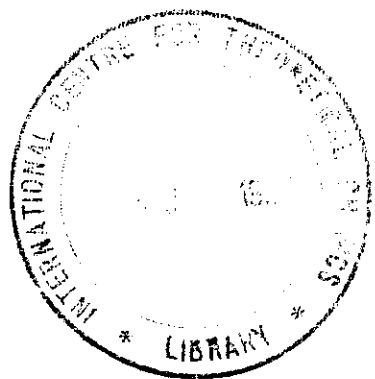


INTERNATIONAL ATOMIC ENERGY AGENCY
UNITED NATIONS EDUCATIONAL, SCIENTIFIC AND CULTURAL ORGANIZATION



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SPRING COLLEGE ON PLASMA PHYSICS

(25 May - 19 June 1987)

ROTATIONAL INSTABILITIES
IN THE
FIELD - REVERSED CONFIGURATION

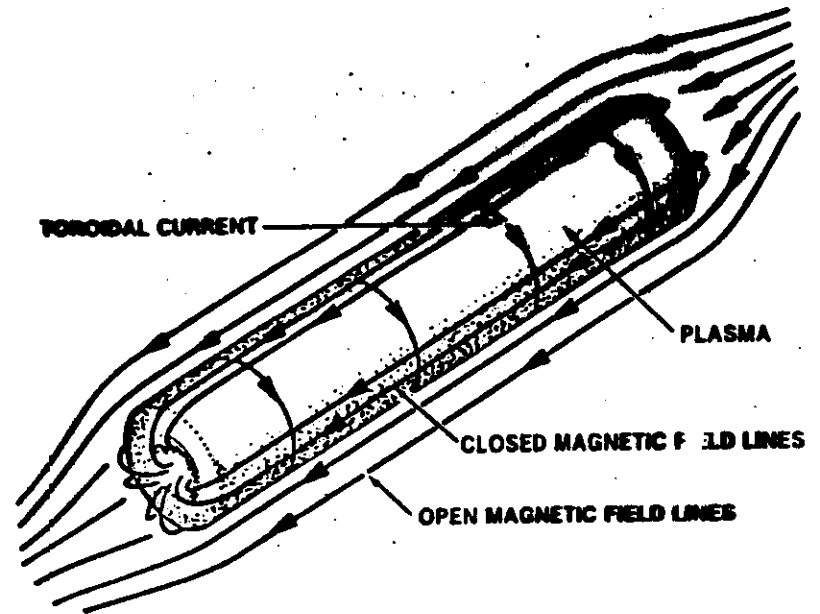
D.S. HARNED
S.A.I.C.
San Diego
U.S.A.

ROTATIONAL INSTABILITIES IN THE FIELD-REVERSED CONFIGURATION

D.S. Harned - SAIC, San Diego

Acknowledgment: FRX experimental group, Los Alamos Natl. Lab.

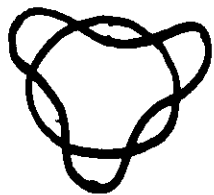
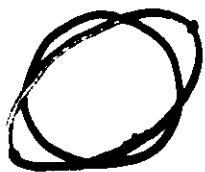
FIELD-REVERSED CONFIGURATION (FRC)



EXPERIMENTS ARE GENERALLY OBSERVED TO:

- begin rotating
- are disrupted by $n=2$ rotational instability (without multipole stabilization)

$n=2$



$n=3$

MHD Theory predicts:

— instability for any amount of rotation

— large n grows fastest, all n unstable $\gamma = \sqrt{n-1} \Omega$

(Taylor, 1962)

Define: $\alpha = \frac{-\Omega_i}{\Omega_e - \Omega_i} = \frac{-\Omega_i}{\Omega_*}$

Finite Larmor radius theory predicts:

— $n > 2$ modes stable

— $n=2$ instability for $\alpha > 1.6$
(Freidberg & Pearlstein, 1978)

But, experiments observe instability onset
at $\alpha = 0.4$ to 1.0 ³

FLR theory:

— good approximation for non-reversed theta pinch

— poor approximation for FRC due to field null

Need kinetic model for ions

∴ use quasineutral hybrid computational model

QUASINEUTRAL HYBRID MODEL

(Byers et al, 1978; Hameiri, 1982)

- Assume quasineutrality
- Neglect transverse displacement current (Darwin)
- Ion particles
- Electrons - inertialess fluid
- Fully nonlinear

EQUATIONS :

Ampere's law and electron momentum eq. :

$$\vec{E} = \frac{1}{4\pi n_e} (\nabla \times \vec{B}) \times \vec{B} - \frac{1}{n_e e} \vec{j} \times \vec{B} - \frac{1}{n_e} \nabla (n_e T_e) + \frac{c^2}{4\pi} \nabla \times \vec{B}$$

Faraday's law :

$$\frac{\partial \vec{B}}{\partial t} = -c \nabla \times \vec{E}$$

Particle equations of motion :

$$\frac{d\vec{v}}{dt} = \frac{q}{m} \left(\vec{E} + \frac{\vec{v}}{c} \times \vec{B} \right)$$

$$\frac{d\vec{x}}{dt} = \vec{v}$$

→ Forces on particles and source EXAMPLE:

terms from linear weighting (PIC)

$$\alpha = 1.0, \beta_0 = 0.49, \Omega_* = 0.05 \omega_{ci}$$

→ Advance fields with predictor-corrector method

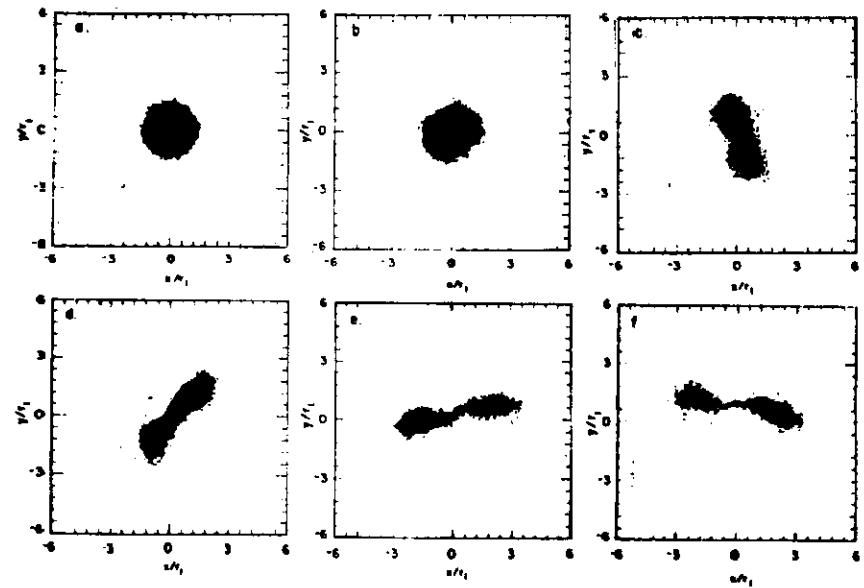
$$S = \rho/p_0 = 8.3 \quad (\text{like FRX-R})$$

Should be stable from FLR theory

→ Vacuum fields matched across a transition region

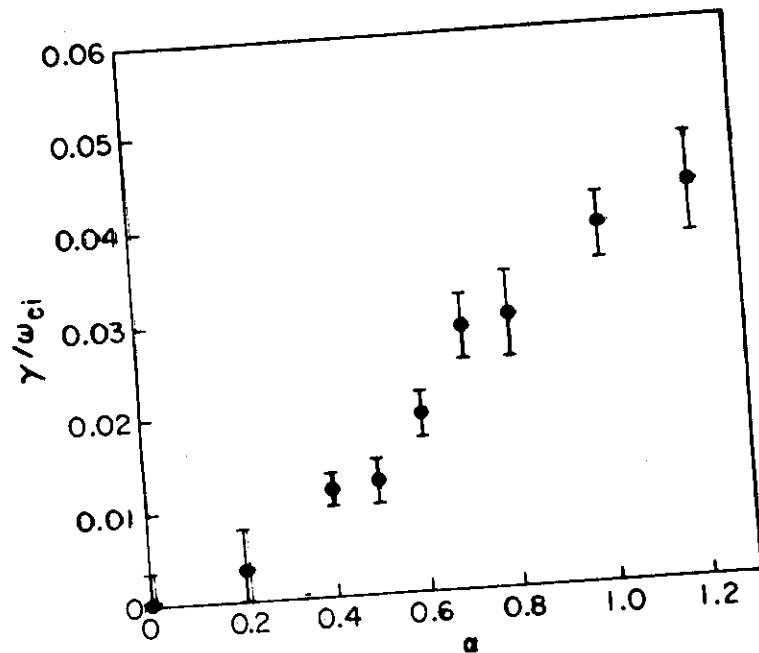
50,000 particles, 100x100 grid

$$\rightarrow 2\frac{1}{2}d, \frac{d}{\delta z} = 0$$



Growth rates

- slightly less than MHD
- $n > 2$ stable

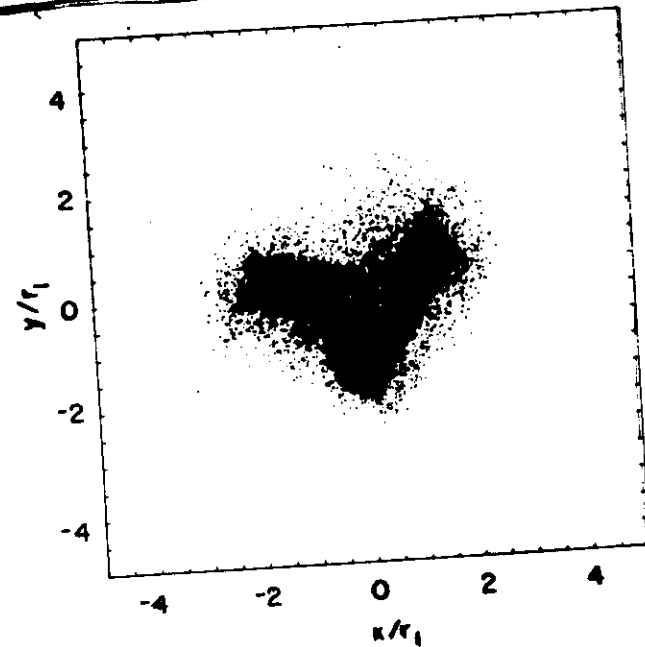


- No threshold observed as in FLR

For strongly reversed cases higher
 n modes can be unstable

$$\Omega_* = 0.054, \quad \beta_0 = 0.23, \quad \alpha = 1.1$$

$n=3$ instability



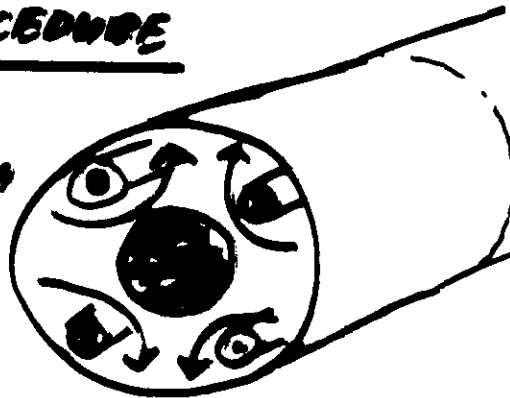
MULTIPOLE FIELDS

Can stabilize $n=2$ mode

(Ohi et al, 1983; Hoffman et al, 1983;
Siemon et al, 1986)

EXPERIMENTAL PROCEDURE

- Form FRC plasma
- Apply multipole fields, which are excluded from the main part of the plasma on experimental timescales
- $n=2$ instability is not observed if the multipole field strength exceeds a threshold value



THEORY

MHD - (Ishimura, 1984)

- assumed circular cross-section
- ignored mode coupling
- Predicts stability if quadrupole field at plasma edge satisfies
$$B_q^2 > r_s^2 \Omega^2 \mu_0 \rho_0$$
- B_q is a factor of ≈ 4 larger than experimental observation

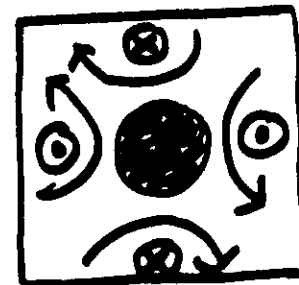
MUST CONSIDER:

- 1.) Modification of equilibrium due to quadrupole field
- should be stabilizing
(Spencer + Tuszewski, 1985)
- 2.) Kinetic treatment of ions
- substantially changes mode and growth rates (Hornat, 1983)

SIMULATION PROCEDURE

- 1.) $t=0$, initialize rotating truncated rigid rotor distribution ($\rho \rightarrow 0$ outside r_s) with no quadrupole field

- 2.) Turn on Ioffe bars, increasing linearly in time until $B_{quad} = B_z$ at the plasma edge



- 3.) Quadrupole field does not penetrate appreciably inside separatrix because plasma is highly conducting on these time scales

FRX-C CROSS-SECTION WITH QUADRUPOLE FIELDS (End-on holography) (K. McKenna)



Equilibrium has a large amplitude
 $n=4$ distortion

FRX-C/T Experimental

Parameters: (Roj et al, 1981)

$$B_0 = 4 \text{ kG}, \beta_s = 0.31, x_s = 0.6, S = \frac{R}{\beta_0} = 13.1$$

$$\alpha = -\frac{D_i}{\sqrt{r_s}} = 0.81$$

Simulations use:

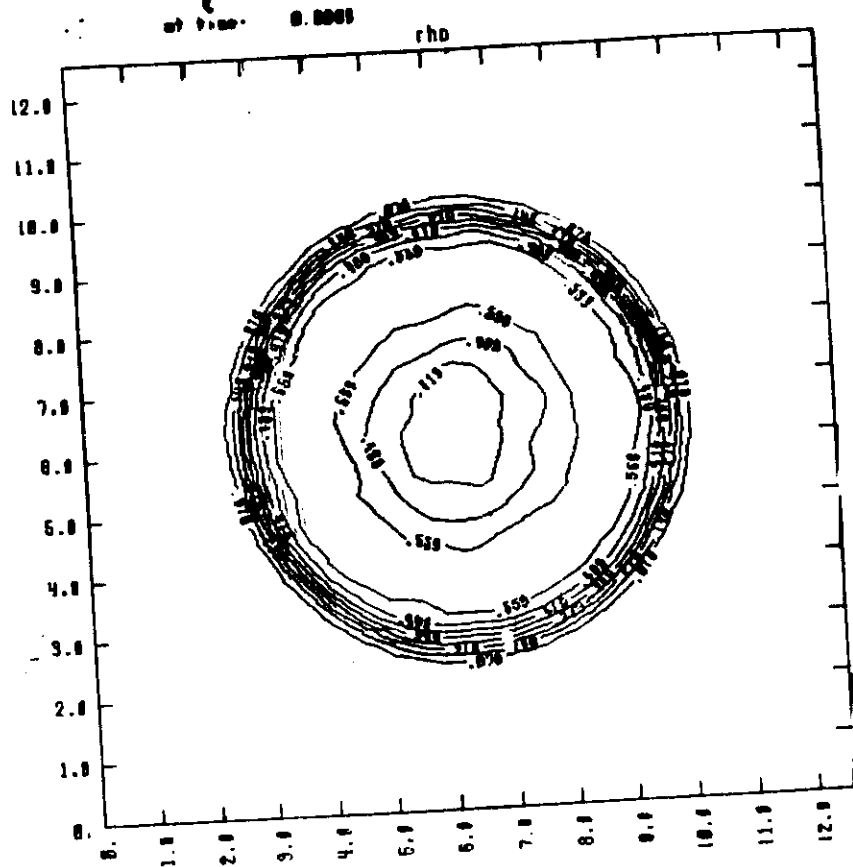
50,000 - 100,000 particles

80x80 grids

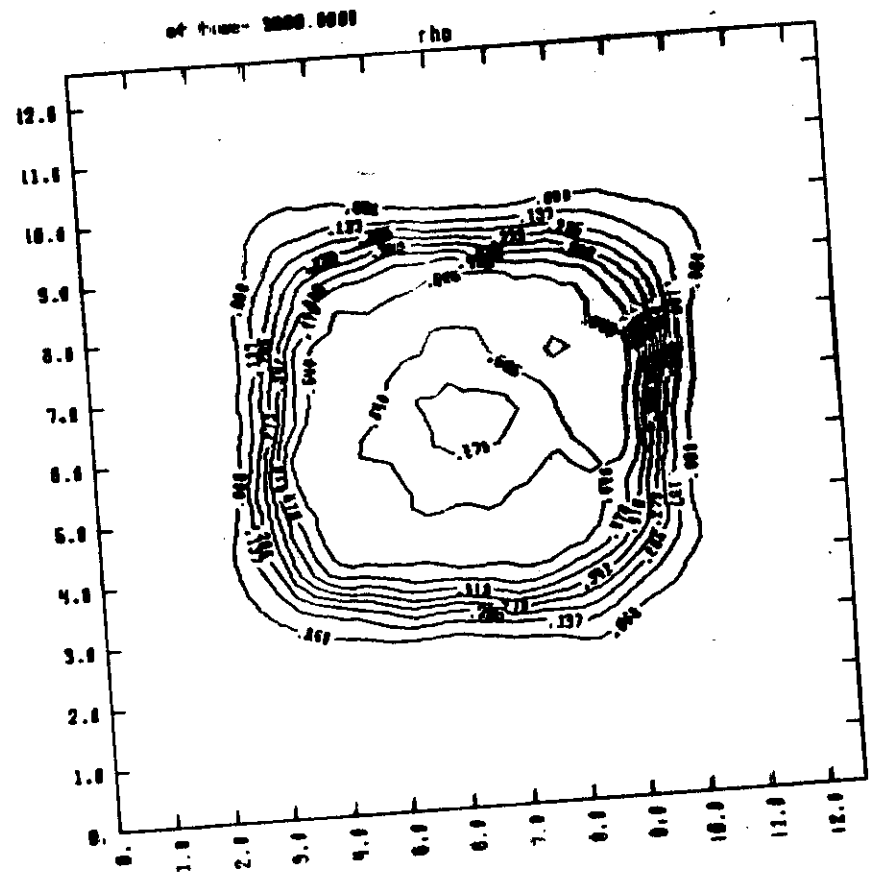
≡ Vary quadrupole field strength, Q_z ,
applied to this axisymmetric
equilibrium

EXAMPLE:

Density contours, initial axisymmetric equilibrium



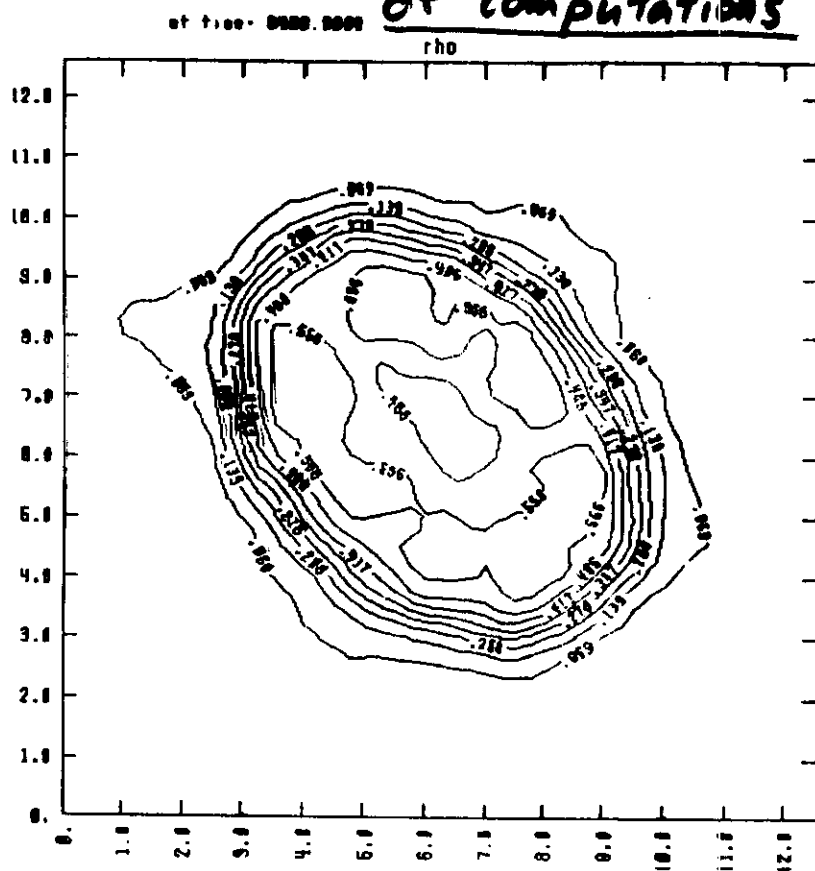
- Quadrupole field has been applied
- Equilibrium develops a large amplitude $n=4$ distortion



(Note: coils have been rotated 45° from experimental positions)

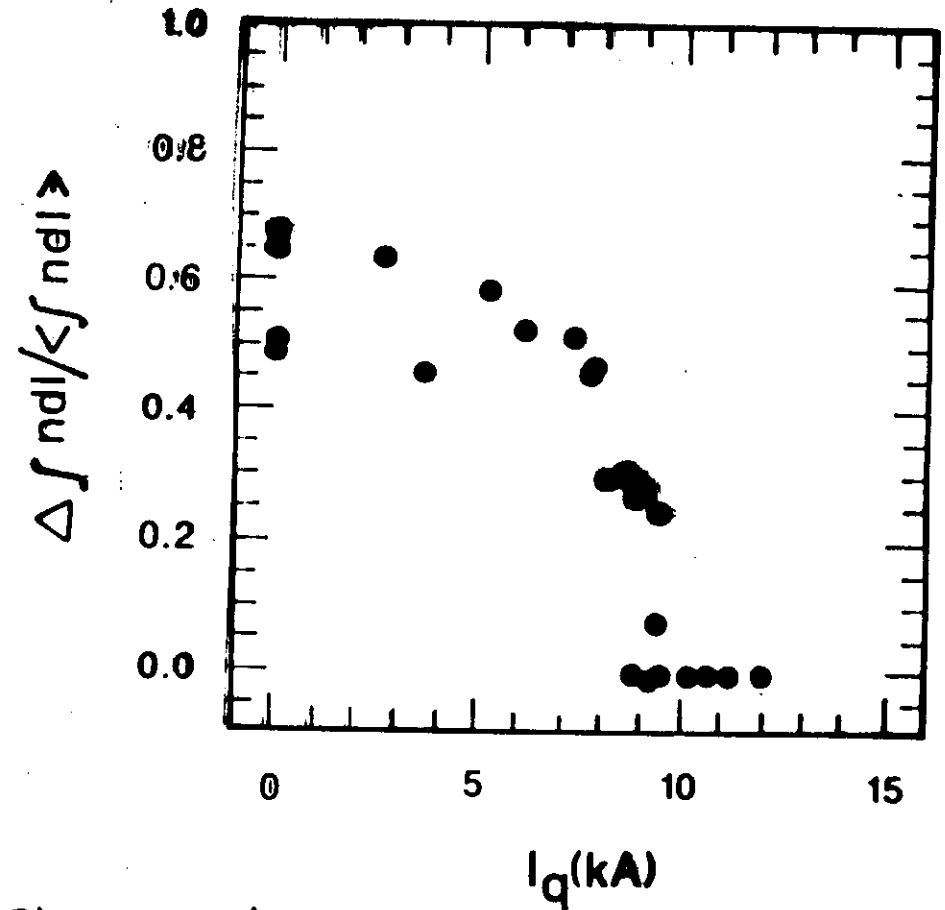
$$B_q/B_0 = 0.13$$

= just below stability threshold
of computations



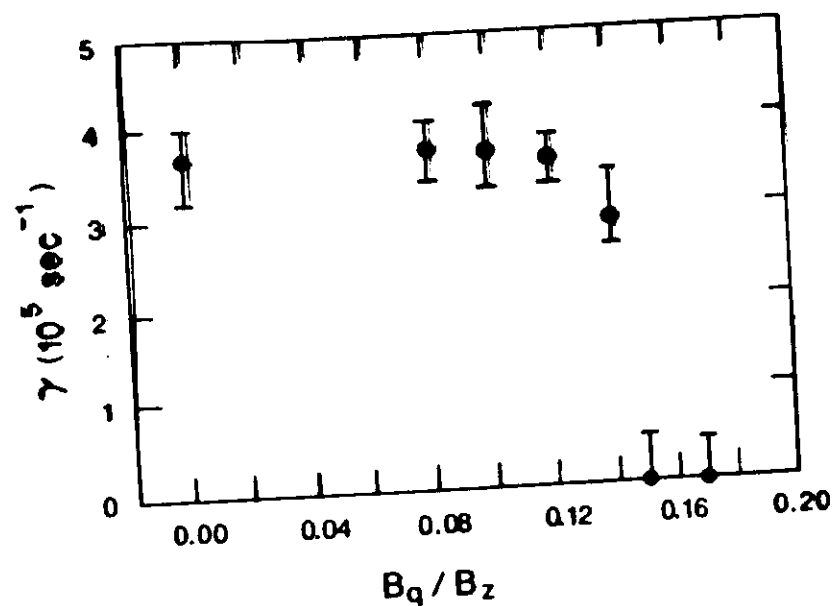
$n=2$ instability appears

FRX-C/T Integrated line density
measurements



Shows sharp threshold for
quadrupole field to stabilize the
 $n=2$ mode

Growth rates from hybrid simulations



BOTH EXPERIMENT AND SIMULATIONS SHOW:

- Sharp threshold for stabilization of $n=2$ mode, well below axial field strength
- Mode grows essentially unaffected by quadrupole field for B_q below threshold
- Equilibria develop large amplitude $n=4$ distortions

DIFFERENCE:

- Experimental threshold $B_q \approx 0.08 B_0$
- Simulation threshold $B_q \approx 0.15 B_0$

POSSIBLE SOURCES OF DISCREPANCY

- 1.) Equilibrium in expt. is probably not a rigid rotor, since
 $\beta_s = 0.31$ and $x_s = 0.6$ do not satisfy the avg. beta condition,
 $\langle \beta \rangle = 1 - x_s^2 / 2$ for a rigid rotor
- 2.) Equilibrium in simulation should have more time to relax to the presence of the quadrupole field

CONCLUSIONS

- 1.) Kinetic effects are essential to model FRC rotational instability (FRC is NOT sufficient)
- 2.) No rotation threshold for $n=2$ instability
- 3.) Simulations show large amplitude $n=4$ equilibrium distortions due to quadrupole field
- 4.) Simulations show stabilization of $n=2$ by quadrupole field above threshold
($B_q > 0.15 B_z$ for FRC-CIT-like parameters)

