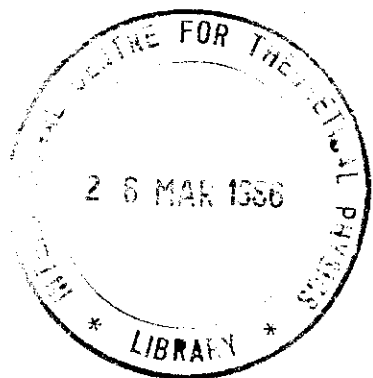




INTERNATIONAL ATOMIC ENERGY AGENCY
UNITED NATIONS EDUCATIONAL, SCIENTIFIC AND CULTURAL ORGANIZATION



INTERNATIONAL CENTRE FOR THEORETICAL PHYSICS
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WORKSHOP ON OPTICAL FIBER COMMUNICATION
(24 February - 21 March 1986)

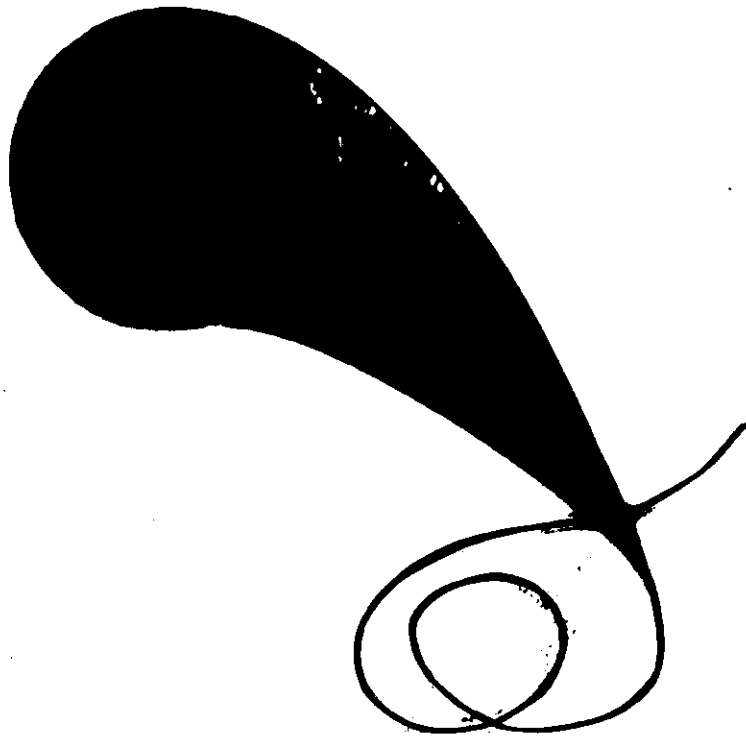
PROPAGATION THEORY IN OPTICAL FIBERS

presented by

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Centro Studi e Laboratori
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Italy

These are preliminary lecture notes, intended only for distribution to participants.

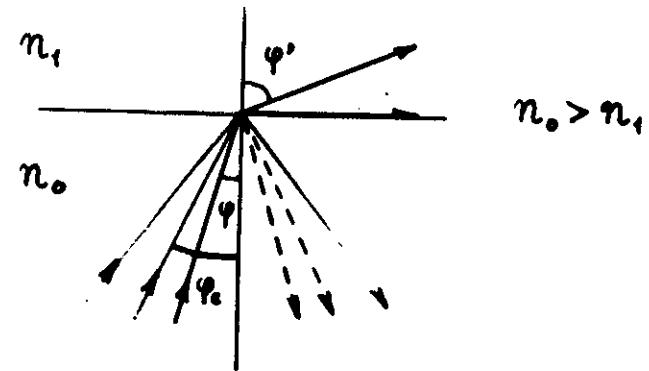
THE OPTICAL FIBRE



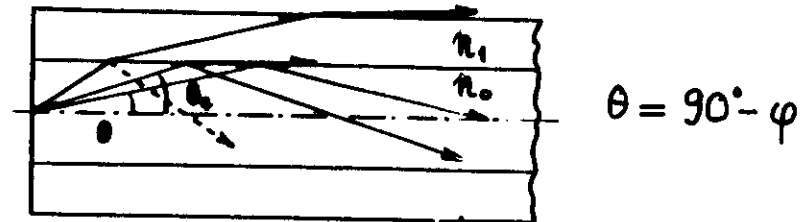
	REFRACTIVE INDEX	DIAMETER
--	------------------	----------

CLADDING	$n_1 \approx 1.46$	$\sim 125 \mu\text{m}$
CORE	$n_1 + 5-10\%$	$\sim 5-10; 50 \mu\text{m}$

TOTAL INTERNAL REFLECTION



- * SNELL'S LAW : $n_0 \sin \varphi = n_1 \sin \varphi'$
- * LIMITING CONDITION: $\varphi' = 90^\circ \Rightarrow n_0 \sin \varphi_c = n_1$
- * TOTAL REFLECTION FOR $\varphi > \varphi_c$



LIMITING ANGLE FOR GUIDED RAYS: $\theta_c = 90^\circ - \varphi_c \Rightarrow \cos \theta_c = \frac{n_1}{n_0}$

GUIDED RAY (TOT. INT. REFLECTION) FOR $\theta < \theta_c$

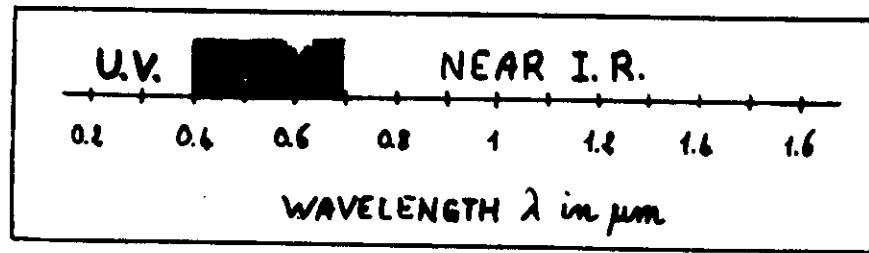
NUMERICAL APERTURE: $A = n_0 \sin \theta_c$

$$A = \sqrt{n_0^2 - n_1^2} \approx \sqrt{2n_1(n_0 - n_1)}$$

($A \sim 0.1-0.2$ typic.)

LIGHT: AN ELECTROMAGNETIC WAVE

WITH $\nu \sim 10^{16} \text{ Hz}$



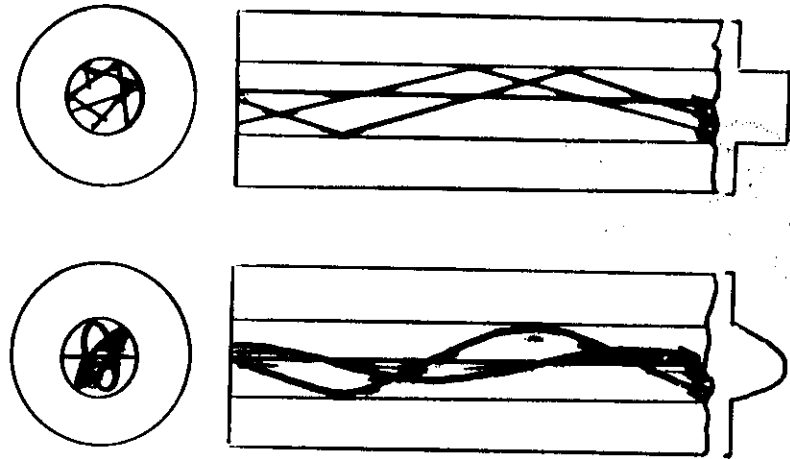
⇒ INTERFERENCE AND DIFFRACTION PHENOMENA ⇒

THE E.M. ENERGY PROPAGATES THROUGH THE FIBRE ACCORDING TO WELL DEFINED WAVES:

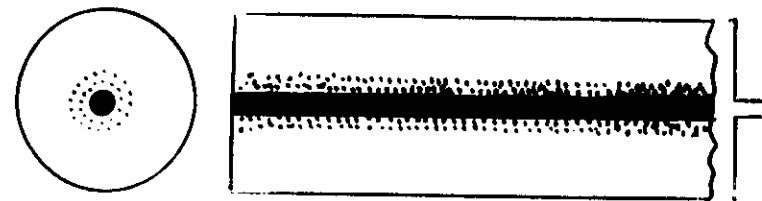


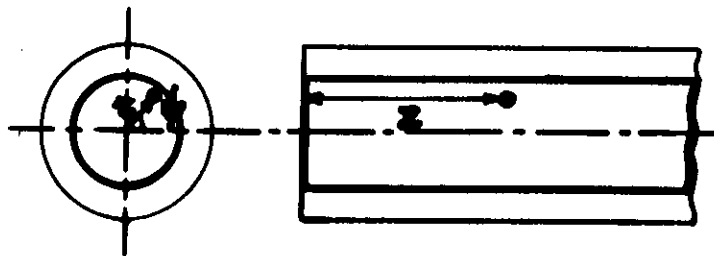
FIBRE TYPES

MULTIMODE FIBRE



SINGLE-MODE FIBRE





$$\begin{cases} \vec{E}(r, \psi, z, t) = \vec{E}(r) \exp[-i(\nu\psi + \beta z - \omega t)] \\ \vec{H}(r, \psi, z, t) = \vec{H}(r) \exp[-i(\nu\psi + \beta z - \omega t)] \end{cases}$$

● PROGRESSIVE WAVE : $\exp[-i(\beta z)]$ ($\beta > 0$)

● AZIMUTHAL PERIODICITY: $\exp[-i(\nu\psi)]$
 $(\nu = 0, 1, 2, \dots)$

RADIAL PART of the FIELD:

$$\begin{cases} \vec{E}(r) = \vec{E}_r(r) + \hat{z} E_z \\ \vec{H}(r) = \vec{H}_r(r) + \hat{z} H_z \end{cases}$$

SCALAR WAVE EQUATION: $\vec{E}_t \rightarrow \phi(r)$

$$\phi''(r) + \frac{\phi'(r)}{r} + \left[\left(\frac{2\pi}{\lambda} \right)^2 n^2(r) - \beta^2 - \left(\frac{\nu \pm 1}{r} \right)^2 \right] \phi(r) = 0$$

$$\begin{cases} E_z = -\frac{i}{k} \left(E'_r \pm \frac{\nu \pm 1}{r} E_r \right) \\ H_z = -\frac{i}{k} \left(H'_r \pm \frac{\nu \pm 1}{r} H_r \right) \end{cases} \quad [k(r) = \frac{2\pi}{\lambda} n(r)]$$

$$E_z(r) = \begin{cases} A_0 \phi_0(r) & , z < a \\ A_1 \phi_1(r) & , z > a \end{cases}$$

CONTINUITY CONDITIONS at $z = a$:

$$\frac{\phi'_0(a)}{\phi_0(a)} = \frac{\phi'_1(a)}{\phi_1(a)} \quad \text{ELECTROMAGNETIC CHARACTERISTIC EQUATION}$$

FOR EACH ν -VALUE $\rightarrow$ TWO FAMILIES OF SOLUTIONS:

$[\nu \pm 1]$



(UPPER SIGN)



(LOWER SIGN)

• $\nu=0$ $\begin{cases} EH_{0\delta} \equiv TE_{0\delta} : E_z = E_r = H_\phi = 0 \\ HE_{0\delta} \equiv TM_{0\delta} : H_z = H_r = E_\phi = 0 \end{cases}$

• $\nu \neq 0$ $\begin{cases} EH_{\nu\delta} : H_z = -i\sqrt{\frac{\epsilon}{\mu}} E_r; E_r = -iE_\phi; H_\phi = iH_r \\ HE_{\nu\delta} : H_z = i\sqrt{\frac{\epsilon}{\mu}} E_r; E_r = iE_\phi; H_\phi = iH_r \end{cases}$

$H_r = -\sqrt{\frac{\epsilon}{\mu}} E_\phi; H_\phi = \sqrt{\frac{\epsilon}{\mu}} E_r$

$|H_z| \ll |H_r|, |H_\phi|; |E_z| \ll |E_r|, |E_\phi|$

$\vec{E} = \sqrt{E_r} \hat{r} + \sqrt{E_\phi} \hat{\phi}$

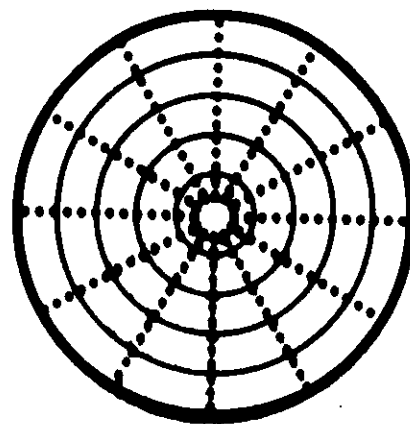
LINES OF FORCE OF $\vec{E} : z(\psi)$

(Real Parts) $\frac{E_r}{E_\phi} = \frac{z'}{z} = \begin{cases} -\tan \nu\psi \rightarrow EH_{\nu\delta} \\ +\tan \nu\psi \rightarrow HE_{\nu\delta} \end{cases}$

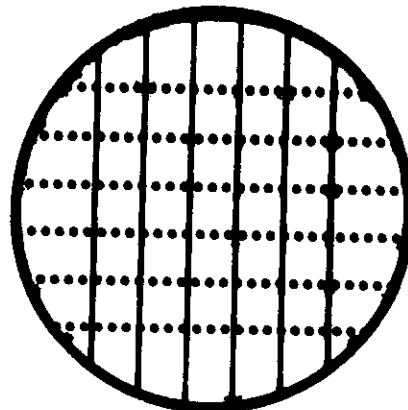
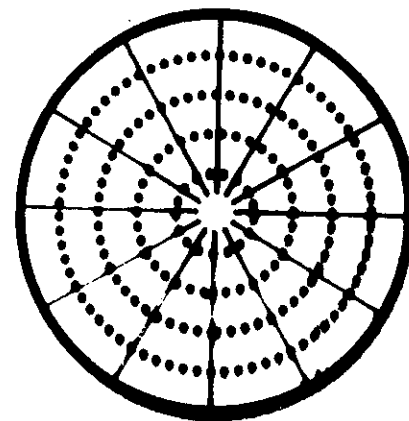
HE

$\frac{H}{E}$

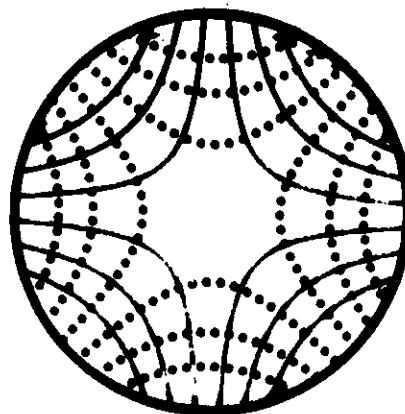
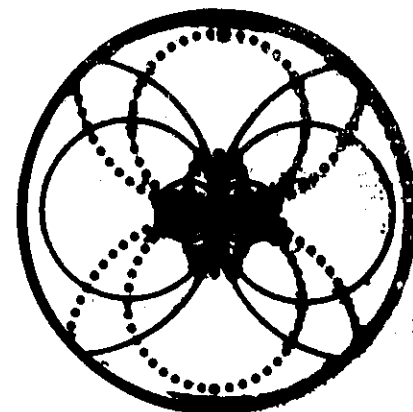
EH



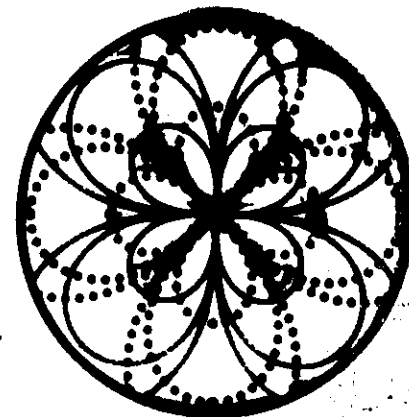
$\nu=0$



$\nu=1$



$\nu=2$



MODE DEGENERACY: LP MODES

- $HE_{\nu+1,\delta}$ and $EH_{\nu-1,\delta}$ FOLLOW THE SAME EQUATION
- ➔ THEY HAVE THE SAME β and PROPERTIES



COMPOSED MODE:



LINEARLY POLARIZED

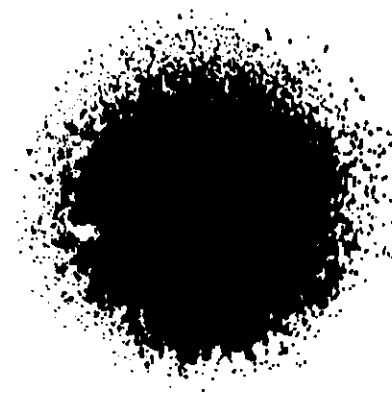
- $\nu=0$ ONLY ONE MODE $LP_{0\delta} \equiv HE_{1\delta}$

- $\nu=1$ THREE MODES COMBINED:

$$LP_{1\delta} \equiv \begin{cases} HE_{2\delta} + TE_{0\delta} \\ HE_{2\delta} + TM_{0\delta} \end{cases}$$

PHYSICAL MEANING OF ν and δ

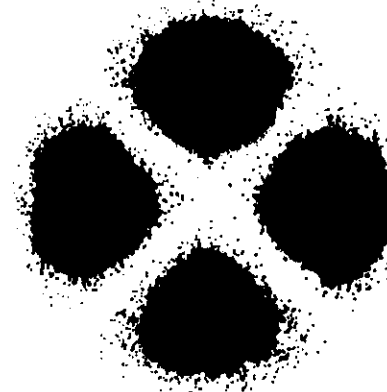
- 2ν NUMBER OF MAXIMA OF THE FIELD INTENSITY IN A TURN AROUND THE FIBRE AXIS
- δ NUMBER OF MAXIMA OF THE FIELD INTENSITY IN RADIAL DIRECTION



LP_{01}
(HE_{11})



LP_{11}
($TE_{01} + TM_{01} + HE_{21}$)



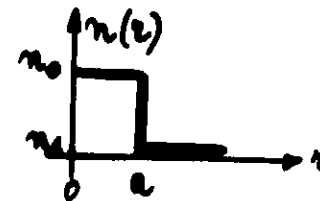
LP_{21}
($EH_{11} + HE_{31}$)



LP_{02}
(HE_{12})

STEP INDEX FIBRES

$$\bullet \quad k^2(z) = \left(\frac{2\pi}{\lambda}\right)^2 n^2(z) \rightarrow n^2(z) = \begin{cases} n_0^2, & z < a \\ n_1^2, & z > a \end{cases}$$



• $LP_{\nu s}$:

$$E_1 = \begin{cases} \phi_0(z) = J_\nu(kz), & z < a & - [k = \sqrt{(\frac{2\pi}{\lambda})^2 - \gamma^2}] \\ \phi_1(z) = K_\nu(\gamma z), & z > a & - [\gamma = \sqrt{(\frac{2\pi}{\lambda})^2 - k^2}] \end{cases}$$

$$\bullet \quad a^2(k^2 + \gamma^2) = V^2$$

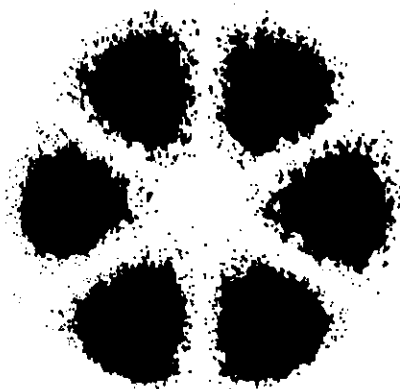
$$V = \frac{2\pi a}{\lambda} \sqrt{n_0^2 - n_1^2}$$

NORMALIZED
FREQUENCY

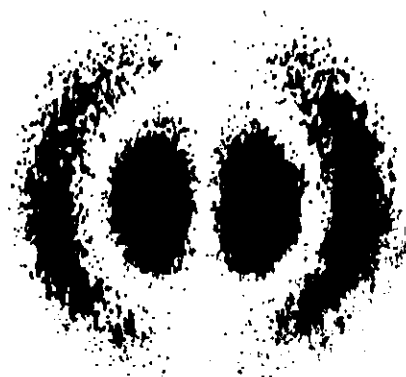
CHARACTERISTIC EQUATION

$$\frac{\phi_0'(a)}{\phi_0(a)} = \frac{\phi_1'(a)}{\phi_1(a)}$$

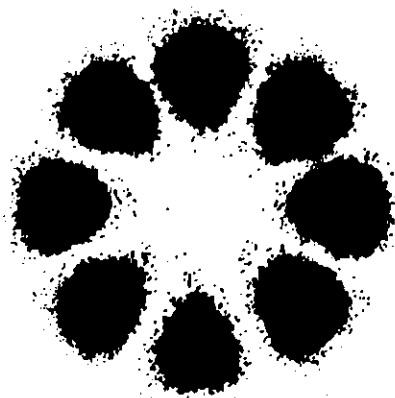
► $\beta(\gamma)$ for each $LP_{\nu s}$ mode



LP_{31}
($EH_{21} + HE_{41}$)



LP_{12}
($TE_{02} + TM_{02} + HE_{22}$)

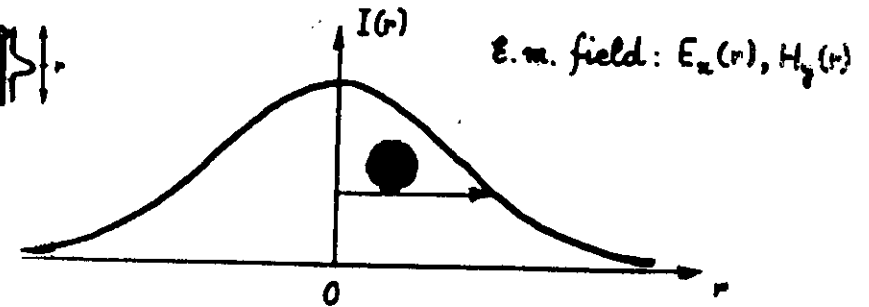


LP_{41}
($EH_{31} + HE_{51}$)



LP_{22}
($EH_{12} + HE_{32}$)

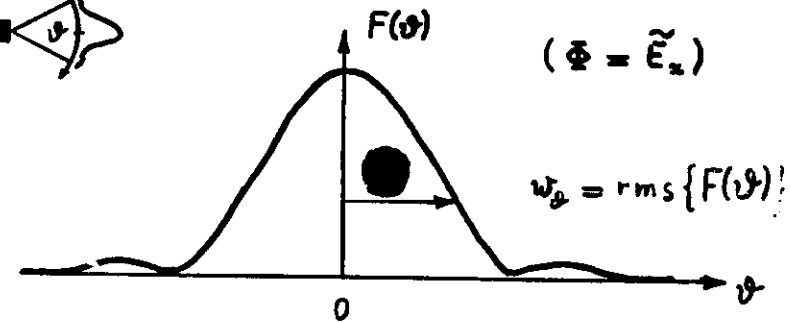
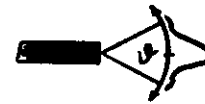
NEAR-FIELD INTENSITY: $I(r) \propto E_x^2(r), H_y^2(r)$



* Mode (field) radius (spot-size) 1.



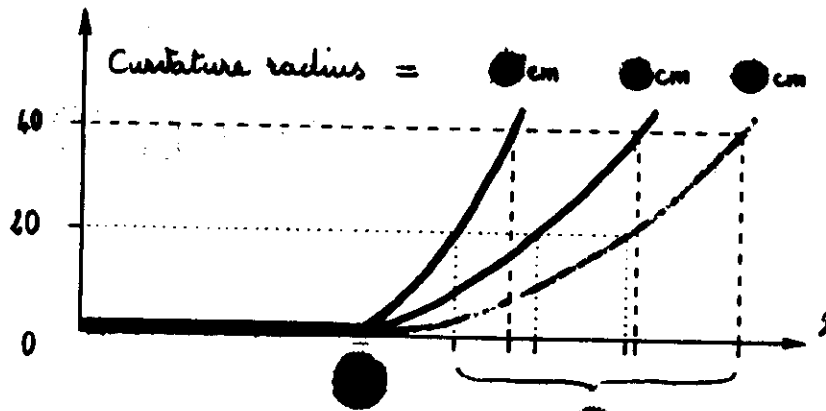
FAR-FIELD INTENSITY: $F(\vartheta) \propto |\Phi(\vartheta)|^2$



* Mode radius (spot-size) 2.



Higher-order radiation loss (dB/m)



CCITT Normalization:

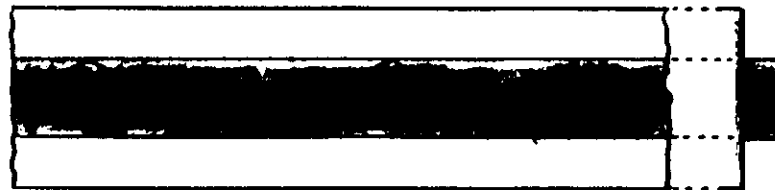
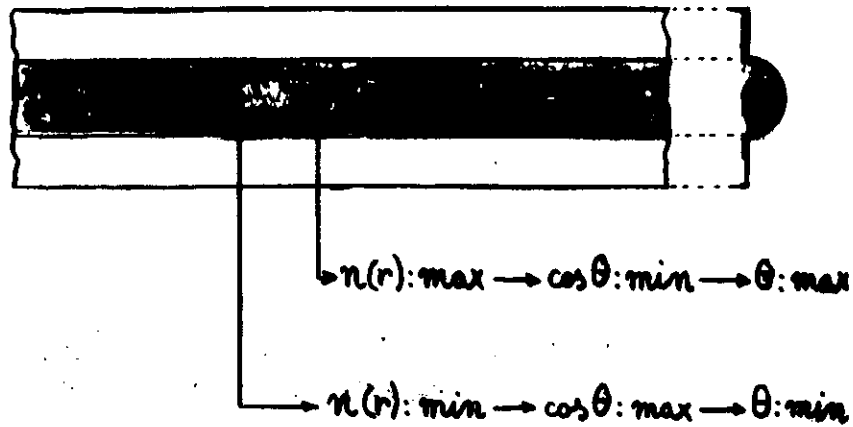
Curvature radius = $\bullet$ cm

Sample length = $\bullet$ m

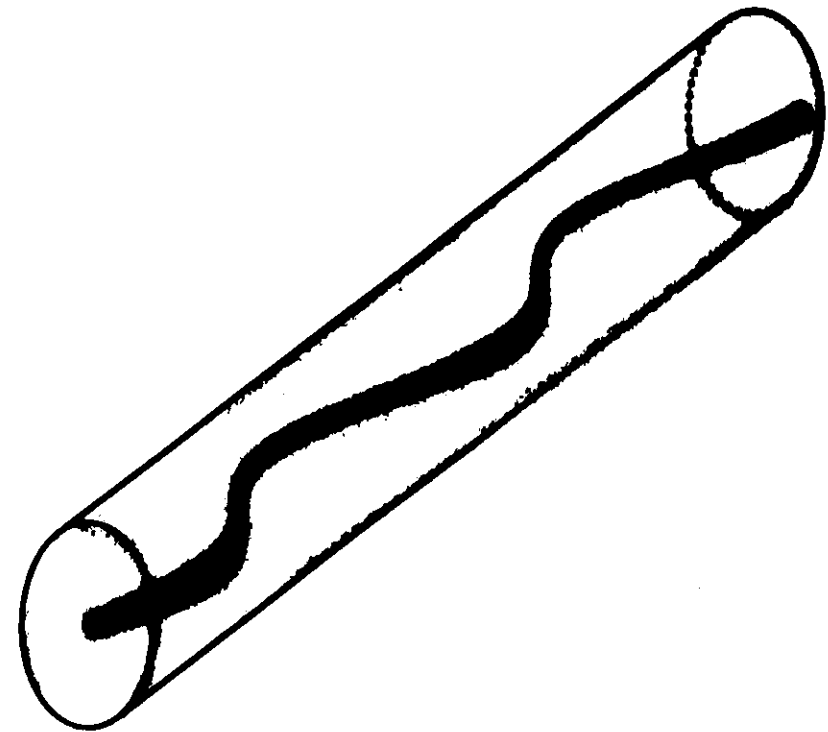
MULTIMODE FIBRES

18

$$\beta = \frac{2\pi}{\lambda} n(r) \cos \theta = \text{constant}$$



$$n(r) = \text{const.} \rightarrow \theta = \text{const.}$$



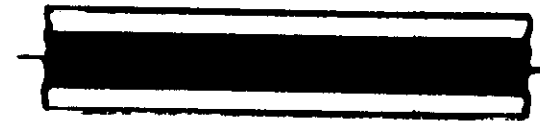
RAGGIO MERIDIANO IN FIBRA
A PROFILO PARABOLICO

CONGRUENZA DI RAGGI

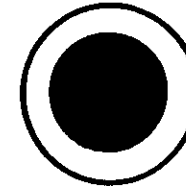
(INSIEME DEI RAGGI CON STESSI v e β)

● OGNI RAGGIO OTTENUTO DA UN ALTRO MEDIANTE:

• TRASLAZIONE LUNGO L'ASSE DELLA FIBRA

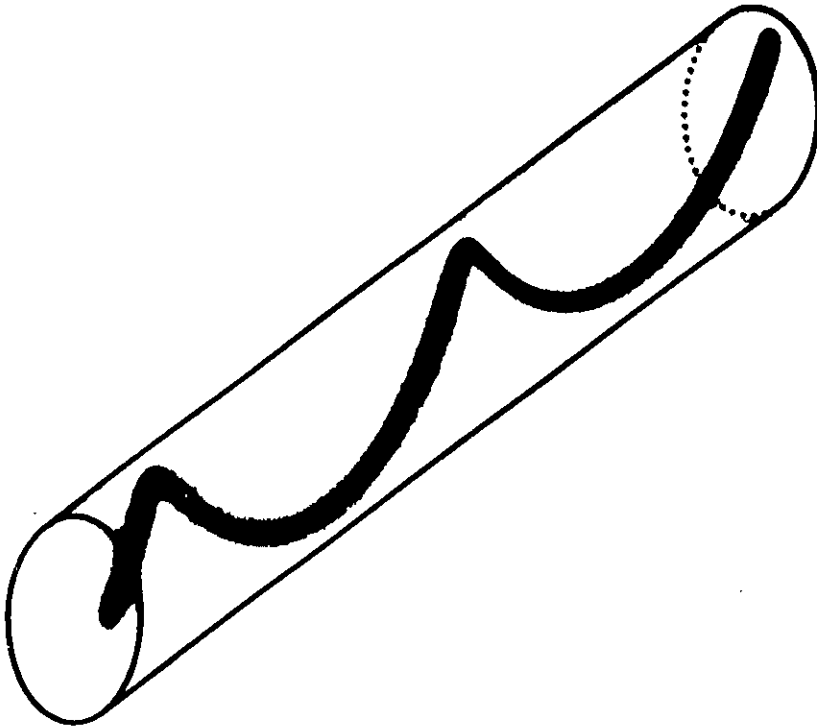


• ROTAZIONE ATTORNO ALL'ASSE DELLA FIBRA

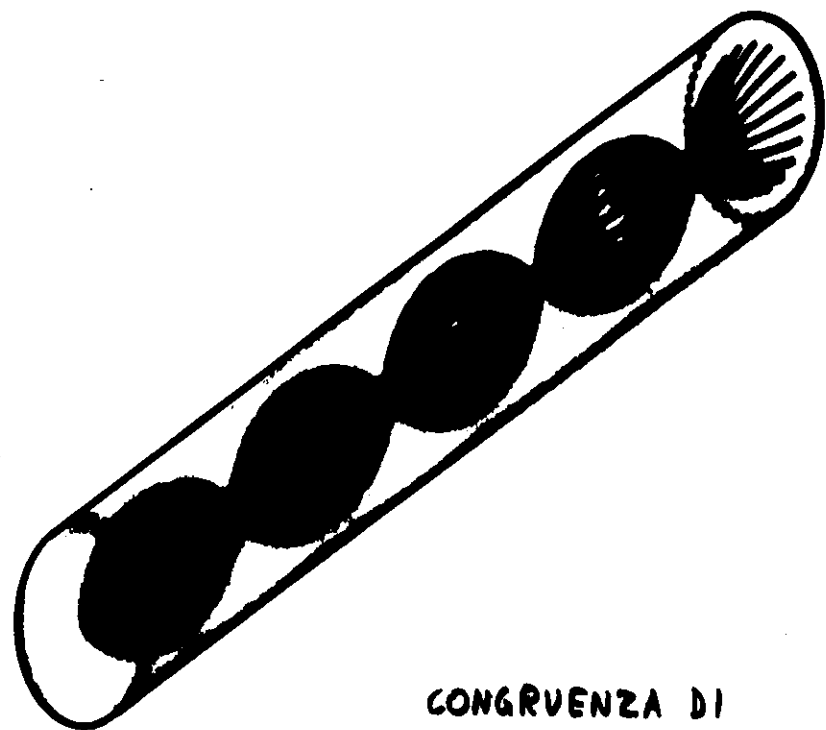


● TALI RAGGI HANNO TUTTE
LE STESSA PROPRIETA':

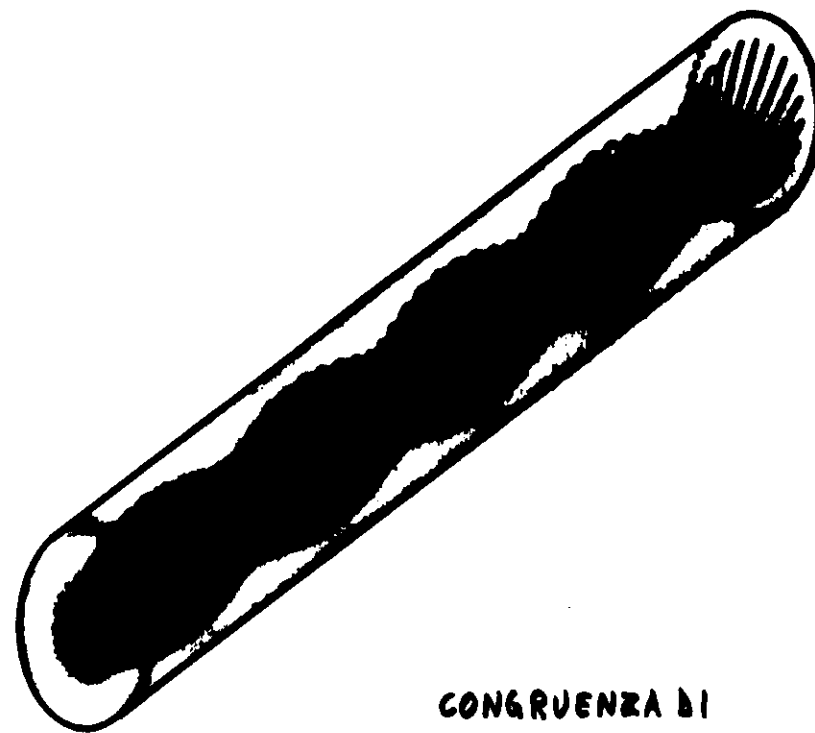
▷ PER $\lambda \ll a$ ($V \gg 1$) CORRISPONDE:
A UN MODO $\leftrightarrow$ UNA CONGRUENZA DI RAGGI



RAGGIO SGHEMBO IN FIBRA
A PROFILO PARABOLICO



CONGRUENZA DI
RAGGI MERIDIANI IN FIBRA
A PROFILO PARABOLICO



CONGRUENZA DI
RAGGI SGHEMBI IN FIBRA
A PROFILO PARABOLICO

FIBRE NUMERICAL APERTURE

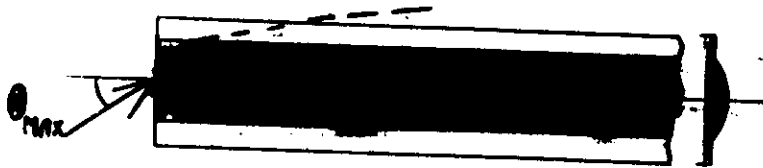
$$A = n \cdot \sin \theta_{\max}$$

$$(\text{CUTOFF COND.: } \beta = \frac{2\pi}{\lambda} n(r) \cos \theta_{\max} = \frac{2\pi}{\lambda} n_{\min})$$

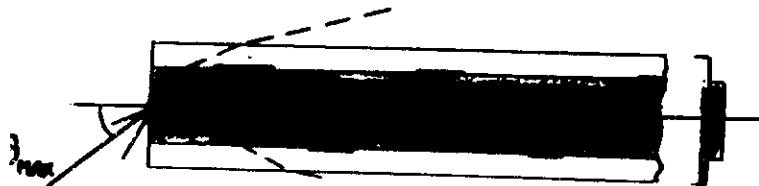


$$A(r) = \sqrt{n^2(r) - n_{\min}^2}$$

$$(A_M = \sqrt{n_{\max}^2 - n_{\min}^2} = \begin{cases} \sim 0.1 \text{ SINGLE-MODE} \\ \sim 0.2 \text{ MULTIMODE} \end{cases})$$



GRADED-INDEX PROF. FIBRES: CONTIN. INTERN. REFRACTION



STEP-INDEX PROF. FIBRES: TOTAL INTERN. REFLECTION



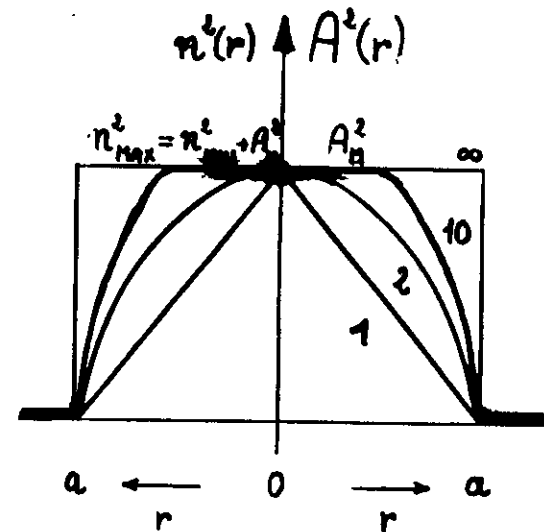
$$V = \frac{2\pi}{\lambda} a A_M$$

* "α"-PROFILES:

$$n^2(r) = n_{\min}^2 + A_M^2 \left[1 - \left(\frac{r}{a} \right)^\alpha \right] \quad r \leq a$$

• $\alpha = 2$: PARABOLIC PROFILE

• $\alpha \rightarrow \infty$: STEP PROFILE

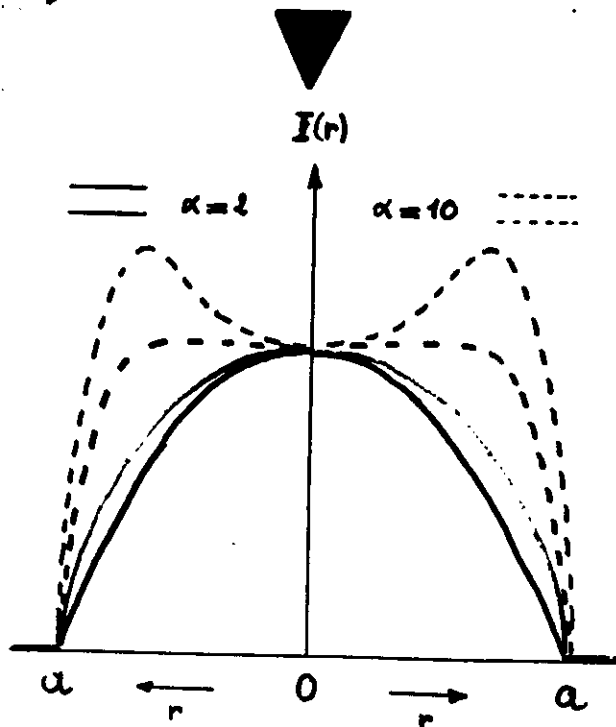


NEAR-FIELD INTENSITY (MULTIMODE FIBRES WITH $R(x) = R_0$)

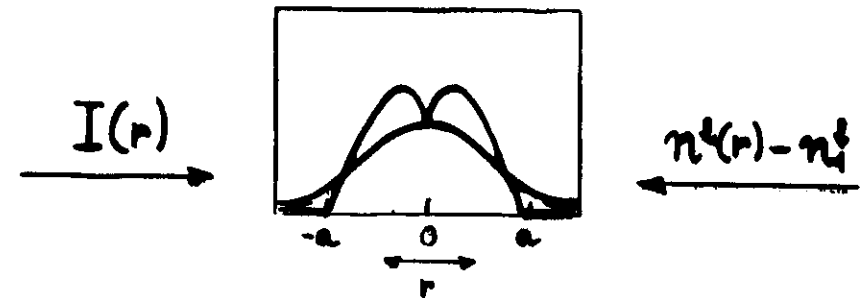
• $I(r) \propto [n^2(r) - n_{\text{MIN}}^2]$: GUIDED RAYS ONLY

• $I(r) \propto \left(\frac{n^2(r) - n_{\text{MIN}}^2}{\sqrt{1 - r^2/a^2}} \right)$: LEAKY RAYS TOO

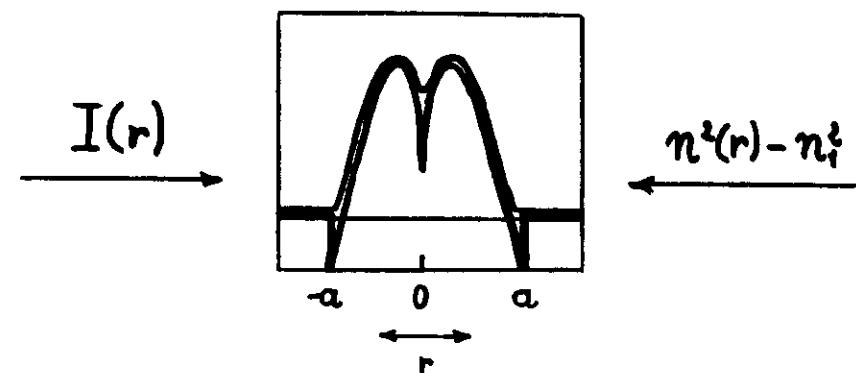
➡ MEASUREMENT OF $n^2(r)$



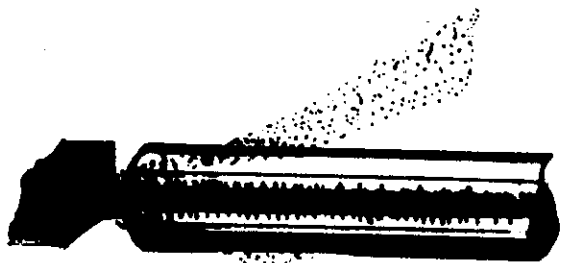
* FEW-MODE FIBRES



* FIBRES OF PARTICULAR PROFILES



LAUNCHING AND COUPLING PROBLEMS



$$W_e = W_o + W_i$$

EMITTED POWER $\rightarrow$ GUIDED POWER + RADIATED (LOST) POWER

* LAUNCHING EFFICIENCY: $\Lambda = W_o / W_e$ (% typ.)

* COUPLING LOSS: $L = W_i / W_o$ (dB typ.)

SINGLE-MODE FIBRES:

$I_e(r)$ - Emitted power radial distribution

Λ DEPENDS ON THE "LIKENESS" OF $I_e(r)$ AND $I(r)$

$$\Lambda = [\text{Overlap integral between } I_e \text{ and } I]^2$$

* Typical values:

SOURCE-FIBRE COUPLING

* LAUNCHING IN MULTIMODE FIBRES

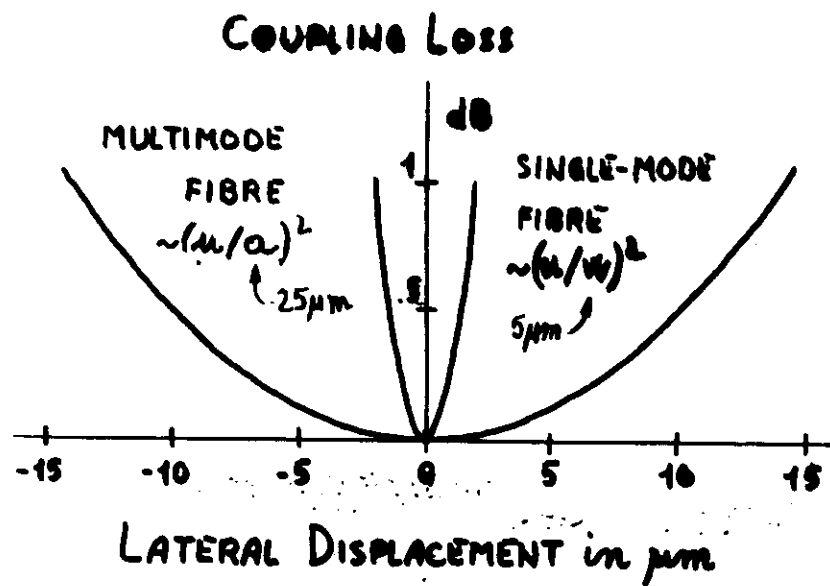
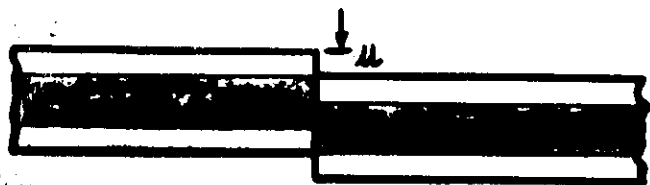
$$\Lambda \propto \frac{A_m^2}{A_{\text{source}}^2} \cdot R_e \cdot f(\text{PROF.}) \cdot a^2$$

$$f(\text{PROF.}) = \begin{cases} 1/2 & \text{PARABOLIC PROFILE} \\ 1 & \text{STEP PROFILE} \end{cases}$$

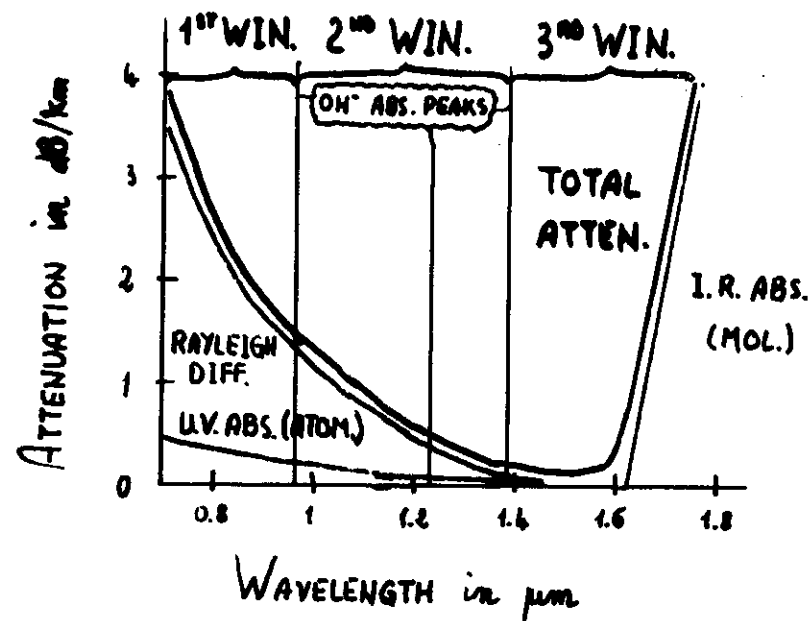
* TYPICAL VALUES:

$$\Lambda = \begin{cases} \sim 2\% & \text{LED SOURCES} \\ \sim 50\% & \text{SEMICONDUCTOR LASER SOURCES} \end{cases}$$

COUPLING PROBLEMS



ATTENUATION INTRINSIC FACTORS



ATTENUATION EXTRINSIC FACTORS

* ION ABSORPTION:

OH^- , Fe^{++} , Ni^{+++} , Cu^{++} , Cr^{+++} , Mn^{+++} , ...

(Loss: 1-5 dB/km per part per billion)

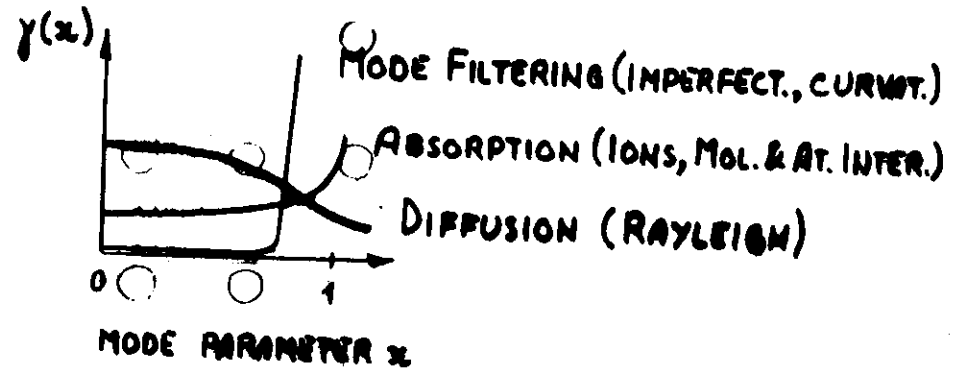
* DIFFUSION BY OPTICAL INHOMOGENEITIES

* MODE-FILTERING BY PHYSICAL IMPERFECTIONS

* MODE-CONVERSION BY MICROBENDING

POWER ATTENUATION

* EACH MODE IS ATTENUATED IN A DIFFERENT WAY: $\gamma(x)$ - MODE ATTENUATION CONSTANT

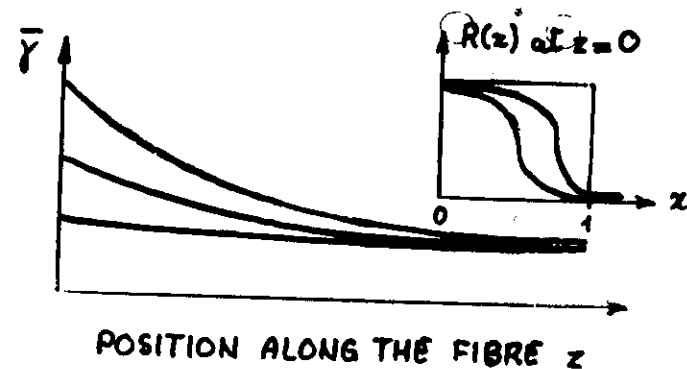


* AVERAGE ATTENUATION: $\bar{\gamma} \equiv \langle \gamma(x) \rangle$ IN $R(z)$

• $\bar{\gamma}$ DEPENDS ON LAUNCHING CONDITIONS: $R(z)$ AT $z=0$

• $\bar{\gamma}$ DEPENDS ON z -COORDINATE TOO

for large z : $\bar{\gamma} \rightarrow \gamma_{\min}(x)$



MODE CONVERSION

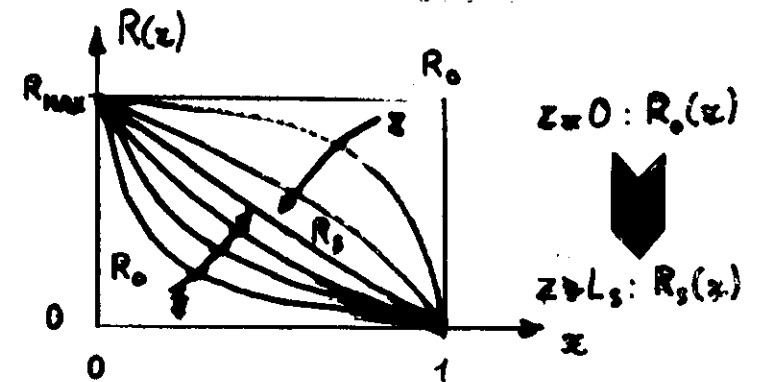
- CAUSED BY MICROBENDING,
SMALL IMPERFECTIONS,
.....
- PRODUCES: OPTICAL LOSSES,
(CONVERS. TO RADIAT. MODES)
MODIFICATION ON DISTRIB.
(CONVERS. VS. GUIDED MODES)

SINGLE-MODE FIBRES

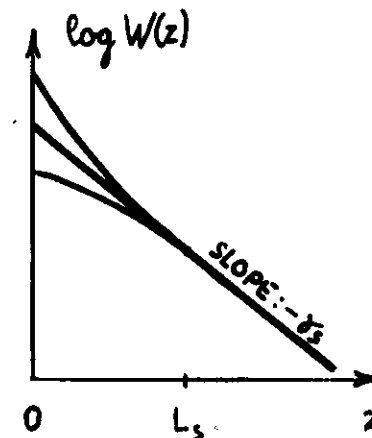
- MODE CONVERSION PRODUCES LOSSES ONLY
- LOSS COEFFICIENT $\propto W^p$ ($p \approx 2-4$)
(INCREASING FUNCTION OF MODE RADIUS)

MULTIMODE FIBRES

- * for every launched $R(z)$, a steady state is reached after a length L_s : $R_s(z)$, where each mode has the same attenuation: γ_s



- * $R_s(z)$: steady state mode distribution
- * L_s : steady state length
- * γ_s : steady state attenuation coefficient



for $z \gg L_s$:

$$R(z) \sim R_s(z) \cdot e^{-\gamma_s z}$$



$$W(z) \sim e^{-\gamma_s z}$$

INFORMATION TRANSMISSION

CAPACITY

PULSE DISPERSION

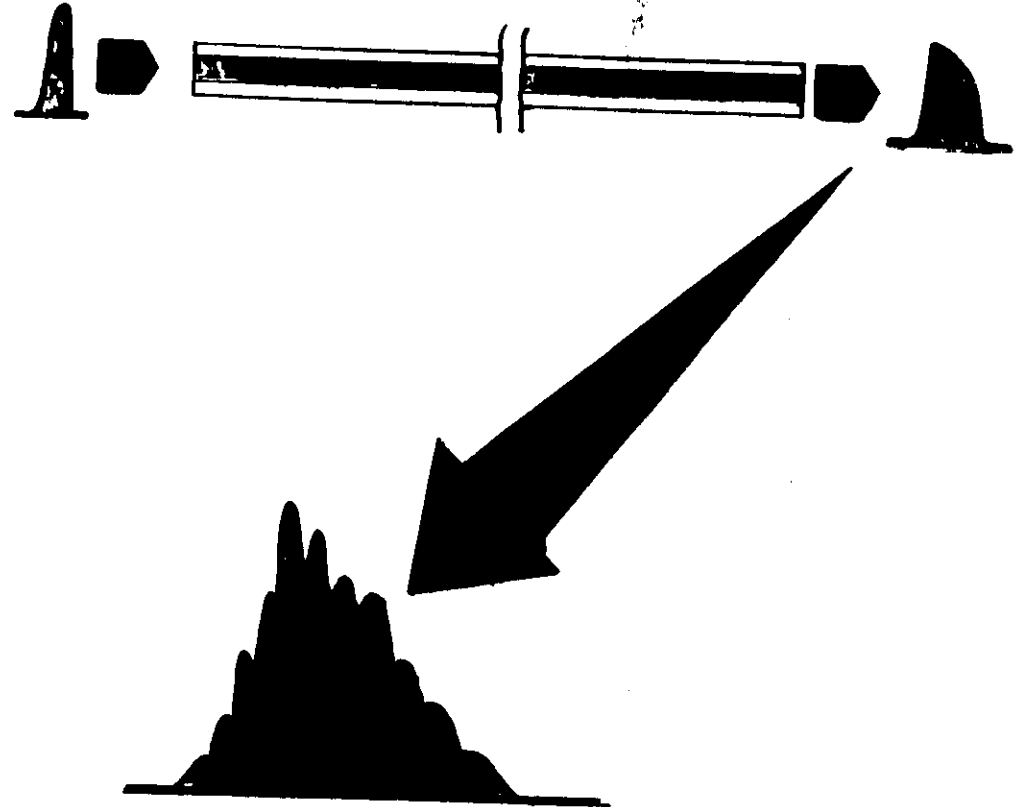
* MODAL DISPERSION

- MODE COUPLING
- MODE FILTERING
- OPTIMIZED PROFILE

* CHROMATIC DISPERSION

- WAVEGUIDE DISPERSION
- MATERIAL DISPERSION
- ZERO-DISPERSION

PULSE DISTORTION



* travelling time depends on mode:

► INTERMODAL DISPERSION

* travelling time depends on λ :

► INTRAMODAL DISPERSION

CHROMATIC DISPERSION

MODE TRAVELLING TIME DEPENDS ON λ



τ : TRAVELLING TIME IN 1km OF FIBRE

$D = d\tau/d\lambda$: CHROMATIC DISPERSION

PULSE BROADENING:

$$\Delta t = L \cdot \Delta \lambda \cdot D$$

L : FIBRE LENGTH
 $\Delta \lambda$: SPECTRA WIDTH OF THE OPT. SOURCE

BANDWIDTH OF SINGLE-MODE FIBRES

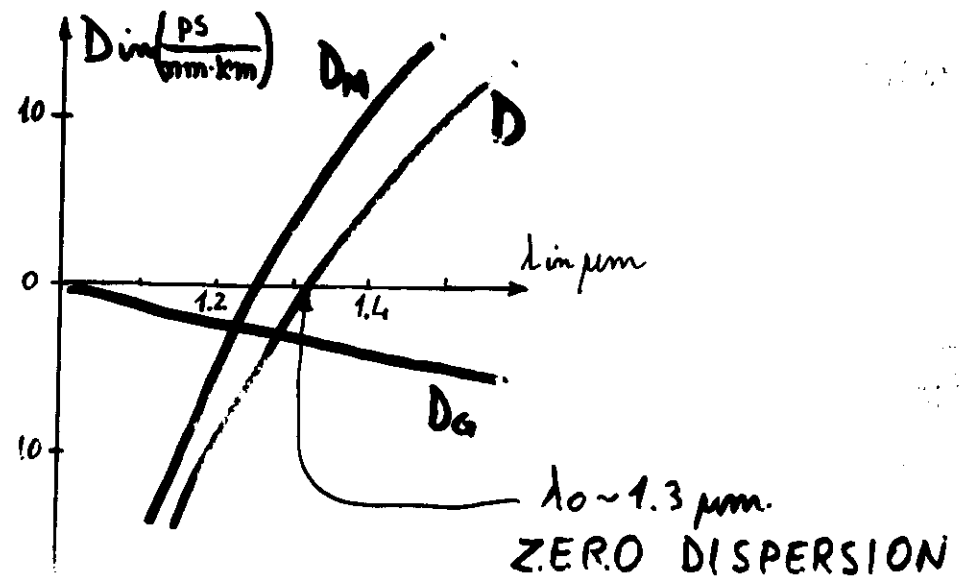
	λ in μm	D in $\frac{\text{ps}}{\text{nm} \cdot \text{km}}$	BW in $\text{GHz} \cdot \text{km}$
LTIMODE SER: $\Delta \lambda \sim 2 \text{ nm}$	1.3	~ 1	~ 100
	1.5	~ 10	~ 10
RIF-MCMF CFB: $\Delta \lambda \sim 1.4 \text{ nm}$	1.3	~ 1	~ 2000
	1.5	~ 10	~ 200

• τ DEPENDS ON λ :

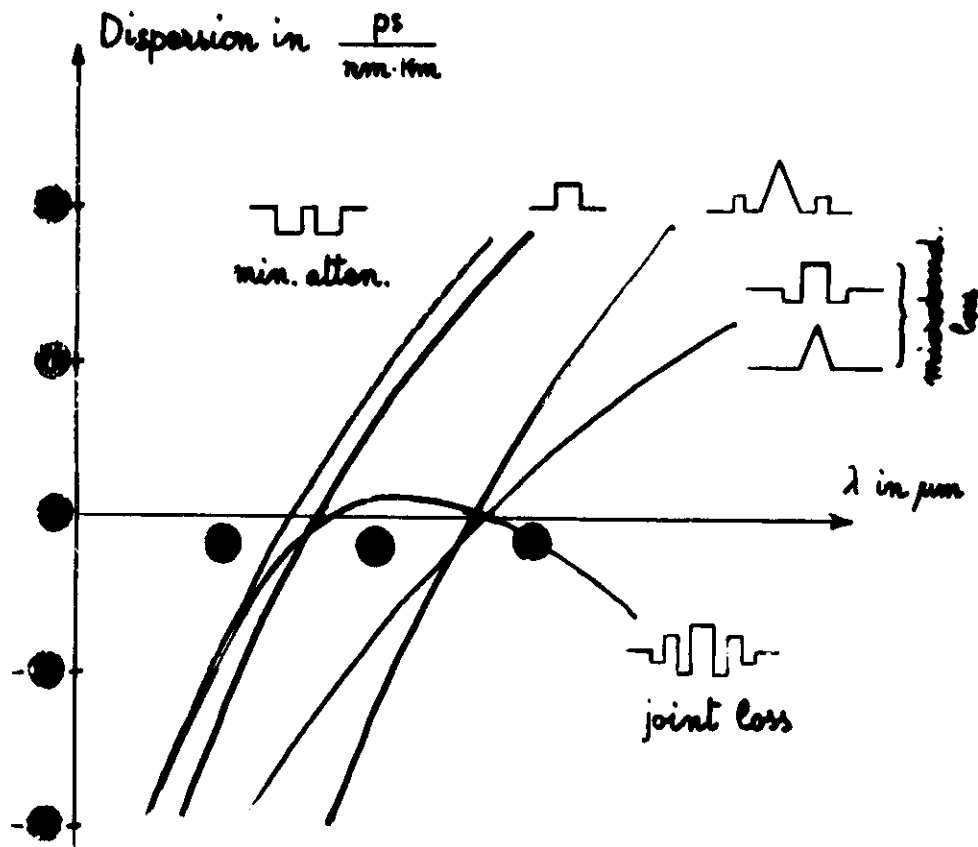
1. THROUGH THE GUIDE EFFECT
2. BECAUSE THE MATERIAL IS DISPERSIVE ($n = n(\lambda)$)

➡ D HAS TWO CONTRIBUTIONS:

1. GUIDE DISPERSION: D_g
2. MATERIAL DISPERSION: D_m

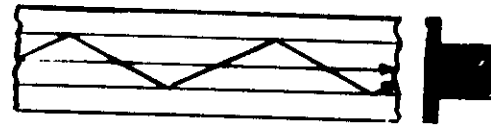


DISPERSION OPTIMIZATION



INTERMODAL DISPERSION (TRAVELLING TIME DEPENDS ON MODES)

* STEP PROFILE

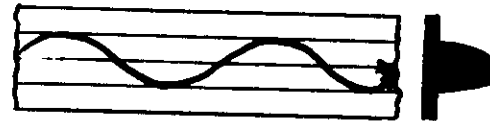


$$\Delta t \propto (\text{NUM. APERT.})^2$$

$$\sim 40 \text{ ns/km}$$

$$BW \sim 5 \text{ MHz} \cdot \text{km}$$

* PARABOLIC PROFILE



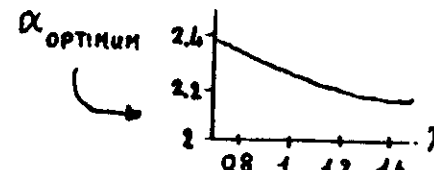
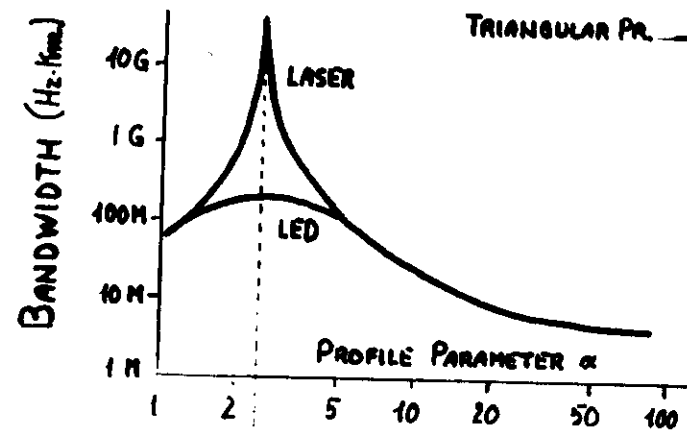
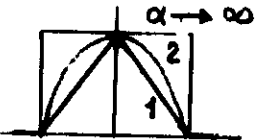
$$\Delta t \propto (\text{NUM. APERT.})^4$$

$$\sim .4 \text{ ns/km}$$

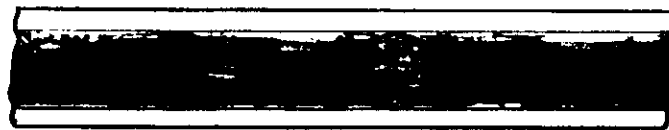
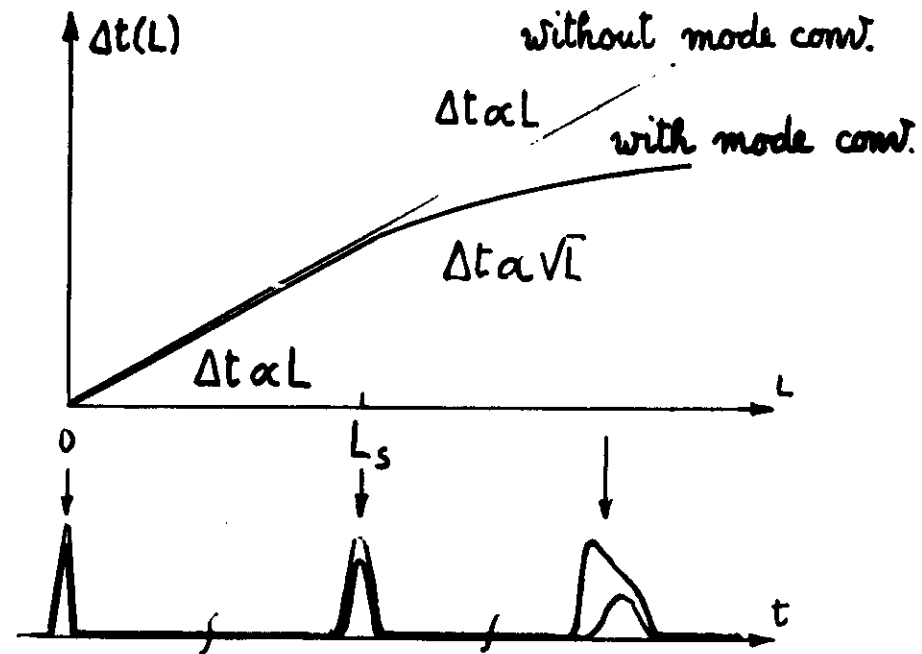
$$BW \sim 500 \text{ MHz} \cdot \text{km}$$

* "α" - PROFILES

STEP PROFILE
PARABOLIC PROF.
TRIANGULAR PR.



MODE CONVERSION EFFECTS ON PULSE DISTORTION



* Statistical equalization of optical paths