



INTERNATIONAL ATOMIC ENERGY AGENCY
UNITED NATIONS EDUCATIONAL, SCIENTIFIC AND CULTURAL ORGANIZATION



INTERNATIONAL CENTRE FOR THEORETICAL PHYSICS
34100 TRIESTE (ITALY) - P.O.B. 588 - MIRAMARE - STRADA COSTIERA 11 - TELEPHONE: 2240-1
CABLE: CENTRATOM - TELEX 460392 - I

SMR/389 - 2

**WORKING PARTY ON
MODELLING THERMOMECHANICAL BEHAVIOUR OF MATERIALS
(29 May - 16 June 1989)**

"REVIEW OF THERMOELASTICITY"

C. TOME'
Universidad Nacional de Rosario
Instituto de Fisica
Bv. 27 de Febrero 210 (Bis)
2000 Rosario
Argentina

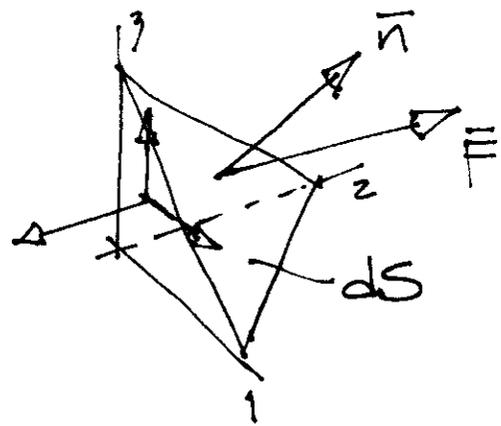
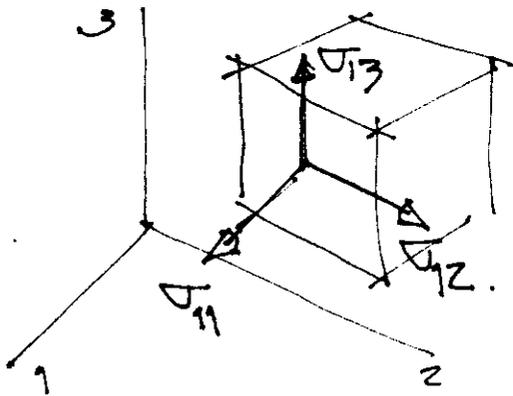
These are preliminary lecture notes, intended only for distribution to participants.

WORKING PARTY on
"MODELLING of THERMOMECHANICAL PROPERTIES"

REVIEW of THERMOELASTICITY (*) (C. TOMÉ)

1- STRESS.

Define σ_{ij} : forces per u. area acting on the 3-coordinate directions on planes parallel to the coordinate planes.



The force \vec{F} acting on a surface element dS with normal \vec{n} is given by

$$F_i = n_j \sigma_{ji} = \sum_{j=1}^3 n_j \sigma_{ji}$$

(sum over repeated indices !!).

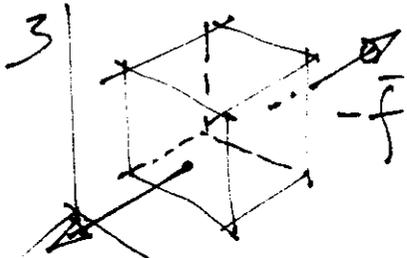
(*) see Y.C. FUNG, "FOUNDATIONS OF SOLID MECHANICS"
PRENTICE-HALL (1965).

STRESS DEVIATOR.

$$\sigma'_{ij} = \sigma_{ij} - \frac{\sigma_{kk}}{3} \delta_{ij} \quad (\text{with } \sigma_{kk} = \text{tr} \sigma)$$

EQUILIBRIUM EQUATION. (STATIC).

Equating to zero the total force acting on an elementary cube we get:



$$\frac{\partial \sigma_{ji}}{\partial x_j} + F_i = 0.$$

$$\vec{f} + \frac{\partial \vec{f}}{\partial x_i} dx_i$$

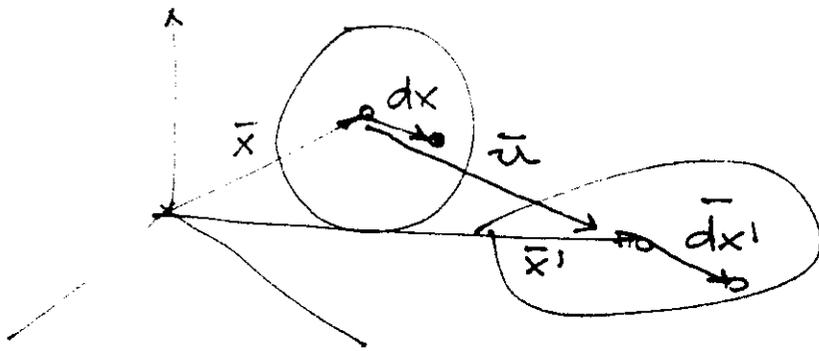
or

$$\sigma_{ji,j} + F_i = 0.$$

F_i = force per unit volume.

Equating to zero the moment of the forces (and assuming null moment per unit volume) we get

$$\sigma_{ij} = \sigma_{ji}$$



$$ds^2 = \sum_i dx_i^2$$

$$ds'^2 = \sum_i dx_i'^2$$

$$x_i' = x_i'(x_1, x_2, x_3) \quad \text{univocal \& continuous.}$$

$$dx_i' = \frac{\partial x_i'}{\partial x_j} dx_j$$

and defining: $\bar{x}' = \bar{x} + \bar{u}$ (displacement)

$$dx_i' = \left(\delta_{ij} + \frac{\partial u_i}{\partial x_j} \right) dx_j$$

$$\# \quad ds'^2 - ds^2 = \sum_{jk} 2 \epsilon_{jk} dx_j dx_k + \mathcal{O}(u_{ij}^2)$$

with $\boxed{\epsilon_{jk} = \frac{1}{2} (u_{j,k} + u_{k,j})}$ (SYMMETRIC)

being the infinitesimal strain tensor related to changes in length of the differential segment \bar{ds}

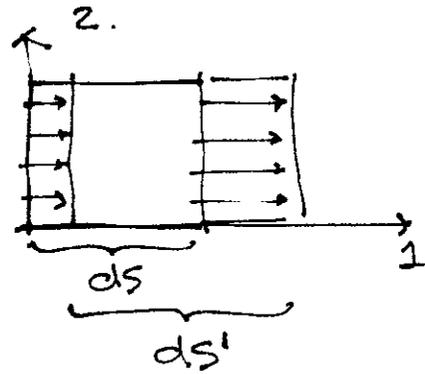
Small deformations $\rightarrow \frac{\partial u_i}{\partial x_j} \ll 1 \quad !!$

INTERPRETATION OF STRAIN COMPONENTS

choosing $\bar{dS} = (dx_1, 0, 0)$

$$\Rightarrow \epsilon_{11} = \frac{dS' - dS}{dS} = \frac{\Delta x_1}{dx_1}$$

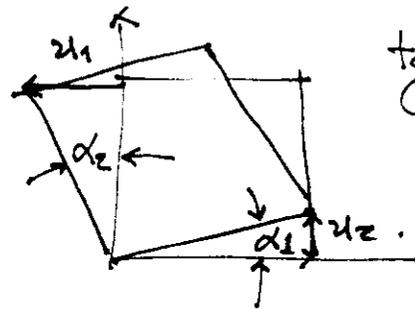
elongation.



$$\Rightarrow \epsilon_{12} = \frac{1}{2} \left(\frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right)$$

$$= \frac{1}{2} (-|\alpha_2| + |\alpha_1|)$$

$= \frac{1}{2}$ angular distortion of element.

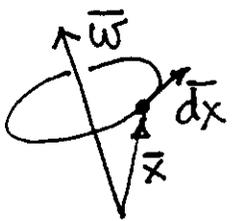


$$\tan \alpha_2 = \frac{u_1}{x_2} - \alpha_2$$

if $|\alpha_2| = |\alpha_1| \Rightarrow \epsilon_{12} = 0$.

but $w_{12} = \frac{1}{2} \left(\frac{\partial u_1}{\partial x_2} - \frac{\partial u_2}{\partial x_1} \right) = |\alpha_1| =$

$=$ rigid rotation of element.



$$w_{ij} = \frac{1}{2} (u_{ij} - u_{ji}) = \text{rotation tensor (ANTISYMMETRIC)}$$

$$w_k = \frac{1}{2} \epsilon_{kij} w_{ij} = \frac{1}{2} (\nabla \times \bar{u})_k = \text{rotation vector}$$

\hookrightarrow Levi-Civita tensor.

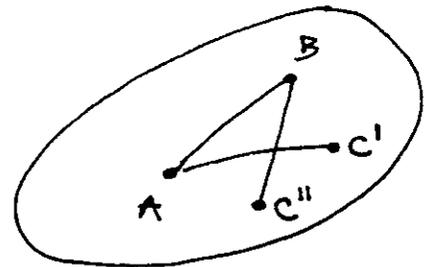
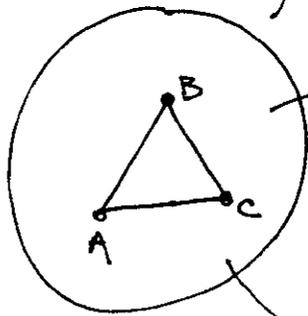
COMPATIBILITY EQUATIONS.

$$\epsilon_{ij} = \frac{1}{2} (u_{i,j} + u_{j,i}).$$

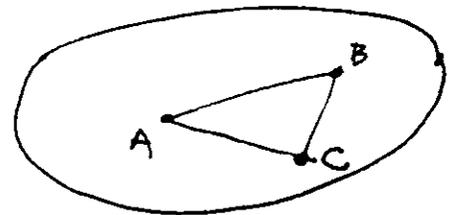
$$\Rightarrow \epsilon_{ij,kl} + \epsilon_{kl,ij} - \epsilon_{ik,jl} - \epsilon_{jl,ik} \equiv 0.$$

a total of 6 independent equations
due to redundancy in definition
of strain tensor.

Necessary for avoiding:



and getting:



CONSTITUTIVE EQUATION.

(Hooke + Cauchy law).

$$\sigma_{ij} = C_{ijkl} \epsilon_{kl}.$$

with C_{ijkl} = elastic stiffness tensor.
(81 constants).

$$\sigma_{ij} = \sigma_{ji} \Rightarrow C_{ijkl} = C_{jilk}$$

$$\epsilon_{kl} = \epsilon_{lk} \Rightarrow C_{ijkl} = C_{ijlk}$$

$$\frac{\partial W}{\partial \epsilon_{nm}} = \sigma_{nm} \Rightarrow C_{ijkl} = C_{klij}$$

(21 indep. constants).

CONTRACTED VOIGT NOTATION.

$$11 \rightarrow 1$$

$$22 \rightarrow 2$$

$$33 \rightarrow 3$$

$$23 \rightarrow 4$$

$$13 \rightarrow 5$$

$$12 \rightarrow 6$$

$$\sigma_{ij} = C_{ijkl} \epsilon_{kl} \quad (i, j, k, l = 1, 2, 3)$$

$$\sigma_n = C_{nm} \epsilon_m \quad (n, m = 1, \dots, 6).$$

but pseudo-tensors !!

Do not transform as vectors !!

$$\sigma_{ij}^T = a_{is} a_{jt} \sigma_{st} \quad \text{but} \quad \sigma_n^T \neq a_{nm} \sigma_m !!$$

CUBIC SYMMETRY.

$$C = \begin{bmatrix} C_{11} & C_{12} & C_{12} & | & 0 \\ & C_{11} & C_{12} & | & 0 \\ & & C_{11} & | & 0 \\ \hline & & & | & C_{44} \\ & 0 & & | & C_{44} \\ & & & | & C_{44} \end{bmatrix}$$

3 indep. constants ✓

HEXAGONAL & TRIGONAL SYMMETRY.

$$C = \begin{bmatrix} C_{11} & C_{12} & C_{13} & | & 0 \\ & C_{11} & C_{13} & | & 0 \\ & & C_{33} & | & 0 \\ \hline & & & | & C_{44} \\ & 0 & & | & C_{44} \\ & & & | & \frac{C_{11}-C_{12}}{2} \end{bmatrix}$$

5 indep. cts.

ISOTROPIC ELASTICITY.

$$C = \begin{bmatrix} \lambda + 2\mu & \mu & \mu & | & 0 \\ & \lambda + 2\mu & \mu & | & 0 \\ & & \lambda + 2\mu & | & 0 \\ \hline & & & | & \mu \\ & 0 & & | & \mu \\ & & & | & \mu \end{bmatrix}$$

2 indep. cts ✓

(λ, μ : Lamé's coefficients).

4 - ISOTROPIC ELASTICITY EQUATIONS (STATIC)

a) $\sigma_{ij} = \lambda \epsilon_{kk} \delta_{ij} + 2\mu \epsilon_{ij}$ (HOOKE'S).

b) $\sigma_{ij,j} + F_i = 0$ (EQUILIBRIUM).

c) $\epsilon_{ij} = \frac{1}{2} (u_{i,j} + u_{j,i})$

a+b+c \Rightarrow
$$2\mu u_{ij,j} + (\lambda + \mu) u_{j,ji} + F_i = 0.$$
 NAVIER EQ.

d) BOUNDARY CONDITIONS

$f_i^s = \sigma_{ij} n_j = (\lambda u_{kk} \delta_{ij} + \mu (u_{i,j} + u_{j,i})) n_j$
on the surface. or

u_i^s imposed on the surface.



Observe that compatibility eqs. are not used.

(≠ STATIONARY)

$$a) \sigma_{ij} = C_{ijkl} \epsilon_{kl} - \beta_{ij} (T - T_0)$$

where $\beta_{ij} = \beta_{ji} \rightarrow$ thermal moduli

and $T_0 \rightarrow$ reference state (stress free).

and $T = T(x_1, x_2, x_3)$.

$$b) h_i = -k_{ij} \frac{\partial T}{\partial x_j} = \text{heat flux.}$$

where $k_{ij} =$ heat conduction coefficients.

$$c) \nabla \cdot \bar{h} = h_{i,i} = 0 \quad (\text{conservation}).$$

ISOTROPIC THERMOELASTICITY:

$$a) \sigma_{ij} = \lambda \epsilon_{kk} \delta_{ij} + 2\mu \epsilon_{ij} - \beta (T - T_0) \delta_{ij}$$

with $\beta = (3\lambda + 2\mu) \alpha$ (linear coeff thermal. exp.)

$$b) h_i = -k \frac{\partial T}{\partial x_j} \quad (\text{or } \bar{h} = -k \nabla T)$$

(uncoupled elastic and thermal behavior)