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H4-SMR 393/6

## **SPRING COLLEGE ON PLASMA PHYSICS**

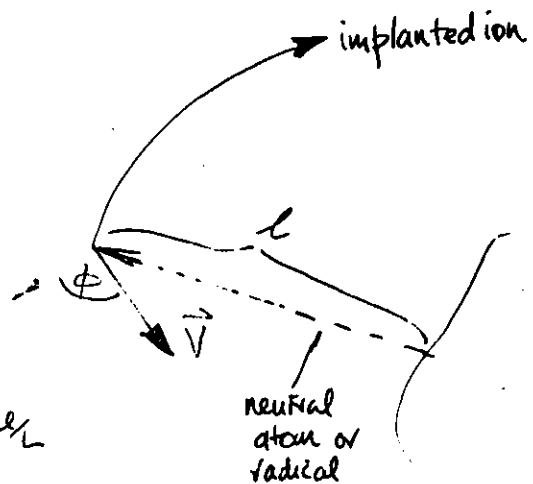
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### **ALFVENIC TURBULENCE IN BEAM-PLASMA SYSTEMS (I)**

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# THE FUNDAMENTAL COMETARY PLASMA PROCESS.



$$P(l) = e^{-l/L}$$

$$L = \frac{V_e}{\Sigma} \quad \begin{array}{l} \text{expansion velocity} \sim 1 \text{ km/s} \\ \text{ionisation rate} \sim 10^{16} \text{ s}^{-1} \end{array} \approx 10^6 \text{ km}$$

Neutral density

$$n(r) = \frac{Q}{4\pi V_e r^2} \exp\left[-\frac{r}{L}\right]$$

## FLUID EQUATIONS FOR COMETARY FLOW

See also: *Journal of Solar Wind Physics* 1, 45+, 1967

Conservation of

$$\text{number} \quad \frac{\partial n}{\partial t} + \nabla \cdot n \vec{u} = A$$

$$\text{mass} \quad \frac{\partial \rho}{\partial t} + \nabla \cdot \rho \vec{u} = B$$

$$\text{momentum} \quad \frac{\partial}{\partial t} (\rho \vec{u}) + (\vec{u} \cdot \nabla) \rho \vec{u} + \rho \vec{u} (\nabla \cdot \vec{u}) + \nabla p = C$$

$$\text{energy} \quad \frac{\partial}{\partial t} \left( \rho \frac{\vec{u}^2}{2} + \frac{p}{\gamma-1} \right) + \nabla \cdot \vec{u} \left( \rho \frac{\vec{u}^2}{2} + \frac{\gamma p}{\gamma-1} \right) = D$$

Source Terms

- A local ionisation rate
- B ionisation rate x mass difference
- C loss of momentum by charge exchange
- D loss of energy by charge exchange

Assumption

Instantaneous assimilation into solar wind flow of cometary ions.

## COMETARY ION SOURCE

### Assumptions

- (a) constant gas production rate  $\mathcal{Q}$  molecules/s  $\sim 10^{30}$
- (b) radially symmetric outflow at constant velocity  $v_e \sim 1 \text{ km/s}$
- (c) constant ionisation rate  $\nu \sim 10^{-6} \text{ s}^{-1}$

Number density  $n_e$

$$n_e = \frac{\mathcal{Q}}{4\pi v_e r^2} \exp\left(-\frac{\nu r}{v_e}\right)$$

Ionisation rate

$$\frac{\partial n_i}{\partial t} = \frac{\mathcal{Q} \nu}{4\pi v_e r^2} \exp\left(-\frac{\nu r}{v_e}\right)$$

Ion flux

$$n_i v_i = \int_{x_0}^{\infty} \frac{\mathcal{Q} \nu}{4\pi v_e r^2} \exp\left(-\frac{\nu r}{v_e}\right) dS$$

Define scale length  $L = v_e/\nu \sim 10^6 \text{ km}$

$$n_i v_i = \frac{\mathcal{Q}}{4\pi L^2} \int_{x_0}^{\infty} \frac{1}{(x^2 + y_0^2)^{3/2}} \exp\left[-\frac{(x^2 + y_0^2)^{1/2}}{L}\right] dx$$

## Ion Pickup Trajectories

Coordinate system

$\vec{v}_s$  in x-z plane at angle  $\alpha$  to  $\vec{B}$  along  $\hat{z}$

Pickup trajectories

$$x = v_{\perp} \left[ t - \frac{1}{\Omega_i} \sin(\Omega_i t + \epsilon) \right] + x'$$

$$y = -\frac{v_{\perp}}{\Omega_i} \cos(\Omega_i t + \epsilon) + y'$$

$$z = v_{\parallel} t + z' \approx z'$$

$$\Omega_i = eB/m_i$$

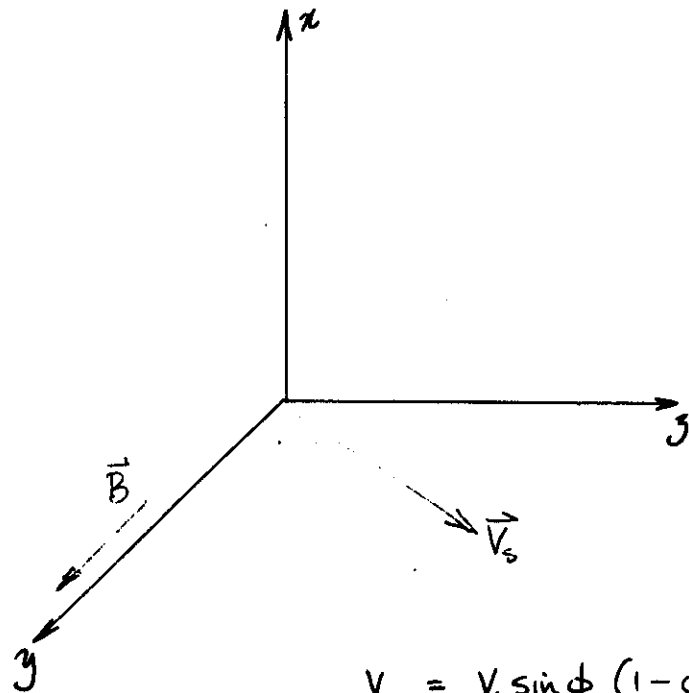
$$v_{\perp} = v_s \sin \alpha$$

$$v_{\parallel} \approx v_e \approx 0$$

Motion is (gyration at  $v_{\perp}$ ) + (drift at  $v_{\perp}$ )

all perpendicular to  $\vec{B}$

# Ion pickup perpendicular to $\vec{B}$



$$\begin{aligned} V_z &= V_s \sin \phi (1 - \cos \omega_c t) \\ V_y &= 0 \\ V_x &= V_s \sin \phi \sin \omega_c t \end{aligned}$$

Maximum Velocity  $2V_s \sin \phi$

Maximum Energy  $4E \sin^2 \phi$

where  $E = \frac{1}{2} m V_s^2$

Max. momentum  $m V_s \sin \phi$  along  $z$

Max. energy  $\frac{1}{2} m V_s^2 \sin^2 \phi$

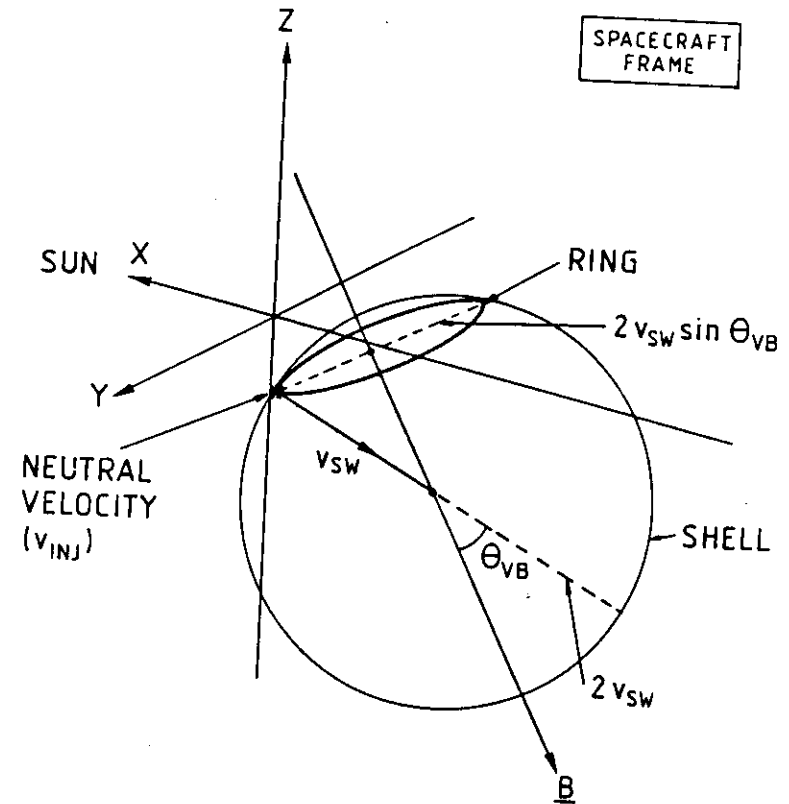


Figure 1 - Ion pickup velocity space geometry in the spacecraft frame, for an arbitrary orientation of the magnetic field  $\vec{B}$ . Neutral particles are injected at point indicated by the star ( $\underline{v}_{inj} \approx (0, 0.15, -68.4) \text{ km/s}$  in this frame) and their initial trajectory is in a ring around  $\vec{B}$ . Later they scatter onto a shell centred on the solar wind speed to a first approximation.

## ION VELOCITY DISTRIBUTION

Distribution a ring of radius  $v_s \sin \alpha$  in velocity space centred on the solar wind speed

maximum energy in ring  $4 \frac{m_i}{m_p} E_p \sin^2 \alpha$

$E_p$  energy of solar wind proton

If scattered into shell

maximum energy  $4 \frac{m_i}{m_p} E_p$

For instant assimilation we require

either  $\alpha = 90^\circ$

or rapid scattering over the shell

## One dimensional equations

Wallis + Ong Plan. Sp. Sci 23, 713, 1975  
Galeev et al ApJ 289, 807, 1985

Coordinate system

$x$  parallel to flow  
 $z$  parallel to  $\vec{B}$

Conservation equations

$$\frac{d}{dx} [\rho_i u f(\mu, u)] = \frac{m_i Q \nu}{4\pi v_e r^2} \exp\left(-\frac{\nu r}{v_e}\right) \delta\left(\mu - \frac{m_i u^2}{2B}\right)$$

$$\frac{d}{dx} (\rho u) = \frac{m_i Q \nu}{4\pi v_e r^2} \exp\left(-\frac{\nu r}{v_e}\right)$$

$$\frac{d}{dx} \left( \rho u^2 + p_\perp + \frac{B^2}{8\pi} \right) = 0$$

Solution

$$\frac{u_\infty}{U} = 2 - \left[ 4 - 3 \left( \frac{\rho u}{\rho_\infty u_\infty} \right) \right]^{1/2}$$

$\infty$  refers to upstream, undisturbed values

reversal when  $\frac{\rho u}{\rho_\infty u_\infty} = \frac{4}{3}$

Shock must form.

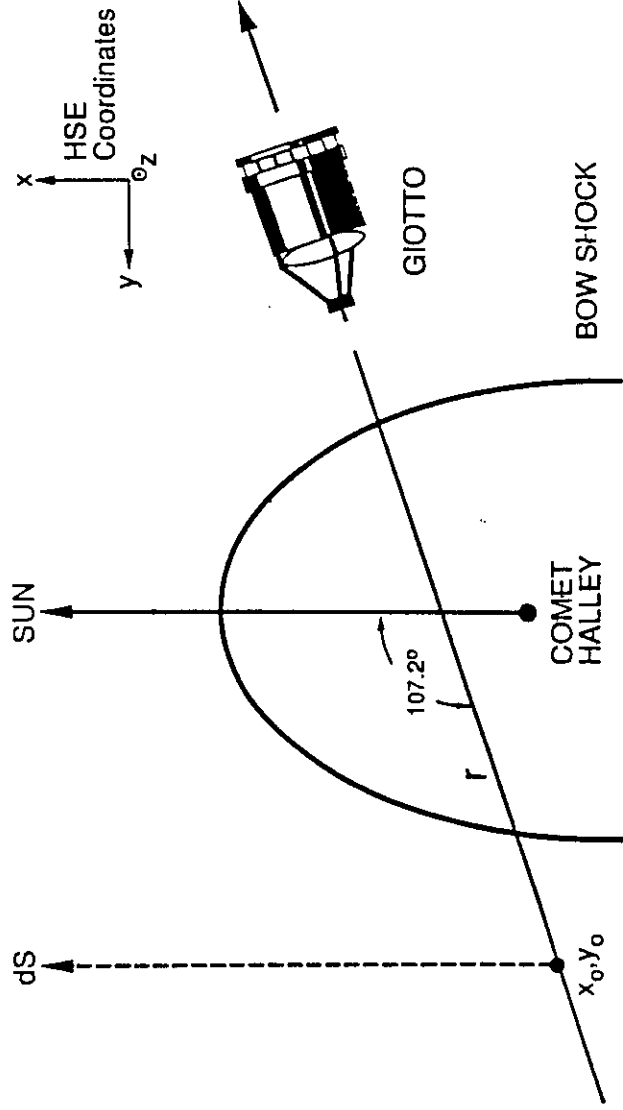


Figure 1 - Schematic representation of the Giotto fly-by geometry in Halley-centred Solar Ecliptic (HSE) coordinates, showing the integration path,  $dS$ , for estimation of the implanted cometary ion flux at a spacecraft position  $(x_0, y_0)$ .

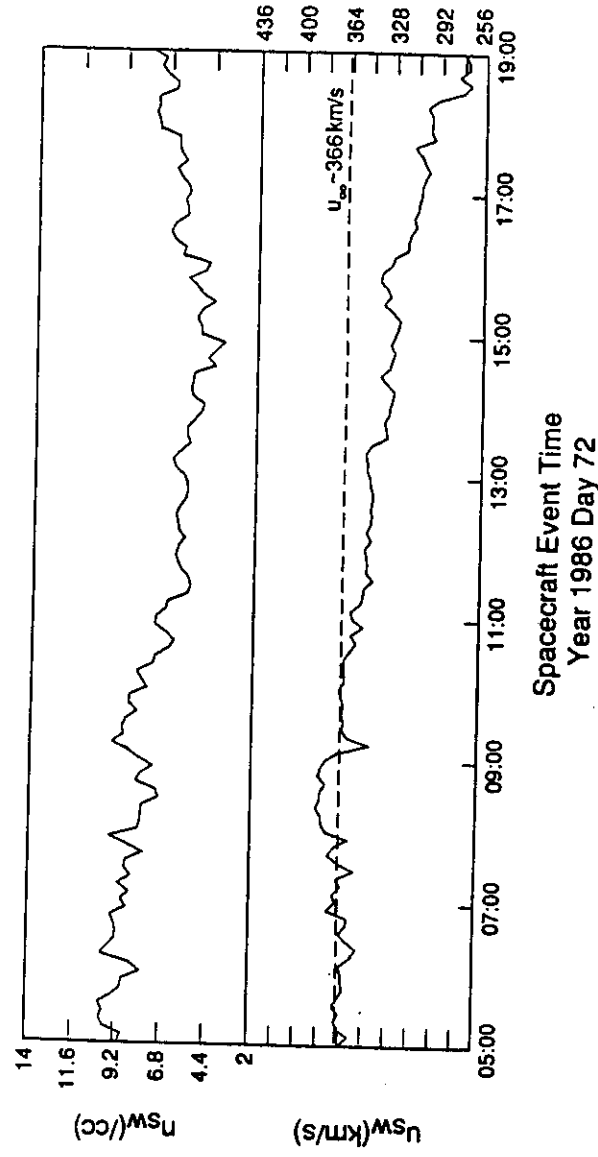


Figure 3 - Smoothed solar wind proton density and velocity profiles for the same period as in Figure 2. An upstream velocity of  $u_\infty \sim 366 km/s$  is overlaid.

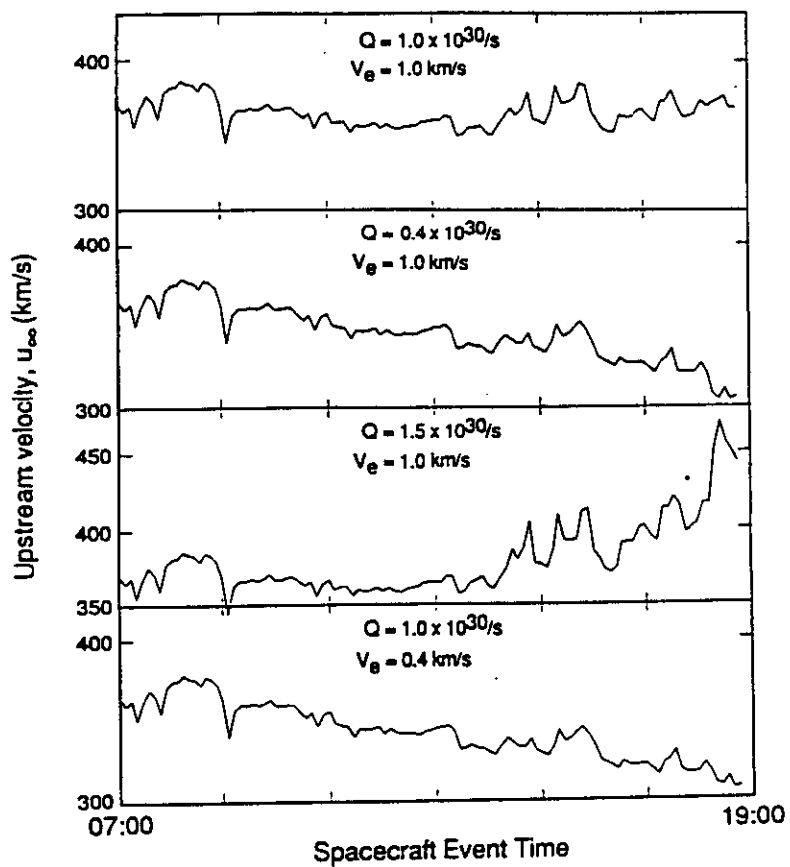


Figure 5 - Effects of differing gas emission values on the  $u_{\infty}$  profiles inferred.

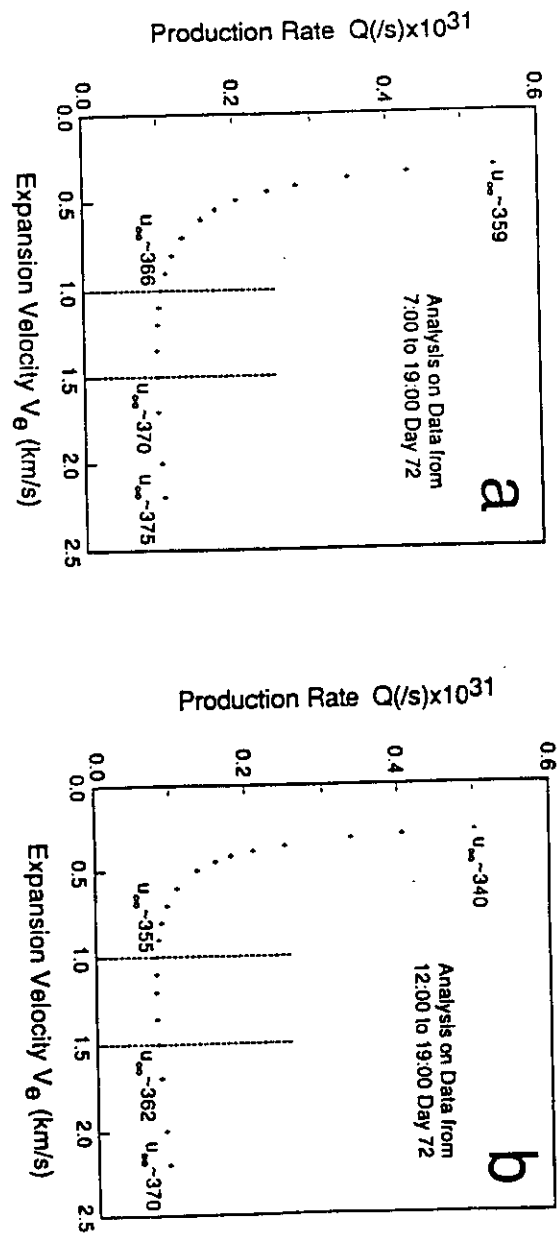
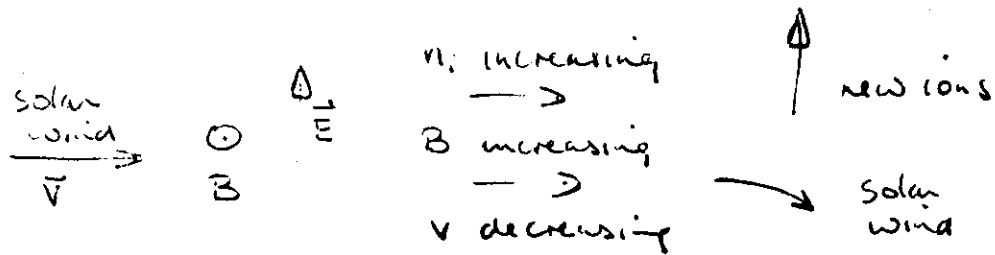


Figure 7 - Set of possible solutions resulting from analysis on the data periods (a) 07:00-19:00 hrs, and (b) 12:00-19:00 hrs.

## PERPENDICULAR PICKUP MECHANISM



- ⊗ Momentum balance perpendicular to flow between new ions and solar wind
- ⊗ Solar wind decelerated by electric field
- ⊗  $\vec{E} \neq -\vec{V}_s \times \vec{B}$

Is the ring distribution stable?

Wu + Davidson JGR 77, 5399, 1972  
Lee + Ip JGR 92, 11041, 1987

Dispersion relation for electromagnetic waves propagating parallel to the magnetic field

$$k^2 c^2 = \omega^2 + 4\pi^2 \omega \sum_j \frac{q_j^2}{m_j} \int dv_{\parallel} dv_{\perp} v_{\perp}^2 \frac{G F_j}{\omega - kv_{\parallel} + \Omega_j}$$

$$G \equiv \frac{\partial}{\partial v_{\perp}} + \frac{k}{\omega} \left( v_{\perp} \frac{\partial}{\partial v_{\parallel}} - v_{\parallel} \frac{\partial}{\partial v_{\perp}} \right)$$

- ⊕ work in the solar wind frame
- ⊕ cold sw electrons
- ⊕ cold sw protons
- ⊕ ring beam of cometary ions
- ⊗ assume  $\omega - kv_{\parallel} + \Omega_j$  more sharply peaked than  $F_j$

Then  $\omega_p^2 = k^2 V_A^2$

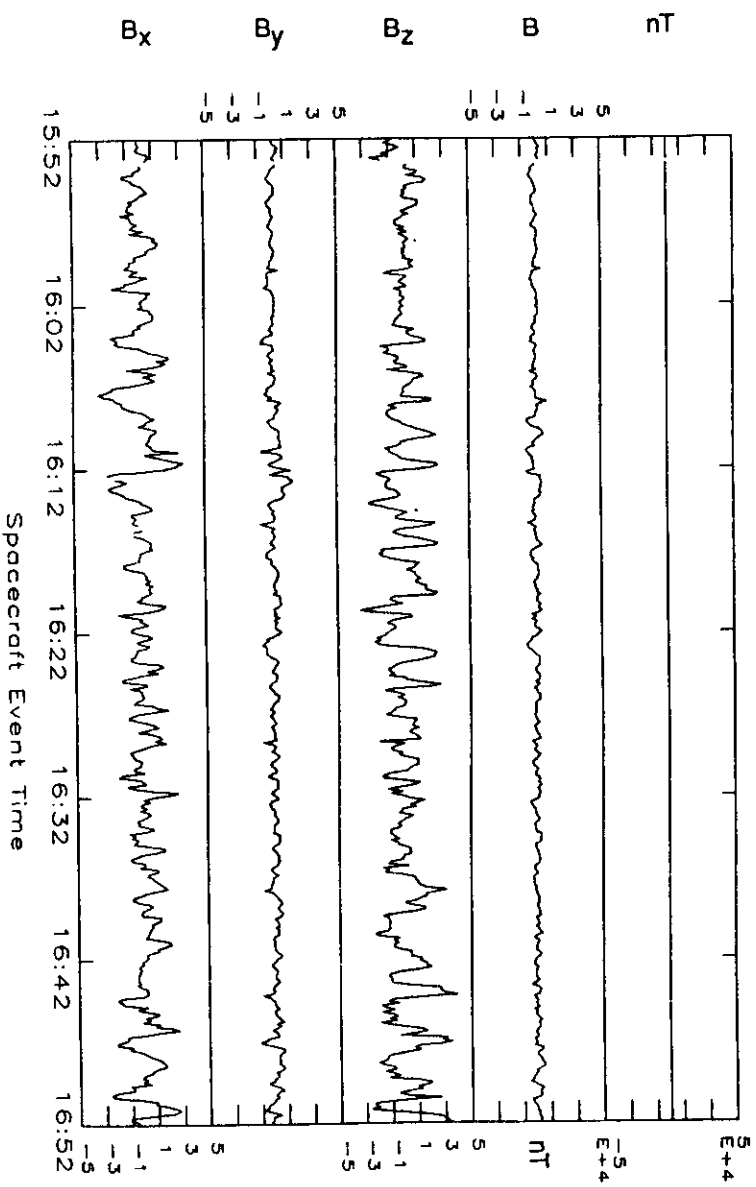
$$\gamma = \frac{\pi^2}{2} \frac{V_A^2}{c^2} \frac{\omega_p^2}{n_i} \int dv_{\parallel} dv_{\perp} v_{\perp}^2 \delta(\omega - kv_{\parallel} + \Omega_i) G$$

Put  $F_i = \frac{n_i}{2\pi^{3/2} v_{\perp 0}} \delta(v_{\perp} - v_{\perp 0}) \exp\left(-\frac{(v_{\parallel} - v_{\parallel 0})^2}{\sigma^2}\right)$

Obtain maximum  $\gamma = 1/(2\pi)^{1/2} V_A \omega_p^2 v_{\perp 0}^2 v_{\parallel 0}$

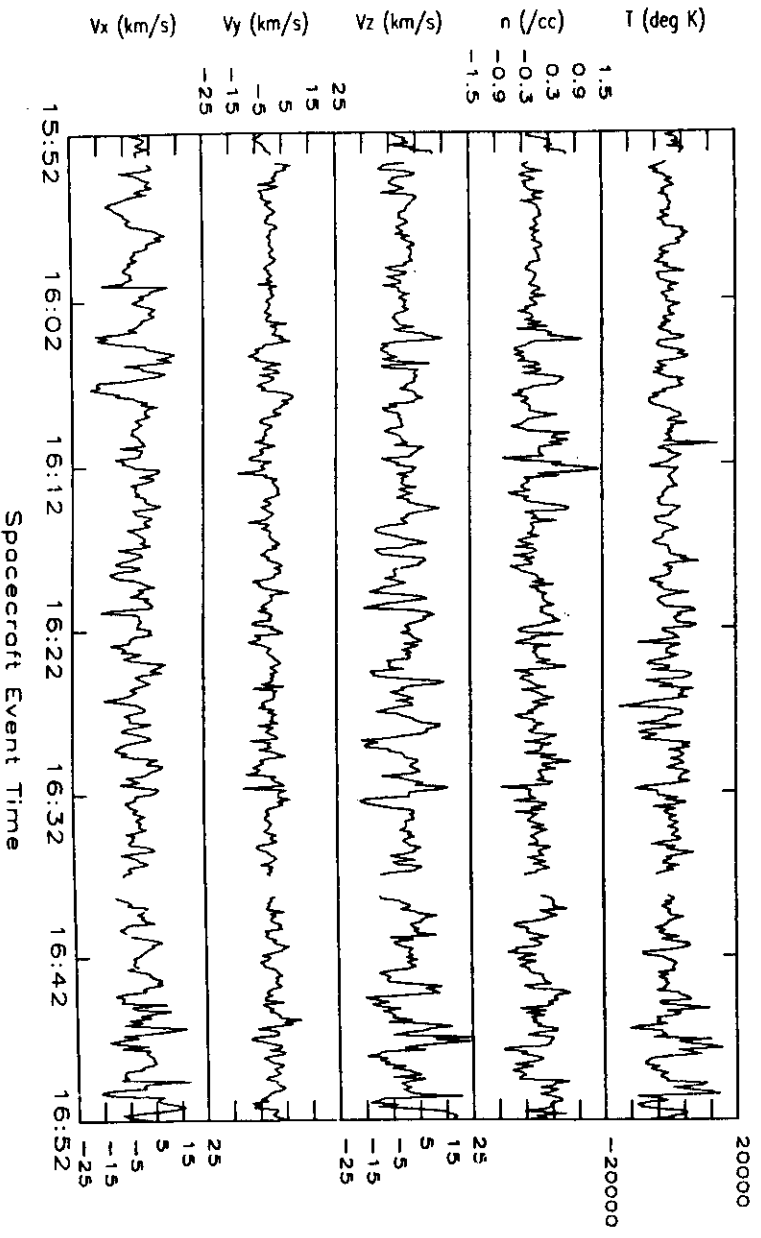
# Giotto MAG Magnetic Field Fluctuations

Year 1986 Day 72 FLD coordinates



## Fluctuations

GIOTTO JPA FIS Solar Wind Proton Parameters  
Year 1986 Day 72 FLD coordinates



WAVES 3mHz < f < 50mHz  
 GIOTTO JPA FIS Solar Wind Proton Parameters  
 Year 1986 Day 72 FLD coordinates

