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AUTUMN COURSE ON APPLICATIONS OF ANALYSIS TO MECHANICS

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FUNCTIONAL ANALYTIC METHODS IN ELASTIC-PLASTIC
TORSION OF CYLINDRICAL BARS

(Part II)

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These are preliminary lecture notes intended for participants only.
Copies are available outside the Publications Office (T-floor) or from
room 112.

II. A general boundary value problem with Hencky's law

II.1 Recapitulation

We want to resolve several problems about elasto-plastic torsion of cylindrical bars:

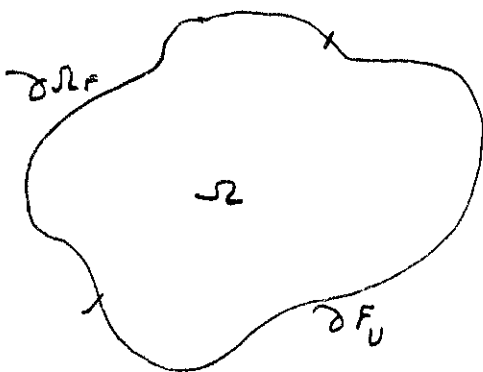
Homogeneous bars with simply connected cross sections
multiply connected cross sections.

Fiber reinforced bars with simply connected cross sections

Let us consider these problems like some particular cases of a more general one that we are going to study in this Chapter

For reasons of simplicity we shall adopt the Hencky's constitutive law - that will be justified later on by the fact that for increasing loads (extensively) we shall be able to prove that, in the case of the torsion problem, the solutions obtained by Hencky's and Prandtl-Reuss's laws are the same.

II.2 Description of the problem (A)



Ω is here an arbitrary, bounded open set of $\mathbb{R}^3$ with a "regular" boundary $\partial\Omega$

(In fact we can allow $\partial\Omega$ to be piecewise continuously differentiable)

The data are:

$$\begin{cases} f, & \text{body forces in } \Omega. \\ \bar{F}_h & \text{on } \partial\Omega_{F_h} \subset \partial\Omega. \\ U_e & \text{on } \partial\Omega_{U_e} \subset \partial\Omega. \end{cases}$$

(f , F_h , $\partial\Omega_{F_h}$ and $\partial\Omega_{U_e}$ being such that the problem is well posed)

We shall call $\partial\Omega_F$ the union of the parts of $\partial\Omega$ on which some F_h is given and $\partial\Omega_U$ the union of the parts of $\partial\Omega$ on which some U_e is given.

Then we have to resolve the following problem.

Problem (A):

To find the couple (σ, u) such that:

$$\left. \begin{aligned} (15) \quad \sigma_{ij} &= \sigma_{ji} \text{ in } \Omega \\ (16) \quad \sigma_{ij,j} + f_i &= 0 \text{ in } \Omega \end{aligned} \right\} \text{ (equilibrium)}$$

$$\left. \begin{aligned} (17) \quad \sigma_{hj} n_j &= \bar{F}_h \text{ on } \partial\Omega_h \\ (18) \quad u_L &= u_e \text{ on } \partial\Omega_e \end{aligned} \right\} \text{ Boundary condition.}$$

$$\left. \begin{aligned} (19) \quad \mathcal{F}(\sigma) &\leq 0 \\ (20) \quad \varepsilon_{ij}(u) &= A_{ijkl} \sigma_{kl} + \lambda_{ij} \\ (21) \quad \lambda_{ij} (z_{ij} - \sigma_{ij}) &\leq 0 \quad \forall \sigma \in R^9 \text{ s.t.} \\ &\quad z_{ij} = \sigma_{ji} \text{ and } \mathcal{F}(\sigma) \leq 0 \end{aligned} \right\} \text{ Hensky's constitutive law}$$

With:

$$(22) \quad \mathcal{F}: R^9 \longrightarrow R, \text{ convex and continuous (Elasticity criterion)}$$

$$(23) \quad A_{ijkl}: \Omega \longrightarrow R, \text{ bounded on } \Omega$$

$$(24) \quad A_{ijkl} = A_{jikh} = A_{khij}.$$

$$(25) \quad \exists \beta > 0 \text{ constant such that}$$

$$A_{ijkl} x_{ij} x_{kl} \geq \beta |x|^2.$$

$$\forall x \in R^9 \text{ with } x_{ij} = x_{ji}$$

components of the elasticity tensor.

$$(26) \quad \varepsilon_{ij}(u) = \frac{1}{2} (u_{ij} + u_{ji}) = \varepsilon_{ji}(u), \text{ linearized strain tensor.}$$

II. 3. Principle of Haas Karman and variational formulation for (A)

It is very easy to prove the following statement.

Principle of Haas Karman: if the problem (A) has a sufficiently regular solution (σ^0, u^0) then σ^0 minimizes the functional

$$(27) \quad J(\sigma) = \frac{1}{2} \int_{\Omega} A_{ijkl} \sigma_{ij} \sigma_{kl} dx - \int_{\partial\Omega_h} \sigma_{hj} n_j u_e d\Gamma.$$

To prove that, we have to write down the difference $J(\sigma) - J(\sigma_0)$, using the properties (15) (24) (25) and also the fact that

$$\begin{aligned} \int_{\Omega} \varepsilon_{ij}(u_0) (\sigma_{ij} - \sigma_{ij}^0) dx &= \int_{\Omega} u_{ij}^0 (\sigma_{ij} - \sigma_{ij}^0) dx = \int_{\partial\Omega} u_i^0 (\sigma_{ij} - \sigma_{ij}^0) n_j d\Gamma \\ &\quad - \int_{\Omega} u_i^0 (\sigma_{ij,j} - \sigma_{ij,j}^0) dx \\ &= \int_{\partial\Omega_0} u_i (\sigma_{ij} - \sigma_{ij}^0) n_j d\Gamma \end{aligned}$$

because the properties (18) for u_i^0 and (15) (16) (17) (19) for σ and σ^0 . We have also to take account of the fact that u_0 and σ_0 , solution of (A) are related by (20) and (21)

We have now to choose a functional framework in order to give a meaning to $J(\sigma)$

• (23) implies that $A_{ijkl} \in L^\infty(\Omega)$ therefore,

• $\int_{\Omega} A_{ijkl} \sigma_{ij} \tau_{kl} dx$ will have a meaning if $\sigma_{ij} \in L^2(\Omega) \quad \forall i, j = 1, 2, 3$

• Then, (20) suggests to look for $\varepsilon_{ij}(u) \in L^2(\Omega)$ and also $\lambda_{ij} \in L^2(\Omega)$.

• Hence, u_i can be reasonably supposed in $H^1(\Omega)^*$ and then $\gamma_0 u_i \in H^{1/2}(\partial\Omega)^*$.

• The second integrals of $J(\sigma)$ will have now a meaning if and only if $\sigma_{ij} n_j|_{\partial\Omega} \in H^{-1/2}(\partial\Omega)$. But we assumed $\sigma_{ij} \in L^2(\Omega)$ and σ is supposed to satisfy in particular (16) that is to say

$$\sigma_{ij,j} = -f_i$$

Now, we know from the theorem 7 (lecture notes H. Lanchon p 20) that if $\sigma_{ij} \in L^2(\Omega)$ and $\sigma_{ij,j} \in L^2(\Omega)$, then $\sigma_{ij} n_j|_{\partial\Omega} \in H^{-1/2}(\partial\Omega)$. It is therefore sufficient to make the following assumptions on the data:

$$(28) \quad f_i \in L^2(\Omega)$$

$$(29) \quad \bar{F}_R \in H^{-1/2}(\partial\Omega)$$

$$(30) \quad u_i \in H^{1/2}(\partial\Omega),$$

Recalls (cf lecture notes H. Lanchon p 14 and following)

$$H^1(\Omega) = \{v / v \in L^2(\Omega), v_{,j} \in L^2(\Omega) \text{ for } j=1, 2, 3\}$$

$$H^{1/2}(\partial\Omega) = \text{space of restrictions at } \partial\Omega \text{ of all the elements of } H^1(\Omega) \left. \begin{array}{l} \text{Sobolev} \\ \text{Space} \end{array} \right\}$$

to obtain a convenient functional scope for the problem (A)

Let us introduce some more definitions

(31) $\mathcal{K} = \{ \sigma / \sigma \in [L^2(\Omega)]^9 \text{ verifying (15)(16)(17)(19)} \}$ is the set of allowed stresses for the problem (A)

(32) $a(\sigma, \tau) = \int_{\Omega} A_{ijkl} \sigma_{ij} \tau_{kl} dx$ is a bilinear form over $[L^2(\Omega)]^9$

We are now able to write down clearly a variational principle for the stress field.

Variational principle: If (σ^0, u^0) is a solution of (A) Then σ^0 is solution of a new problem that we shall call "(A₁)".

(33) $\sigma^0 \in \mathcal{K}$ and $J(\sigma^0) \leq J(\sigma) \quad \forall \sigma \in \mathcal{K}$ with

(34) $J(\sigma) = a(\sigma, \sigma) - 2 \int_{\partial \Omega_0} \sigma_{ij} n_j u_i dP$

II.4 Results and conclusions about the problem (A₁)

We have already studied the general problem of minimizing a functional and this preceding one is already prepared to be interpreted as a variational Inequality (or "Inequation"?) - We have just to check now the properties of $\mathcal{K}$, $a(\cdot, \cdot)$ and $\int_{\partial \Omega_0} \sigma_{ij} n_j u_i dP$ to be able to apply the theorem 15 (lecture notes of H. Lancoson p41)

i) $\mathcal{K}$ is a closed convex set of $[L^2(\Omega)]^9$, the norm of this space being defined by the scalar product:

$$(\varphi, \psi) = \int_{\Omega} \varphi_{ij} \psi_{ij} dx$$

In fact, let $\sigma^{(m)}$ be a Cauchy sequence in $\mathcal{K}$ then

• $\sigma_{ij,j}^{(m)}$ converges then toward $\sigma_{ij,j} \in \mathcal{D}'(\Omega)$ because the continuity of the operator $\frac{\partial}{\partial x_j}$ in $\mathcal{D}'(\Omega)$ (cf lecture notes H.L p 13)

moreover: $\sigma_{ij,j}^{(m)} = -f_i$ implies $\lim \sigma_{ij,j}^{(m)} = \sigma_{ij,j} = -f_i \in L^2(\Omega)$

• now, we already saw above that $\sigma_{ij} \in L^2(\Omega)$ and $\sigma_{ij,j} \in L^2(\Omega)$ implies $\sigma_{ij} n_j \in H^{-1/2}(\partial\Omega)$ - we have now to check that:

$$\sigma_{ij} n_j = \bar{F}_i \text{ over } \partial\Omega_{F_h}$$

For that, let ϕ be an element of $H^{1/2}(\partial\Omega)$; we know that $\exists \psi \in H^1(\Omega)$ unique such that:

$$\begin{cases} \Delta \psi = 0 \text{ in } \Omega \\ \psi|_{\partial\Omega} = \phi \end{cases} \quad (\text{cf lecture notes H.L p 28})$$

$$\text{and then } \int_{\partial\Omega} (\sigma_{ij}^{(m)} n_j - \sigma_{ij} n_j) \phi \, dP = \int_{\Omega} (\sigma_{ij,j}^{(m)} - \sigma_{ij,j}) \psi \, dx + \int_{\Omega} (\sigma_{ij}^{(m)} - \sigma_{ij}) \psi_{,j} \, dx$$

$$\text{whence } \lim_{m \rightarrow \infty} \int_{\partial\Omega} (\sigma_{ij}^{(m)} n_j - \sigma_{ij} n_j) \phi \, dP = 0 \quad \forall \phi \in H^{1/2}(\partial\Omega).$$

$$\text{and therefore } \lim_{m \rightarrow \infty} \sigma_{ij}^{(m)} n_j = \sigma_{ij} n_j \text{ almost everywhere over } \partial\Omega.$$

$$\Rightarrow \sigma_{ij} n_j = \bar{F}_i \text{ over } \partial\Omega_{F_h}.$$

• The convergence of $\sigma_{ij}^{(m)}$ toward σ_{ij} in $L^2(\Omega)$ implies the convergence almost everywhere in Ω of a subsequence $\sigma_{ij}^{(p)}$ toward σ_{ij} .
then $\lim_{p \rightarrow \infty} \mathcal{F}(\sigma^{(p)}) = \mathcal{F}(\sigma) \leq 0$ almost everywhere in Ω .

• The convexity of $\mathcal{K}$ is obvious

ii) $a(\sigma, \tau)$ is a continuous, concave ^{symmetric} bilinear form over the subspace of symmetric element in $(L^2(\Omega))^9$

That is obvious. from (23) (24) (25)

iii) $\int_{\partial\Omega_u} \sigma_{ij} n_j u_i \, dP$ is a continuous linear form over $\mathcal{K}$

The linearity is obvious. Let u be a fixed element of $(H^1(\Omega))^3$ satisfying (18)

hence:
$$\int_{\partial\Omega_0} \sigma_{ij} n_j u_i dP = (\sigma, \varepsilon(u)) + \text{constants} \quad \forall \sigma \in \mathcal{K}$$

hence the continuity expected.

All the hypothesis of the theorem 15 (mentioned above) are satisfied and we can state:

Theorem 1: Let us suppose that $\mathcal{K} \neq \emptyset$, then.

$\exists \sigma_0 \in \mathcal{K}$ unique such that

(B4) $J(\sigma_0) \leq J(\sigma) \quad \forall \sigma \in \mathcal{K}$.

and σ^0 is characterized by the variational inequality.

(B5) $\sigma_0 \in \mathcal{K} \quad a(\sigma^0, \sigma - \sigma^0) \geq \int_{\partial\Omega_0} (\sigma_{ij} - \sigma_{ij}^0) n_j u_i dP, \quad \forall \sigma \in \mathcal{K}.$

II.5. Recapitulation about the problem (A)

What have we got now?

- i) If σ_0 is solution of (A), then σ_0 is the unique solution of (A₁); whence the uniqueness of the stress field solution of (A)
- ii) We shall be able to obtain the existence of a solution for (A), if we can show that $\mathcal{K}$ is not empty and also that $\exists u \in H^1(\Omega)$ satisfying (18) and associated to σ^0 by (20) and (21)

Remark 1: We obtained also a variational principle for the displacement field but, this one is much more complicated and we have just obtained from this principle a very weak solution. It is therefore now to particularize the problem (A) in order to go ahead.

II.6 Bibliography.

About Sobolev spaces.

J.L. Lions and E. Magenes

"Non homogeneous boundary value problems"

Springer Verlag 1972

III Elastic-plastic torsion of cylindrical bars.

III 1. Application of the theorem 1 for the torsion of an homogeneous bar.

Problem B

Ω is a cylinder with a cross section Σ and a height $h=1$. (cf p. 2)

$$6) \mathcal{K} = \left\{ \sigma / \sigma \in [L^2(\Omega)]^9; \sigma_{ij} = \sigma_{ji}; \sigma_{i2,1} = 0; \mathcal{F}(\sigma) \leq 0 \text{ a.e. in } \Omega \text{ and } \right. \\ \left. \sigma_{ij} n_j = 0 \text{ a.e. over } \partial\Omega_1; \sigma_{33} = 0 \text{ a.e. over } \partial\Omega_0 \text{ and } \partial\Omega_1 \right\}$$

$$7) \mathcal{F}(\sigma) = \frac{1}{2} \sigma_{ij} \sigma_{ij} - \frac{1}{6} (\sigma_{ii})^2 - g^2 = \frac{1}{2} \sigma_{ij}^0 \sigma_{ij}^0 - g^2 \text{ with } \sigma^0 = \sigma - \frac{h a c c \sigma}{3} \mathbf{1}$$

$$8) J(\sigma) = \frac{1}{2} \int_{\Omega} \left\{ \frac{1+\nu}{E} \sigma_{ij} \sigma_{ij} - \frac{\nu}{E} \sigma_{ii}^2 \right\} d\alpha - \alpha \int_{\partial\Omega_1} (x_1 \sigma_{13} - x_2 \sigma_{23}) d\alpha_1 d\alpha_2.$$

$$9) A_{ijkl} x_{ij} x_{kl} > \frac{1-2\nu}{E} |\lambda|^2; \quad 0 \in \mathcal{K} \Rightarrow \mathcal{K} \neq \emptyset. \text{ Then } \exists \sigma_0 \in \mathcal{K} \text{ unique s.t.}$$

$$J(\sigma_0) \leq J(\sigma) \quad \forall \sigma \in \mathcal{K}.$$

σ^0 is the only possible stress tensor field solution for the problem B

Proposition 1: σ^0 is such that:

$$1/ \sigma_{ij}^0 = 0 \text{ a.e. in } \Omega \text{ except } \sigma_{13}^0 = \sigma_{31}^0 \text{ and } \sigma_{23}^0 = \sigma_{32}^0.$$

$$2/ \sigma_{i2,3}^0 = 0 \text{ a.e. in } \Omega.$$

Idea of the proof: Suppose the proposition is not true - then let us introduce $\bar{\sigma}$:

$$\bar{\sigma}_{ij} = 0 \text{ except } \bar{\sigma}_{p3} = \bar{\sigma}_{3p} \quad (p=1,2) \text{ defined as distribution on } \Sigma \text{ by.}$$

$$9) \langle \varphi, \bar{\sigma}_{p3} \rangle = \int_{\Omega} \varphi(x_1, x_2) \sigma_{p3}^0(x_1, x_2, x_3) d\alpha \quad \forall \varphi \in \mathcal{D}(\Sigma)$$

$\bar{\sigma} \in \mathcal{K}$; in fact:

$$1/ |\langle \varphi, \bar{\sigma}_{p3} \rangle| \leq \| \sigma_{p3}^0 \|_{L^2(\Omega)} \| \varphi \|_{L^2(\Sigma)} = \| \sigma_{p3}^0 \|_{L^2(\Omega)} \| \varphi \|_{L^2(\Sigma)}$$

$$\Rightarrow \bar{\sigma}_{p3} \in \text{dual } L^1(\Sigma) \equiv L^2(\Sigma)$$

ii) $\sigma_{h,j} = 0$ in Ω because:

for $h = j = 1$ or 2 .

$$\langle \varphi, \bar{\sigma}_{p,j} \rangle = - \langle \varphi, \bar{\sigma}_{p,3} \rangle = 0 \quad \forall \varphi \in \mathcal{D}(\mathbb{R})$$

for $h = 3$

$$\begin{aligned} \langle \varphi, \bar{\sigma}_{3,j} \rangle &= - \int_{\Omega} \varphi_{,p} \sigma_{3p}^{\circ} d\alpha = - \int_{\partial\Omega_1} \varphi \sigma_{33}^{\circ} d\alpha, d\alpha_1 + \int_{\partial\Omega_2} \varphi \sigma_{33}^{\circ} d\alpha, d\alpha_2 \\ &\quad - \int_{\partial\Omega_L} \varphi (\sigma_{3p}^{\circ} n_p) d\alpha = 0 \quad \forall \varphi \in \mathcal{D}(\mathbb{R}) \end{aligned}$$

$$\Rightarrow \sigma_{h,j} = 0 \quad \text{for } h = 1, 2, 3 \text{ a.e. on } \Omega.$$

iii) $\bar{\sigma}_{3p} n_p = 0$ on $\partial\Omega_L$ because.

$$\bar{\sigma}_{h,j} \in L^2(\Omega) \text{ and } \sigma_{h,j} = 0 \in L^2(\Omega) \Rightarrow \sigma_{h,j} n_j \in H^{-1/2}(\partial\Omega)$$

$$\text{but also: (40) } \bar{\sigma}_{p3} = \int_0^1 \sigma_{p3}^{\circ} d\alpha_3 \quad \text{by definition.}$$

$$\Rightarrow \bar{\sigma}_{31} n_1 + \bar{\sigma}_{32} n_2 = \int_0^1 (\sigma_{13}^{\circ} n_1 + \sigma_{23}^{\circ} n_1) d\alpha_3 = 0.$$

$$\text{IV. } \mathcal{J}(\bar{\sigma}) = \bar{\sigma}_{13}^2 + \bar{\sigma}_{23}^2 - g^2 \leq 0 \quad \text{because.}$$

$$\begin{aligned} \bar{\sigma}_{13}^2 + \bar{\sigma}_{23}^2 &= \int_0^1 (\sigma_{13}^{\circ} \bar{\sigma}_{13} + \sigma_{23}^{\circ} \bar{\sigma}_{23}) d\alpha_3 \\ &\leq (\bar{\sigma}_{13}^2 + \bar{\sigma}_{23}^2)^{1/2} \left[\int_0^1 (\sigma_{13}^{\circ 2} + \sigma_{23}^{\circ 2}) d\alpha_3 \right]^{1/2}. \end{aligned}$$

$$\Rightarrow \sigma_{13}^2 + \sigma_{23}^2 \leq g^2.$$

• $\mathcal{J}(\bar{\sigma}) \leq \mathcal{J}(\sigma^{\circ})$; in fact:

$$\mathcal{J}(\sigma) = \frac{1}{2} a(\sigma, \sigma) - \alpha \int_{\partial\Omega_1} (\alpha_1 \sigma_{23} - \alpha_2 \sigma_{13}) d\alpha, d\alpha_2$$

we are going to prove that $\int_{\partial\Omega_1} (\alpha_1 \bar{\sigma}_{23} - \alpha_2 \bar{\sigma}_{13}) d\alpha, d\alpha_2 = \int_{\partial\Omega_1} (\alpha_1 \sigma_{23}^{\circ} - \alpha_2 \sigma_{13}^{\circ}) d\alpha, d\alpha_2$

and also that $a(\bar{\sigma}, \bar{\sigma}) \leq a(\sigma^{\circ}, \sigma^{\circ})$.

$$(41) \quad \int_{\partial\Omega_1} (\alpha_1 \bar{\sigma}_{23} - \alpha_2 \bar{\sigma}_{13}) d\alpha, d\alpha_1 = \int_{\Omega} (\alpha_1 \sigma_{23}^0 - \alpha_2 \sigma_{13}^0) d\alpha.$$

now $\forall \sigma \in \mathcal{K}$ and $u \in (H^1(\Omega))^3$ satisfying $u_1 = u_2 = 0$ over $\partial\Omega_0$ and $u_1 = -\alpha \alpha_2, u_2 = \alpha \alpha_1$ over $\partial\Omega_1$, we have.

$$\int_{\Omega} \sigma_{ij} \varepsilon_{ij}(u) d\alpha = \alpha \int_{\partial\Omega_1} (\sigma_{23} \alpha_1 - \sigma_{13} \alpha_2) d\alpha, d\alpha_1.$$

in particular for $u^* = (-\alpha \alpha_2 \alpha_3, \alpha \alpha_1 \alpha_3, 0)$

$$\int_{\Omega} \sigma_{ij} \varepsilon_{ij}(u^*) d\alpha = \alpha \int_{\Omega} (\sigma_{23} \alpha_1 - \sigma_{13} \alpha_2) d\alpha = \alpha \int_{\partial\Omega_1} (\sigma_{23} \alpha_1 - \sigma_{13} \alpha_2) d\alpha, d\alpha_1.$$

then for $\sigma_{ij} = \sigma_{ij}^0$, $\int_{\Omega} (\sigma_{23}^0 \alpha_1 - \sigma_{13}^0 \alpha_2) d\alpha = \int_{\partial\Omega_1} (\sigma_{23}^0 \alpha_1 - \sigma_{13}^0 \alpha_2) d\alpha, d\alpha_1.$

Whence (41)

In an other end, from (40) and the inequality $\nu < \frac{1}{2}$, we obtain.

$$a(\sigma^0, \sigma^0) - a(\bar{\sigma}, \bar{\sigma}) \geq \frac{1+\nu}{2E} \int_{\Omega} |\sigma_{ij}^{0,D} \sigma_{ij}^{0,D} - \bar{\sigma}_{ij}^D \bar{\sigma}_{ij}^D| \geq 0.$$

$\Rightarrow J(\bar{\sigma}) \leq J(\sigma^0) \Rightarrow \bar{\sigma} = \sigma^0$ because the uniqueness of σ^0 and then the Prop. 1

Remark 2: This proposition is still true in the two following cases:

a) Ω is a cylinder with some cylindrical cavities $\Omega_1, \dots, \Omega_q$ with same direction of generatrix: we have then to apply the preceding proof to the multiply connected domain.

$$\Omega^* = \left[\overbrace{\Omega}^{\cup \Omega_i} \right]_{i=1, \dots, q}$$

b) Ω is a cylinder reinforced by cylindrical fibers $\Omega_1, \dots, \Omega_q$ with same direction of generatrix. In this case, ν and E are discontinuous.

$$\nu \text{ (and } E) = \begin{cases} \nu_{\text{fib}}(E_f) \text{ on the fibers } \Omega_1, \dots, \Omega_q. \\ \nu_m \text{ or } (E_m) \text{ on the matrix } \Omega^* \end{cases}$$

Remark 3: The problem can be reduce now to a bidimensional one

III.2 new formulation of the problem (A₁) in the particular case of Torsion

$$\mathcal{K}_T = \left\{ \sigma / \sigma \in \mathcal{K}, \sigma_{ij} = 0 \text{ except } \sigma_{p3} = \sigma_{3p}; p=1,2 \right\} \quad (\text{T or Torsion})$$

We know now that $\sigma^0 \in \mathcal{K}_T$ and that $J(\sigma^0) \leq J(\sigma) \quad \forall \sigma \in \mathcal{K}_T \subset \mathcal{K}$

Now:

Lemma 1: $\forall \sigma \in \mathcal{K}_T, \exists \theta \in H_0^1(\Sigma)$ unique such that.

i) $\theta_{,2} = \sigma_{13}; \theta_{,1} = -\sigma_{23} \quad \text{a.e. on } \Sigma \quad (\text{non section of } \Omega)$

ii) Furthermore $|\text{grad } \theta| \leq g$.

Proof: • $\sigma_{ij} \in L^2(\Omega); \sigma_{12,2} = 0$ and $\sigma_{33} = 0$ on $\partial\Omega_0$ and $\partial\Omega_1 \Rightarrow$

$\exists \varphi \in H^1(\Sigma)$ such that $\varphi_{,2} = \sigma_{13}$ and $\varphi_{,1} = -\sigma_{23}$.

• $\sigma_{ij} n_j = 0$ on $\partial\Omega_L \Leftrightarrow \overline{\text{grad } \varphi} \wedge \vec{n} = 0$ on $\partial\Sigma$.

$$\Leftrightarrow \varphi = h = \text{constants on } \partial\Sigma.$$

• Let $\theta = \varphi - h$, then $\theta \in H_0^1(\Sigma)$

• $|\text{grad } \theta|^2 = \sigma_{13}^2 + \sigma_{23}^2$ and $\|\theta\|_{H_0^1(\Sigma)}^2 = \int_{\Sigma} |\text{grad } \theta|^2 d\alpha$.

imply the uniqueness of such a θ in $H_0^1(\Sigma)$

• $J(\sigma) \leq 0 \Rightarrow |\text{grad } \theta| \leq g$.

It is easy to see now that $\exists$ a bijection between $\mathcal{K}_T$ and $\mathcal{K}$.

$$(42) \quad \mathcal{K} = \left\{ \varphi / \varphi \in H_0^1(\Sigma) \text{ s.t. } |\text{grad } \varphi| \leq g \right\};$$

for two corresponding elements σ and θ .

$$(42) \quad 2\mu J(\sigma) = J_\alpha(\theta) = \int_{\Sigma} |\text{grad } \theta|^2 d\alpha - 4\mu d \int_{\Sigma} \theta d\alpha \quad \text{with } \mu = \frac{E}{2(1+\nu)}$$

Therefore, in the particular case of the torsion of a homogeneous, elastic-plastic cylindrical bar, the problem (A₁) is replaced by the equivalent problem (B₁) Find θ_α such that

$$(43) \quad \theta_\alpha \in \mathcal{K}; \quad J_\alpha(\theta_\alpha) \leq J_\alpha(\varphi) \quad \forall \varphi \in \mathcal{K}$$

Lemma 2 : $\exists f_\alpha \in H_0^1(\Sigma)$, unique such that

$$(45) \quad \int_{\Sigma} \varphi \, d\alpha = \langle \varphi, f_\alpha \rangle \quad \forall \varphi \in H_0^1(\Sigma)$$

with (45) $\langle \varphi, \psi \rangle = \int_{\Sigma} \text{grad} \varphi \cdot \text{grad} \psi \, d\alpha$, scalar product on $H_0^1(\Sigma)$

This lemma is nothing else than a direct application of the Riesz theorem (cf Lecture notes H.L. p 7)

Then we got a new expression of the functional.

$$(47) \quad J_\alpha(\varphi) = \|\varphi - f_\alpha\|^2 - \|f_\alpha\|^2.$$

and because f_α is a fixed element of $H_0^1(\Sigma)$, we obtain a new form for the problem B_α , because θ_α is the element which minimizes the distance $\|\varphi - f_\alpha\|$ of f_α to the closed convex K of $H_0^1(\Sigma)$. We obtain then the following conclusion (cf example 6 p 39 in lecture notes of H.L)

The solution θ_α of the problem (B_α) is exactly the projection on the closed convex K of the element f_α defined by (45). θ_α is characterized by the variational inequality.

$$(48) \quad \theta_\alpha \in K; \quad \langle \theta_\alpha - \varphi, \theta_\alpha - f_\alpha \rangle \leq 0 \quad \forall \varphi \in K$$

III.3 First Properties of the stress function θ_α

To obtain this projection, it is natural to study first the general properties of K .

III.3.1. Properties of K (cf references given p 7)

• K is the set of Lipschitz functions of $H_0^1(\Sigma)$ with modulus g , that is to say.

$$(49) \quad |\varphi(x) - \varphi(x')| \leq g|x - x'| \quad \forall \varphi \in K \text{ and } \forall x \text{ and } x' \in \Sigma$$

Therefore these functions are uniformly continuous on $\bar{\Sigma}$

- K is compact for the norm of uniform convergence.

Proof: That come from Ascoli theorem which says that: if X and Y are two compact ^{metric} spaces, if E is a set of equicontinuous mapping of X into Y , then the closure of E in $\mathcal{C}(X, Y)$, set of continuous mapping from X into Y , is compact.

here: $X = \bar{\Sigma}$ (compact of $\mathbb{R}^2$); $Y = \bigcup_{\varphi \in K} \varphi(\bar{\Sigma})$ (compact of $\mathbb{R}$)

K is a set of equicontinuous mapping from X into Y .

$\Rightarrow$ if $\mathcal{C}(X, Y)$ is supplied with the norm of the uniform convergence then $\bar{K}$ is uniformly compact and it is easy to prove that with this topology $K = \bar{K}$.

- if φ_1 and $\varphi_2 \in K$ then $\sup(\varphi_1, \varphi_2)$ and $\inf(\varphi_1, \varphi_2) \in K$

III 3.2 Regularity of θ_α

From H. Brezis and G. Stampacchia (cf. reference) we have:

$\theta_\alpha \in C^{1,1}(\bar{\Sigma})$ / set of continuous functions with Lipschitz derivatives in the two following cases

Σ , convex set of $\mathbb{R}^2$ with Lipschitz boundary.

Σ , non convex set but with more regular boundary.

III 4 Behaviour of θ_α for α "small": elastic solution

Proposition 2: (50) Let $\gamma = \frac{f_\alpha}{2\mu\alpha}$ then:

a) γ is the solution of the Dirichlet problem.

$$(51) \quad \Delta \gamma + 2 = 0 \text{ on } \Sigma \text{ and } \gamma = 0 \text{ on } \partial \Sigma$$

b) $\gamma \in W^{2,p}(\Sigma) \cap C^\infty(\Sigma) \cup C^{1,1}(\bar{\Sigma})$ $\forall 2 \leq p < \infty$.

c) $\gamma > 0$ in Σ

d) $|\text{grad } \gamma|$ is bounded on $\bar{\Sigma}$ and $\sup |\text{grad } \gamma|$ must be reached only on the boundary $\partial \Sigma$

Remark 4: the practical consequence of this result is the following:

(52) Let $M = \sup_{\bar{Z}} |\text{grad } \gamma|$ from b) we know that M is bounded

from d) we know that $|\text{grad } \gamma| < M$ on Σ

Let us introduce then (53) $\alpha_0 = \frac{g}{\mu M}$ then for $0 \leq \alpha \leq \alpha_0$, $|\text{grad } f_\alpha| \leq g$.

$\Rightarrow f_\alpha \in K$ and then $\partial \alpha = f_\alpha$

then from a)
$$\begin{cases} \Delta \partial \alpha + 2\mu \alpha = 0 & \text{in } Z. \\ \partial \alpha = 0 & \text{on the boundary.} \end{cases}$$

$|\text{grad } \partial \alpha| < g$ on Σ .

We can recognize this the classical solution of elastic torsion for a homogeneous cylindrical bar.

Moreover, from d) we know that the plasticity will appear first on the boundary.

Proof of the Proposition 2:

a) $\langle \gamma, \gamma \rangle = 2 \int_Z \gamma \, d\alpha \quad \forall \gamma \in H_0^1(Z)$ and in particular.

$\forall \gamma \in \mathcal{D}(Z)$ then

$$\int_Z (\text{grad } \gamma \cdot \text{grad } \gamma - 2\gamma) \, d\alpha = 0 \quad \forall \gamma \in \mathcal{D}(Z)$$

$$\int_Z -(\Delta \gamma + 2) \gamma \, d\alpha = 0 \quad \forall \gamma \in \mathcal{D}(Z).$$

$\Rightarrow \begin{cases} \Delta \gamma + 2 = 0 & \text{in } Z. \\ \gamma = 0 & \text{on } \partial Z. \end{cases} \quad \text{homogeneous Dirichlet problem. (Lecture notes H.2)}$

$\Delta \gamma \in L^p \quad \forall p \geq 2$ and $\gamma \in H_0^1(Z) \Rightarrow \gamma \in W^{2,p}(Z) \cap H_0^1(Z)$

$\Rightarrow \gamma \in C^{1,1}(\bar{Z})$

(general results of regularity mentioned orally in the first lecture)

$\gamma \in C^0(Z)$ is a consequence of the constant coefficients of (51)

b) c and d are direct consequences of the maximum principle:

In fact: $\Delta \gamma = -2 < 0$ shows that γ is superharmonic and therefore,

γ can reach its minimum only on ∂Z ; $\gamma = 0$ on ∂Z implies then $\gamma > 0$ on Z .

Otherwise: $\gamma \in C^{1,1}(\bar{Z})$ implies that $|\text{grad } \gamma|$ is bounded on $\bar{Z}$

$$\Delta |\text{grad } \gamma|^2 = 2 \text{grad } \Delta \gamma \cdot \text{grad } \gamma + 2 \gamma_{,i} \gamma_{,i} = 2 \gamma_{,i} \gamma_{,i} \geq 0 \text{ on } Z$$

