
SCHOOL ON SYNCHROTRON RADIATION

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*The application of multilayer coatings
in synchrotron radiation research*

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School on Synchrotron Radiation
ICTP, November 22nd, 2000

The application of multilayer coatings in synchrotron radiation research

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I will try to give answers to the following questions:

- a) Why do we need them?
- b) How do we make them?
- c) Can we simulate their performance?
- d) How do we validate them?
- e) Do experiment and simulation really agree?
- f) Where do we finally use them?

We all know that using visible light

(red = $\lambda = 800 \text{ nm} = 1.5 \text{ eV}$;

violet = $\lambda = 400 \text{ nm} = 3 \text{ eV}$)

- we can take an image with a camera by use of lenses or mirror optics!

- we can make an enlarged image in a microscope once more using either lenses or less often mirror systems!

- we can make simple polarizers or phase retarders by use of birefringent material!

What happens if we want to use these devices in the soft x-ray range (10 eV - 2000 eV)?

With increasing photon energy the light starts

- a) to be absorbed in lenses
- b) to not be dispersed anymore in lenses
- c) to find little anisotropy in lenses/filters
- d) to not be reflected anymore in normal incidence

Why this:

let's take some examples:

n = index of refraction

a) focal length of lens: $(1/f) = (n-1) ((1/r_1) + (1/r_2))$ $r_1, r_2 = \text{radii}$

b) normal incidence reflectivity: $R = ((n-1)/(n+1))^2$

The index of refraction varies as follows for glass (SiO₂):

photon energy:	3 eV	30 eV	100 eV	1000 eV
$n =$	1.5	0.9	0.985	0.9987
$\lambda =$	e.g. 0.2 m	- 1 m	- 6.6 m	-80 m
$R =$	0.04	0.0028	$6 \cdot 10^{-5}$	$4 \cdot 10^{-7}$
$R_{\text{int}} =$	0.37	0.08	0.0014	$1 \cdot 10^{-6}$

So what can we do?

a) Not much as far as classical lenses are concerned!

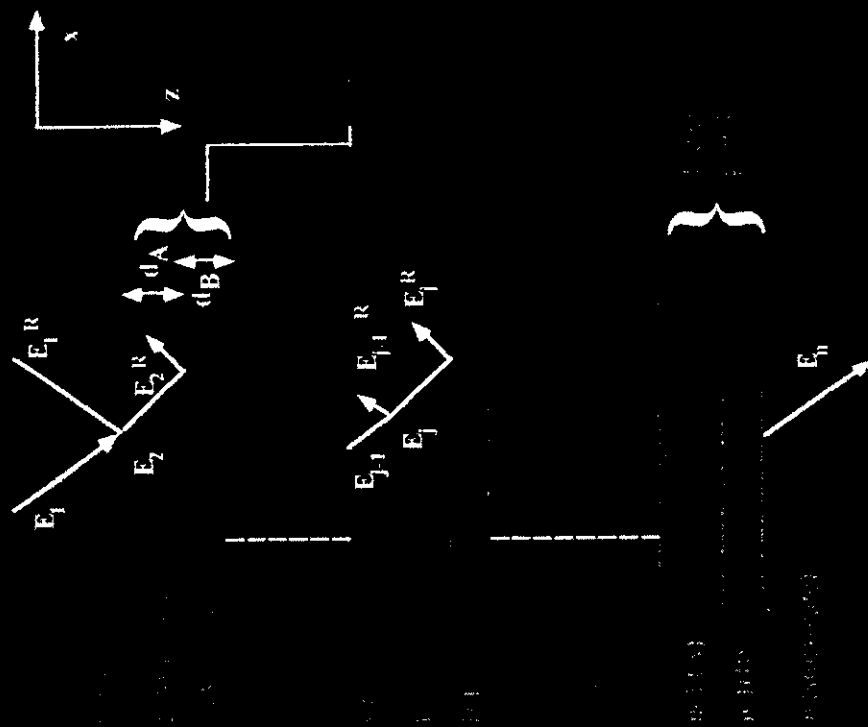
b) However, for m interfaces in a sufficiently transparent structure we can get $R_{\text{total}} = m R_{\text{int}}$, which becomes interesting for heavier materials with better R_{int} than glass.

But is it really so simple?

Not really, the different waves need to collaborate in phase, which is affected by the index of refraction n of a layer and which can undergo additional changes at any reflexion at and any transmission through an interface.

So let us write a program!

E. M. field propagation in film systems:



Solution of Parratt (Phys. Rev. 95, 359 (1954))

$E_{\text{lang}} = \text{const}$ at interface

$$a_j = \text{amplitude factor} \quad a_j = \exp\left(-i\pi \frac{g_j d_j}{\lambda}\right)$$

$d_j = \text{thickness}, \quad \lambda = \text{wavelength},$

$$g_j = \tilde{n}_j \sin \theta_j \quad \tilde{\epsilon}_j = \tilde{n}_j^2 \quad \tilde{n}_j = 1 - \delta + i\beta$$

$$a_j E_j + a_j^{-1} E_j^R = a_{j+1}^{-1} E_{j+1} + a_{j+1} E_{j+1}^R$$

$$g_j (a_j E_j - a_j^{-1} E_j^R) = g_{j+1} (a_{j+1} E_{j+1} - a_{j+1} E_{j+1}^R)$$

Solution $\Rightarrow$ Recursion equation

$$R_{j,j+1} = a_j^4 \left[\frac{R_{j+1,j+2} + J_{j,j+1}}{R_{j+1,j+2} J_{j,j+1} + 1} \right] \quad R = \frac{I}{I_0} = |R_2|^2$$

$$J_{j,j+1}(s - \text{pol}) = \frac{E_j^R}{E_j} \frac{g_j - g_{j+1}}{g_j + g_{j+1}}$$

$$R_{j,j+1} = a_j^2 \frac{E_j^R}{E_j} \quad J_{j,j+1}(p - \text{pol}) = \frac{g_j / \tilde{n}_j^2 - g_{j+1} / \tilde{n}_{j+1}^2}{g_j / \tilde{n}_j^2 + g_{j+1} / \tilde{n}_{j+1}^2}$$

Do not be afraid!

The job is done already by the Center for X-ray Optics at LBNL at Berkeley, CA (USA):

at

http://www-cxro.lbl.gov/optical_constants

and at

http://www-cxro.lbl.gov/optical_constants/multi2.html

And if you want to analyse data, ask the ESRF for the XOP software package, which contains the program IMD of David Windt.

Production:

is it more difficult? Are there any problems?

NO!!

Evaporation:

- repeatability
- cluster evaporation
- sublimation
- monitoring of thickness
- homogeneity over large sample sizes

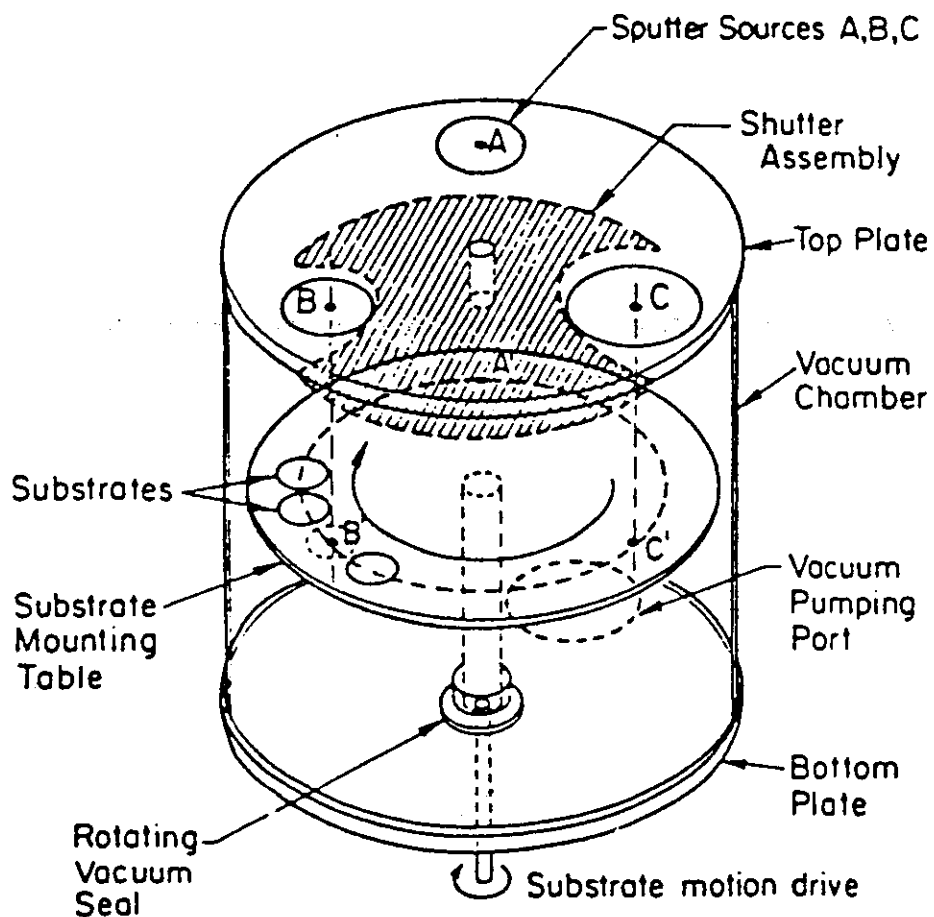
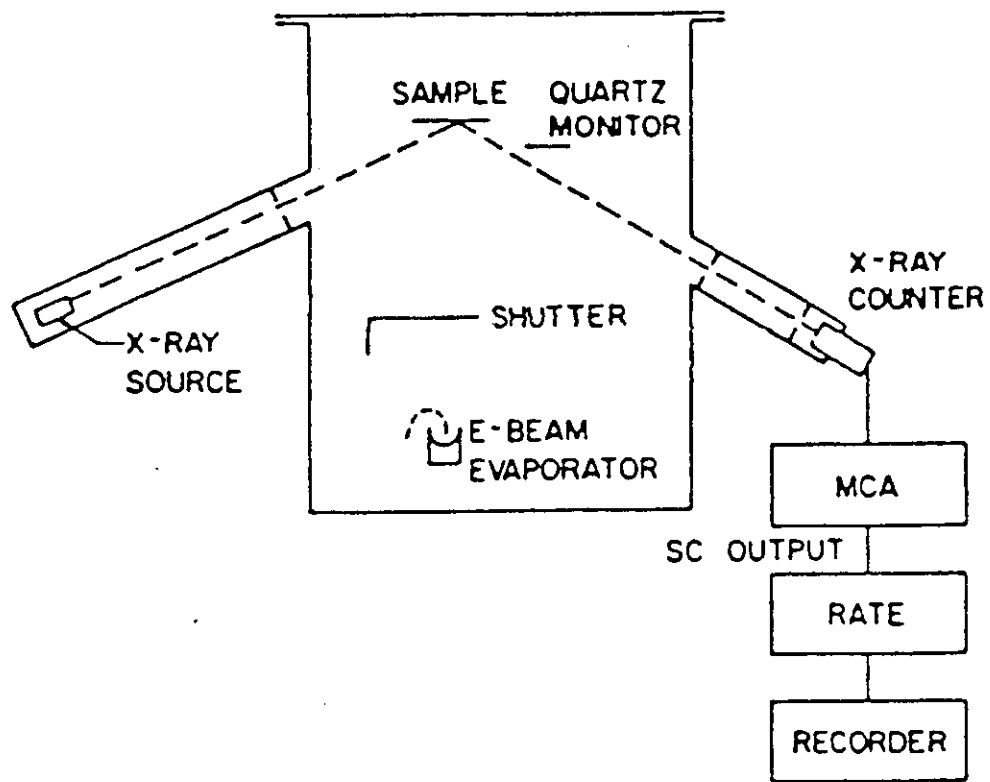
Sputtering:

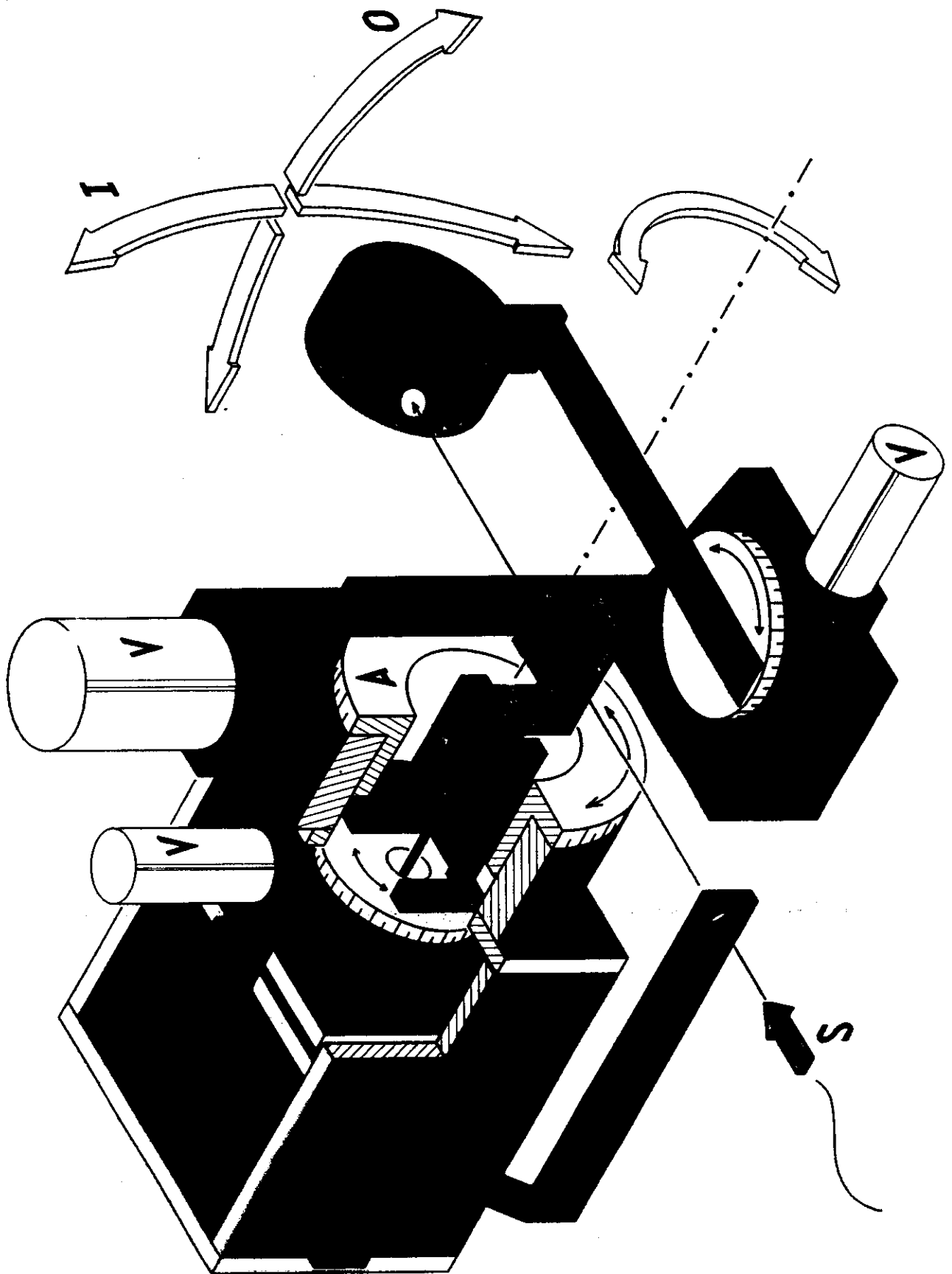
- homogeneity over large sample sizes
- simultaneous bombardment of sample with plasma electrons
- plasma gas inclusions ($p \approx 1 - 100 \text{ mbar}$)
- reactions with restgas of vacuum

but: easy monitoring and very repeatable due to plasma stability and reproducibility

⇒ monitoring by timing

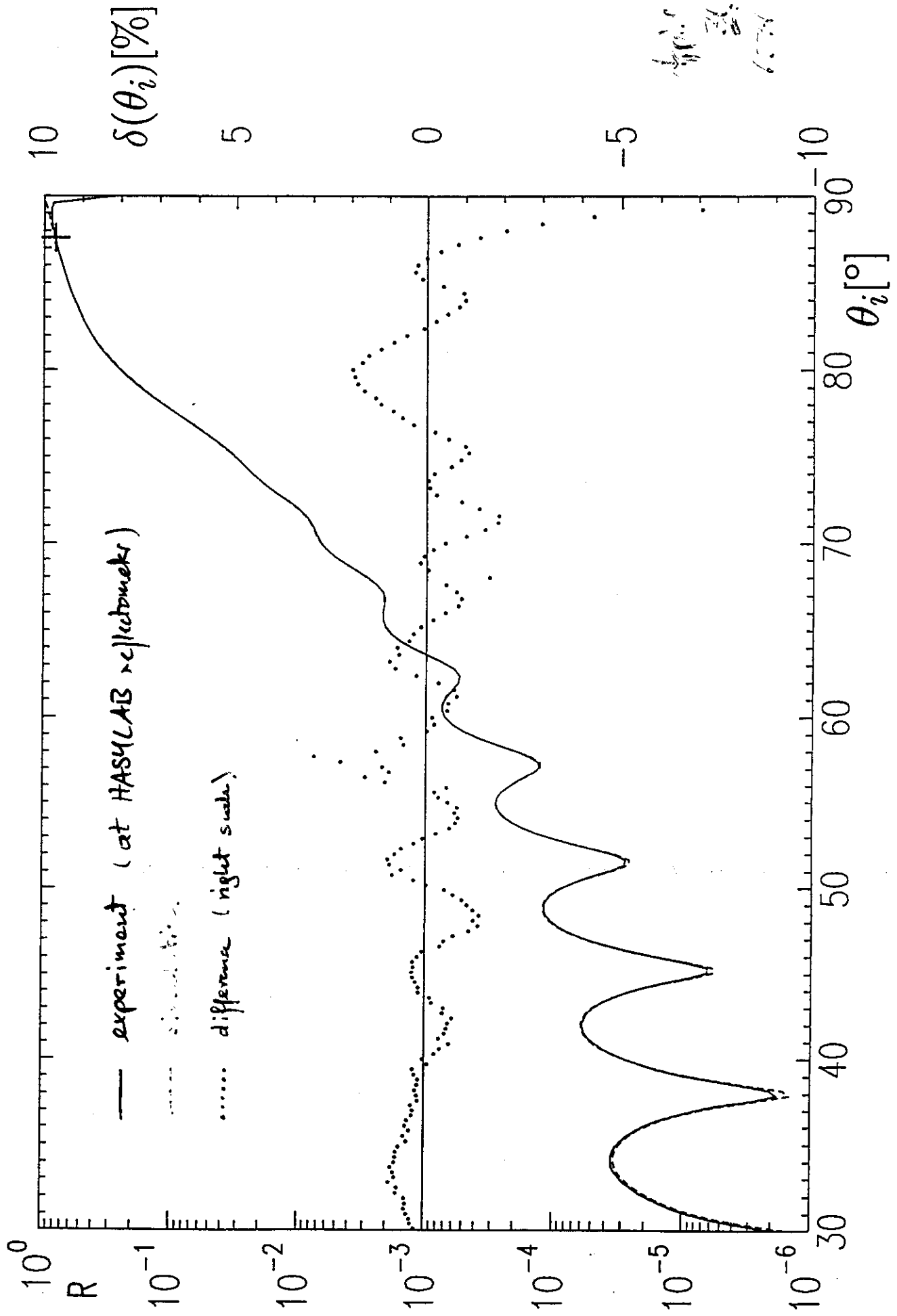
are there any problems?

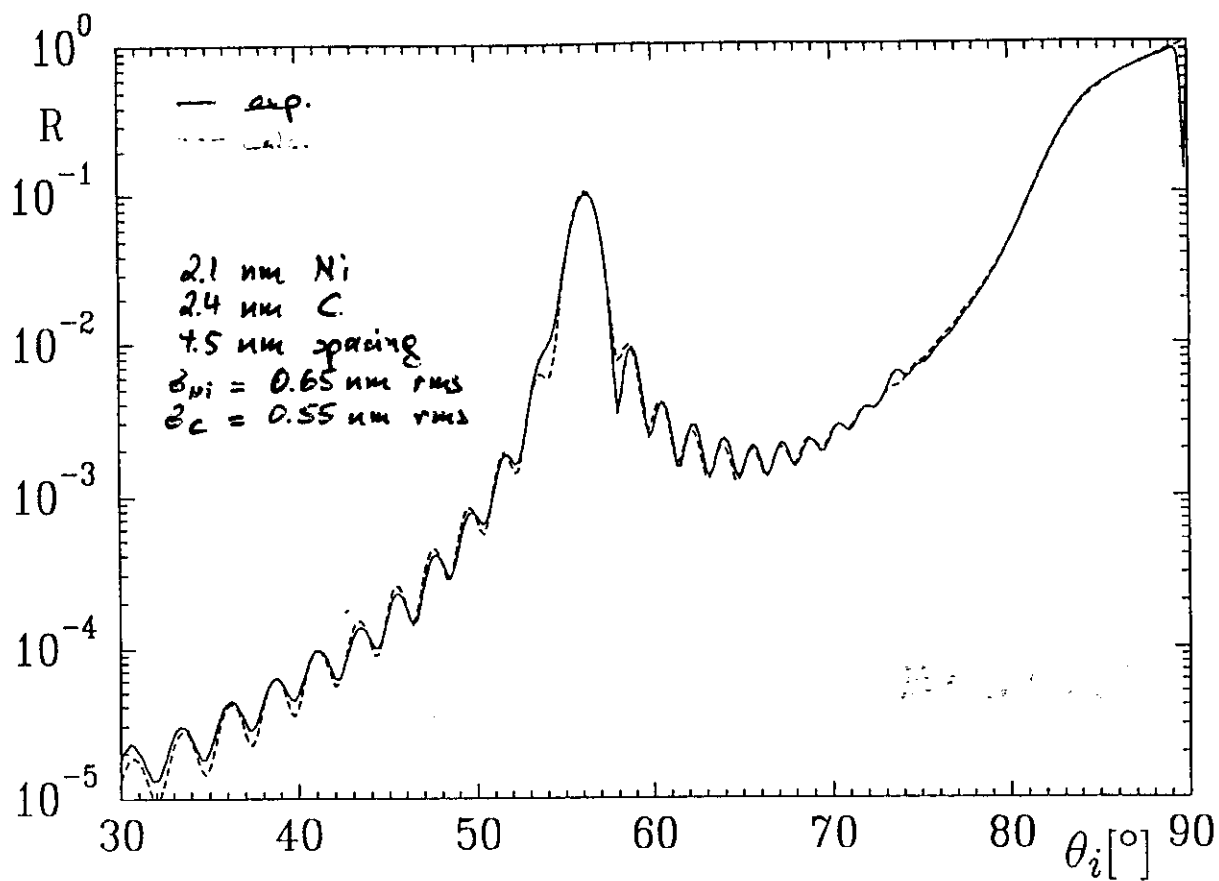




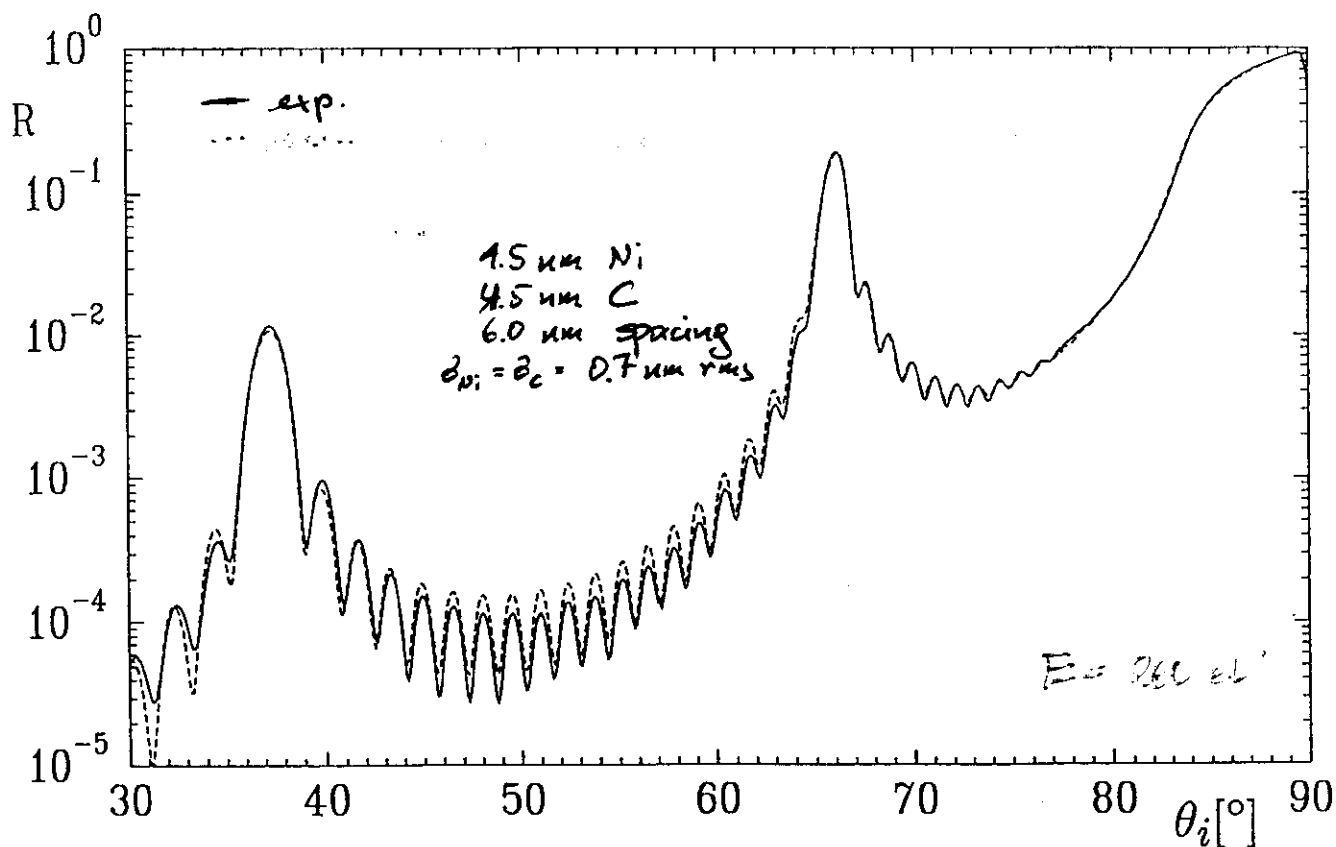
sample description: a simple Ni-film on thick substrate

$d = 50.4 \text{ nm}$ $E = 266 \text{ eV}$



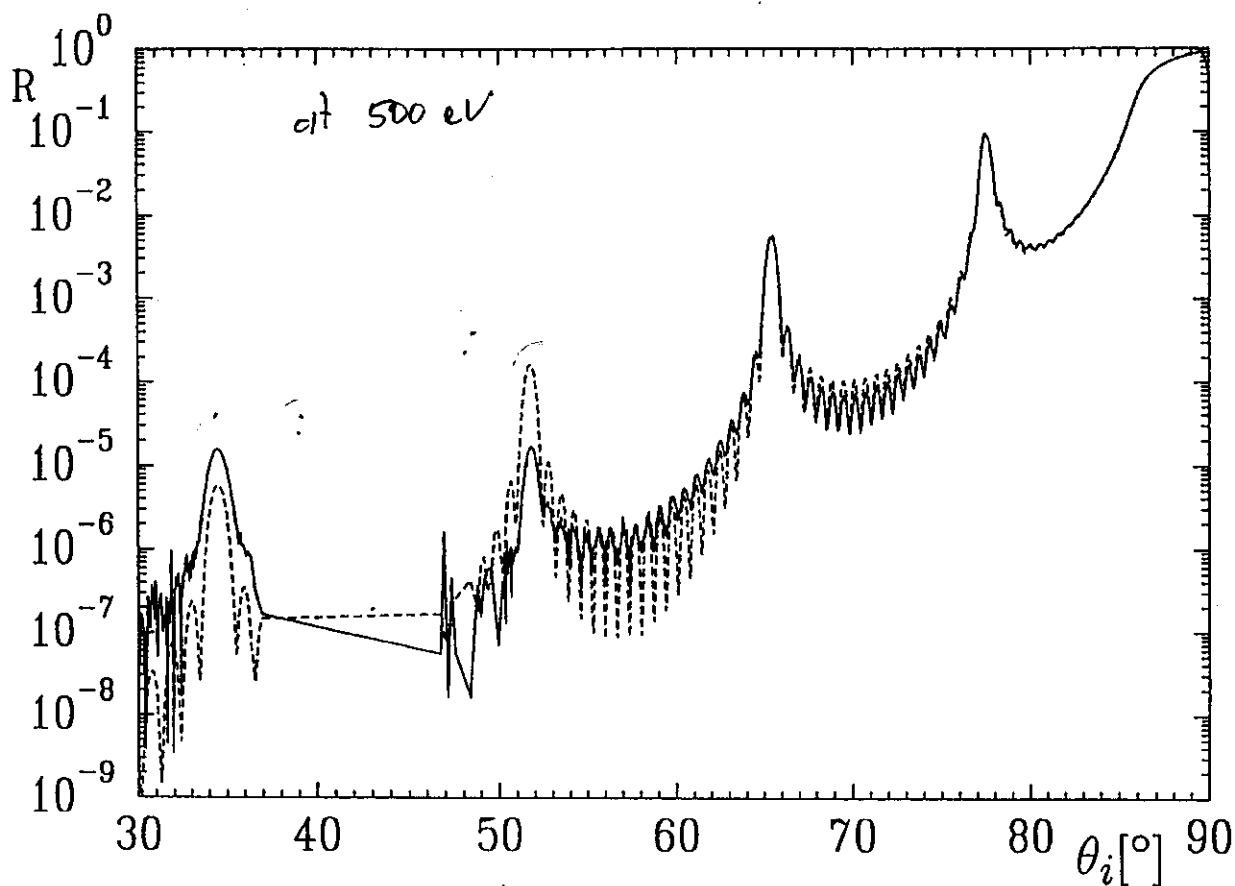


$N = 20 \text{ periods}$ Ni/C

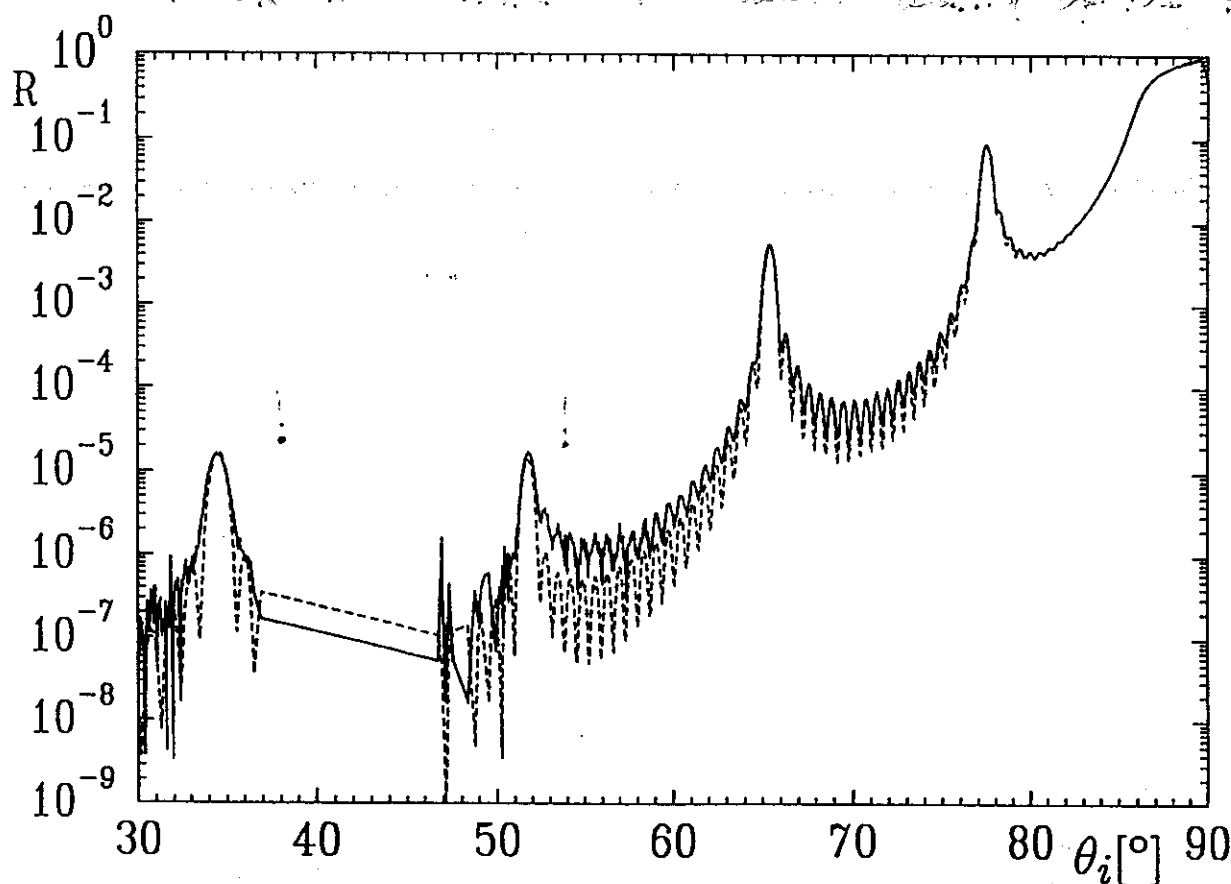


$E = 260 \text{ eV}$

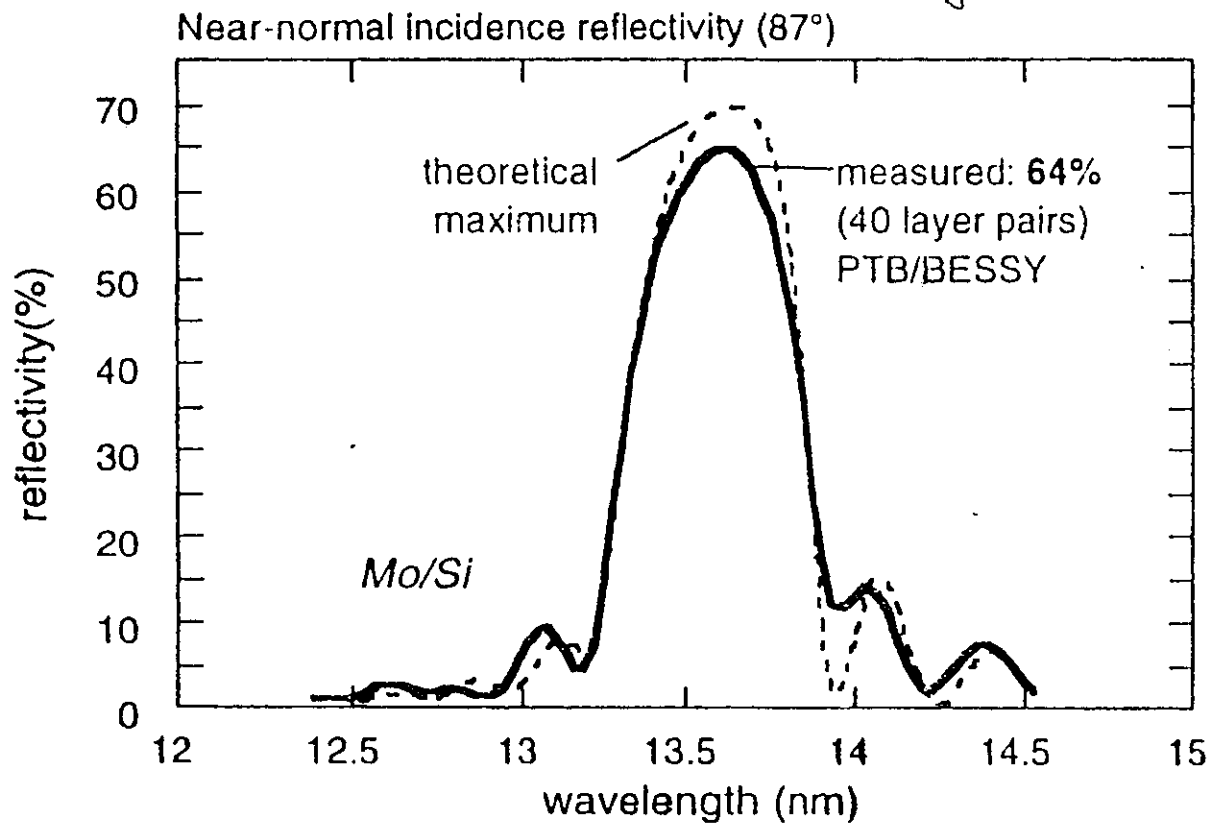
for. ratio 26, 4328 (1987)



we do have $d_{vi} = 1.5 \text{ nm}$
 with interface roughness of 0.7 nm (RHS)



top. (p. 36, 6328 (1987))



method

- *e-beam evaporation + in-situ monitoring*
 - reproducibility: 0.1% run-to-run
- *ion-beam smoothening of interfaces*
 - no accumulation of roughness through stack
 - ability to correct thickness errors
 - extra research tool: e.g. H implantation
- *substrate temperature controlled deposition*

this Mo/Si mirror

- *polished: Si, 300 eV, Kr⁺, 45°*
- ➡ *>90% of maximum theoretical reflectivity achievable*

Elliptical Wiggler

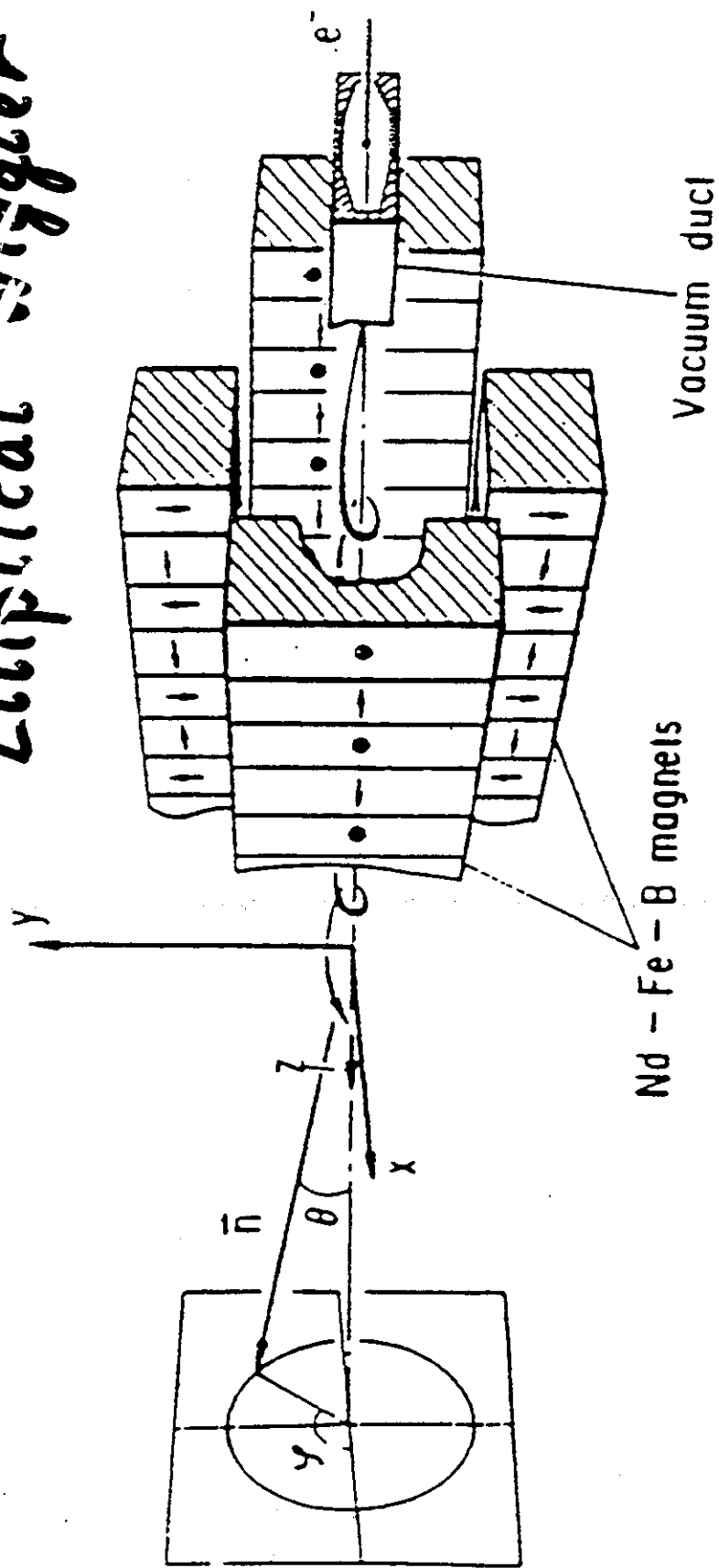
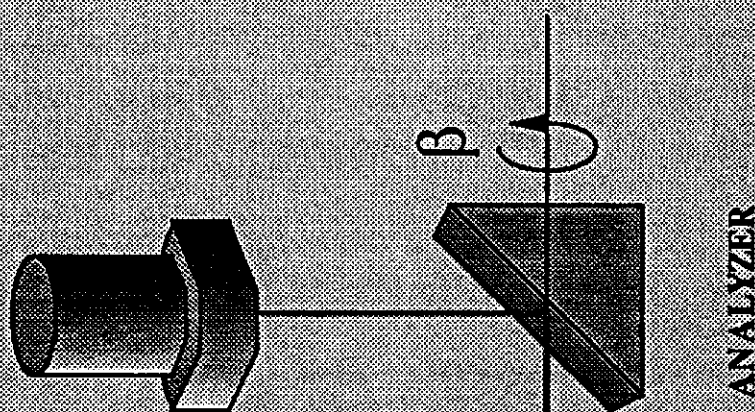


FIG. 2. Schematic illustration of the mechanism of the present insertion devices. Arrows denote the magnetization direction of each magnet, which totally forms the magnetic field given by Eq. (1) on the axis of the device.

$$\mathbf{B} = \pm \mathbf{e}_x B_{x0} \cos(2\pi z/\lambda_w) - \mathbf{e}_y B_{y0} \sin(2\pi z/\lambda_w) \quad (1)$$

REQUIRED DETECTION SCHEME

DETECTOR



PHOTON BEAM

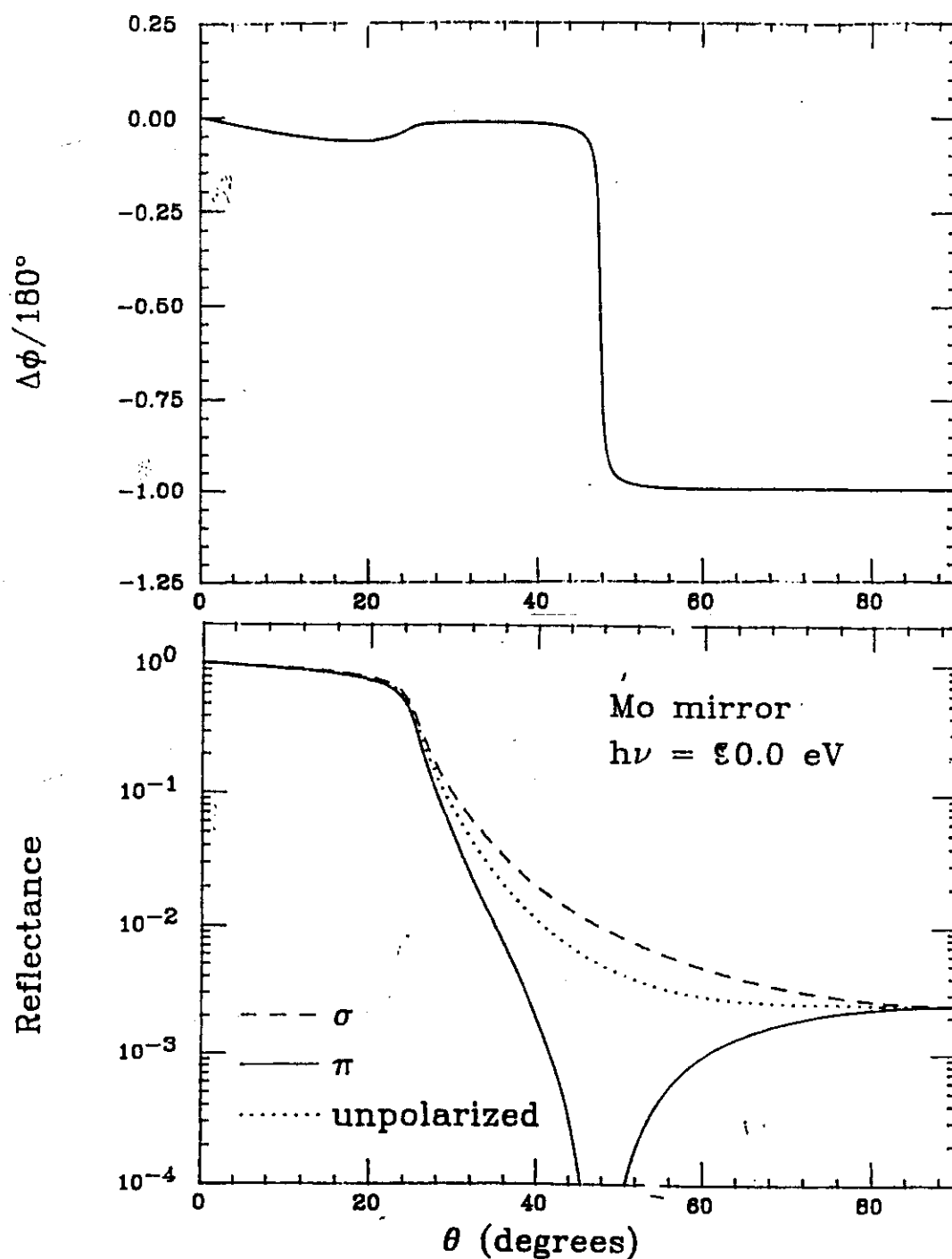
COLLIMATOR

POLARIZER

90 eV PHASE RETARDATION in SIMPLE mirrors:

J.B. Kortright and J.H. Underwood

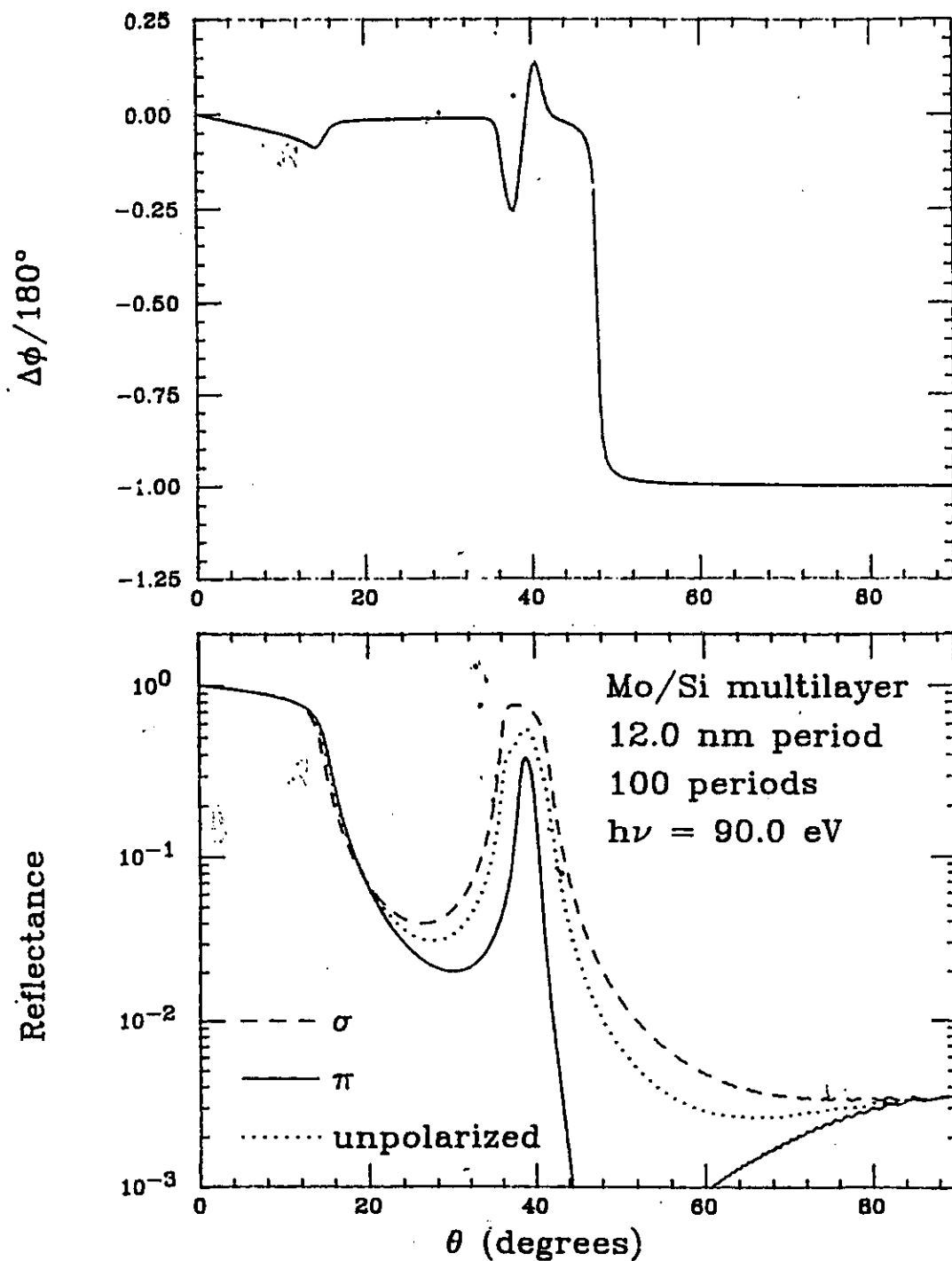
Nucl. Instrum. Methods A291, 272 (1990)



90 eV PHASE RETARDATION in MULTILAYER mirrors:

J.B. Kortright and J.H. Underwood

Nucl. Instrum. Methods A291, 272 (1990)



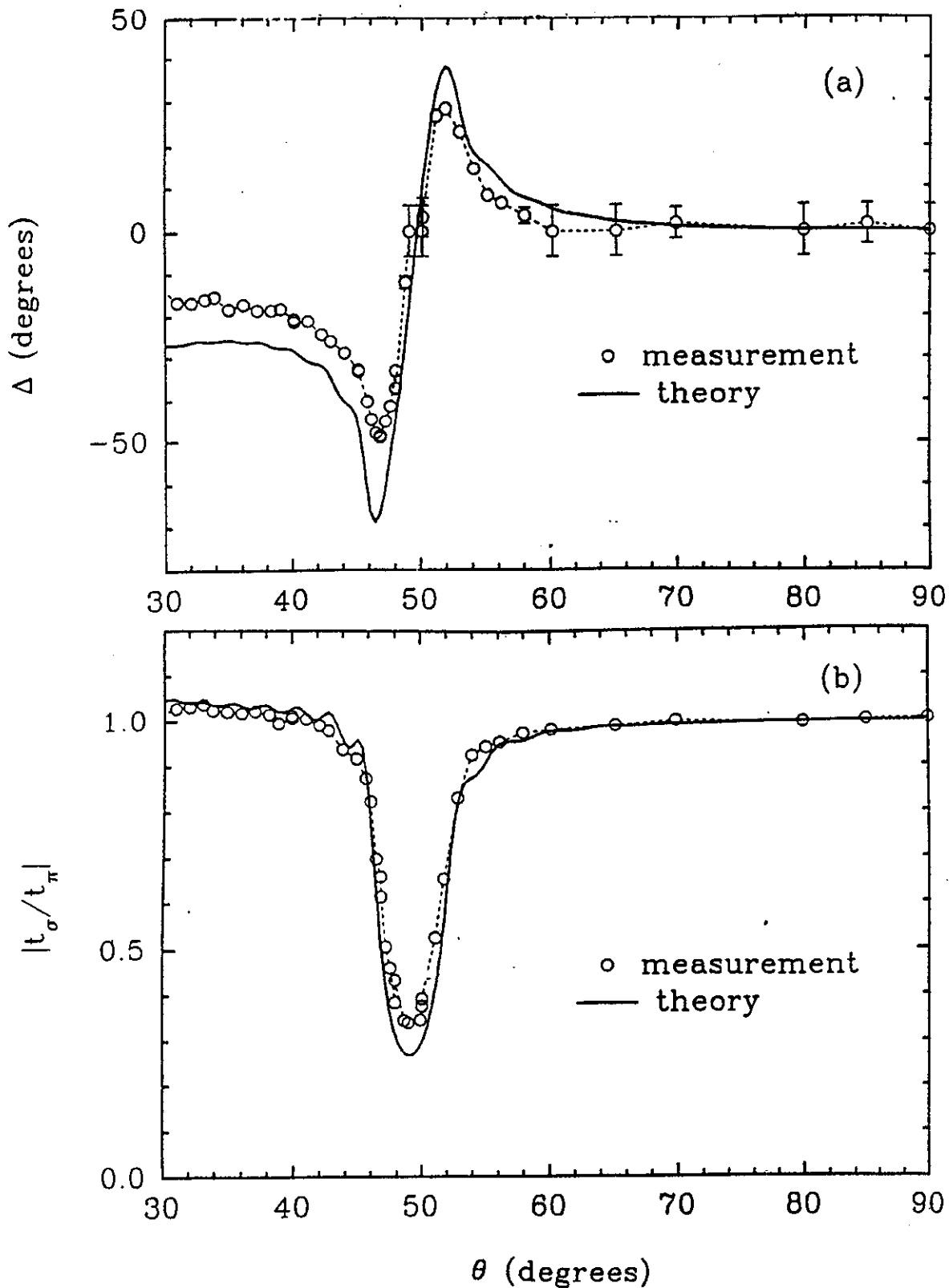
100 periods Mo/Si with $d = 12$ nm (2/3 Si and 1/3 Mo)

97 eV PHASE RETARDATION:

PLASMA SLOWED

J.B. Kortright, H. Kimura, V. Nikitin, K. Mayama, M. Yamamoto and M. Yanagihara

Appl. Phys. Lett. 60, 2963 (1992)



20 periods Mo/Si with $d = 8.75$ nm (2/3 Si and 1/3 Mo)

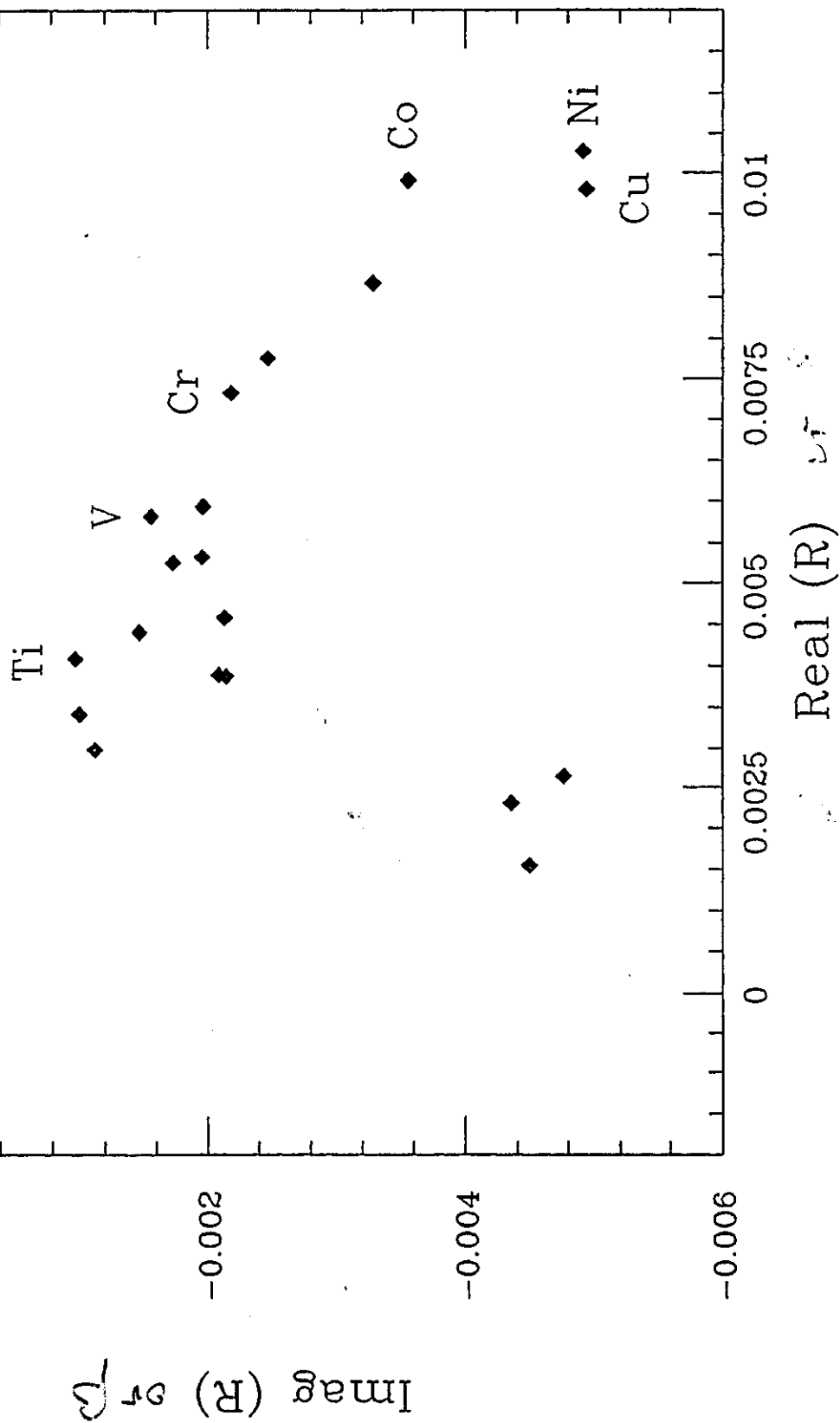
effect 2/3 of theory with 20% transmission

Components of reflexion coefficient

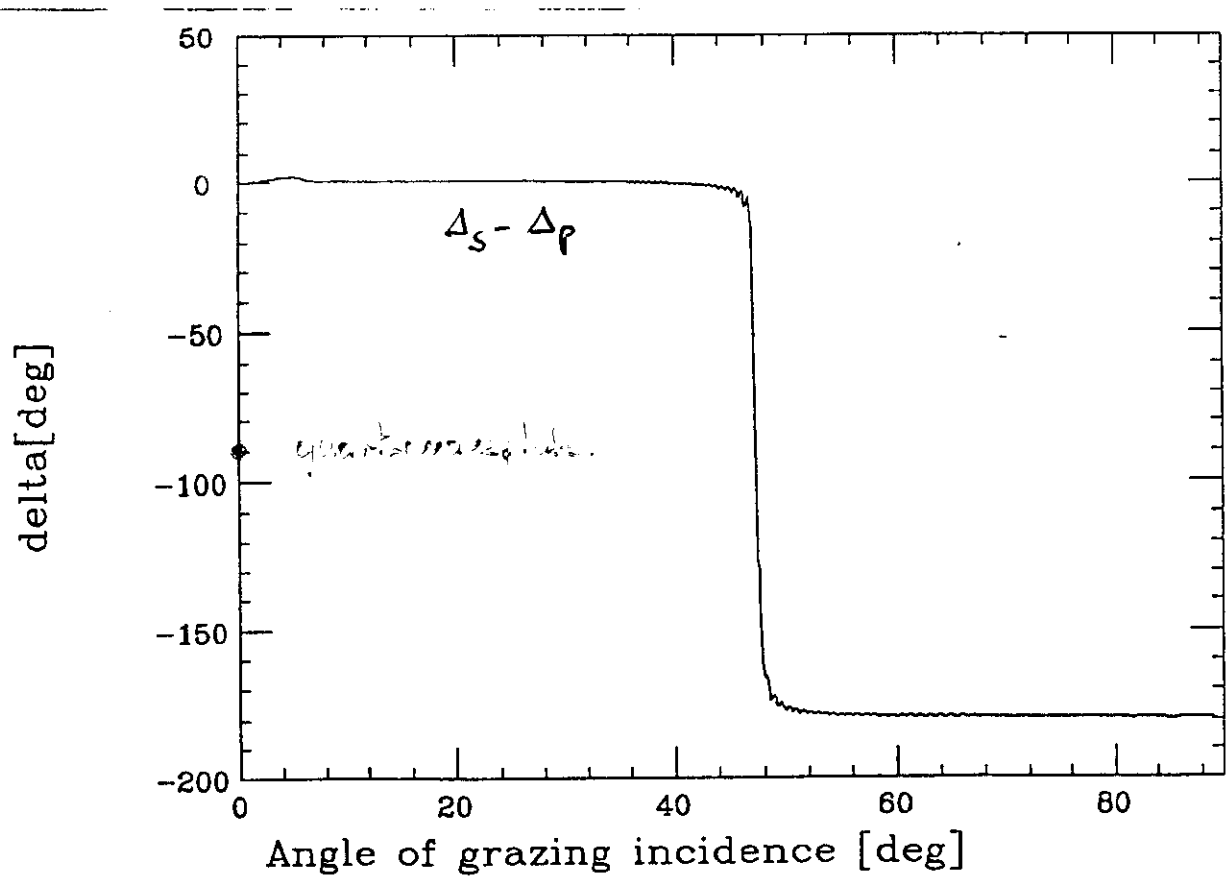
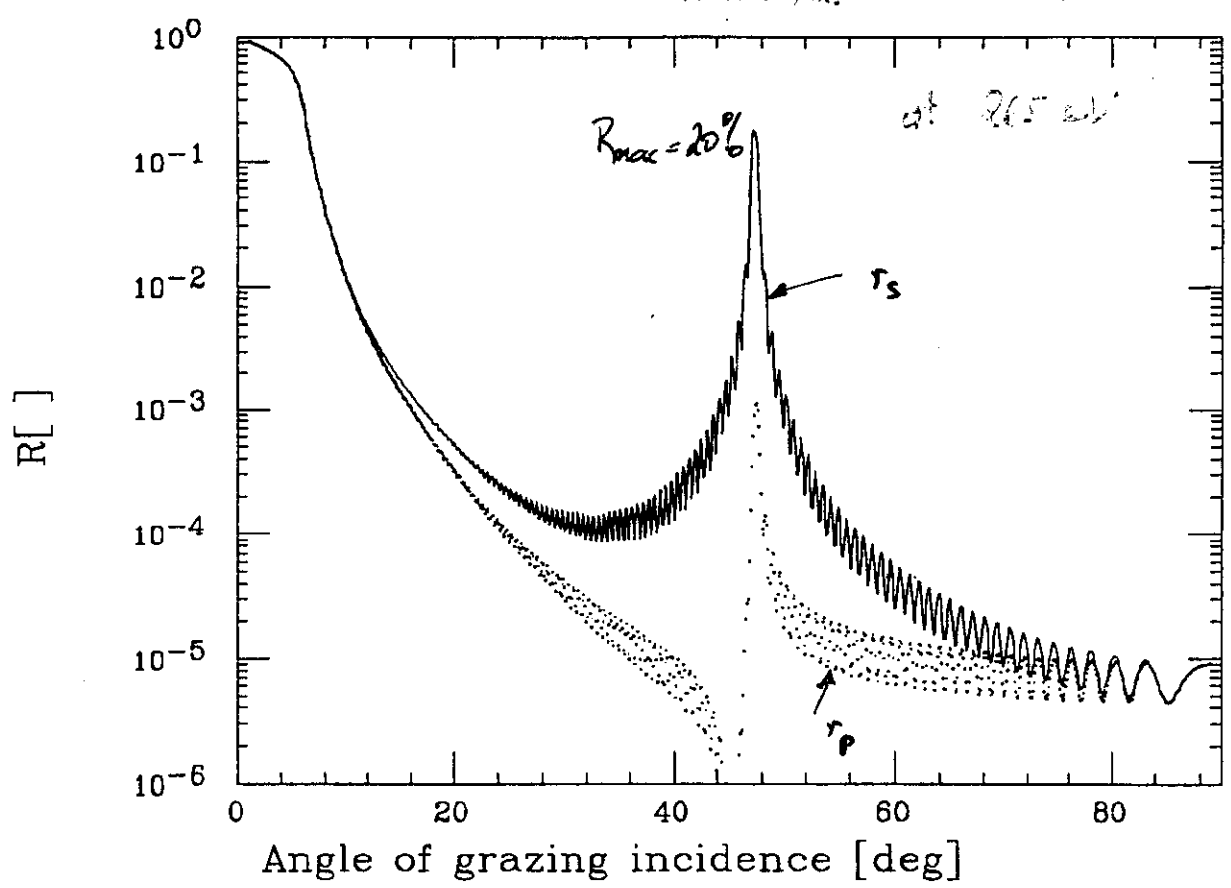
$$n = A - \delta + i\epsilon$$

$$R = \left(\frac{n-1}{n+1} \right)^2$$

for 283 eV



Calculated: Δ_s and Δ_p of ultra-thin film
 $d = 0.1 \mu\text{m}$ $n = 1.5$ $n_s = 1.0$ $n_p = 1.5$



2.1.1.1. for the given objects:

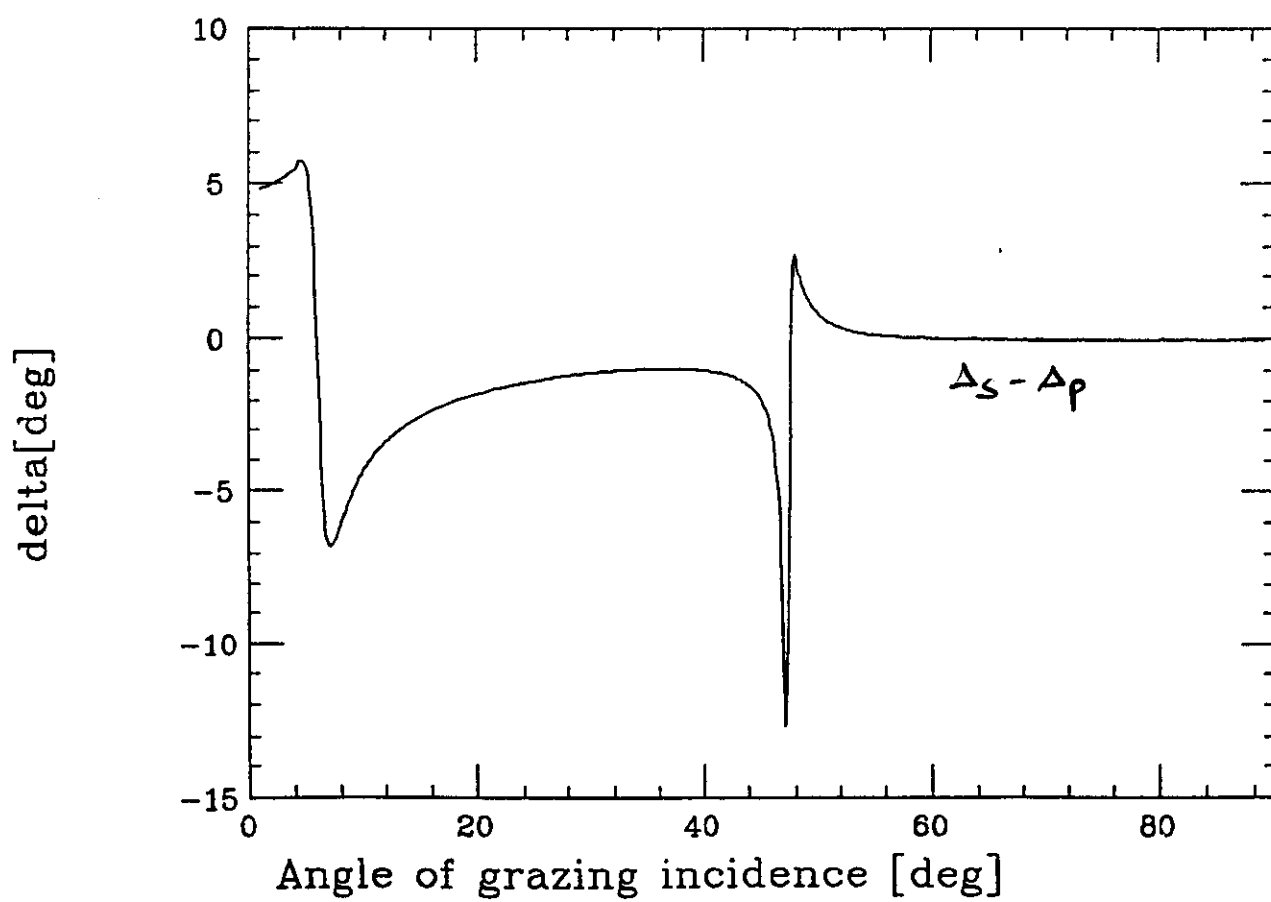
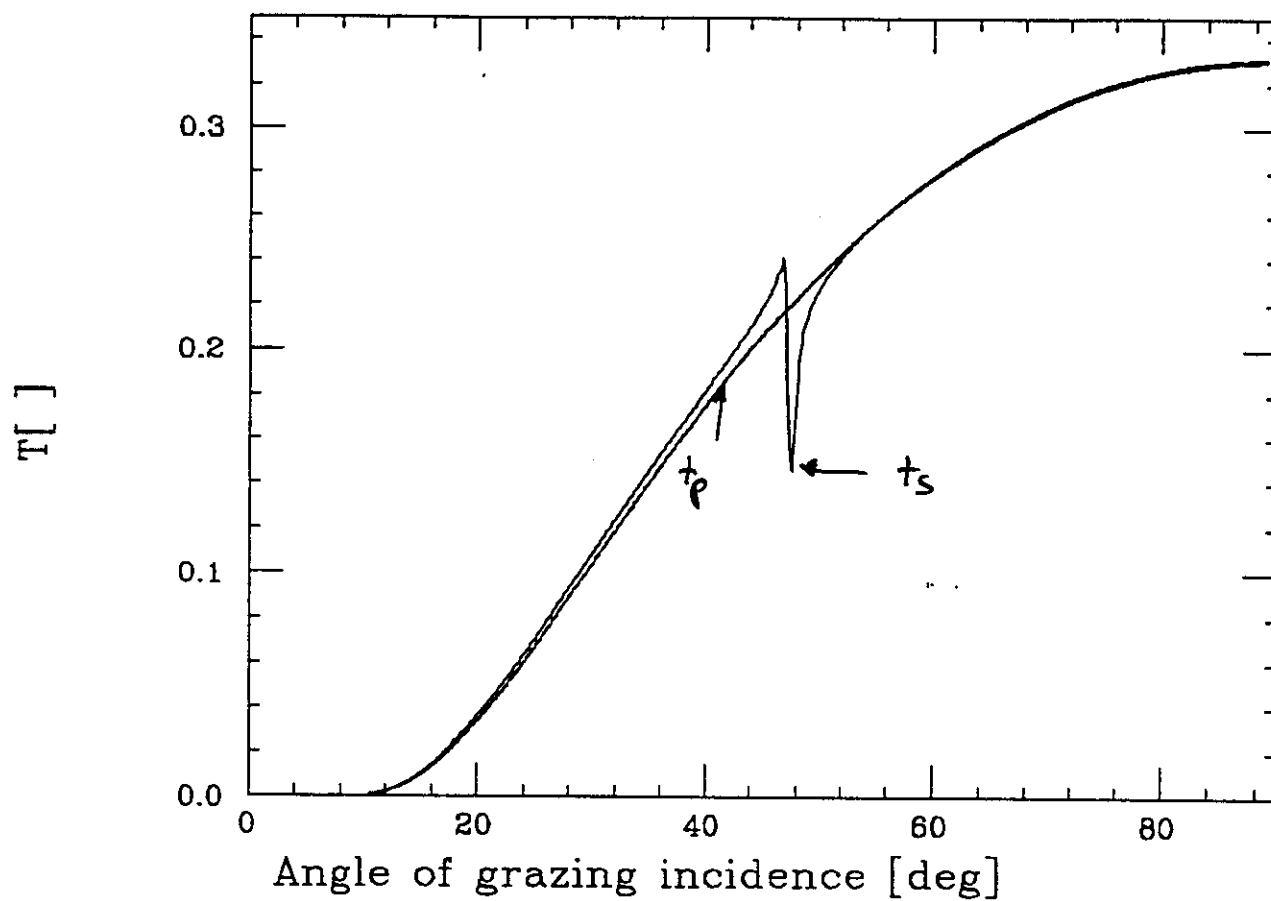
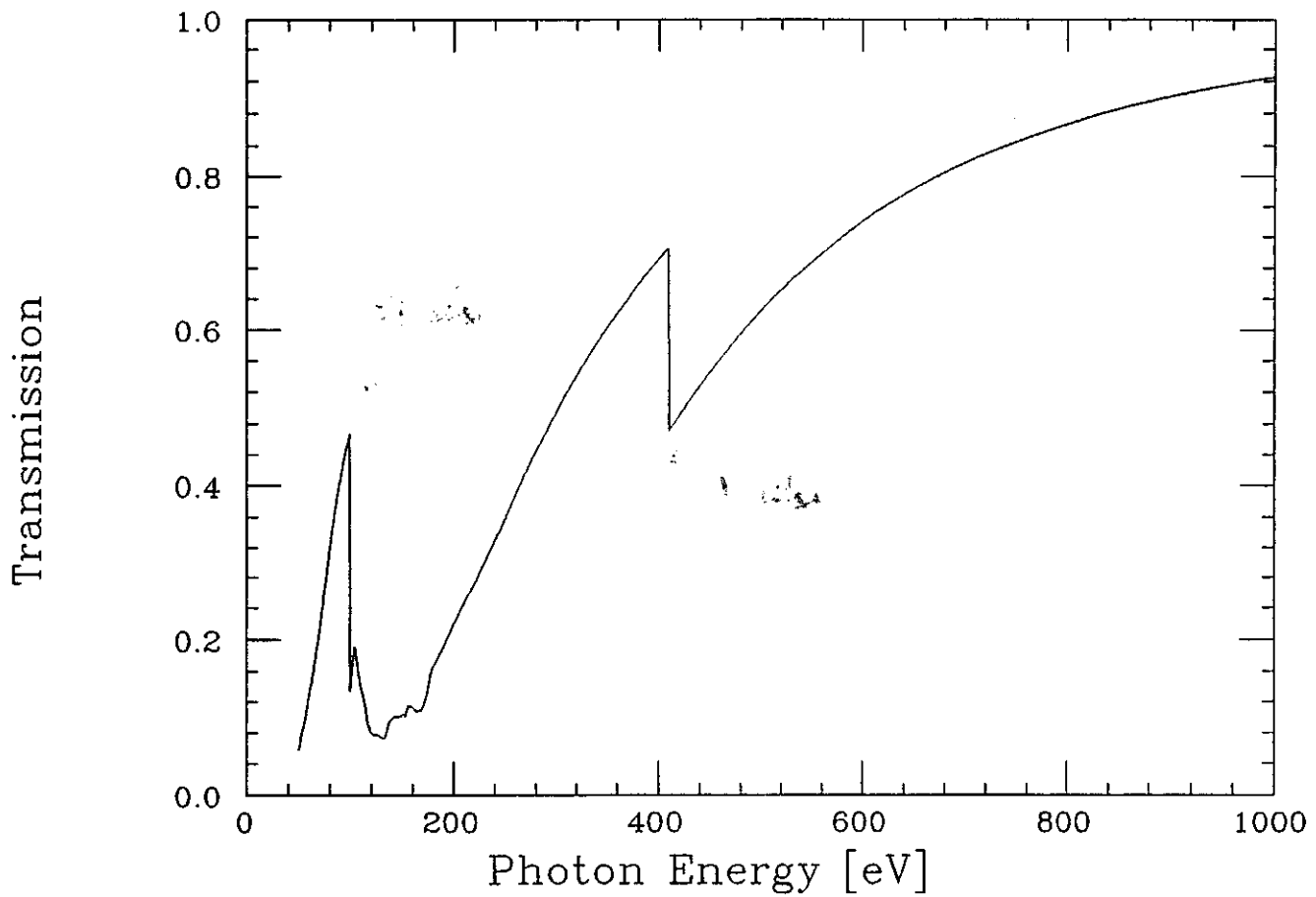


Figure 1. Calculated transmission spectra

calculated with recent collection of ϵ_1 and ϵ_2 data

Silicon Nitride Filter: $d=100$ nm



*Collaboration project for polarization
characterization*

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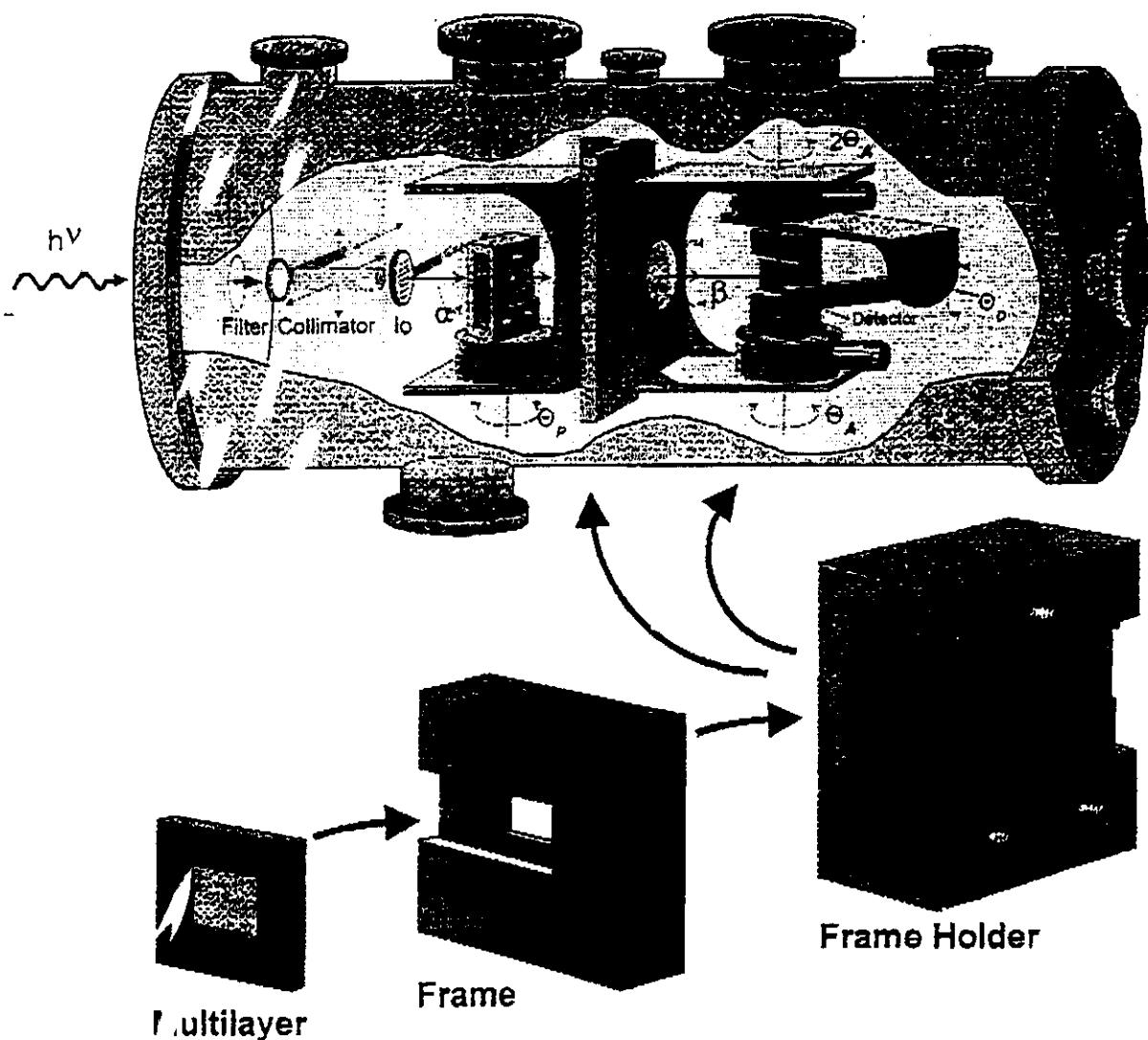
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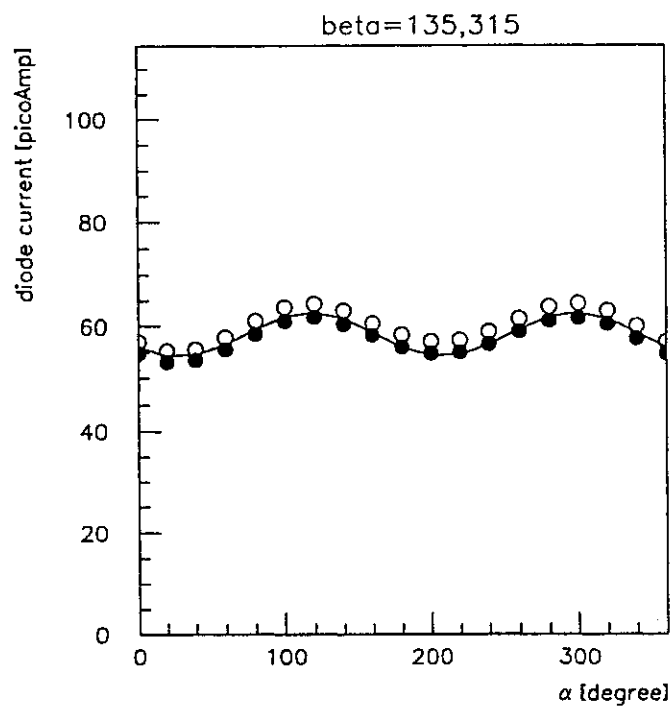
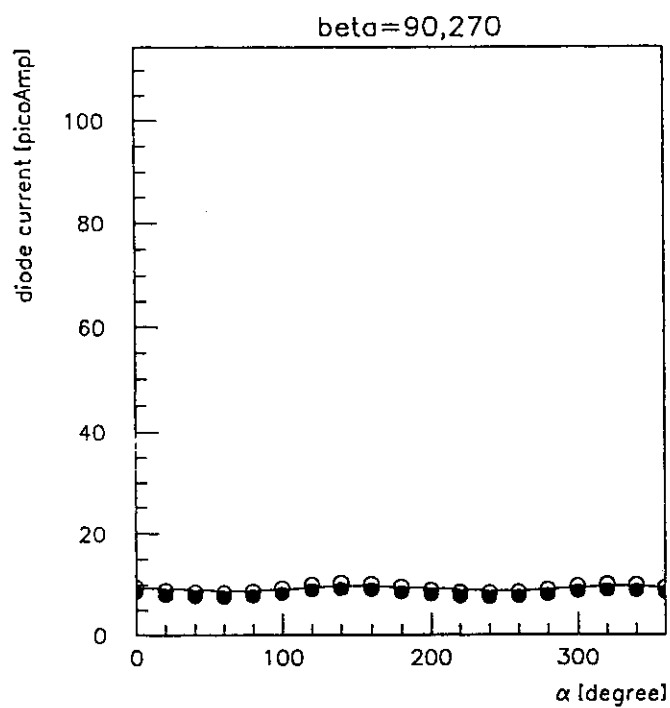
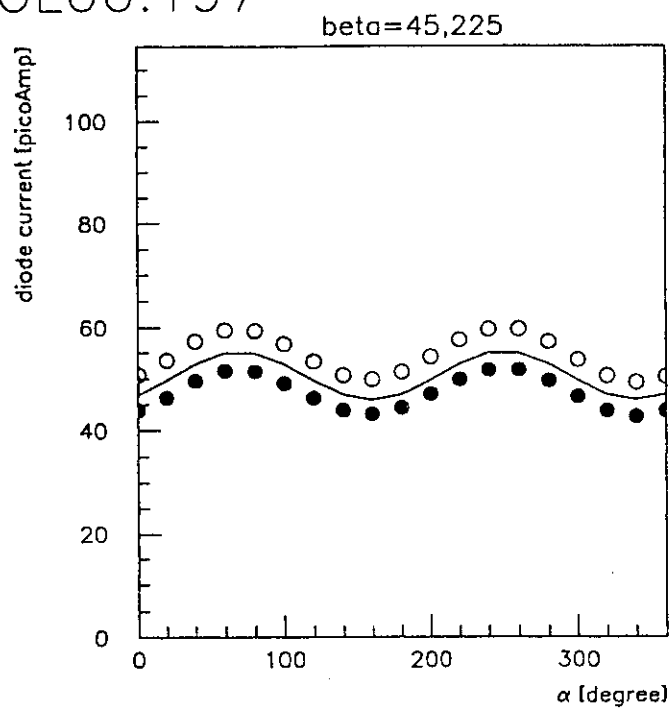
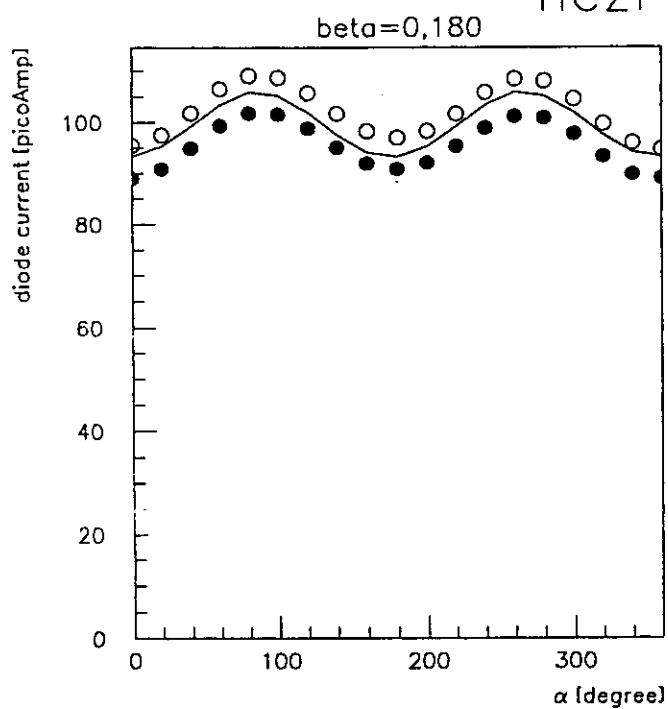
BESSY Soft-X-Ray Polarimeter



Applications / features	Technical features
• Determination of all Stokes parameters S_0 (Intensity), S_1, S_2, (lin.pol.), S_3 (circ.pol.) • Optical elements - Polarizer: transmission multilayers, $\lambda/4$-plates (MgF_2) - Analyzer: reflection multilayers, mirrors, (Mo/Si, Cr/C, Cr/Sc, Ni/Ti) • Working energy: visible to soft X-ray range determined by optical elements in use • Characterization of optical elements: - Multilayers, transmission foils, crystals - Reflectance, transmission - Phase retardation - s- and p-polarization geometry	• 8-axes polarimeter: - Collimator (pinholes $\varnothing$: 0.5 – 2 mm) Translation in x, y ± 12 mm - Polarizer / analyzer: Azimuthal angle: $0^\circ \leq \alpha, \beta \leq 370^\circ$ Incidence angle: $0^\circ \leq \Theta_p, \Theta_A \leq 90^\circ$ - Detector (GaAsP diode): scan range: in plane $0^\circ \leq 2\Theta_A \leq 180^\circ$ off plane $-27^\circ \Theta_D \leq +27^\circ$ • UHV compatible, $p \leq 10^{-8}$ mbar • UHV-goniometers, motors f. positioning ($>0.001^\circ$) • Sample size: $10 \times 10 \times 0.5 \text{ mm}^3$ to $140 \times 50 \times 11 \text{ mm}^3$ • Magazine for 10 samples, alignable in situ • Filter for higher order: Be, C_6H_8, Ti, Cr, Fe, Cu

18/12/97 22.36

hczPOL00.197



The BESSY polarimeter

(A. Gaupp and M. Mast: Rev. Sci. Instrum. **60**, 2213 (1989))

$$\begin{aligned} t_{s1} &= |t_{s1}| \exp(i\delta_{s1}) & r_{s2} &= |r_{s2}| \exp(i\delta_{s2}) \\ t_{p1} &= |t_{p1}| \exp(i\delta_{p1}) & r_{p2} &= |r_{p2}| \exp(i\delta_{p2}) \end{aligned}$$

$$\Delta_1 = \delta_{p1} - \delta_{s1} \qquad \Delta_2 = \delta_{p2} - \delta_{s2} = 0$$

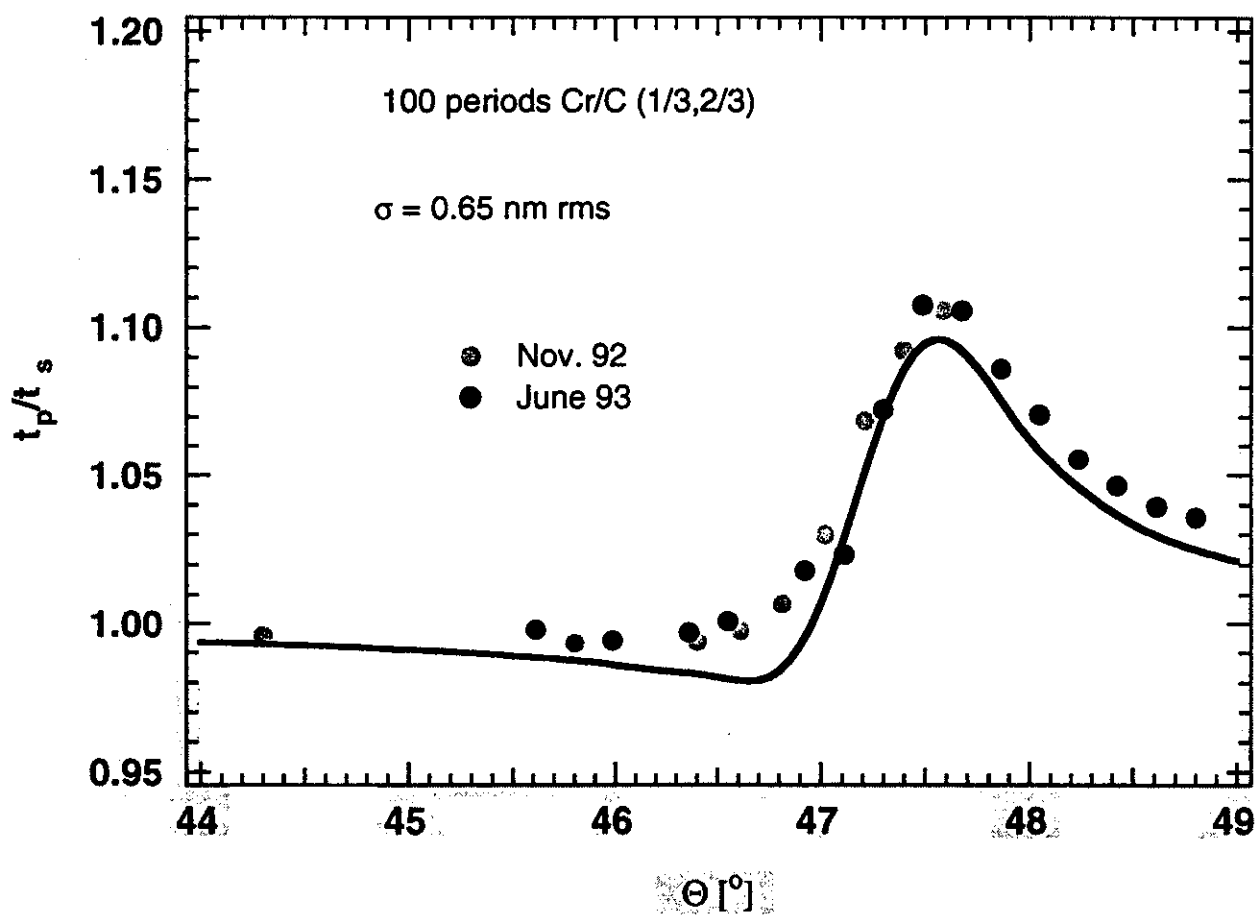
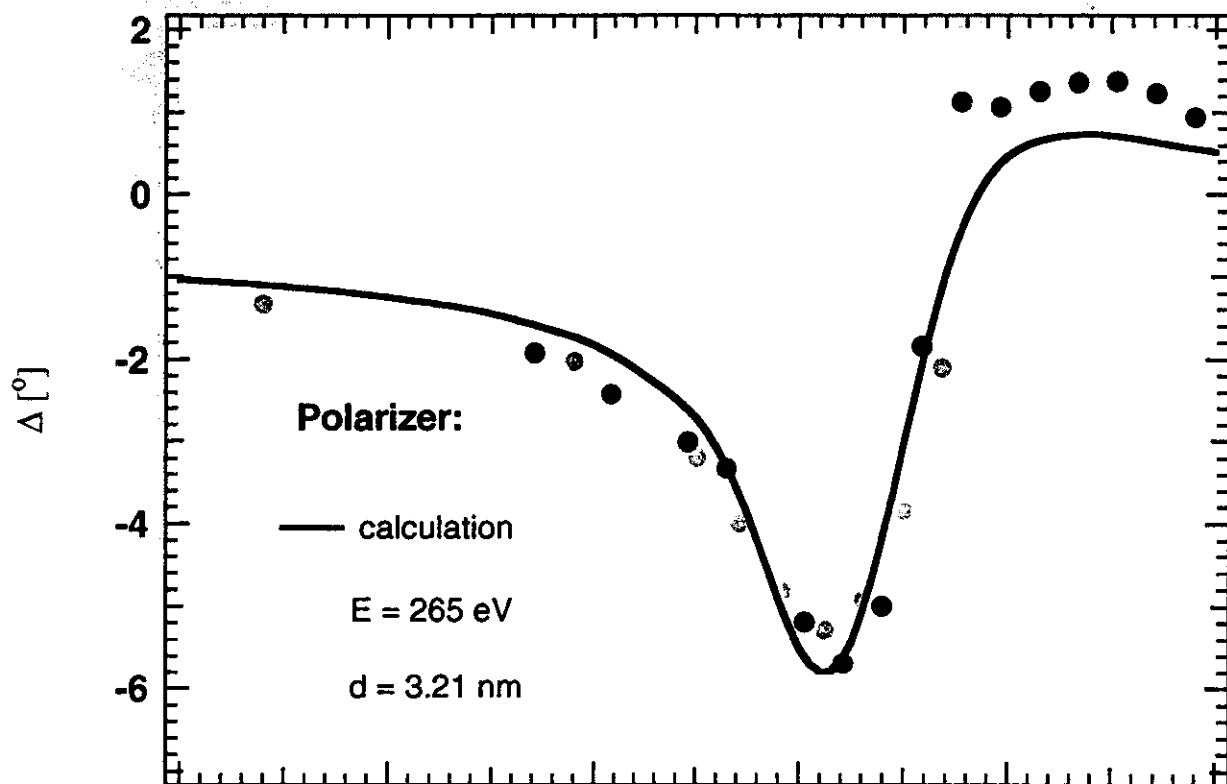
$$\tan \Psi_1 = |t_{p1}|/|t_{s1}| \qquad \tan \Psi_2 = |r_{p2}|/|r_{s2}|$$

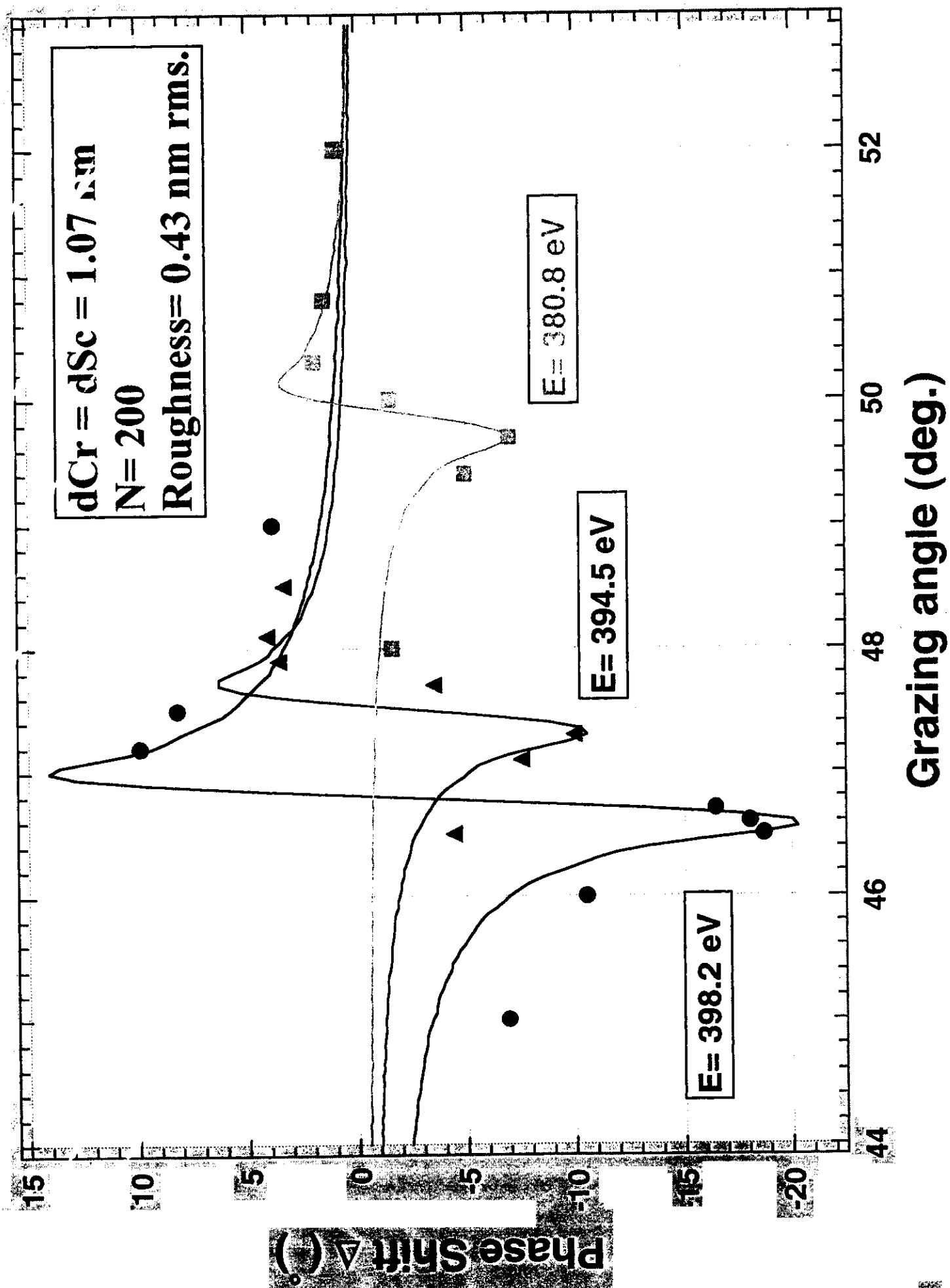
Stokes-parameters:

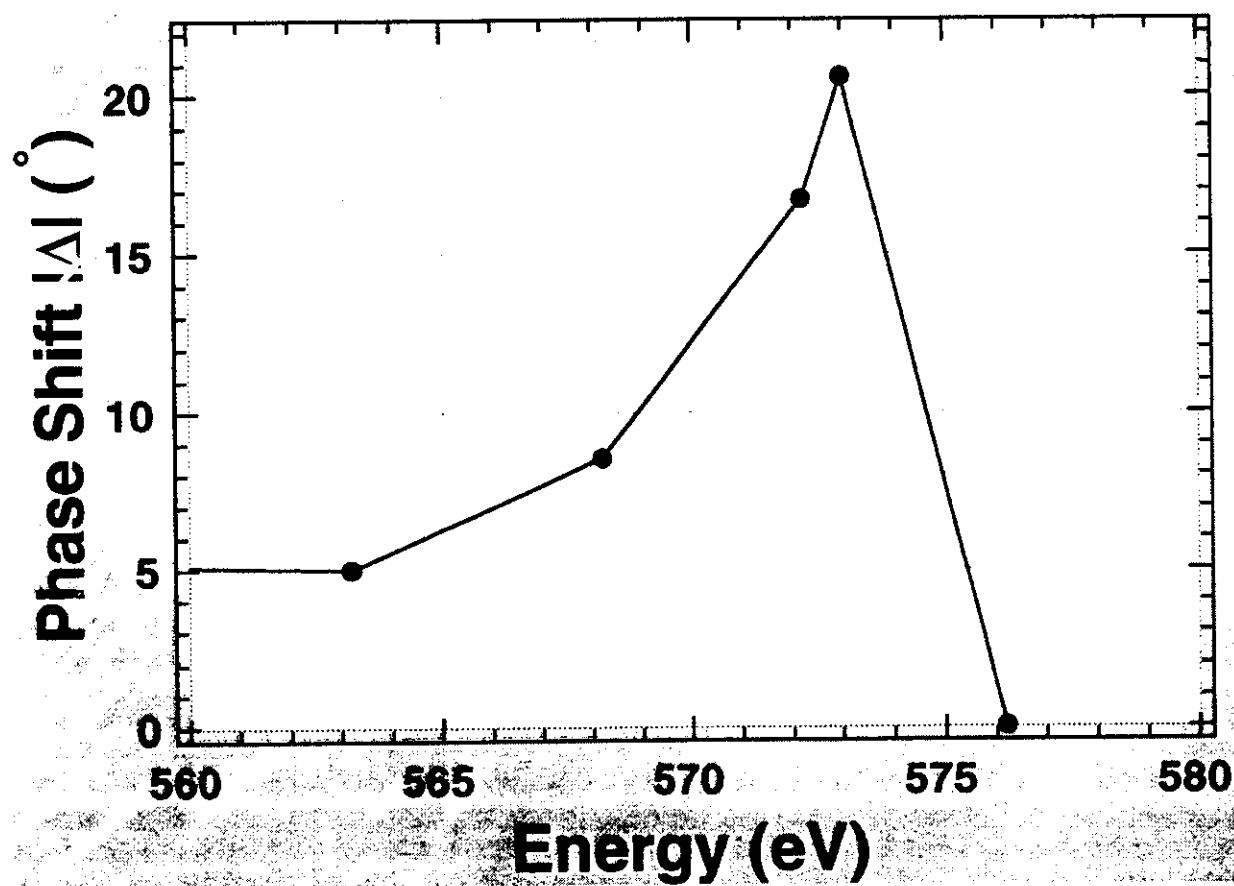
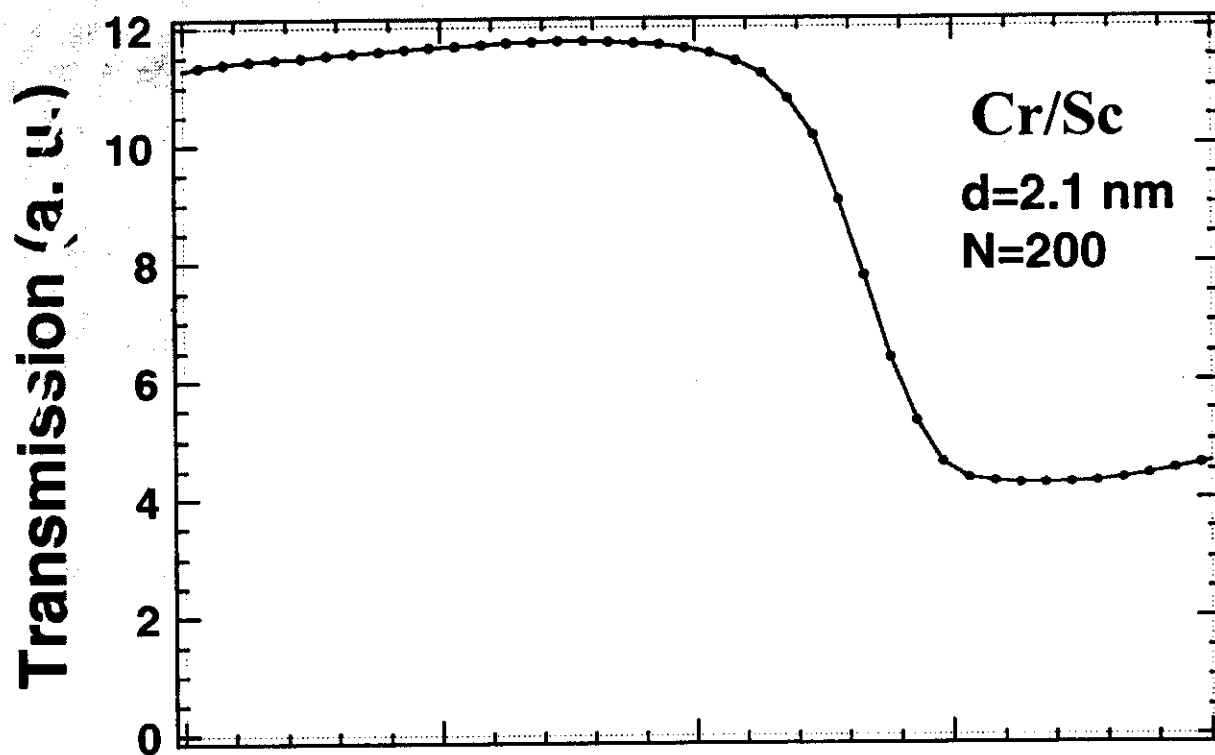
S_0, S_1, S_2 and S_3

$$I_{\text{pass}} = \frac{1}{2} (t_{s1}^2 + t_{p1}^2) * \frac{1}{2} (r_{s2}^2 + r_{p2}^2) *$$

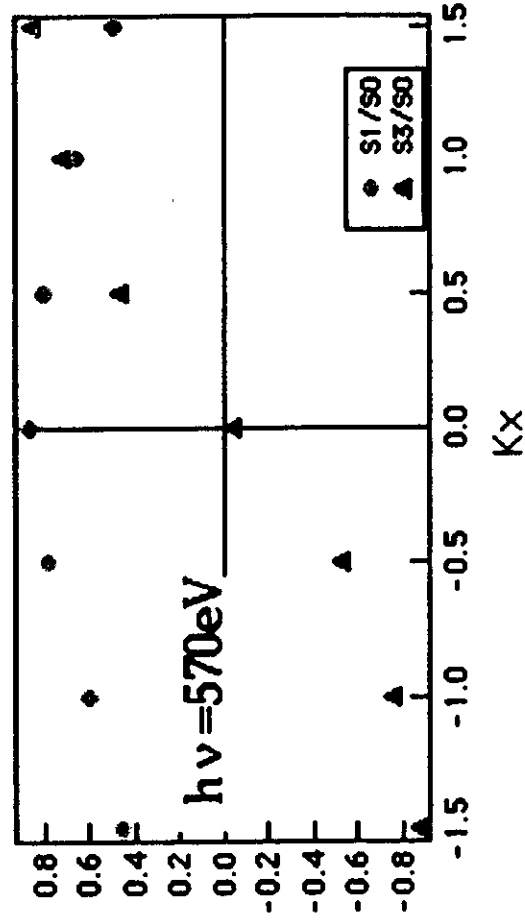
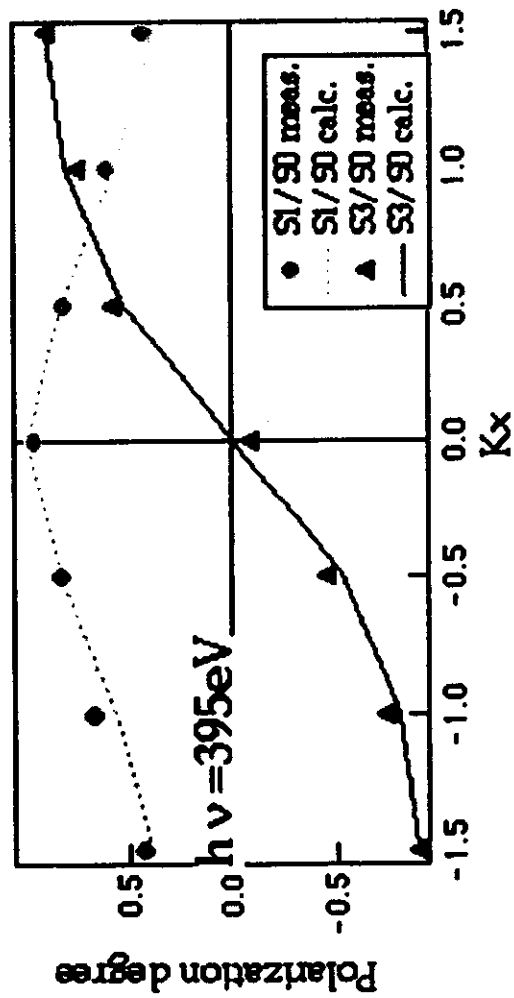
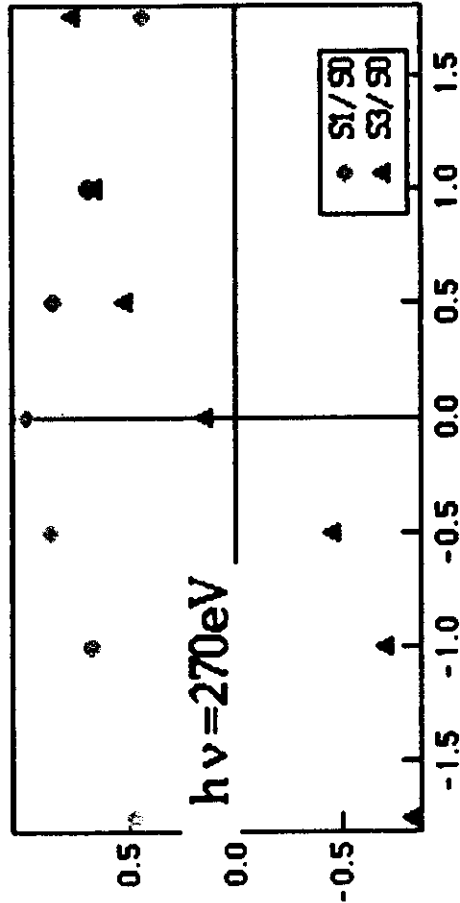
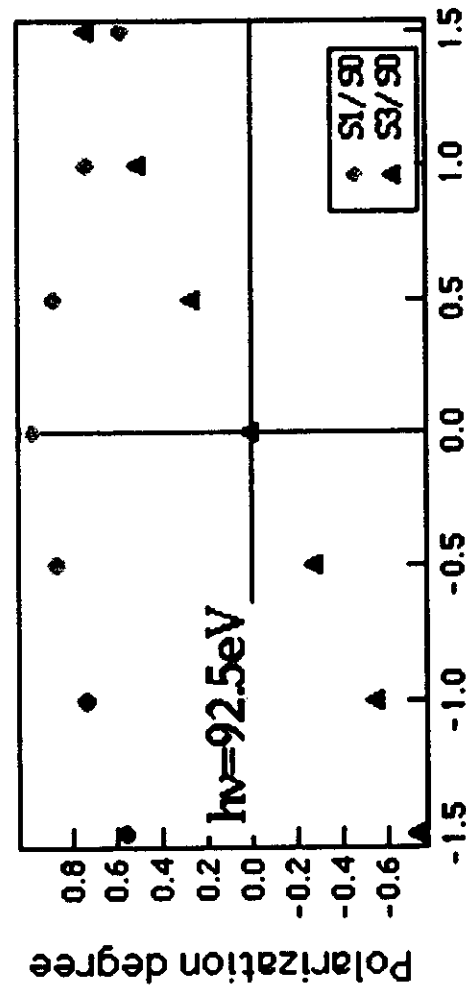
$$\begin{aligned} & \{ \\ & + \cos 2\alpha \qquad \qquad \qquad [-S_1 \cos 2\Psi_1] \qquad \qquad \qquad + \sin 2\alpha \qquad \qquad \qquad [-S_2 \cos 2\Psi_1] \\ & + \qquad \cos 2\beta \qquad \qquad \qquad [-S_1 \cos 2\Psi_2 * (1 + \sin 2\Psi_1 \cos \Delta_1)/2] \qquad \qquad \qquad + \qquad \sin 2\beta \qquad \qquad \qquad [-S_2 \cos 2\Psi_2 * (1 + \sin 2\Psi_1 \cos \Delta_1)/2] \\ & + \cos 2\alpha \cos 2\beta \qquad \qquad \qquad [+S_0 \cos 2\Psi_1 * \cos 2\Psi_2] \qquad \qquad \qquad + \sin 2\alpha \cos 2\beta \qquad \qquad \qquad [+S_3 \sin 2\Psi_1 * \cos 2\Psi_2 \sin \Delta_1] \\ & + \cos 2\alpha \sin 2\beta \qquad \qquad \qquad [-S_3 \sin 2\Psi_1 * \cos 2\Psi_2 \sin \Delta_1] \qquad \qquad \qquad + \sin 2\alpha \sin 2\beta \qquad \qquad \qquad [+S_0 \cos 2\Psi_1 * \cos 2\Psi_2] \\ & + \cos 4\alpha \cos 2\beta \qquad \qquad \qquad [-S_1 \cos 2\Psi_2 * (1 - \sin 2\Psi_1 \cos \Delta_1)/2] \qquad \qquad \qquad + \sin 4\alpha \cos 2\beta \qquad \qquad \qquad [-S_2 \cos 2\Psi_2 * (1 - \sin 2\Psi_1 \cos \Delta_1)/2] \\ & + \cos 4\alpha \sin 2\beta \qquad \qquad \qquad [+S_2 \cos 2\Psi_2 * (1 - \sin 2\Psi_1 \cos \Delta_1)/2] \qquad \qquad \qquad + \sin 4\alpha \sin 2\beta \qquad \qquad \qquad [-S_3 \cos 2\Psi_2 * (1 - \sin 2\Psi_1 \cos \Delta_1)/2] \} \end{aligned}$$





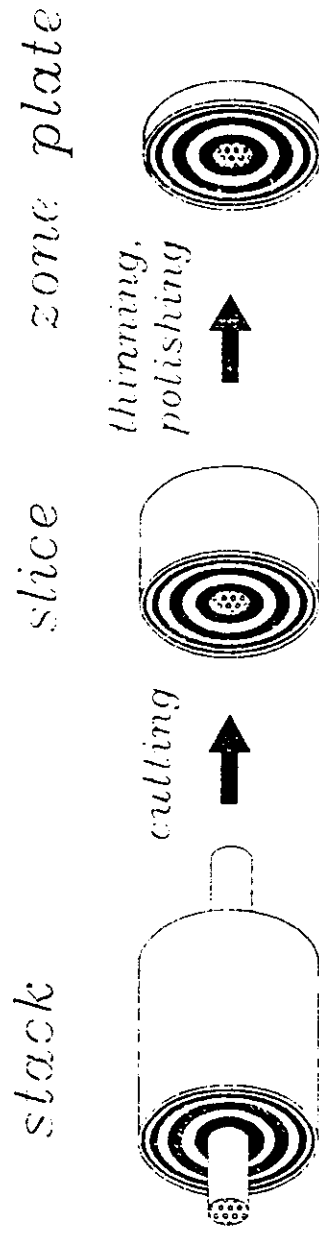


Wiggler mode
Mo/Si 92.5 eV (edge of Si)
Cr/C 270 eV (edge of C)
Cr/Sc 395 eV (edge of Sc)
Cr/ Sc 574 eV (edge of Cr)





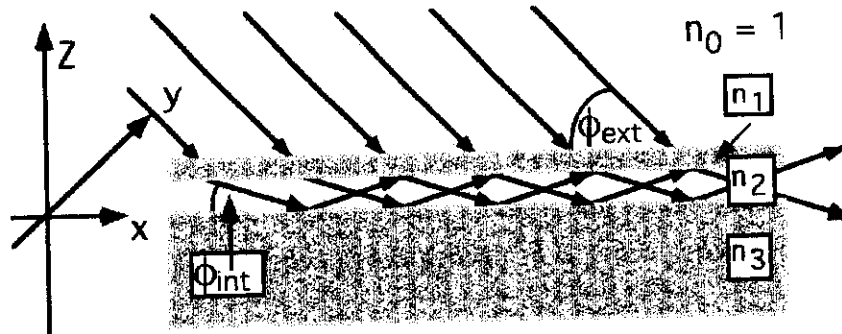
Sputtered Sliced Zone Plate Technology: [11, 8]



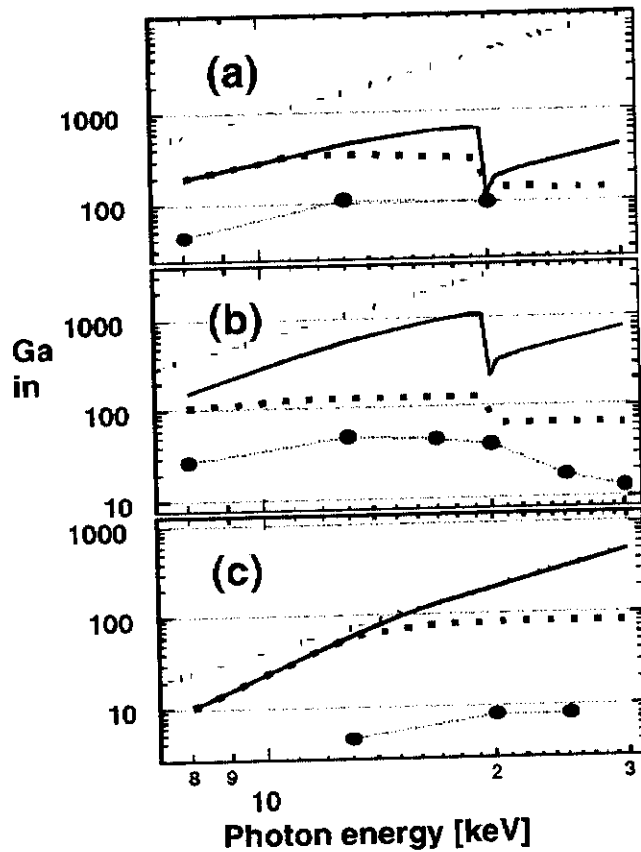
- combination of highest aspect ratios and high resolution is possible
- optimal thickness of zone plates can be used for high efficiencies
- no restriction by reactive ion etching methods if the spacer material has not to be removed
- the limitation in outermost zone width is given by roughness and interdiffusion of the layers

Field intensity enhancement in thin film

W. Jark et al, to be publ.



enhancement = gain if $n_2 > n_1, n_3$



Dots: ideal interfaces, no absorption in top and bottom layer
Line: rigorous calculation
Dashes: rigorous calculation correct for finite beam divergence

The data is presented as function of photon energy for different materials and thickness:

A) 74 nm Be between 5.5 and 20 nm Mo.

B) 110 nm Be between 5.5 and 20 nm Mo.

C) 160 nm C (with relatively 1% inclusions of Ar) between 4 and 20 nm Cr.

