



INTERNATIONAL ATOMIC ENERGY AGENCY
UNITED NATIONS EDUCATIONAL, SCIENTIFIC AND CULTURAL ORGANIZATION



INTERNATIONAL CENTRE FOR THEORETICAL PHYSICS
34100 TRIESTE (ITALY) - P.O.B. 586 - MIRAMARE - STRADA COSTIERA 11 - TELEPHONE: 3340-1
CABLE: CENTRATOM - TELEX 460392 - I

HA.SMR/285 - 36

WINTER COLLEGE ON
LASER PHYSICS: SEMICONDUCTOR LASERS
AND INTEGRATED OPTICS

(22 February - 11 March 1986)

INTEGRATED OPTICAL PASSIVE DEVICES

W. Sohler
Angewandte Physik
Universität-GH-Paderborn
F.R. Germany

INTEGRATED OPTICAL PASSIVE DEVICES

WAVEGUIDES

COUPLERS

LENSES

PRISMS

POLARIZERS

TE/TM - MODE SPLITTER

BEAM SPLITTER

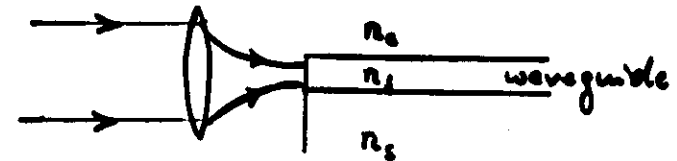
GRATING DEVICES (FILTER ETC.)

MIRRORS

.....

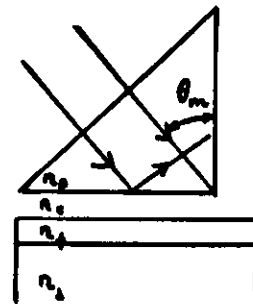
Beam couplers

i) transversal coupler (end-fire coupling)

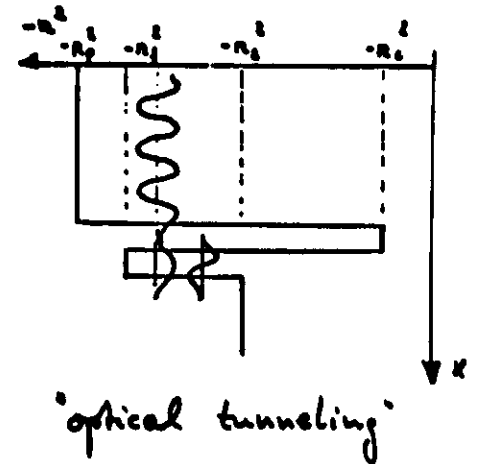


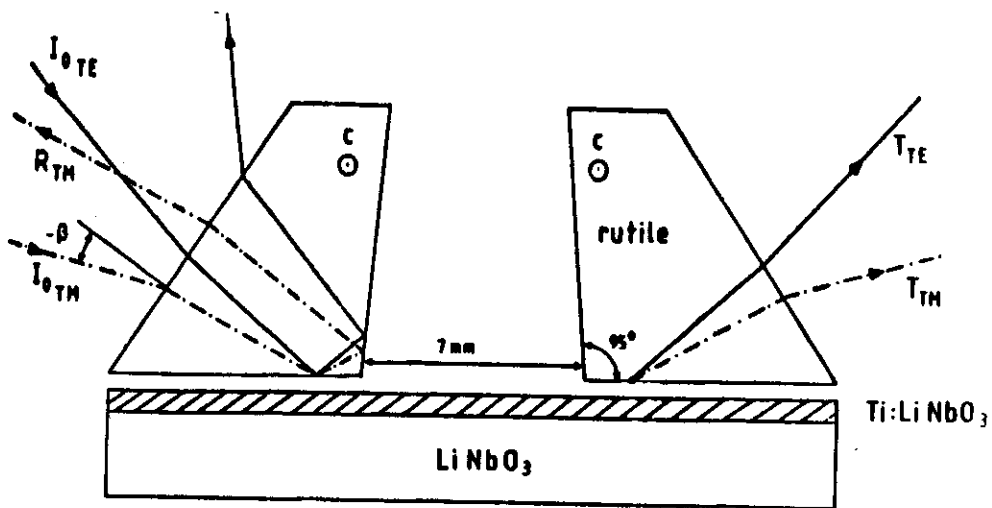
$$\eta = \frac{\left[\int E_i(x,y) E_{wg}^*(x,y) dx dy \right]^2}{\int |E_i(x,y)|^2 dx dy \int |E_{wg}(x,y)|^2 dx dy}$$

ii) prism coupler



$$\beta_m = \frac{\omega}{c} n_p \sin \theta_m$$





iii) grating coupler
(after Tein)

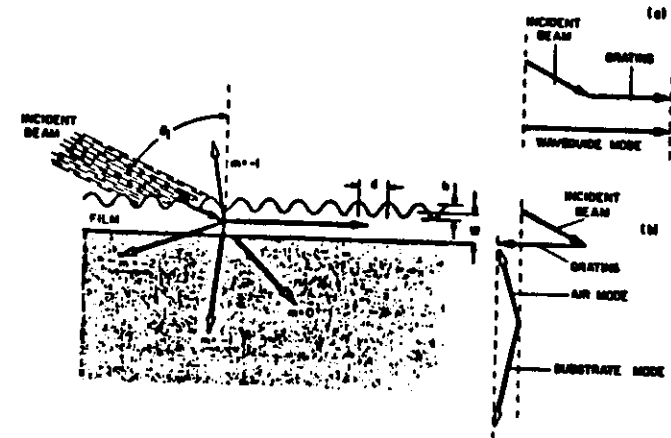
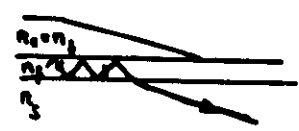


FIG. 13. The figure at the left shows how a grating can diffract an incident wave into many orders. A zigzag wave and thus a "waveguide" mode is excited in a waveguide only if one of the diffraction orders is in phase with the zigzag wave. The diagram in (a) shows the vector relation required for such a grating coupler. On the other hand, other diffraction orders of the grating could match the phases of the "substrate" or "airspace" modes and thus introduce losses to the coupling by diverting part of the incident energy into radiations in substrate or air-space.

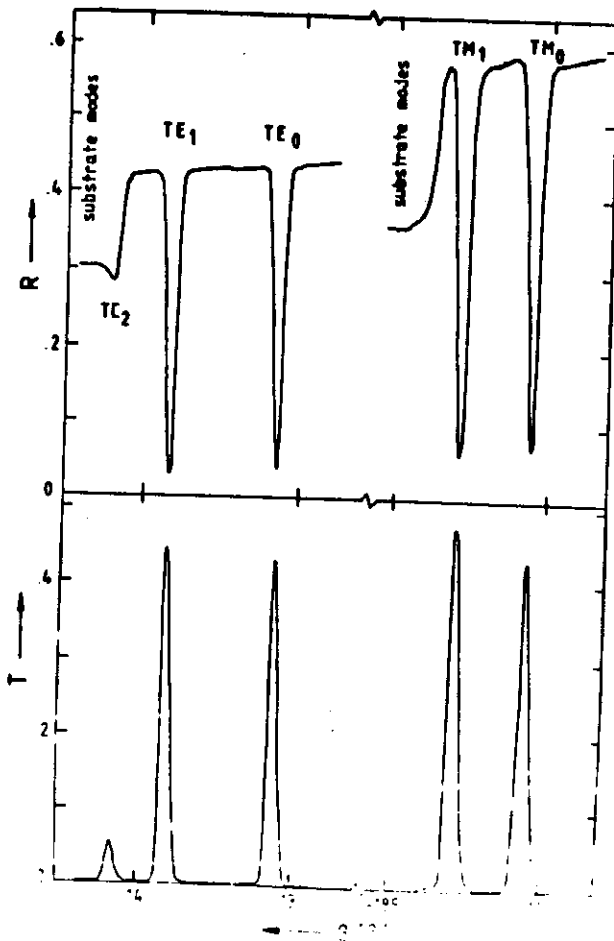
iv) taper coupler



(Tein)

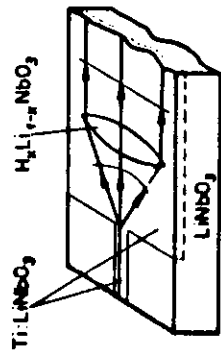


(Soller)

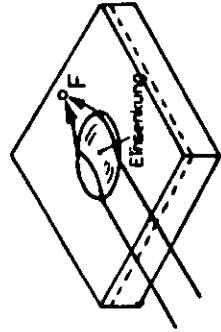


after: Soller + Sude

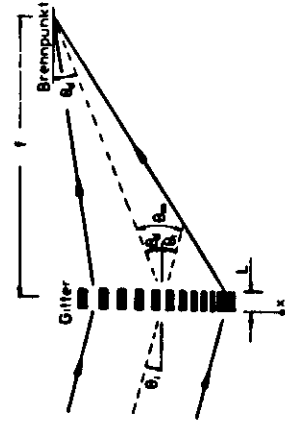
Integriert optische Linsen



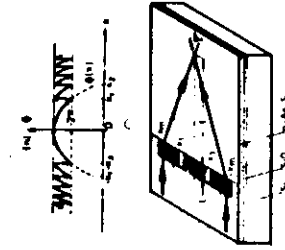
Bikonvexe Linse



Geodätische Linse

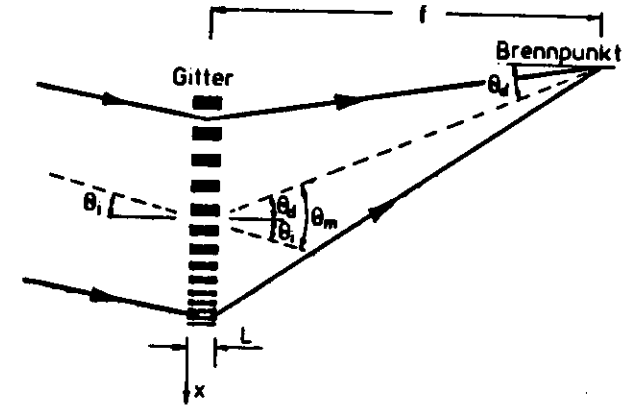


Gitterlinse

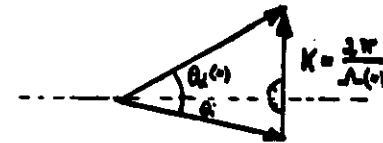


Fresnel-Linse

grating lens:



wave vectors:



$$\frac{2\pi}{\lambda} n_{eff} (\sin \theta_i + \sin \theta_d(\omega)) = \frac{2\pi}{\Lambda(z)}$$

$$\Rightarrow \Lambda(z) = \frac{\lambda}{n_{eff} (\sin \theta_i + \sin \theta_d(\omega))}$$

with $x = f \tan \theta_d$

experimental example:
(after Yao et al.)

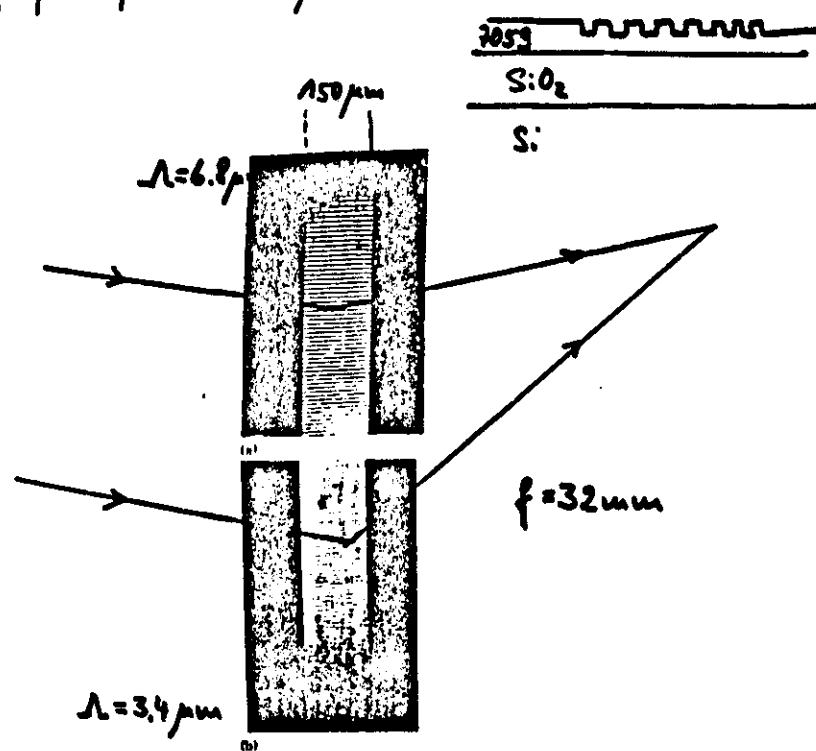


FIG. 2. The chromium master mask for the chirp-grating lens. (a) Low-frequency portion; (b) High-frequency portion.

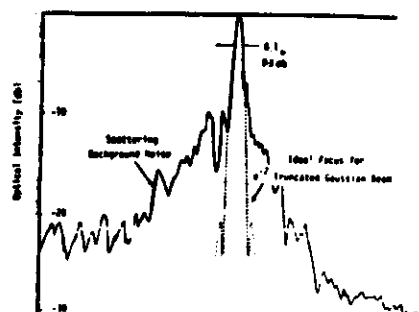


FIG. 4. Measured focal spot of the chirp-grating lens.

Fresnel Lens

convex. lens: $\phi(x) \approx -\frac{2\pi}{\lambda} n \frac{1}{2f} x^2$

Fresnel lens: $\phi_p(x) \approx -\frac{2\pi}{\lambda} n_{eff} \frac{1}{2f} x^2 + m \cdot 2\pi$

with $-2\pi < \phi_p(x) < 0$

$m = 1, 2, \dots$

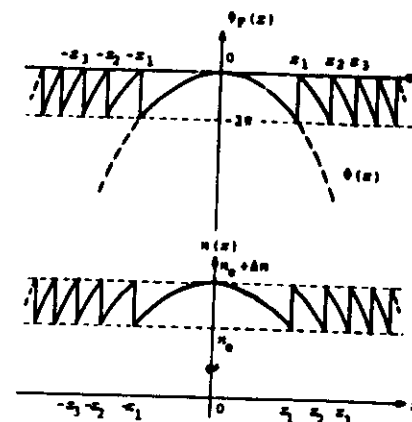


Fig. 2. Phase modulation $\phi(x)$, $\phi_p(x)$, and mode-index distribution $n(x)$ of a GRIN Fresnel waveguide lens.

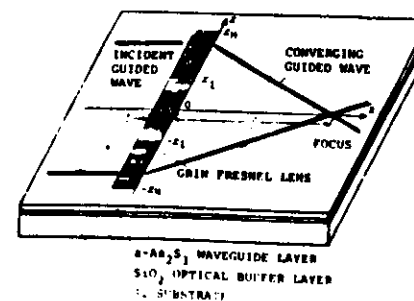


Fig. 1. Graded index Fresnel waveguide lens configuration.

experimental example: after Valette et al.

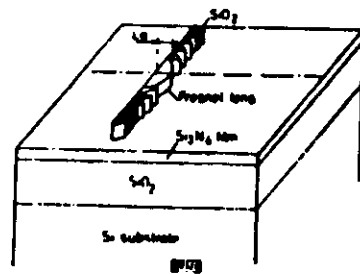
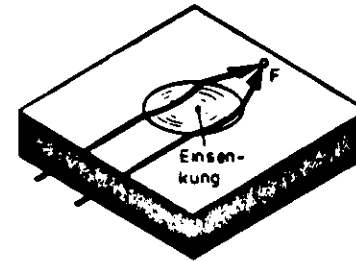


Fig. 1 Schematic drawing of integrated analog Forward lens

geodesic lenses:



after Schulz + Voges

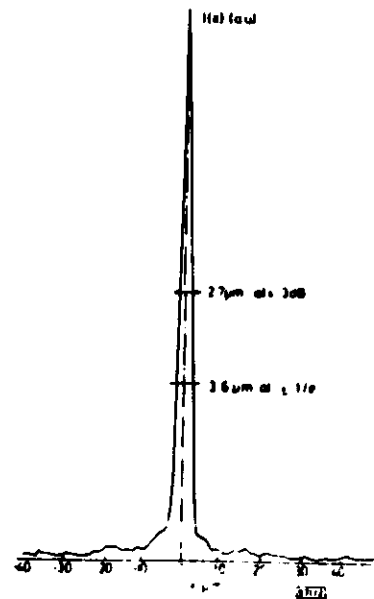
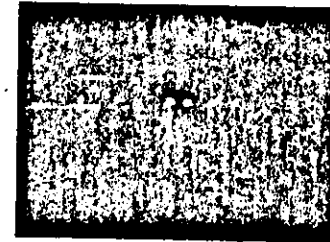


Fig. 2 Typical light distribution in dB as a function of angle θ of a lens on input beam

after Righini et al.



Fig. 3 Micrograph of a lens element on input beam

Luneburg-lens:

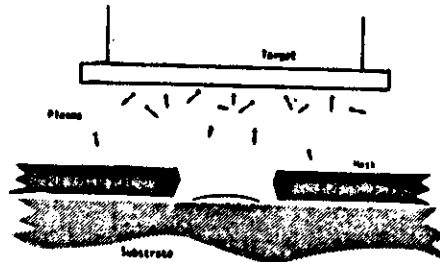


FIG. 1. Shadow mask sputtering configuration.

experimental example: after Yao et al.
 Ta_2O_5 lens on 7059 waveguide on SiO_2 substrate

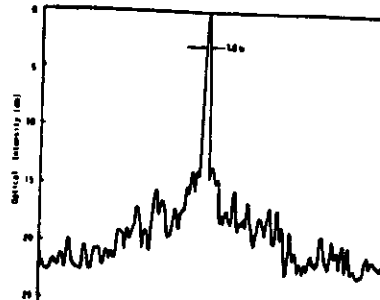
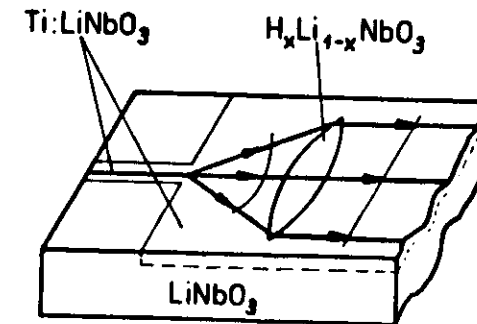


FIG. 3. Measured focal intensity pattern for an F/4 waveguide Luneburg lens with a 5-mm lens diameter and a 5-mm focal length, using a 2-mm aperture.

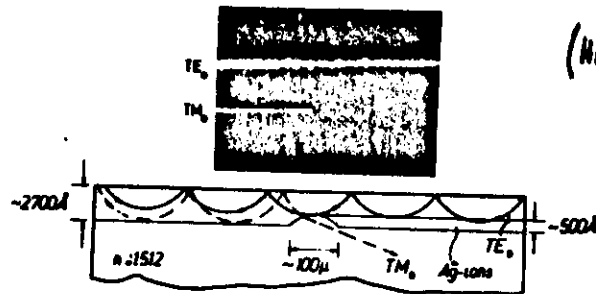
"Conventional" lens:

in LiNbO_3 by proton exchange
 (see Tsai et al.)



prism: see slide

polarizer:

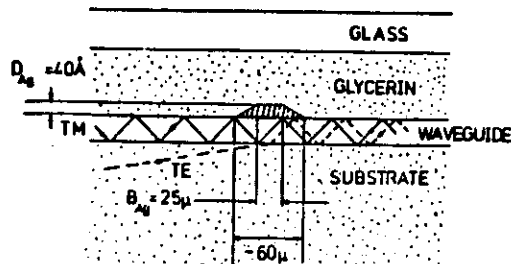


(Hinkov + Sohle)

reduced depth polarizer

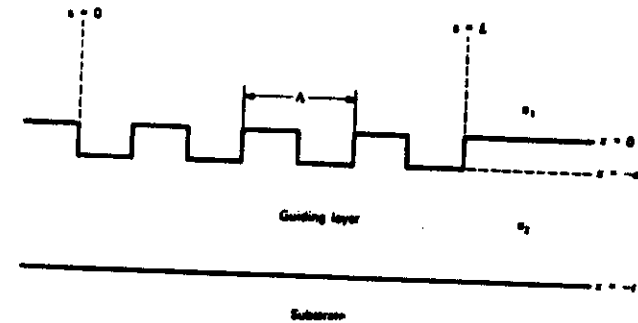


(Rollke + Sohle)



metal stripe polarizer

grating devices:



wave equation:

$$\nabla^2 E - \frac{n^2(x,y)}{c^2} \frac{\partial^2 E}{\partial t^2} = \underbrace{\mu_0 \frac{\partial^2 \rho}{\partial t^2}}_{\frac{\Delta n^2}{c^2} \frac{\partial^2 E}{\partial t^2}}$$

ansatz: $E = \sum A_m(z) E_m(x,y) \exp[i(\omega t - \beta_m z)]$

coupling of two modes:

$$\begin{aligned} A_1'(z) &= \kappa_{12} A_2(z) e^{-i\Delta\beta z} \\ A_2'(z) &= \kappa_{21}^* A_1(z) e^{+i\Delta\beta z} \end{aligned}$$

coupled mode theory
(see: Yariv)

with: $\Delta\beta = \beta_1 - \beta_2 - \kappa$

$$\kappa_{12} \sim \int \Delta n^2(x,z) E_1^*(x,y) E_2(x,y) dx dy$$

$$\beta \approx \frac{\pi}{\Lambda} \pm i \left[\kappa^2 - \left(\frac{n_{eff}}{c} \right)^2 (\omega - \omega_0)^2 \right]^{1/2}$$

with $\omega_0 = \frac{c}{\Lambda n_{eff}}$

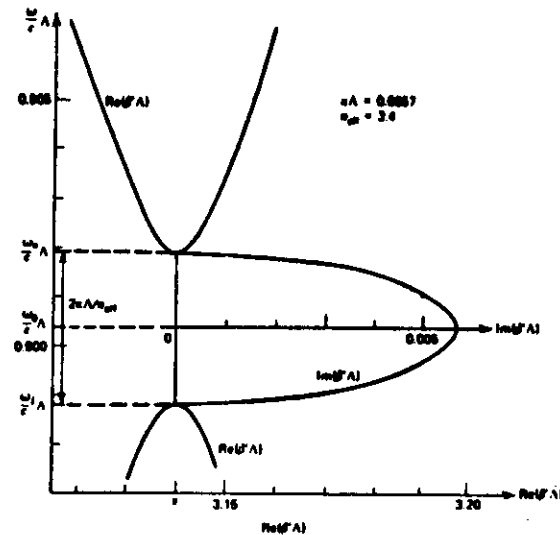


Figure 19.6 The dependence of the real and imaginary parts of the mode propagation constant, β , of the modes in a periodic waveguide near the first "forbidden" gap ($\Gamma = 1$). At frequencies $\omega_1 < \omega < \omega_2$, $\text{Im}(\beta) \neq 0$ and the modes are evanescent. At these frequencies, $\text{Re } \beta = \pi/\Lambda$. For $\omega \neq 1$ the abscissa should be taken as $\text{Re}(\beta\Lambda/\omega)$.

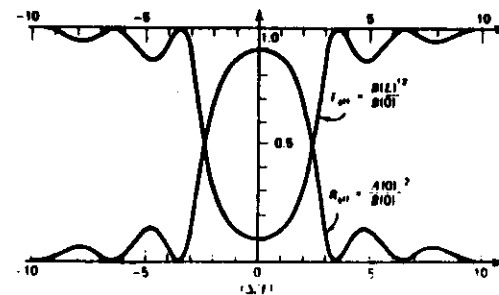


Figure 19.7 The transmission and reflection characteristics of a corrugated section of length l , as a function of the detuning $\Delta\beta L = (\beta - \beta_0)l$, $\beta_0 = 2\pi/\Lambda$, $l/L = 1.84$.

solutions for $\Delta\beta = 0$:

$$A_1(z) = A_2(0) \frac{\kappa_{12}}{|\kappa_{12}|} \frac{\sinh[\kappa(z-L)]}{\cosh(\kappa L)}$$

$$A_2(z) = A_2(0) \frac{\cosh[\kappa(z-L)]}{\cosh(\kappa L)}$$

$\kappa = |\kappa_{12}|$

reflection: $\beta_2 = -\beta_1 \Rightarrow 2\beta_1 = \kappa = \frac{2\pi}{\Lambda}$

$$\Rightarrow \Lambda = \frac{\lambda}{2n_{eff}}$$

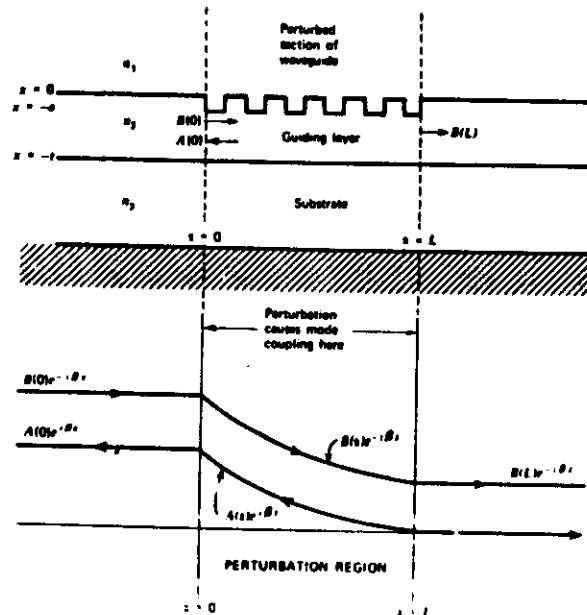
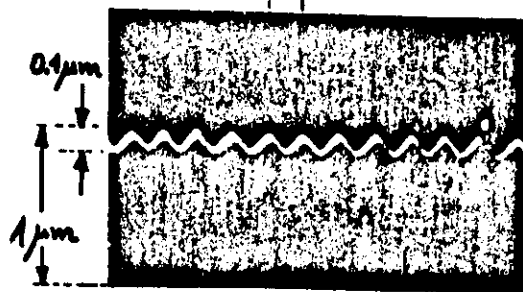


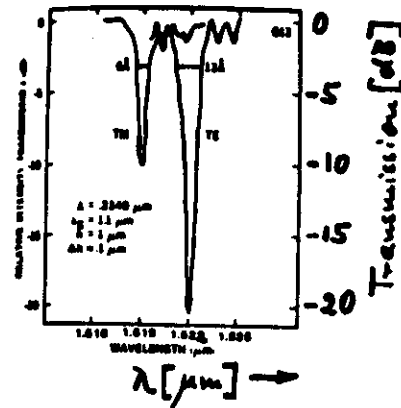
Figure 19.5 Upper: A corrugated Section of a dielectric waveguide. Lower: The incident and reflected fields.

planarer WL: LPE, InGaAsP mit $\lambda_g = 1.1 \mu\text{m}$
chem. Ätzen: 234 nm $6-7 \text{ dB/cm bei } \lambda = 1.5 \mu\text{m}$

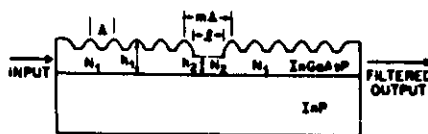


$\leftarrow 1 \mu\text{m} \rightarrow$

nach Alferness + ... (Bell), 1984

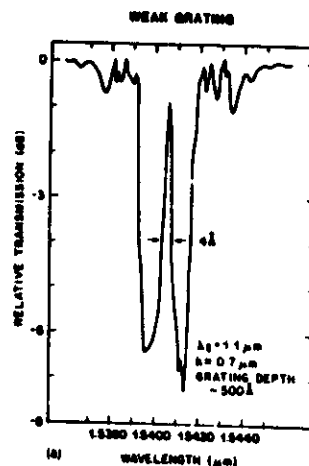


Resonator - Filter:



$(1 \text{ Å} \approx 13 \text{ GHz})$

nach Alferness + ... (Bell), 1986



multiplexer / demultiplexer:

(after: Rottmann + Kges)

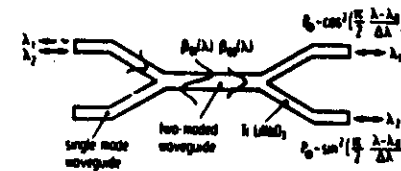


Fig. 1
Configuration of an
integrated-optical
dual-channel multi-
plexer/demultiplexer

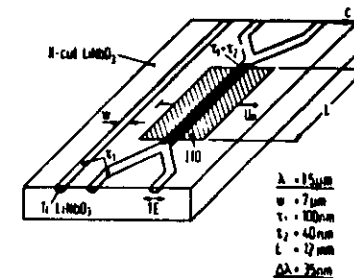


Fig. 2
Geometry of fabri-
cated MDM-devices

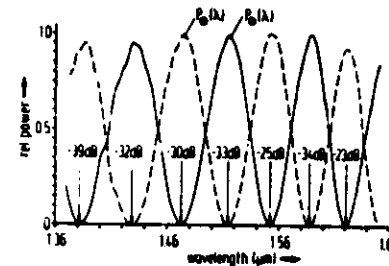


Fig. 3
Transmission charac-
teristic of a dual-
channel multiplexer/
demultiplexer
*1: strip width 7 μm,
t1 thickness
100 nm ± 40 nm,
L = 12 nm, TE pol.

mirrors:

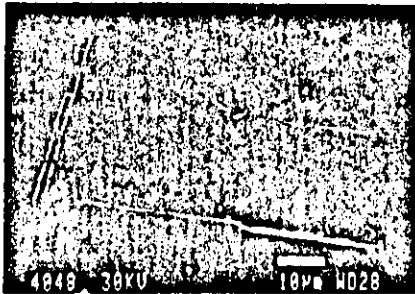


Fig.3 Reactive ion etched rib waveguides with totally reflecting corner mirror

(n/n^+ - GaAs waveguide)

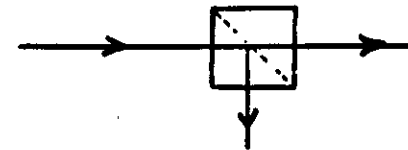
etched mirror
(Budmann et al.)

evaporated end face mirror (see slide)

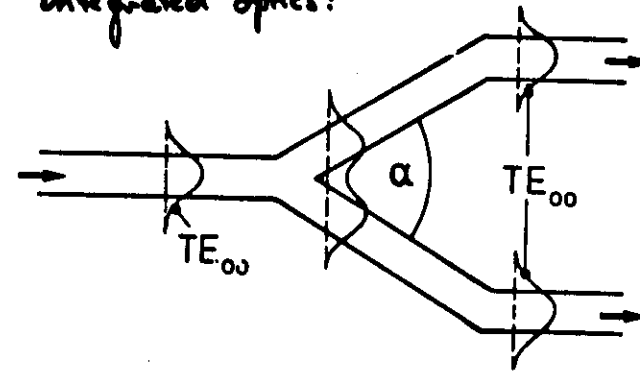
resonators: see modulators

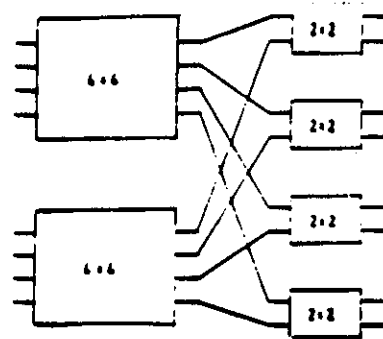
beam splitter:

bulk optics:



integrated optics:





Star coupler:

Fig. 1: Architecture of the 8x8 star coupler

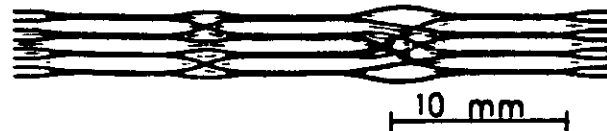


Fig. 2: 8x8 star coupler design

(after: Döldissen et al.)

