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**LASERS IN SURFACE SCIENCE**

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*Eigenstate-Resolved Measurements of  
Methane Dissociation on Ni(100)*

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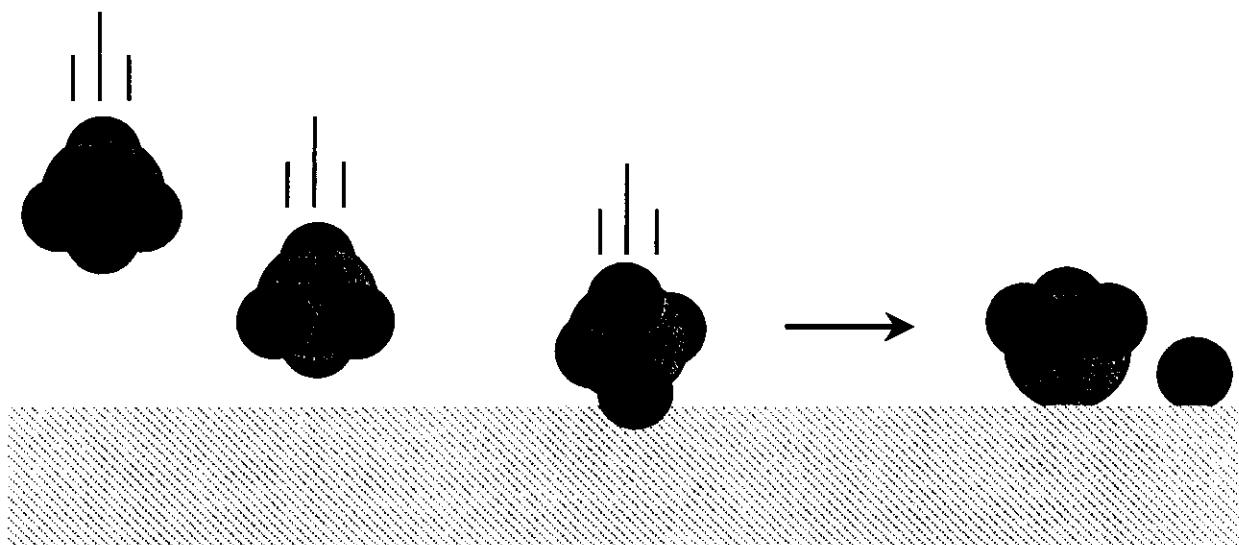
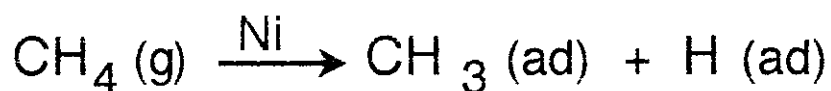


# Eigenstate-Resolved Measurements of Methane Dissociation on Ni(100)

- Overview
  - Motivation for State-Resolved Studies
  - Methane Surface Chemistry
- Experimental Approach
  - Instrumental Considerations
  - Doppler Mapping Detection
- Experimental Results
  - Efficacy of C-H stretch Excitation ( $v=1$ ,  $\nu_3$ )
  - Evidence for Mode Specificity ( $v=3$ ,  $\nu_4$ )
  - Rotational Effects
    - Efficacy of Rotational Energy
    - Reactivity of Coriolis-Coupled  $\nu_3$  Sublevels
    - Dynamical Effects
- Summary

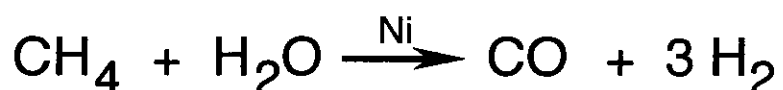
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## Why study methane dissociation on nickel?



- Prototype for C-H bond activation
- Benchmark for theoretical calculations
- Permits comparison of distinct vibrational modes
- Methane conversion is important industrial process

### Steam Reforming Reaction



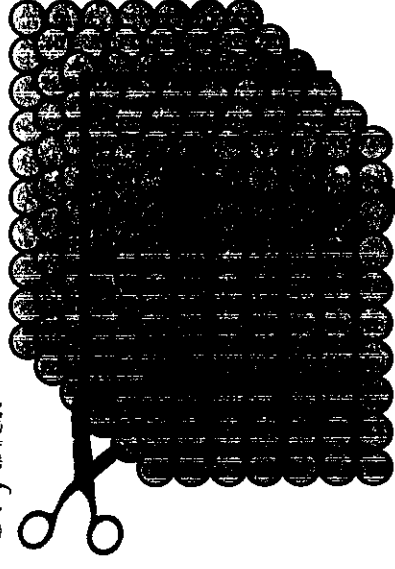
- Chief industrial pathway for  $\text{H}_2$  production
- $\text{CH}_4$  dissociation step is rate limiting

# Measuring State-Selected ethane Dissociative reabsorption Probabilities

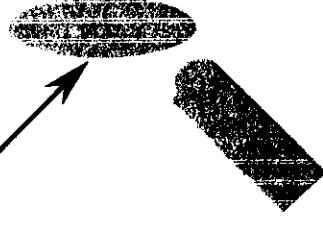
Supersonic Beam  
 $\Delta E/E \approx 10\%$



Ni Single Crystal



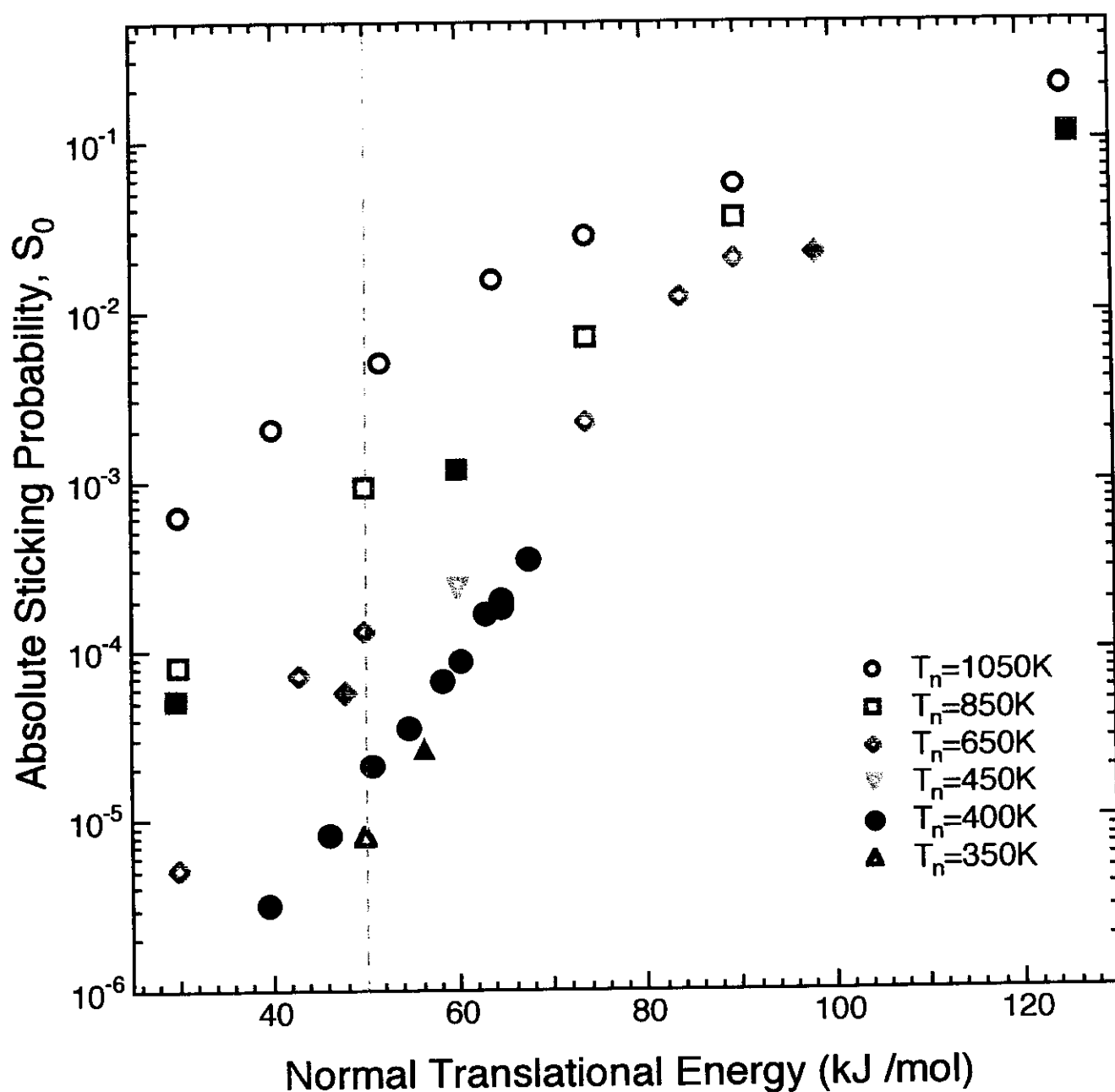
Ni(100) face



Auger Electron  
Spectroscopy

$$\frac{\left. \begin{array}{l} \text{Measure C surface coverage} \\ \text{Auger Electron Spectroscopy} \end{array} \right\} \text{ \# of Molecules Dissociated / cm}^2}{\left. \begin{array}{l} \text{Molecular Beam Flux} \times \text{time} \end{array} \right\} \text{ \# of Incident Molecules / cm}^2} = S_0 = \text{Reaction Probability}$$

## Some More Experimental Data for CH<sub>4</sub>/Ni(100)



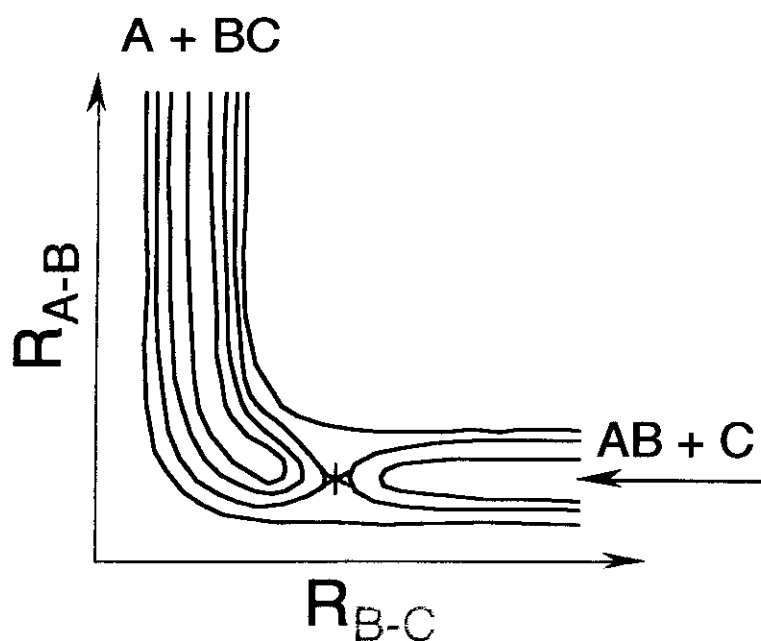
- $S_0$  increases exponentially with  $E_{trans}$
- Nozzle Temperature ( $T_n$ ) affects  $S_0$  at fixed  $E_{trans}$
- $T_n$  effect due to enhanced reactivity of thermally populated vibrations in the beam

Open Symbols: Holmblad et al., J. Chem. Phys., **102**, 8255, (1995). Solid Symbols: This Work

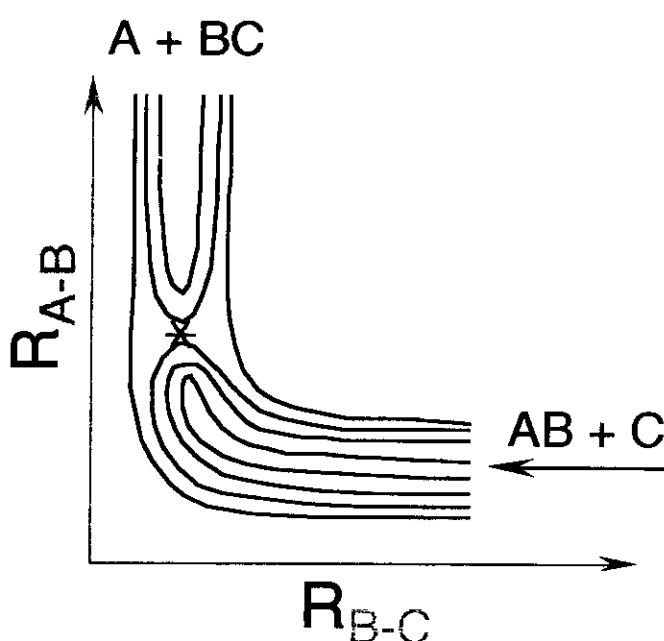
## 2-Dimensional Potential Energy Surface



(Collinear geometry)



*EARLY BARRIER*



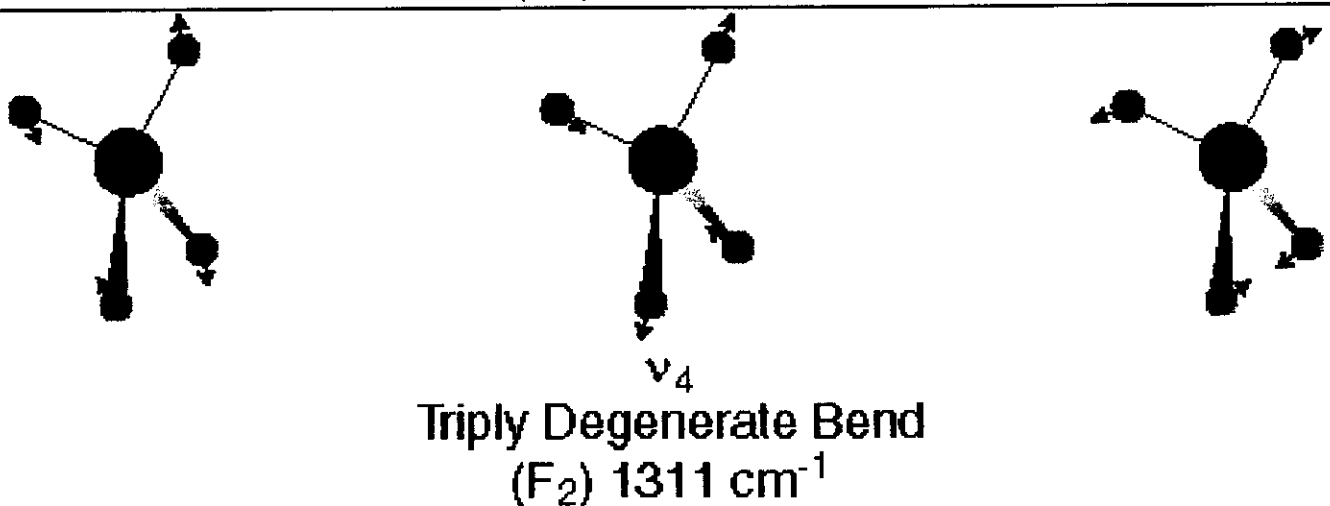
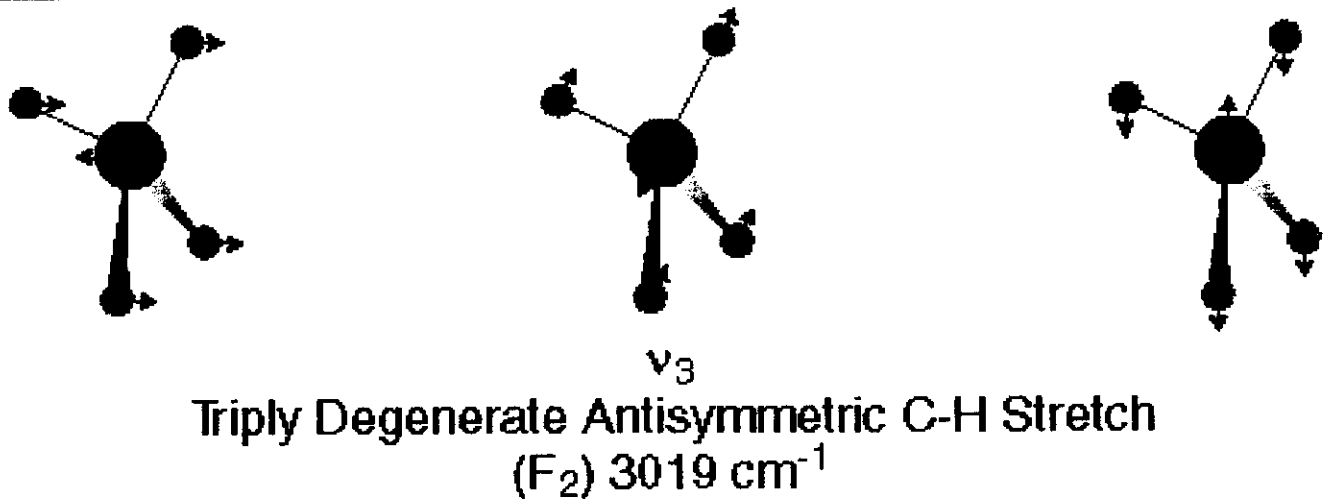
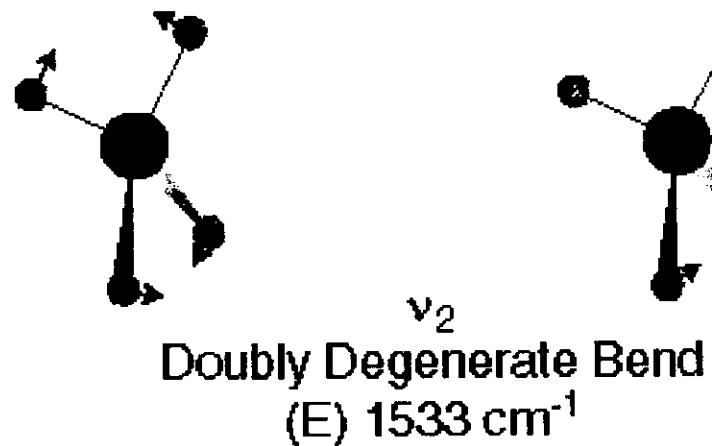
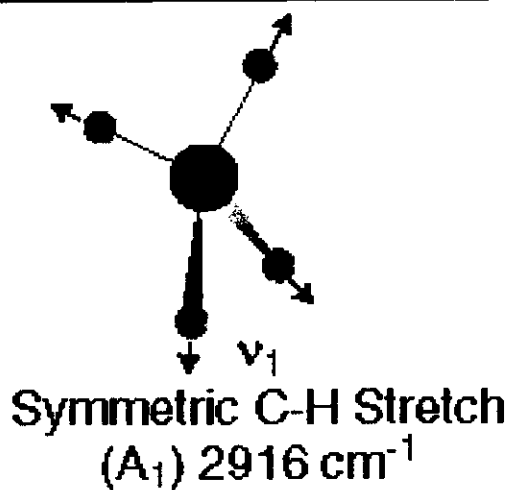
*LATE BARRIER*

- Polanyi's Rules
  - Translational energy for early barrier
  - Vibrational energy for late barrier
- Real PES has many more degrees of freedom
  - Rotations
  - Vibrations ( $3N-6$ )
  - Translations ...

How do we probe high dimensionality  
Gas-Surface PES?

# Methane Vibrations

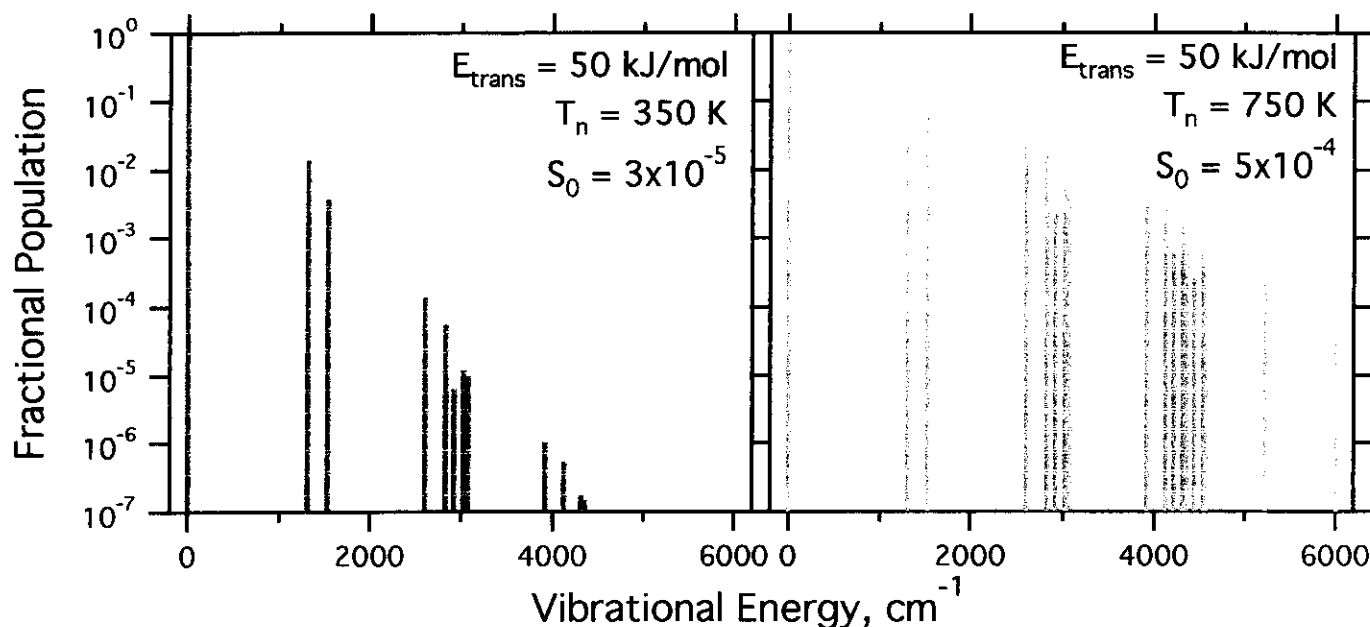
- Methane has four vibrational normal modes



- These fundamentals, plus their overtones and combination are present in a thermal distribution of vibrational states

# The Problem

- “Hot” vibrational distribution is more reactive, but...



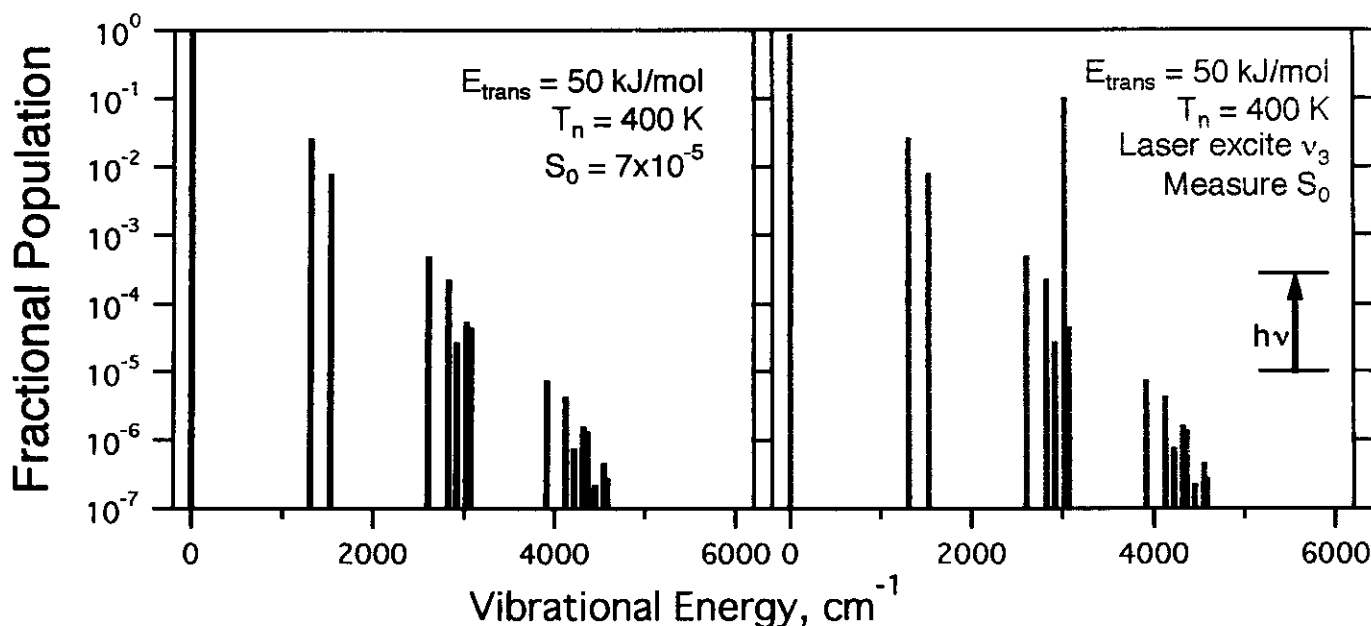
- Many vibrations populated in thermal distribution
  - Different modes
  - Different levels of excitation
- Thermal averaging obscures identity of most reactive mode
- Absolute number of highly excited molecules remains small

Fractional populations calculated.

Dissociation probabilities from P. M. Holmblad, J. Wambach, and I. Chorkendorff, *J. Chem. Phys.*, **102**, 8255 (1995) ( $T_n = 350$  and  $750\text{ K}$ ) and this work ( $T_n = 400\text{ K}$ ).

# A Solution

- Infrared laser excites single rovibrational eigenstates (e.g.  $\nu_3$ )



Laser Off: Average  $S_0$  has contributions from  $v=0$  and  $v>0$

$$S_0^{\text{LaserOff}} = f_{v=0} S_0^{v=0} + (1 - f_{v=0}) S_0^{v>0}$$

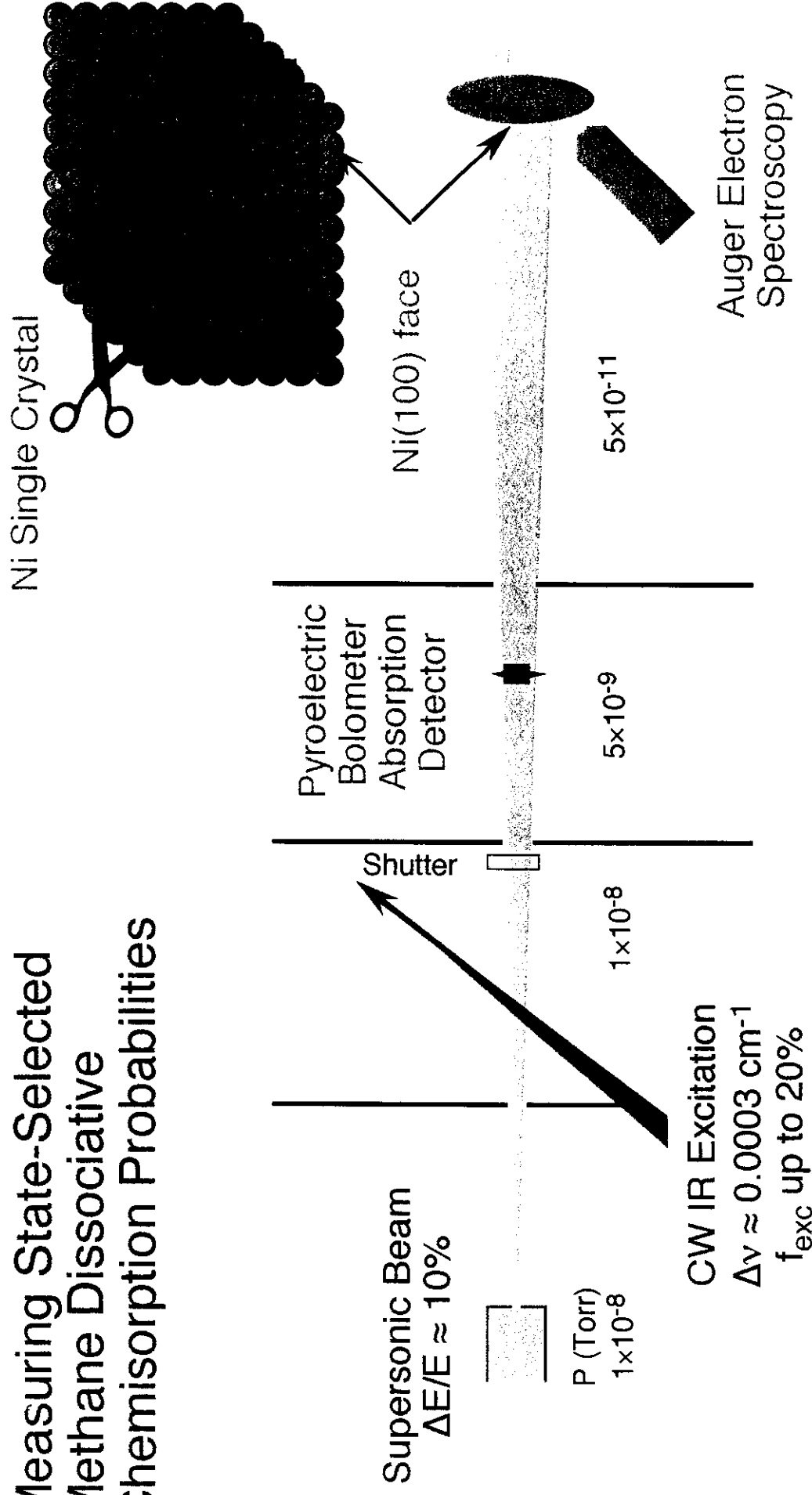
Laser on: Average  $S_0$  has contributions from  $v=0$ ,  $v>0$ , and  $\nu_3$

$$S_0^{\text{LaserOn}} = (f_{v=0} - f_{\text{exc}}) S_0^{v=0} + (1 - f_{v=0}) S_0^{v>0} + f_{\text{exc}} S_0^{\nu_3}$$

$$S_0^{\nu_3} = \frac{S_0^{\text{Laser On}} - S_0^{\text{Laser Off}}}{f_{\text{exc}}} + S_0^{v=0}$$

- Measure  $S_0^{\text{LaserOff}}$  and  $S_0^{\text{LaserOn}}$
- Optical Saturation Measurement yields  $f_{\text{exc}}$  (12 - 16%)
- $S_0^{v=0} \ll S_0^{\nu_3}$
- Extracted  $S_0$  independent of vib. state distribution in beam

# Measuring State-Selected Methane Dissociative Chemisorption Probabilities



Measure C surface coverage  
Auger Electron Spectroscopy

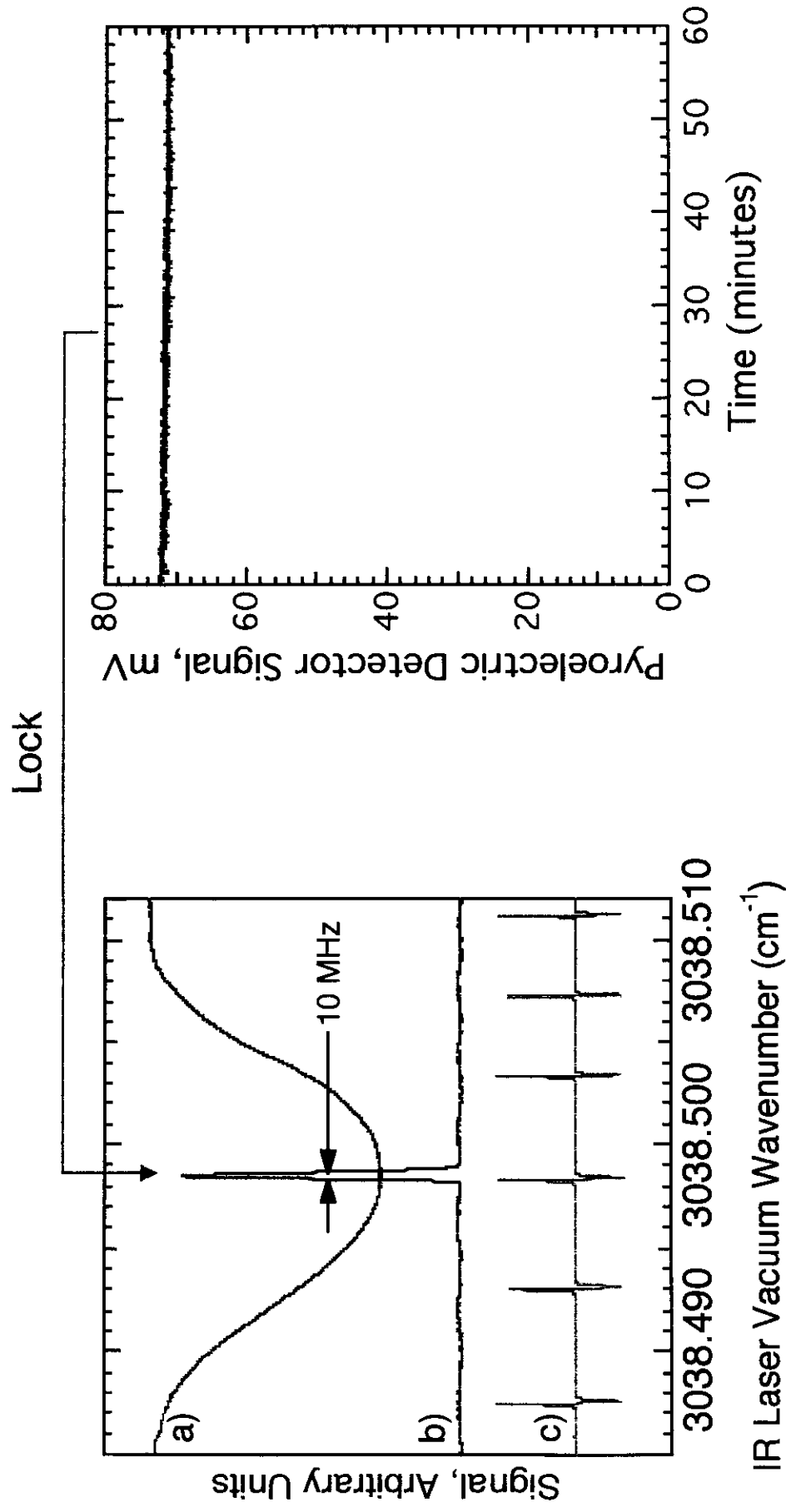
} # of Molecules Dissociated /  $\text{cm}^2$

Molecular Beam Flux  $\times$  time  
 $\times$  IR Excitation Fraction

} # of Incident Molecules /  $\text{cm}^2$

$$= S_0 \quad \text{Reaction Probability}$$

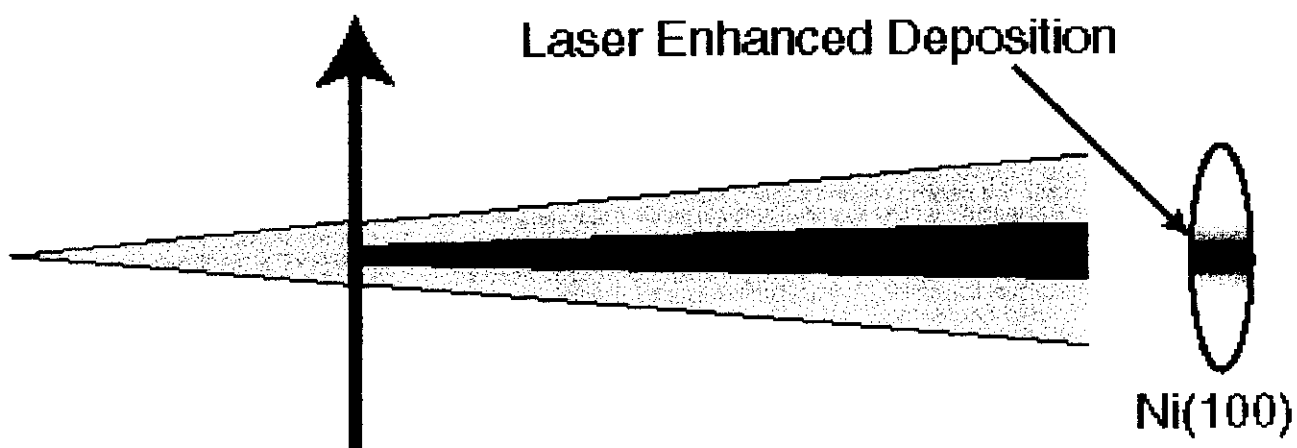
# Controlling Reagent Internal Energy



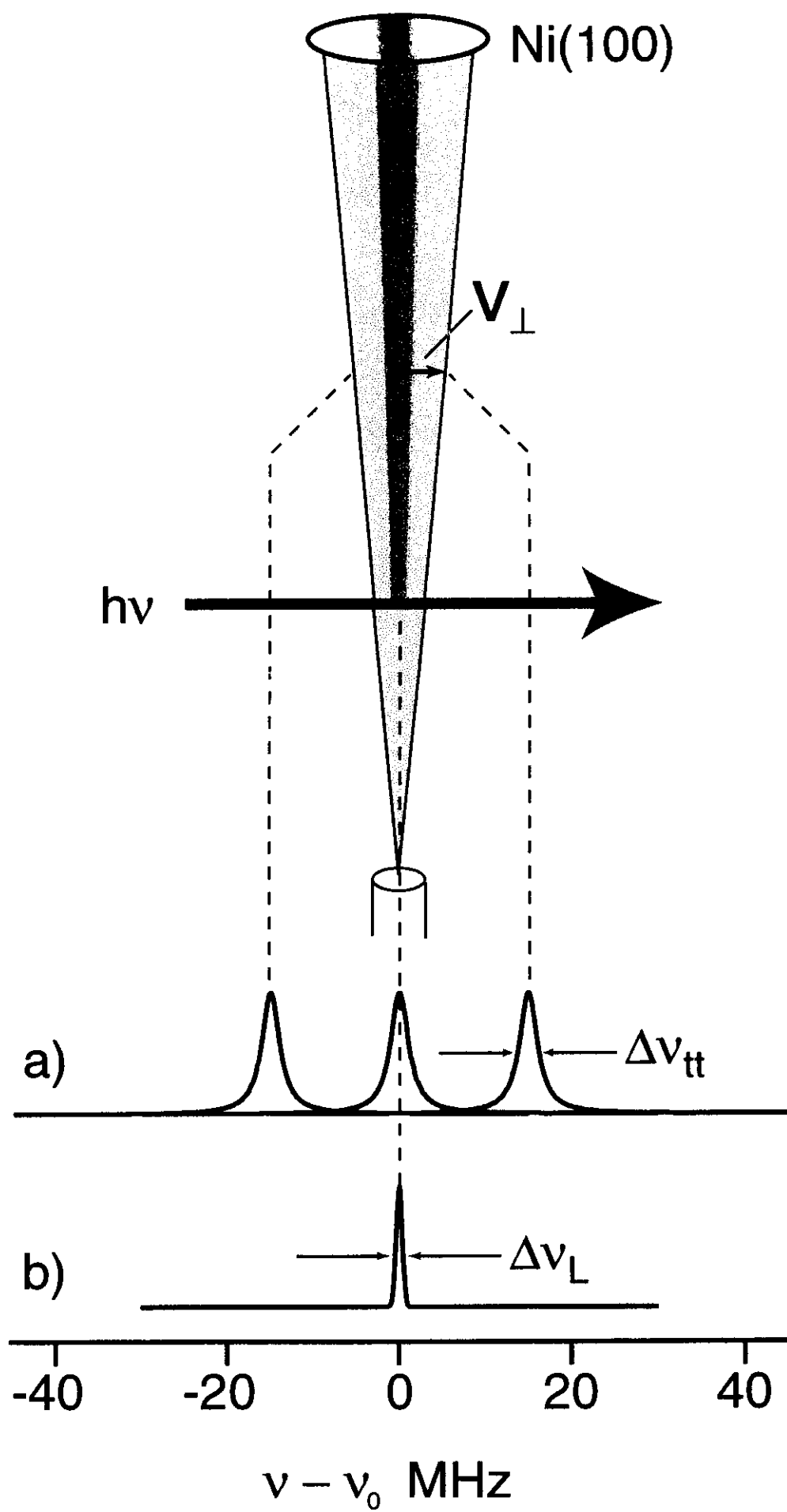
- a)  $\text{CH}_4$  absorption in room temperature cell.
- b)  $\text{CH}_4$  absorption in molecular beam.
- c) Derivative signal of light transmitted through 150 MHz FSR Fabry-Perot etalon.
- Transition shown excites  $J=2$  of  $\nu_3$ , the antisymmetric C-H stretch of  $^{12}\text{CH}_4$ , via  $R(1)$ .
- Lock laser frequency to zero crossing of Fabry-Perot reference cavity.

## Doppler Mapping Detection

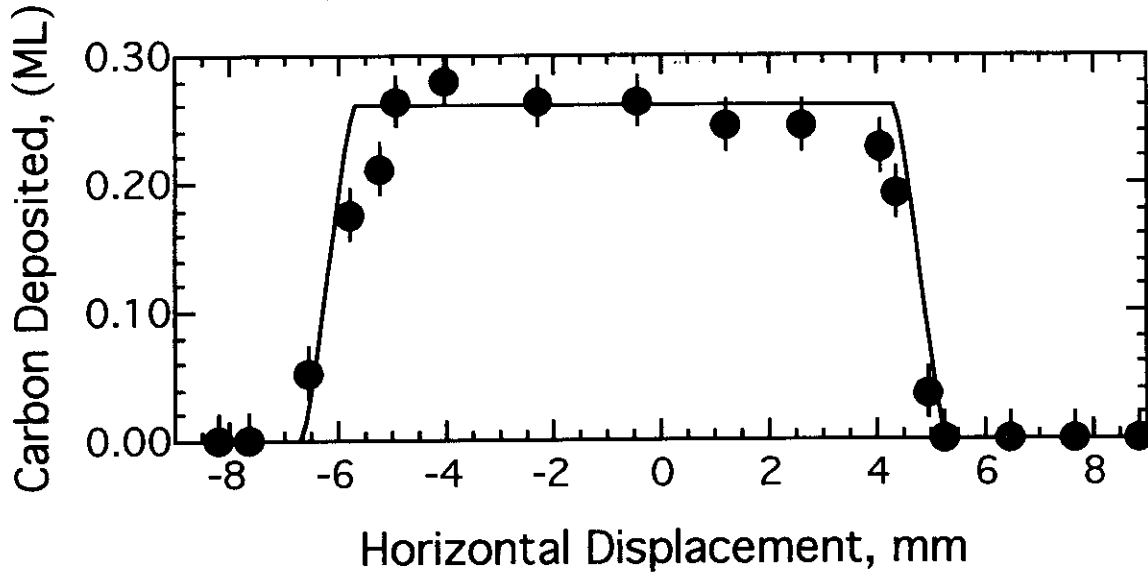
- CW single mode infrared laser linewidth  $< 1\text{MHz}$ .
- Transit-time broadening dominates lifetime broadening:  
 $\Delta\nu_L \approx 1\text{MHz}$
- Doppler broadening dominates absorption profile due to spread of transverse velocities in diverging molecular beam.
- Narrow bandwidth laser excites subset of the inhomogeneous Doppler profile in the molecular beam.



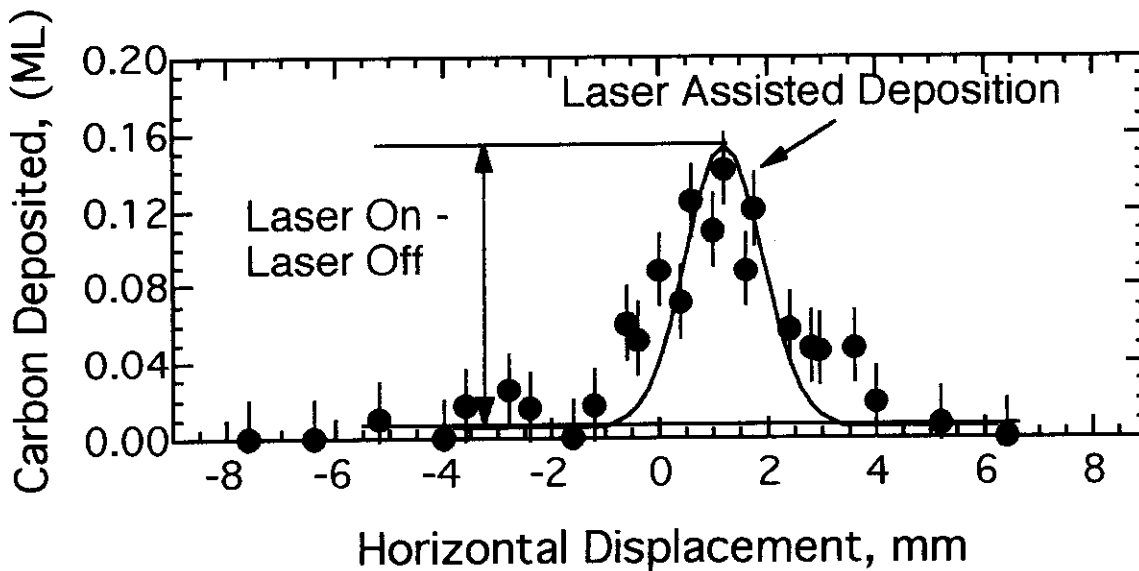
Laser-enhanced deposition is localized on the surface



## Doppler Mapping Detection

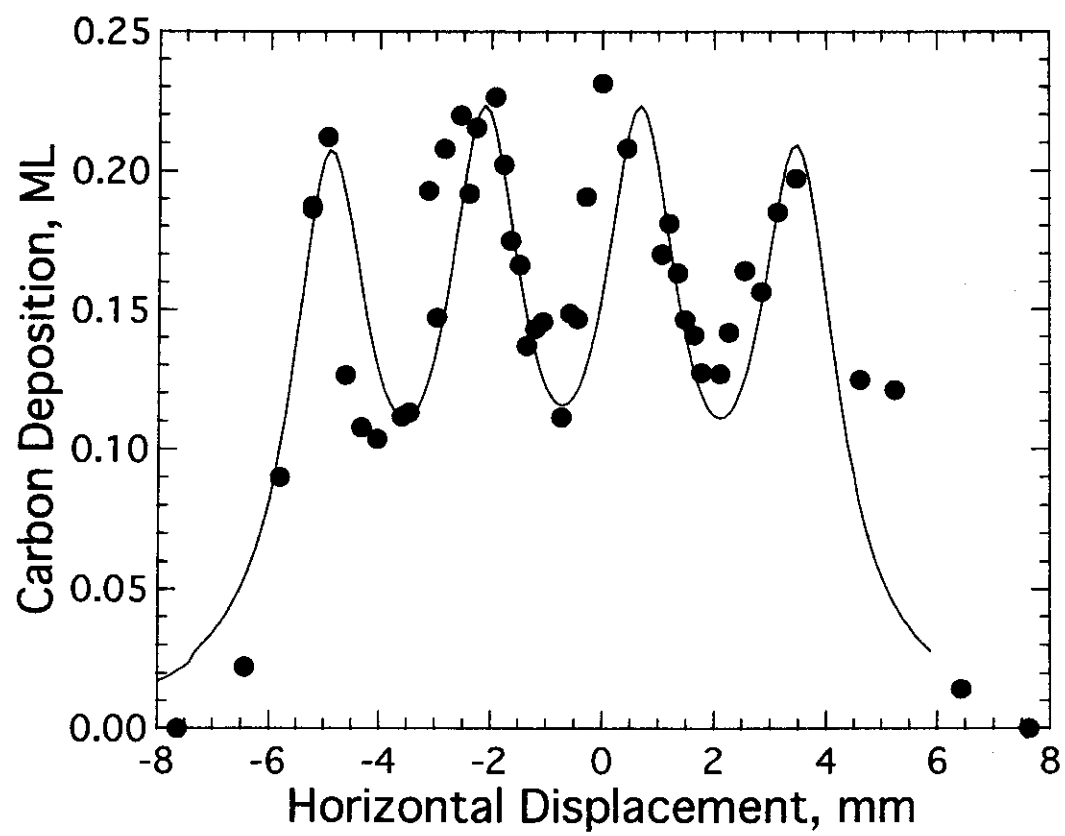
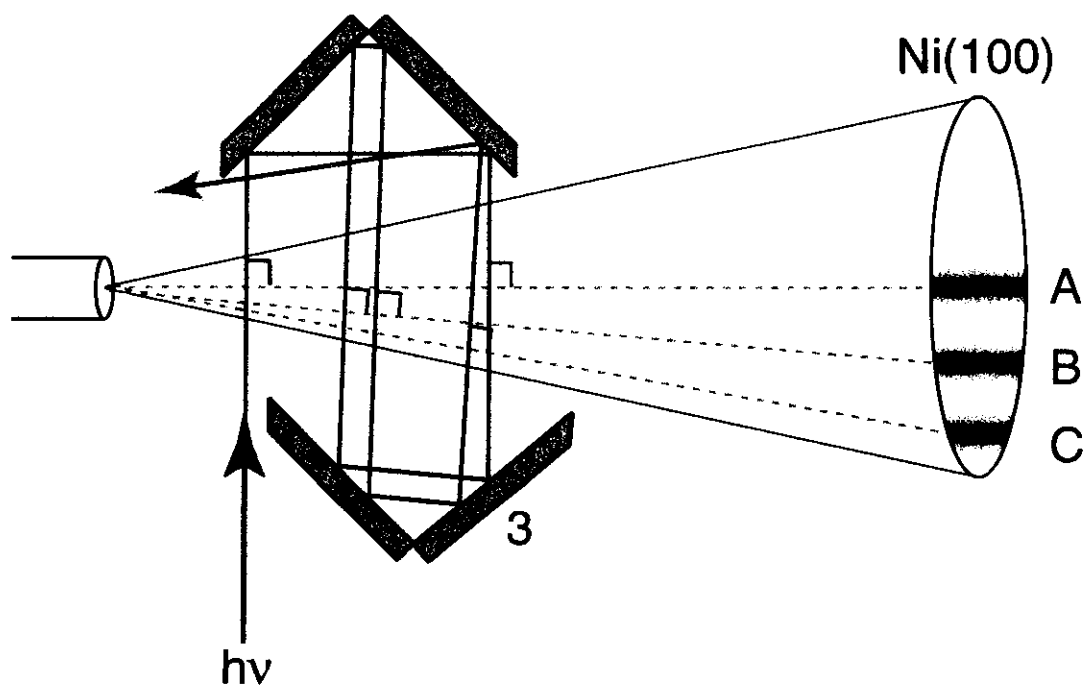


- Dose with Laser OFF.  $t = 1200$  s
- Carbon uniformly deposited across 10 mm diameter crystal

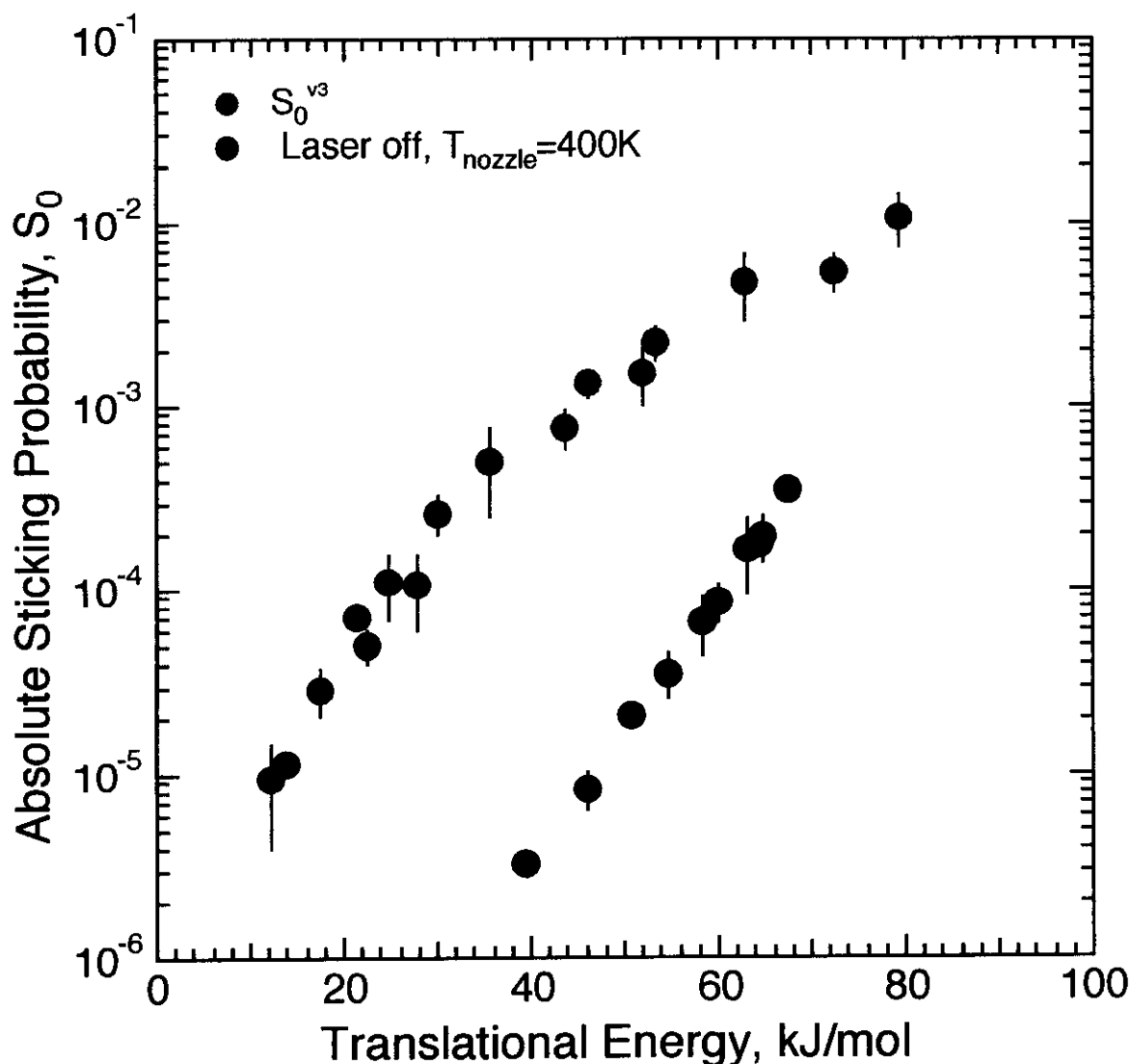


- Dose with Laser ON.  $t = 50$  sec
- Laser-excited molecules localized near crystal center
- Nonzero baseline due to "Laser Off" reactivity
- Simultaneous measurement of signal and background enhances sensitivity.

# Effect of Misaligned Mirror



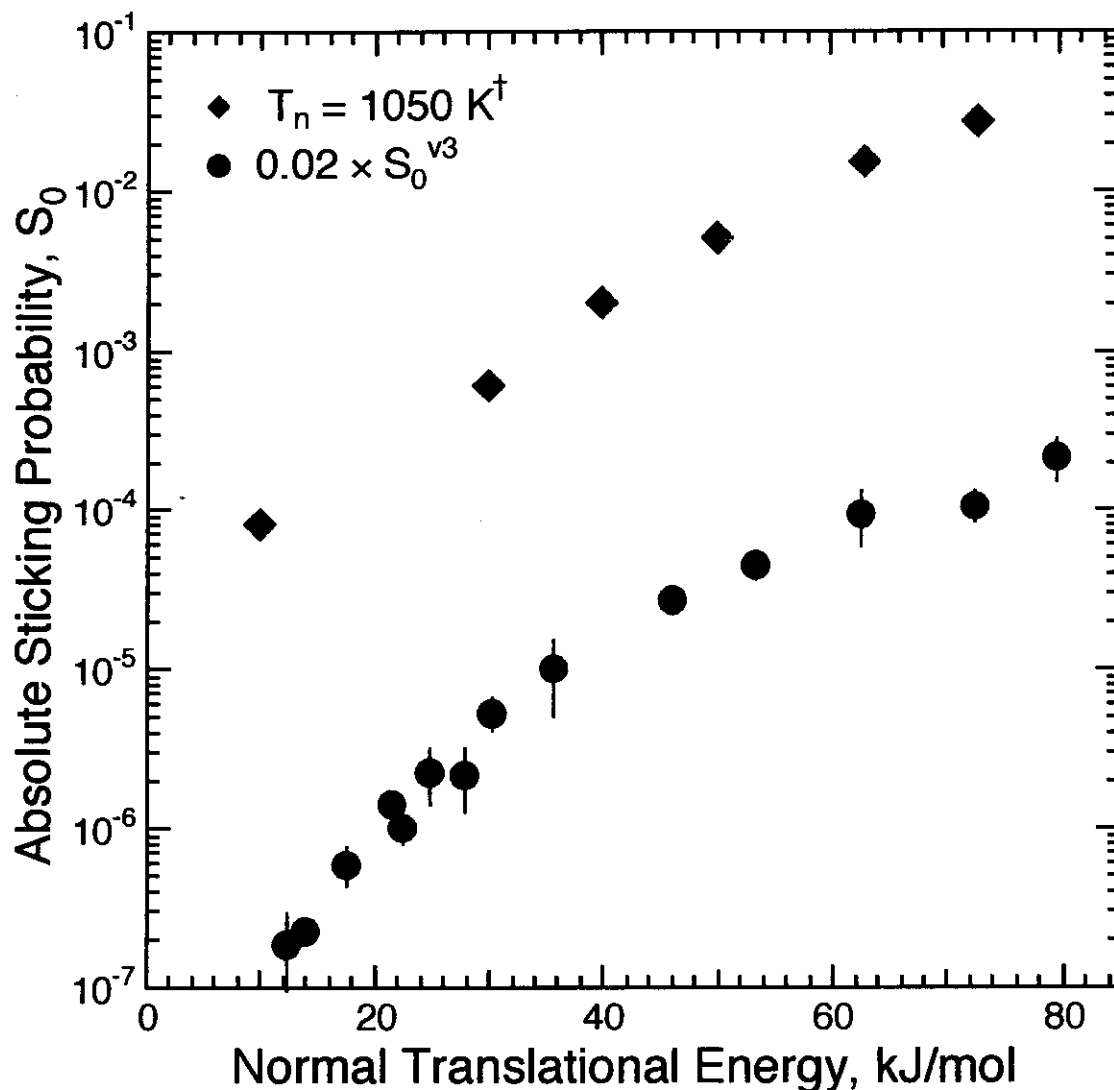
# State-Resolved Sticking Probabilities for CH<sub>4</sub> ( $\nu_3$ , J=2) on Ni(100)



- Excitation of  $\nu_3$  J=2 deposits 36 kJ/mole of internal energy and increases  $S_0$  by 500-fold or more at low  $E_{\text{trans}}$
- $S_0$  increases monotonically from  $E_{\text{trans}} = 13\text{-}80$  kJ/mole
- Despite having  $E_{\text{trans}} + E_{\text{vib}} = 116$  kJ/mole,  $S_0$  @  $E_{\text{trans}} = 80\text{kJ/mol}$  is 0.01

# The Role of $\nu_3$ in Thermal Processes

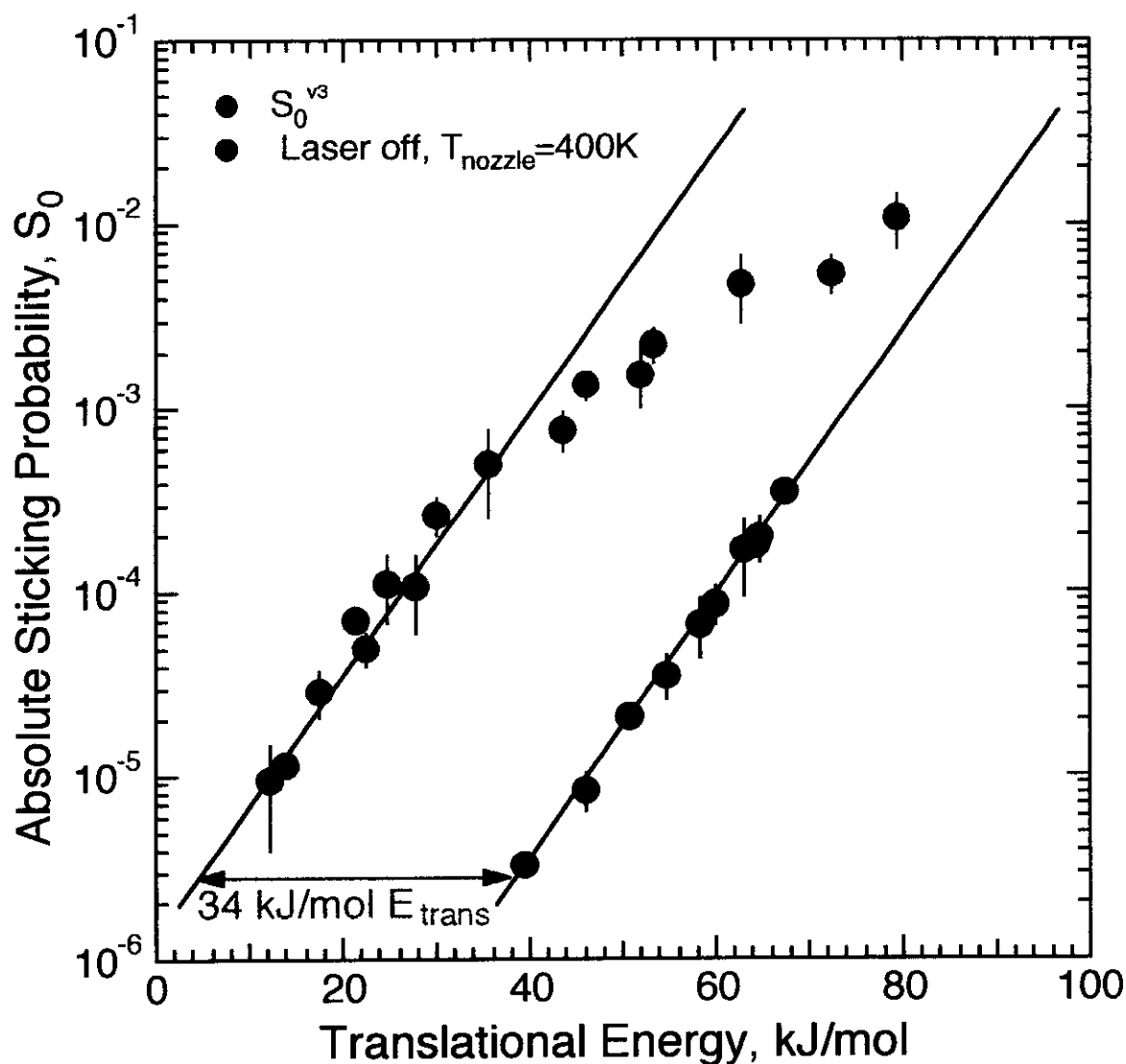
- State-resolved data reveal the contribution of molecules in the  $\nu_3$  eigenstate to thermally averaged reaction probabilities.
- Holmblad et al.<sup>†</sup> report reaction probabilities averaged over a thermal vibrational state distribution at 1050K.



- At 1050K, the fractional population of  $\nu_3$ ,  $v=1$  is  $f_{\nu_3} = 0.02$
- $f_{\nu_3} \times S_0^{\nu_3} \approx 0.02 S_0^{T_{\text{vib}}=1050\text{K}}$
- $\nu_3$ ,  $v=1$  only accounts for  $\approx 2\%$  of  $S_0$  at  $T_{\text{nozzle}} = 1050\text{K}$ .

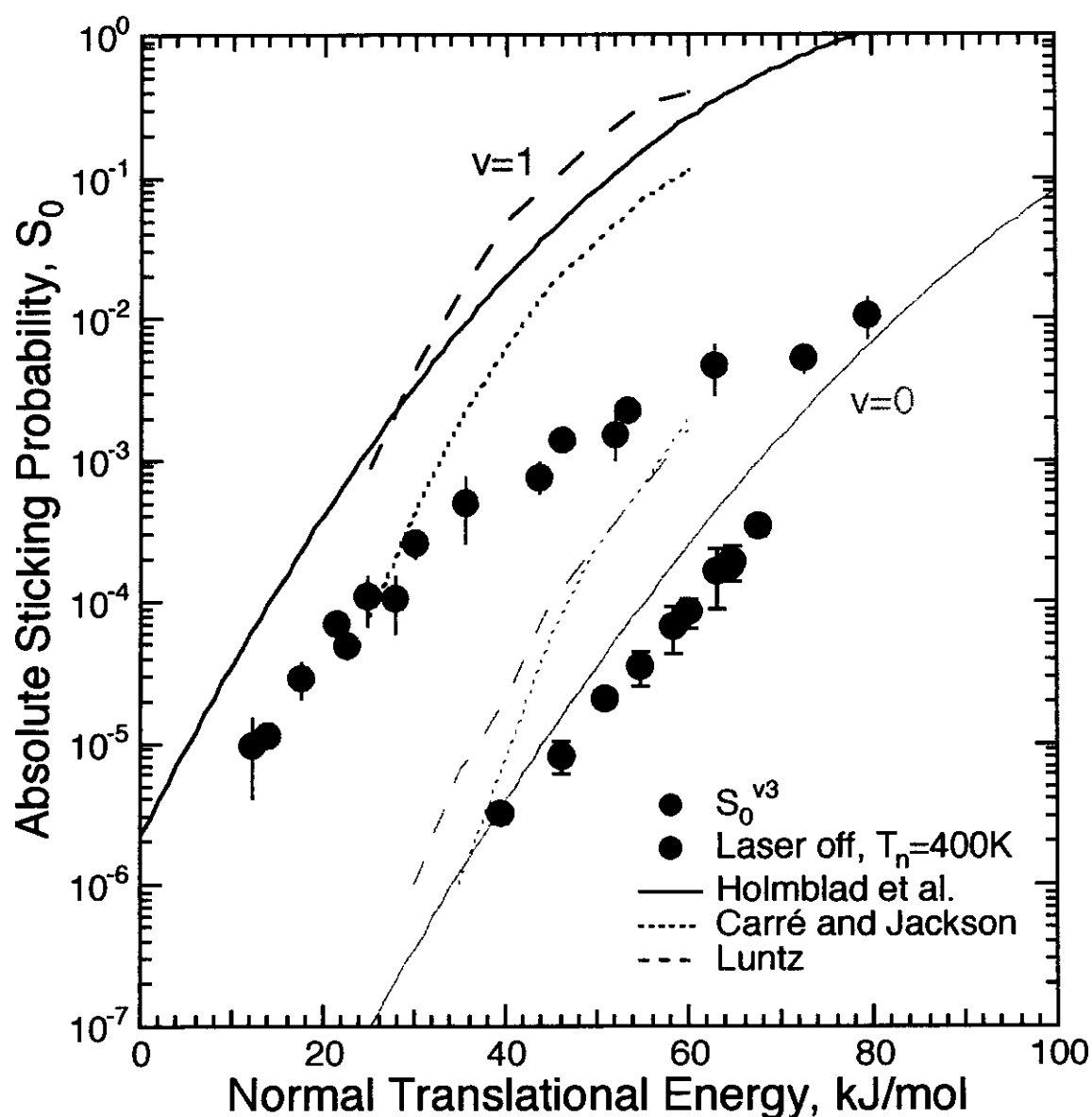
<sup>†</sup> Data from P. M. Holmblad, J. Wambach, and I. Chorkendorff, *J. Chem. Phys.*, **102**, 8255 (1995).

# Testing a Statistical Model of Reactivity



- Excitation of  $\nu_3$  deposits 36 kJ/mol of internal energy and increases  $S_0$  by a factor of  $\approx 200$ .
- At low  $S_0$ ,  $\nu_3$  excitation shifts the  $E_{\text{trans}}$  threshold  $\approx 34$  kJ/mol
- These data do not rule out a statistical model of methane dissociation on Ni(100).

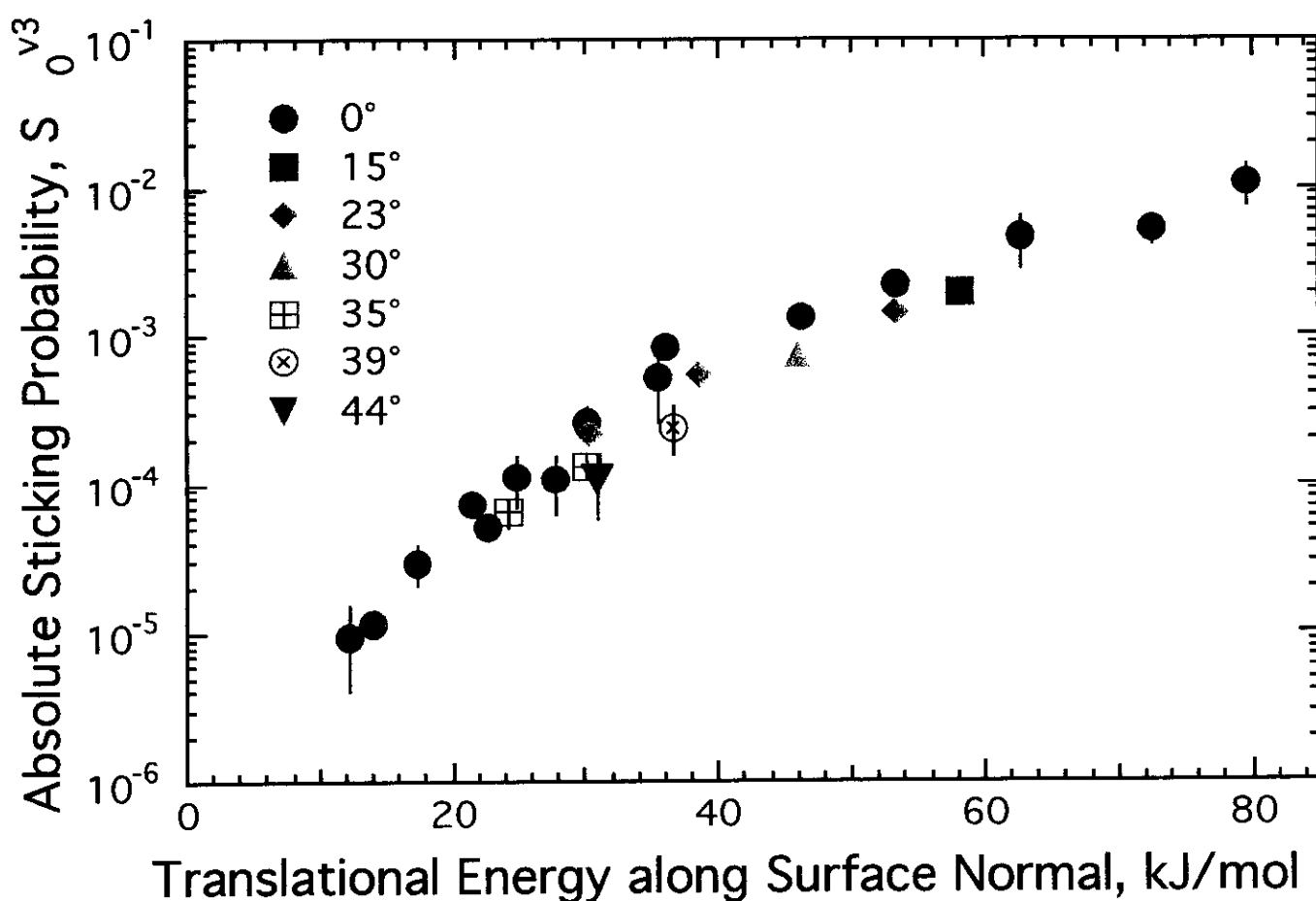
## Comparison with Current Theories & Models



- Carré & Jackson: Quantum Dynamics on ab initio PES
- Luntz: Quantum Calculation on Empirical PES.
- Holmblad et al.: Empirical model of data assuming only active vibration is C-H stretch.
- Models and calculations overestimate  $S_0$  at high  $E_{trans}$

# Normal Energy Scaling

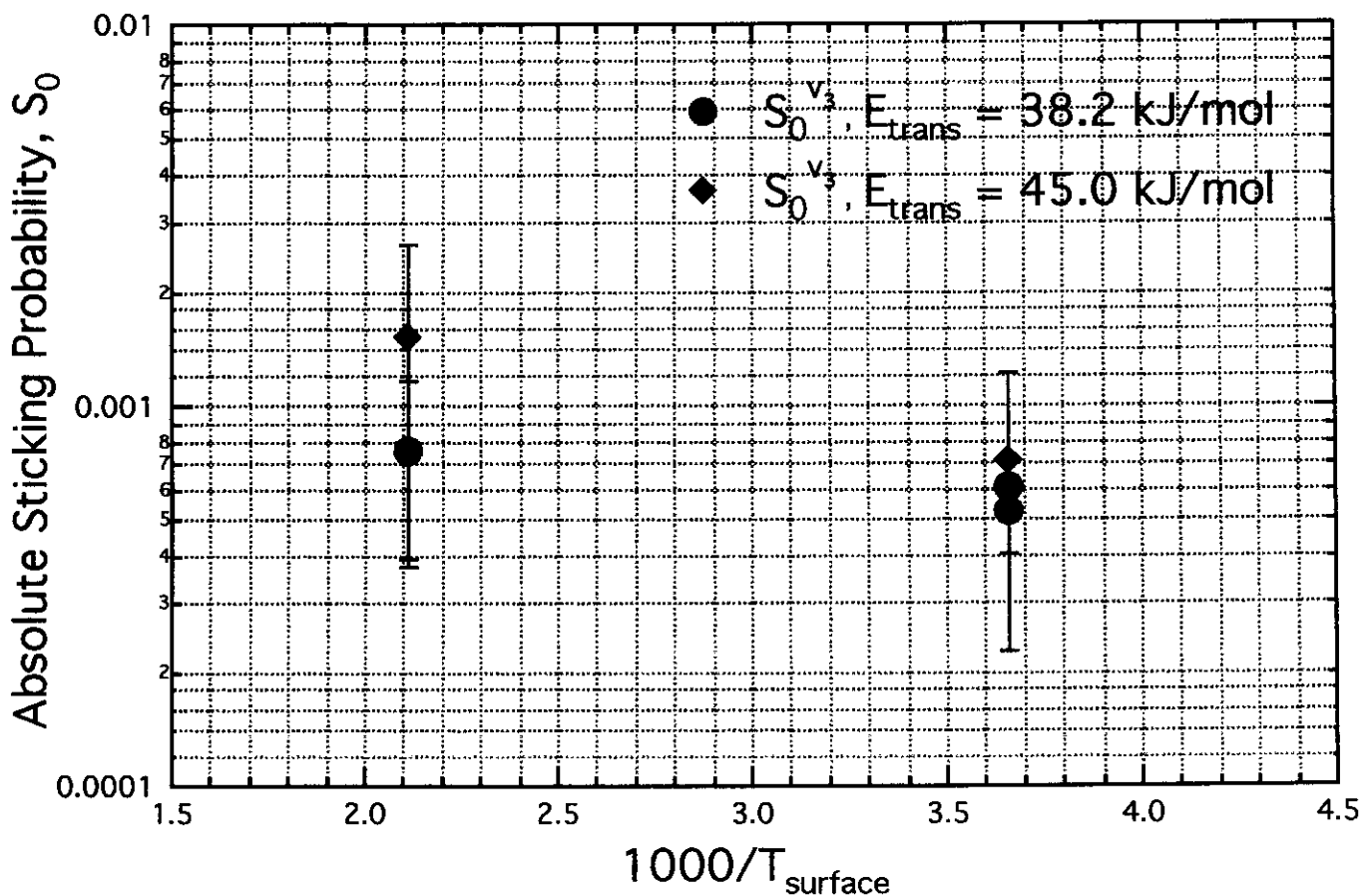
- Methane dissociation on many surfaces, including Ni(100), exhibits normal energy scaling in which only kinetic energy directed along the surface normal promotes reactivity.
- Adherence to normal energy scaling implies that translational motion along the surface normal is coupled to the reaction coordinate while translation parallel to the surface is not.
- We have measured  $S_0^{v_3}$  for several translational energies and incident angles. Values are plotted vs. the normal kinetic energy,  $E_{\text{trans}} (\cos^2\theta)$ .



- The data show that  $\text{CH}_4$  ( $v_3$ ) nearly follows normal energy scaling over the kinetic energy range investigated.
- Excitation of  $v_3$  does not eliminate the need for translational energy directed along the surface normal.

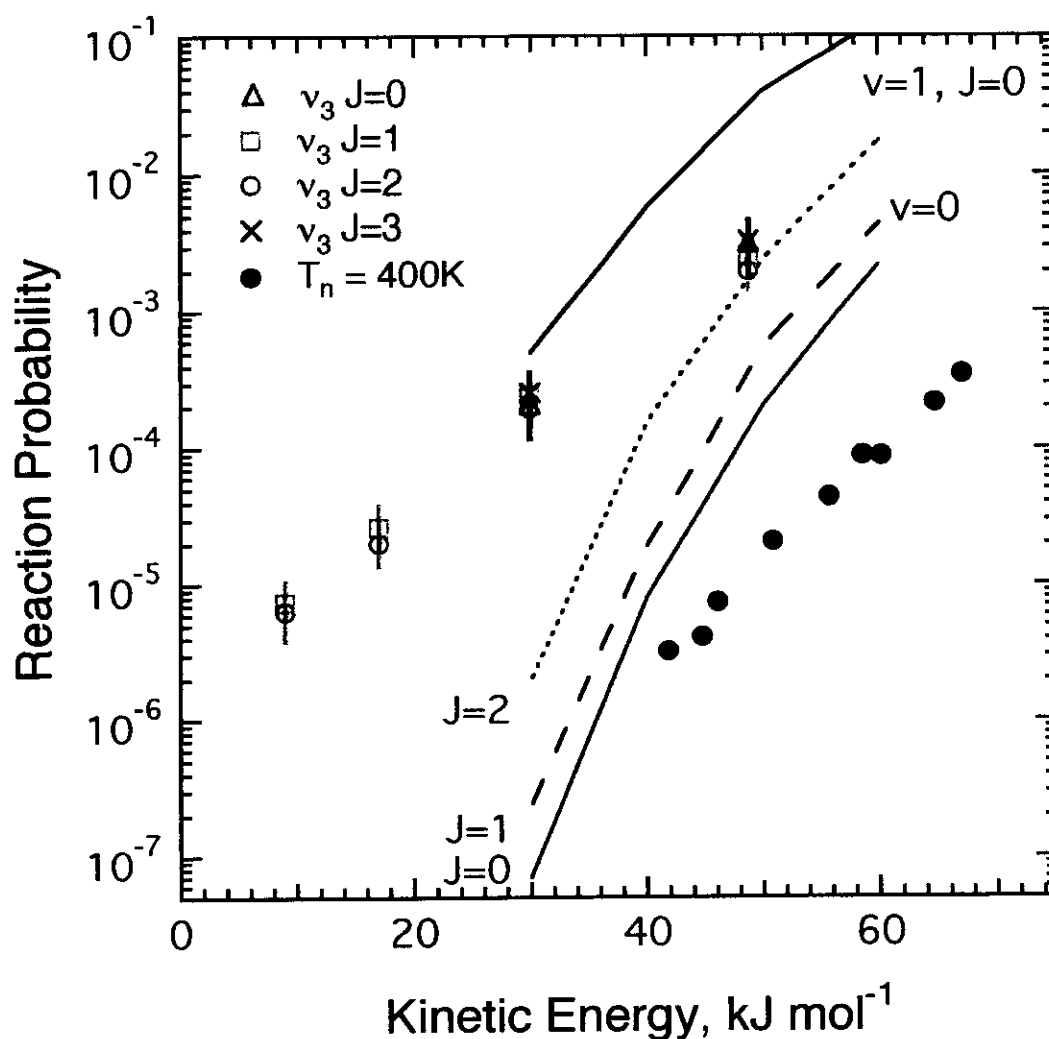
## Surface Temperature Dependence

- Methane dissociation on some transition metal surfaces is strongly influenced by surface temperature.
- Such dependence suggests that degrees of freedom associated with the surface are coupled to the reaction coordinate.
- We have measured  $S_0$  for  $\text{CH}_4$  in  $\nu_3$ ,  $v=1$ ,  $J=2$  at different surface temperatures and incident translational energies.



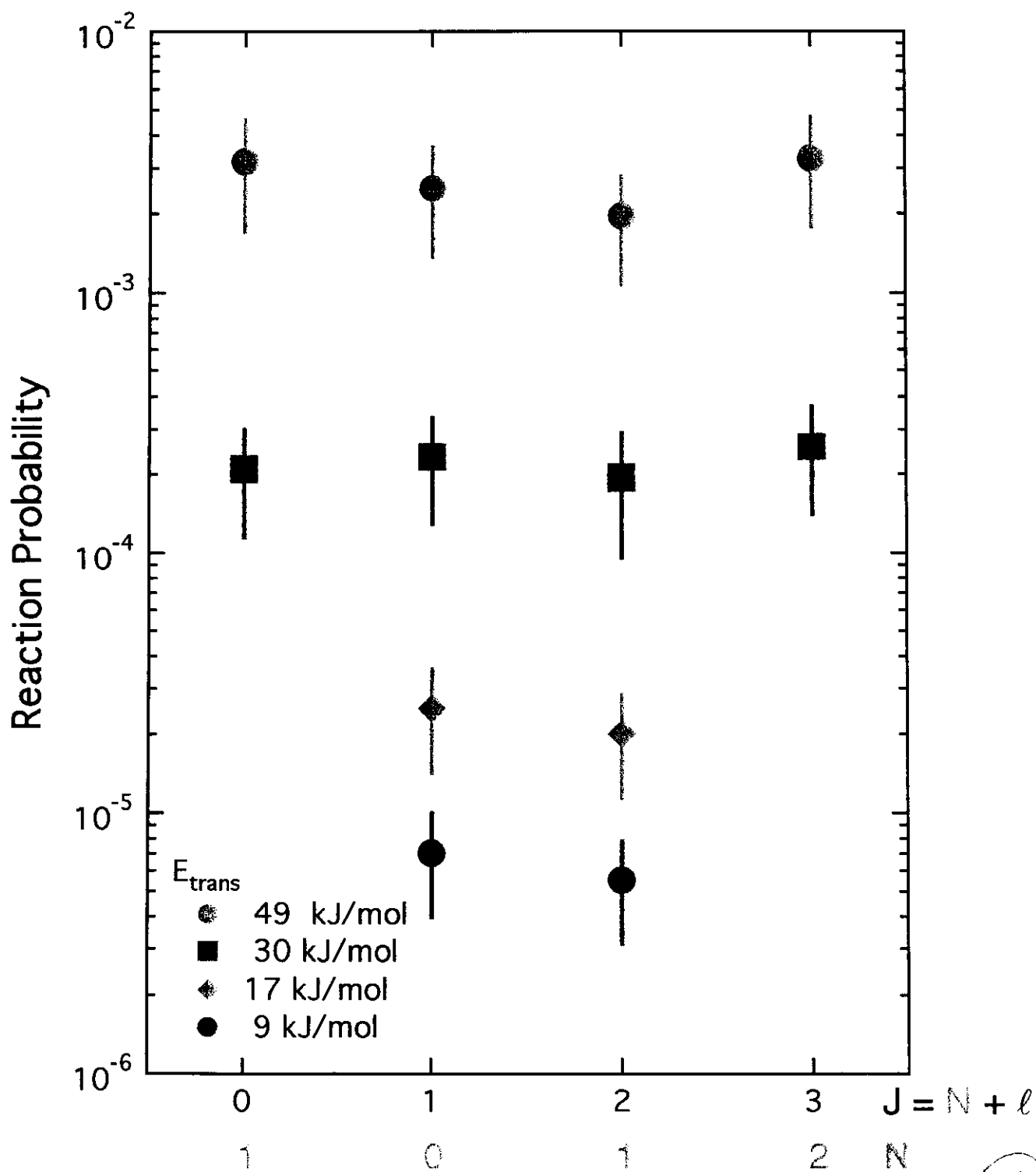
- The data show that  $S_0^{\nu_3}$  depends only weakly on  $T_{\text{surface}}$ .

## Rotational-State Resolved Sticking Probabilities for CH<sub>4</sub>, $\nu_3$ ( $v=1$ )

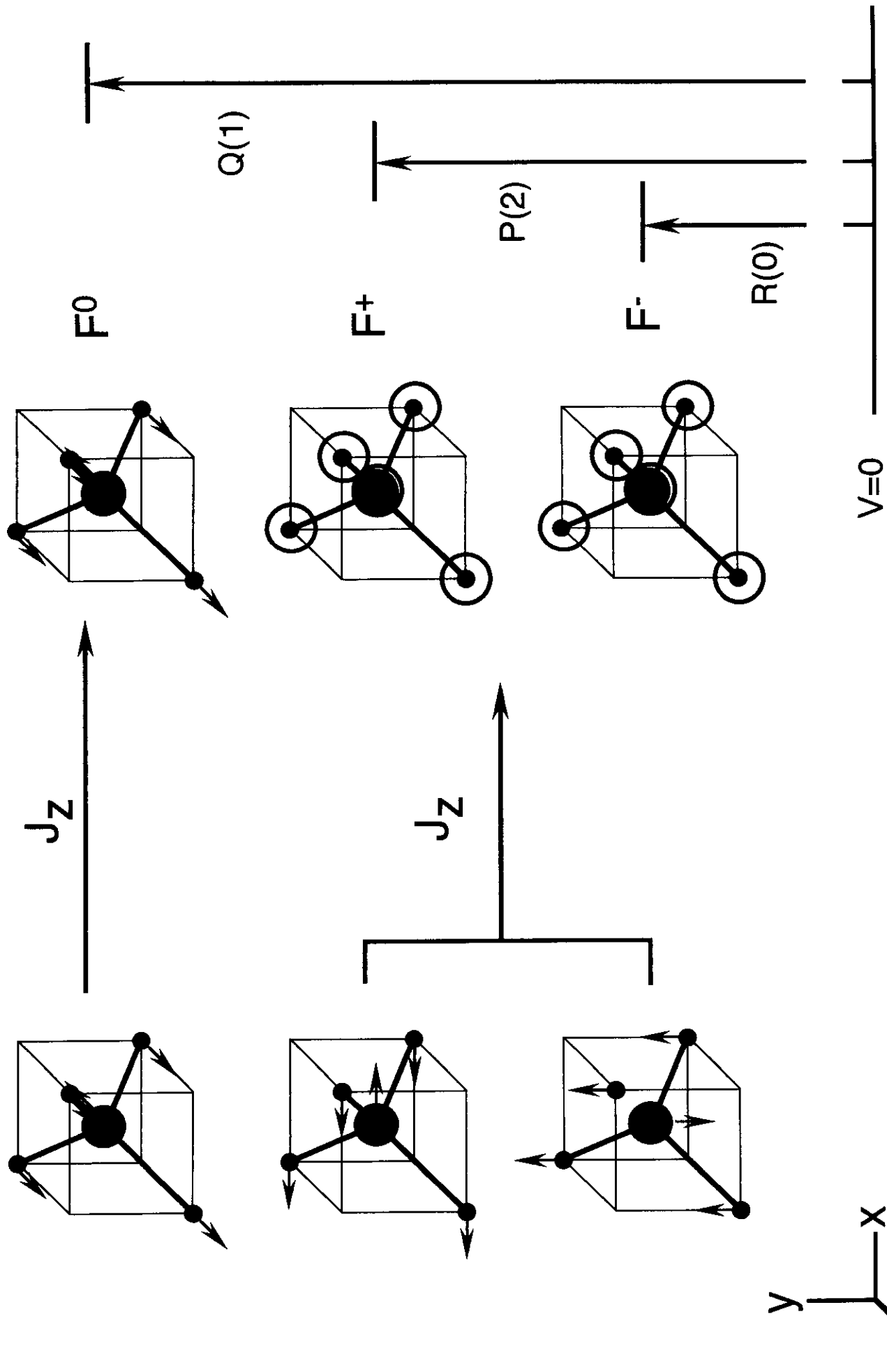


- Excitation of the P(1), R(0), R(1) and R(2) transitions prepares CH<sub>4</sub> in selected rotational levels of  $\nu_3$ ,  $v=1$ .
- Rotational excitation does not dramatically alter CH<sub>4</sub> reactivity over the range of  $J = 0$  to 3.
- Results of Carré and Jackson's quantum dynamics calculation appear for comparison.

# Rotational-State Resolved Sticking Probabilities for $\text{CH}_4$ , $\nu_3$ ( $v=1$ )



# Coriolis Mixing in $v_3$



## Reactivity of the $\nu_3$ Vibrational Sublevels

|      | v = 0 |          |     |                |            |    | v = 1   |    |                |                |       | fractional<br>excitation | total internal<br>energy (cm <sup>-1</sup> ) | Sticking<br>probability |
|------|-------|----------|-----|----------------|------------|----|---------|----|----------------|----------------|-------|--------------------------|----------------------------------------------|-------------------------|
| T    | N''   | $\ell''$ | J'' | RS             | population | N' | $\ell'$ | J' | RS             | VS             |       |                          |                                              |                         |
| R(0) | 0     | 0        | 0   | A <sub>1</sub> | 0.313      | 0  | 1       | 1  | A <sub>2</sub> | F <sup>-</sup> | 0.154 | 3028.7522                | 0.0025 ± 0.0012                              |                         |
| Q(1) | 1     | 0        | 1   | F <sub>1</sub> | 0.563      | 1  | 0       | 1  | F <sub>2</sub> | F <sup>0</sup> | 0.110 | 3029.3058                | 0.0021 ± 0.0010                              |                         |
| P(2) | 2     | 0        | 2   | E              | 0.125      | 2  | -1      | 1  | E              | F <sup>+</sup> | 0.016 | 3030.5024                | 0.0019 ± 0.0013                              |                         |

- Vibrational Sublevels share similar reactivity
- Nuclear motions differ in the phase of the C-H motions
- Reactivity appears to depend on the excitation of a single C-H oscillator and not its phase relative to the other C-H oscillators in the molecule
- Detailed results for  $\nu_3$ , J=2, F<sup>-</sup> state are representative of all three vibrational sublevels.

## Conclusions

- Supersonic Beam + IR Excitation + UHV Techniques
  - Experiments are feasible and widely applicable
  - State resolved data reveal efficacy of energy in different coordinates
- Mechanistic Insights
  - C-H stretch enhances reactivity
  - Efficacy of C-H stretch and  $E_{\text{trans}, \perp}$  comparable
  - $3\nu_4$  bend less effective than  $\nu_3$  C-H stretch
  - Other modes contribute to thermal reactivity
  - Stretch + Bend combination vibrations likely resemble reaction coordinate most
- Comparison with theory
  - Current models overestimate the reactivity of the C-H stretching state.
  - Experimental data point to the importance of both stretching and bending excitation.
  - Statistical model inconsistent with  $3 \nu_4$  reactivity
- Future work
  - Investigate reactivity of combination modes
  - "Tuning" molecular eigenstates

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