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H4-SMR 393/11

SPRING COLLEGE ON PLASMA PHYSICS

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QUASILINEAR THEORY

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QUASILINEAR THEORY

Ronchi, Similon, Sudan (1989)

$$n = n_L + n_s$$

$$\varphi = \varphi_L + \varphi_s$$

$$\langle \varphi_s \rangle = \langle n_s \rangle = 0$$

$$\frac{\partial n_L}{\partial t} + \frac{c}{B_0} \hat{y} \times \nabla \varphi_L \cdot \nabla n_L = \nabla \cdot [\mu_{Le} n_L \nabla \varphi_L + D_e \nabla n_L] - \nabla \cdot \Gamma + Q - \alpha n_L^2$$

$$\Gamma = -\mu_{Le} \langle n_s \nabla \varphi_s \rangle - \mu_{He} \langle n_s \nabla \varphi_s \times \hat{y} \rangle$$

$$\begin{aligned} \frac{\partial n_s}{\partial t} + \frac{c}{B_0} \hat{y} \times \nabla \varphi_L \cdot \nabla n_s + \frac{c}{B} \nabla n_L \times \hat{y} \cdot \nabla \varphi_s \\ = \nabla \cdot [\mu_{Le} n_L \nabla \varphi_s + D_e \nabla n_s] \end{aligned}$$

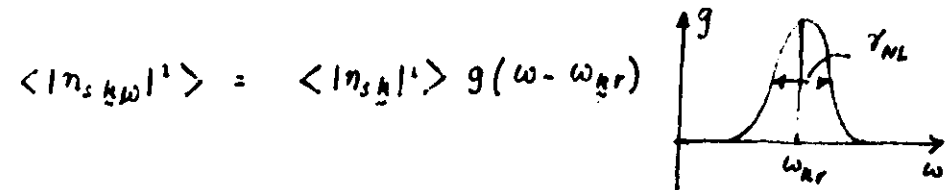
Neglect non-linear terms in Eqn. for n_s

$$\varphi_{s, k\omega} = \frac{\omega + \mu_{He} \nabla \varphi_L \times \hat{y} \cdot \underline{k} + i k^2 D_e}{\mu_{He} (\nabla n_L \times \hat{y}) \cdot \underline{k} - i k^2 \mu_{Le} n_L} n_{s, k, \omega}$$

$$\mu_{He} = \frac{c}{B_0}$$

$$\langle n_s \nabla \varphi_s \rangle$$

$$= \int d^3k d\omega \underline{k} \langle |n_{s, k\omega}|^2 \rangle \frac{k^2 D_e \frac{c}{B_0} \underline{k} \times \hat{y} \cdot \nabla n_L - k^2 \mu_{Le} n_L (\omega + \frac{c}{B_0} \nabla \varphi_L \times \hat{y} \cdot \underline{k})}{[\frac{c}{B_0} \nabla n_L \times \hat{y} \cdot \underline{k}]^2 + [\mu_{Le} n_L k^2]^2}$$



$$\langle |n_{s, k\omega}|^2 \rangle = \langle |n_{s, k}|^2 \rangle g(\omega - \omega_{kr})$$

$$\langle n_s \nabla \varphi_s \rangle$$

$$= \int \frac{d^3k}{k^2} \underline{k} \underline{k} \times \hat{y} \langle | \frac{n_s}{n_L} |^2 \rangle \left[\frac{\gamma}{1+\gamma} \frac{\Omega_e}{\omega_e} n_L \nabla \varphi_L - \frac{(2+\gamma+\gamma^2)}{1+\gamma} \frac{\gamma}{c} \frac{\Omega_e}{\omega_e} \nabla n_L \right]$$

$$\begin{aligned} = -\frac{1}{2} \frac{\gamma}{1+\gamma} \frac{\Omega_e}{\omega_e} \langle | \frac{\delta n(z)}{n_L} |^2 \rangle n_L \nabla \varphi_L \times \hat{y} \\ + \frac{1}{2} \frac{(2+\gamma+\gamma^2)}{(1+\gamma)^2} \frac{\gamma}{c} \frac{\Omega_e}{\omega_e} \langle | \frac{\delta n(z)}{n_L} |^2 \rangle \nabla n_L \times \hat{y} \end{aligned}$$

$$\therefore \int \frac{d^3k}{k^2} \langle | \frac{n_s}{n_L} |^2 \rangle \underline{k} \underline{k} \times \hat{y} = -\frac{1}{2} \frac{\gamma}{c} \times \hat{y} \langle | \frac{\delta n(z)}{n_L} |^2 \rangle$$

$$\underline{\Gamma} = -\mu_H^* n_L \nabla \phi_L \times \hat{y} - D_H^* \nabla n_L \times \hat{y} - \mu_L^* n_L \nabla \phi_L - D_L^* \nabla n_L$$

$$\mu_H^* = \frac{1}{2} \frac{\psi}{1+\psi} \left\langle \left| \frac{\delta n}{n_L} \right|^2 \right\rangle_s \mu_{He}$$

$$D_H^* = \frac{1}{2} \frac{(2+\psi+\psi^2)}{(1+\psi)^2} \left\langle \left| \frac{\delta n}{n_L} \right|^2 \right\rangle_s \frac{T}{e} \mu_{He}$$

$$\mu_L^* = \frac{1}{2} \frac{\psi}{1+\psi} \left(\frac{\mu_{He}}{\mu_{Le}} \right)^2 \left\langle \left| \frac{\delta n}{n_L} \right|^2 \right\rangle_s \mu_{Le}$$

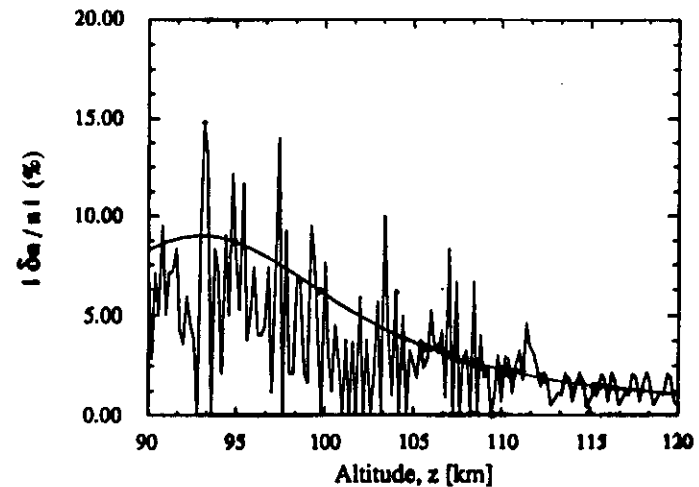
$$D_L^* = \frac{1}{2} \frac{(2+\psi+\psi^2)}{(1+\psi)^2} \left(\frac{\mu_{He}}{\mu_{Le}} \right)^2 \left\langle \left| \frac{\delta n}{n_L} \right|^2 \right\rangle_s D_e$$

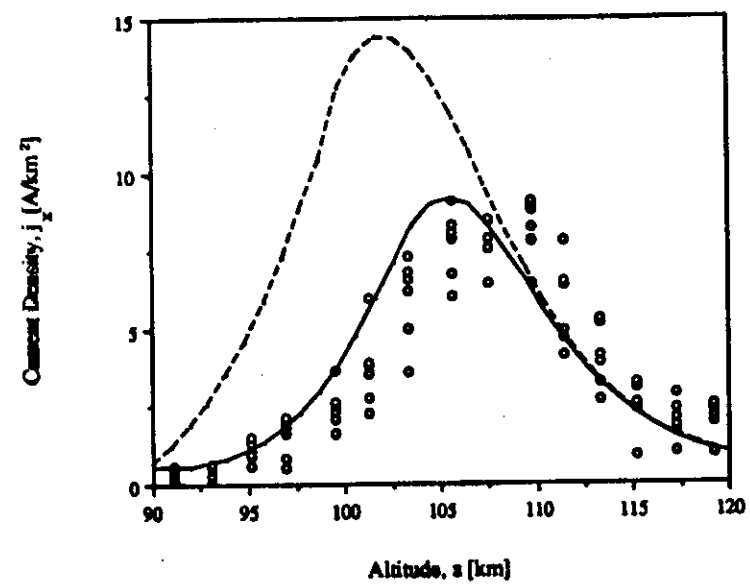
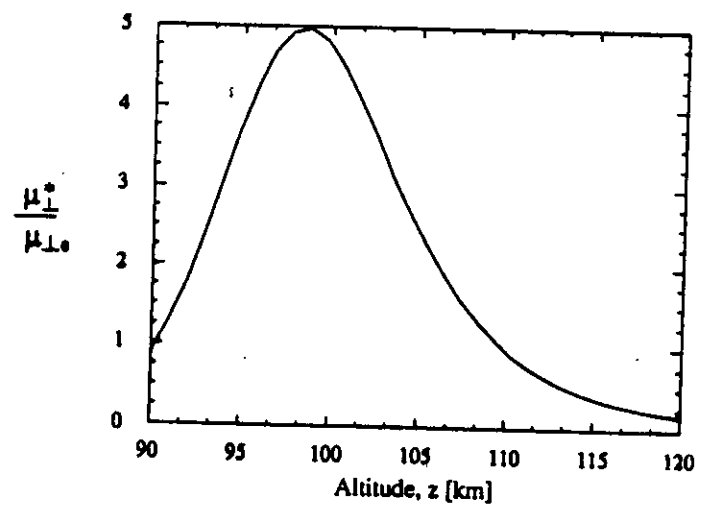
$$\left(\frac{\mu_{He}}{\mu_{Le}} \right)^2 = \frac{\Omega_e^2}{\nu_e^2}$$

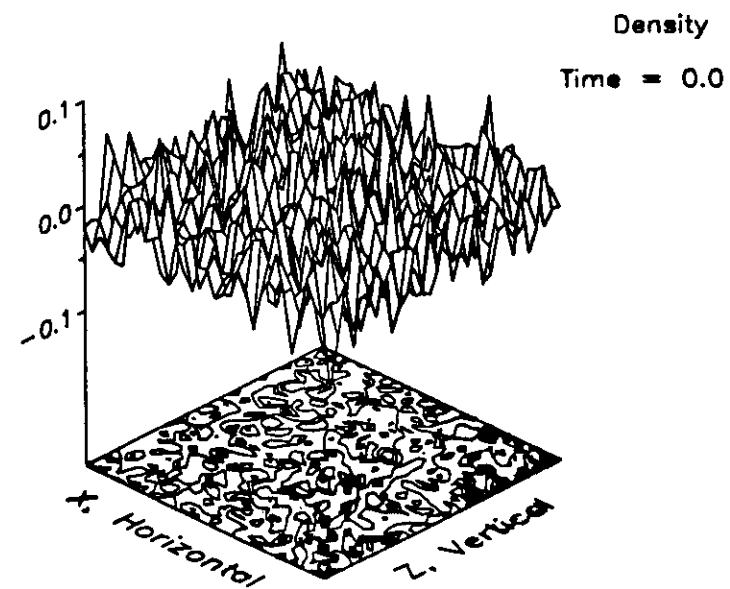
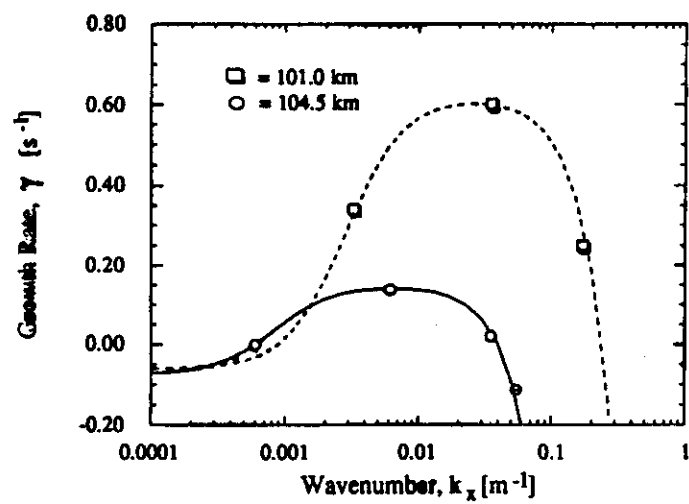
Finally :

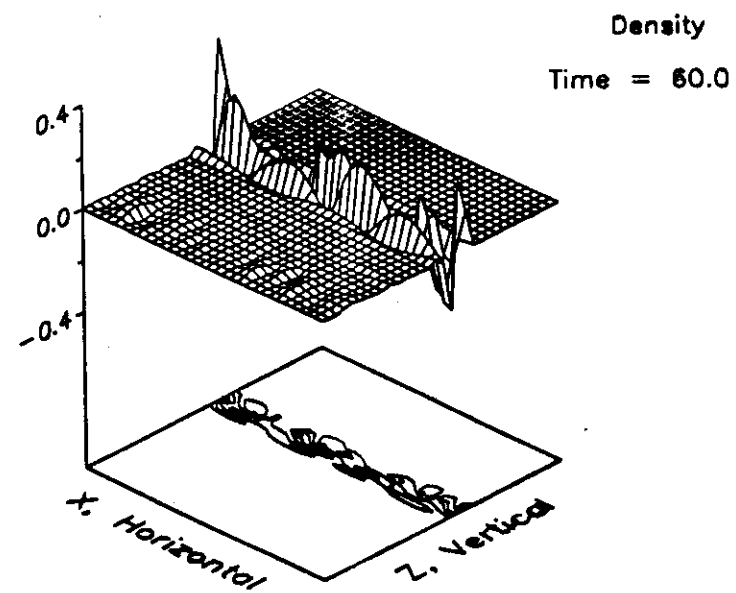
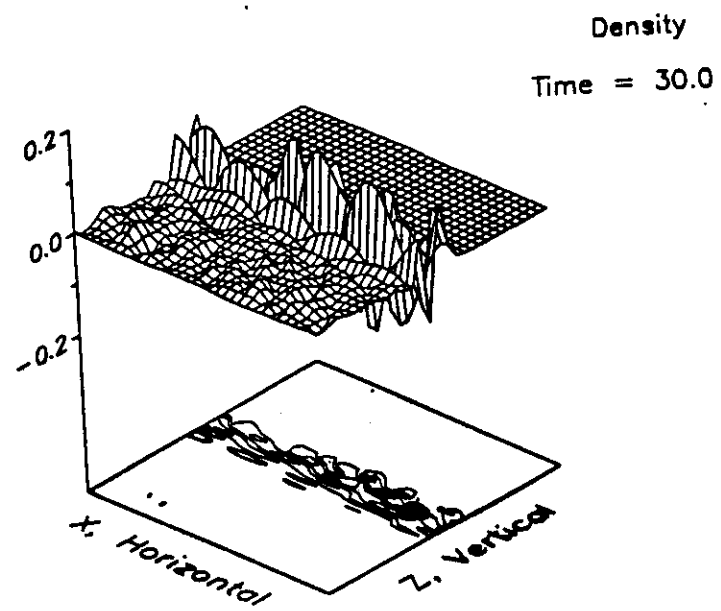
$$\frac{\partial n_L}{\partial t} + \frac{c}{B_0} \hat{y} \times \nabla \phi_L \cdot \nabla n_L = \nabla \cdot [(\mu_{Le} + \mu_L^*) n_L \nabla \phi_L] + \nabla \cdot [(D_e + D_L^*) \nabla n_L] + Q - \alpha n_L^2$$

$$\mu_H^* \nabla n_L \cdot \nabla \phi_L \times \hat{y} + \nabla \cdot [(\mu_L + \mu_L^*) n_L \nabla \phi_L] + \nabla \cdot (D_L - D_e - D_L^*) \cdot \nabla n_L = 0$$









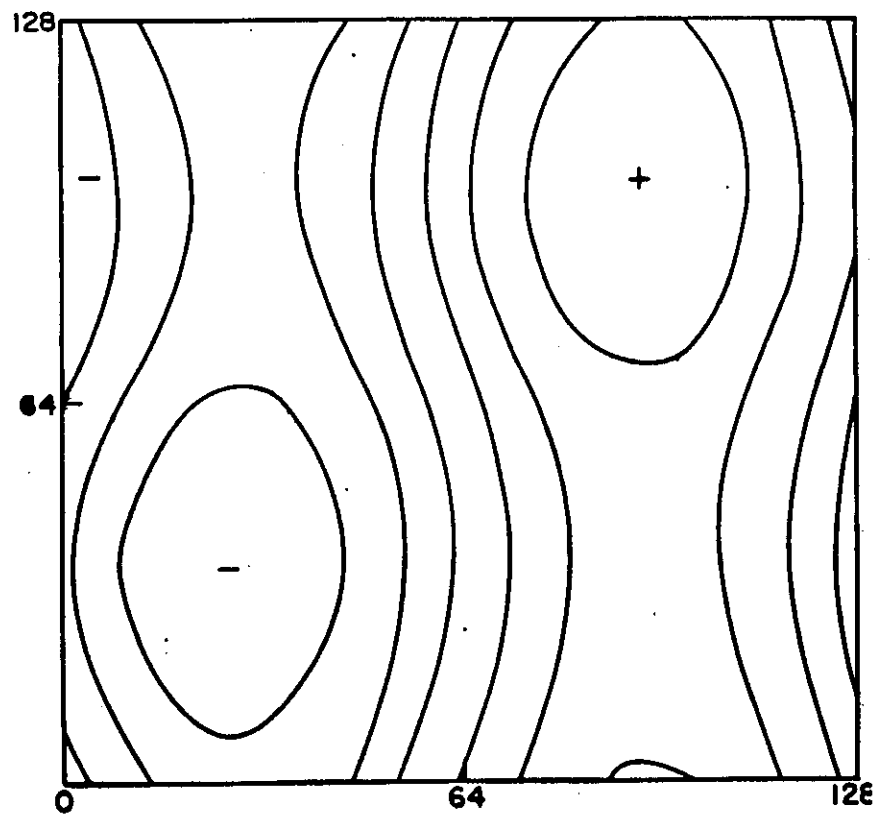


Fig. 3.2a

Keskinen, Sudan, Ferch J.G.R. 24, 1419
(1979)

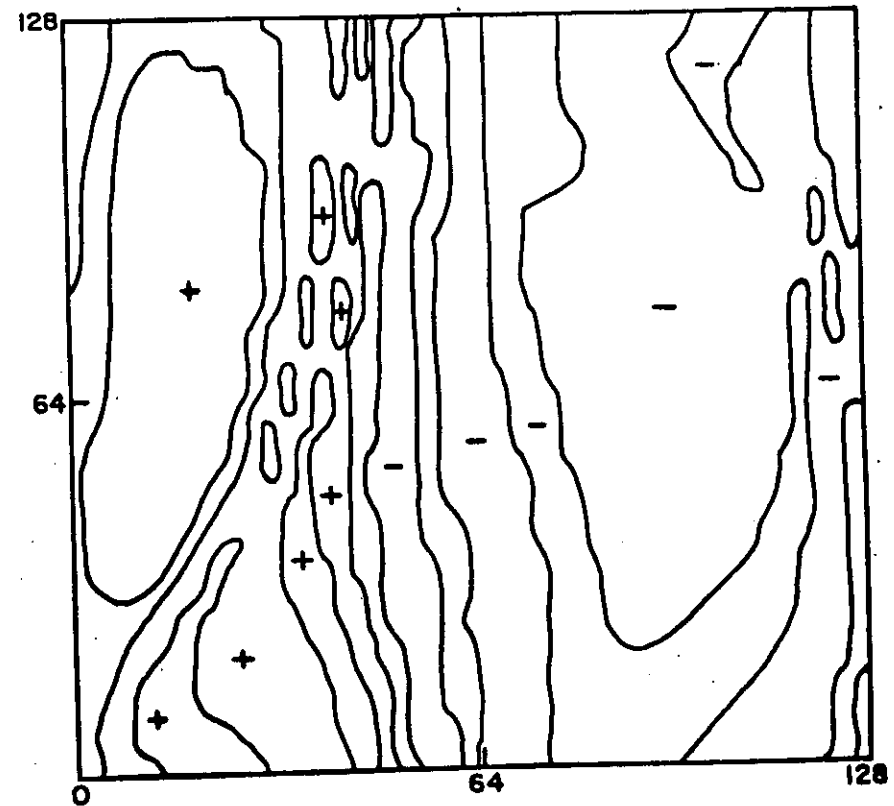


Fig. 3.2b

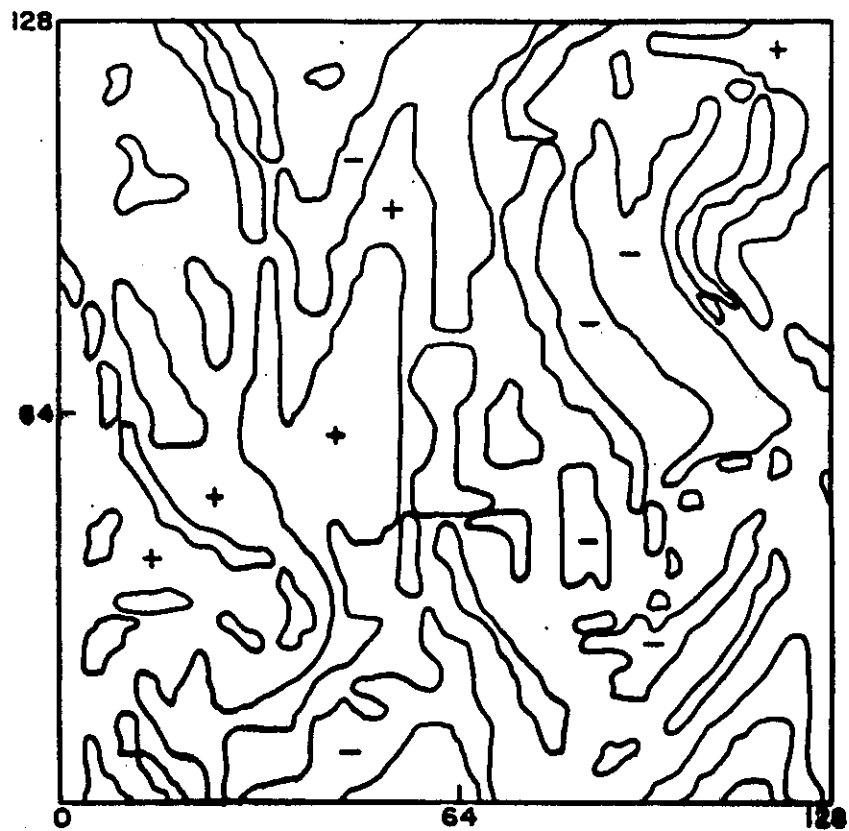
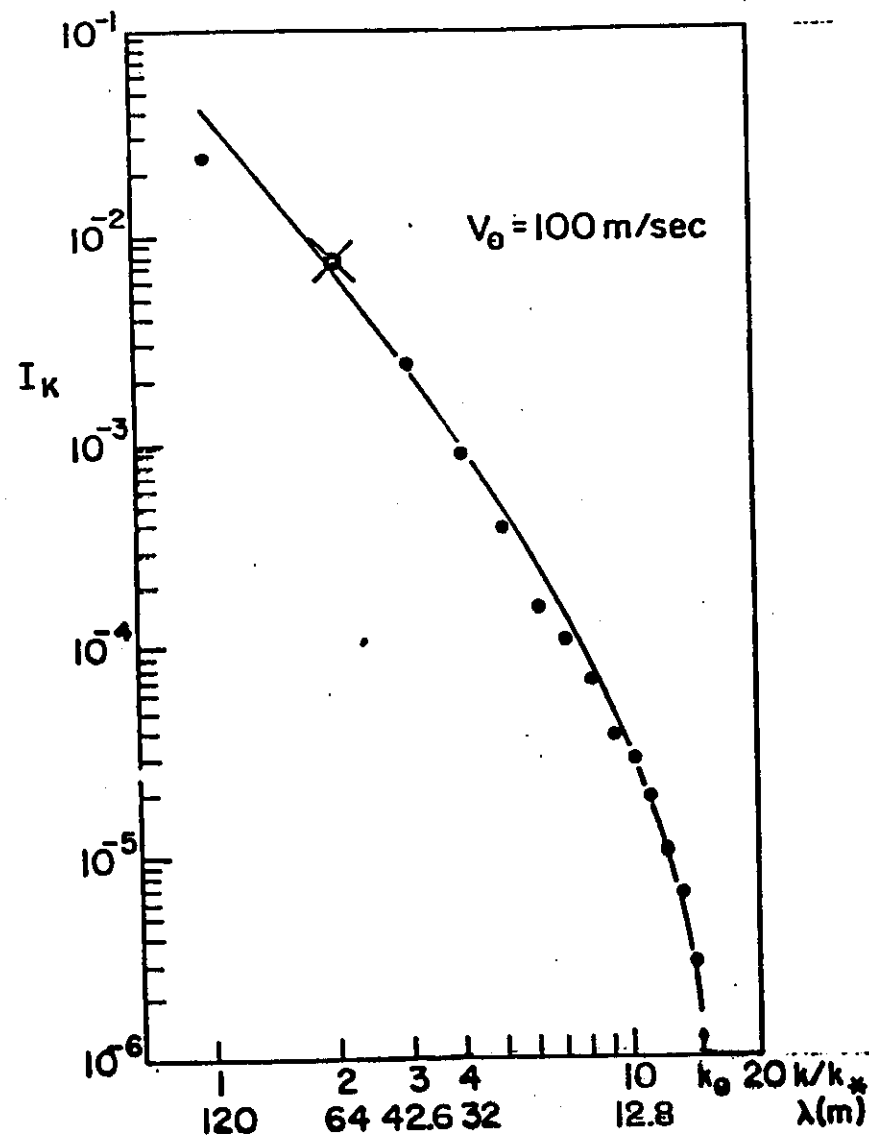


Fig. 5.2r



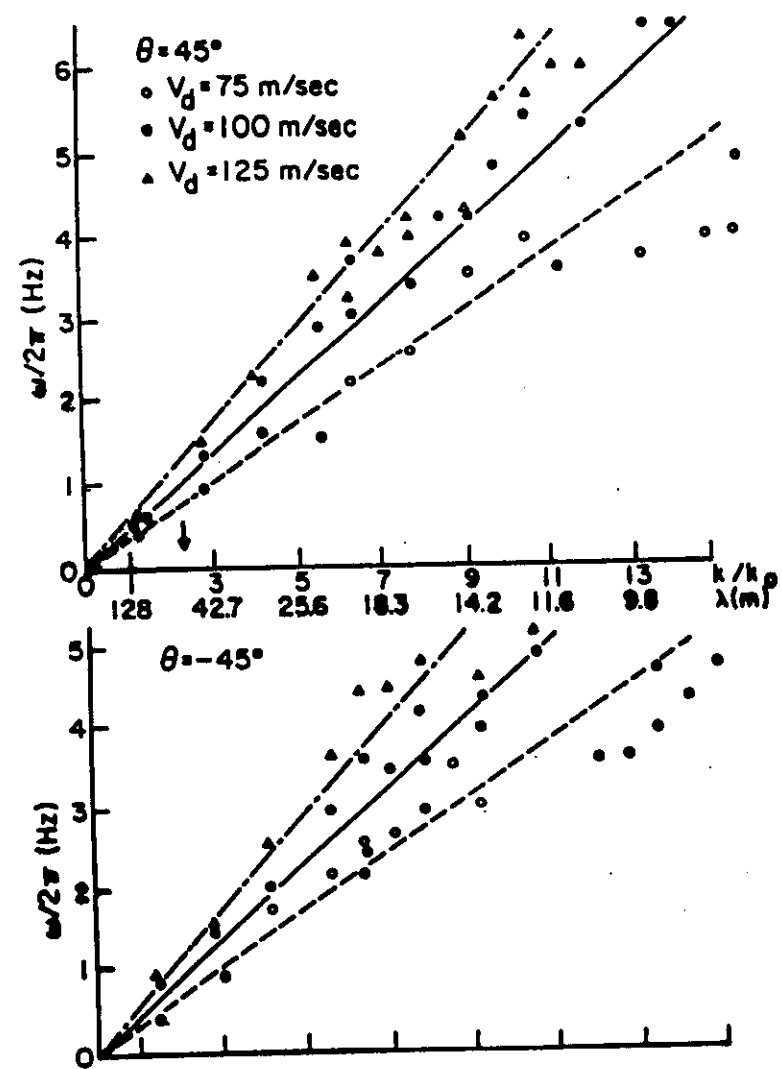
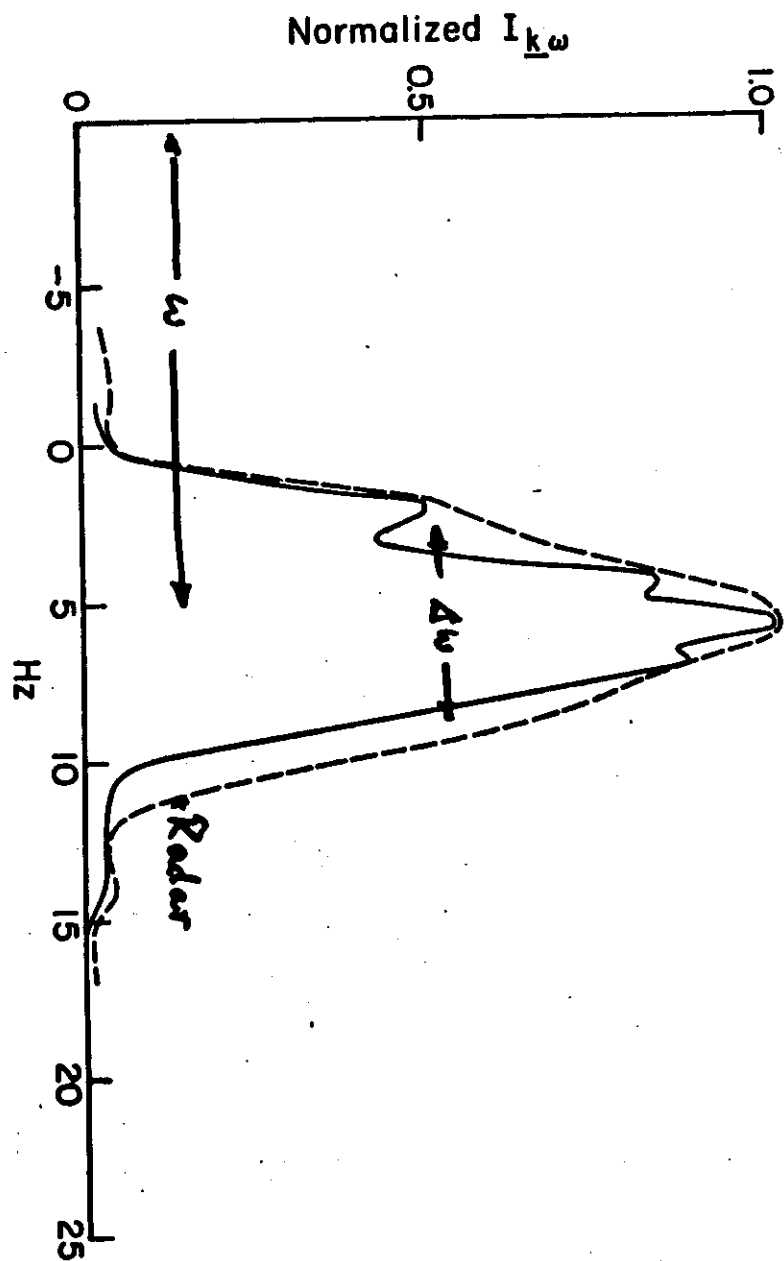


Fig. 5.7b

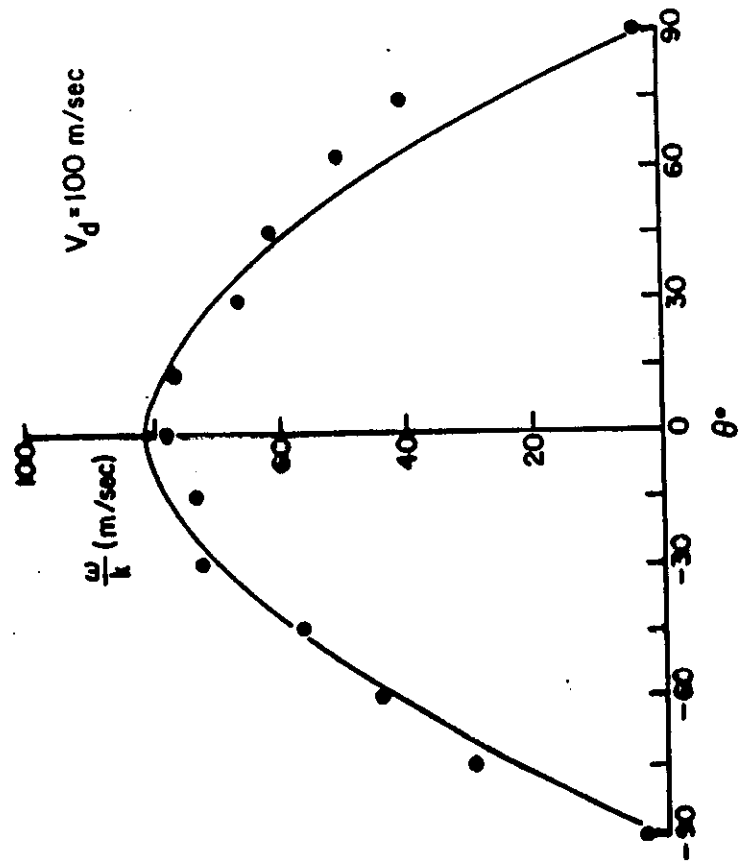


Fig. 5.8

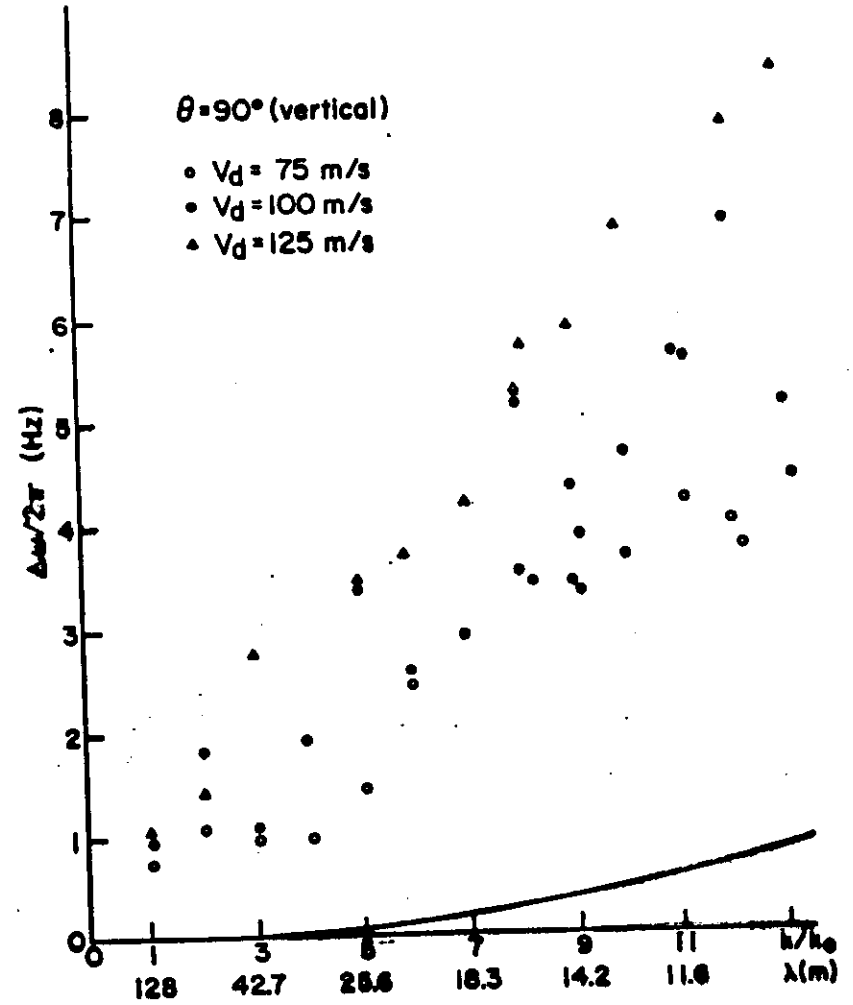


Fig. 5.9a

* Almost 2D ($k_z \ll k_\perp$)

* Low frequency ($\omega \ll \Omega_i$)

* Low beta ($\beta < m_e/m_i$)

I Drift Wave turbulence

$$\frac{\partial \varphi_{\underline{k}}}{\partial t} + i(\omega_{\underline{k}_1} + i\gamma_{\underline{k}}) \varphi_{\underline{k}} = \sum_{\underline{k}'} W_{\underline{k}_1, \underline{k}_1'} \varphi_{\underline{k}'} \varphi_{\underline{k}-\underline{k}'}$$

$$W_{\underline{k}_1, \underline{k}_1'} = \frac{1}{2} \hat{B} \cdot \underline{k}_1' \times \underline{k}_1'' (\underline{k}_1''^2 - \underline{k}_1^2) / (1 + \underline{k}_1^2)$$

$$\omega_{\underline{k}_1} = k_{y1} / (1 + \underline{k}_1^2)$$

$$\underline{k}'' = \underline{k} - \underline{k}'$$

Hasegawa-Mima PRL 39, 205 (1977)

II EXB Turbulence

$$\frac{\partial n_{\underline{k}}}{\partial t} + i(\omega_{\underline{k}_1} + i\gamma_{\underline{k}}) n_{\underline{k}} = \sum_{\underline{k}'} W_{\underline{k}_1, \underline{k}_1'} n_{\underline{k}'} n_{\underline{k}-\underline{k}'}$$

$$W_{\underline{k}_1, \underline{k}_1'} = \alpha \frac{1}{2} \hat{B} \cdot \underline{k}_1' \times \underline{k}_1'' \left(\frac{\underline{k}_1'' \cdot \underline{v}_D}{\underline{k}_1''^2} - \frac{\underline{k}_1' \cdot \underline{v}_D}{\underline{k}_1'^2} \right)$$

$$\omega_{\underline{k}_1} = \frac{\underline{k}_1 \cdot \underline{v}_D}{1 + \gamma}$$

Sudan, Kasaiwan PRL 39, 966 (1977)

$$\gamma = \frac{1}{2} \nu_D = \frac{1}{2} A^2 D$$

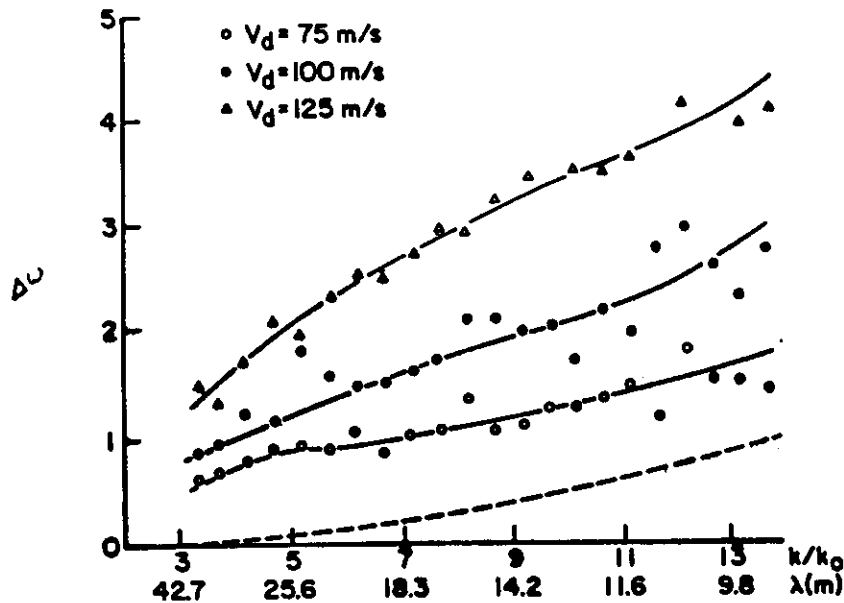
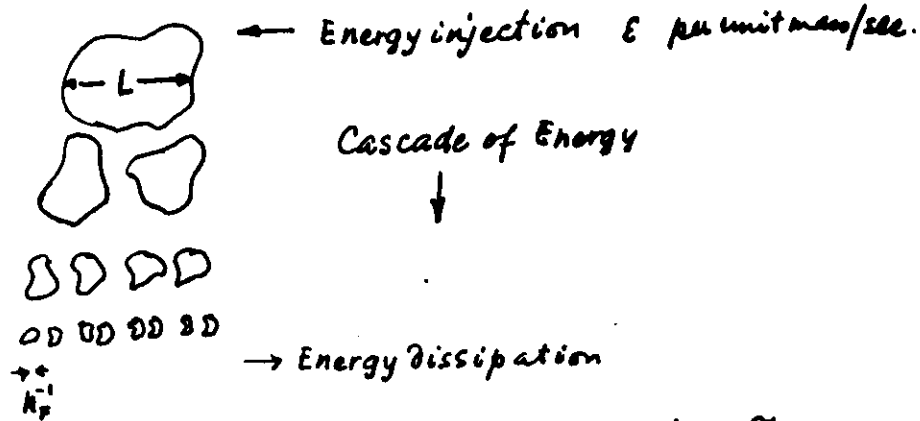


Fig. 6.3

Kolmogorov Theory (1941)



* Eddy velocity of scale size $l \sim k^{-1}$ $\tilde{v}_k$

* Eddy life time $\tau_k \sim l/v_k \sim \frac{1}{k v_k}$

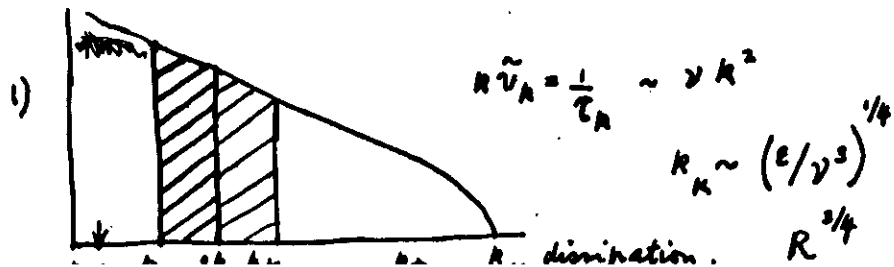
* Rate of energy transfer = $\frac{\tilde{v}_k^2}{\tau_k} = k \tilde{v}_k^3 = \epsilon$ (constant)

$$\therefore v_k \sim \epsilon^{1/3} k^{-1/3}$$

$$\langle v^3 \rangle = \int dk \epsilon(k)$$

$$\text{Reynolds no: } \frac{v_k L}{\nu}$$

$$\tilde{v}_k^2 = k \epsilon(k) \Rightarrow \epsilon(k) \sim \epsilon^{2/3} k^{-5/3}$$



$\underline{n}_k$ is a stochastic variable

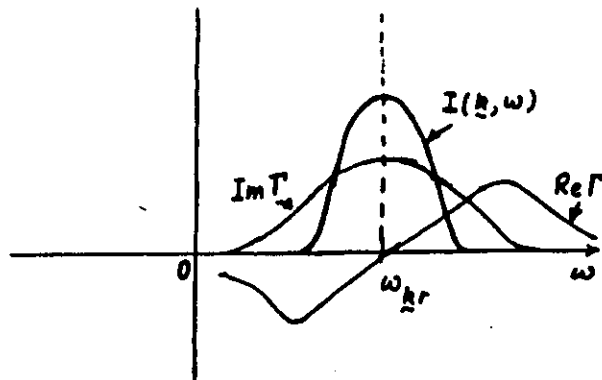
$$\langle \frac{n_{k,\omega} n_{k',\omega'}^*}{n_0} \rangle = I_{k,\omega} \delta(k-k') \delta(\omega-\omega')$$

$$\langle \frac{n(\underline{x}, t) n(\underline{x} + \underline{X}, t + T)}{n_0} \rangle = \langle n(0) n(\underline{X}, T) \rangle$$

$$I_{k,\omega} = \int d\underline{X} dT e^{i(\underline{k} \cdot \underline{X} - \omega T)} \langle n(0) n(\underline{X}, T) \rangle$$

$$|\omega - (\omega_k + i\gamma_k + i\Gamma_k)|^2 I_{k,\omega} = \frac{1}{2} \sum_{k',\omega'} |W_{k,k'}|^2 I_{k',\omega'} I_{k'',\omega''}$$

$$\Gamma_{k,\omega} = -i \sum_{k',\omega'} \frac{W_{k,k'} W_{k'',k'}^* I_{k',\omega'}}{(\omega - \omega_{k''} - i\gamma_{k''}) \cdot i\Gamma_{k'',\omega''}}$$



Let

$$I(k, \omega) = I(k) g\left(\frac{\omega - \omega_{pe}}{\tilde{\Gamma}}\right)$$

* Solve for $\Gamma(k, \omega_{pe})$; $\tilde{\Gamma} = i \int \frac{d\omega}{2\pi} \Gamma(k, \omega_{pe})$

* $\int dk \rightarrow \int \frac{dk}{k}$

Result

$$\tilde{\Gamma}(k_1) = \mathcal{O}(\alpha) \beta v_0 k_1^2 I(k_1)$$

Sudan & Keshinen
P.F. 22, 2305 (1979)

$$\beta = \frac{\gamma_0}{\Omega_i (1 + \psi)}$$

Interpretation

from linearized eqn.

electron velocity $v_k \sim \beta v_0 \left(\frac{\hat{b} \times k}{k^2} \right) \frac{n_k}{n_0}$

$$\tilde{\Gamma} \sim \tau^{-1} \sim \frac{\delta v}{\ell} \sim k \delta v \sim k [k^2 \langle |v_k|^2 \rangle]^{1/2} \sim \beta v_0 k^2 I^{1/2}$$

τ : Eddy turnover time

Kolmogorov Cascade

* Assume isotropy perpendicular to B

Energy density $\propto \int_0^\infty dk k_1 I(k_1) = \int_0^\infty \frac{dk_1}{k_1} \frac{k_1^2 I(k_1)}{E(k_1)}$

Energy Conservation

$$\frac{d}{d \ln k_1} \left[\tilde{\Gamma} k_1^2 I(k_1) \right] = 2 \underbrace{\gamma(k_1)}_{\gamma_0 - k_1^2 G} \underbrace{k_1^2 I(k_1)}_{\tilde{\Gamma} \propto k_1^2 I^{1/2}}$$

Kulsrud and Sudan

$$\tilde{\Gamma} \gg \gamma$$

for $\gamma = 0$

$$I(k_1) \propto k_1^{-8/3} ; \tilde{\Gamma}(k_1) \propto k_1^{2/3}$$

$$E(k_1) \propto k_1^{-5/3}$$

$$\hat{I}\left(\frac{k_1}{k_0}\right) = \left(\frac{k_0}{k_1}\right)^{+8/3} \left\{ 1 - \left(\frac{k_1}{k_0}\right)^{-1/3} - \frac{1}{2R} \left[\left(\frac{k_1}{k_0}\right)^{4/3} - 1 \right] \right\}^2$$

Sudan and Pfirsch
P.F. 28, 1702
(1985)

$$\frac{L u}{v}$$

$$\frac{1}{k_0} \frac{\gamma_0}{k_0 G}$$

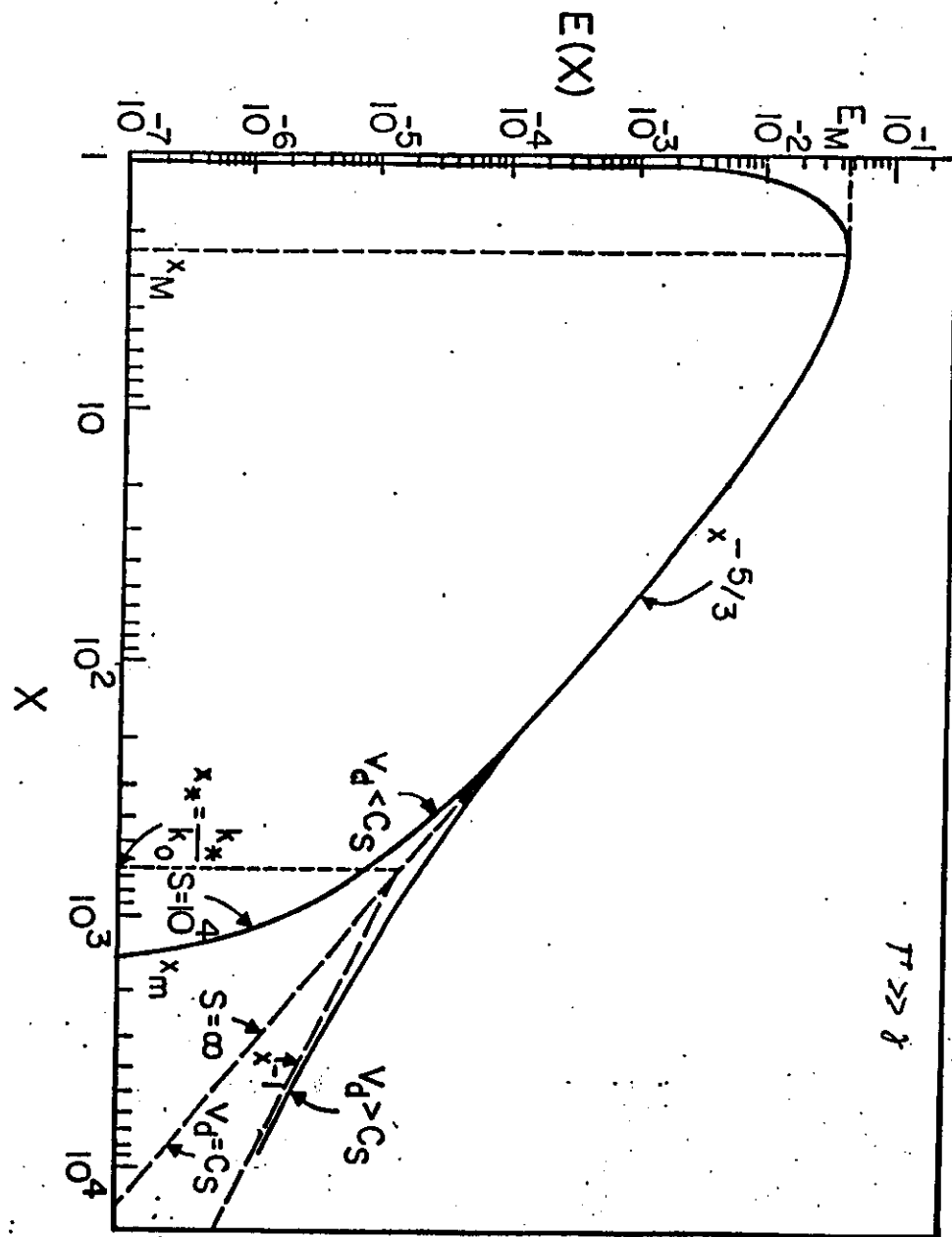
$$\hat{I} = \left(\frac{\gamma_0}{\beta v_0} \right)^2 k_0^2 I(k_1)$$

$R = \gamma_0 / k_0^2 G$ equivalent Reynolds number

Marginally stable $k = k_c = \sqrt{\gamma_0 / G}$; $\frac{k_c}{k_0} = R^{1/2}$

spectrum drops off at $k = \frac{k}{k_0} = \frac{k_K}{k_0} = (2R)^{3/4}$ Kolmogorov cut-off

$$\frac{k_K}{k_c} \sim R^{1/4} ; R \sim 10^4$$



Primary Two-Stream Waves (Upleg)
March 12, 1983 - Punta Lobos, Peru

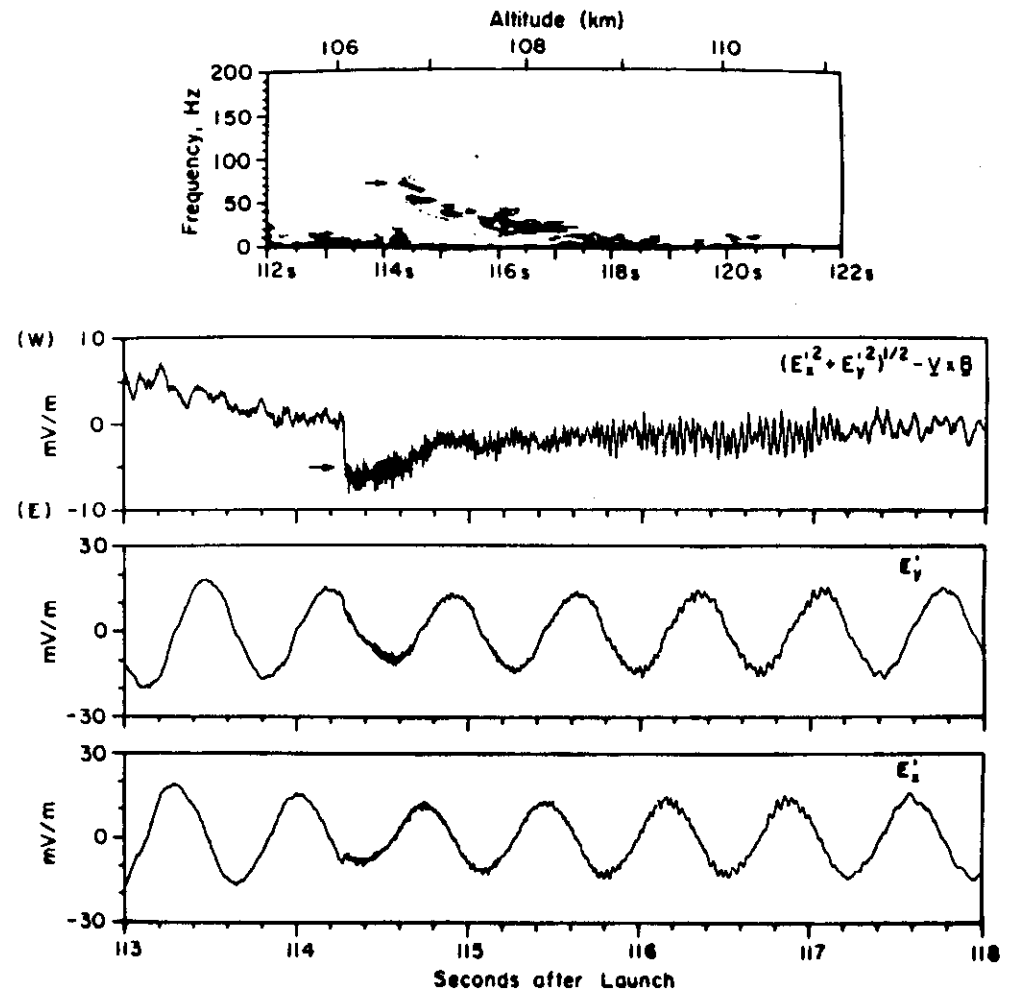


Fig. 4.29. Electric field observations of the two-stream waves for the upleg. The lower panels show the raw DC-coupled data, above which is plotted the square root of the sum of the squares of these waveforms. The upper panel shows a sonogram of these waves. (Note the change in the scale of the time axis.) An arrow indicates the onset of the strong burst of primary two-stream waves. (After Pfaff et al., 1987b. Reproduced with permission of the American Geophysical Union.)

CONDOR
MAR 14 1983

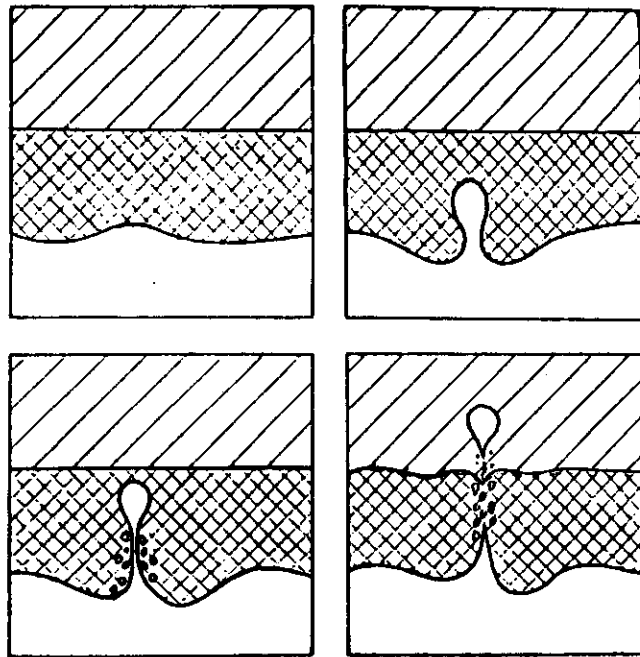
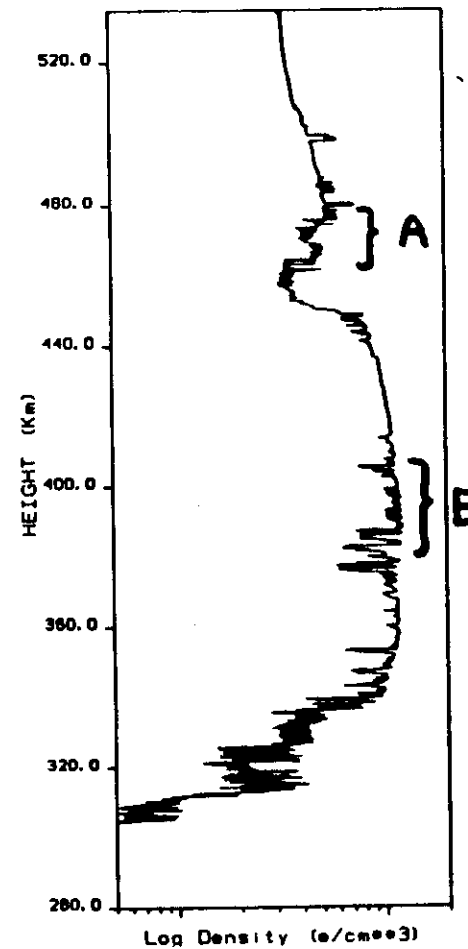
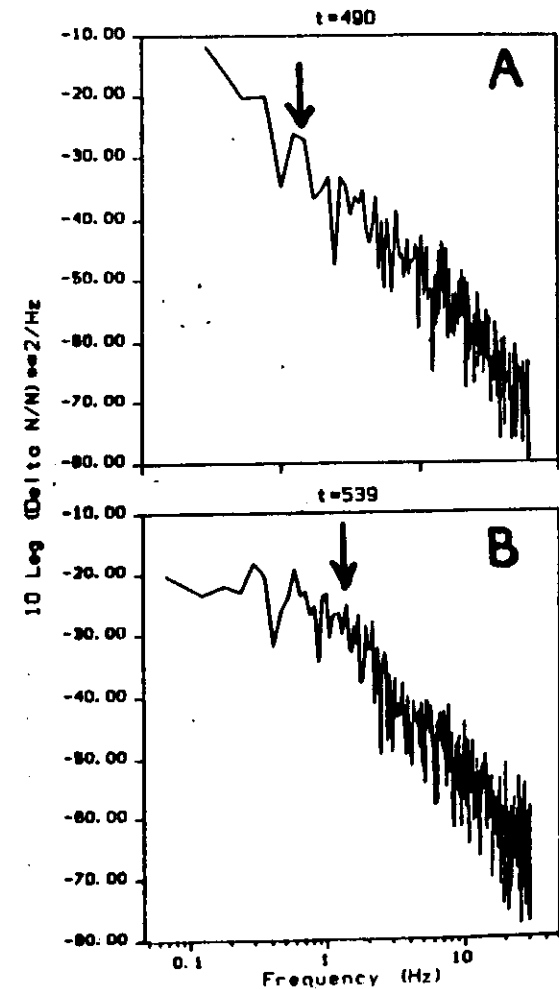


Figure 2.4: The model proposed by Woodman and LaHoz [1976] to explain topside spread F: the bottomside region, unstable to the gravitational Rayleigh-Taylor process, distorts into a pinched-off "bubble" which rises into the topside.



14 March 1983



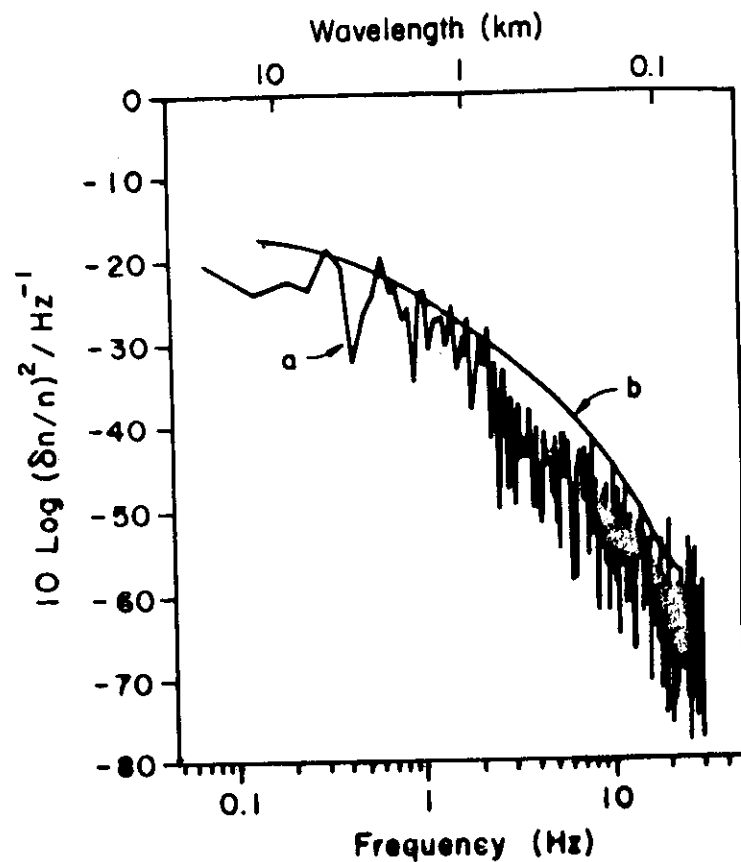


Fig. 4.15. (a) The spectrum of density irregularities measured at 390 km on the upleg of the March 14 rocket flight. (

(b) The spectrum of density irregularities predicted by Sudan and Keskinen (1984). The wave number to frequency transformation has been approximated using the rocket velocity, and the theoretically predicted spectrum has been integrated over wave number to obtain the same units as the measured spectrum. (After LaBelle and Kelley, 1986. Reproduced with permission of the American Geophysical Union.) $D_1 \approx 40 D_c$

E-Region

1. No turbulence where $\nabla n = 0$
However if $v_d > C_s / (1 + \gamma)$ observe Farley-Buneman instability
2. $\nabla n > 0$
 $v_d < \frac{C_s}{1 + \gamma}$ } Observe turbulent fluctuations
spectral index fair agreement.
3. Line width and Doppler shift from radar observations ; good agreement with simulations

F-Region

4. Above F-peak observed spectral index ~ -2.5 is greater than Kolmogorov index ; region of decaying turbulence
5. When turbulent activity is strong spectral index $\rightarrow -1.5$
6. Breaks in spectrum occur at ~ 100 m and 100 m, flatter (k^{-1}) for $\lambda < 100$ m and steeper ($k^{-2.5}$) for $\lambda > 100$ m.

Moderate Turbulence

$$\gamma_0 = T \propto k^2 I_k^{1/2}$$

$$\therefore k I_k \propto k^{-3}$$