



INTERNATIONAL ATOMIC ENERGY AGENCY
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SCHOOL ON
NON-ACCELERATOR PHYSICS
25 April - 6 May 1988

GRAVITATIONAL WAVES (1 & 2)

by

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GRAVITATIONAL WAVES

- Introduction
- Generation of g.w.
- Possible astrophysical sources
- Resonant and non resonant antennas
- Energy absorbed by the resonant antennas
- Noise in a resonant antenna

- Experiments in the world with resonant and non resonant antennas
- Comparison between the two antennas
- The home experiment
- The transducer
- Measurements
- Preliminary home-Stanford coincidences

GRAVITATIONAL WAVES

D. Heaviside 1893
(Electromagnetic theory)

H.A. Lorentz 1900

H. Poincaré 1905
(Circolo Matematico di Palermo)

A. Einstein 22 June 1916

31 January 1918

G. Ricci Curbastro (1853-1925)

T. Levi-Civita (1873-1941)
University of Padova

Any metric theory of gravity, which incorporates Lorentz invariance in its field equations, forces the propagation of gravitational waves $ds^2 = g_{ik} dx^i dx^k$

Theories alternative to G.R.

Scalar-tensor theory

vector-tensor theory

Mosca's bimetric theory

Bravais' theory

BSLL theory

stratified theories

C.M. Will: "Theory and Experiment
in Gravitational Physics"

Cambridge Univ. Press

London 1981

Adimensional coupling constants

Interaction

strong

$$\frac{g^2}{\hbar c} \sim 1$$

electromagnetic

$$\frac{e^2}{\hbar c} = \alpha \sim \frac{1}{137}$$

weak

$$\frac{G \mu m_p^2}{(\hbar c)^3} \sim 10^{-5}$$

gravitational

$$\frac{G m_p^2}{\hbar c} \sim 6 \times 10^{-39}$$

Gravitational wave experiments

1960-1970 J. Weber

1970-1975 Various room temperature experiments

1971 → beginning II generation resonant g.w. antennas

1975 → beginning the development of laser interferometry

Plans, at least, until 1995

15 experimental groups

$$10^{-30} \text{ joule} \simeq 10^{-11} \text{ eV}$$

Energy of a large body

~ 2000 kg

General Relativity

Einstein eq.
in vacuum $R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = 0$

The metric tensor $ds^2 = g_{\mu\nu} dx^\mu dx^\nu$

weak field approx.

$$\left\{ \begin{array}{l} g_{\mu\nu} = g_{\mu\nu}^{(0)} + h_{\mu\nu} \\ |h_{\mu\nu}| \ll 1 \end{array} \right.$$

The wave equation $\frac{\partial^2 h_{\mu\nu}}{\partial x^\alpha \partial x_\alpha} \equiv \square h_{\mu\nu} = 0$

Plane waves propagating along x $h_{\mu\nu} = h_{\mu\nu}(t - x/c)$
light velocity \uparrow

Gravitational waves

Transverse-Traceless gauge
"TT gauge"

Only 4 components of $h_{\mu\nu}$
are different from zero

$$\left. \begin{array}{l} h_{32} = h_{23} = h_x \\ h_{22} = -h_{33} = h_\perp \end{array} \right\} \begin{array}{l} \text{for a wave} \\ \text{propagating} \\ \text{along } x \end{array}$$

- Two polarization states
- Transverse waves

Do gravitational waves carry energy?

(3)

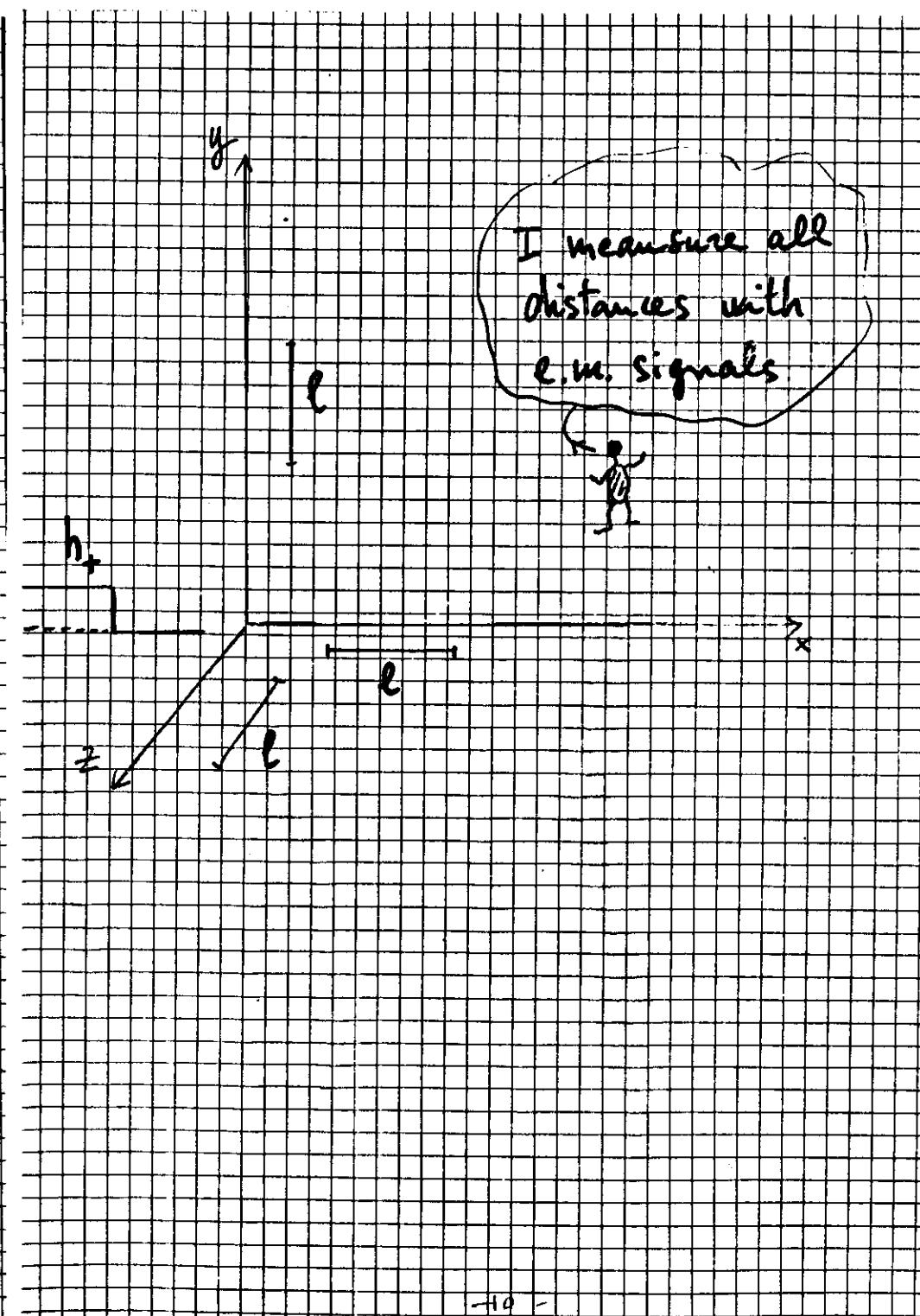
$$W = \frac{c^3}{16\pi G} h^2 \left[\frac{\text{joule}}{\text{m}^2 \text{s}} \right]$$

$$= 8 \cdot 10^{33} h^2$$

For $W = 1400 \frac{\text{joule}}{\text{m}^2 \text{s}}$ (solar e.m. on the Earth)

and $\omega = 5000 \text{ rad/s}$

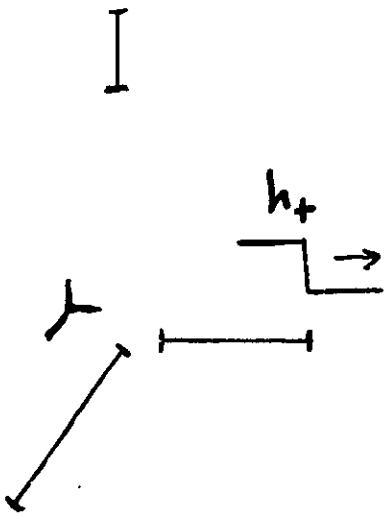
$$h \sim 10^{-19}$$



(2)

We change now point of view:

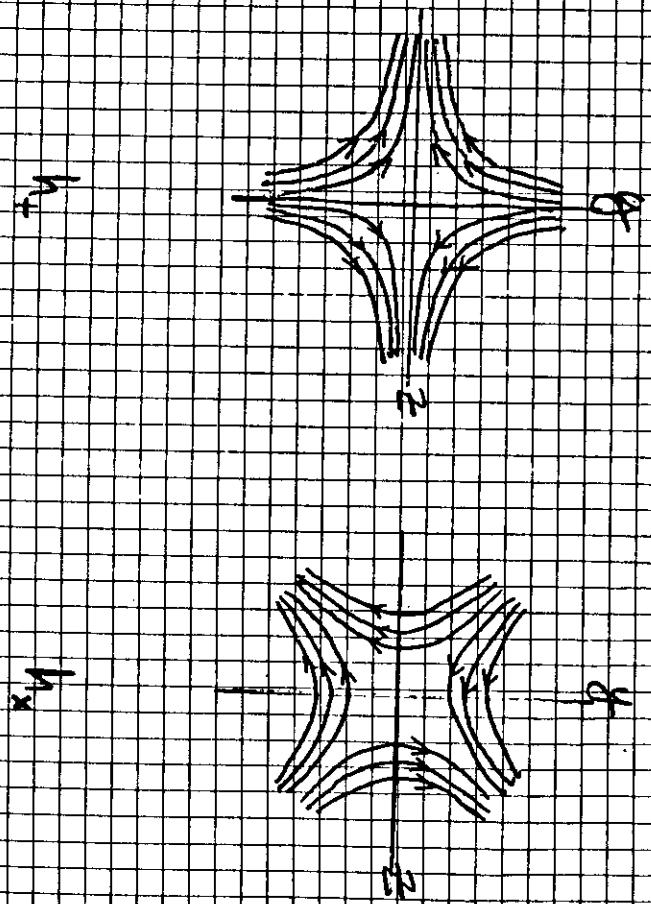
we pass from Einstein
geometro-dynamical description
to normal dynamics based on the
concept of forces:



$$\frac{\delta e}{e} = 0 \quad e \parallel x$$

$$\frac{\delta e}{e} = -\frac{1}{2} h_+ \quad e \parallel y$$

$$\frac{\delta e}{e} = +\frac{1}{2} h_+ \quad e \parallel z$$



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Figure 2 Scheme for a laser interferometer

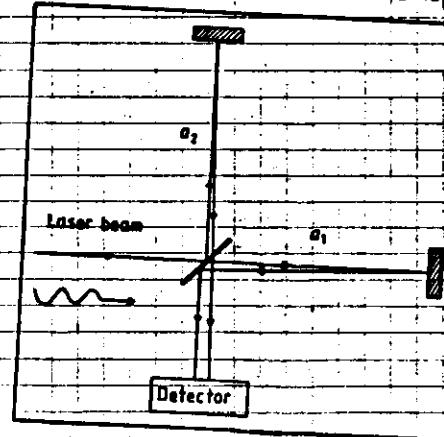
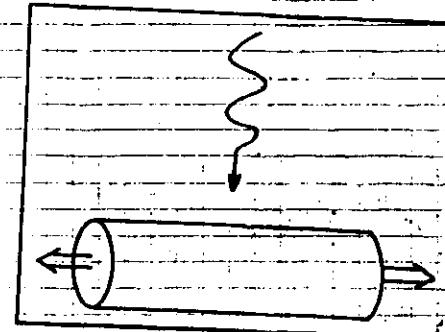


Figure 3 Interaction of a metal cylinder with a gravitational wave



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Generation of gravitational waves

$$I = \frac{G}{36\pi c^5 R_0^2} \left[\left(\frac{\ddot{D}_{11} - \ddot{D}_{22}}{2} \right)^2 + \dot{D}_{33}^2 \right] \frac{\text{Joule}}{m^2 s}$$

$D_{\alpha\beta}$ = mass quadrupole tensor

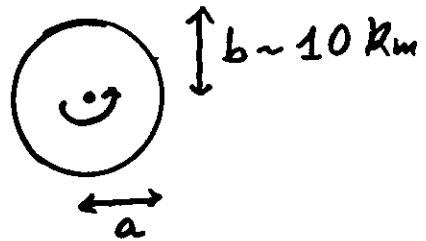
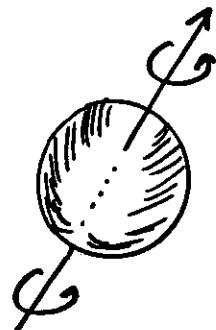
R_0 = distance from source

$$\frac{dI}{dt} = \frac{G}{45c^5} \frac{\dot{D}_{\alpha\beta}}{a^3} \frac{\text{Joule}}{s}$$

$$\frac{G}{45c^5} = 6.1 \times 10^{-55} \text{ m}^{-2} \text{ s}^3 \text{ kg}^{-1}$$

Einstein: "... in all thinkable cases it must have a practically vanishing value".

Pulsar



$$W = \underbrace{\frac{288}{45} \frac{G}{c^5} I^2 \omega^6}_{---} \frac{(a-b)^2}{ab} \text{ watt} \quad ---$$

1937+219
the fast pulsar

on the Earth

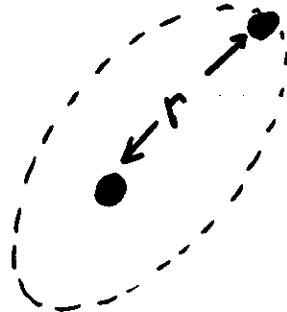
$$2.5 \text{ Kpc}$$

$$\nu_{gw} = 1283,856542 \text{ Hz}$$

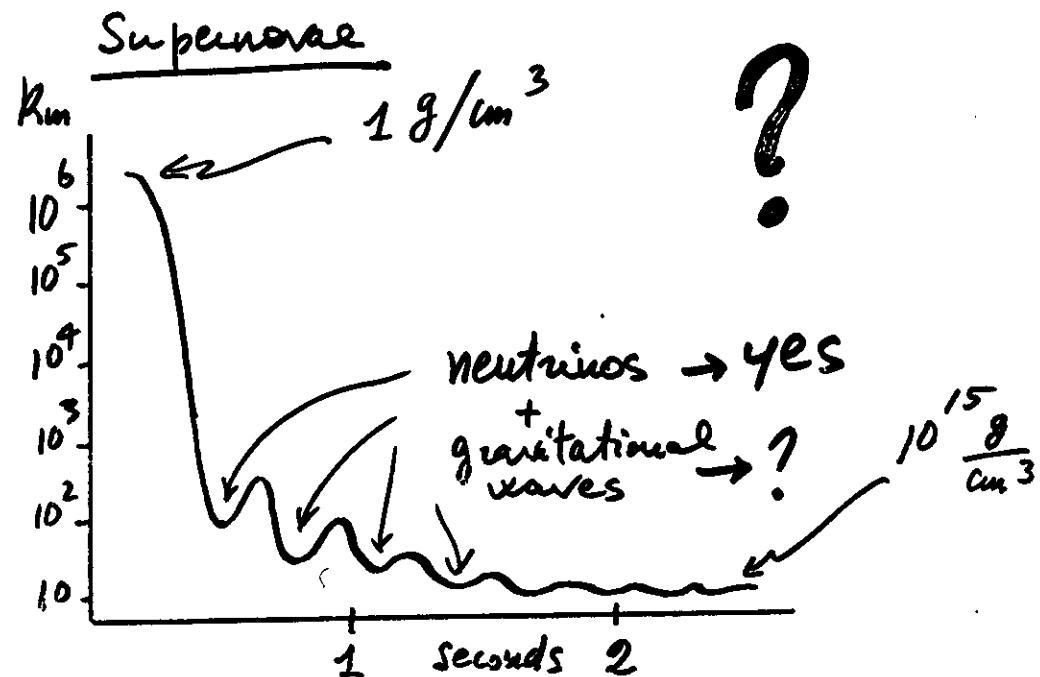
$$W = 2 \times 10^{29} \text{ watt}$$

$$h_0 \approx 3 \times 10^{-27} \text{ m} \quad |a-b| \approx 100 \mu\text{m}$$

Binary system



$$W = \frac{32}{5} \frac{G}{c^5} \left(\frac{m_1 m_2}{m_1 + m_2} \right)^2 r^4 \omega^6$$



$$M \approx 8 M_{\odot} \rightarrow 1.44 M_{\odot}$$

$$\text{e.m. waves} \approx 0.002 M_{\odot}$$

$$\text{cosmic rays} = " "$$

$$\text{neutrinos} = 0.2 M_{\odot} \approx 3 \cdot 10^{53} \text{ erg}$$

$$\text{g. w.} = (0.01 \rightarrow 0) M_{\odot} ??$$

~detected

$$\text{Duration of g. a. bursts } \tau_g = \frac{1}{2} \text{ ms}$$

PSR 1913+16 (Taylor) $T = 7^h 45^m 2^s$

$$W \approx 6.4 \times 10^{23} \text{ watt}$$

$$5 \text{ kpc}$$

$$h_0 \approx 1 \times 10^{-22}$$

Resonant antennas



$$v = \sqrt{\frac{y}{g}} = \\ 5400 \frac{\text{m}}{\text{s}} \text{ at } 4.2\text{K}$$

equivalent to many simple oscillators
with resonance angular frequencies

$$\omega_k = (2k+1) \frac{\pi\omega}{L}$$

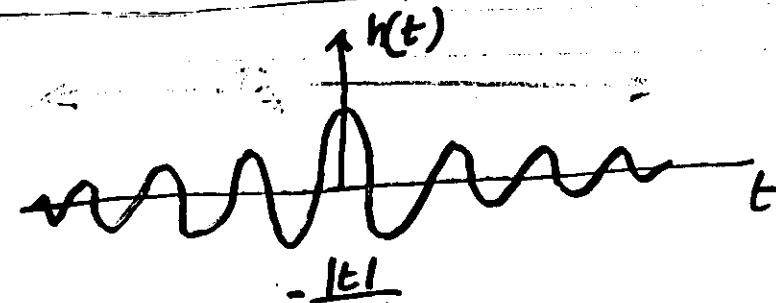
$$\omega_0 = \frac{\pi\omega}{L} \sim 2\pi(1 \div 2 \text{ KHz})$$

$$m = \frac{M}{2}$$

$$\ddot{\xi} + 2\beta_1 \dot{\xi} + \omega_0^2 \xi = \frac{2}{\pi^2} L \ddot{h}$$

$$2\beta_1 = \omega_0 / Q$$

Special wave form



$$h(t) = h_0 e^{-|t|/\tau_g/2} \cos \omega_0 t$$

$$\text{If } \tau_g \gg 2\pi/\omega_0$$

$$H(\omega_0) = \frac{h_0 \tau_g}{2}$$

$$\xi(t) = -\frac{L}{\pi^2} h_0 \tau_g \omega_0 e^{-\beta_1 t} \sin \omega_0 t$$

In general for short bursts

$$\xi(t) = -\frac{2L}{\pi^2} H(\omega_0) \omega_0 e^{-\beta_1 t} \sin \omega_0 t$$

$H(\omega_0)$ is the Fourier transform of $h(t)$
at the resonance

(ii)

Expected h values from gravitational collapses

$$\frac{mc^2}{4\pi R^2 \tau_g} = \left[\frac{\text{joule}}{\text{m}^2 \text{s}} \right] = \frac{c^3}{32\pi G} \omega^2 h^2$$

R = distance of the source

τ_g = duration of g.w. emission

$M^* c^2$ = energy of g.w.

Center of Galaxy

8.5 Kpc

$$h = 4 \times 10^{-17} \sqrt{\frac{M^*}{M_\odot} \frac{1/c_g}{1000 \text{ Hz}}}$$

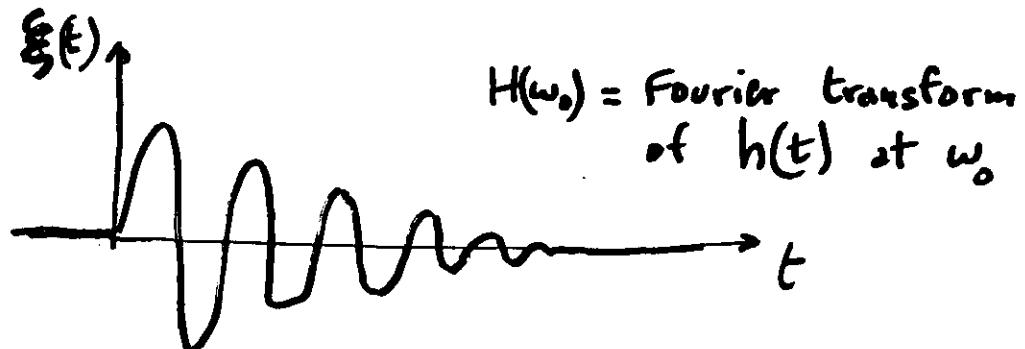
Virgo cluster

19 Mpc

$$h = 1.8 \times 10^{-20} \sqrt{\frac{M^*}{M_\odot} \frac{1/c_g}{1000 \text{ Hz}}}$$

Our goal $h \sim 10^{-21}$

$$(M \sim 10^{12} M_\odot)$$



$\xi(t)$ is the displacement of the bar end faces in absence of noise

$$\xi(t) = -\frac{2L}{\pi^2} H(\omega_0) \omega_0 e^{-\beta_0 t} \sin \omega_0 t$$

spectral energy density

$$f(\omega) = \frac{c^3}{16\pi G} |\omega H(\omega)|^2 \left[\frac{\text{joule}}{\text{m}^2 \text{Hz}} \right]$$

$$H(\omega) \approx \frac{h}{V\sigma_g} \left[\text{Hz}^{-1} \right]$$

the cross section

$$\Sigma = \frac{16}{\pi} \left(\frac{\omega}{c} \right)^2 \frac{G}{c} M \left[\text{m}^2 \text{Hz} \right]$$

$$\left(\Sigma = 8.4 \times 10^{-25} \text{ m}^2 \text{Hz} \right)$$

Prune - CERN

the energy absorbed by the antenna

$$E = \Sigma \cdot f(\omega) \left[\text{joule} \right]$$

Signal in the g.w. antenna

Center of
Galaxy

1 M_\odot

$$\Delta E \sim 9 \times 10^{-24} \text{ Joule}$$

$$= 6.5 \text{ Kelvin}$$

Virgo
cluster

1 M_\odot

$$\Delta E \sim 1.3 \times 10^{-6} \text{ K}$$

$$\approx 10^{-29} \text{ Joule}$$

$$\approx 10^{-10} \text{ eV}$$

LMC

1 M_\odot

$$\Delta E \approx \frac{6.5}{25} \text{ K} \approx 0.25 \text{ K}$$

⑧ A interferometers

- resonant antennas

NOISE

Thermal noise $\sim kT = 1.38 \times 10^{-23}$ joule
at 1 K

With optimum data filtering

$$\text{NOISE} = \frac{2kT}{\beta Q} + 2kT_n$$

T = antenna temperature

β = fraction of energy available in the transducer

Q = merit factor

T_n = electronic amplifier noise

$$T \longrightarrow 0.03 \text{ K}$$

$$\beta = 10^{-2}$$

$$Q = 10^7$$

$$T_n \longrightarrow 10^{-7} \text{ K}$$

Today
↓

$$4.2 \text{ K}$$

$$10^{-2}$$

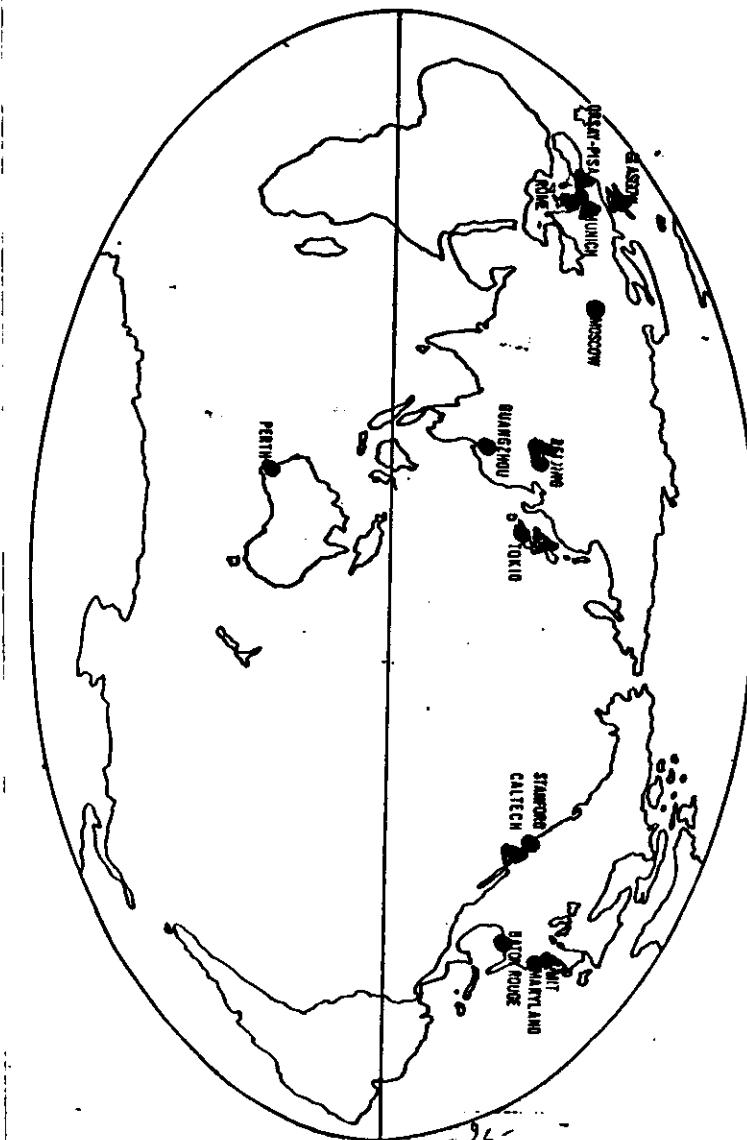
$$7 \cdot 10^6$$

$$10^{-3} \text{ K}$$

$$\text{NOISE} \approx 8 \times 10^{-30} + 2.8 \times 10^{-30} \text{ joule}$$

$$\approx 8 \times 10^{-7} \text{ K}$$

Signal $\sim 1.3 \times 10^{-6} \text{ K}$
(1 Hz from
Virgo)



NON RESONANT DETECTORS(1)

$$\Delta \ell(t) = \frac{1}{2} \ell h_+(t)$$

Shot noise $\delta \ell = \sqrt{\frac{4 \pi c^2 \Delta \nu}{\pi P}}$

Sensitivity $h_+ \approx \frac{2 \Delta \ell}{e} = 2 \sqrt{\frac{4 \pi c^2 \Delta \nu}{\pi P e^2}}$

$$P = 1 \text{ watt}$$

$$\lambda = 0.6 \mu\text{m}$$

$$\Delta \nu = 1000 \text{ Hz}$$

in order to have $h_+ \approx 1.4 \times 10^{-20}$

$$\ell = 300 \text{ km}$$

N reflections $\ell = N \ell_0$

$$N = 100$$

$$\ell_0 = 3 \text{ km}$$

time spent by light in the interferometer
 $\frac{300 \text{ km}}{300 \text{ km}} = 10^{-3} \text{ s}^{-1} \cdot \Delta \nu = 1000 \text{ Hz}$

Noise

Shot noise $\delta \ell = \sqrt{\frac{4 \pi c^2 \Delta \nu}{\pi P}}$

$$\Delta \nu = 1000 \text{ Hz}$$

$$P = 1 \text{ W}$$

$$\lambda = 0.6 \mu\text{m}$$

Electronic + brownian noise
with optimum filter

$$M_0 = \sqrt{\frac{8 k T_u}{\pi \omega_0^2}}$$

$$T_u = 10 \text{ K}$$

$$\eta = 2.3 \times 10^{-25}$$

$$\omega_0 = 50 \text{ rad/s}$$

Signal

$$\Delta \ell \approx \frac{1}{2} \ell h_+$$

$$\ell = 100 \cdot 10 \text{ km}$$

$$\Delta \ell = 10^{-15} \text{ m}$$

$$\eta_0 \approx \frac{\ell}{2} e^{-\beta \ell} \tau_g \omega_0$$

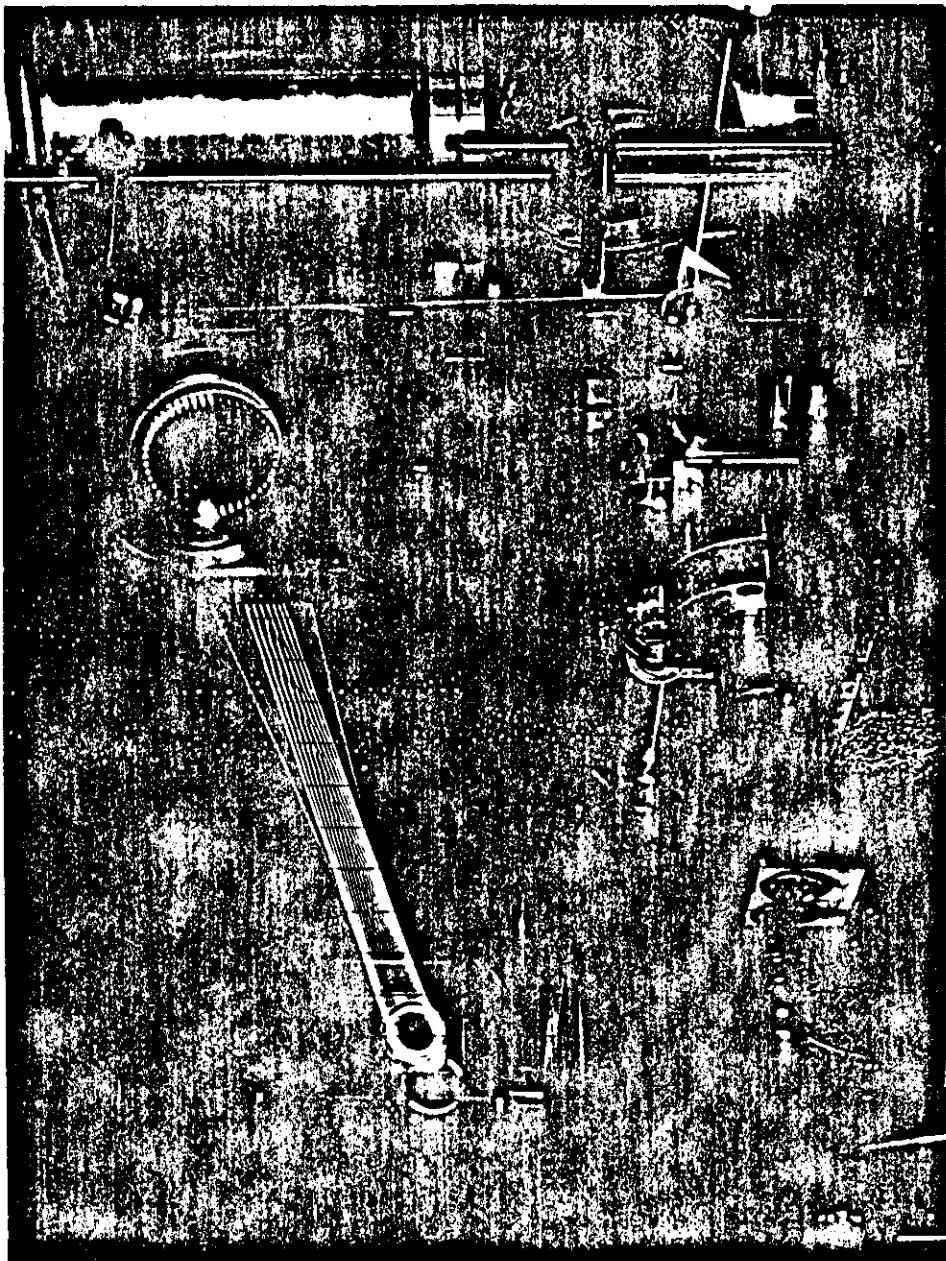
Resonant
antennas

$$\eta_0 \approx 1.5 \times 10^{-25} \text{ m}^{-2}$$

Interferometers

$$h_0 \approx 2 \cdot 10^{-21} \text{ Hz}^{-1/2}$$

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IL NUOVO CIMENTO

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Luglio-Agosto 1986

Preliminary Results on the Operation of a 2270 kg
Cryogenic Gravitational-Wave Antenna with a Resonant
Capacitive Transducer and a d.c. SQUID Amplifier.

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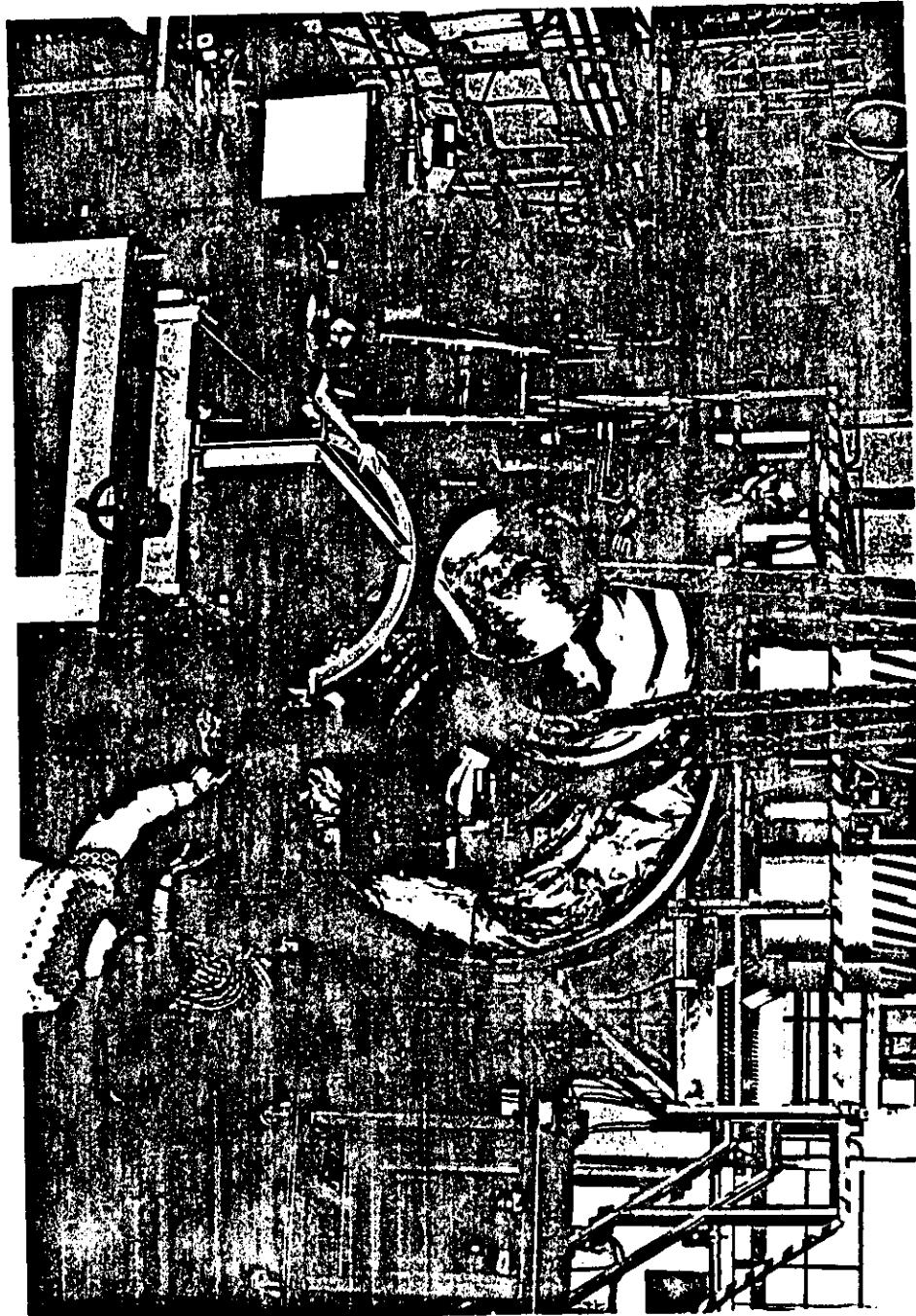
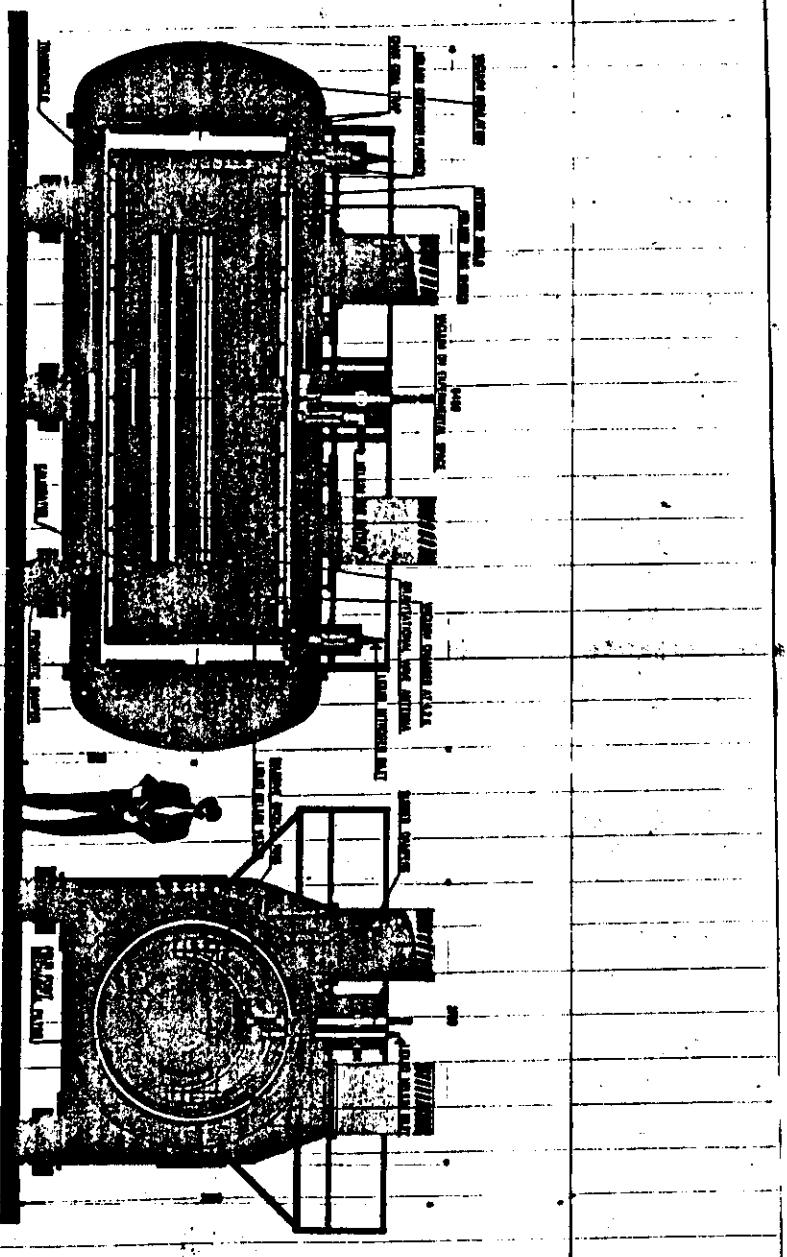
C.E.R.N., European Organisation for Nuclear Research - Genova

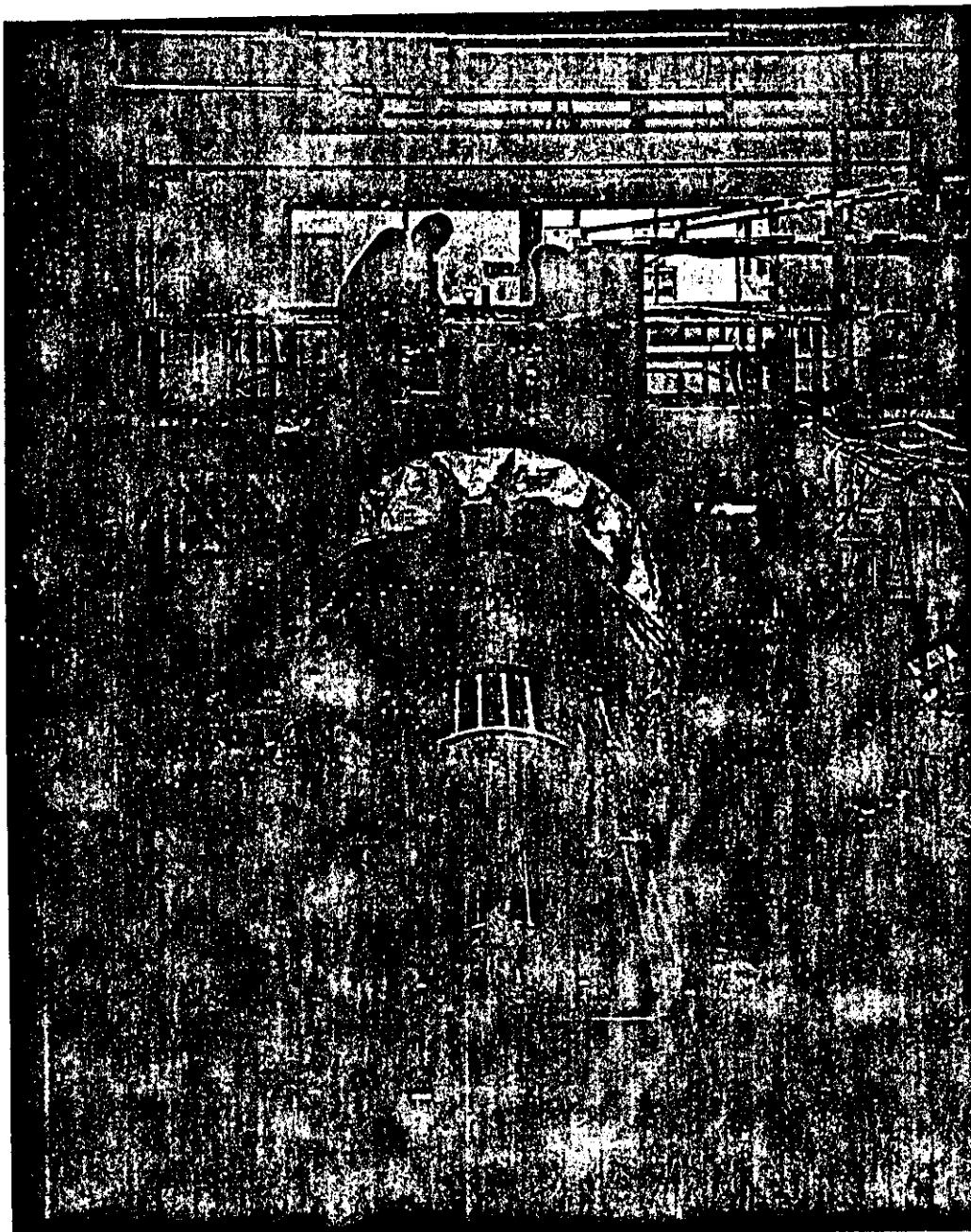
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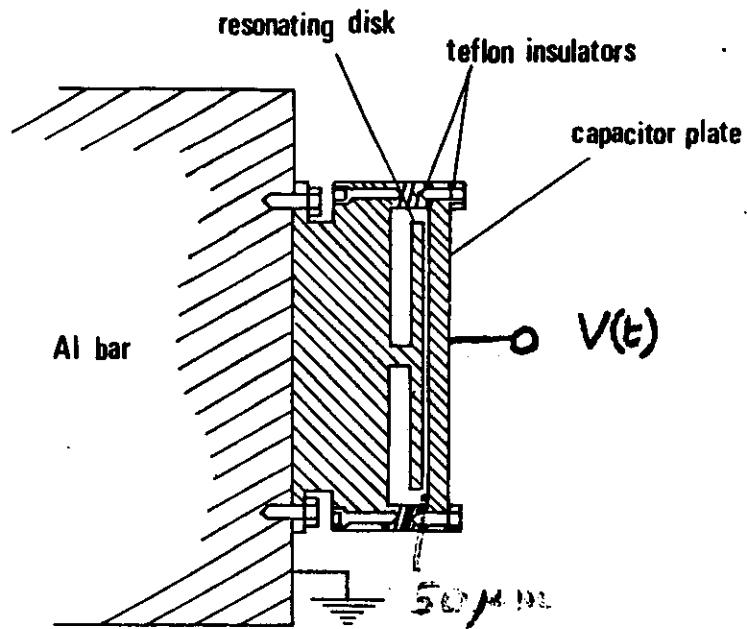
(ricevuto il 12 Maggio 1986)

Summary. — In November 1985 the gravitational-wave antenna of the Rome group, installed at CERN, has started operating. It consists of a 5056 aluminium cylinder 3 m long, 2270 kg heavy, cooled at 4.2 K. The antenna vibrations are detected by means of a resonant capacitive





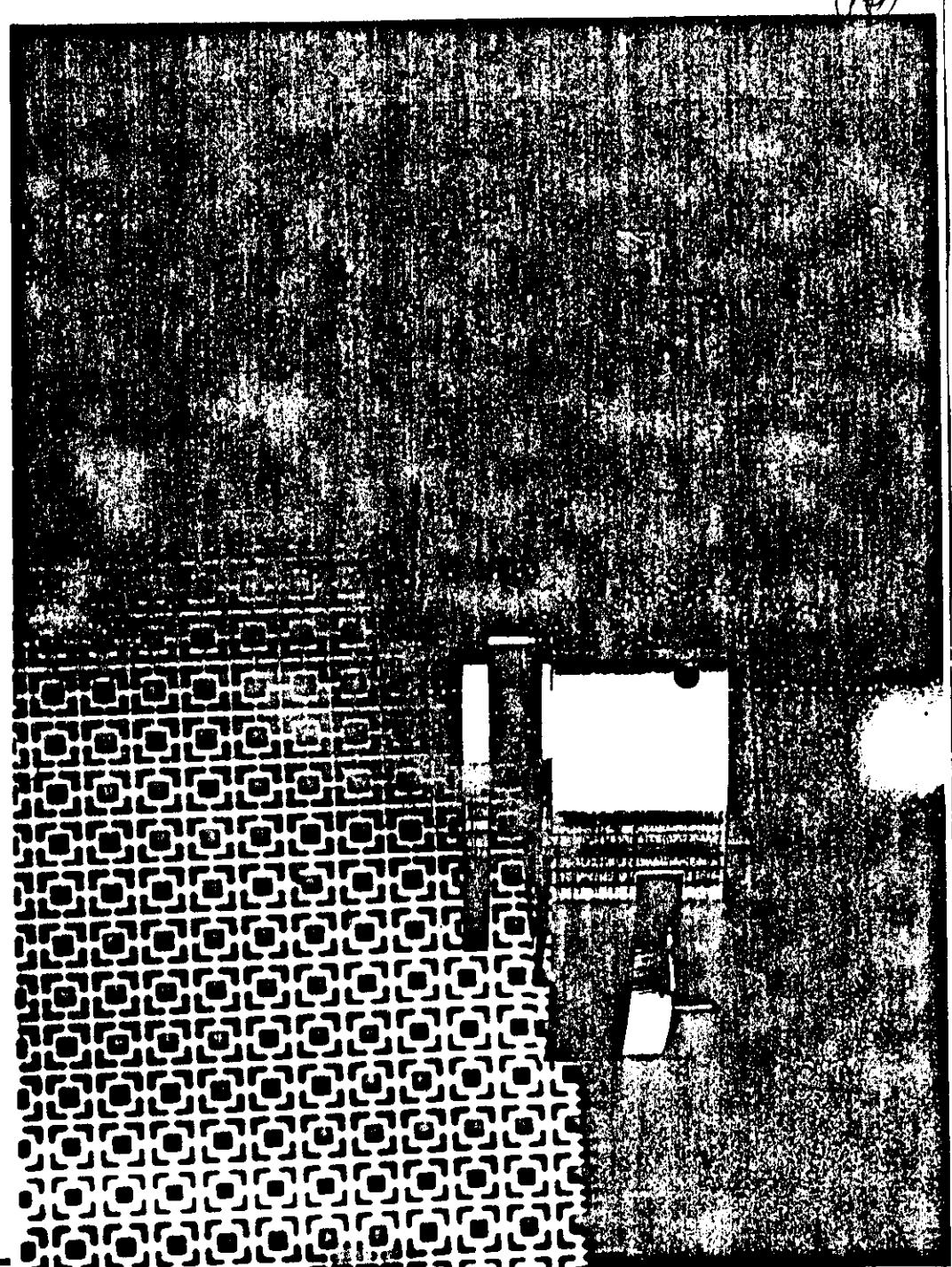
Resonant transducer



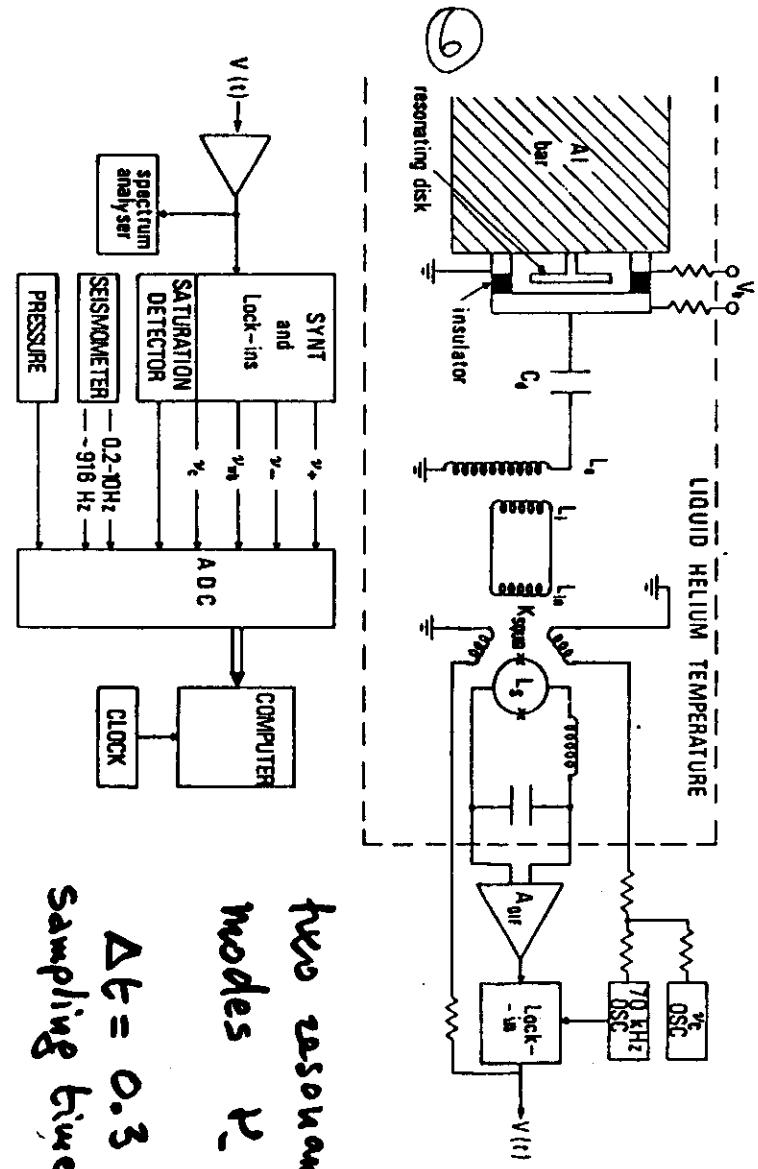
$$V(t) = \frac{Q}{C(t)} = \alpha \cdot \frac{\xi(t)}{t}$$

ξ_t is the transducer displacement

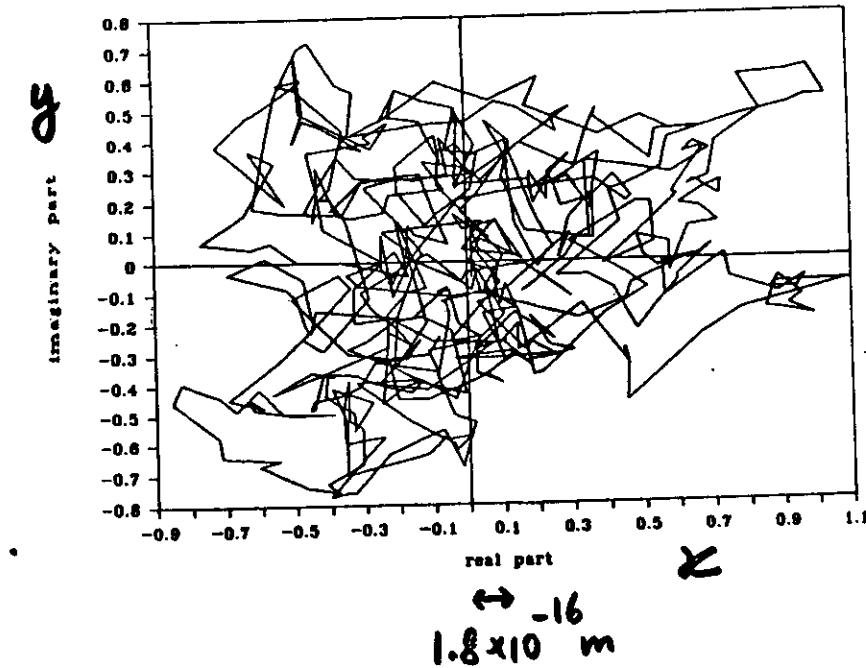
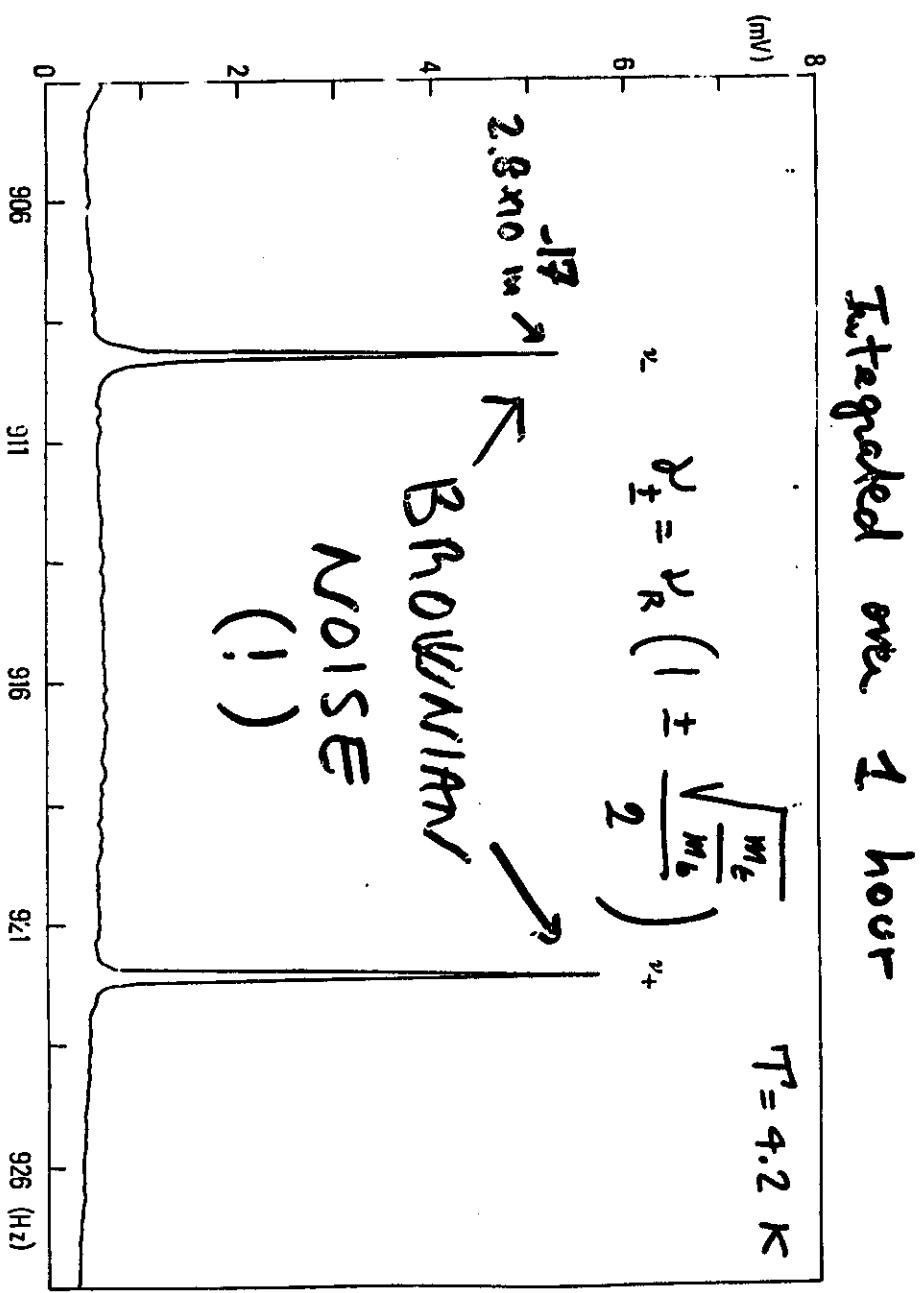


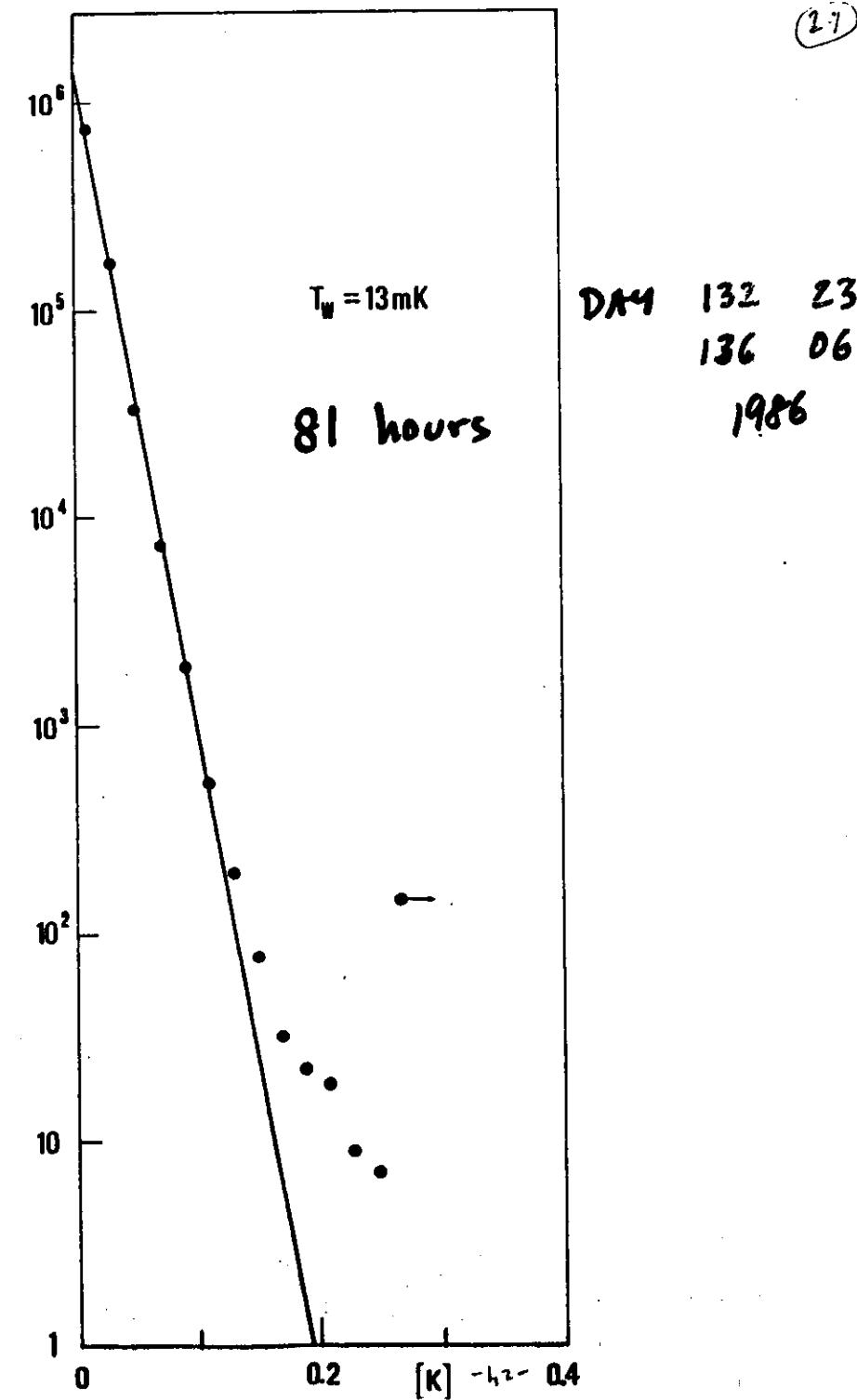
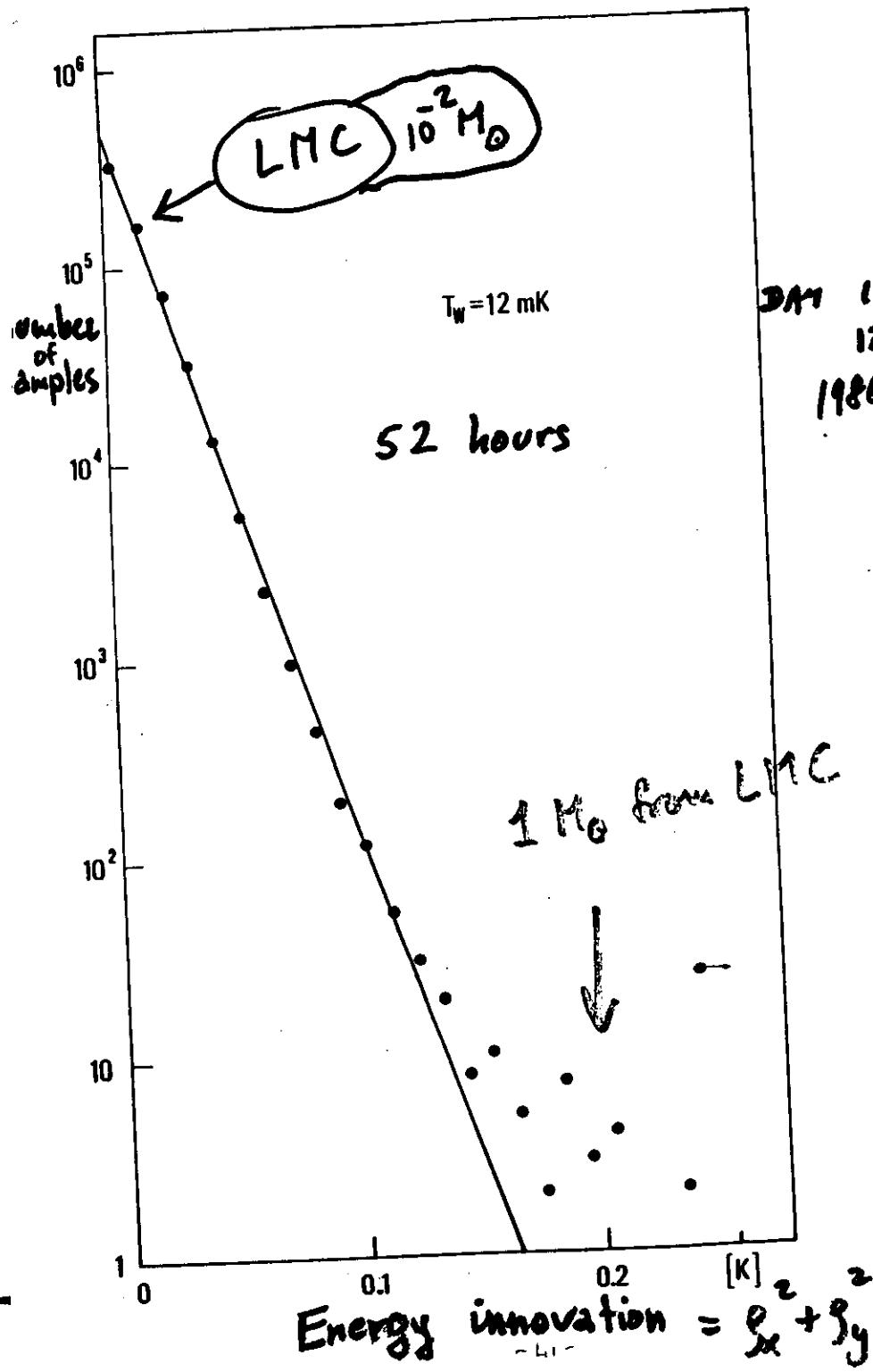


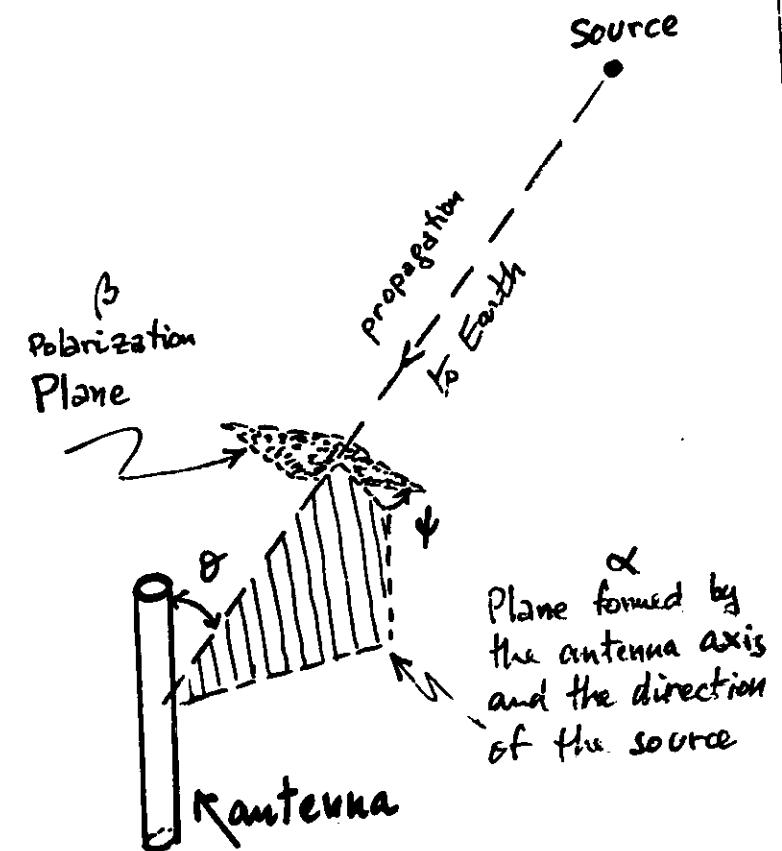
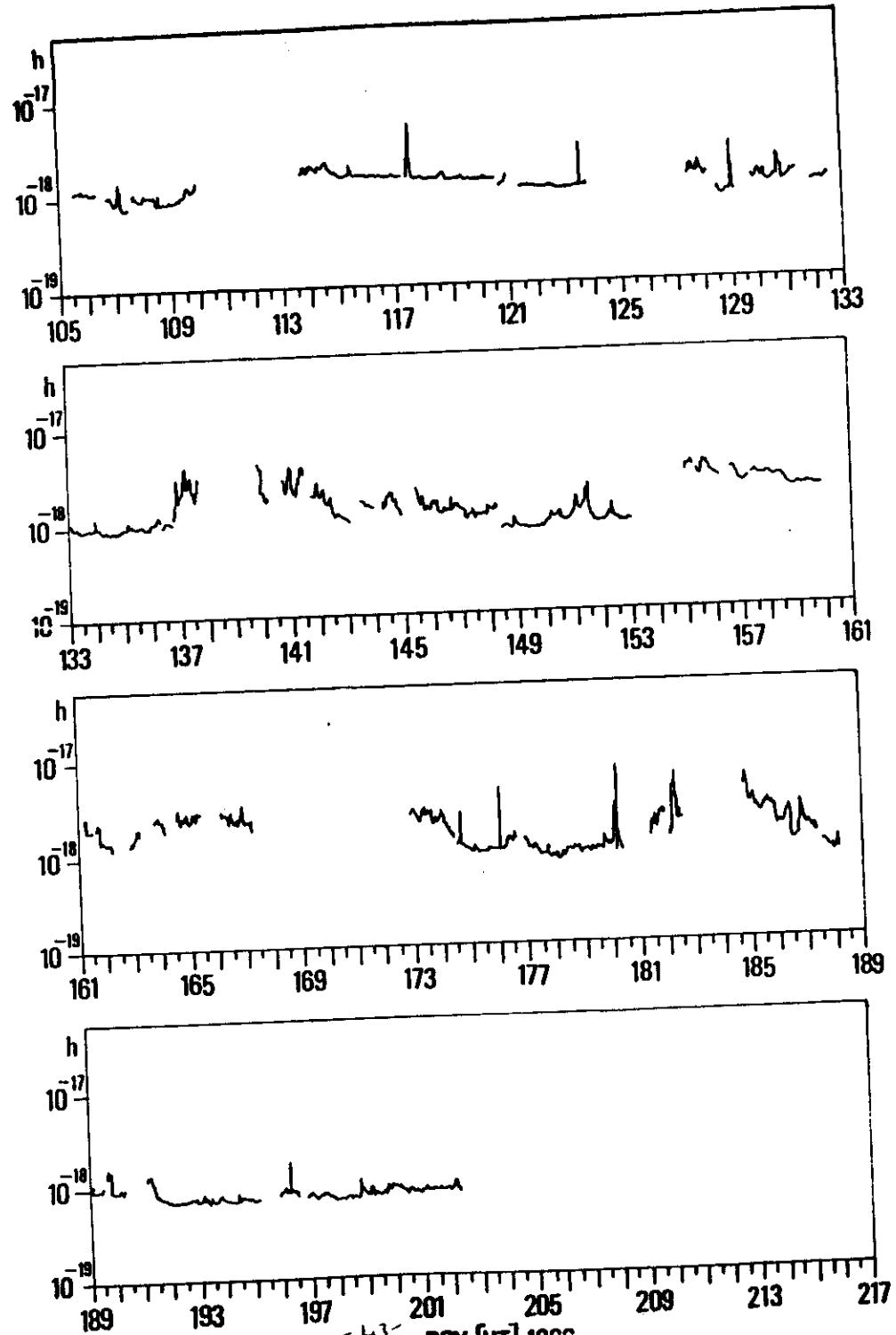
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two resonance
modes ν_+ , ν_-
 $\Delta t = 0.3$ s
Sampling time interval







$$\sum_{\chi}(\theta, \phi) = \sum \sin^4 \theta \left(\frac{1-\chi}{2} + \chi \cos^2 \phi \right)$$

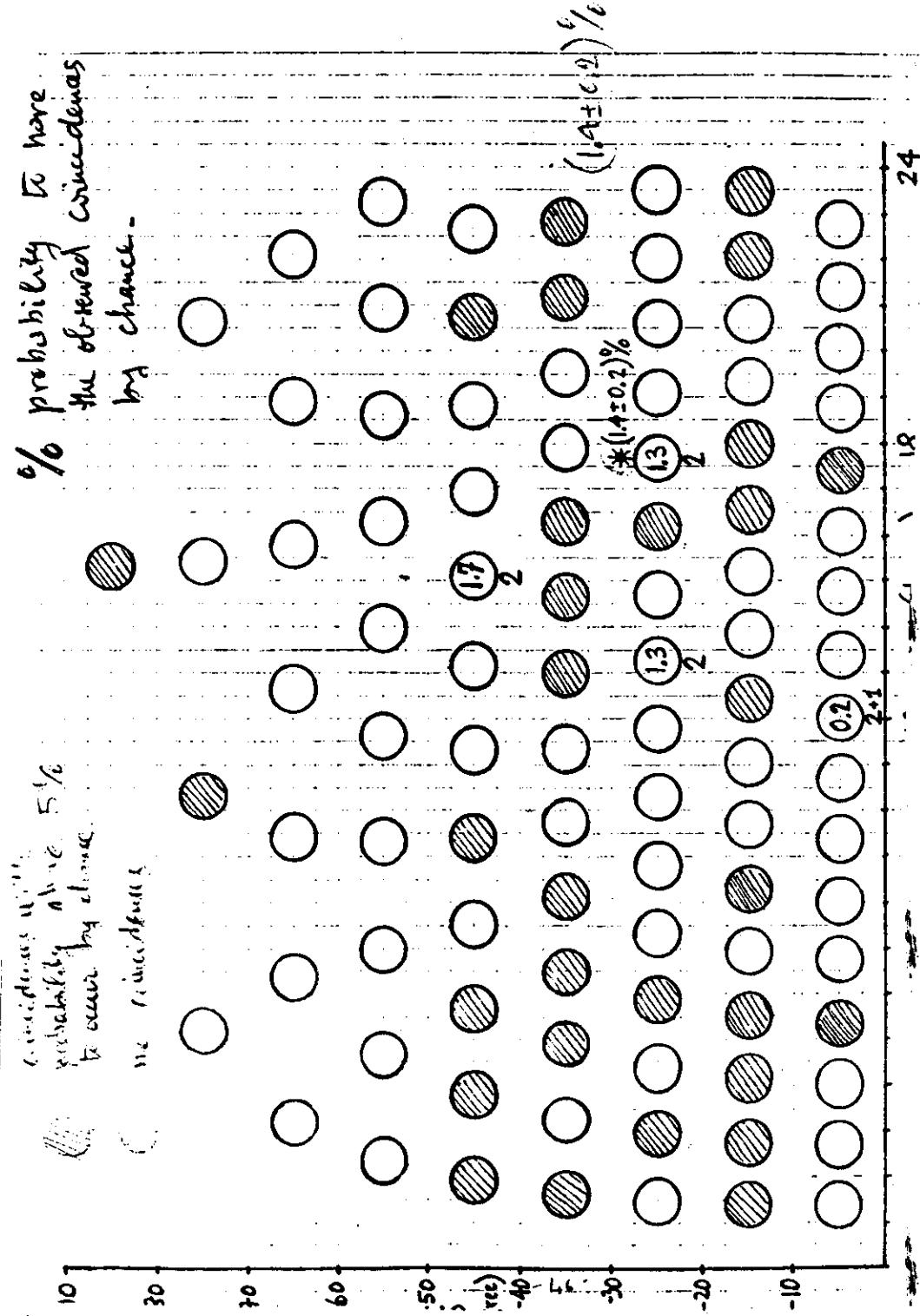
χ = fraction of polarized wave

ϕ = angle between one of the wave asymptotic lines of force and the plane α

Galactic normalization

$$\Delta E(t)_{\text{normalized}} = \frac{\Delta E(t)_{\text{measured}}}{\sin^4 \theta_W}$$

$\delta(t)$ = angle at the time t of measurement between the antenna axis and the direction of the Galactic Center.



imposing the signals
in the free antennas be equal
after normalization

