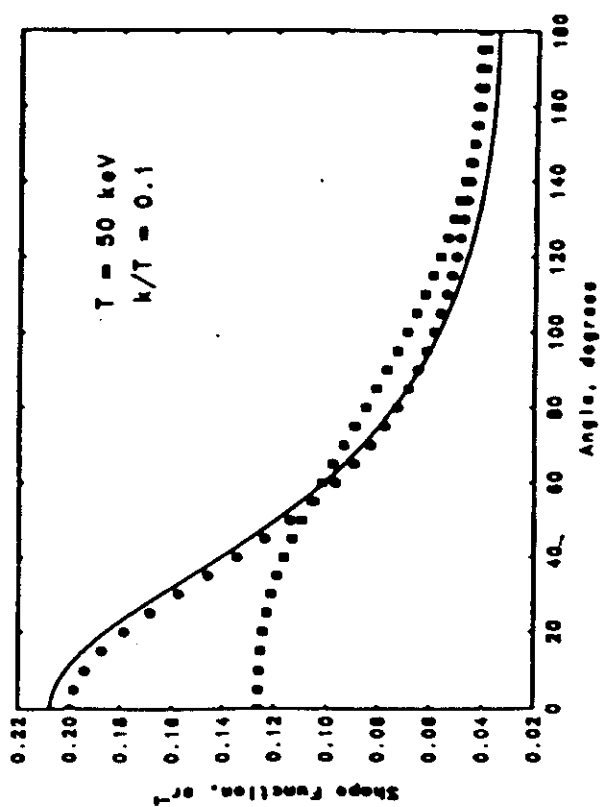
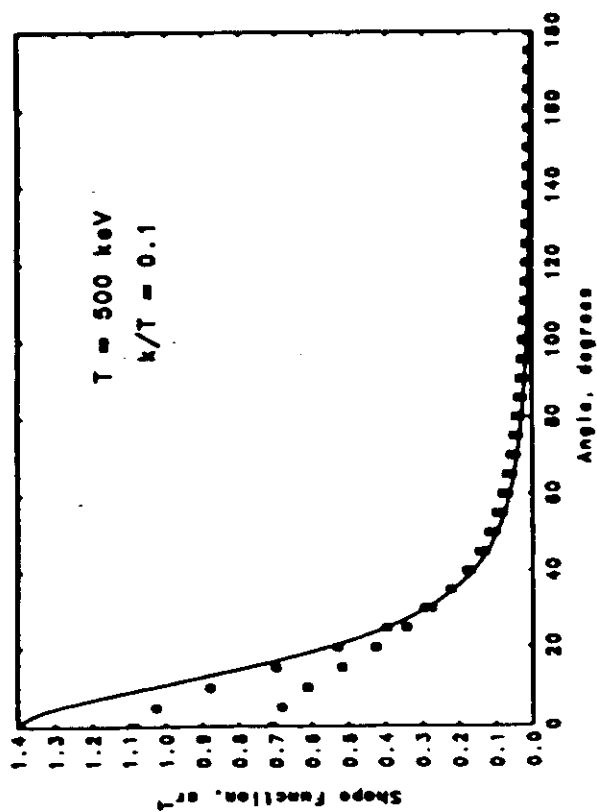
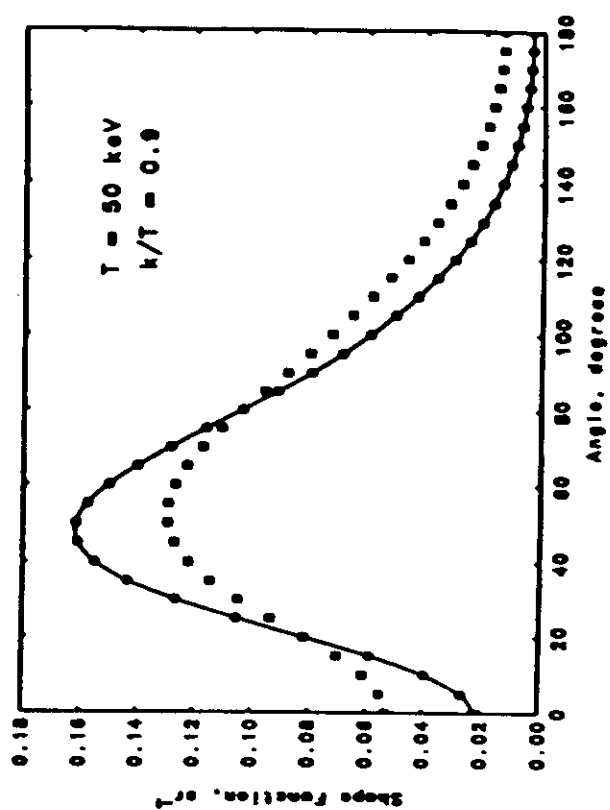
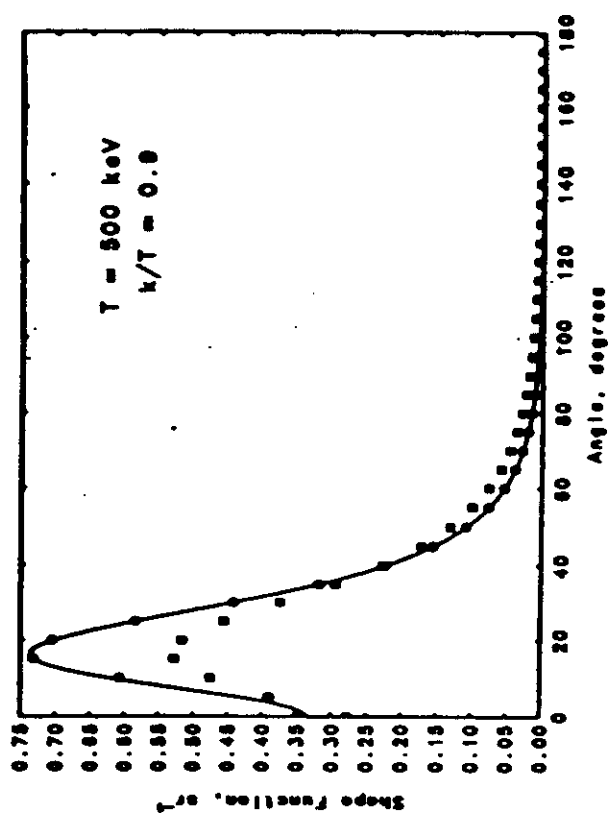
BREMSSTRAHLUNG ANGULAR DISTRIBUTIONS

$$f(\theta) = \frac{d^2\sigma}{d\Omega dk} \bigg/ \frac{d\sigma}{dk}$$

High energies: Olsen-Maximon formula with screening and
Coulumb corrections

Low energies: Sauter formula (no screening)

Should utilize information from Kissel-Pratt phase-shift
calculations, 1 - 500 keV.



OTHER FEATURES

HISTORY TERMINATION

Electrons

$T > T_{\text{cut}}$ follow

$T_{\text{save}} > T > T_{\text{cut}}$ follow if scoring boundary can
be reached (if electron fluence
is scored, $T_{\text{save}} = T_{\text{cut}}$)

$T \leq T_{\text{cut}}$ terminate; residual energy
deposited along straight path of
length = residual range
× detour factor

Photons

$T \leq T_{\text{cut}}$ terminate

VARIANCE REDUCTION OPTIONS

- **To improve statistics for secondary photons, can artificially scale up:**

bremsstrahlung production cross sections

electron-impact ionization cross sections

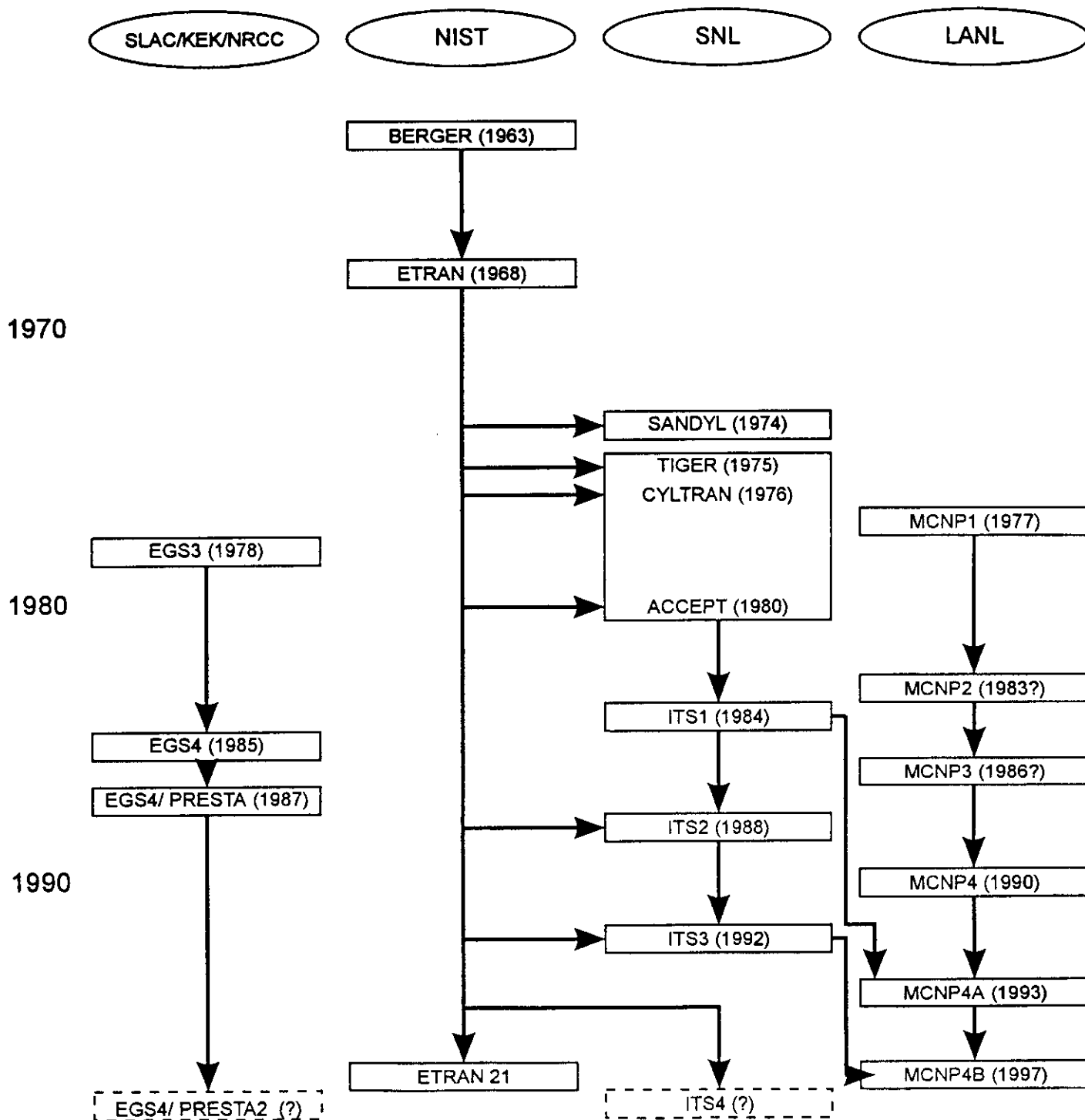
Compensating weights are assigned to the histories

- **Numbers of histories of secondary electrons from the large number of scaled-up photons can be controlled**
- **Emergent photon scoring by collision-density method (i.e., score uncollided escape probability after each scattering)**

LIMITATIONS

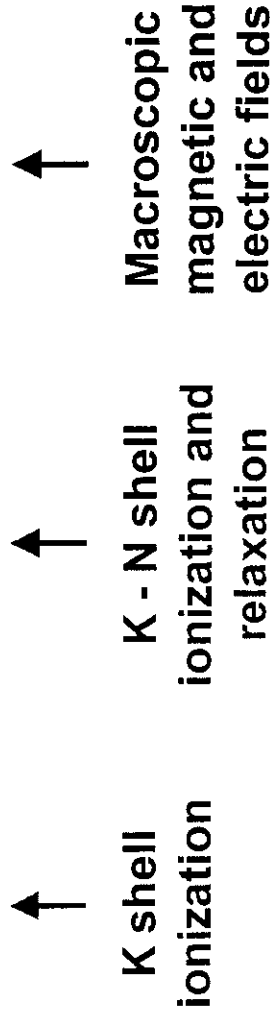
- **Covers energies only down to 1 keV**
- **No electron or photon diffraction, no electron channeling, no intense-beam collective effects**
- **No interactions with fields of nucleons (no electron elastic scattering from nucleons, no electro- and photo-disintegration of nucleus)**
- **Lack of correlation between sampled primary electron energy loss and sampled knock-on electron energy**
- **Short pathlength and low-energy electron worries**
 - **general inapplicability of multiple-scattering treatment**
 - **some approximate boundary-crossing algorithms**
 - **failure of Bethe-theory stopping power**

GENERAL PURPOSE ELECTRON-PHOTON MONTE CARLO CODES



Integrated TIGER Series (ITS): Eight Codes, One Library

	STANDARD	PCODES	MCODES
1D	TIGER	TIGERP	
2.5D	CYLTRAN	CYLTRANP	CYLTRANM
3D	ACCEPT	ACCEPTP	ACCEPTM



PHOTON TRANSPORT

PROCESS	ETRAN	ITS v. 3	MCNP4B	EGS4
Coherent scattering	yes	yes	yes	yes
Incoherent scattering	with binding	with binding	approx. binding	no binding*
Photoabsorption	yes (K relaxation)	yes (complete relax)	yes (K, L relaxation)	yes (K x rays)*
Pair production	yes	yes	yes	yes
Triplet production	included	included	no	?
Photonuclear abs.	no	no	no	no
Cross Sections	current NIST	current NIST	ENDF/B-IV (1975)	Storm&Israel (1975)

*/ Special versions developed at KEK include incoherent binding and atomic relaxation.

ELECTRON TRANSPORT

PROCESS	ETRAN	ITS v. 3	MCNP4B	EGS4
Energy-loss multiple scattering	Class I	Class I	Class I	Class II
Stopping powers	ICRU 37	ICRU 37	old B&S	ICRU 37
Collision straggling	adjusted Landau/BL	adjusted Landau/BL	adjusted Landau/BL	only above Ecut
Brems production spectrum	“exact”	“exact”	old B&S/K-M analytical package	renorm HE limit theory
Brems angular distribution	analytical theory package	analytical theory package	analytical theory package	Schiff theory
Elastic multiple scattering	Goudsmit-Saunderson	Goudsmit-Saunderson	Goudsmit-Saunderson	Moliere
Elastic single scattering	factored (correct σ_i)	“exact”	“exact”	screened Rutherford
Spatial displacement correction	TLC (1 st order)	none	none	PRESTA (0 th order)

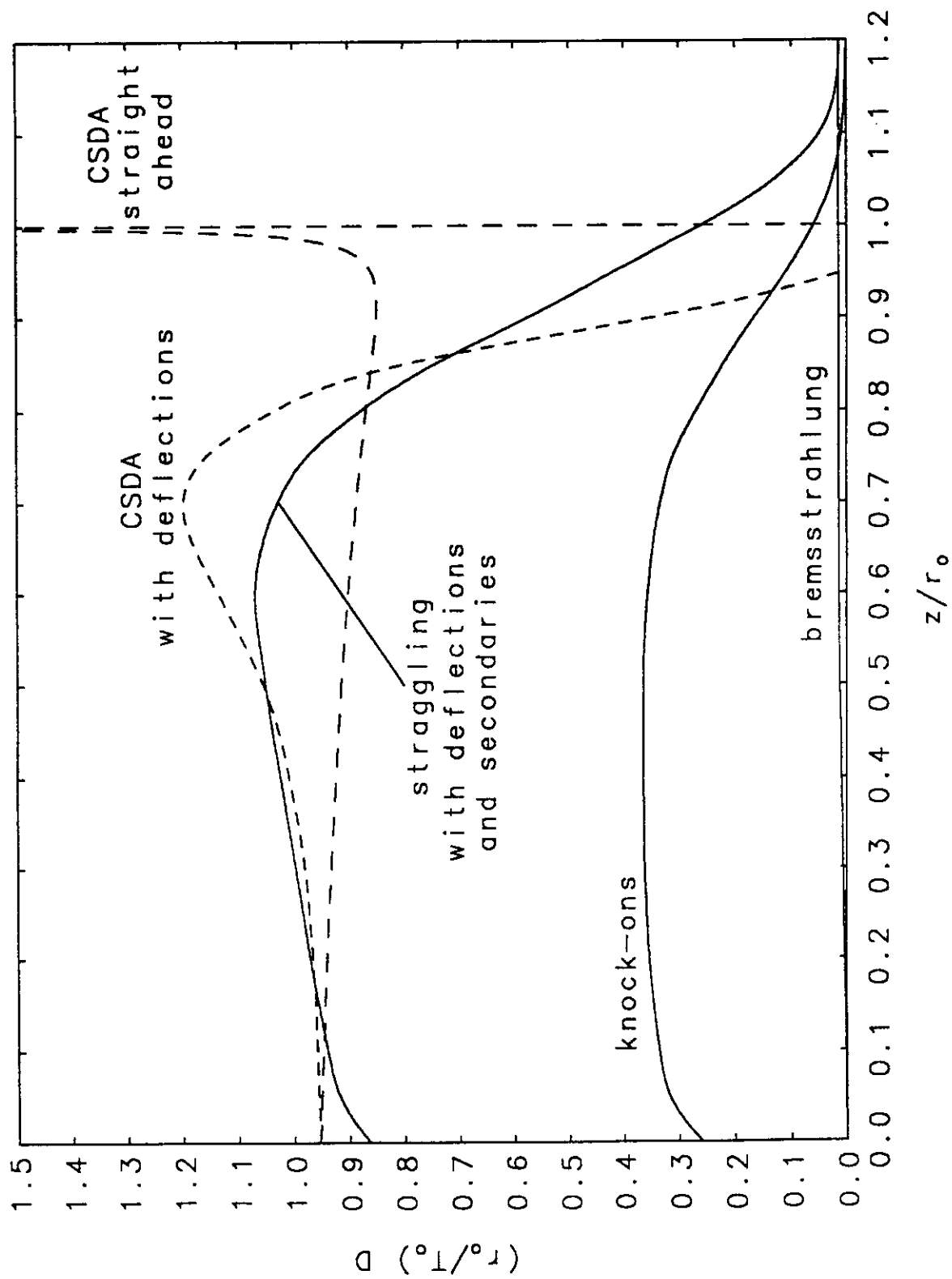
SECONDARY RADIATIONS

Process	ETRAN	ITS v. 3	MCNP4B	EGS4
Brems production spectrum	"exact"	"exact"	old B&S/K-M analytical package	renorm HE limit theory
Brems angular distribution	analytical theory package	analytical theory package	analytical theory package	Schiff theory
Knock-ons	Moller/free	Moller/free	Moller/free	Moller/free
Compton electrons	with binding	with binding	approx. binding	no binding*
Photoelectrons	yes	yes	yes	yes
Electron-positron pairs	both treated as electrons	both treated as electrons	both treated as electrons	both treated as electrons
Annihilation quanta	at rest	at rest	at rest	in flight, at rest
X rays (photoabsorption)	K-shell	all shells	all shells	K-shell
Auger electrons (photo absorption)	K-shell	all shells	all shells (?)	none
X rays (electron impact)	K-shell	all shells	all shells (?)	none
Auger electrons (electron impact)	K-shell	all shells	all shells (?)	none

GEOMETRIES TREATED

	ETRAN	ITS v. 3	MCNP4B	EGS4
Materials	single	multiple	multiple	multiple
Slab geometry	(included)	TIGER	-	available
Cylindrical geometry	simple, but many simultaneously	CYLTRAN	-	available
General geometry	no	ACCEPT	only	*

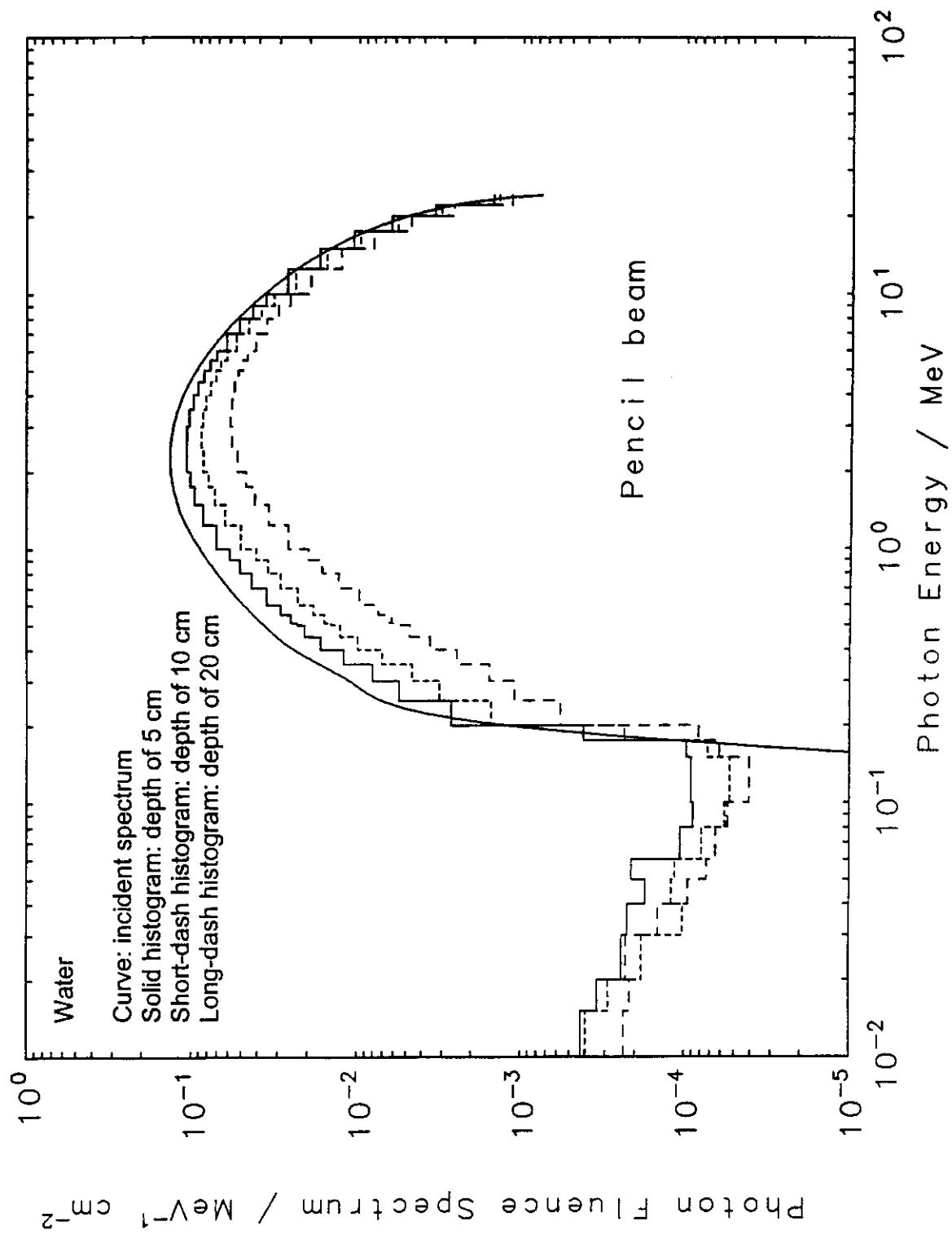
*/ User code available for a number of geometries: rectangular, conical, cylindrical, spherical. In general, geometry is written by user.

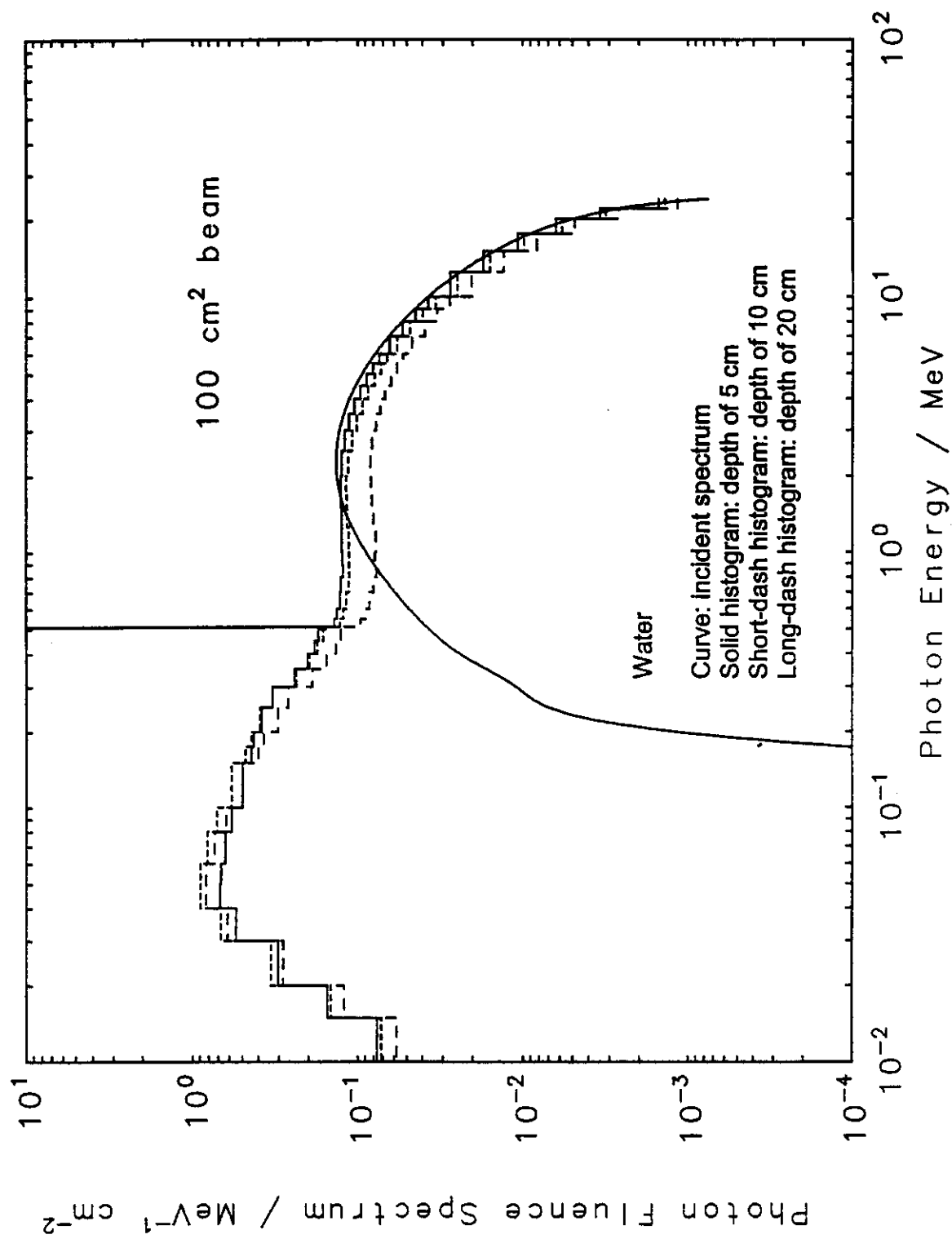


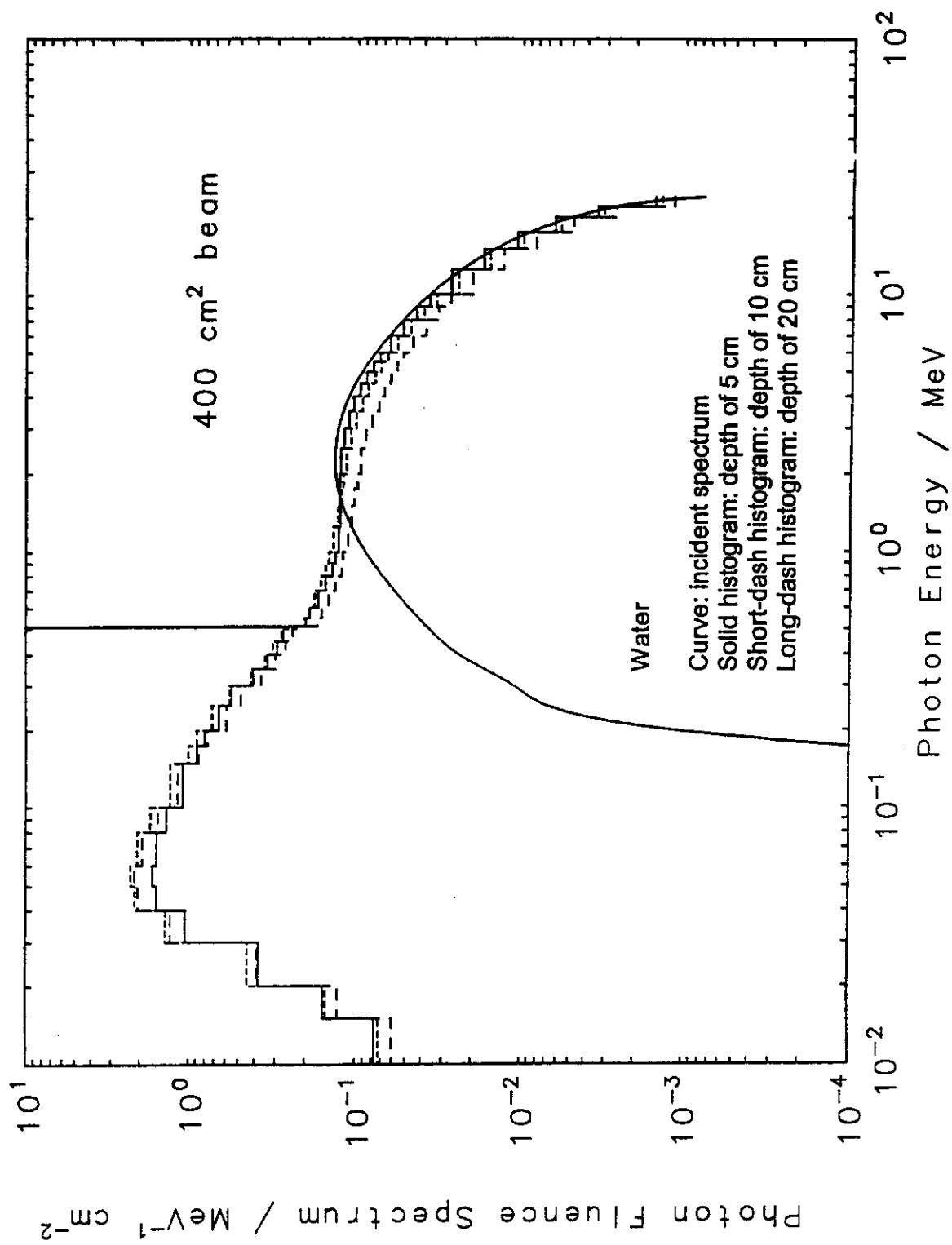
Some Examples

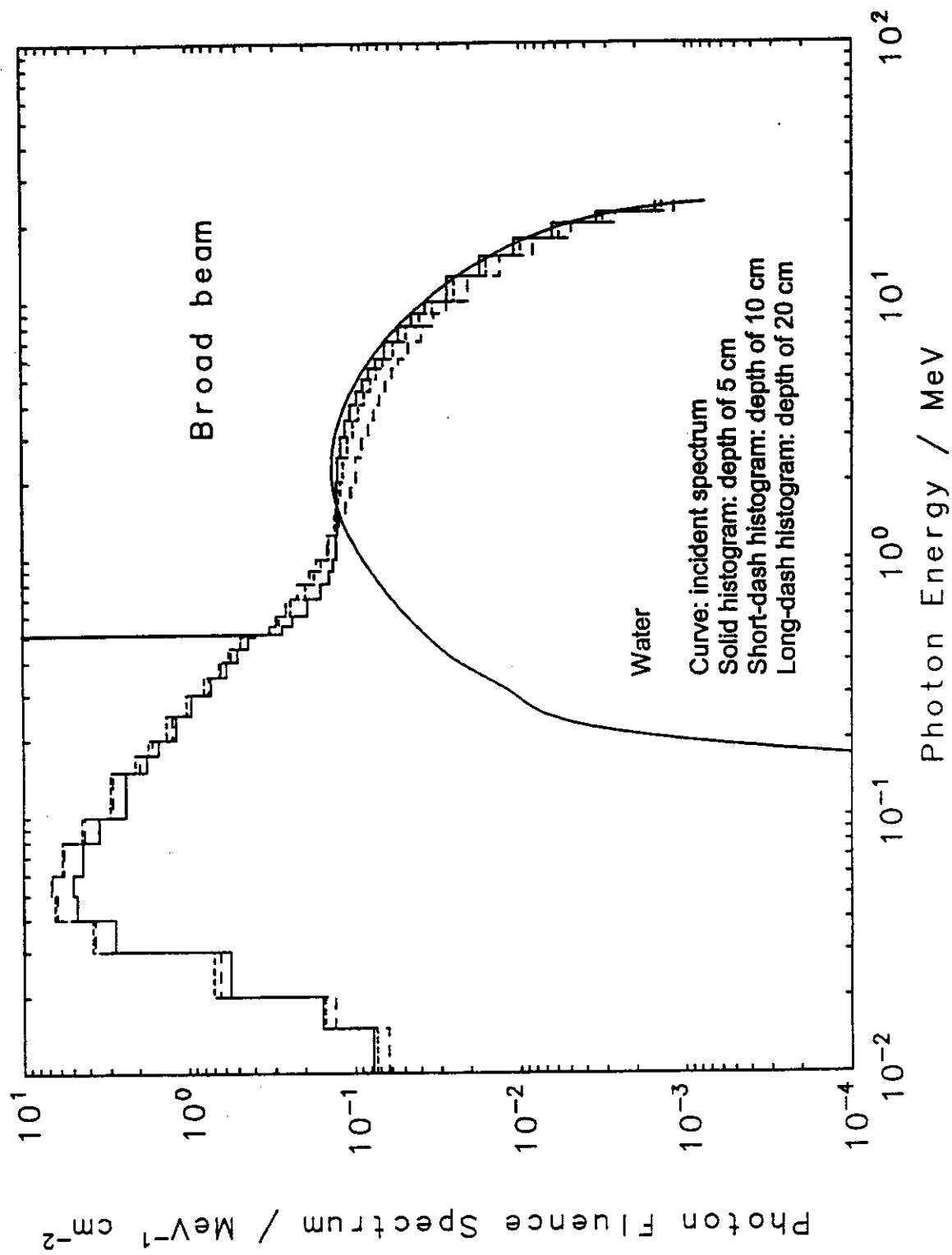
**Example of Monte Carlo calculations for
high-energy photon therapy beams:
scaling theorem**

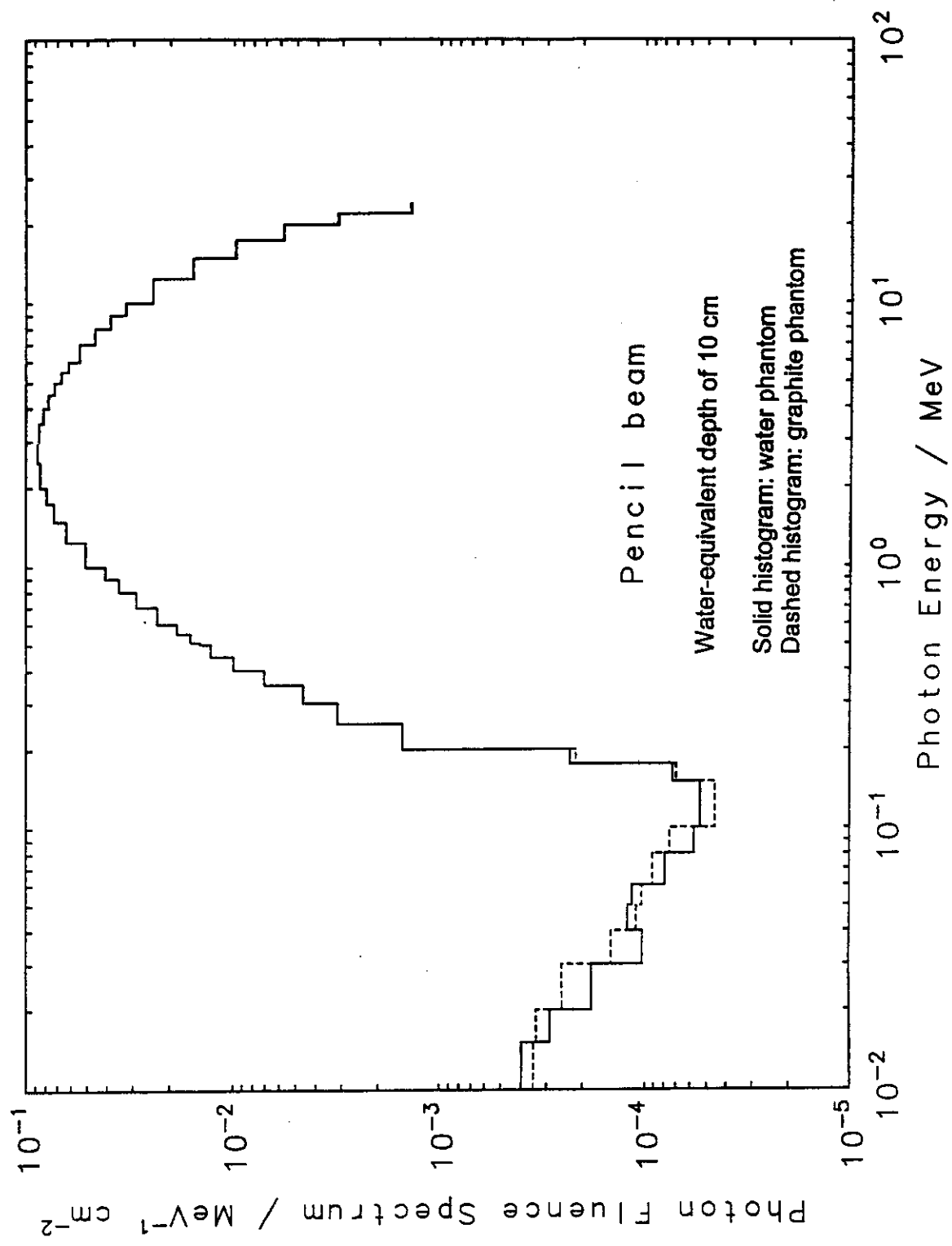
- Scaling: all dimensions scaled by electron density,
 $\varepsilon = \rho N_A Z/A$
 - $x' = (\varepsilon/\varepsilon') x$
 - Holds to extent that cross sections are proportional to Z , such as in Compton scattering
- Example: 24 MV photon beam (no electron contamination)
 - Parallel beams of circular area, 1 cm^2
 - Phantoms of water and of graphite (1.7 g/cm^3), surrounded by vacuum
 - $\varepsilon_{\text{graphite}}/\varepsilon_{\text{water}} = 1.53$
- Monte Carlo calculations with ITSv3.0 (with slight modifications)
 - Results for central axis (0.03 mm^2)

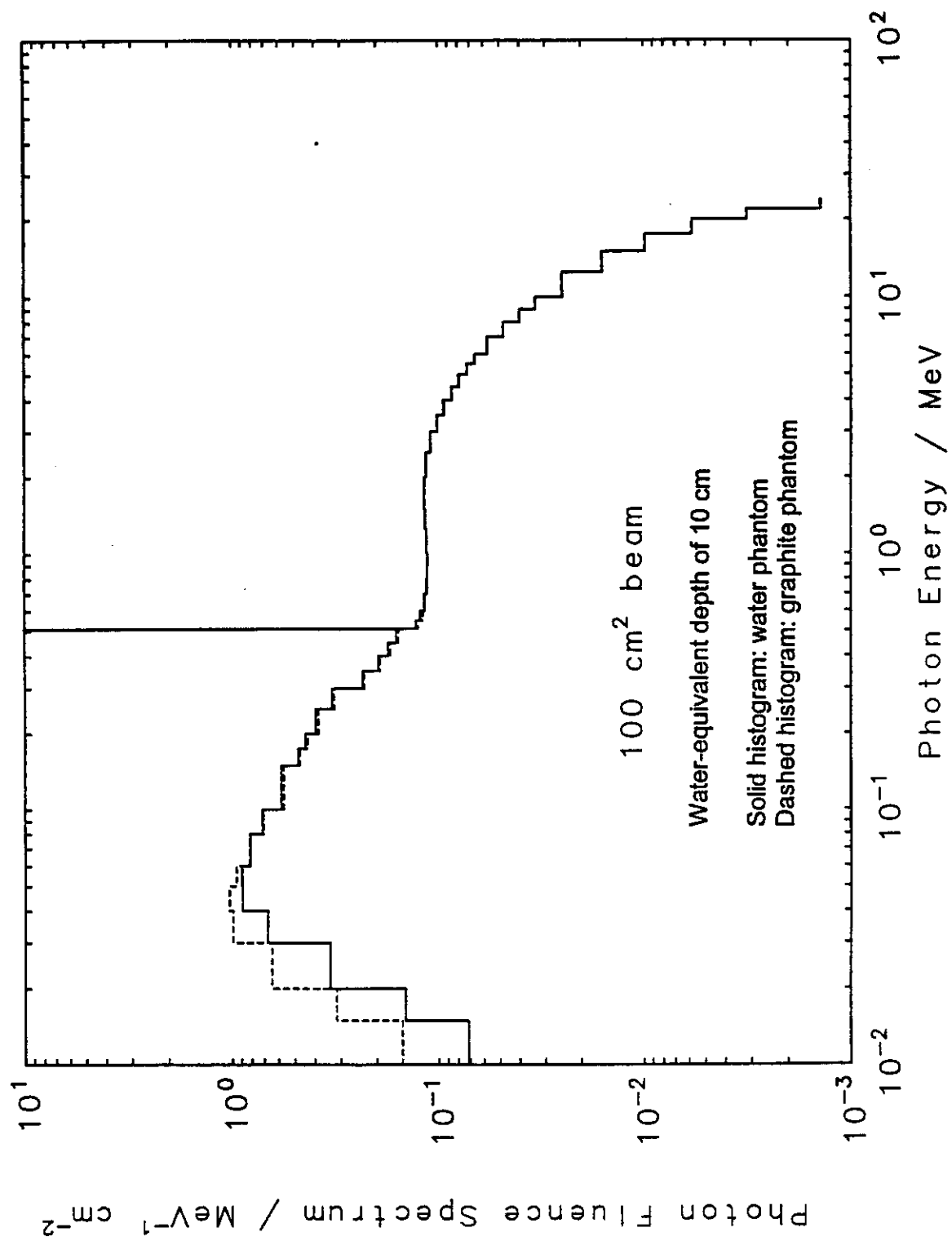


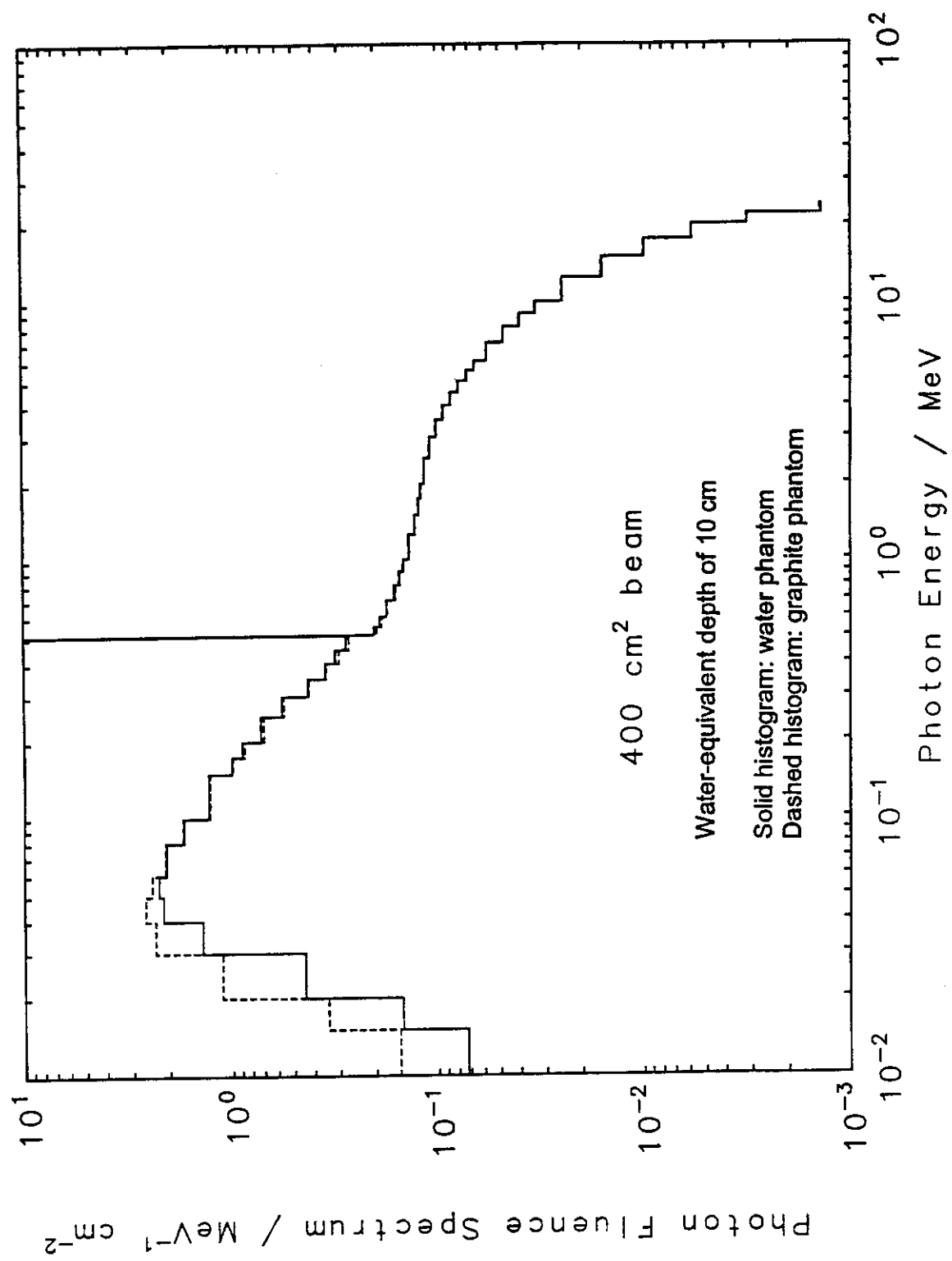


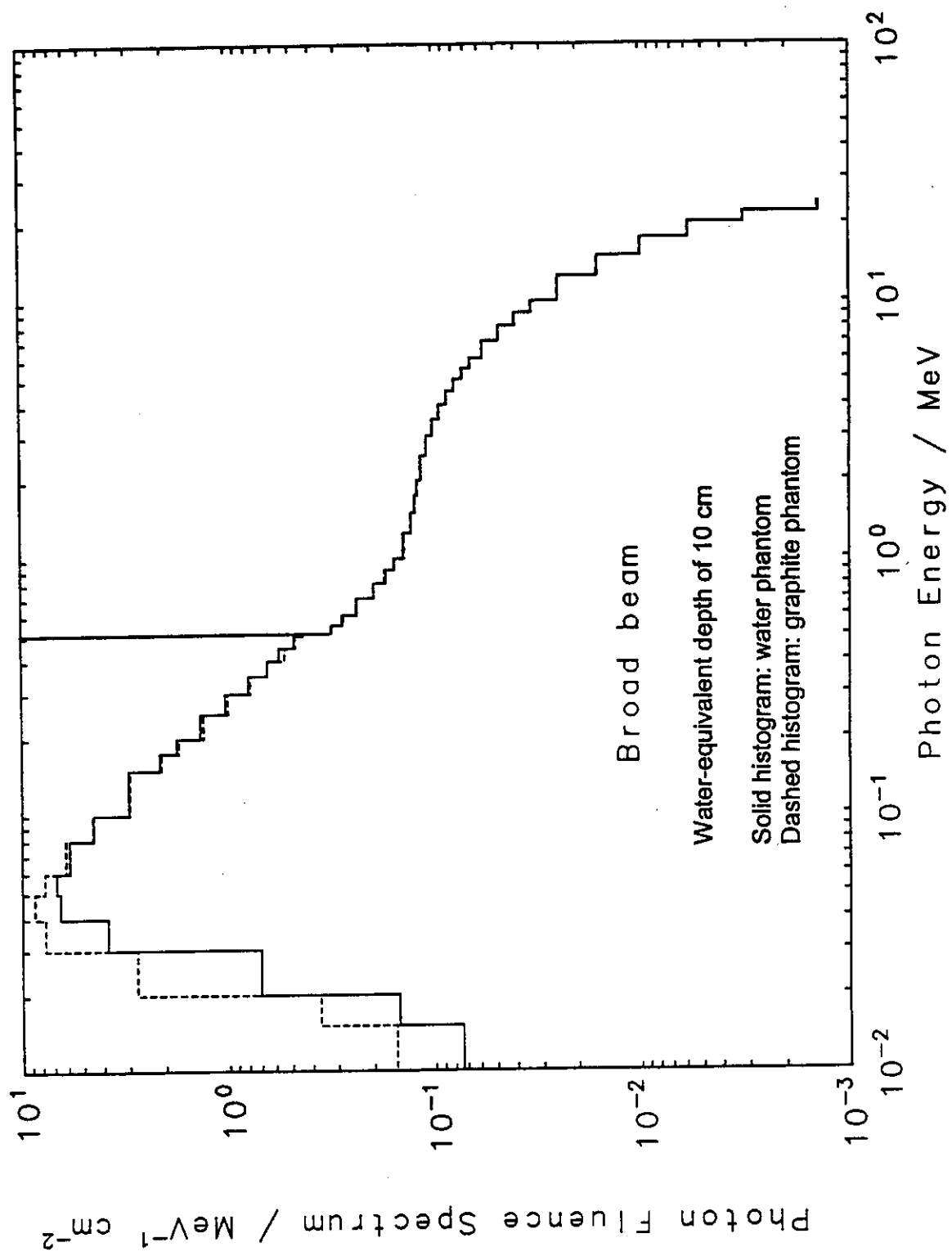


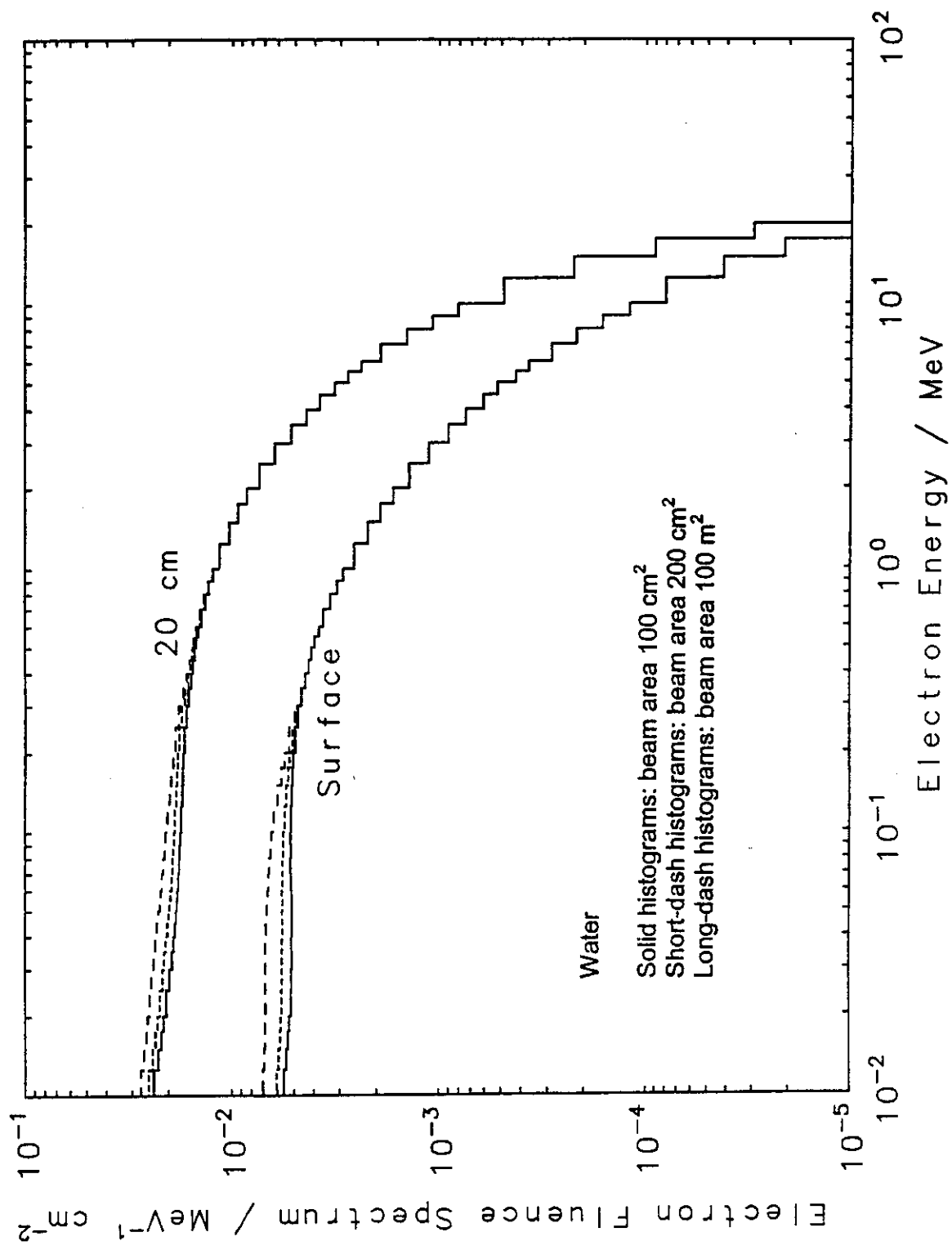


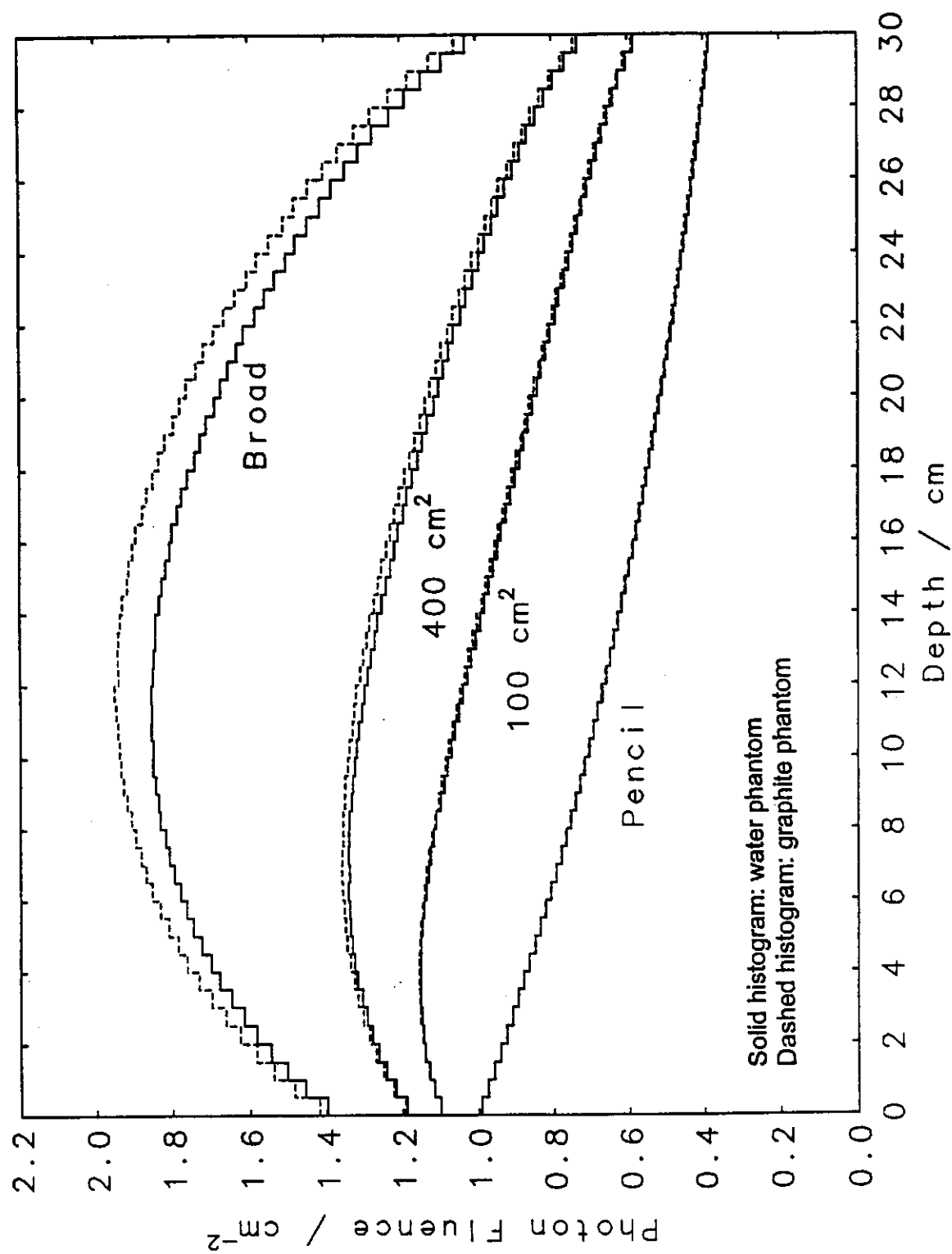


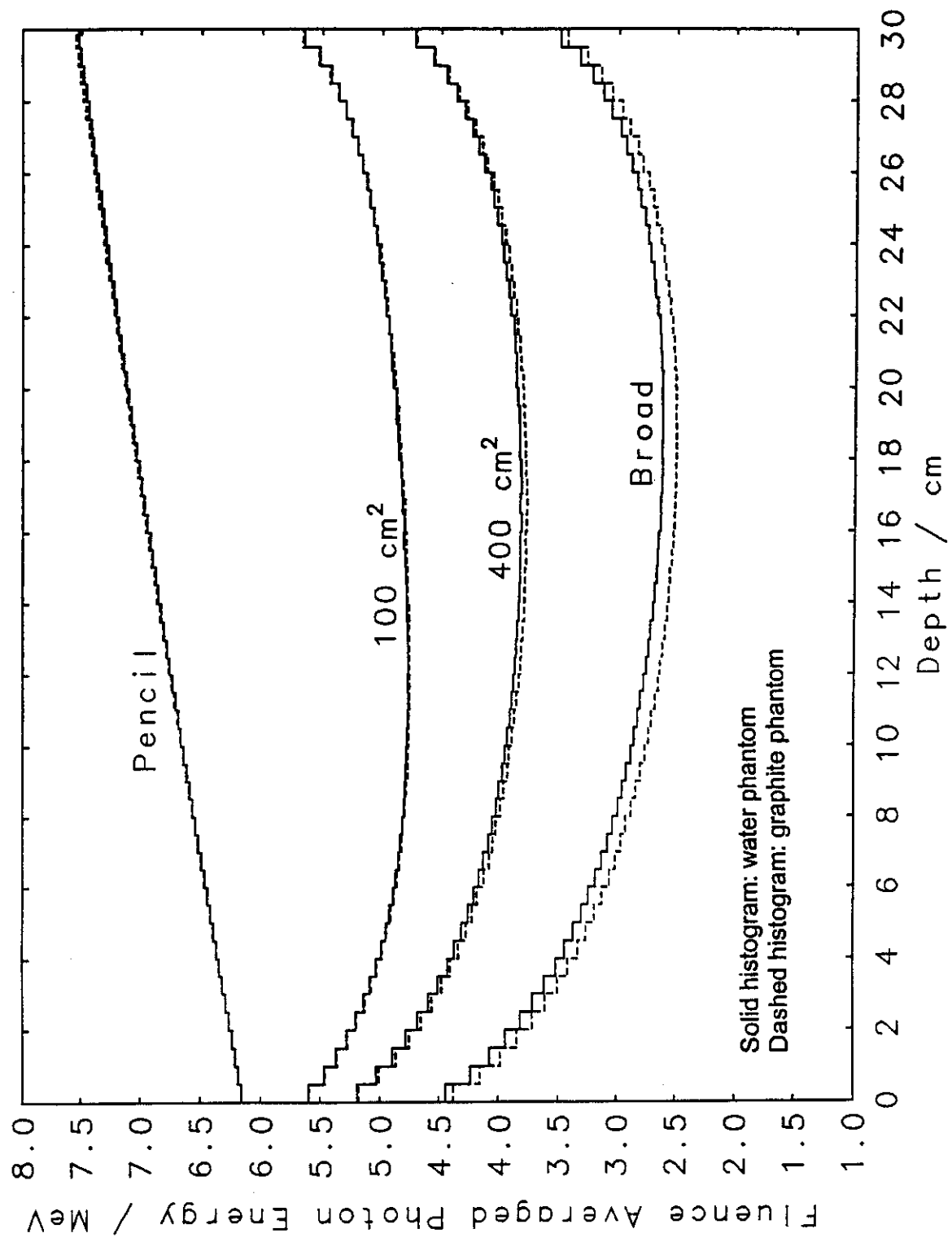


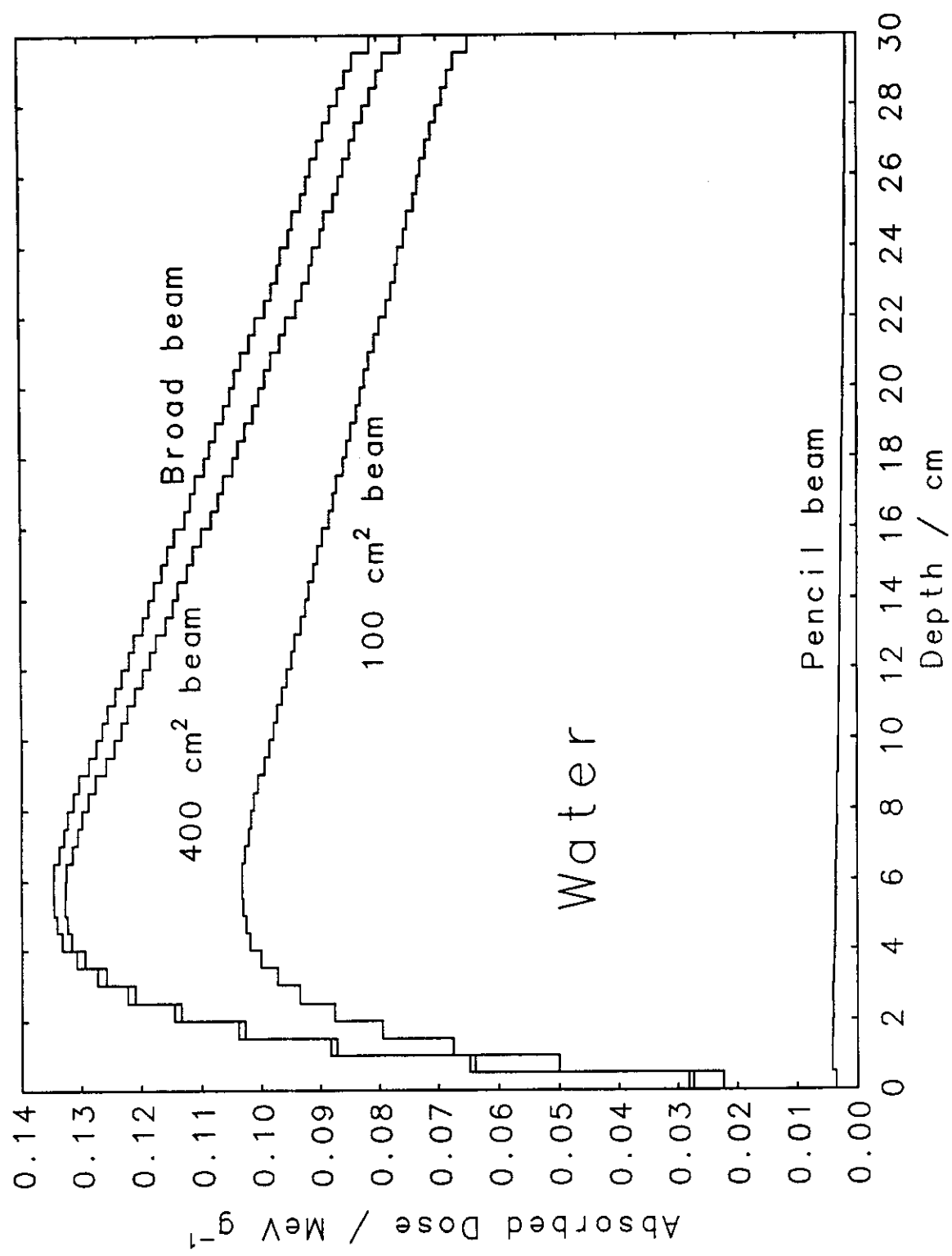




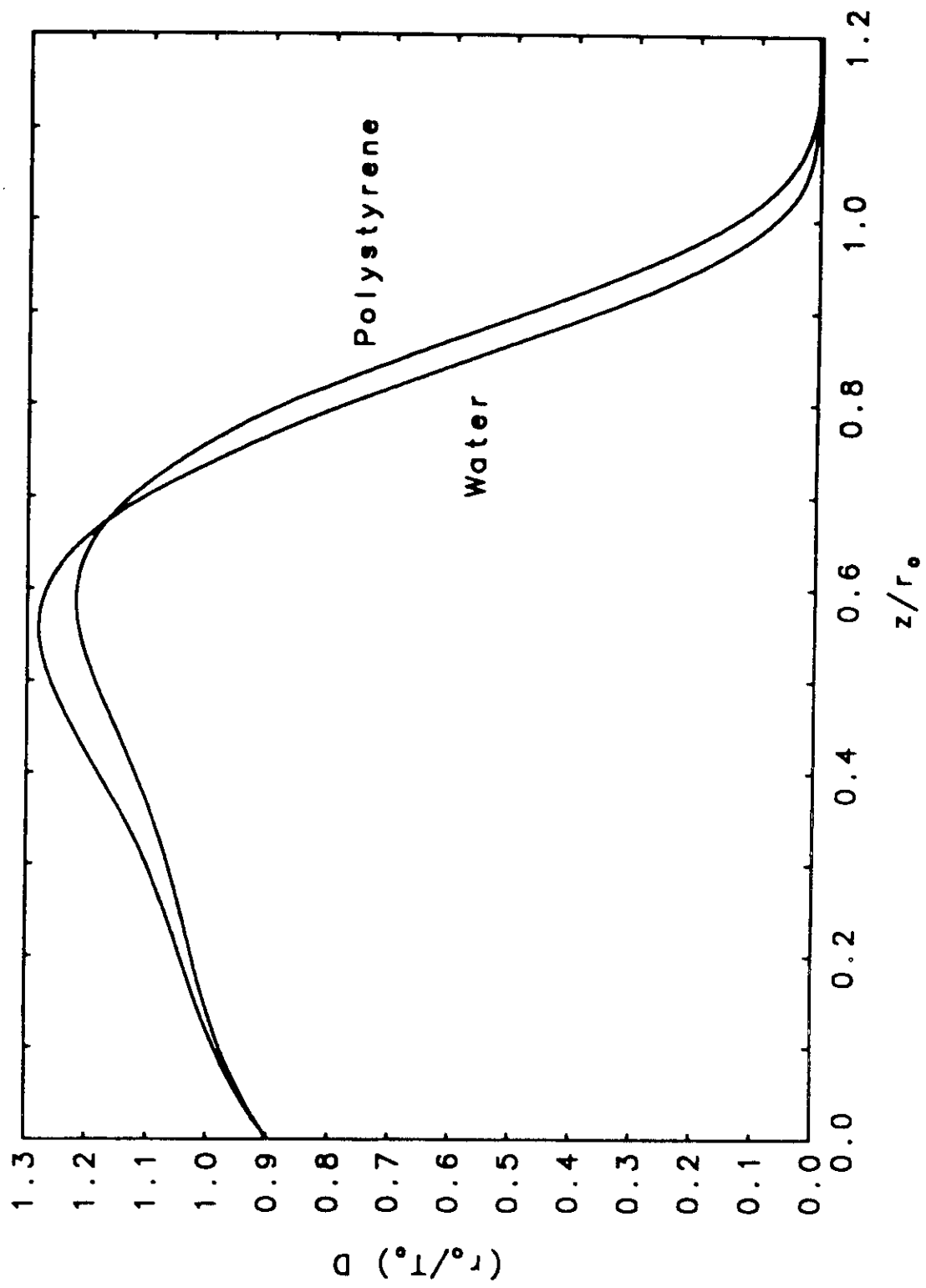


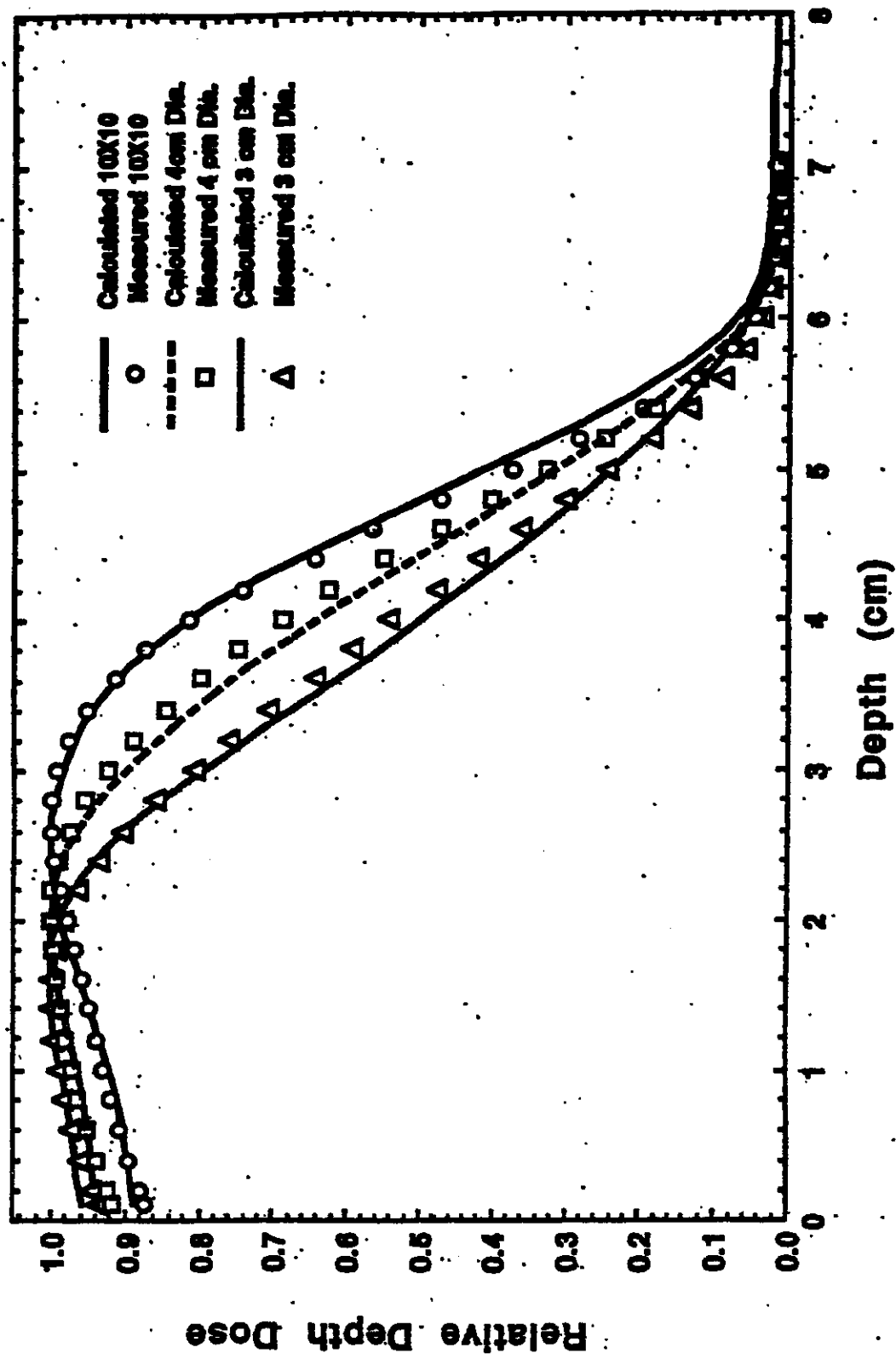




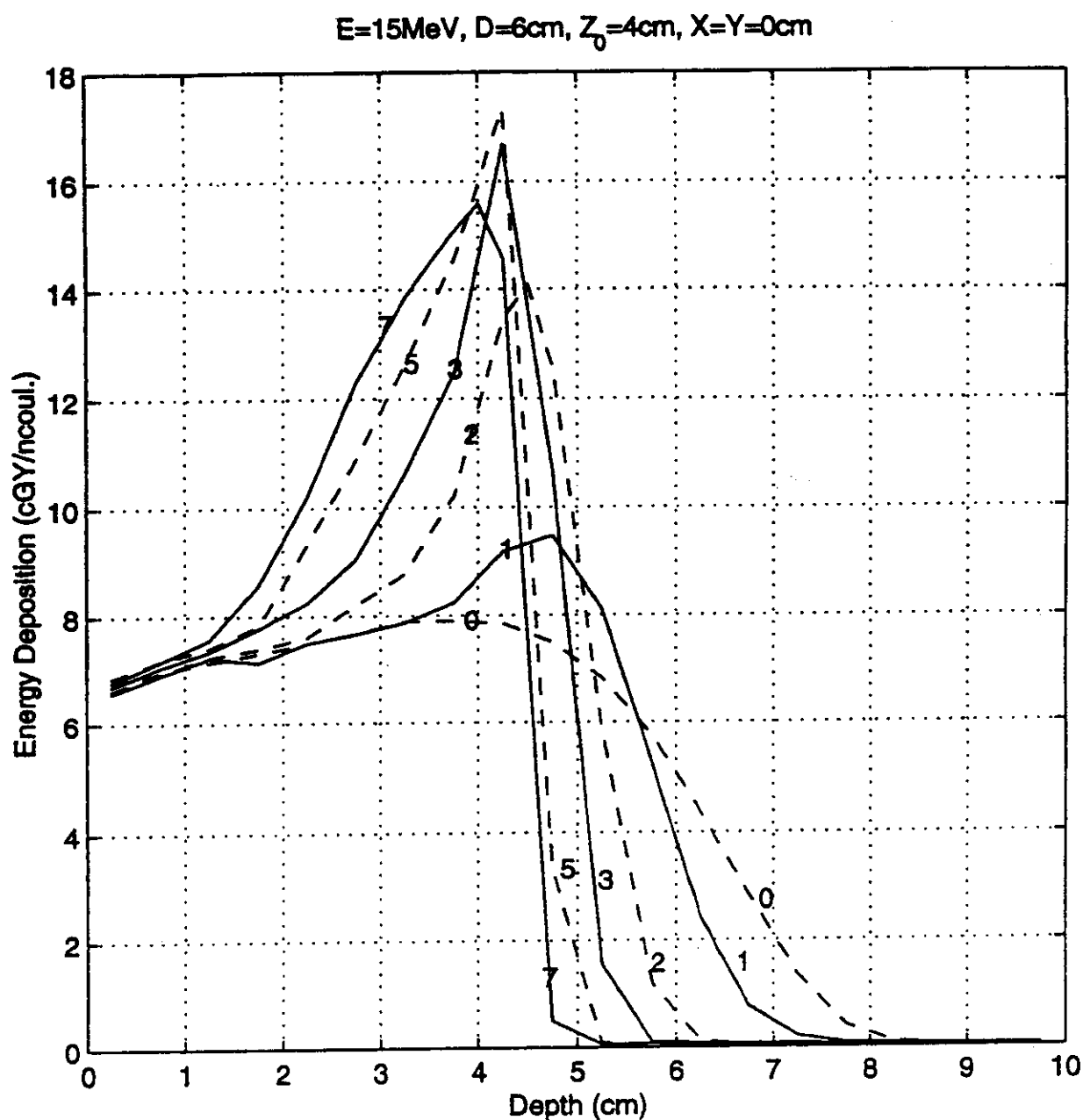


10-MeV Electrons

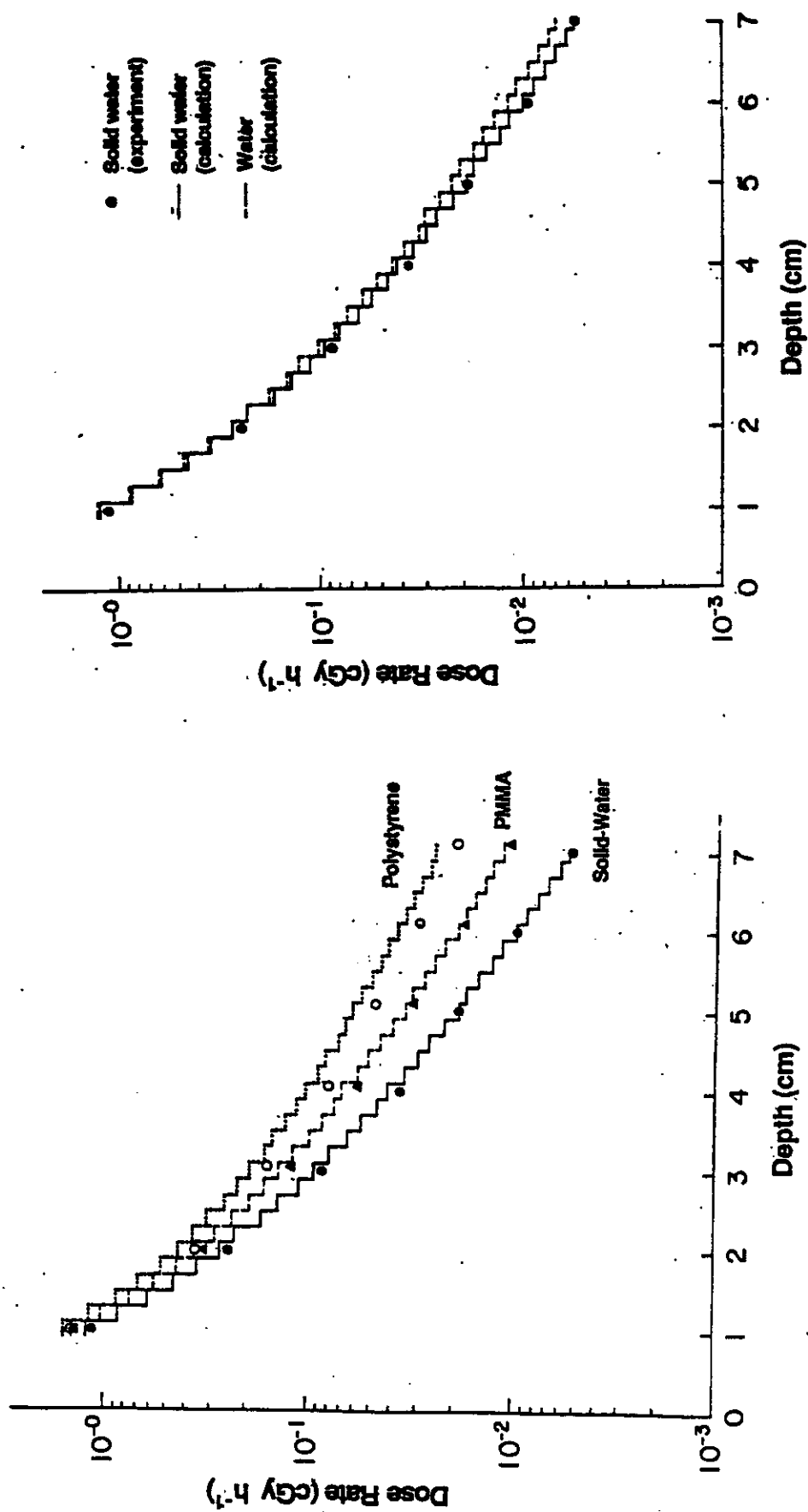




Comparison of central-axis depth-dose distributions from measurement (ion chamber) and Monte Carlo calculations (ITSv3.0) for an electron beam with a nominal energy of 12 MeV (Varian 2100C), after estimation of the actual electron spectrum emerging from the cone applicator. (From Kassaei et al., 1994).



Central-axis depth-dose curves for electron beams incident with an energy of 15 MeV and a beam diameter of 6 cm. The initial 4 cm layer of tissue is field-free, after which there is applied a uniform magnetic field perpendicular to the beam axis. Results are from the ITS v. 3 Monte Carlo code ACCEPTM, and are shown for magnetic fields of 0, 1, 2, 3, 5 and 7 T. From Nardi and Barnea (1999).



Comparison of results from measurement (TLD) and from Monte Carlo calculation (ITSv1.0) for the absorbed-dose rate along the perpendicular bisector in various phantom materials, from a ~1 mCi ¹²⁵I source (model 6702). (From Meigooni et al., 1988)

Endovascular Brachytherapy Sources

Liquid-filled balloon: ^{90}Y , ^{32}P , ^{188}Re

Balloons: 20 to 40 mm length, 2 to 4 mm outside diameter

Polyethylene walls: 0.04 to 0.05 mm thick

$^{90}\text{Sr}/^{90}\text{Y}$ seeds

Emitting core: 2.3 mm length, 0.54 mm diameter
($\rho = 2.6 \text{ g/cm}^3$)

Encapsulation: 2.5 mm length, 0.64 mm diameter
($\rho = 8.1 \text{ g/cm}^3$ steel)

Seed train: 12 seeds, with Au marker seeds at ends

^{32}P wire

Emitting core: 27 mm length, 0.24 mm diameter
($\rho = 1.1 \text{ g/cm}^3$)

Encapsulation: 27 mm length, 0.46 mm diameter
($\rho = 6.5 \text{ g/cm}^3$ NiTi)

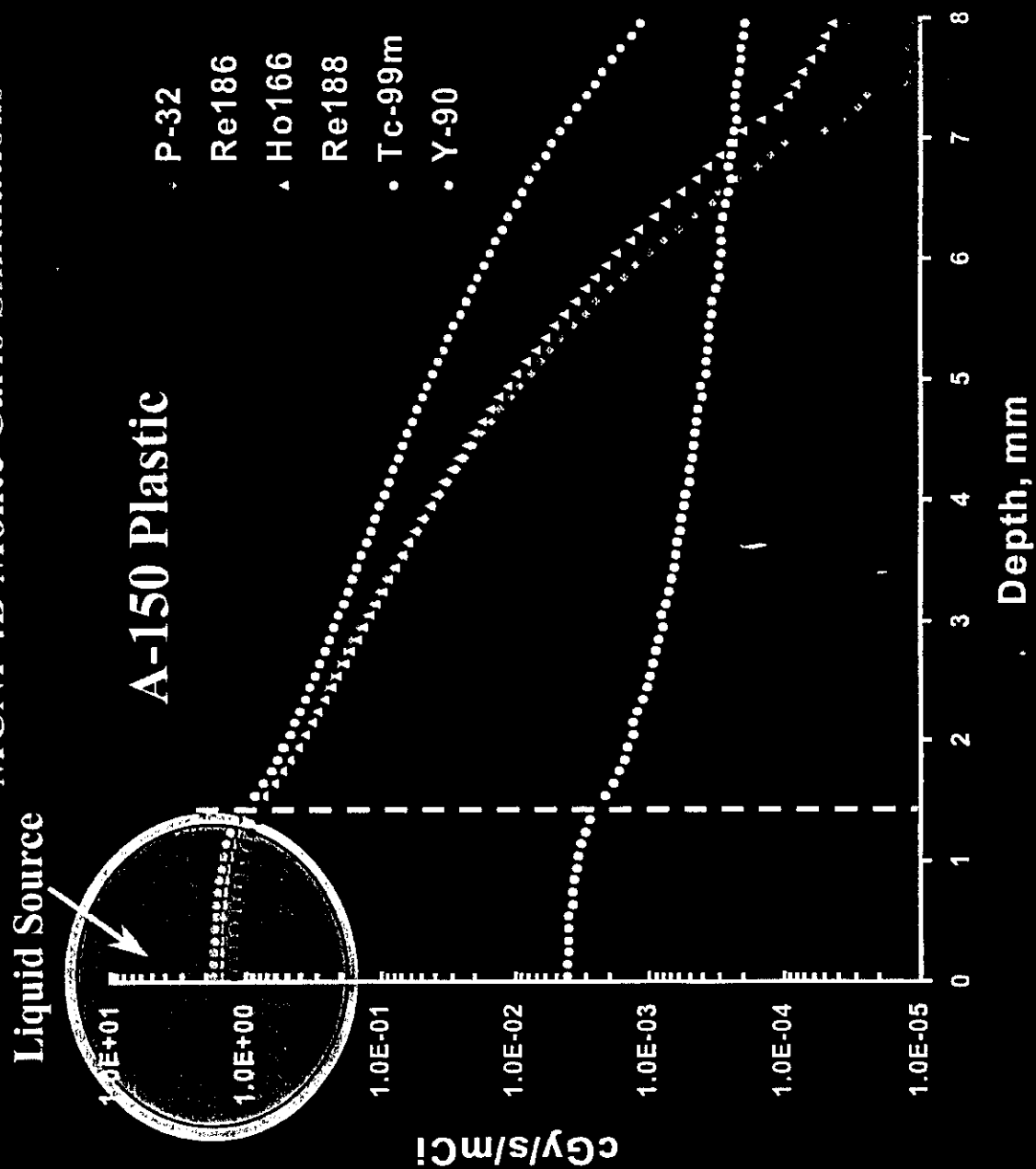
^{192}Ir seeds

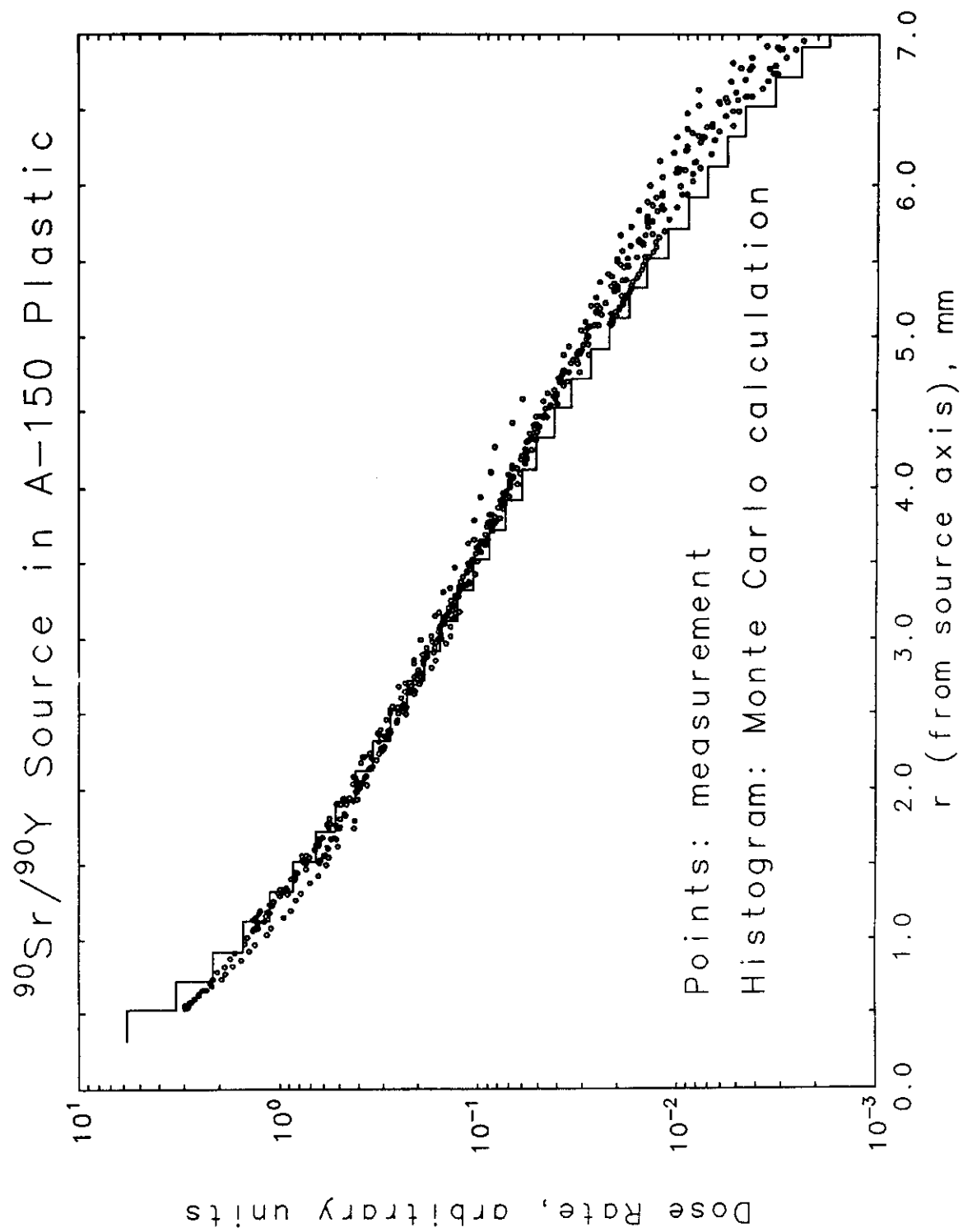
Emitting core: 3 mm length, 0.1 mm diameter
($\rho = 22 \text{ g/cm}^3$ Pt-Ir alloy)

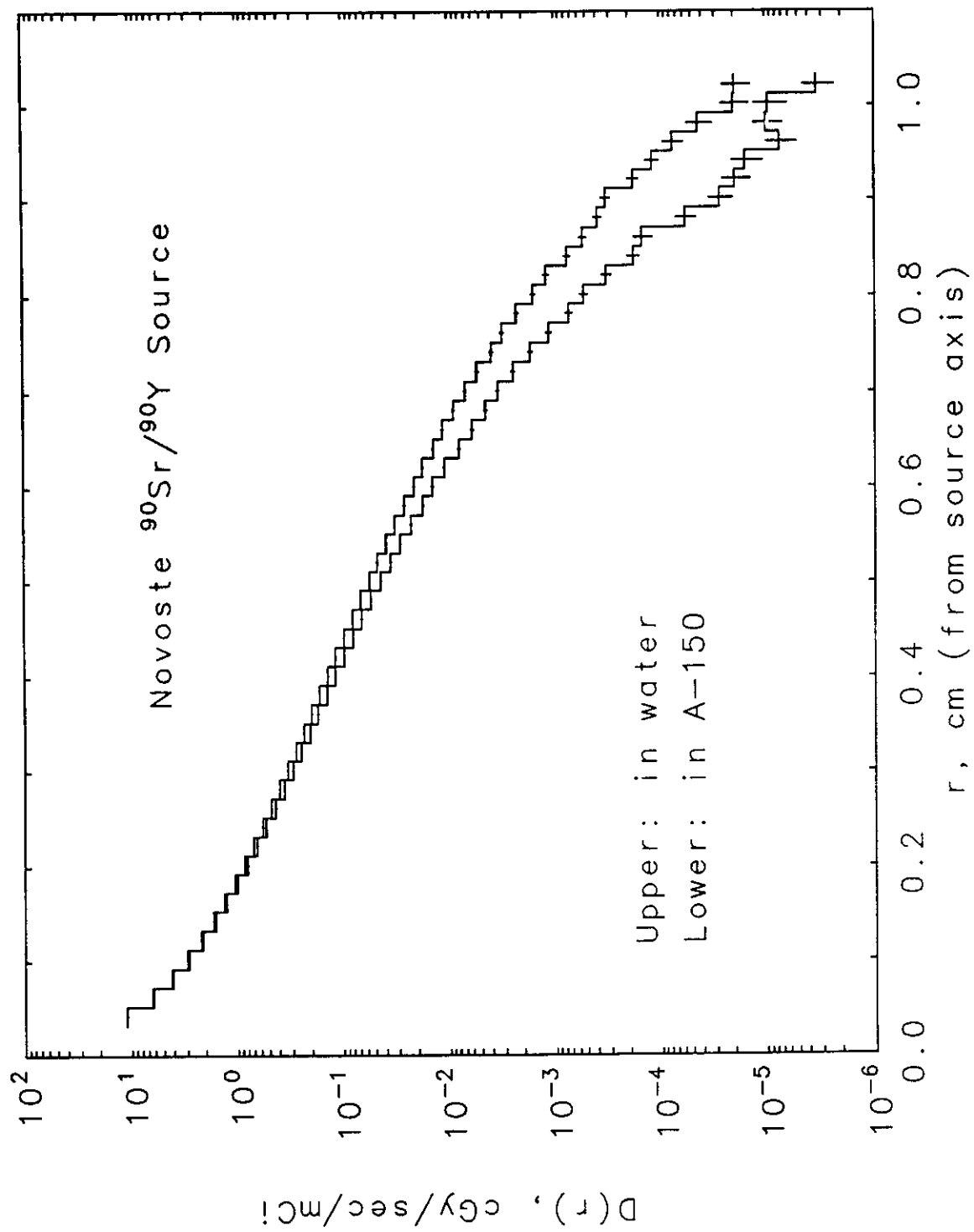
Encapsulation: 3 mm length, 0.5 mm diameter
($\rho = 8.1 \text{ g/cm}^3$ steel)

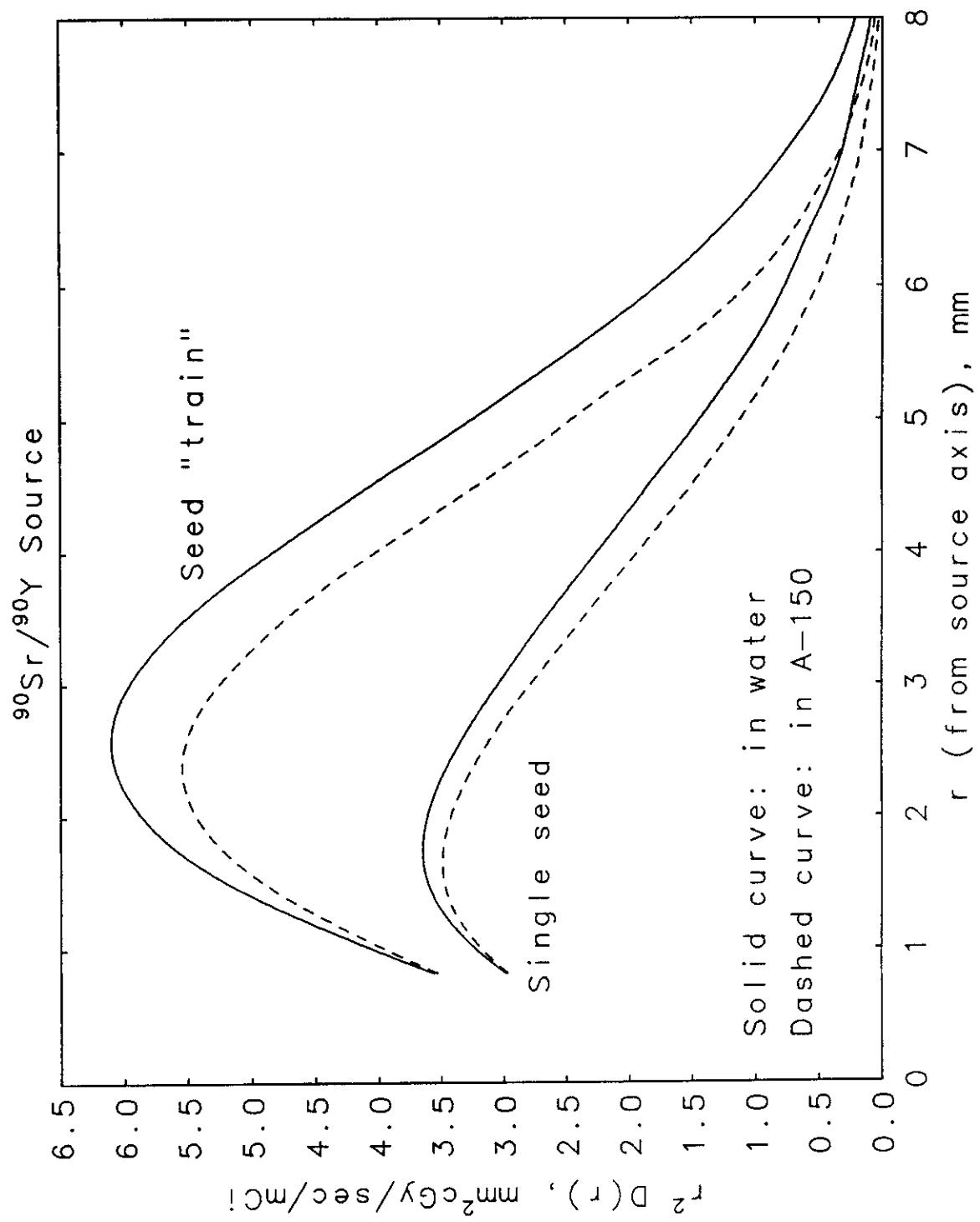
Ribbon: 0.8 mm diameter ($\rho = 1.1 \text{ g/cm}^3$ nylon)

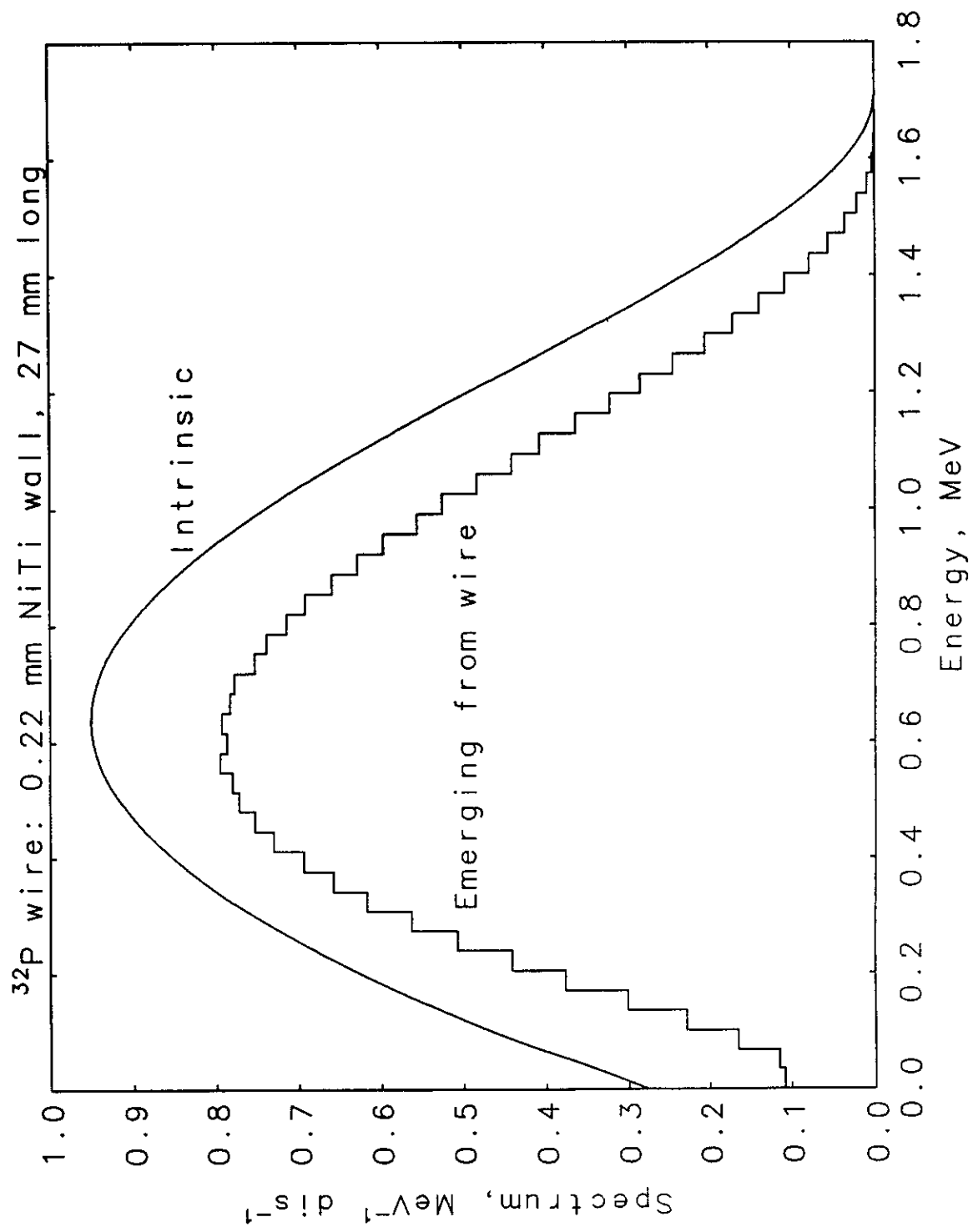
MCNP4B Monte Carlo Simulations

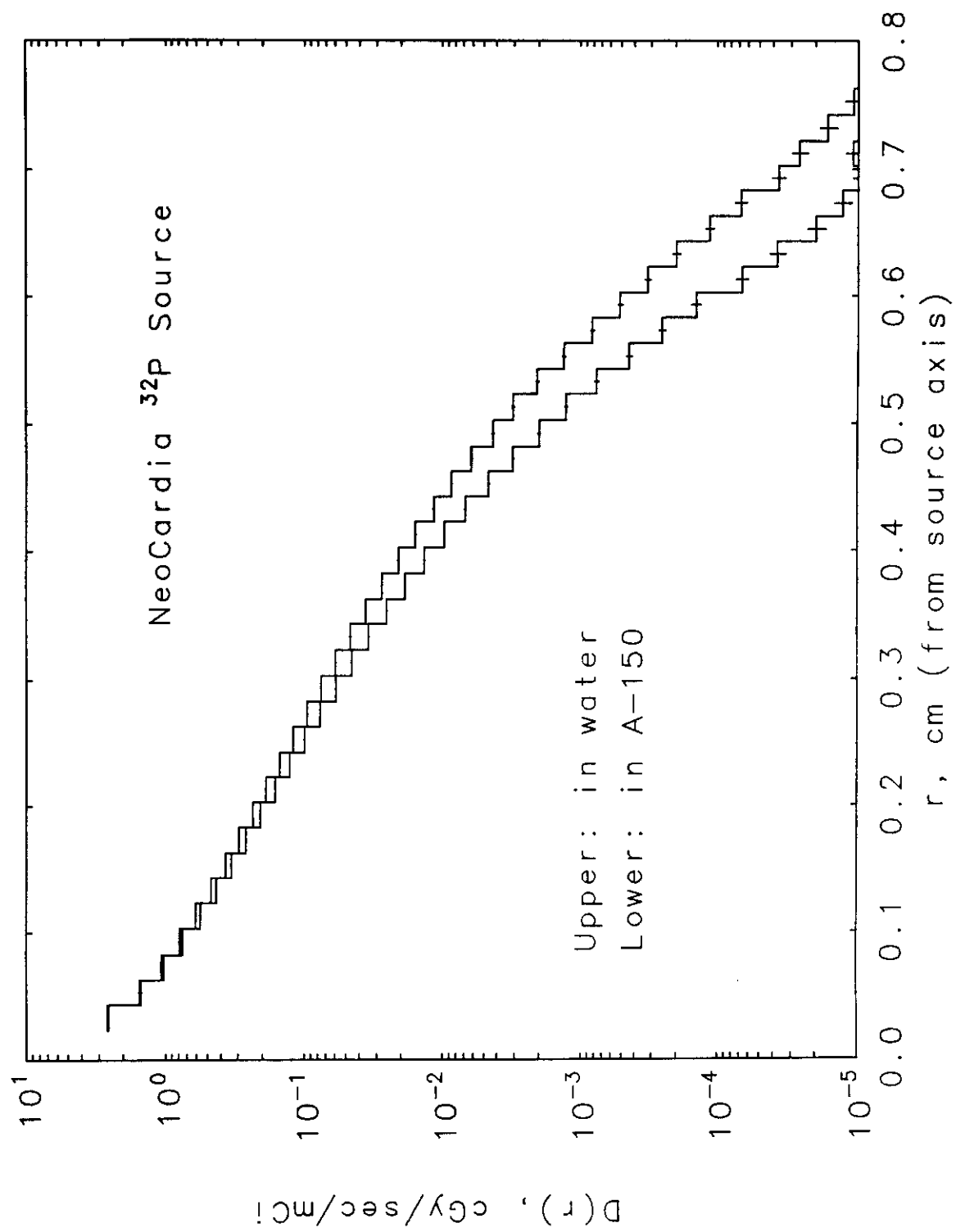












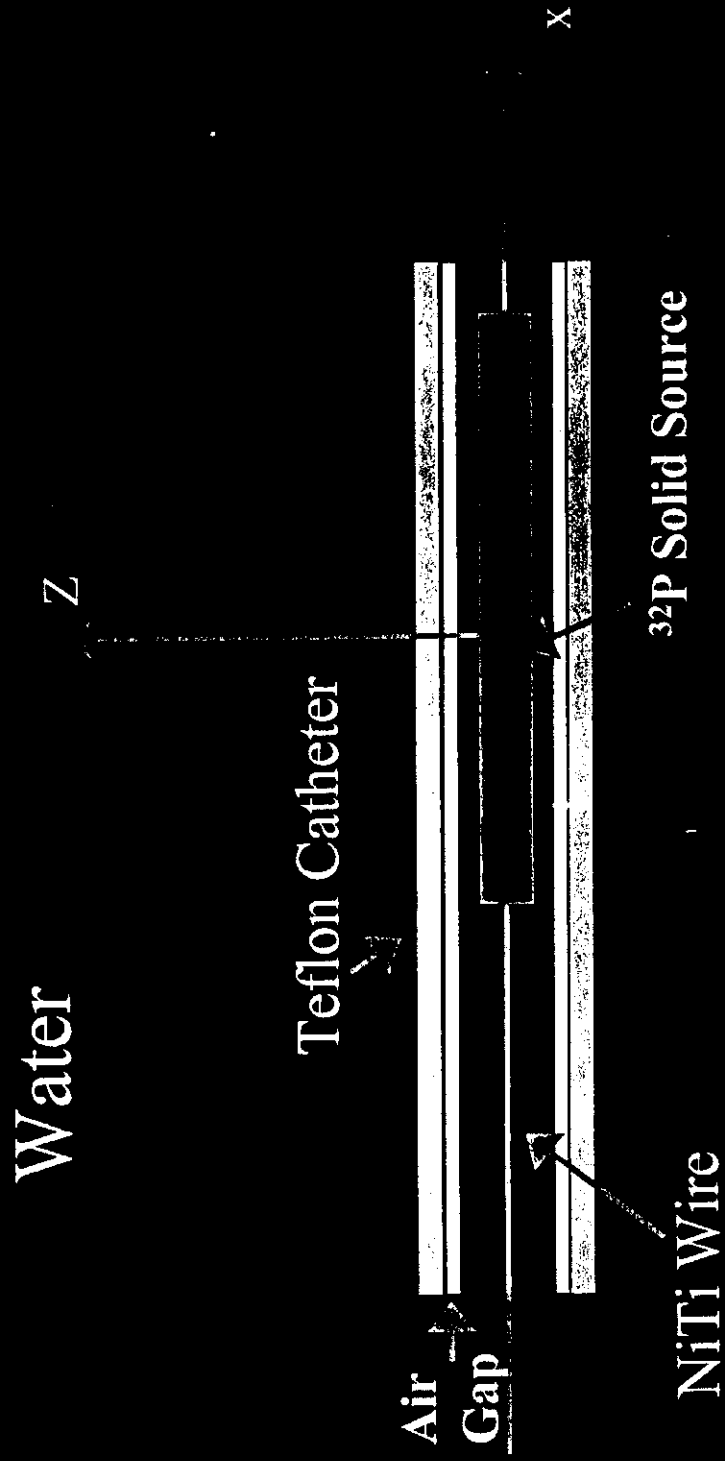
Codes Compared

- Monte Carlo N-Particle transport code (MCNP4B)
- CYLTRAN from the Integrated Tiger Series codes (ITS v.3)
- Electron Gamma Shower code (EGS4)

Codes Common Ground

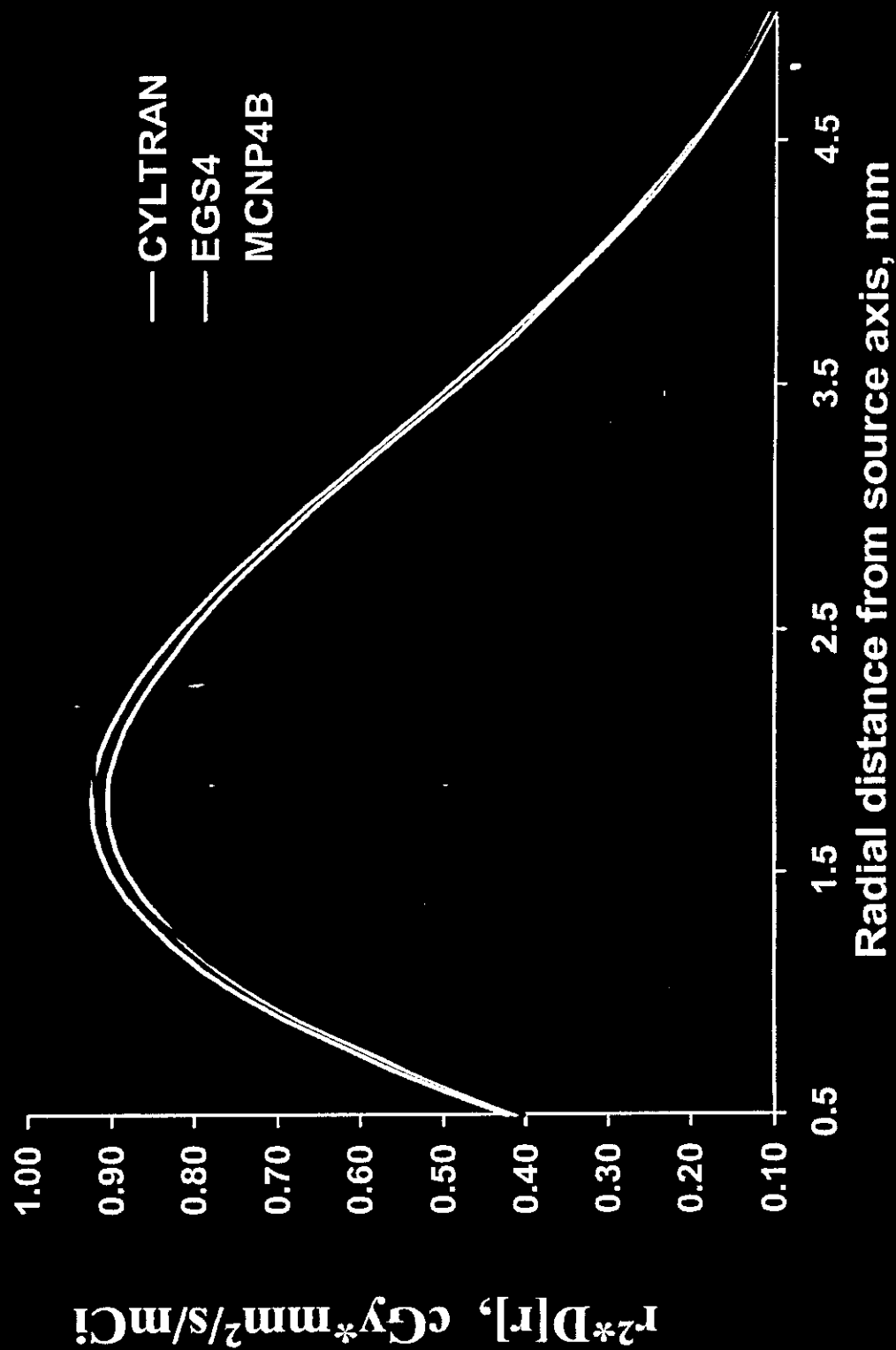
- ^{32}P energy spectrum from data provided by Stephen Seltzer, NIST
- Scoring bins up to 6 mm away from the wire center (0.1 mm thick x 15 mm long)
- Geometry and material composition of wire and catheter provided by Sam Lott, Guidant Inc.

^{32}P Wire Source Geometry



not to scale

Results for ^{32}P wire in water



Point kernels

- describe the distribution of absorbed dose around a point, isotropic emitter (*e.g.*, a radioactive atom),
- in an infinite, homogeneous medium,
- and are usually expressed in some scaled form.
- are a function only of single spatial variable, r

Uses

- by superposition, can apply to source distributed in energy and/or space
- originally developed for nuclear medicine applications

Advantages

- very fast, semi-analytic with no statistical fluctuations

Disadvantages

- does not account for boundaries and inhomogeneities,

Point-Kernel Method: Electrons

- Point kernel for monoenergetic electrons in water

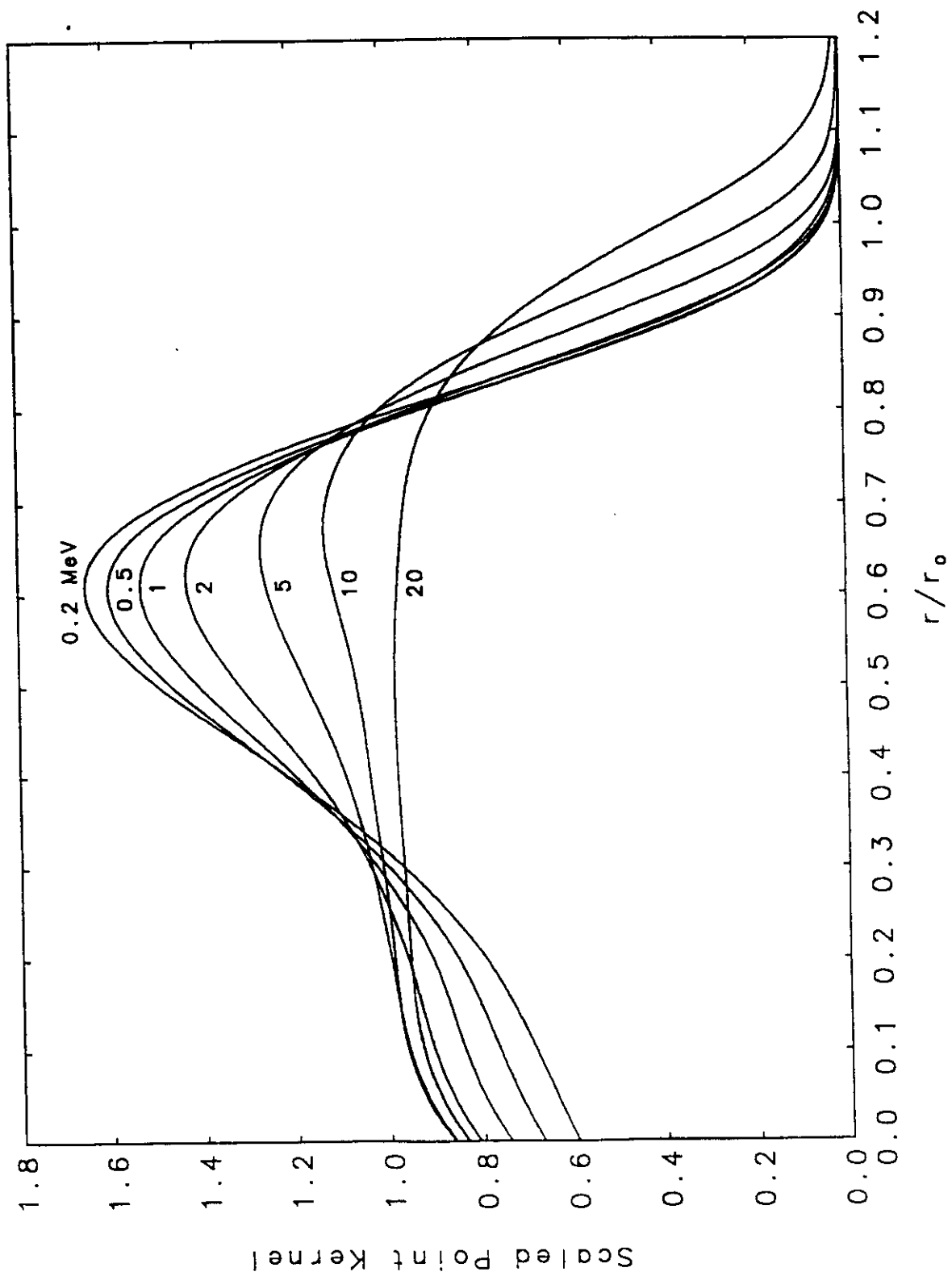
$$F_o(r) = 4\pi r^2 \rho_w D_o(r) / E_o$$

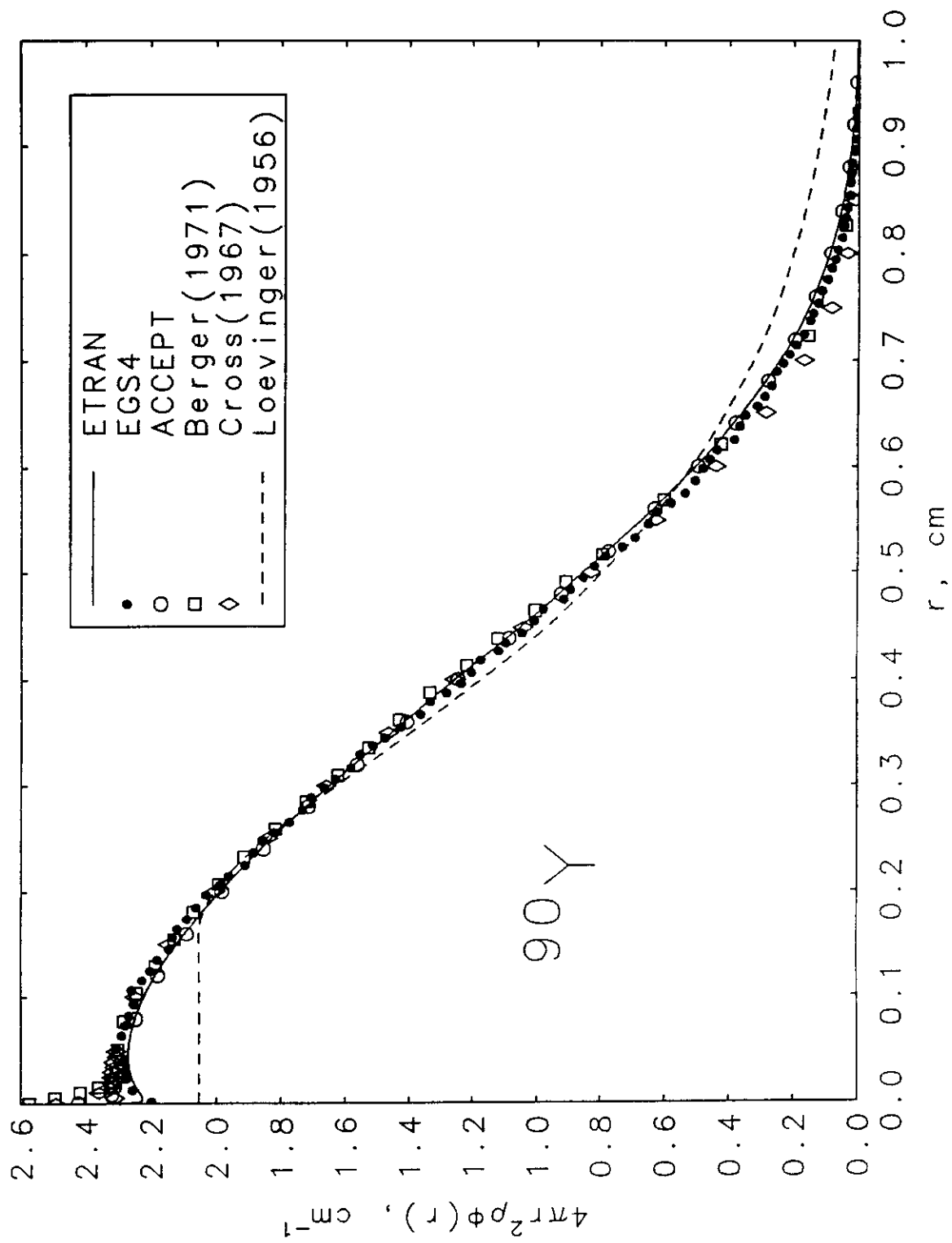
- Point kernel for a beta spectrum in water

$$F(r) = \int_0^{E_{\max}} (E_o / E_{\text{av}}) S(E_o) F_o(r) dE_o$$

- Dose (or dose rate) from uniform volume source (per electron)

$$D(r) = \frac{E_{\text{av}}}{4\pi\rho V} \int_V \frac{F(|\vec{r}-\vec{r}'|)}{|\vec{r}-\vec{r}'|^2} dV$$





Point-Kernel Method: Electrons

- Include non-water materials along each ray r - r' by scaling approximation

$$F(r) = \eta F_w(\eta r \rho / \rho_w)$$

where $\eta = \eta(Z_{\text{eff}})$, and $Z_{\text{eff}} = [\sum \omega_i Z_i^2 / A_i] / [\sum \omega_i Z_i / A_i]$.

- The scaling parameter for a range of applications has been developed by Cross (ICRU Report 56).
- Simple density scaling: $\eta = 1$.

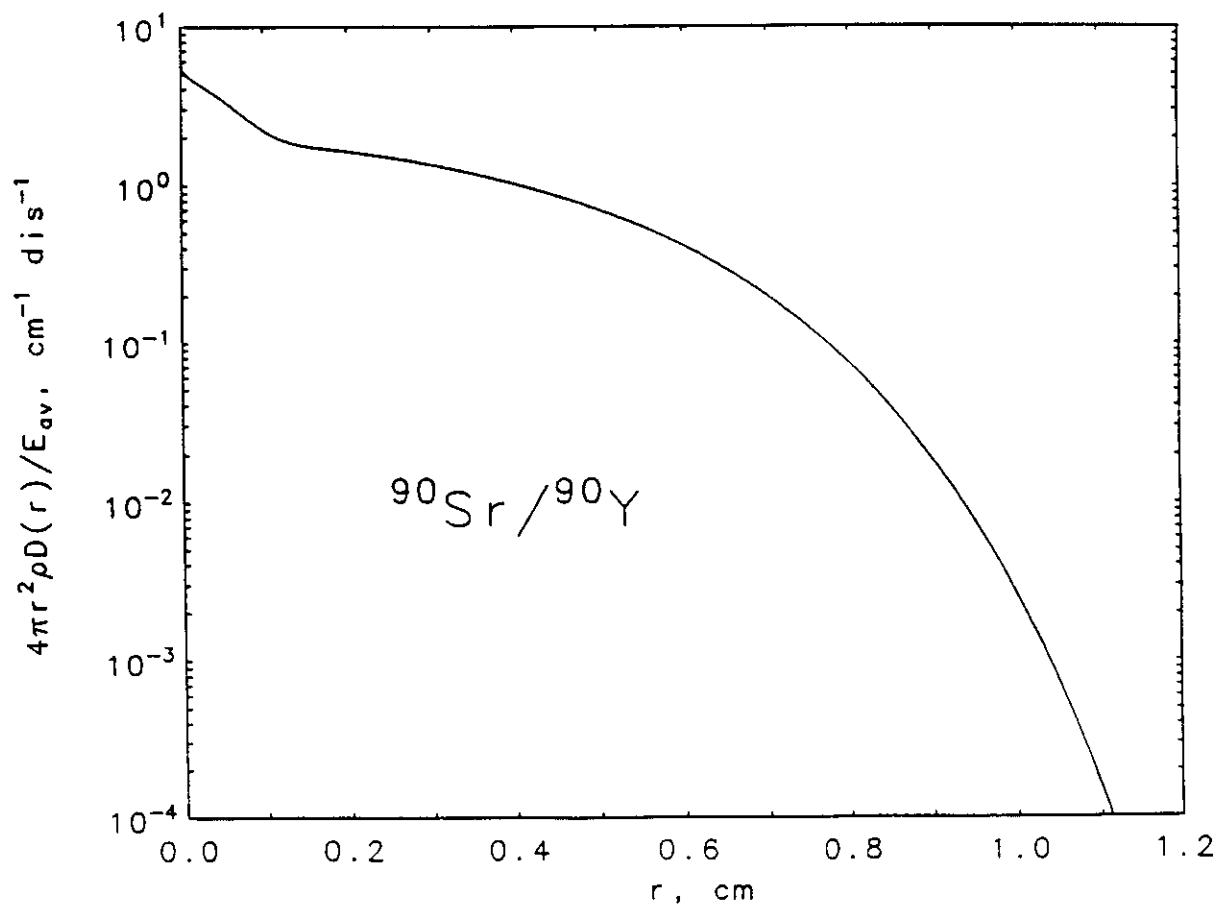
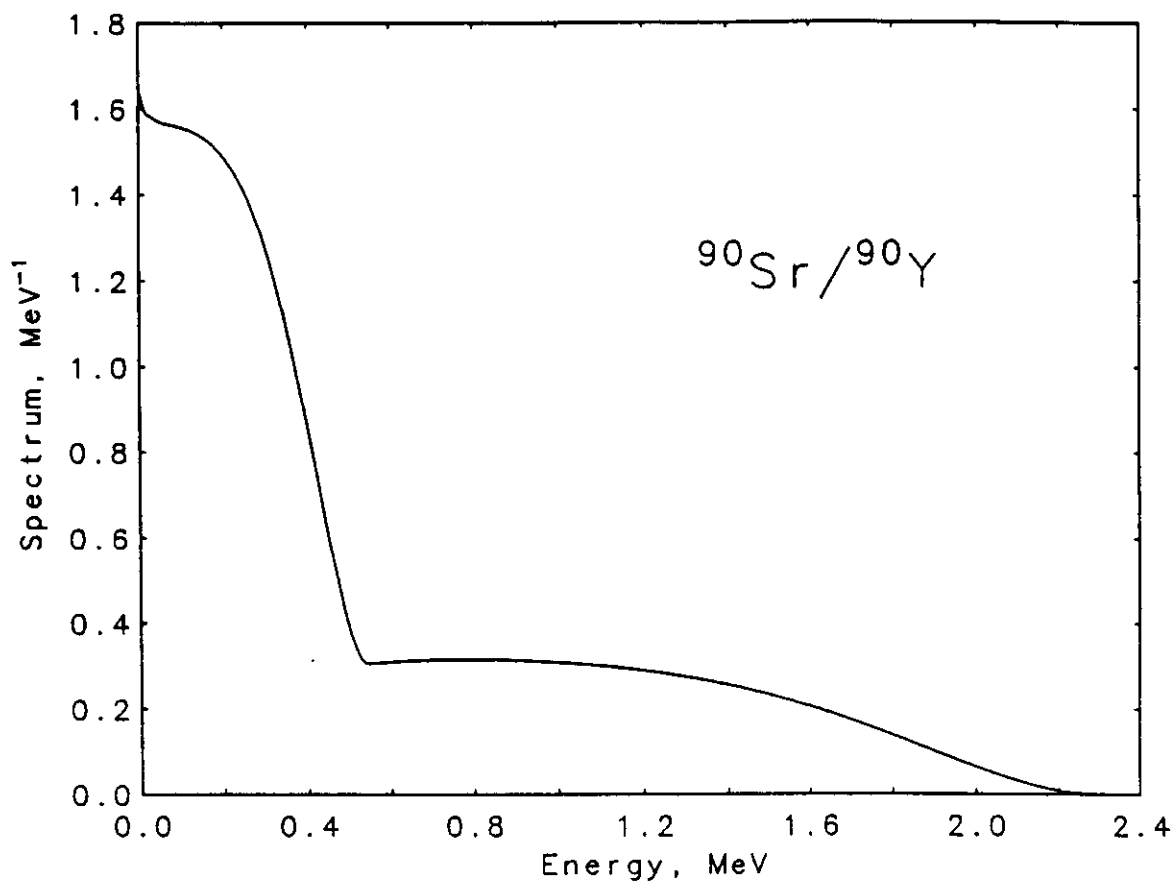
Scaling parameters:

$$\eta_{Cross} = \left[0.777 + \langle Z \rangle (0.03756 - 0.00066 \langle Z \rangle) \right] \frac{r_0^{water}(500 keV)}{r_0^{med}(500 keV)}$$

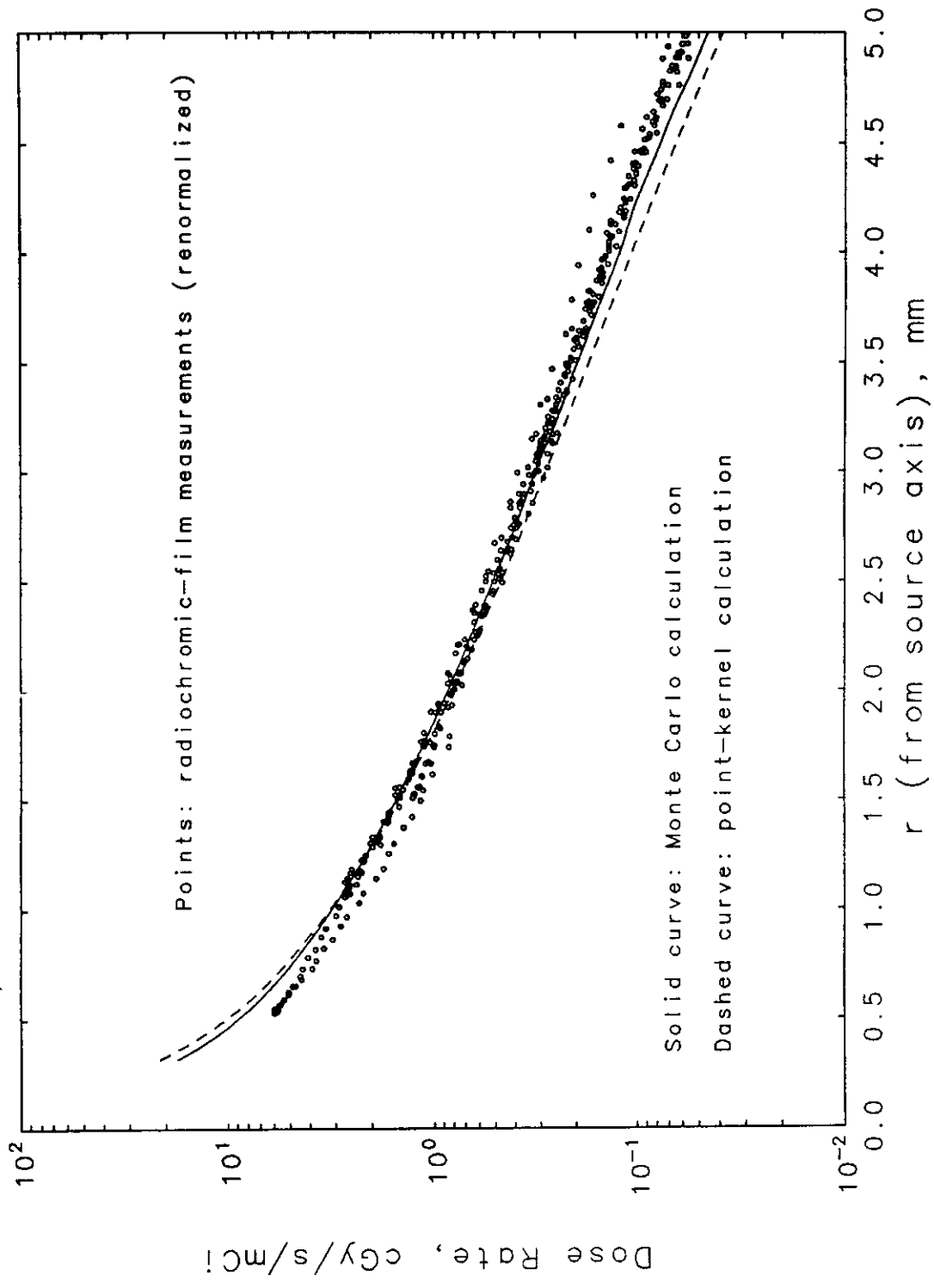
$$\eta_{Berger} = \left[1 + 0.02252 (\langle Z \rangle - 6.60) \right] \frac{S_{med}(200 keV)}{S_{water}(200 keV)}$$

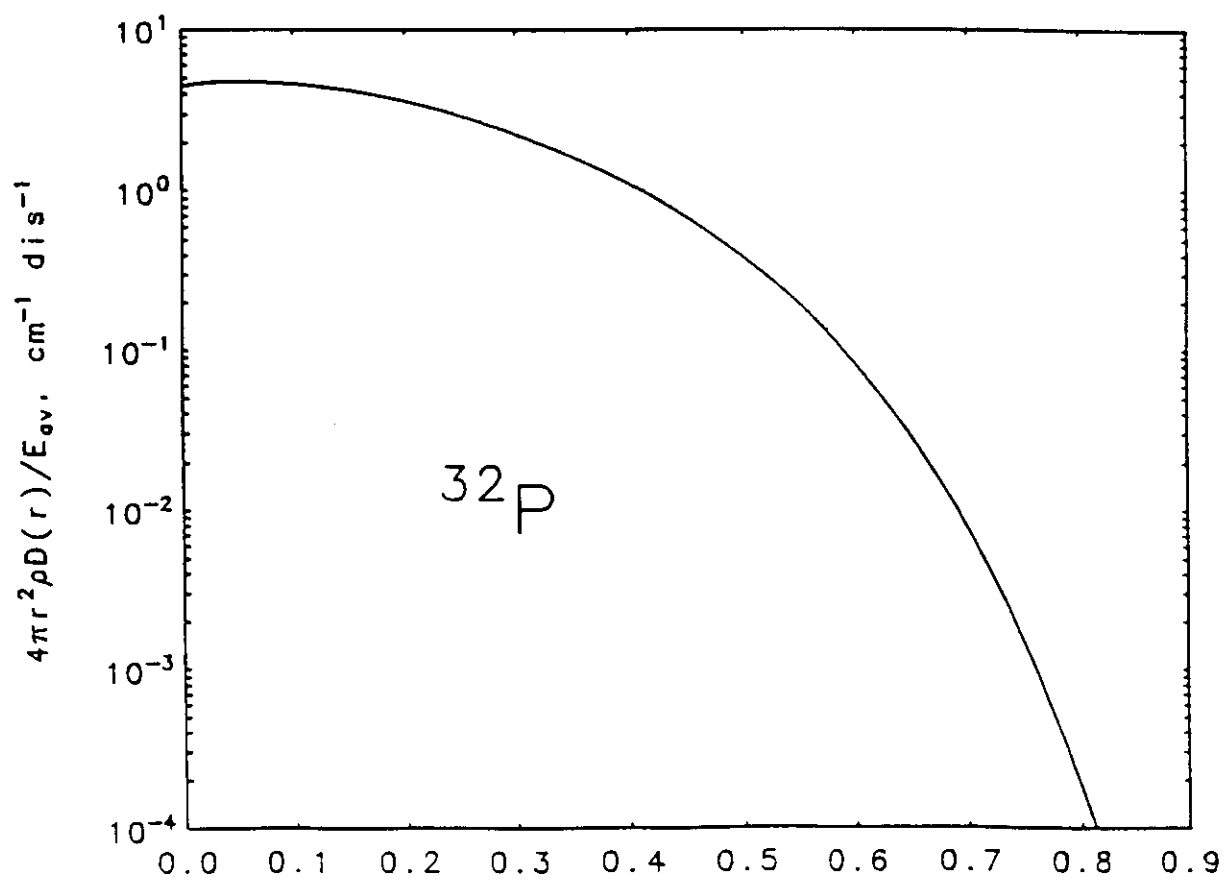
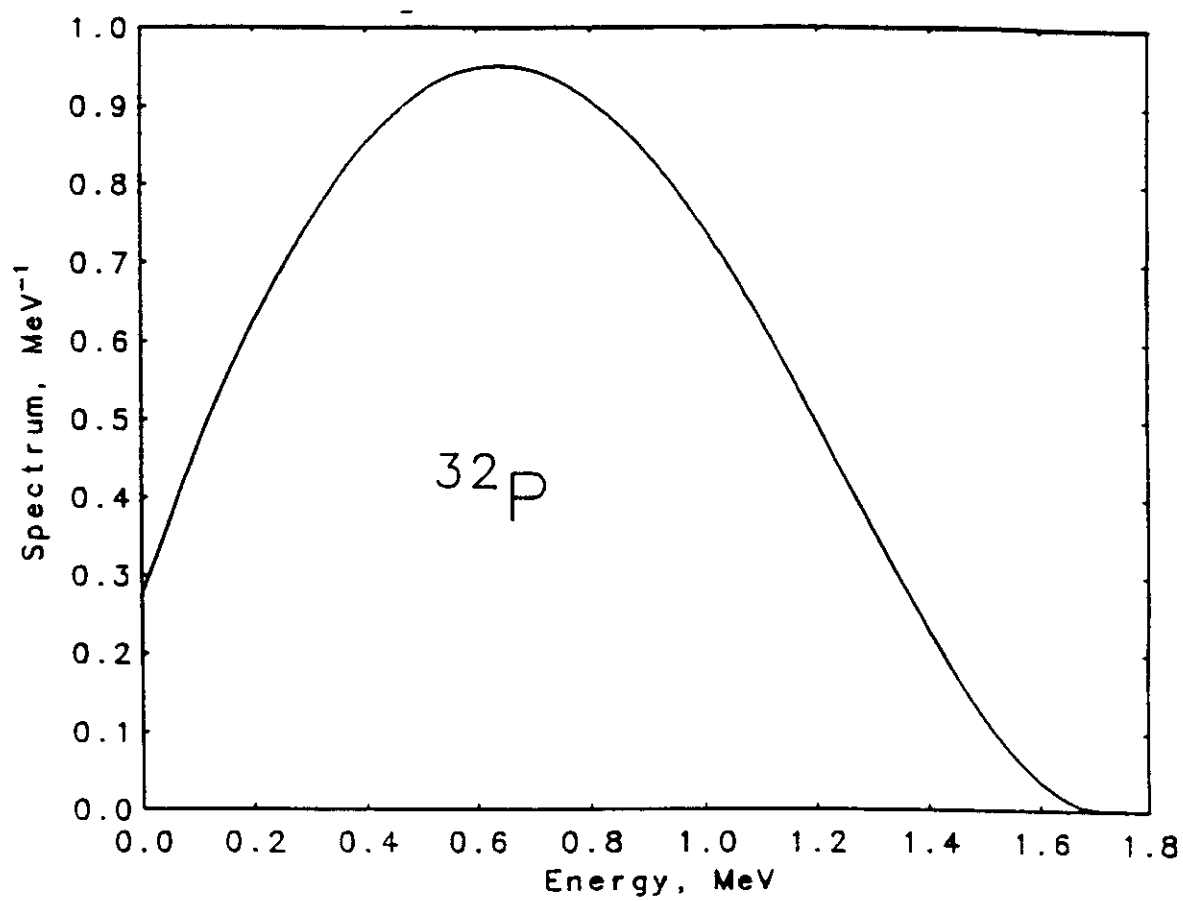
where S is mass collision stopping power

Material	$\langle Z \rangle$	ρ (g/cm ³)	η (Cross)	η (Berger)
polyethylene	4.75	0.924	0.997	1.018
nylon 12	5.16	1.01	0.980	0.997
polystyrene	5.29	1.06	0.936	0.950
A150	5.49	1.127	0.968	0.982
polyurethane	5.74	1.09	0.965	0.975
PMMA	5.85	1.19	0.948	0.957
WT1	5.95	1.02	0.957	0.966
kapton	6.22	1.42	0.903	0.909
mylar	6.24	1.40	0.918	0.924
tissue (ICRU 4 component)	6.50	1.0	0.984	0.988
water	6.60	1.0	1.0	1.0
air	7.36	0.001205	0.901	0.899
5%Xe+95%CO ₂	12.78	0.005992	0.968	0.956
Ni/Ti alloy	25.36	6.5	0.908	0.978
stainless steel (302)	25.71	8.06	0.896	0.970
plaque (0% apatite)	5.01	1.07	0.979	0.997
plaque (10% apatite)	5.82	1.14	0.985	0.993
plaque (20% apatite)	6.65	1.23	0.987	0.992
plaque (30% apatite)	7.50	1.34	0.989	0.988

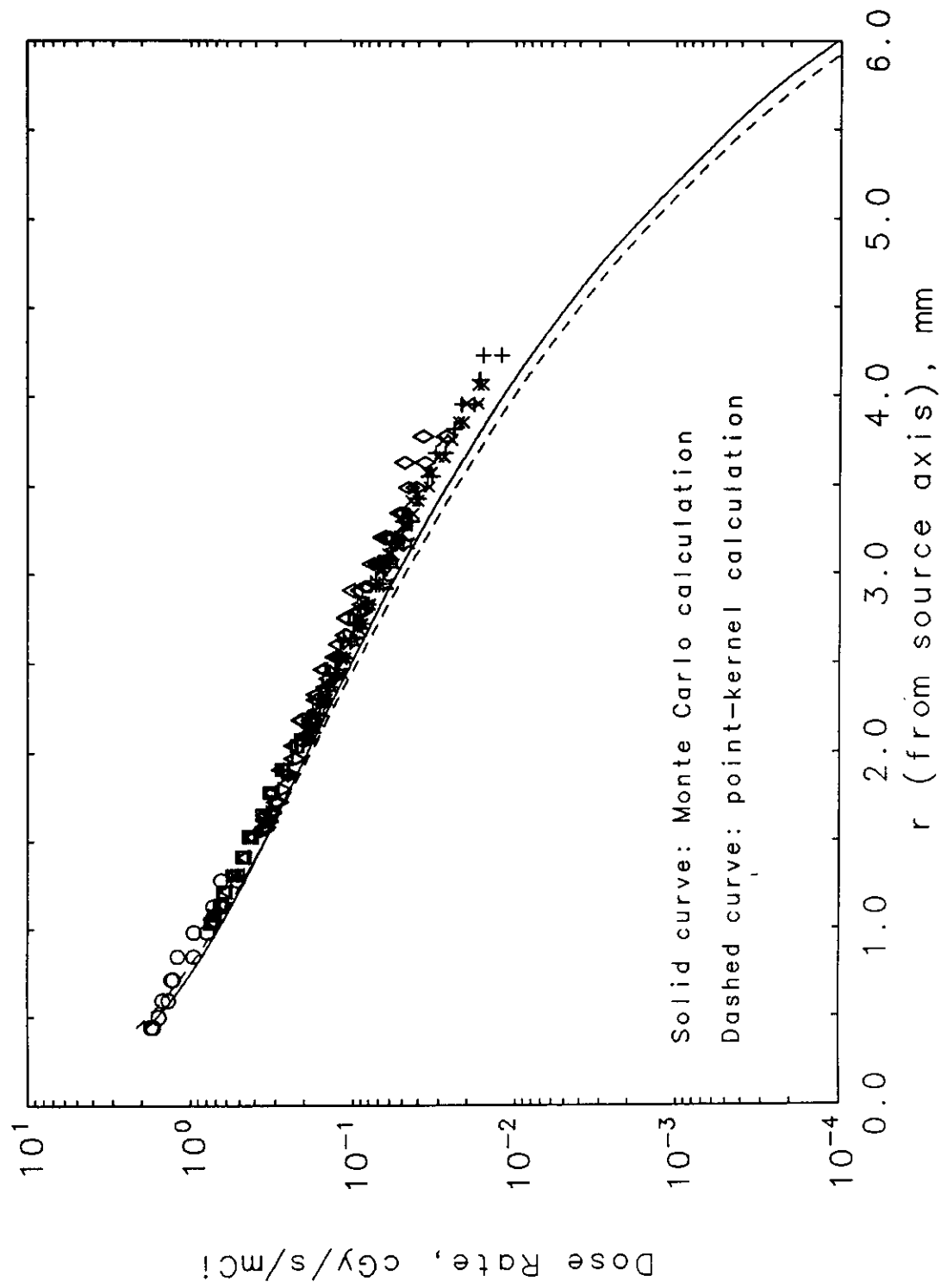


$^{90}\text{Sr}/^{90}\text{Y}$ Source in A-150 Plastic

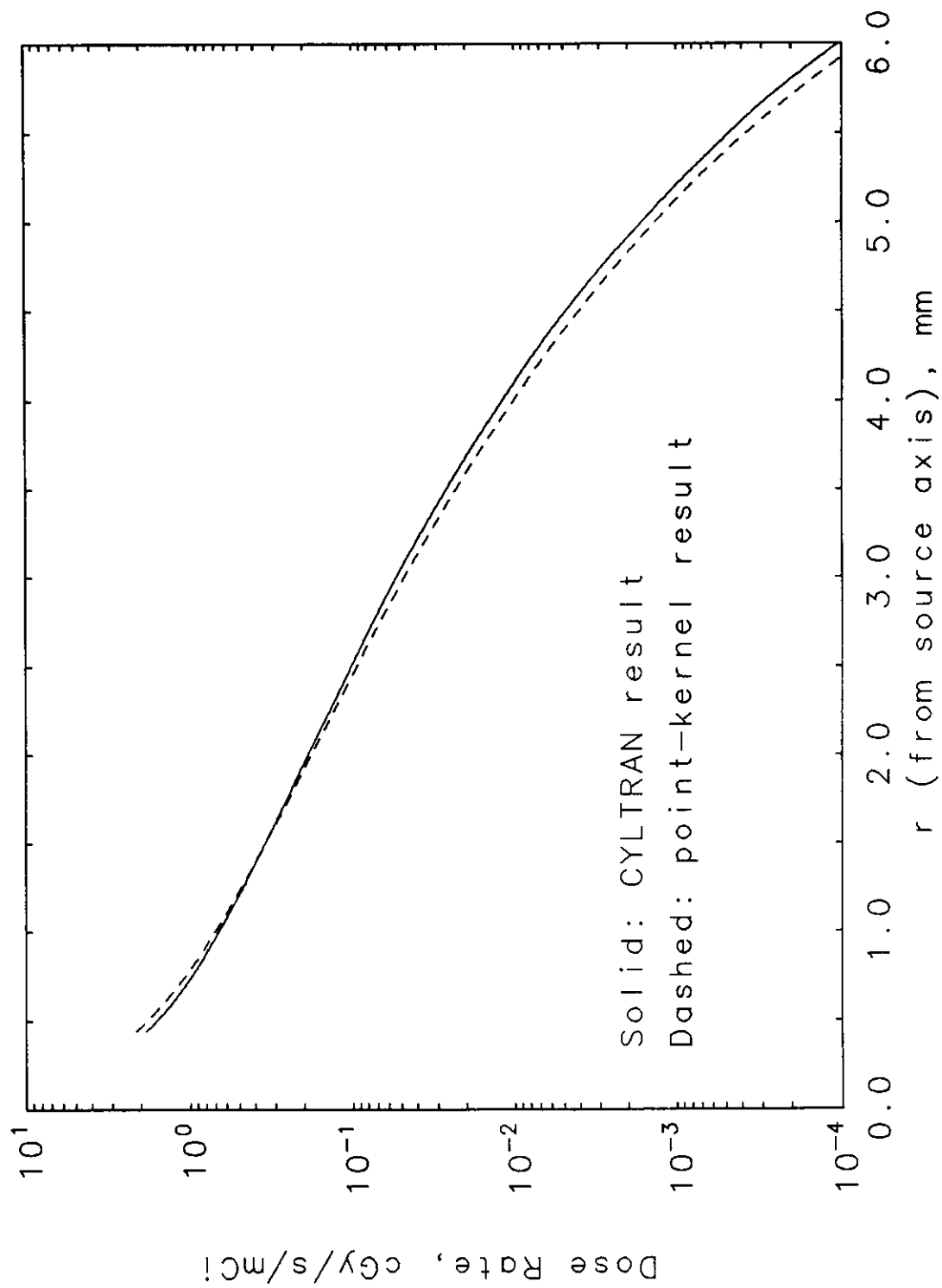




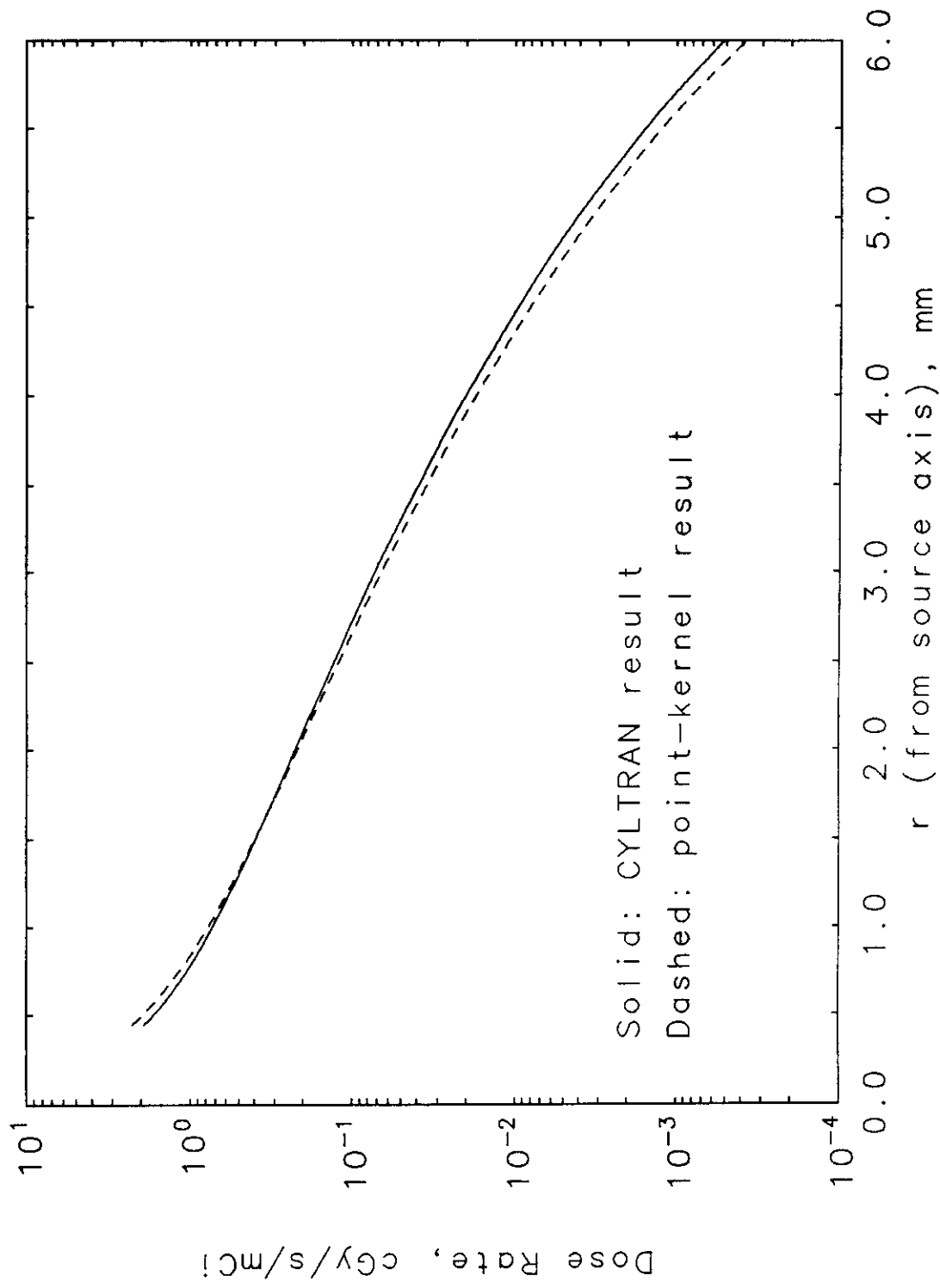
^{32}P Wire Source in A150 Plastic



^{32}P Wire Source in A150 Plastic



^{32}P Wire Source in Water



Point-Kernel Method: Photons

- Dose (or dose rate) from uniform volume source (per photon of energy E)

$$D(r) = \frac{E \mu_{\text{en}}(E)}{4\pi\rho V} \int_V T(E, |\vec{r}-\vec{r}'|) B_{\text{en}}(E, \mu(E) |\vec{r}-\vec{r}'|) \frac{e^{-\mu(E) |\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|^2} dV$$

where B_{en} is the energy-absorption build-up factor for water.

- Include non-water materials along each ray $r-r'$ by taking into account appropriate attenuation coefficient μ

$$\mu(E)r = \sum_j \mu_j(E) r_j$$

and using simple density scaling for $T(r)$, the transition-to-electronic-equilibrium function for water.

- Transition function $T(r)$ is assumed to be a universal function of $r/r_{\text{eff}}(E)$, where $r_{\text{eff}}(E)$ is the weighted average of the maximum csda ranges of the secondary electrons produced in the interaction.

