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*Growth and surface reaction kinetics in thin film
epitaxy studied by oblique-incidence reflectivity
difference*

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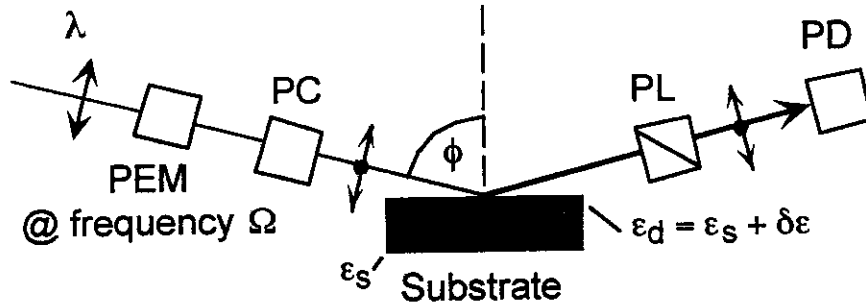
**Growth and surface reaction kinetics
in thin film epitaxy
studied by
*oblique-incidence reflectivity
difference***

X.D. Zhu
University of California at Davis

- **Oblique-incidence reflectivity difference**
- **Growth kinetics in rare gas epitaxy on metals**
 - Xe on Ni(111): commensurate substrate
 - Xe on Nb(110): incommensurate substrate
- **Thermal annealing of ion-sputtered surfaces**
 - Ne⁺ sputtering and thermal annealing of Ni(111)
- **Surface oxidation kinetics in oxide epitaxy**
 - SrTiO₃ and La_{0.67}Ba_{0.33}MnO_{3-δ} on SrTiO₃(001)

(Ariatico Research Conference on Lasers in Surface Science, 2000)

Oblique-incidence reflectivity difference technique



$$\Delta_s \equiv (r_s - r_{s0})/r_{s0}$$

$$\Delta_p \equiv (r_p - r_{p0})/r_{p0}$$

$$r_p = |r_p| e^{i\phi_p} \approx |r_{p0}| e^{i\phi_{p0}} (1 + \text{Re}\Delta_p) e^{i\text{Im}\Delta_p}$$

$$r_s = |r_s| e^{i\phi_s} \approx |r_{s0}| e^{i\phi_{s0}} (1 + \text{Re}\Delta_s) e^{i\text{Im}\Delta_s}$$

$$\Delta_p - \Delta_s \approx i \left\{ \frac{4\pi(\cos\phi)(\sin\phi)^2 \epsilon_s}{\lambda [\epsilon_s^2 (\cos\phi)^2 - \epsilon_s + (\sin\phi)^2]} \right\} \int_0^d \frac{(\epsilon_d - \epsilon_s)(\epsilon_d - 1)}{\epsilon_d} dz$$

Reflected light intensity at 2Ω : (setting $|r_{p0} \cos\theta_{PL}| = |r_{s0} \sin\theta_{PL}|$)

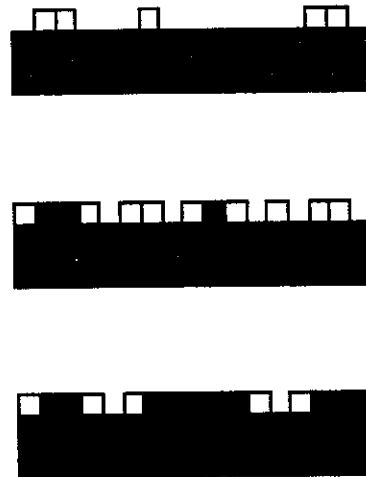
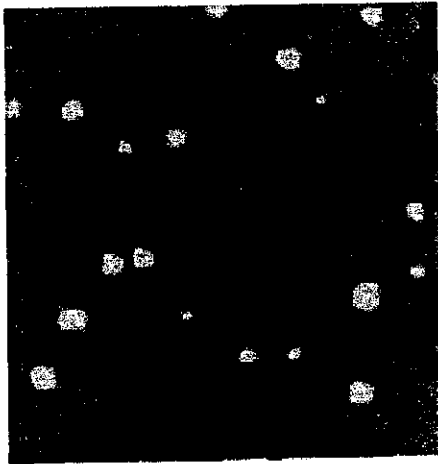
$$I_{2\Omega} \approx I_{inc} |r_{p0} \cos\theta_{PL}|^2 \text{Re}\{\Delta_p - \Delta_s\} \sim \text{Re}\{\Delta_p - \Delta_s\}$$

Reflected light intensity at Ω : (setting $\phi_{PC} + (\phi_{p0} - \phi_{s0}) = 0$)

$$I_{\Omega} \approx I_{inc} |r_{p0} \cos\theta_{PL}|^2 \text{Im}\{\Delta_p - \Delta_s\} \sim \text{Im}\{\Delta_p - \Delta_s\}$$

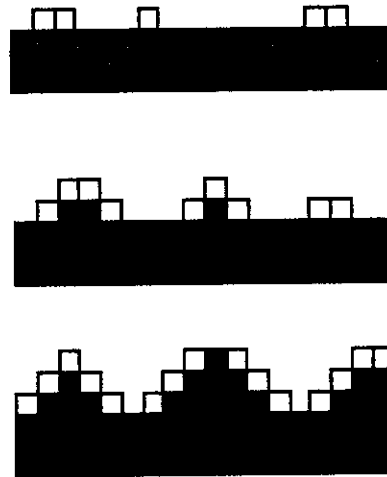
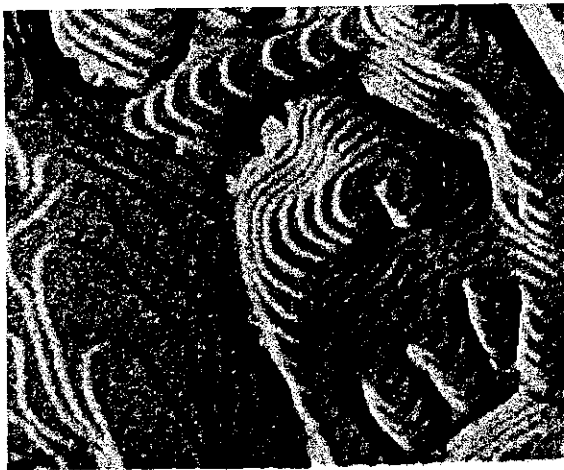
Growth kinetics in rare gas epitaxy on metals

Layer growth (Frank-van der Merwe mode):



ϵ_d changes periodically.

Island growth (Volmer-Weber mode):



ϵ_d changes monotonically.

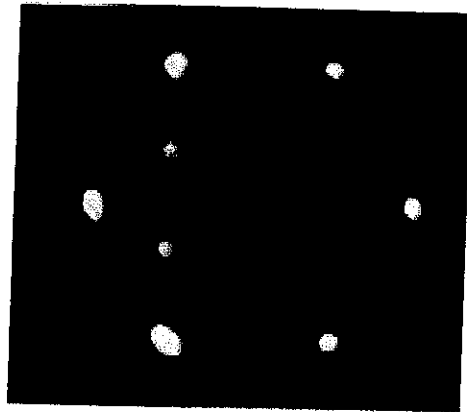
Step-flow growth:

ϵ_d remains constant.

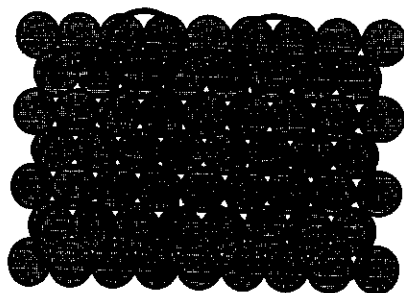
Rare gas growth on commensurate substrate

Xe/Xe(111)/Ni(111)

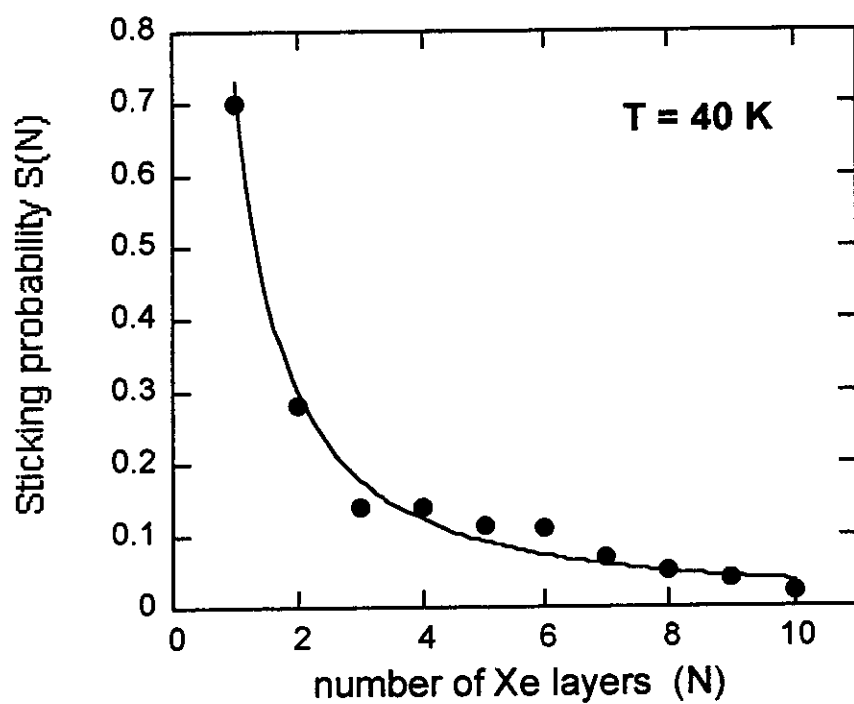
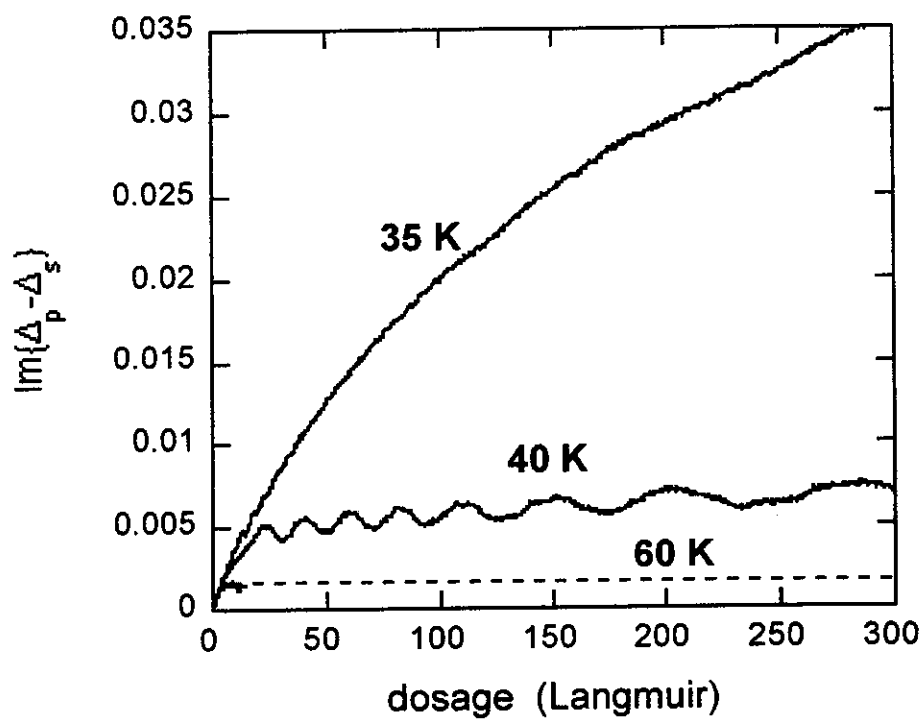
LEED pattern of a commensurate $(\sqrt{3} \times \sqrt{3})R30^\circ$ structure



Real-space structure of a commensurate $(\sqrt{3} \times \sqrt{3})R30^\circ$ structure

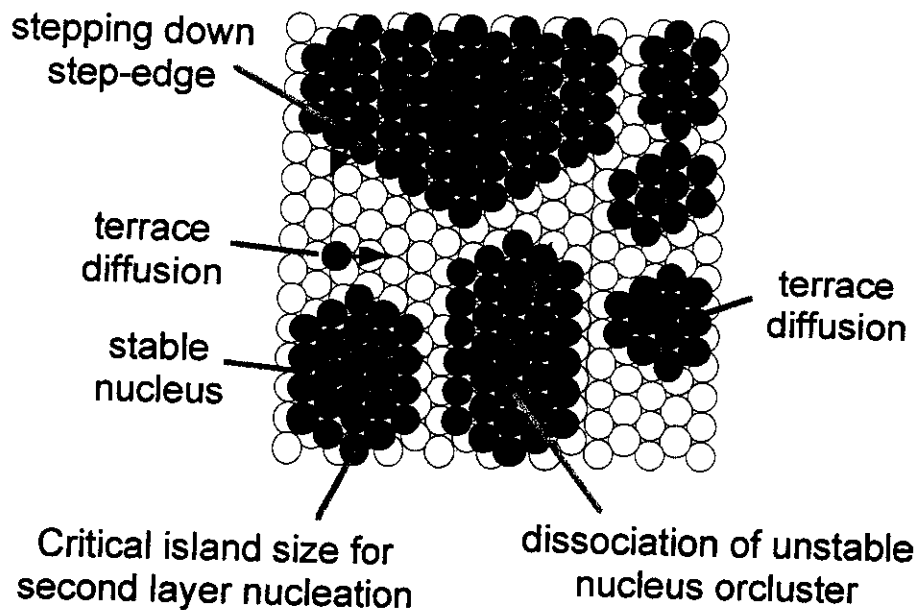


Growth of Xe multilayers on Xe(111)/Ni(111):



Transition from multilayer growth to layer-by-layer growth:

- Rate theory of nucleation and growth and kinetic Monte Carlo (KMC) simulation



Stepping down a step edge:

$$h' = v_{SE} \exp\left[-(E_{SE} + E_d)/k_B T\right]$$

Terrace diffusion:

$$h = v_d \exp(-E_d/k_B T)$$

Critical size i for stable nuclei:

2-D nuclei with $i+1$ atoms stable during deposition

Critical island size R_c for second layer nucleation:

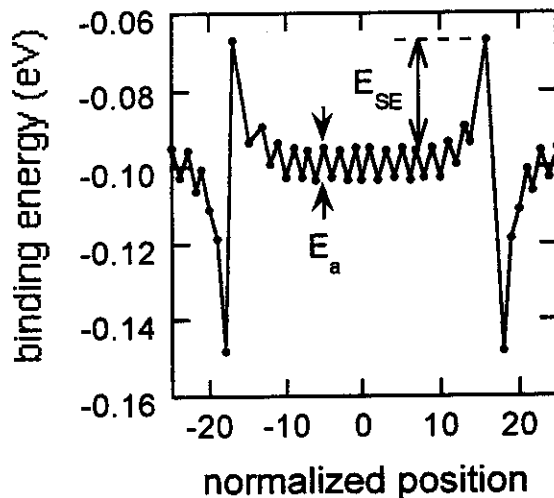
Probability of forming a stable nucleus on 2-D islands of radius R_c is 50%.

Nucleation theory and KMC:

Transition in i and R_c

Binding energies of Xe on Xe(111):

Xe binding energy across a Xe island



Terrace diffusion barrier: $E_d = 8 \text{ meV}$

Extra step edge barrier: $E_{SE} = 28 \text{ meV}$ (Schwoebel-Ehrlich)

Binding energies of Xe 2-D clusters on Xe(111):

Single Xe-Xe bond: $E_1 = 20 \text{ meV}$

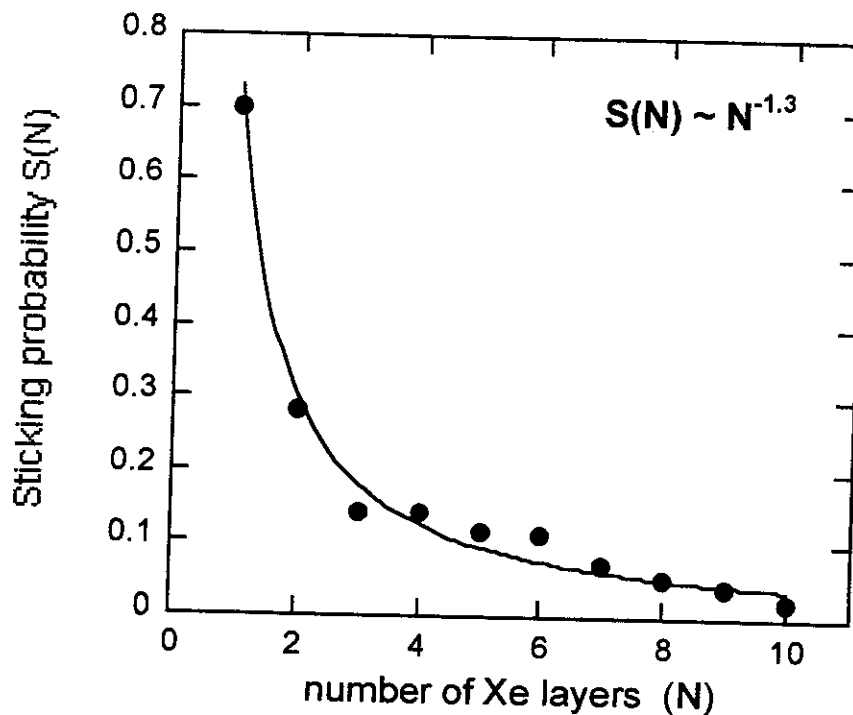
Double Xe-Xe bond: $E_2 = 40 \text{ meV}$

Triple Xe-Xe bond: $E_3 = 60 \text{ meV}$

Kinetics Monte Carlo simulation and nucleation theory:

- (1) For $23 \text{ K} \leq T \leq 40 \text{ K}$, only Xe septamers or larger compact islands ($i \geq 6$) are stable on Xe(111) against dissociation.
- (2) From 35 to 40K, average island size R_{av} before coalescence remains unchanged.
- (3) Critical radius R_c increases by one order of magnitude from 35 to 40 K, causes the transition from multilayer growth at 35 K to layer-by-layer growth at 40 K.

Thickness dependence of Xe sticking probability:

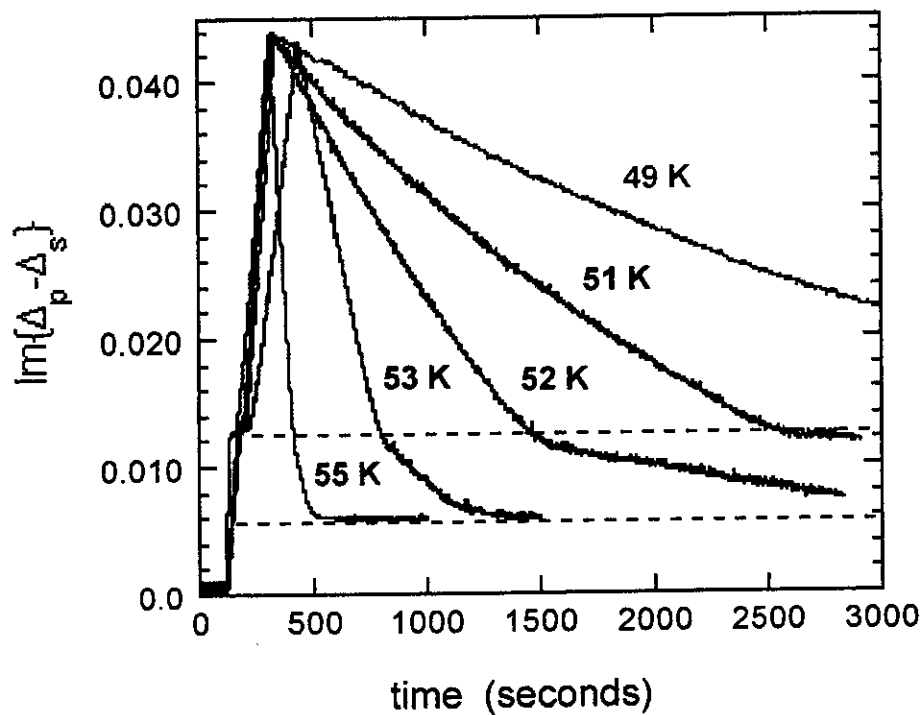
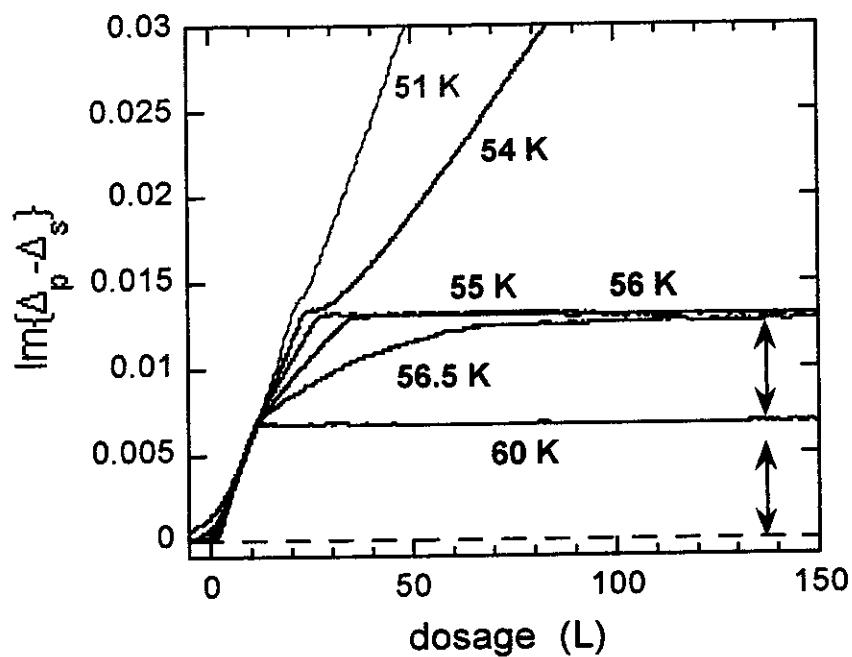


- (1) Heavier Xe ($M_{\text{Xe}}/M_{\text{Ni}} = 2.4$) pushes lighter Ni surface atoms forward without being bounced kinematically and thus the interaction time is longer;
- (2) Xe binding energy is larger near metallic Ni(111) due to additional polarization effects; it is the largest on Ni(111) due to charge transfer.

Both favor larger sticking probability on thinner Xe films.

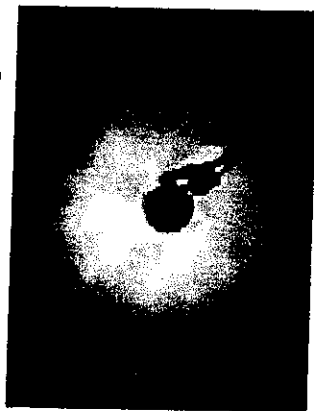
Growth of Xe multilayers on incommensurate Nb(110)

0.1°-miscut, $L_1 \approx 1300\text{\AA}$





After flashing



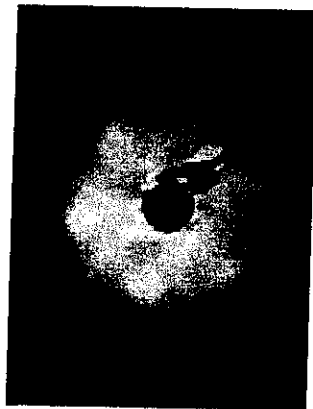
40 L



140 L



10 L



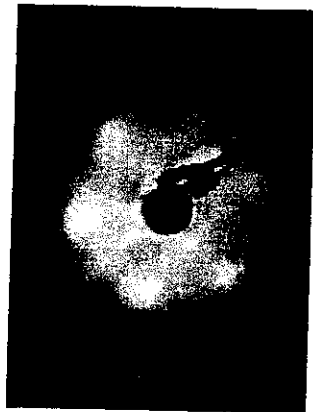
60 L



180 L



20 L

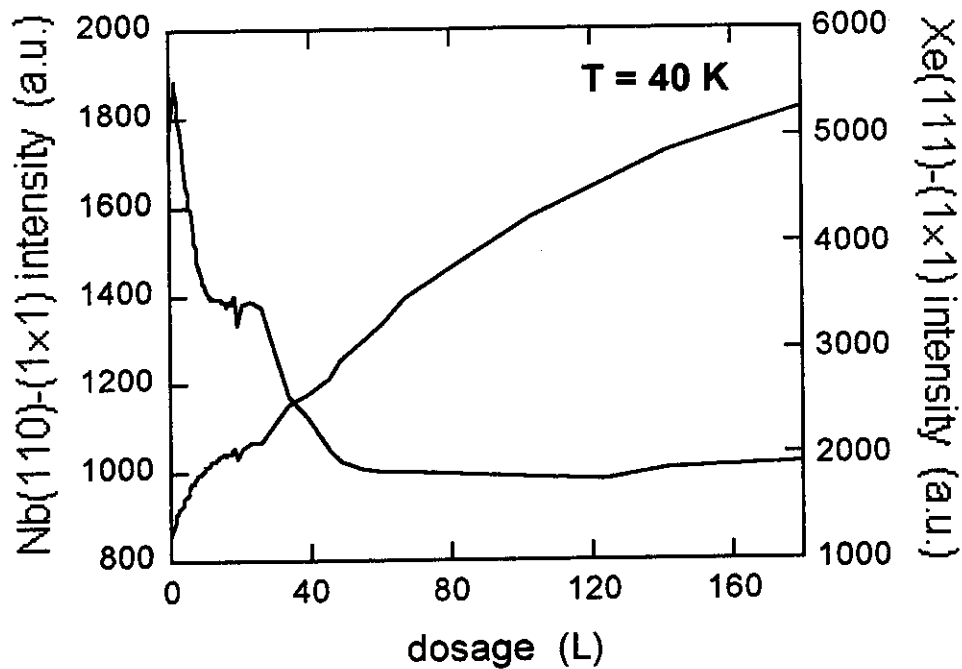


100 L



280 L

Step-flow growth of Xe on 2-ML "buffer" Xe on Nb(110)



Calculated reflectivity difference for step-flow growth:

$$\varepsilon_d = \varepsilon_{\text{Xe}} = 2.0$$

$$\varepsilon_s = \varepsilon_{\text{Nb}} = -0.25 + i 16.13$$

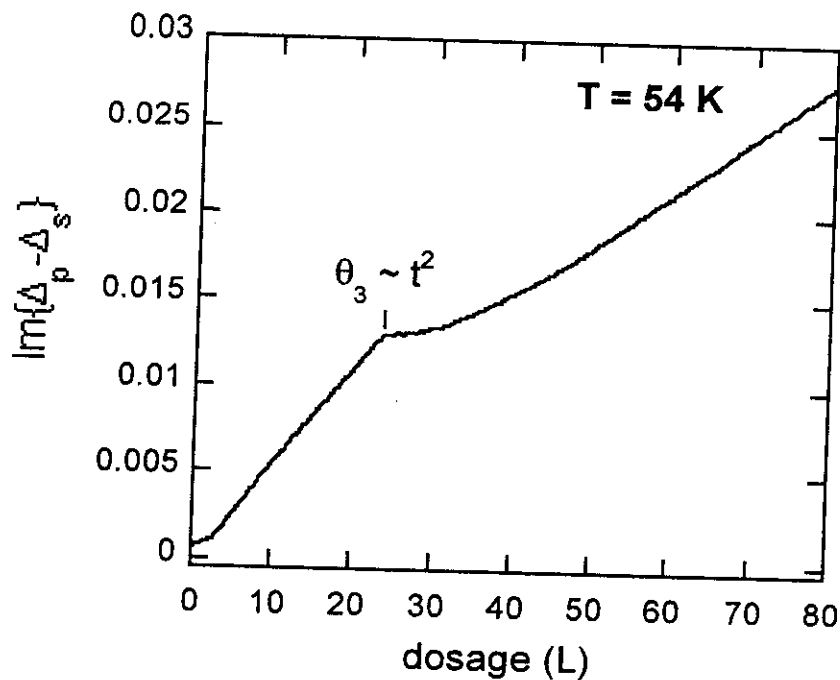
$$\Delta_p - \Delta_s \approx i \left\{ \frac{4\pi \cos \phi \sin^2 \phi \varepsilon_s}{\lambda [\varepsilon_s^2 \cos^2 \phi - \varepsilon_s + \sin^2 \phi]} \frac{(\varepsilon_d - \varepsilon_s)(\varepsilon_d - 1)}{\varepsilon_d} \right\} d$$

$$\approx 5.3 \times 10^{-3} (1 - i 0.12) @ d = 3.54 \text{ \AA}$$

Measured reflectivity difference: (40 K $\leq$ T $\leq$ 51 K)

$$-\frac{\text{Im}\{\Delta_p - \Delta_s\}}{\text{Re}\{\Delta_p - \Delta_s\}} = 0.10 \sim 0.12$$

Adsorption of the 3rd Xe monolayer on 2-ML Xe/Nb(110)
adsorption at edges of stable 2-D clusters



t^2 -dependence: $d\theta_3/dt \sim \sqrt{\theta_3}$

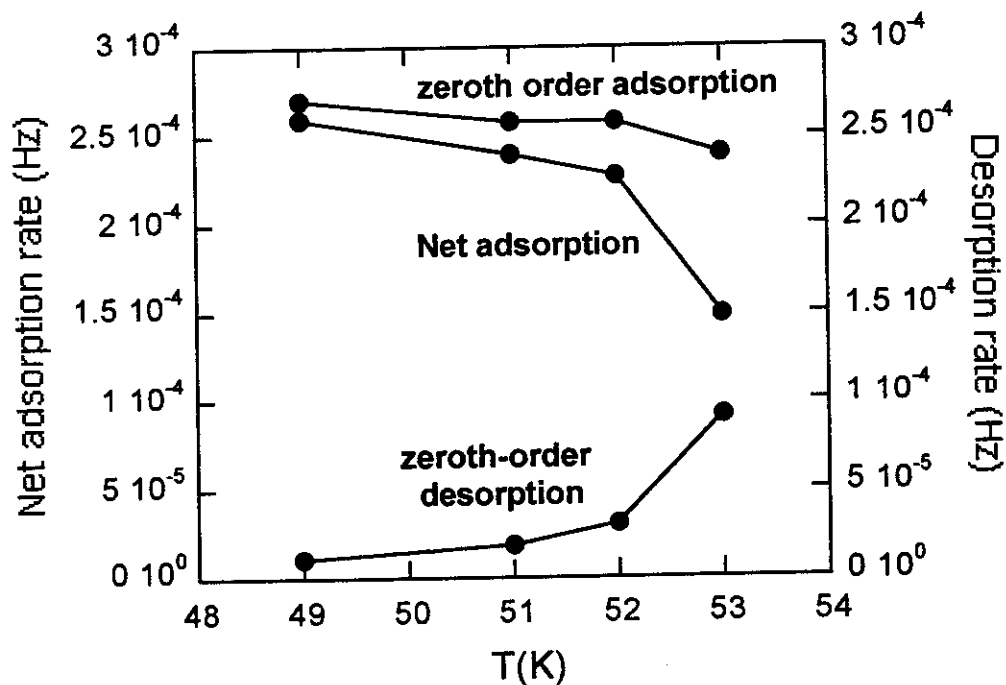
- Xe monomers on 2-ML Xe/Nb(110) are short-lived
- Initially, $d\theta_3/dt \sim$ rate of septamer formation
- Adsorption at edges of septamers or larger clusters dominates

$$d\theta_3/dt \sim \sqrt{\theta_3} \quad \theta_3(t) \sim t^2$$

- Desorption of subsequent Xe layers is rate-limited by detachment from step edges (reversed step-flow)

Zeroth-order desorption and adsorption for overlayers ($n \geq 4$)

Competition between step-flow growth and "step-flow" desorption

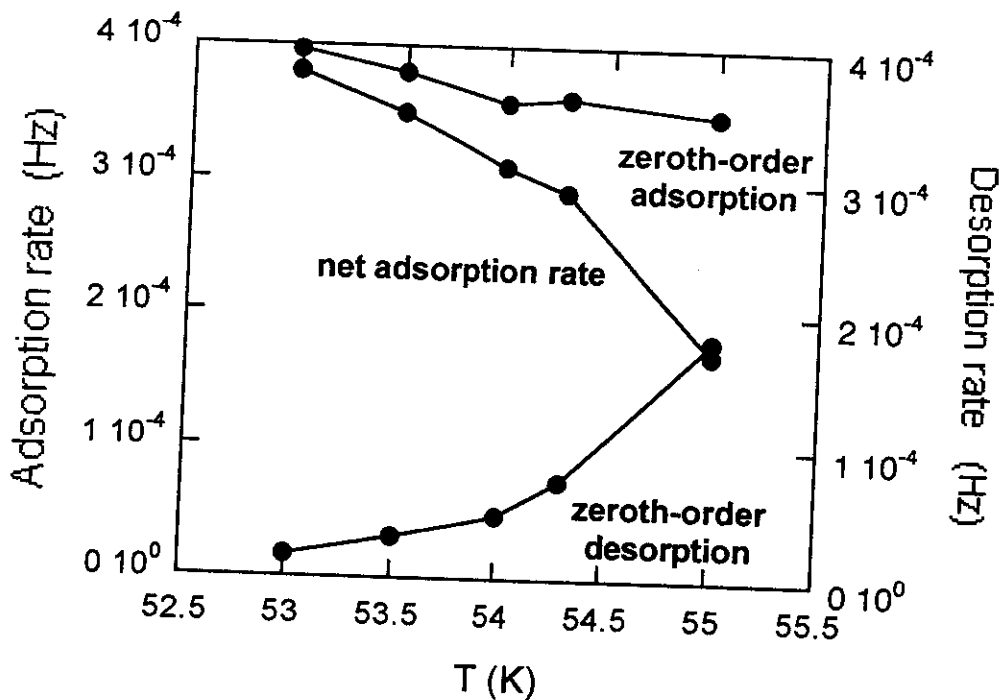


$$dN/dt = S(N)\Gamma - N_{s,Xe} v(T, N), \quad N: \text{number of atoms / cm}^2$$

$v(T, N)$ independent of N : rate-limited by detachment from step edges (reverse step-flow)

$S(N) \approx S_0 \approx 1$: fast diffusion on Xe(111) terraces toward step edges and efficient capture by step edges

Zeroth-order desorption and adsorption for 2nd Xe overlayer

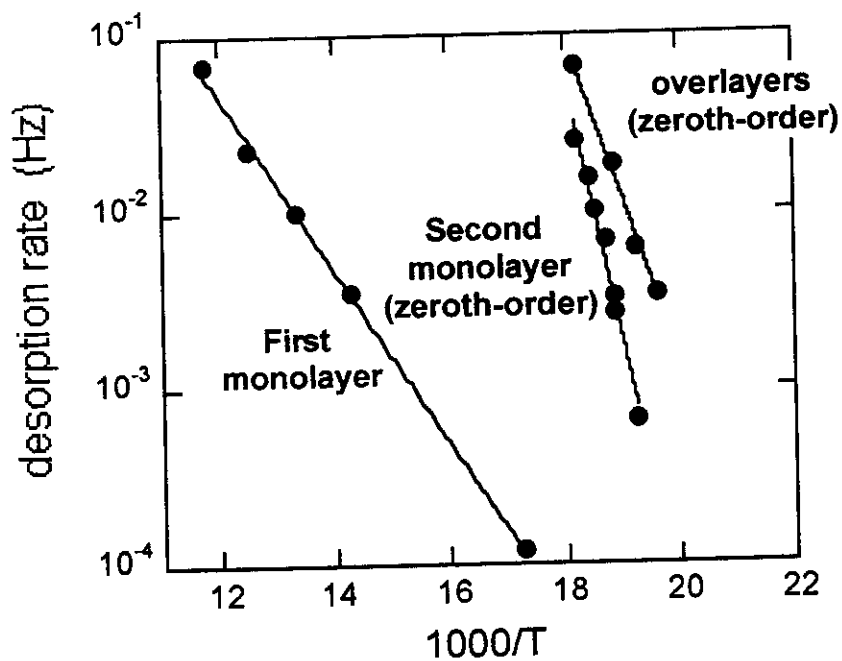


$$dN/dt = S(N)\Gamma - N_{s,Xe} \nu(T, N), \quad N: \text{number of atoms / cm}^2$$

$\nu(T, N)$ weak N -dependence: rate-limited by detachments from island edges or corners, and minorly from terraces directly

$S(N) \approx S_0 \approx 1$: efficient accommodation on terraces of the first Xe ML.

Desorption kinetics of different Xe layers on Nb(110)



- **First monolayer**

$$\nu_1(T) = 4 \times 10^4 \exp(-0.1\theta V/k_B T)$$

- **Second monolayer**

$$\nu_2(T) = 2 \times 10^{26} \exp(-0.3\theta V/k_B T)$$

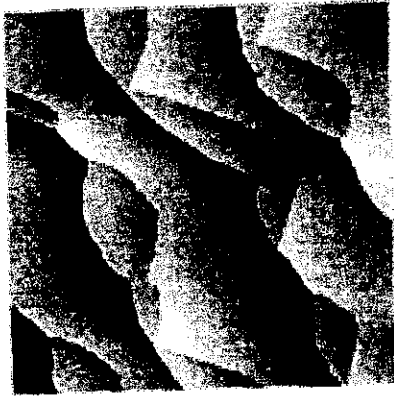
- **Subsequent layers**

$$\nu_{n \geq 3}(T) = 5 \times 10^{15} \exp(-0.19\theta V/k_B T)$$

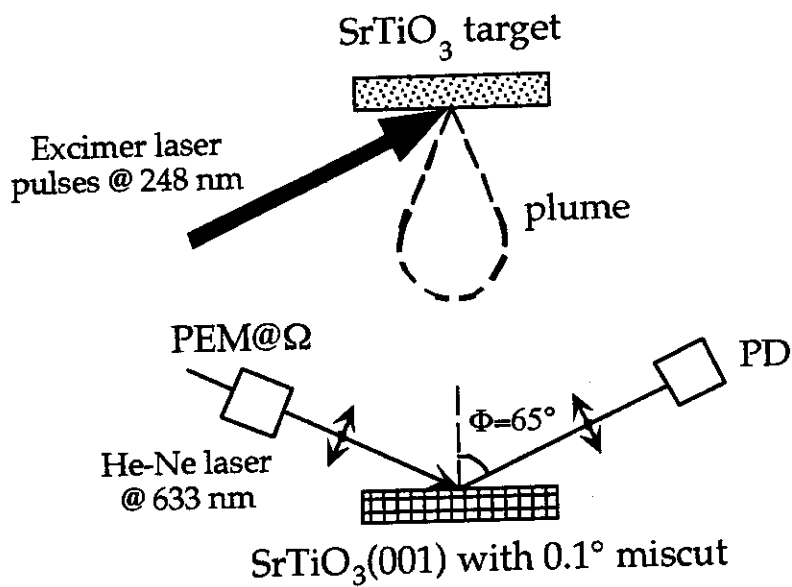
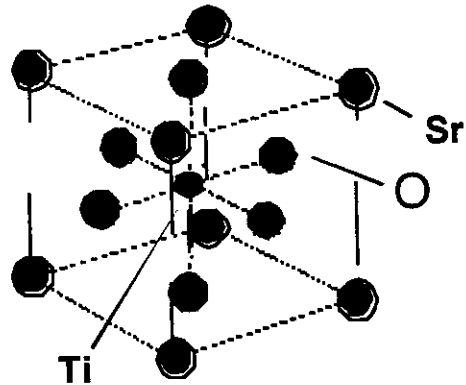
Surface oxidation kinetics in oxide epitaxy

Surface oxidation kinetics in perovskite oxide epitaxy

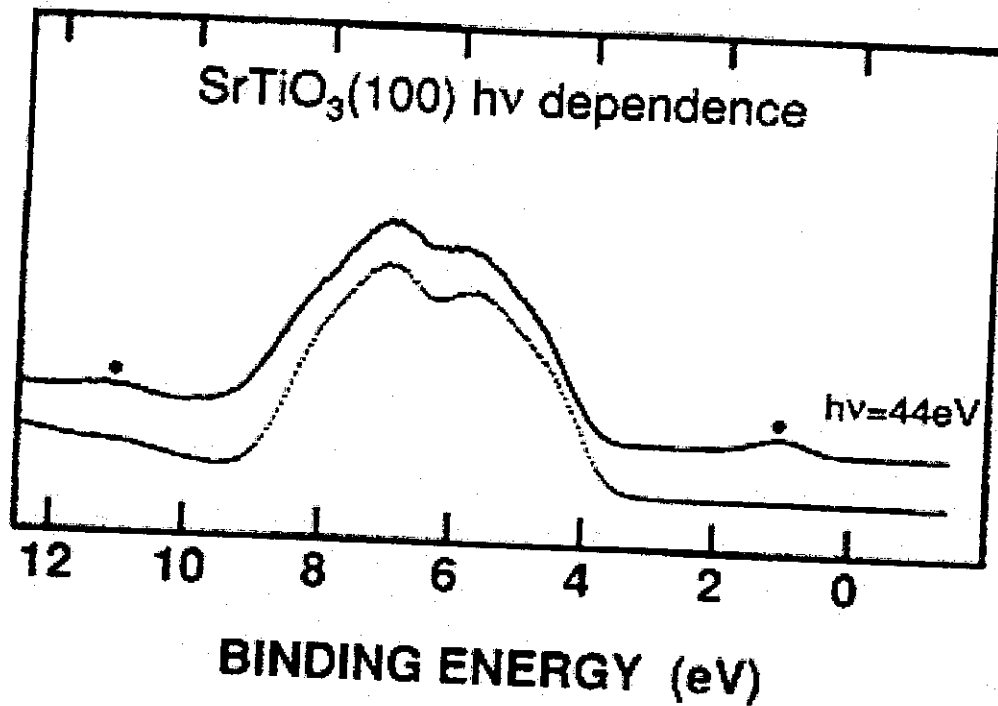
SrTiO₃ on SrTiO₃(001)



1 μm × 1 μm



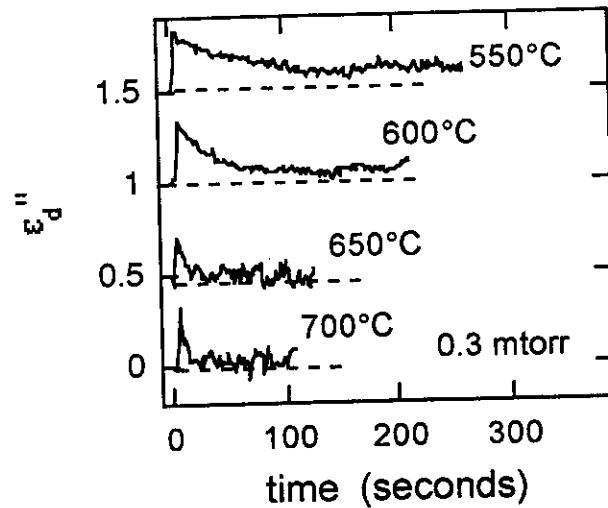
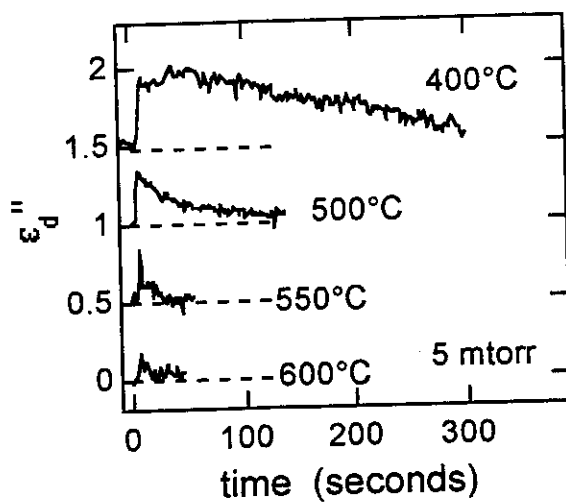
STO Optical constant:	$\epsilon_s = 5.66 + i 0$ @ 633 nm
Temperature:	270 ~ 800 °C (550 ~ 1070 K)
Oxygen pressure:	3×10^{-4} ~ 0.6 torr
PLD rate:	0.3 ~ 0.5 ML/sec



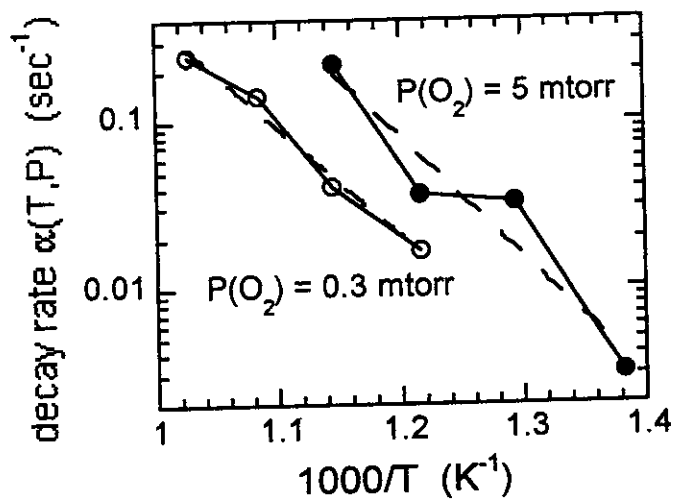
- Oxygen vacancies in TiO₂ planes induce "defect states" at 1.2 eV below the Fermi surface
- Defect states cause optical absorption at $\lambda = 633$ nm with $h\nu = 1.9$ eV
- $\epsilon_d \approx \epsilon_s + \delta\epsilon_d' + i\delta\epsilon_d''$ with $\delta\epsilon_d''$ proportional to oxygen vacancy density in the deposited monolayer

$$\delta\epsilon_d = \delta\epsilon_d' + i\delta\epsilon_d'' \approx (-i) \left\{ \frac{\lambda [\epsilon_s^2 \cos^2 \phi - \epsilon_s + \sin^2 \phi]}{4\pi d \cos \phi \sin^2 \phi (\epsilon_s - 1)} \right\} (\Delta_p - \Delta_s)$$

Optical reflectivity difference measurements ($I_{2\Omega}$):



$$\epsilon_d'' = \delta \epsilon_d'' \sim \exp[-\kappa(T, P_{O_2})t]$$



Oxidation reaction parameters:

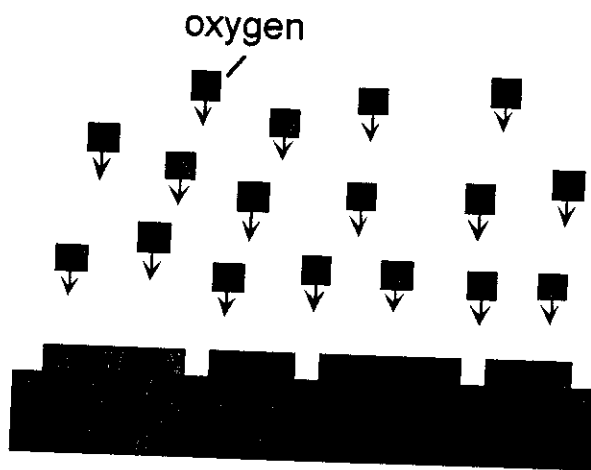
$$\kappa(T, P_{O_2}) = \alpha_0 \exp(-E_a/k_B T) P_{O_2}$$

$$E_a = 1.35 \text{ eV}$$

$$\alpha_0 = 4 \times 10^9 \text{ torr}^{-1} \text{ sec}^{-1}$$

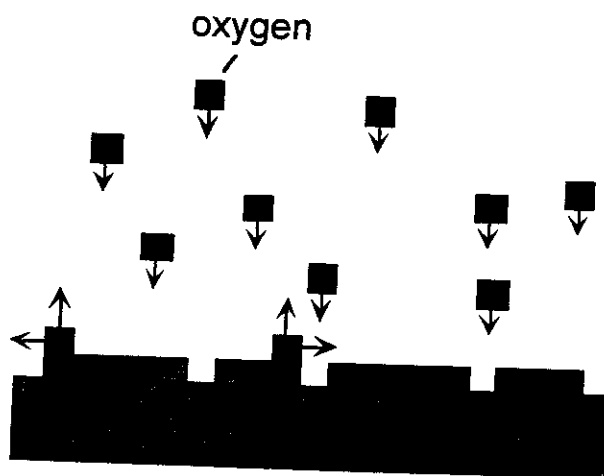
Oxidation reaction in the homoepitaxy of $\text{SrTiO}_3(001)$

Growth of monolayer oxide in pulsed-laser deposition:



ϵ_d deviates from ϵ_s due to incomplete oxidation.

Oxidation reaction during the post-growth treatment:



ϵ_d restores to ϵ_s as the oxidation reaction is completed.

Oxidation reaction parameters:

$$\kappa(T, P_{O_2}) = \alpha_0 \exp(-E_a/k_B T) P_{O_2}$$

$$E_a = 1.35 \text{ eV}$$

$$\alpha_0 = 4 \times 10^9 \text{ torr}^{-1} \text{ sec}^{-1}$$

- (1) Molecular oxygen dissociate near oxygen vacancies (N_v) into oxygen adatoms in a precursor state (N_i);
(T-dependence)
- (2) Precursor-state oxygen either desorb (E_{des}) or react with nearby oxygen vacancies (E_R);

$$\frac{dN_i}{dt} = S\Gamma_p \left[\frac{(N_{v0} - N_v) - N_i}{N_s} \right] - \nu_{des} N_i \exp(-E_{des}/k_B T) - \nu_R N_i \exp(-E_R/k_B T)$$

$$\frac{dN_v}{dt} = -\nu_R N_i \exp(-E_R/k_B T)$$

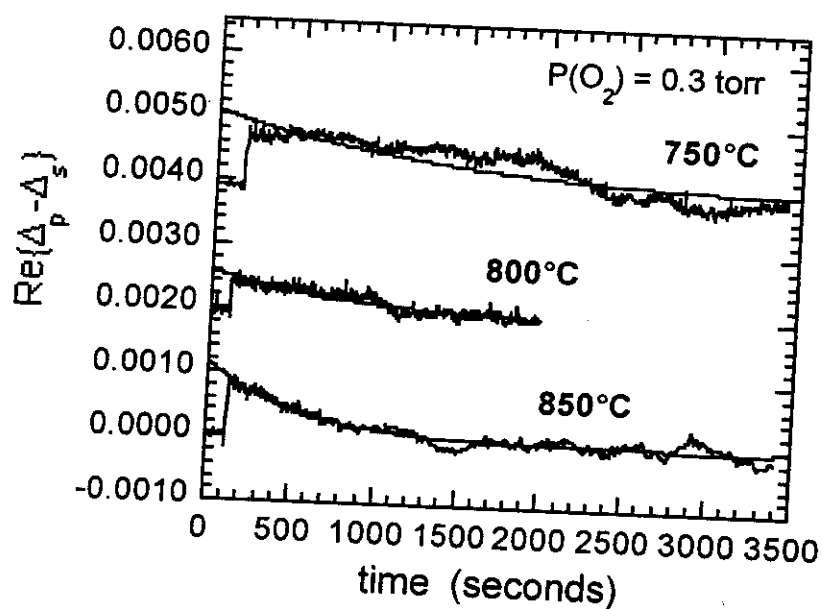
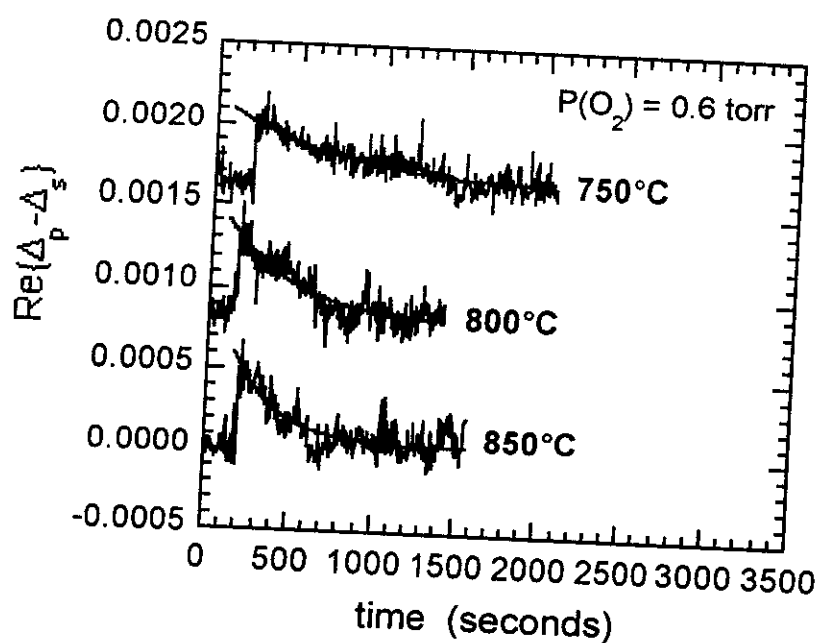
- (3) Desorption rate is large compared to reaction rate such that $N_i \approx [S\Gamma_p(N_{v0} - N_v)/N_s \nu_{des}] \exp(E_{des}/k_B T)$. Consequently,
(P-dependence)

$$N_v \approx N_{v0} \exp[-\alpha(P_{O_2}, T)t]$$

$$\alpha(P_{O_2}, T) = (S\Gamma_p/N_s)(\nu_R/\nu_{des}) \exp[-(E_R - E_{des})/k_B T]$$

Activation energy barrier E_a : $E_a = E_R - E_{des}$

La_{0.67}Ba_{0.33}MnO₃ (LBMO) heteroepitaxy on SrTiO₃(001):



Oxidation reaction parameters: $\kappa(T, P_{\text{O}_2}) = \alpha_0 \exp(-E_a/k_B T) P_{\text{O}_2}$

$$E_a = 1.0 \text{ eV}$$

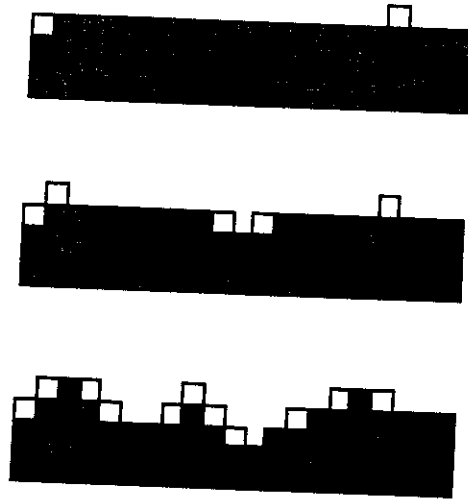
$$\alpha_0 = 180 \text{ torr}^{-1} \text{ sec}^{-1}$$

Thermal annealing of ion-sputtered metal surfaces

Ion sputtering and thermal annealing of a crystalline surface:

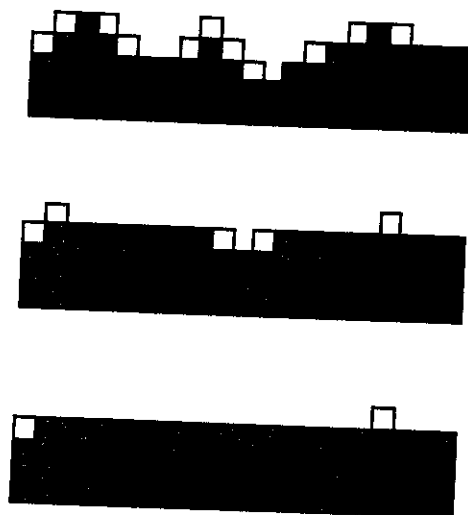
ϵ_d is proportional to the edge atom density

Sputtering:



ϵ_d changes from ϵ_s monotonically.

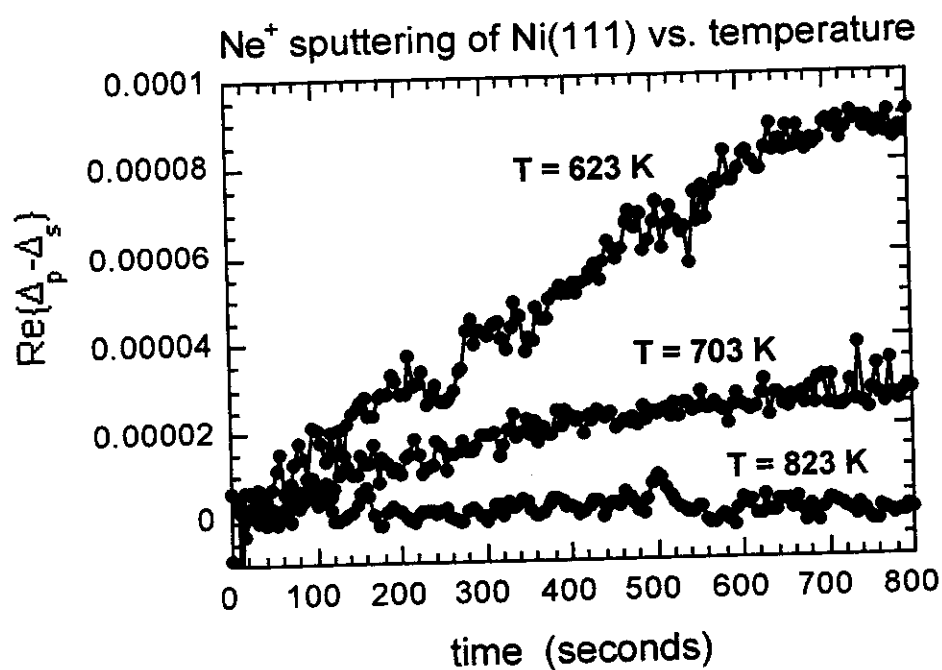
Thermal annealing:



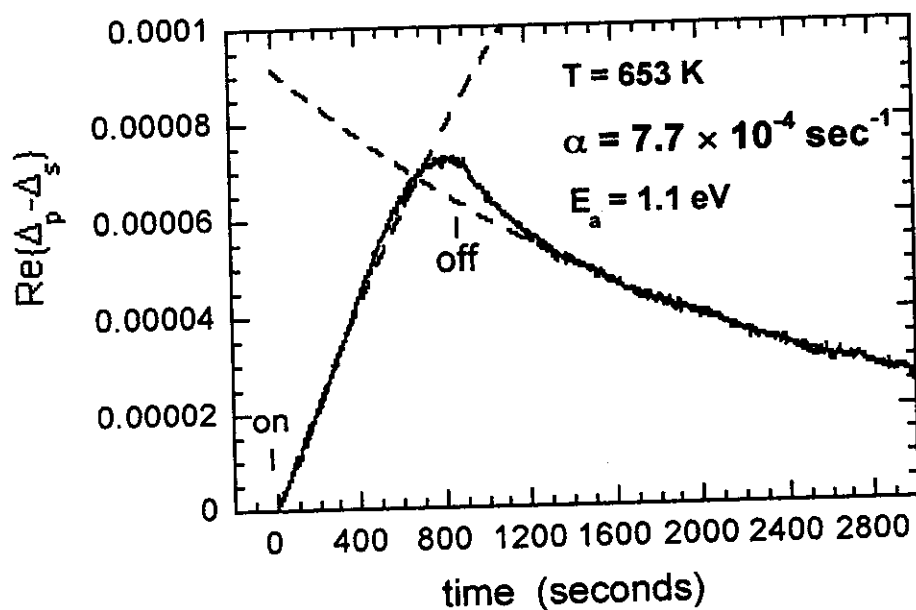
ϵ_d restores to ϵ_s monotonically.

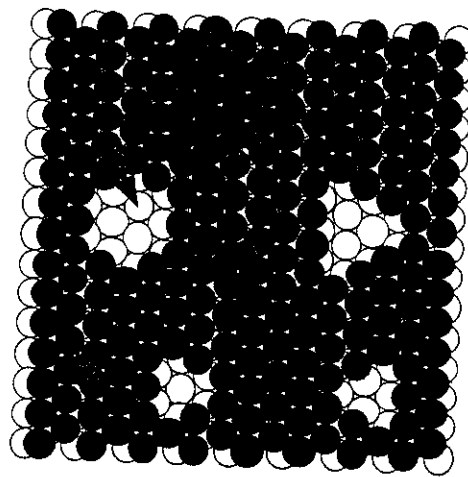
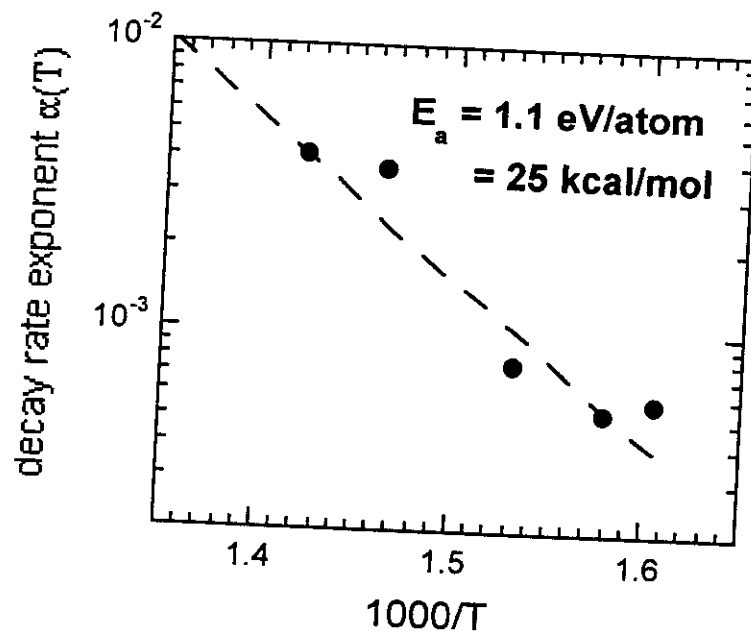
Ne⁺ ion sputtering and thermal annealing of Ni(111):

Beam energy = 1000 eV, Beam current ~ 2 μ A



Ne⁺ sputtering and thermal annealing of Ni(111)





Rate-limiting processes to the thermal annealing of Ni(111):

- Diffusion on Ni(111) terraces: E_{diff}
- **Stepping down B-type edges:** $E_{\text{se}} = E_{\text{step}} - E_{\text{diff}}$
- Evaporation from corners: E_{c}
- **Evaporation from kinks:** E_{kink}

$$E_a \approx E_{\text{kink}} + E_{\text{diff}} + E_{\text{se}} \quad (\sim 1 \text{ eV for Cu on Cu(111)})$$

[Icking-Konert et al., Surf. Sci (1998); Giesen et al., Surf. Sci. (1999); Feibelman, Phys. Rev. Lett. (1998)]

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