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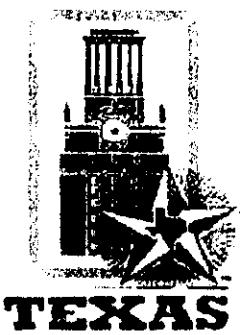
ADRIATICO RESEARCH CONFERENCE on
LASERS IN SURFACE SCIENCE

11-15 September 2000

Miramare - Trieste, Italy

*Second harmonic spectroscopy of Si surfaces with
H, Ge, and B adsorbates:
Experiment meets Microscopic Theory*

Michael Downer
University of Texas at Austin
United States of America



Adriatico Conference on
Lasers in Surface Science:
September 13, 2000



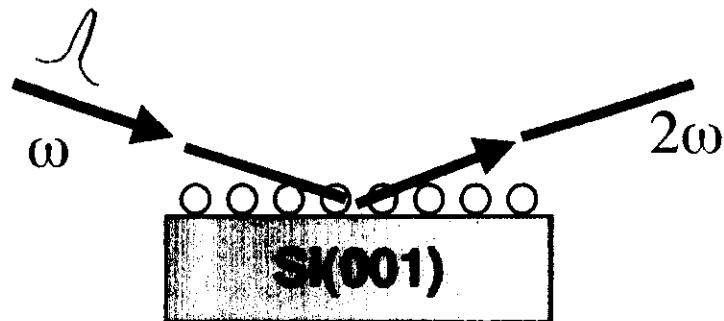
Second harmonic spectroscopy of Si surfaces with H, Ge, and B adsorbates:

Experiment *meets* *Microscopic Theory*

**M. C. Downer, J. G. Ekerdt,
D. Lim, P. Parkinson**
U. Texas at Austin

V. Gavrilenko, R. Q. Wu
California State U.- Northridge

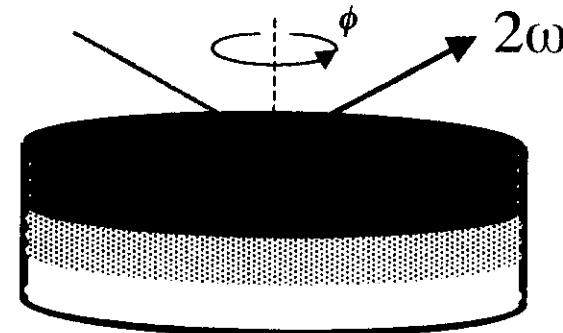
N. Arzate, B. S. Mendoza
*Centro de Investigaciones en Optica,
León, México*



Tom *et al.*, Phys. Rev. Lett. **51**, 1983 (1983)
 Aktsipetrov *et al.*, Sov. Phys. JETP **64**, 167 (1986)
 Sipe *et al.*, Phys. Rev. B **35**, 1129 (1987)

MACROSCOPIC THEORY OF SHG at Si(001)

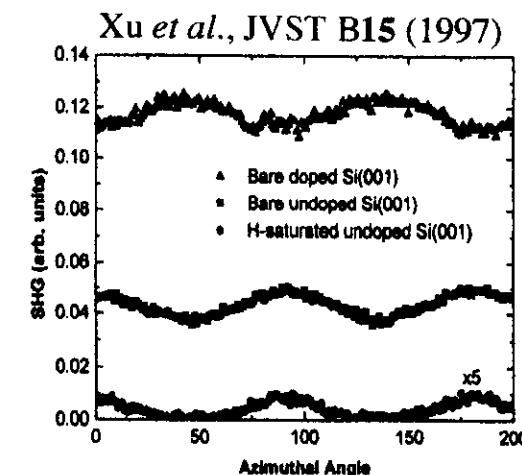
$$P_i(2\omega) = \chi_{ijk}^{(2)surf} E_j(\omega) E_k(\omega) \delta(z) + \chi_{ijkl}^{(2)bulk} E_j(\omega) \nabla_k(\omega) E_l(\omega) + \chi_{ijkz}^{(3)EFISH} E_j(\omega) E_k(\omega) E_z(0)$$

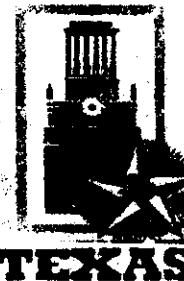


$$\begin{aligned} I_{PP}(2\omega) &= | s_{PP}(\chi_{zz}, \chi_{zxx}, \chi_{xzx}) + b_{PP} + b_{PP}E_z(0) + b_{PP} \cos 4\phi |^2 \\ I_{SP}(2\omega) &= | s_{SP}(\chi_{zxx}) + b_{SP} + b_{SP}E_z(0) + b_{SP} \cos 4\phi |^2 \\ I_{QS}(2\omega) &= | s_{QS}(\chi_{xzx}) + b_{QS} + b_{QS}E_z(0) + b_{QS} \cos 4\phi |^2 \end{aligned}$$

field-dependent

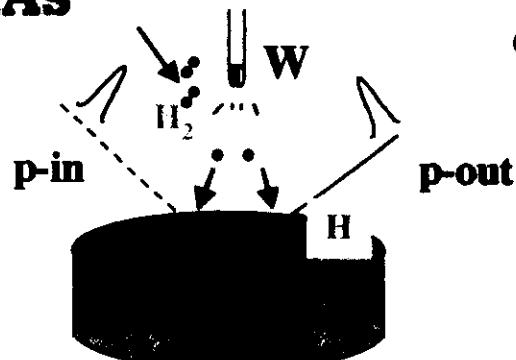
isotropic *anisotropic*



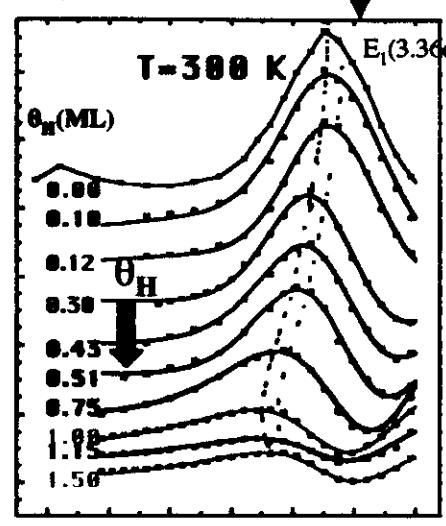


SHG SPECTRA are ADATOM-SPECIFIC

H/Si(001)

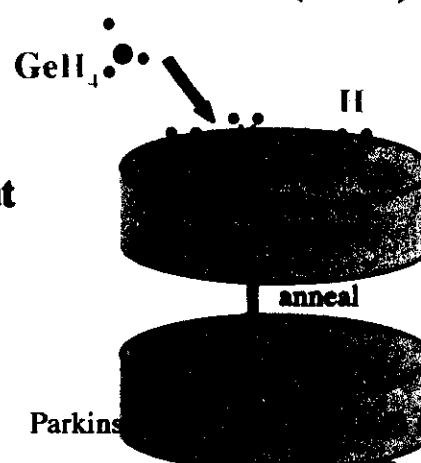


Dadap et.al., PRB 56, 13367(1998)

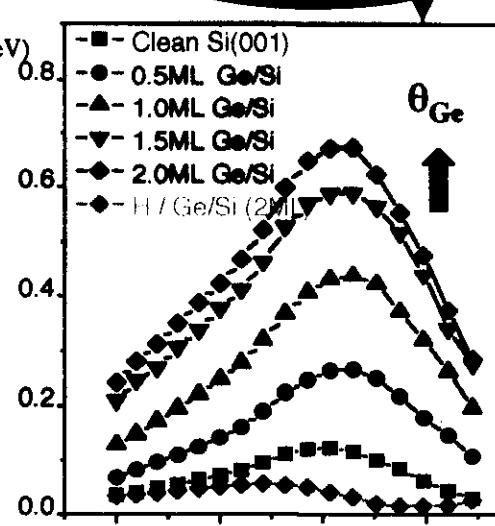


Two Photon Energy(eV)

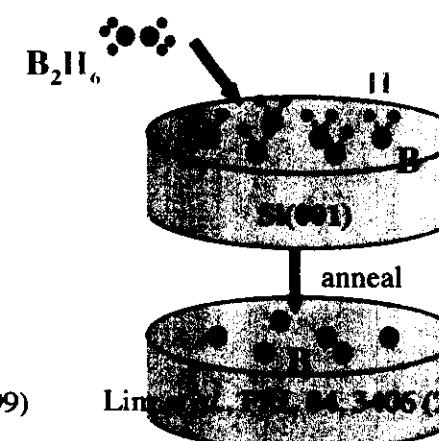
Ge/Si(001)



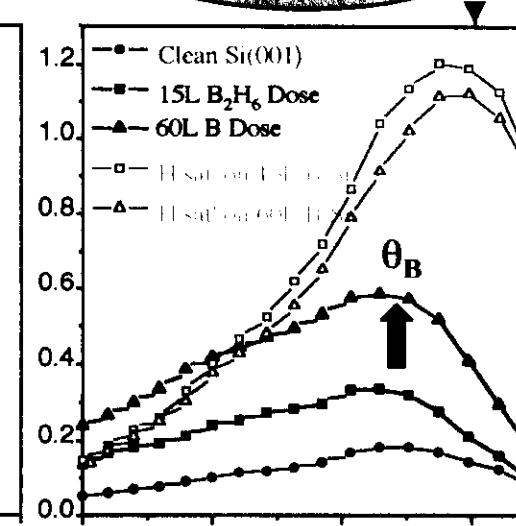
Parkins et.al., PRB 60, 14001(1999)



B/Si(001)



Lindau et.al., PRB 62, 13406(2000)



Two Photon Energy(eV)

Reasons for adsorbate-induced alterations of $\chi^{(2)}$ _{surface}

STRUCTURAL

- adsorbate-induced changes
in near-surface equilibrium
atomic positions
&
strain gradients

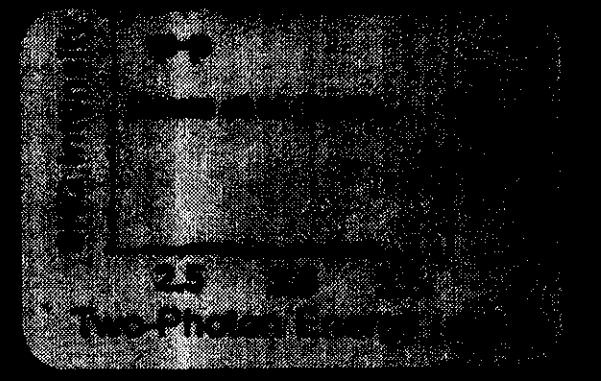
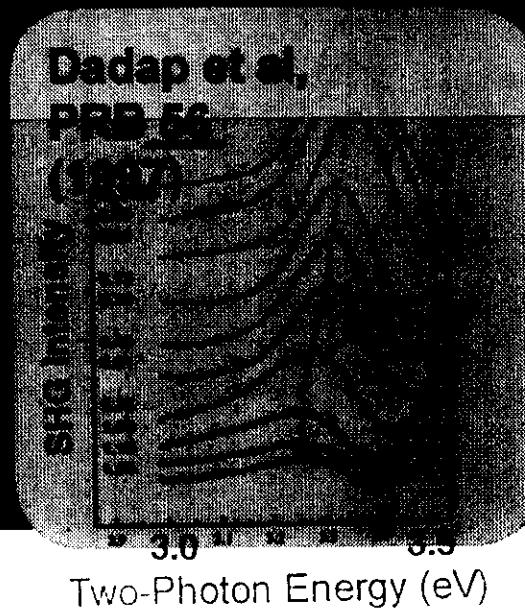
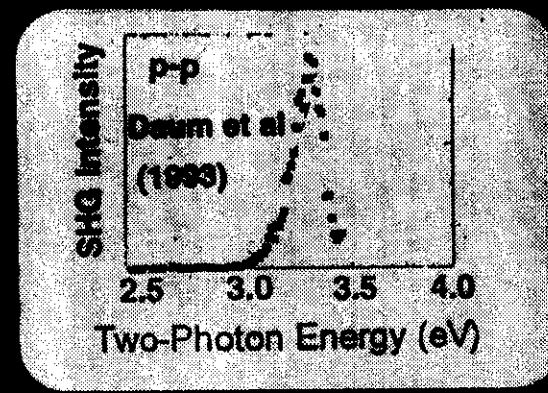
CHEMICAL

- adsorbate-induced
hybridization of surface
atomic orbitals
&
charge transfer

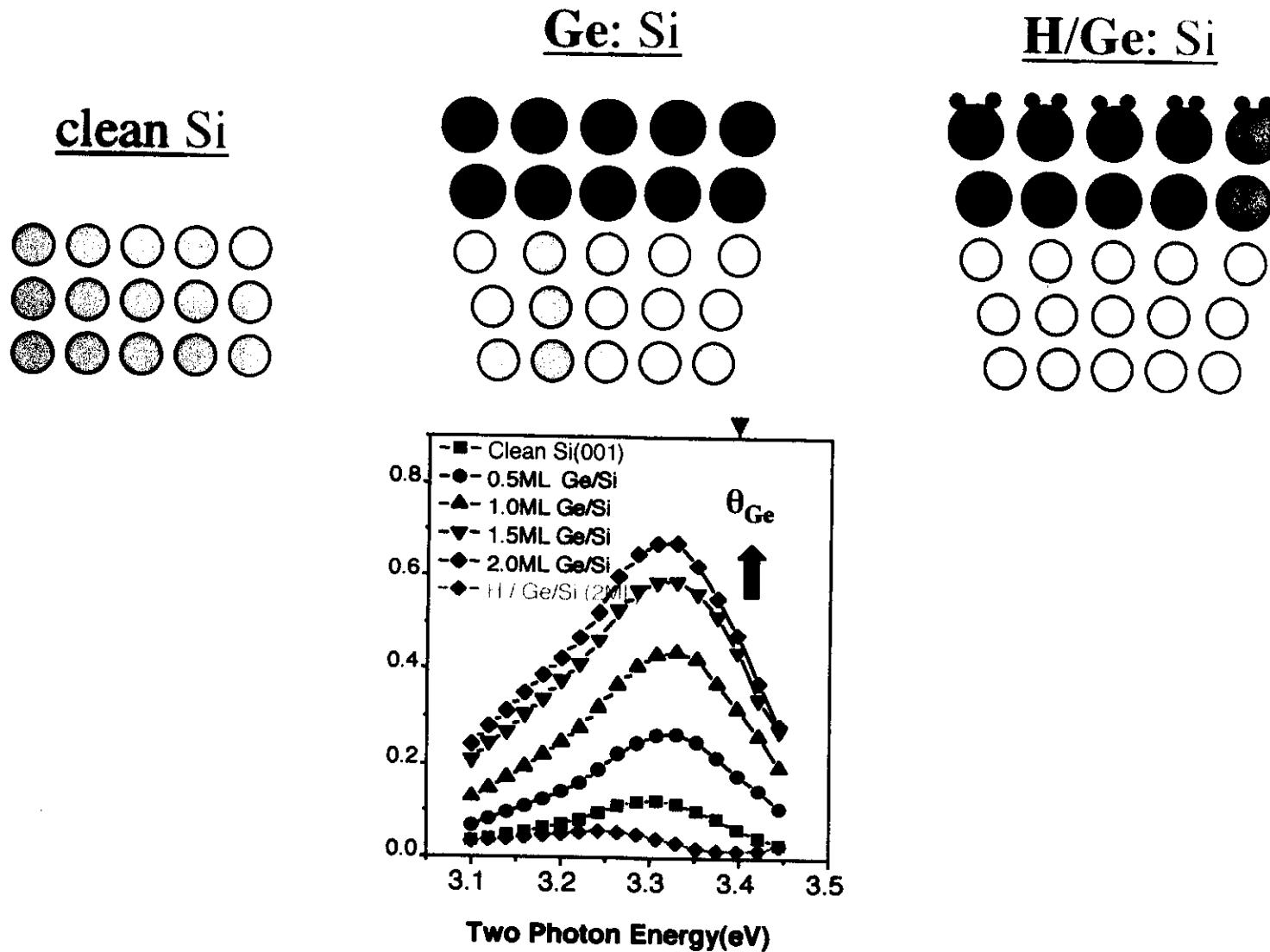
U_T

Physics of SHG at Si(001): 3 types of surfaces occur during CVD

buckled
dimer



SHG trends at H/Ge: Si(001) are uncorrelated with lattice mismatch strain

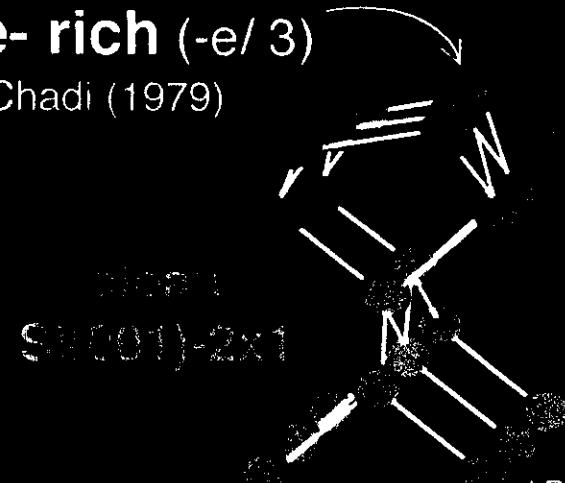


Dimer is strongly POLARIZED on clean surface

Surface structures optimized for 12 ML slab,
using Molecular Dynamics with *ab initio* pseudopotentials

e- rich (-e/3)

Chadi (1979)

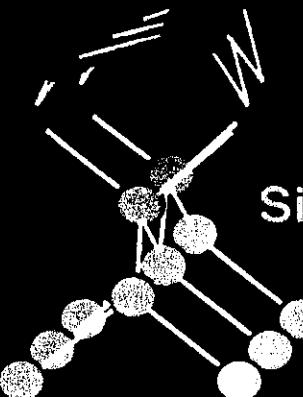


buckled
polarized
dimers

0.7
0.81 Å

e- richer

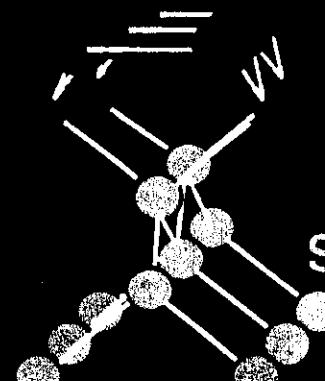
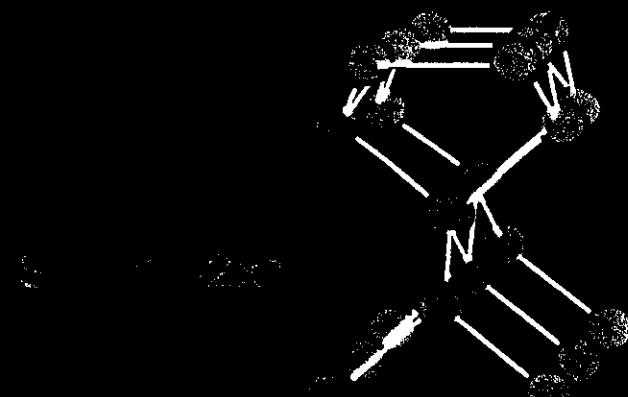
Cho et al.(1994);
Jenkins (1996)



Si(001)-2x1

J.R. Power et al., PRL 80, 3133 (1998)

symmetric
unpolarized
dimers

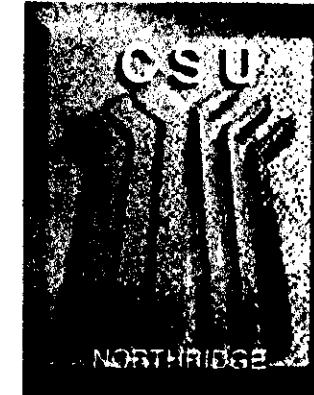


Si(001)-2x1



CIO - Léon

CALCULATION of SURFACE OPTICAL PROPERTIES



INPUT: Optimized Structure

KEY QUANTITY: 2nd order surface susceptibility

$$\chi_{ijk}^{(2)}(-2\omega, \omega, \omega) = -\frac{\pi}{2} \left(\frac{e\hbar}{m}\right)^3 \sum_{g, n, n'} \left\{ \frac{(p_i)_g n \{(p_j)_{n,n'}, (p_k)_{n'g}\}}{(2\omega - \omega_{ng} + i\Gamma_{ng})(\omega - \omega_{n'g} + i\Gamma_{n'g})} + \dots \right\}$$

TO CALCULATE $\chi^{(2)}$, ONE NEEDS: Eigen-values, -states, matrix elements

Semi-Empirical Tight Binding (SETB)

- 5-Si-orbital (sp^3s^*) basis
 - no adsorbate orbitals
 - $\langle sly|p_x \rangle, \langle slx|p_x \rangle$ empirical parameters derived from fits to bulk optical properties

Vogl *et al.*, JPC Solids 44, 365 (1983)

Selloni *et al.*, PR B 33, 8885 (1986)

• DFT

• potentials

• B

• B

• B

Microscopic Theory of Second Harmonic Generation at Si(100) Surfaces

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²*Istituto Nazionale per la Fisica della Materia-Dipartimento di Fisica, II Università di Roma Tor Vergata, Rome, Italy*

(Received 13 May 1998)

SETB with Si sp^3s^* basis

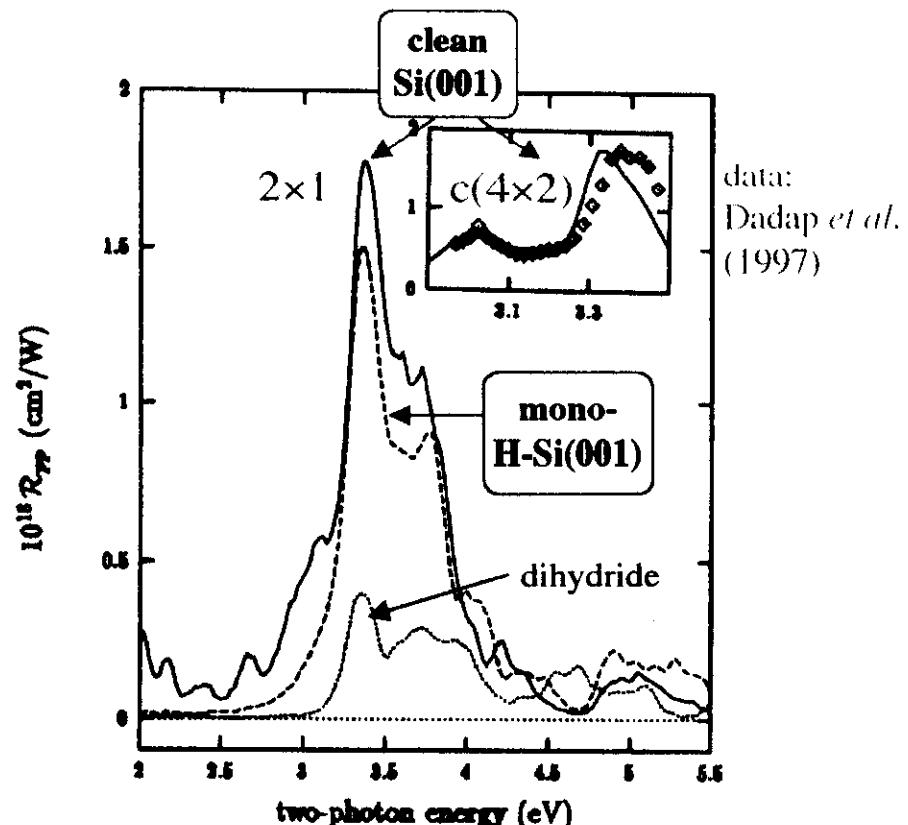
- successful in calculating spectra of **CLEAN** Si surfaces

RDS: Noguera et al., PRL **76**, 4923 (1996)
Power et al., PRL **80**, 3133 (1998)

SHG:

L. Reining et al., Phys. Rev. B **50**, 8411 (1994)
Gavrilenko et al., Appl. Phys. A **60**, 143 (1995)
Mendoza et al., PRL **81**, 3781 (1998)

- less successful in calculating adsorbate effects



Surface nonlinear optics is a new frontier for *ab initio* calculations

I. Ground State Properties:

Density Functional Theory ($E_{gs} \propto f[\rho(\mathbf{r})]$)

Hohenberg & Kohn, *Phys. Rev.* **136**, B864 (1964)

Bernholc, *Physics Today* **52**, 30 (1999).

predict w. few % accuracy:

- lattice constants
- atomic positions
- elastic properties
- phonon frequencies

II. Bulk Linear Optical Properties: *involve excited states*

- quasi-particle & XC corrections for excited states
- Local Field corrections (~15%)
- Excitonic effects

review of 1990's progress:
Lambrecht & Rashkeev, *Phys. Stat. Sol (b)* **217**, 599 (2000)

III. Surface Nonlinear Optical Properties ($\chi^{(2)surf}$)

- No periodic BC's \Rightarrow finite slabs (10-50 ML) \Rightarrow many \mathbf{k} -points,
- Local Fields vary strongly
across surface SHG source region \Rightarrow no solution yet
unit cells \leq diamond

Absolute $\chi^{(2)surf}$ calculations not accurate yet.

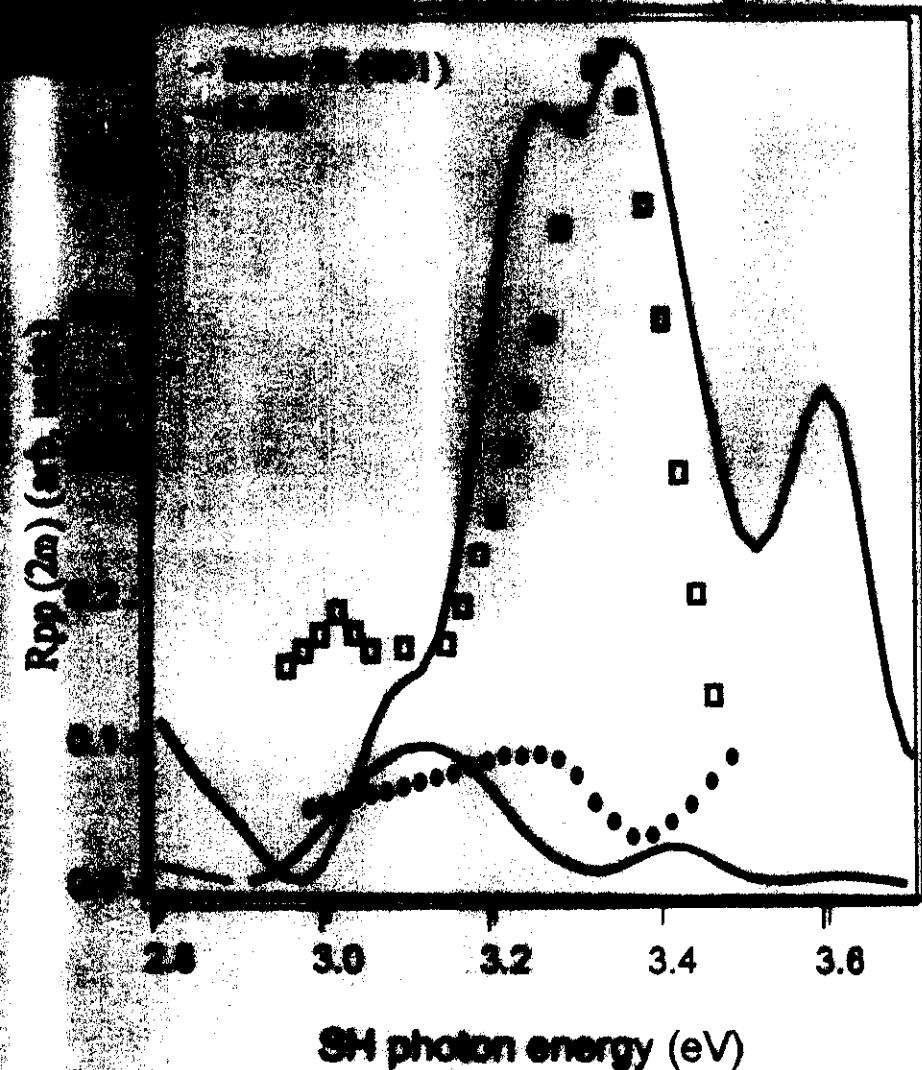
\Rightarrow settle for calculating *changes* in $\chi^{(2)surf}$

\Rightarrow use semi-empirical approaches (e.g. SETB)

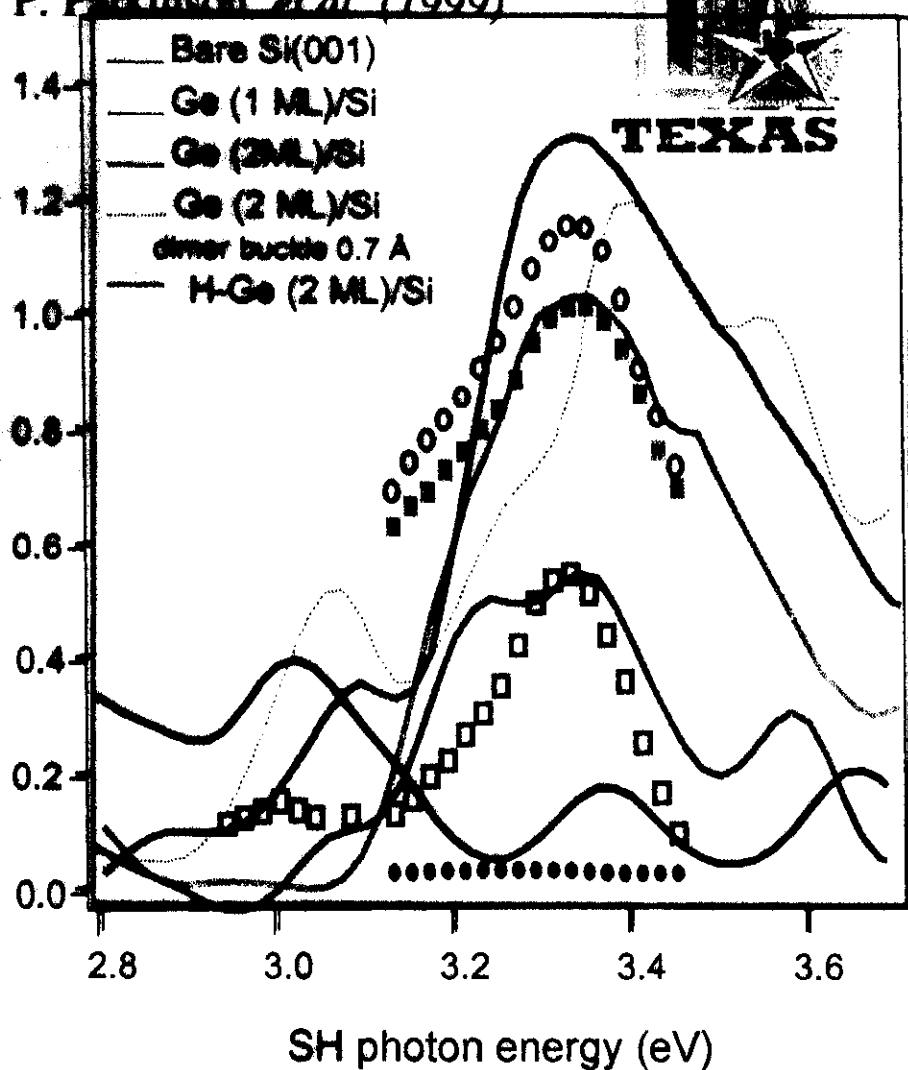
Comparison of ab initio results with experiment

6

M. Gartland and R. Q. Wu



J. Dadap et al. (1997)
P. Parkinson et al. (1999)



Calculated SHG sensitive to:

- eigenstate quality
- realistic description of charge transfer on surface
- convergence of self-consistent calculations

insensitive to:

- small deviations from equilibrium structure

Details of calculation

Ab initio pseudopotentials for
Si, Ge, H generated by scheme of
Bockelie-Hansen-Schiltter (1992)

12 ML slab
adequate for convergence

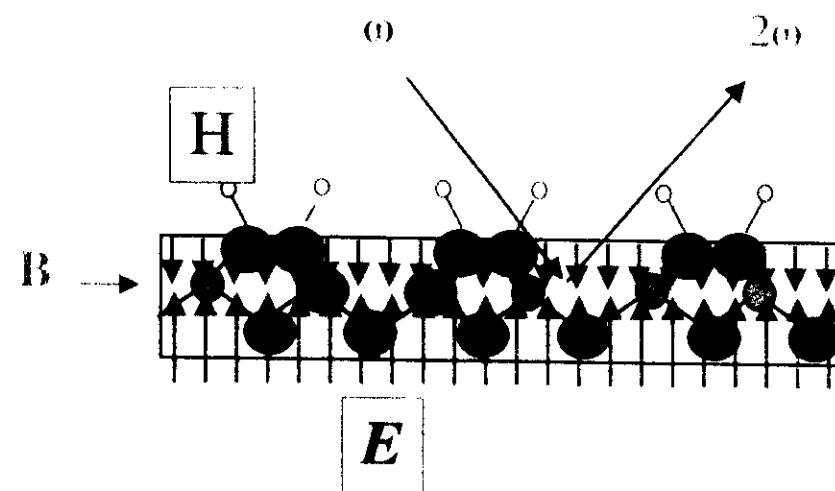
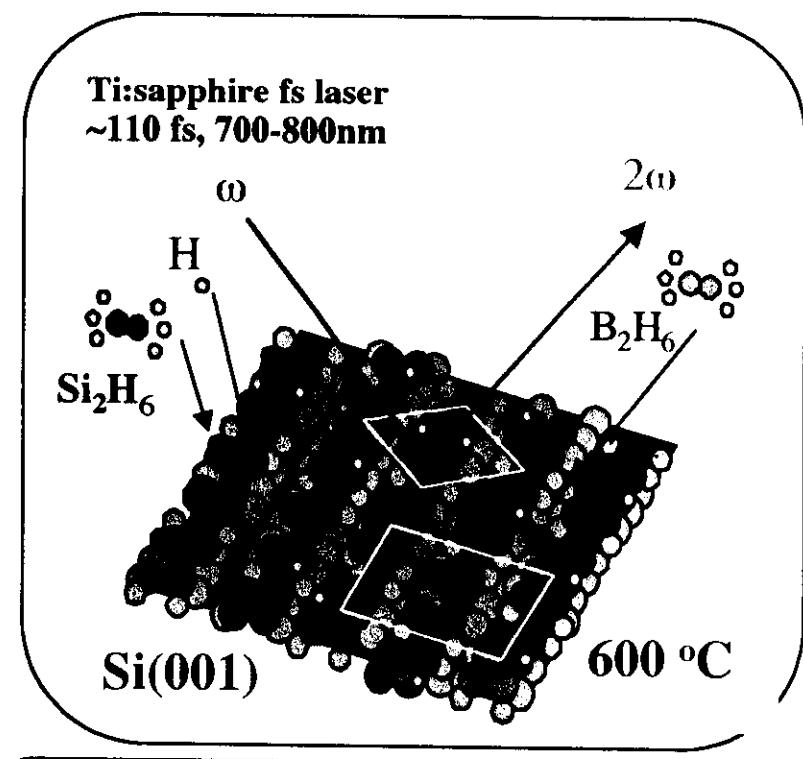
Optical calculations used:
*Eigen-value,-functions from direct
diagonalization of Hamiltonian
after full convergence of $\mu(r)$ calc's*

**48 k-points in irreducible 2D BZ
 $E_{\text{cut}} = 15$ Ry**

Stable 2D Boron layers can be fabricated at Si surfaces

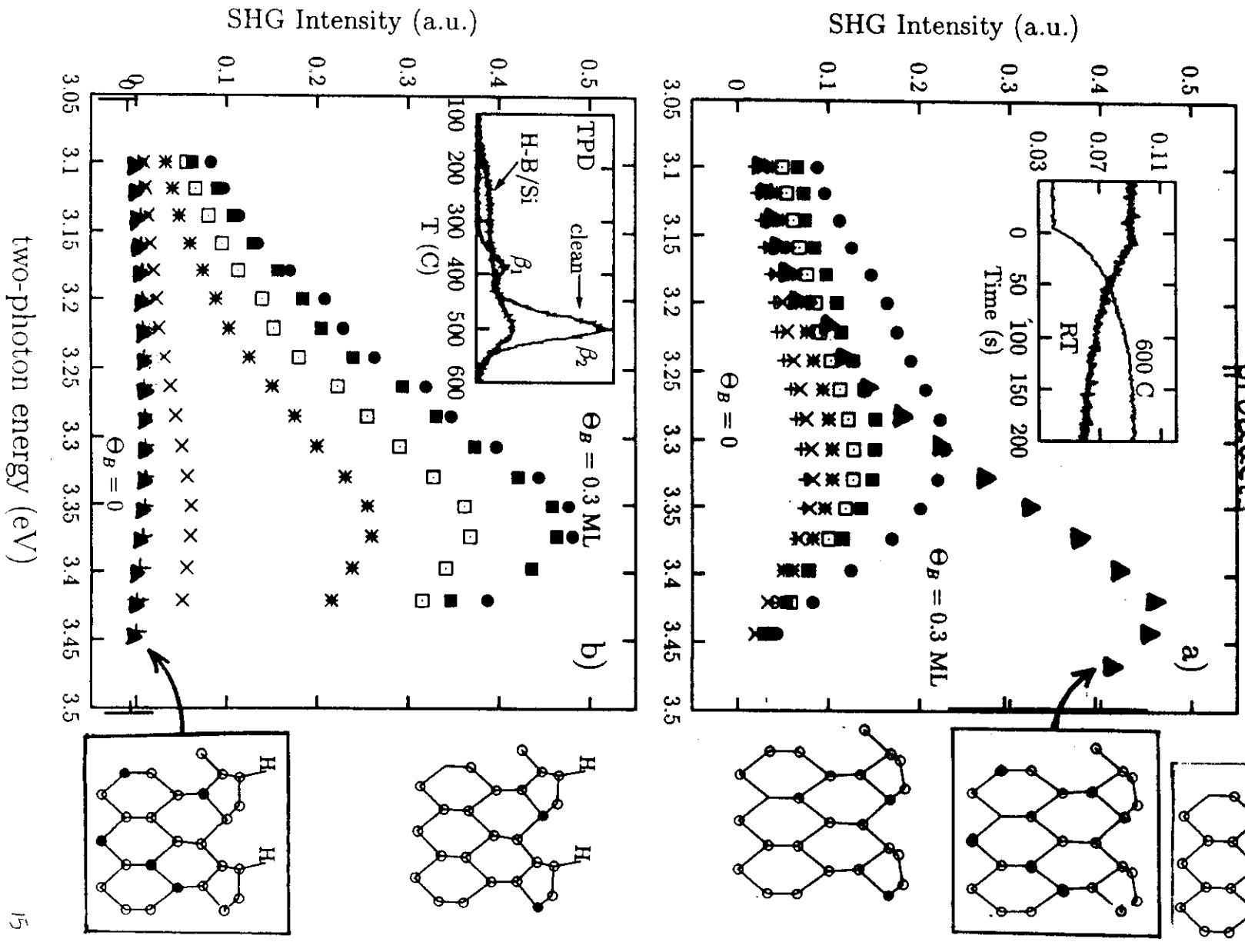
Headrick *et al.*, PRL **63**, 1253 (1989)

Wang *et.al.*, PRL **74**, 403 (1995)



- B occupies 2nd layer & electrically active
- local c(4X4) reconstruction on $Si(001)$

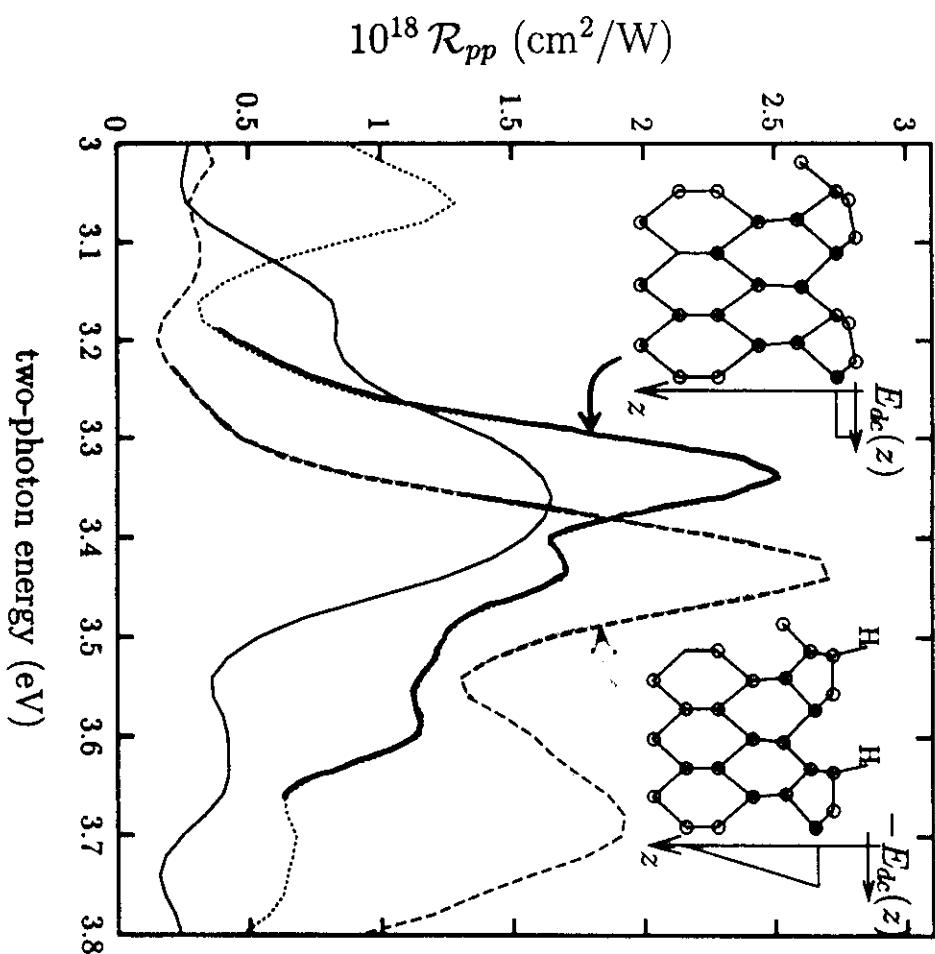
1) 2nd layer B incorporation
2) bulk B doping
3) surface B_2H_6 dissociation
product





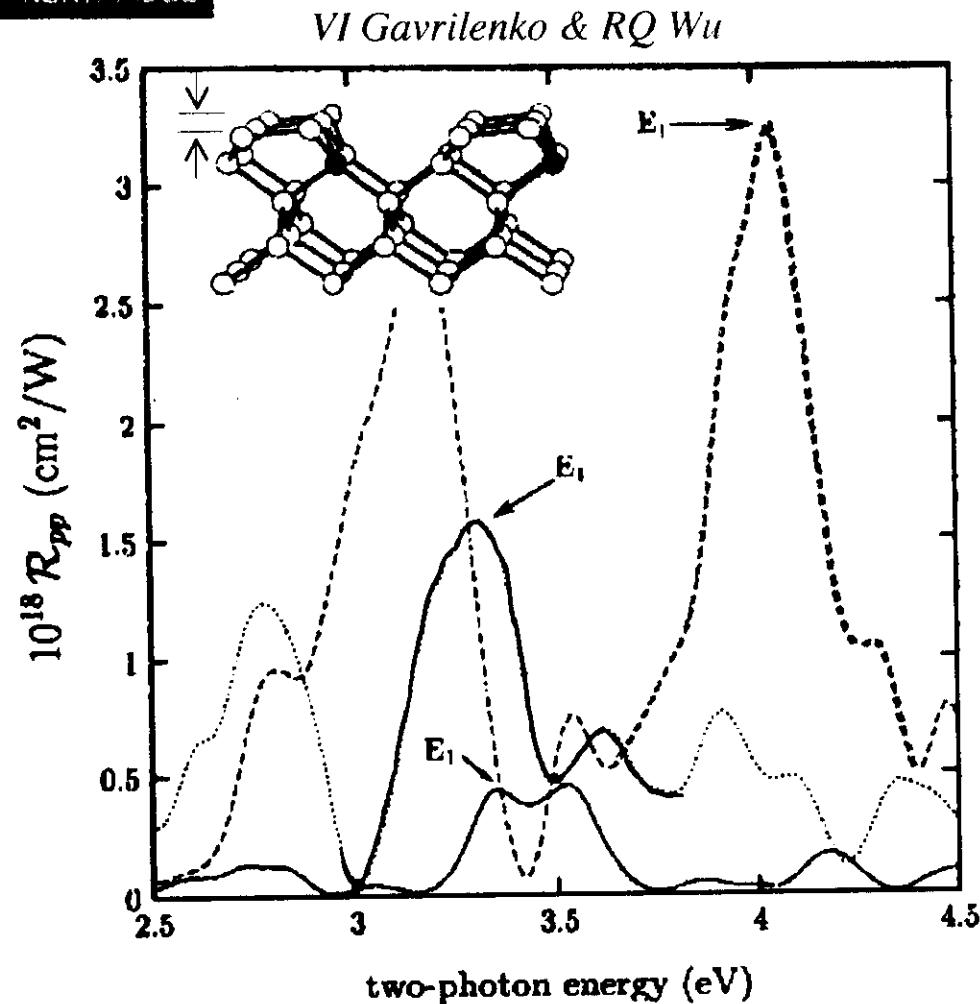
Semi-Empirical Tight Binding (SETB) Model of SHG at B-Si(001) and H-B-Si(001)

- input structure from DFT energy minimization (16 ML slab)
- replace B with Si
- account for effect of B via surface field $E_{dc}(z)$
- CB shifted by $\Delta E = 0.48$ eV





Ab initio calculation of reflected SHG for clean, B- and H-B-Si(001)



Details of calculation

Ab initio pseudopotentials generated by schemes of

- Bachelet-Hamann-Schlüter, PRB (1982) for Si,H
- Troullier, Martins, PRB **43**, 1993 (1991) for B.

Slab of 8 atomic (001) layers

- Quantum shifts negate need for QP corrections
- E_1 occurs near 3.4 eV.

Equilibrium atomic positions

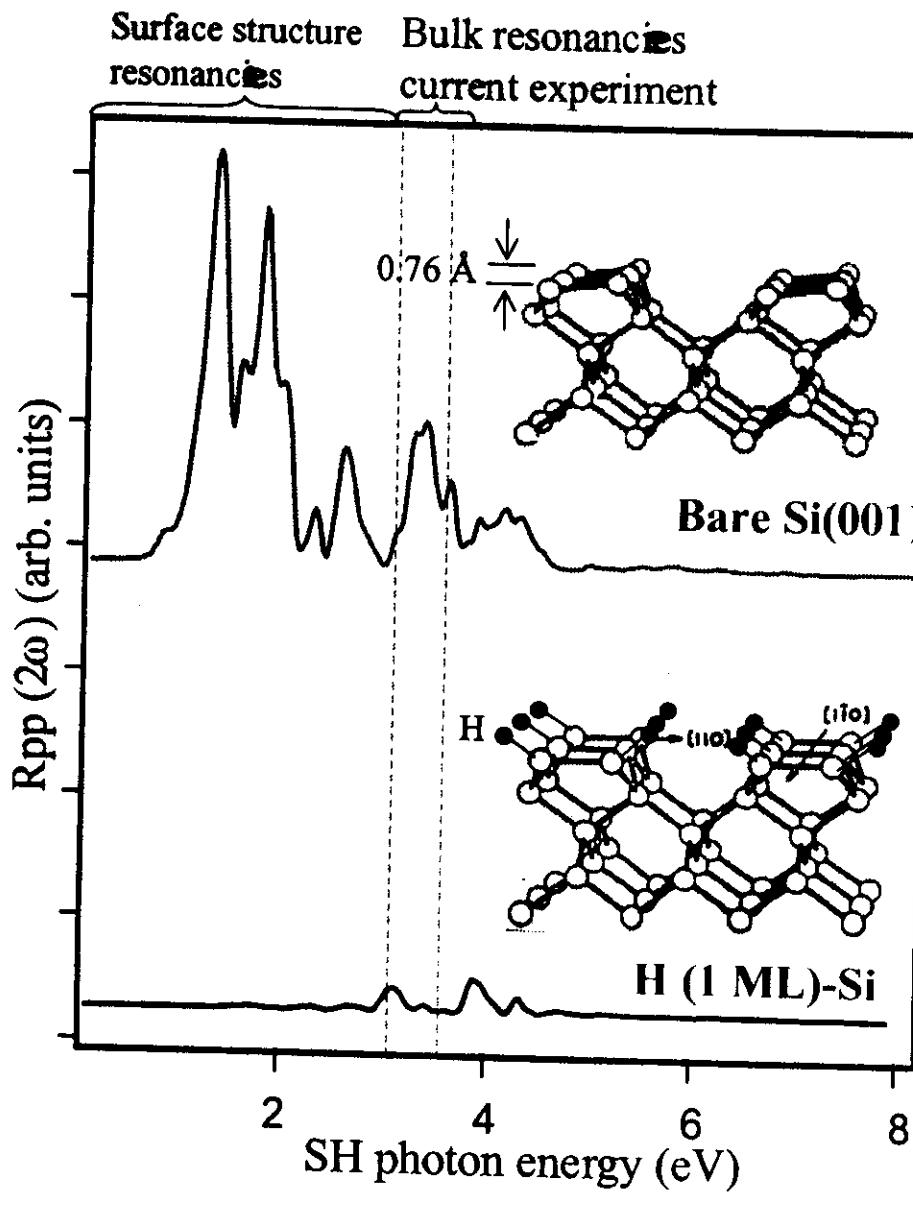
- total energy minimization using $E_{\text{cut}} = 17$ Ry, molecular dynamics method.
- clean Si: 2x1 unit cell
- B-Si: c(4x4), but 2x1 with one B per unit cell used to make calculation tractable

Optical calculations

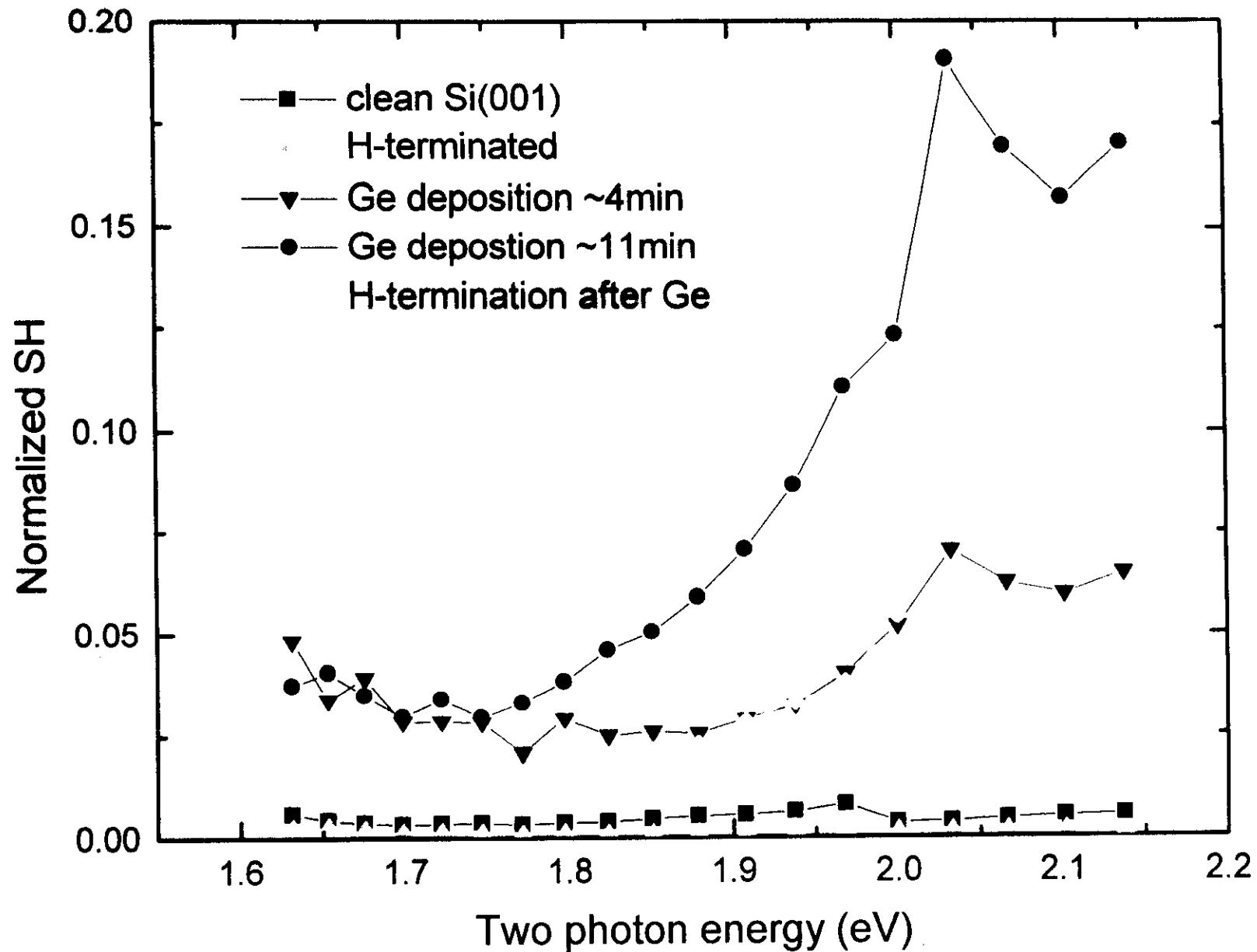
- used $E_{\text{cut}} = 31$ Ry

Ab initio results:

1. Predictions for future experiments



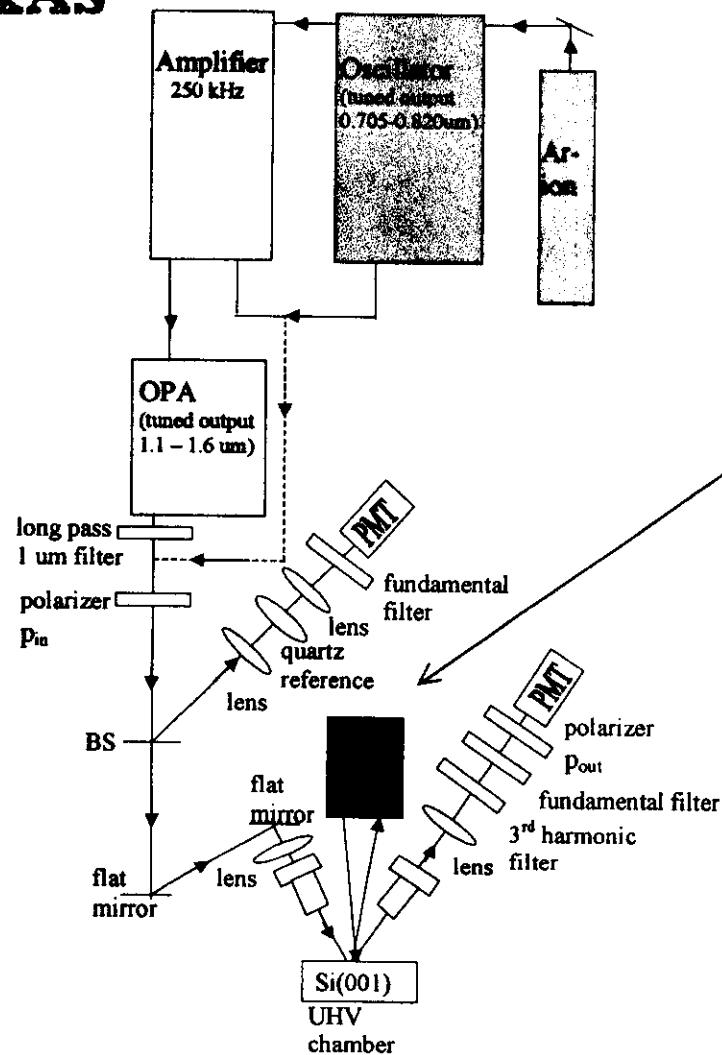
IR-SHG spectroscopy





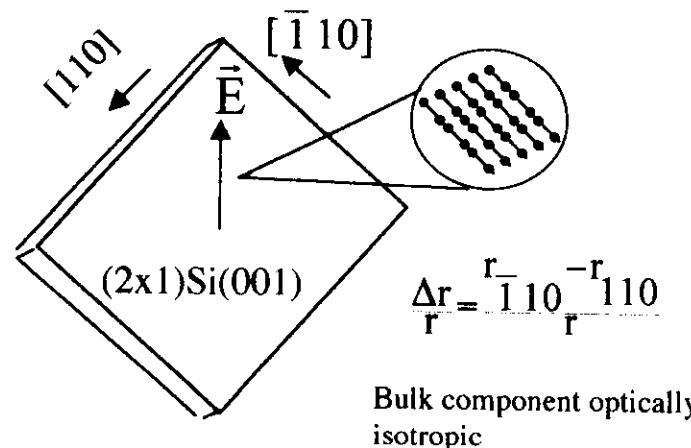
Simultaneous spectroscopic SHG and Reflectance-Difference Spectroscopy (RDS)

Experimental setup



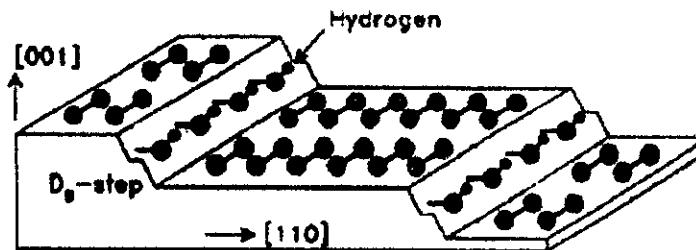
L. Mantesse, D. Lim

RDS is highly surface-sensitive



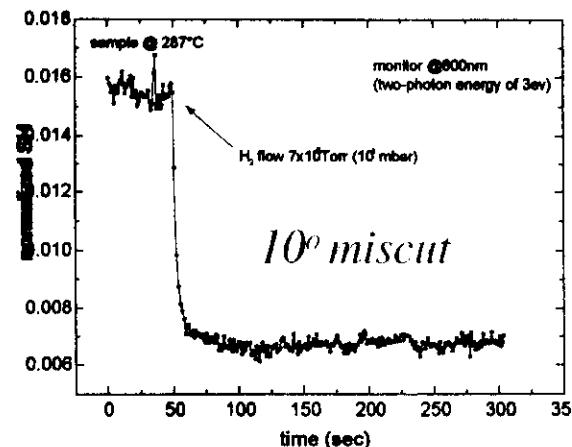
Bulk component optically isotropic

SHG vs. RDS response to H₂ adsorption on vicinal Si(001)



H₂ adsorbs selectively on D_B steps

I. SHG

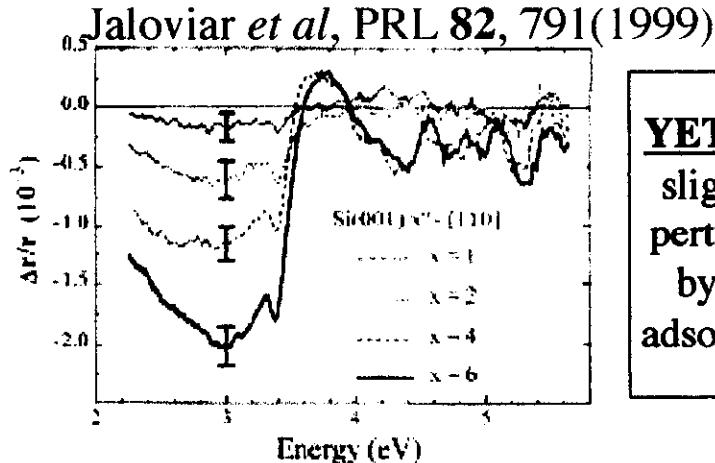


Kratzer *et al.*, PRL 81, 5596 (1998)
Raschke & Höfer, PRB 59, 2783 (1999)

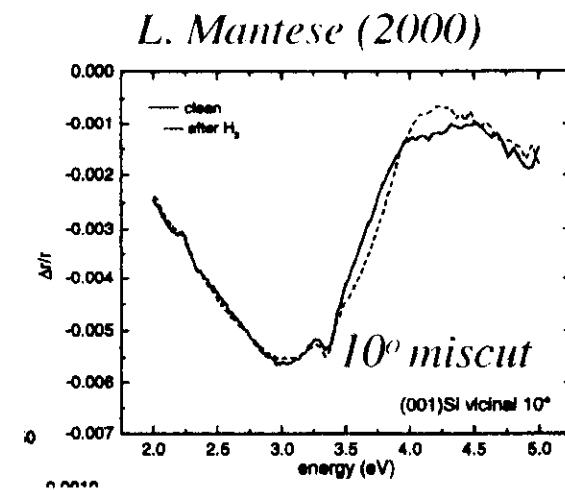
strong step-induced SHG contribution
quenched by H₂ adsorption

II. RDS

strong step-induced optical anisotropy

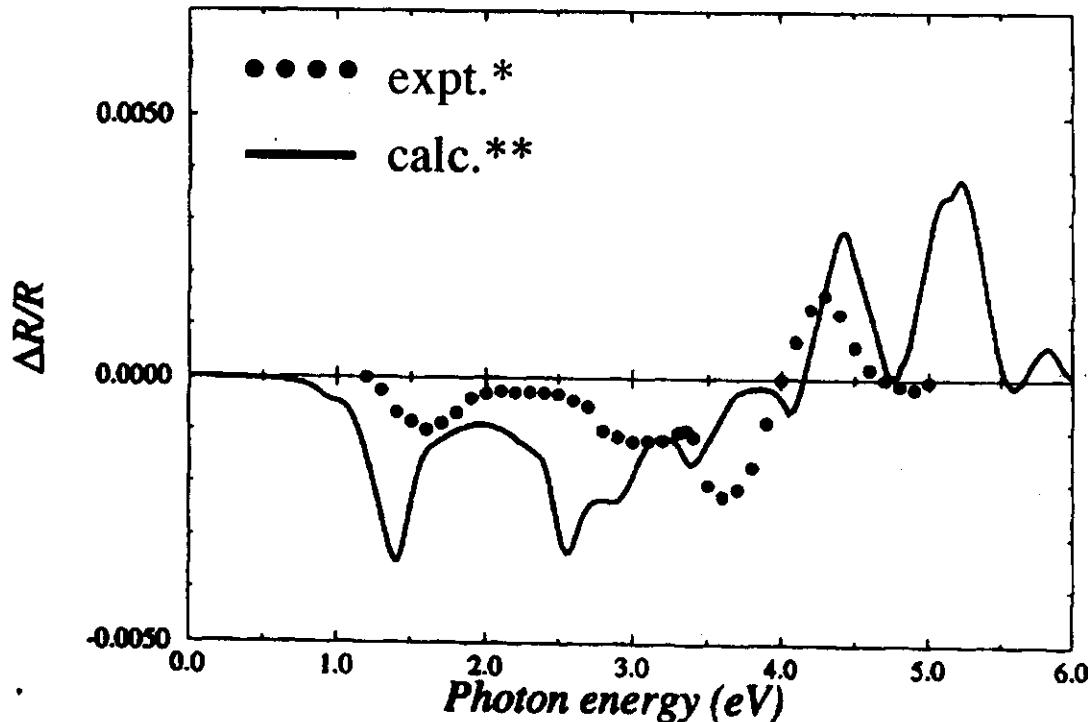


YET only slightly perturbed by H₂ adsorption



RDS of single-domain Si(001)(2x1): experiment* & calculation**

$$\updownarrow \frac{\Delta R(\omega)}{R_0} = \frac{R[110] - R[1\bar{1}0]}{R_0}$$



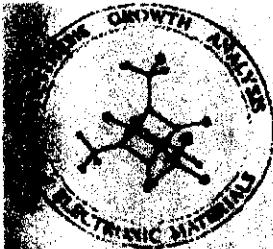
Calculated RDS converges only for $N_{sl} > 12$

(vs. $N_{sl} < 8$ for SHG)

- 64 k-points in irreducible BZ
- $E_{cutoff} = 15$ Ry

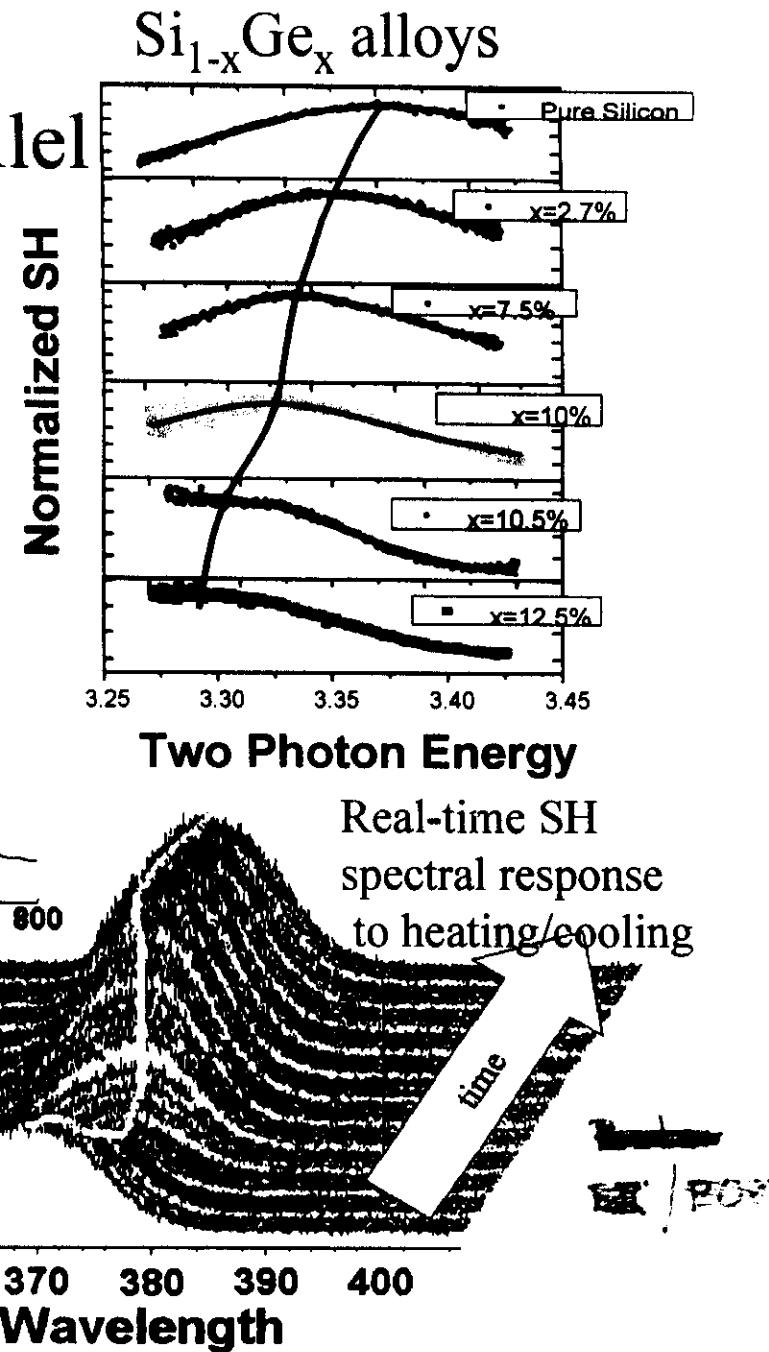
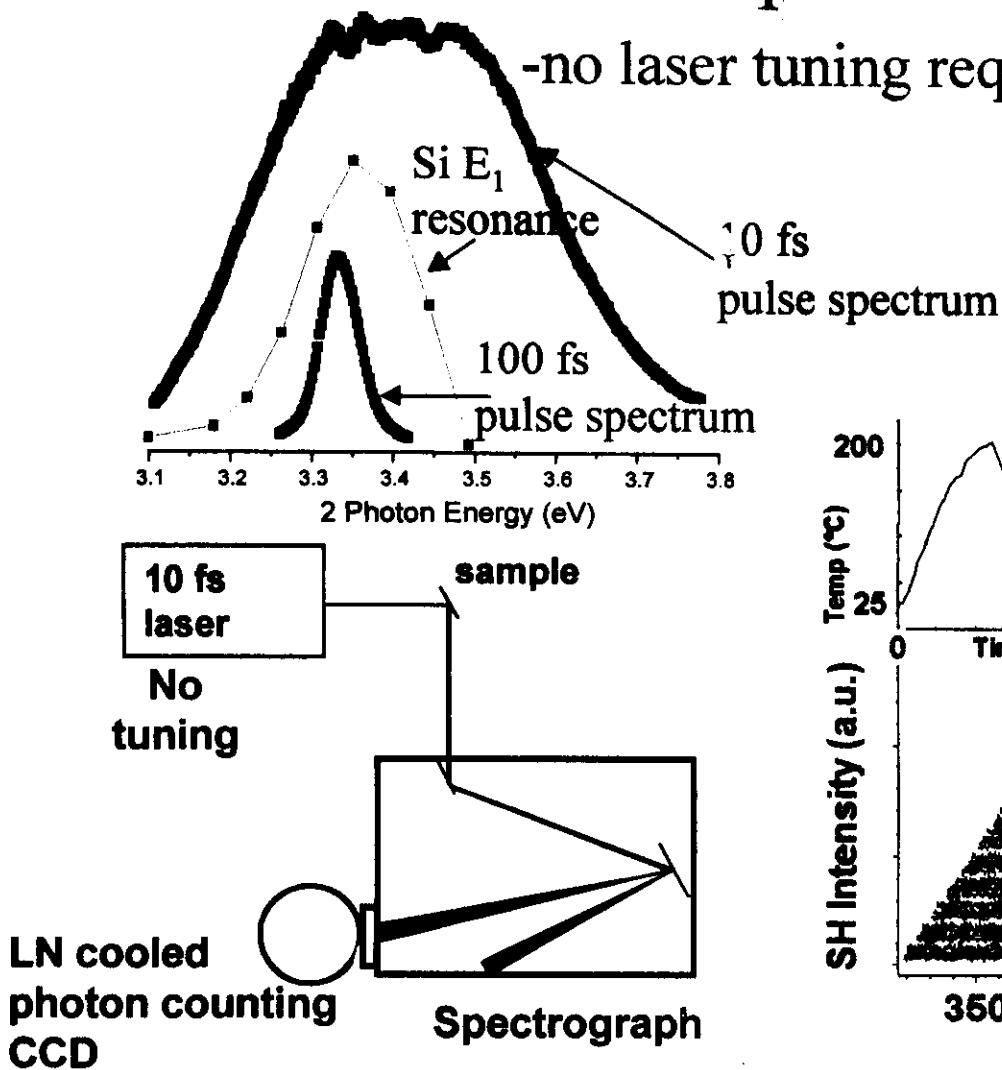
* R. Shioda & J. van der Weide, PRB **57**, R6823 (1998).

** V. Gavrilenko & R. Q. Wu (2000).

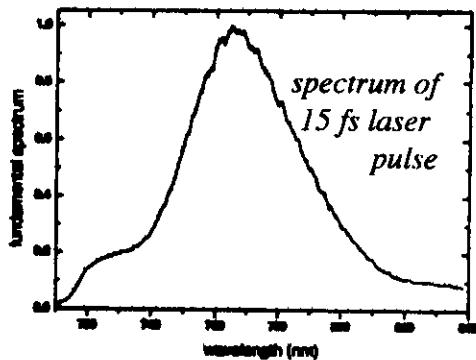


4. SHG in a Box

Broadband Real-time parallel Harmonic spectroscopy

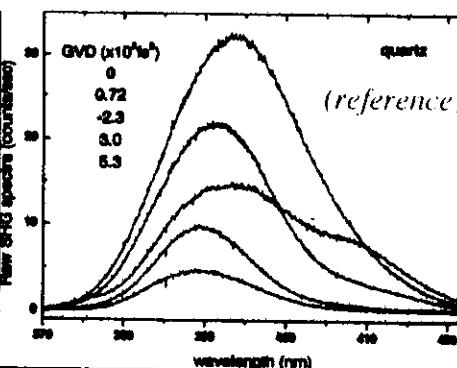
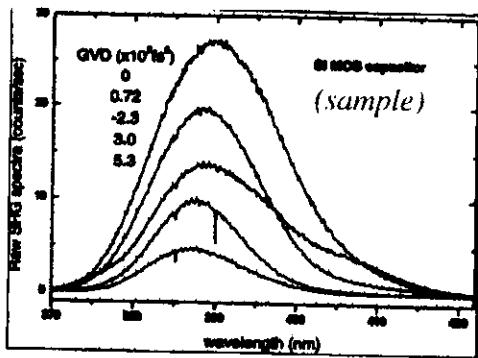


P. T. Wilson, Y. Jiang, O. A. Aktsipetrov, E. D. Mishina M. C. Downer, submitted to Optics Letters(1999)



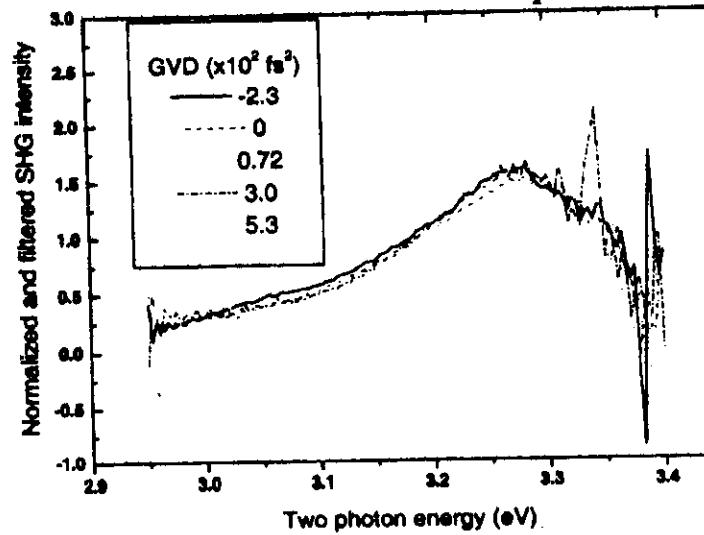
Normalized SHG intensity spectra are chirp-independent

Chirp-dependent raw SHG spectra



- Pulse chirp varied by adding glass to beam path
- Raw spectra change drastically with chirp

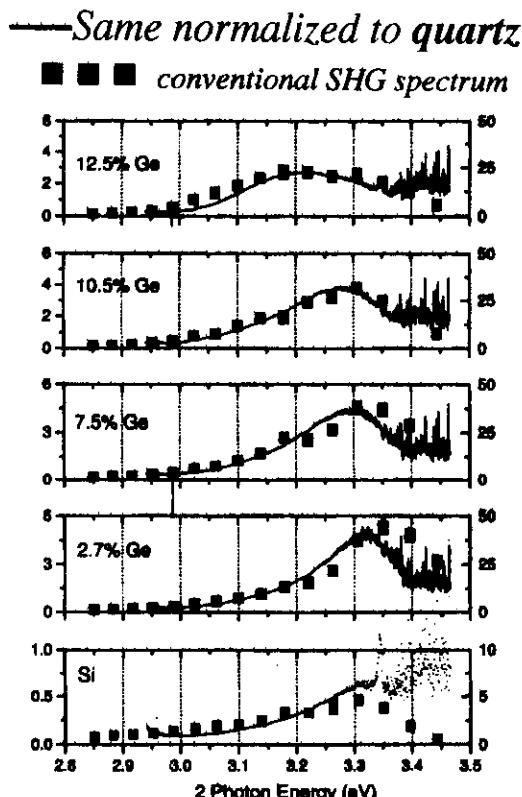
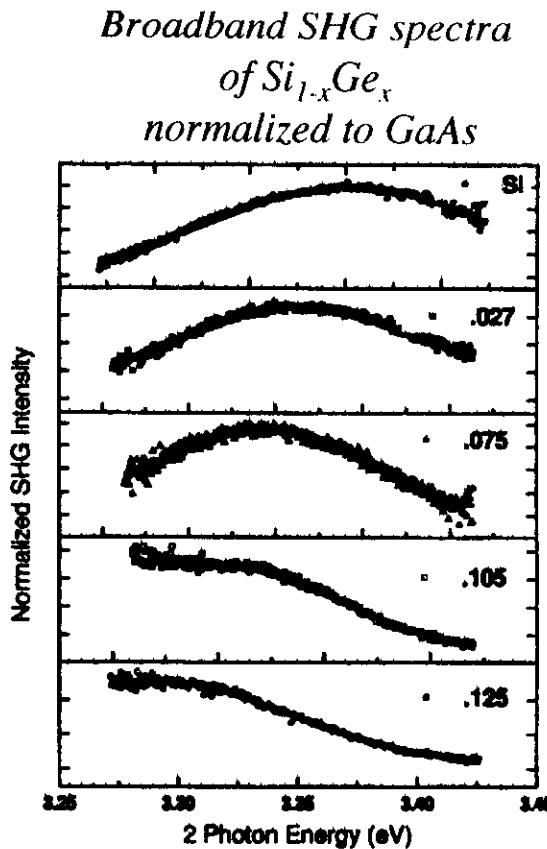
Normalized SHG spectrum



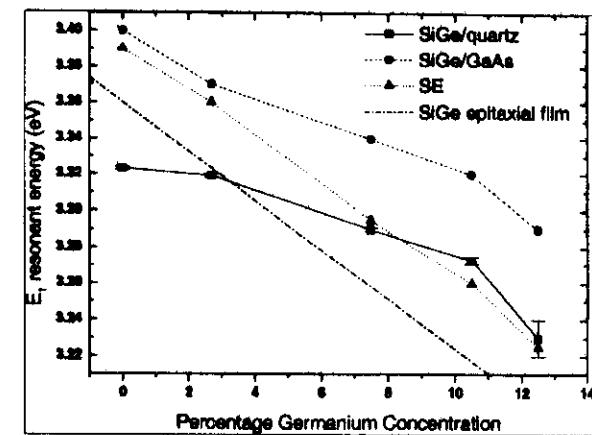
- Normalized spectra agree quantitatively in shape and amplitude

The reference sample must be spectrally flat

Chemical shifts of SiGe E₁ peak



E₁ peak position



- Accurate normalization enables quantitative measurement of material properties

- Broadband & conventional spectra agree quantitatively
- Broadband spectra more robust in middle of range, noisier at extremes

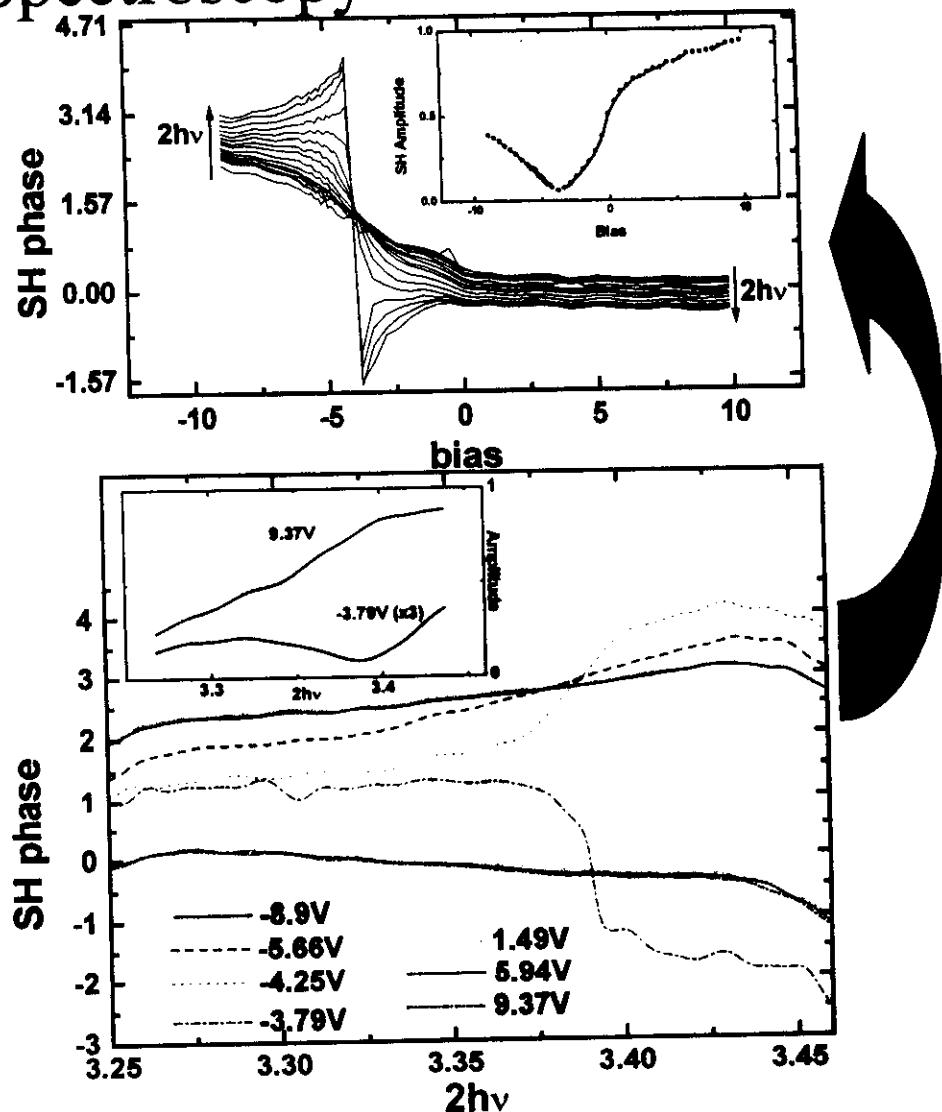
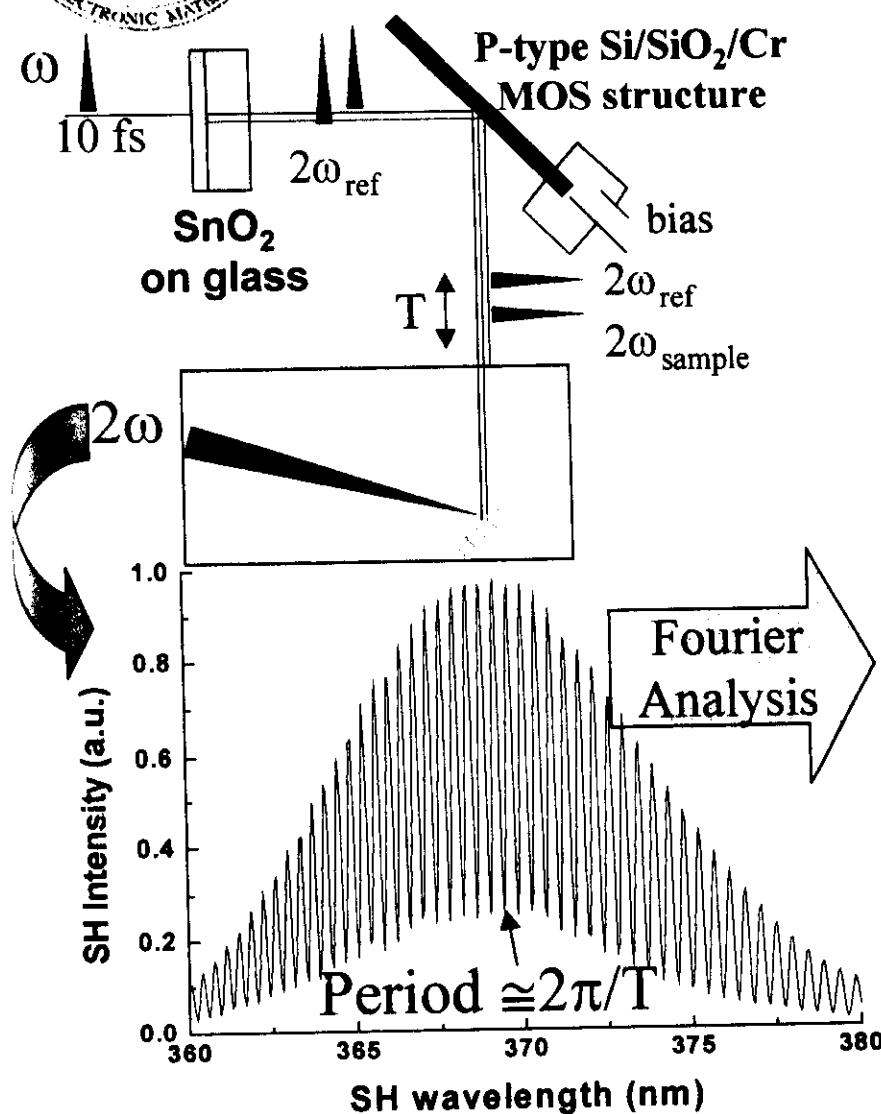
SHG phase measurements

- Phase of SHG is sensitive to polar orientations (Kemnitz, 1986), and can be necessary for distinguishing resonant and background contributions (Stolle, 1996).
- Time-domain measurements require narrow bandwidth tuning and delay variation
- Kramers-Kronig relations allow better lineshape fitting as in spectroscopic ellipsometry



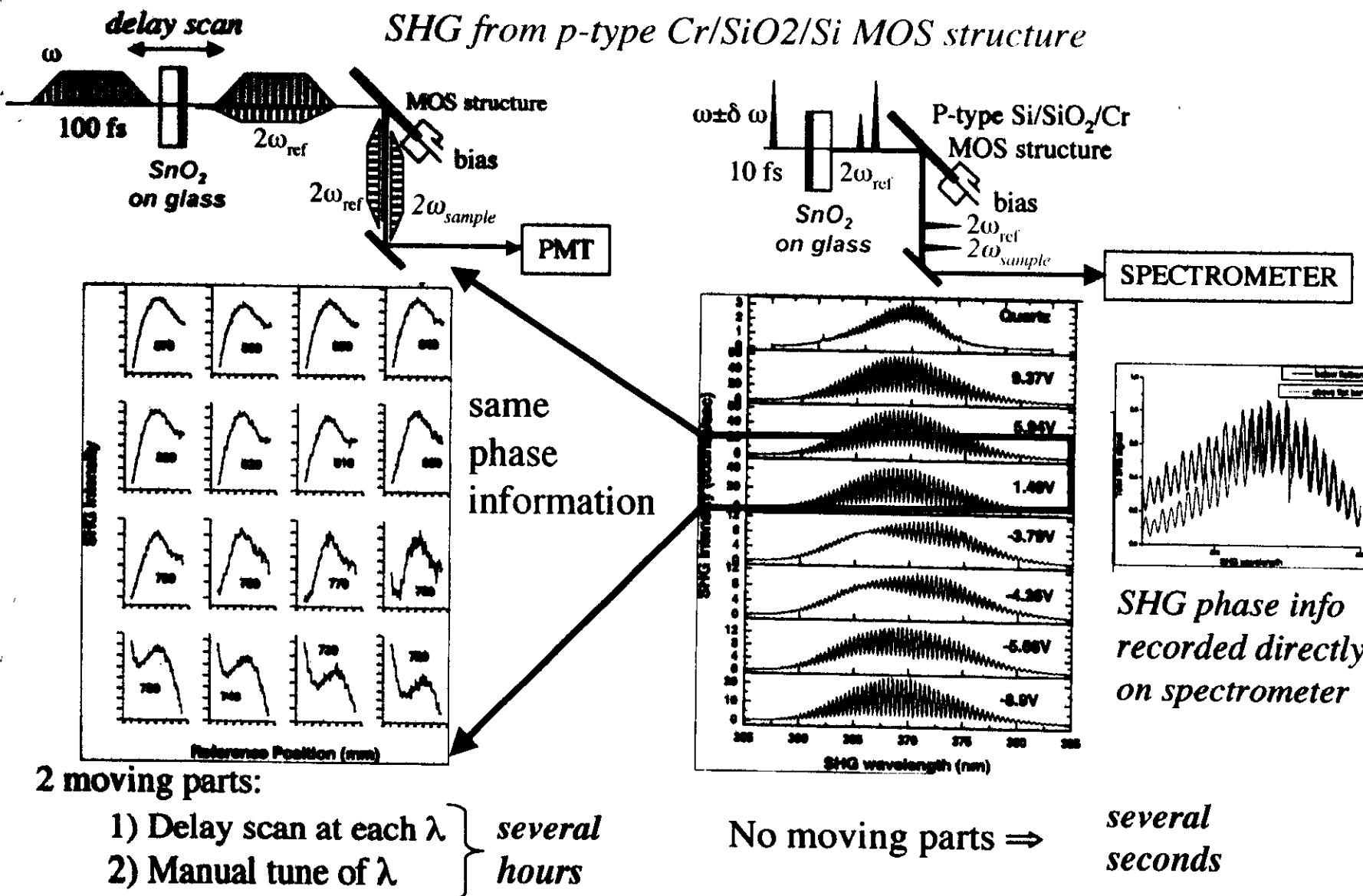
4. SHG in a Box

Frequency Domain Interferometric Second Harmonic Spectroscopy



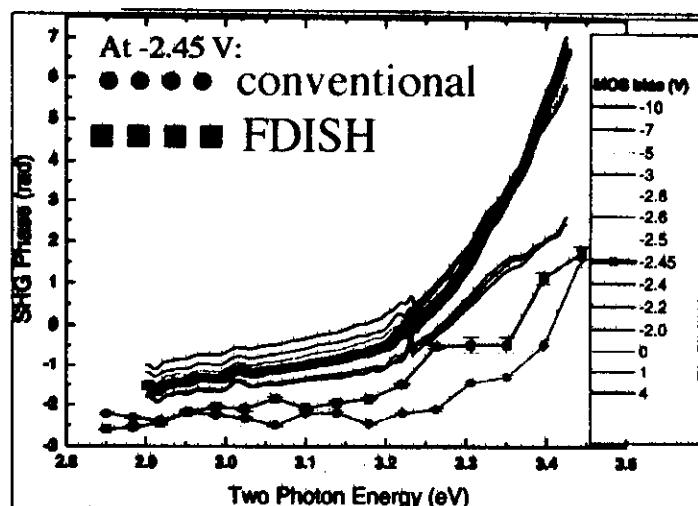
P. T. Wilson, Y. Jiang, O. A. Aktsipetrov, E. D. Mishina M. C. Downer, submitted to Optics Letters(1999)
accepted

Conventional vs. FDISH spectral phase measurements



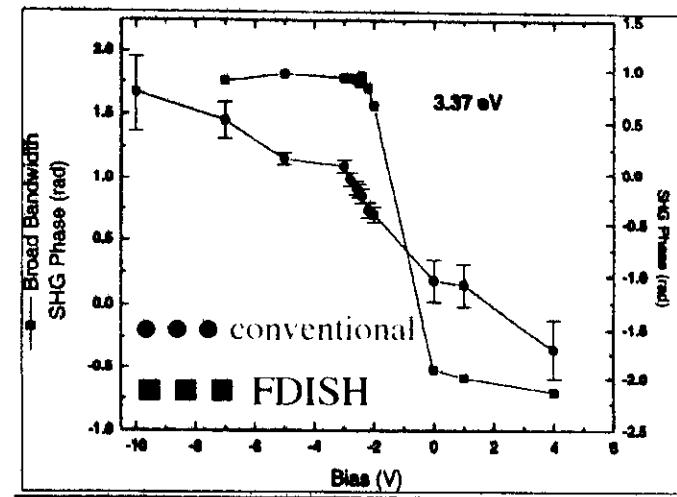
Conventional vs. FDISH spectral phase measurements: results

SHG spectral phase



- conventional measurement based on only 1 oscillation cycle.
- FDISH sensitive to ± 1 fs errors in time delay T.

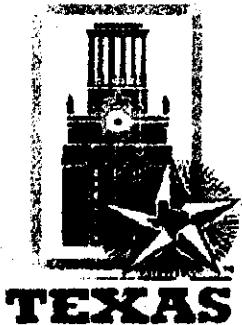
Bias dependence of SHG phase



Sharper bias dependence in FDISH data
 \Rightarrow less carrier screening in 10 fs

$$\Delta\phi_{conventional} = [\varphi_{sample}^X + \varphi_{sample}^G] - [\varphi_{reference}^X + \varphi_{reference}^G + \varphi_{standard}^{reflection}] + \frac{\omega}{c} [n_{air}(\omega) - n_{air}(\omega/2)]$$

$$\Delta\phi_{FDISH} = [\varphi_{sample}^X - \varphi_{sample}^G - \varphi_{sample}^{reflection}] - [\varphi_{standard}^X + \varphi_{standard}^G - \varphi_{standard}^{reflection}]$$



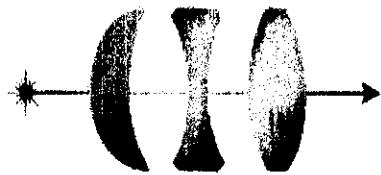
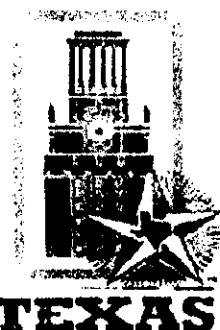
CIO - Léon



- **SHG spectroscopy directly probes charge transfer accompanying**
 - *buckling of surface dimers*
 - *Ge chemisorption*
 - *H passivation*
 - *2nd layer B incorporation*
- **The SHG signatures for each surface-adsorbate system are UNIQUE**
- **The SHG signatures have been microscopically modelled by**
 - *SETB*
 - *first principles calculations*
- *New surface spectroscopic method that exploits 10 fs lasers: FDISH*
- *EFISH microscope images dc fields with submicron resolution.*
- *SHG probes Si nanoparticles embedded in SiO₂.*

D. Lim *et al.*, Phys. Rev. Lett. **84**, 3406 (2000).

P. T. Wilson *et al.*, Opt. Lett. (1999).



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CONCLUSIONS

- SHG spectroscopy directly probes charge transfer accompanying
 - buckling of surface dimers
 - Ge chemisorption
 - H passivation
 - 2nd layer B incorporation
- at Si(001).
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- The SHG signatures have been microscopically modelled by
 - SETB
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