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H4-SMR 393/7

SPRING COLLEGE ON PLASMA PHYSICS

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ALFVENIC TURBULENCE IN BEAM-PLASMA SYSTEMS (II)

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DIFFUSION PATH IN VELOCITY SPACE

$$v_{\perp}^2 + (v_{\parallel} \pm v_A)^2 = \text{constant}$$

- ⊗ particle energy conserved in the wave frame
- ⊗ energy and momentum conserved in wave-particle interaction

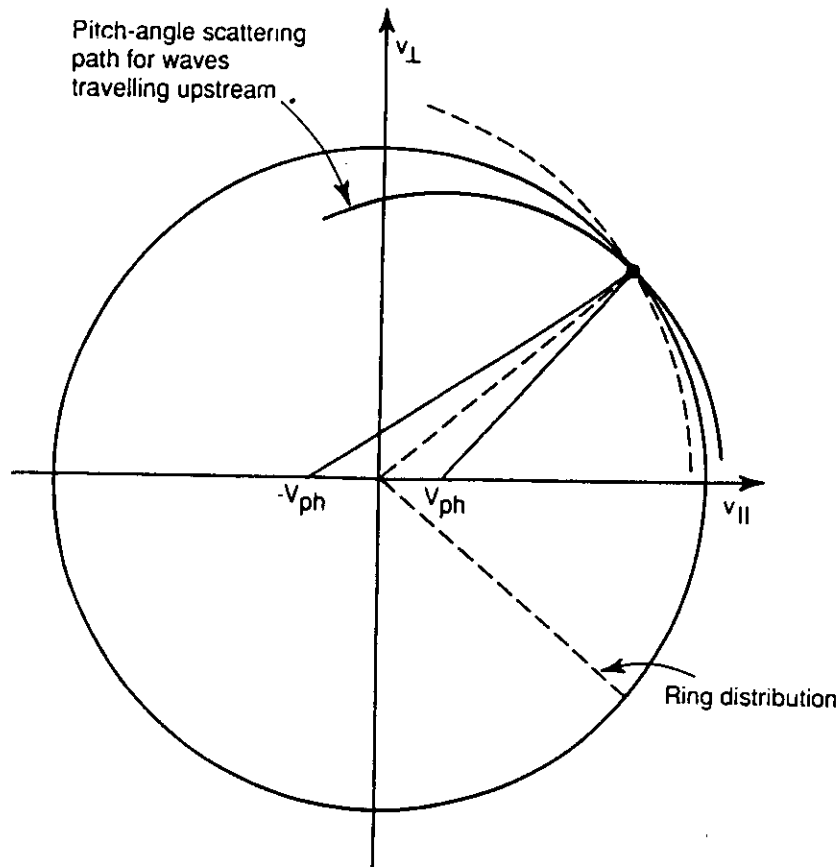
How does energy and momentum change in diffusion process?

Ring	bulk kinetic energy	$\frac{1}{2} \rho_i v^2 \cos^2 \alpha$
	thermal energy	$\frac{1}{2} \rho_i v^2 \sin^2 \alpha$
Shell	bulk kinetic energy	$\frac{1}{2} \rho_i v_A^2$
	thermal energy	$\frac{1}{2} \rho_i (v^2 \sin^2 \alpha + (v \cos \alpha - v_A)^2)$

$$E_{\text{ring}} - E_{\text{shell}} = \frac{1}{2} \rho_i [v_A v \cos \alpha - v_A^2] \Rightarrow \text{"free energy"}$$

$$P_{\text{ring}} - P_{\text{shell}} = \rho_i (v \cos \alpha - v_A) = \frac{\Delta E}{v_A}$$

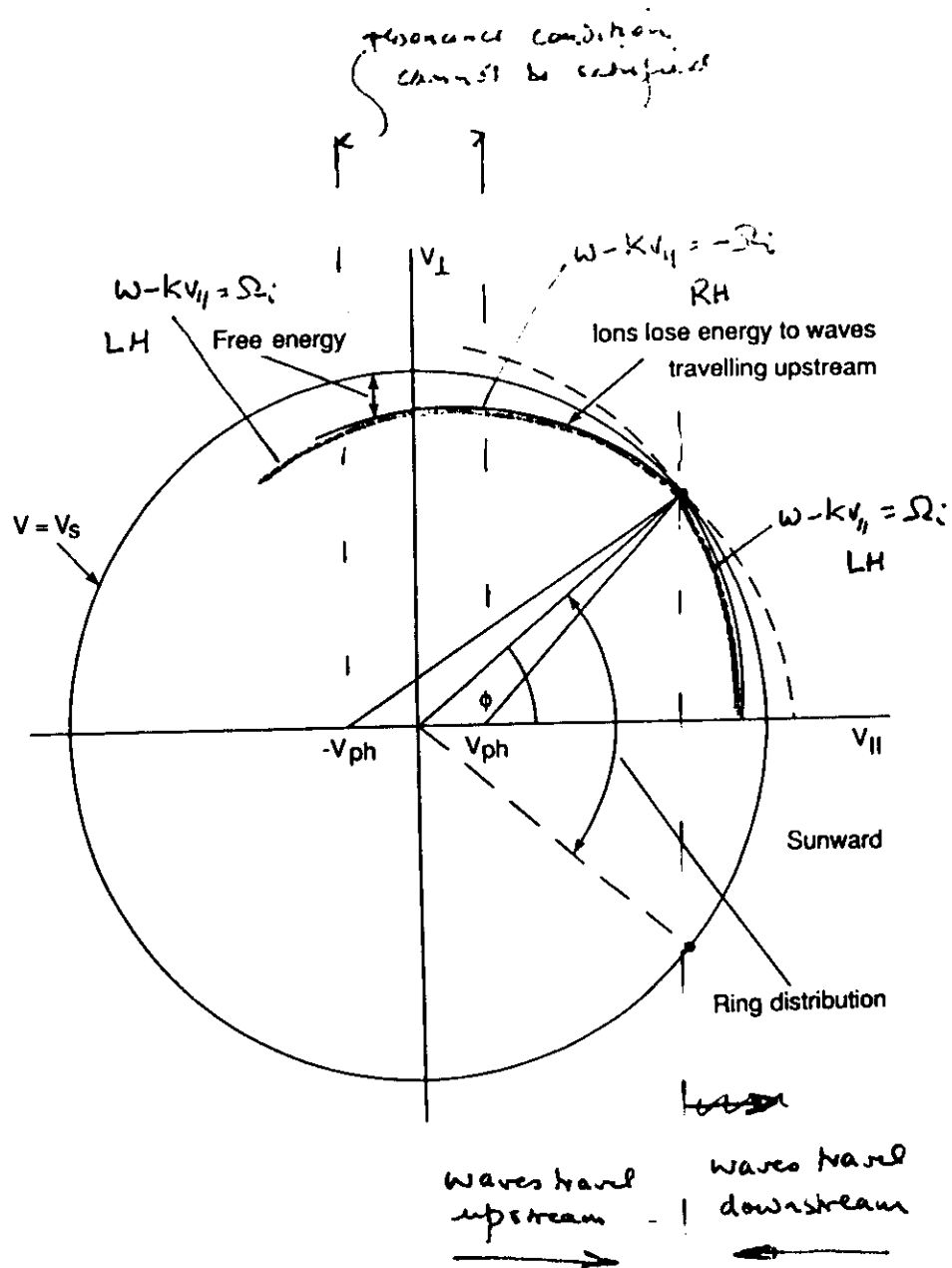
Acceleration requires waves travelling in both directions



$$v_{\perp}^2 + (v_{\parallel} \pm v_{ph})^2 = \text{const}$$

e.g. Alfvénic turbulence $v_{ph} = v_A$

Figure 1 - Schematic diagram showing the pitch-angle scattering path of injected ions from their initial ring distribution to a shell centred on the wave phase velocity.



POYNTING VECTOR $\vec{S}$

$$\vec{S} = \vec{E}_1 \times \vec{H}_1 = \vec{E}_1 \times \vec{B}_1 / \mu_0$$

$$\vec{B} = \vec{B}_0 + \vec{B}_1$$

B_0 — IMF

B_1 — wave field

$$\vec{V} = \vec{V}_0 + \vec{V}_1$$

V_0 — solar wind velocity

V_1 — wave velocity.

$$\vec{E}_1 = -\vec{V}_1 \times \vec{B}_0$$

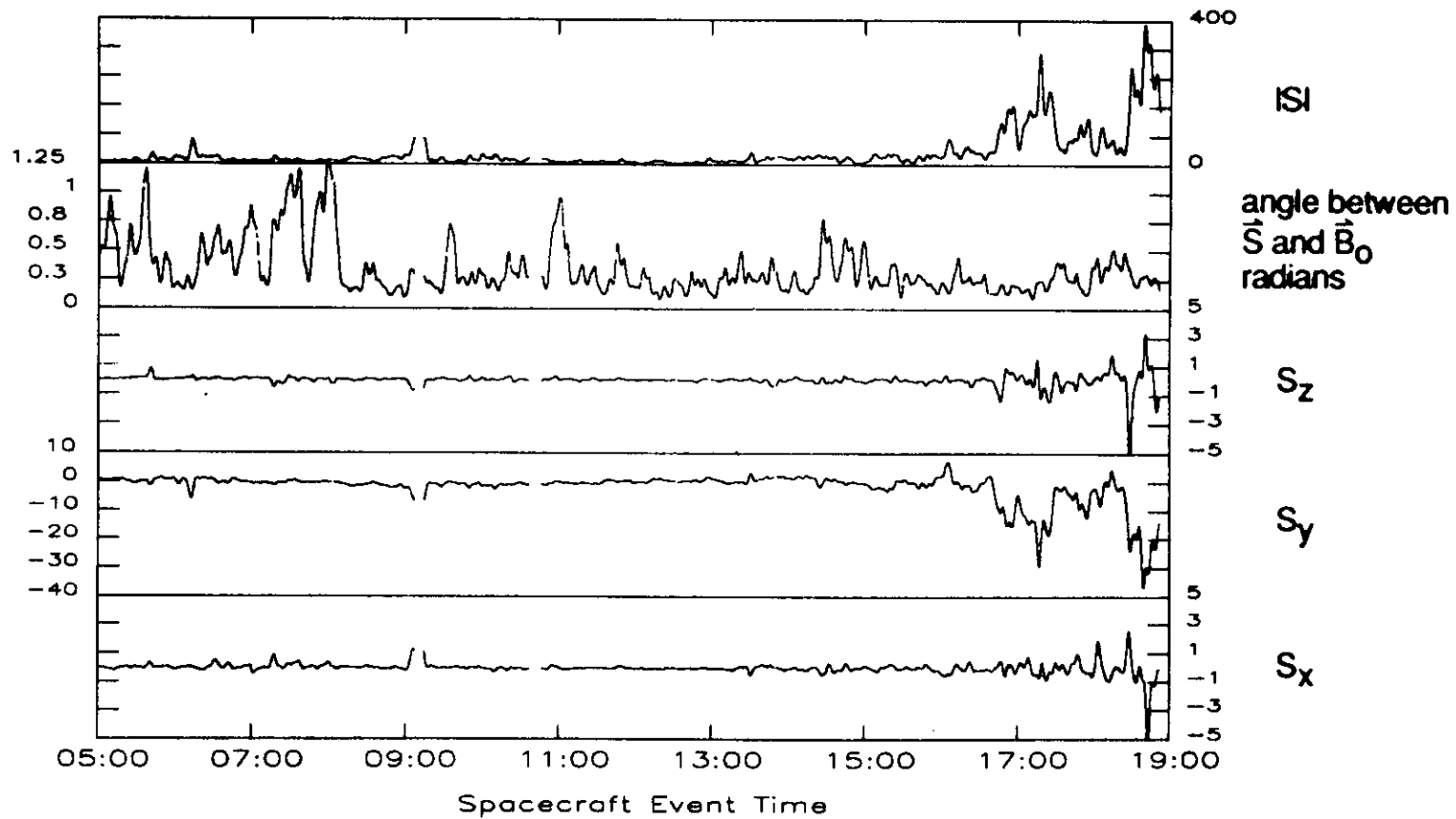
$$\vec{S} = (\text{wave energy density}) \times (\text{propagation velocity}) \times \vec{k}$$

$$\text{wave energy density} \quad \frac{1}{2} \rho \langle v_1^2 \rangle + \frac{1}{2} \langle B_1^2 / \mu_0 \rangle$$

free energy for ion beam instability

$$\frac{1}{2} \rho_i v_0^2 \cos^2 \phi$$

POYNTING VECTOR $\vec{S}$ ($\mu\text{J}/\text{m}^3 \times \text{km}/\text{s}$)
 Year 1986 Day 72 FLD coordinates



TRANSPORT EQUATIONS

For particles, the Boltzmann equation

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \nabla f + \frac{\vec{F}}{m} \cdot \nabla_v f = \left[\frac{\partial f}{\partial t} \right]_{\text{collisions}} + \text{Source} - \text{Loss.}$$

For waves, we require a suitable kinetic equation which relates to the particle distribution

QUASILINEAR THEORY

Lee + Ip JGR 92, 11041, 1987

Grueter et al JGR 93, 5470, 1988

$$\frac{\partial \bar{F}_i}{\partial t} = \frac{q_i^2}{2m_i^2 v_\perp} R_i \cdot \int d\mathbf{k} \sum_{p \neq i} \frac{(\omega_p)^2}{c^2 k^2} Q^p \left[\frac{v_\perp I_p(k) Q^p F_i}{(\omega_p - \mathbf{k} \cdot \mathbf{v}_i + \Omega_i)} \right]$$

$$\frac{\partial I_p}{\partial t} = 2 \gamma_p I_p$$

If waves are propagating in one direction can transform to wave frame

- ⊕ diffusion path is circle centred on origin
- ⊗ Q is simplified
- ⊗ pure pitch angle diffusion

$$\frac{\partial \bar{F}_i}{\partial t} = \frac{\partial}{\partial \mu} \left[\frac{\pi q_i^2 (1 - \mu^2)}{2m_i^2 c^2 v |\mu|} I_+ \left(\frac{\Omega_i}{v \mu} \right) \frac{\partial \bar{F}_i}{\partial \mu} \right]$$

$$\frac{\partial I_+}{\partial t} = I_+ + \frac{4\pi^3 v_A q_i^2}{|k| c^2 m} \int dv d\mu v^2 (1 - \mu^2) \delta\left(\mu - \frac{\Omega_i}{kv}\right) \frac{\partial \bar{F}_i}{\partial \mu}$$

$$\mu = \frac{v_{||} - v_A}{v}$$

WAVE ENERGY CONSIDERATIONS

Integrating over k , the total energy in the magnetic waves is

$$\frac{\langle \vec{B} \cdot \vec{B} \rangle}{8\pi} = \frac{1}{2} n_i m_i v_A v_0 \mu_0$$

Total wave energy is twice this because equal amount in kinetic energy of the particles

Also $v_0 \mu_0 = v \cos \alpha - v_A$ in solar wind frame

Hence
$$E_{\text{waves}} = n_i m_i v_A (v \cos \alpha - v_A)$$

is
$$E_{\text{waves}} = E_{\text{ring}} - E_{\text{shell}}.$$

TIME ASYMPTOTIC DISTRIBUTIONS

See + I_p JGR 92, 11041, 1987.

our
$$F_i = \delta(v - v_0) f(\mu, t)$$

Integrate both equations over time

$$f(\mu, \infty) - f(\mu, 0) = \frac{\partial}{\partial \mu} \left[A(\mu) \int_0^\infty dt I_+ \left(\frac{\Omega_i}{kv_\mu} \right) \frac{\partial f}{\partial \mu} \right]$$

$$I_+(k, \infty) - I_+(k, 0) = B(|k|) \int_{-1}^1 d\mu (1-\mu^2) \delta\left(\mu - \frac{\Omega_i}{kv_0}\right) \int_0^\infty dt I_+ \frac{\partial f}{\partial \mu}$$

$$A(\mu) = \frac{\pi}{2} \left(\frac{q}{mc} \right)^4 \frac{(1-\mu^2)}{v_0 |\mu|}$$

$$B(|k|) = \frac{4\pi^3 v_A q^3 v_0^2}{|k| mc^2}$$

$$f(\mu, 0) = \frac{n_i}{2\pi v_0^2} \delta(\mu - \mu_0) \quad f(\mu, \infty) = \frac{n_i}{4\pi v_0^2}$$

Then

$$I_+(k, \infty) - I_+(k, 0) = \frac{2\pi n_i m_i v_A |\Omega_i|}{k^2} \left[\frac{\Omega_i}{kv_0} - \frac{(\frac{\Omega_i}{kv_0} - \mu_0)}{|\frac{\Omega_i}{kv_0} - \mu_0|} \right]$$

INTERPRETATION OF WAVE FREQUENCIES

① The resonance condition is

$$\omega - kV_{||0} \pm \Omega_i = 0$$

where $V_{||0}$ is the ring beam velocity parallel to the solar wind.

② Waves are propagating parallel to B in the moving plasma and are Doppler shifted in the spacecraft frame

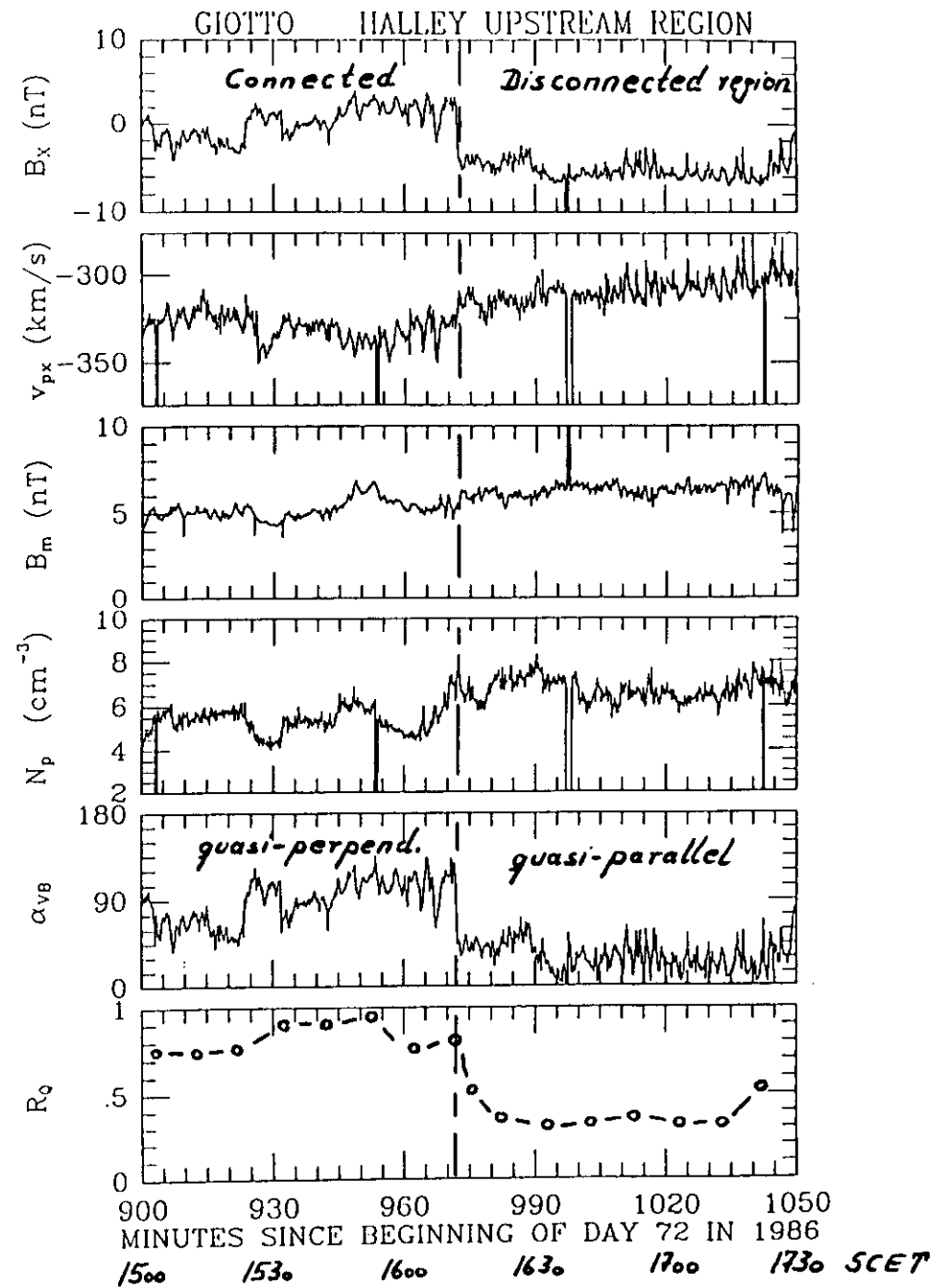
③ Spacecraft velocity is also $V_{||0}$ in solar wind frame

Waves are $\therefore$ seen at $\omega - kV_{||0} = \Omega_i$

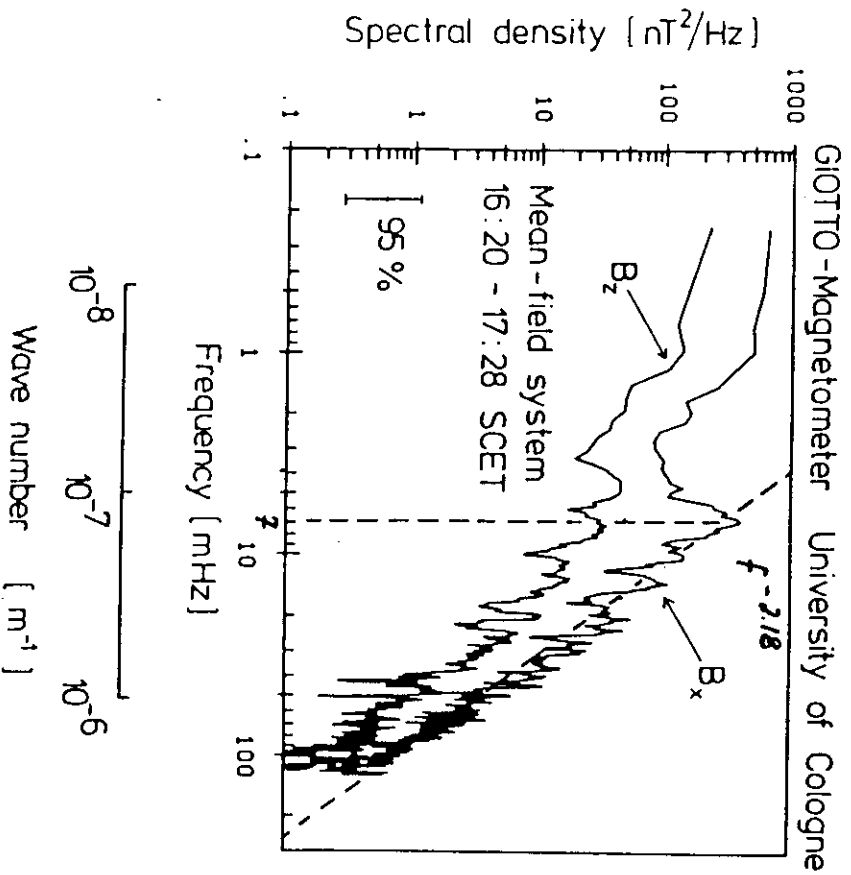
④ Relation between ω_{sc} (s/c frame) and k is

$$\omega = (V_{||0} + V_A)k$$

for waves propagating upstream



Disconnected, quasi-parallel region

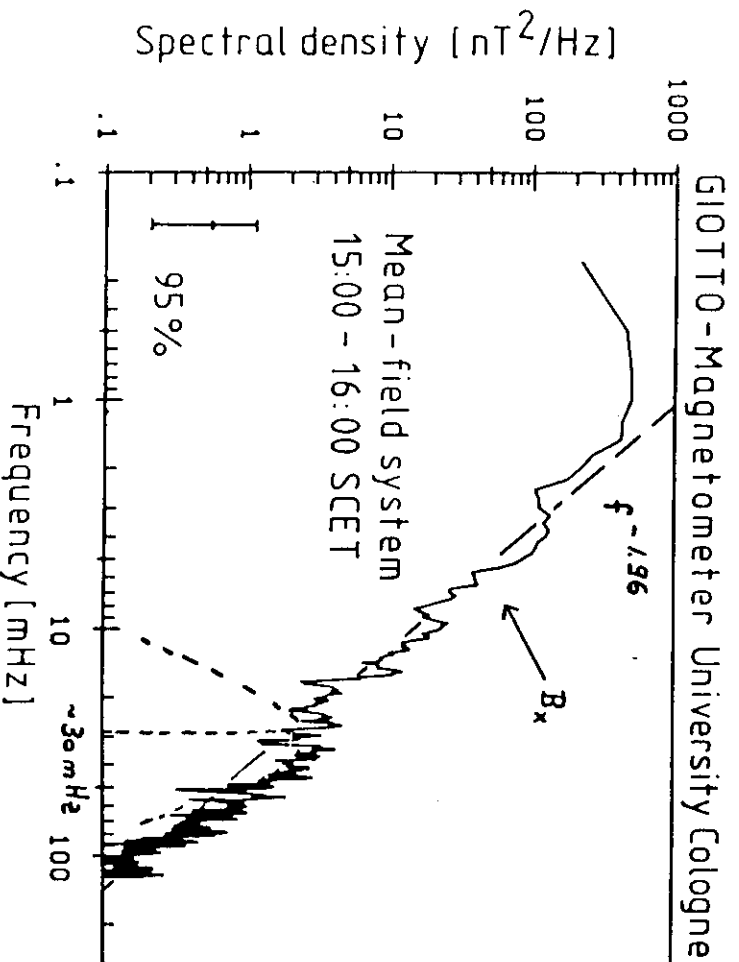


$$\langle \vec{B} \rangle =$$

$$(\sim 5.37, 0.73, 1.39)$$

$$|\vec{B}| = 5.6 \text{ nT}$$

Connected, quasi-perpendicular region



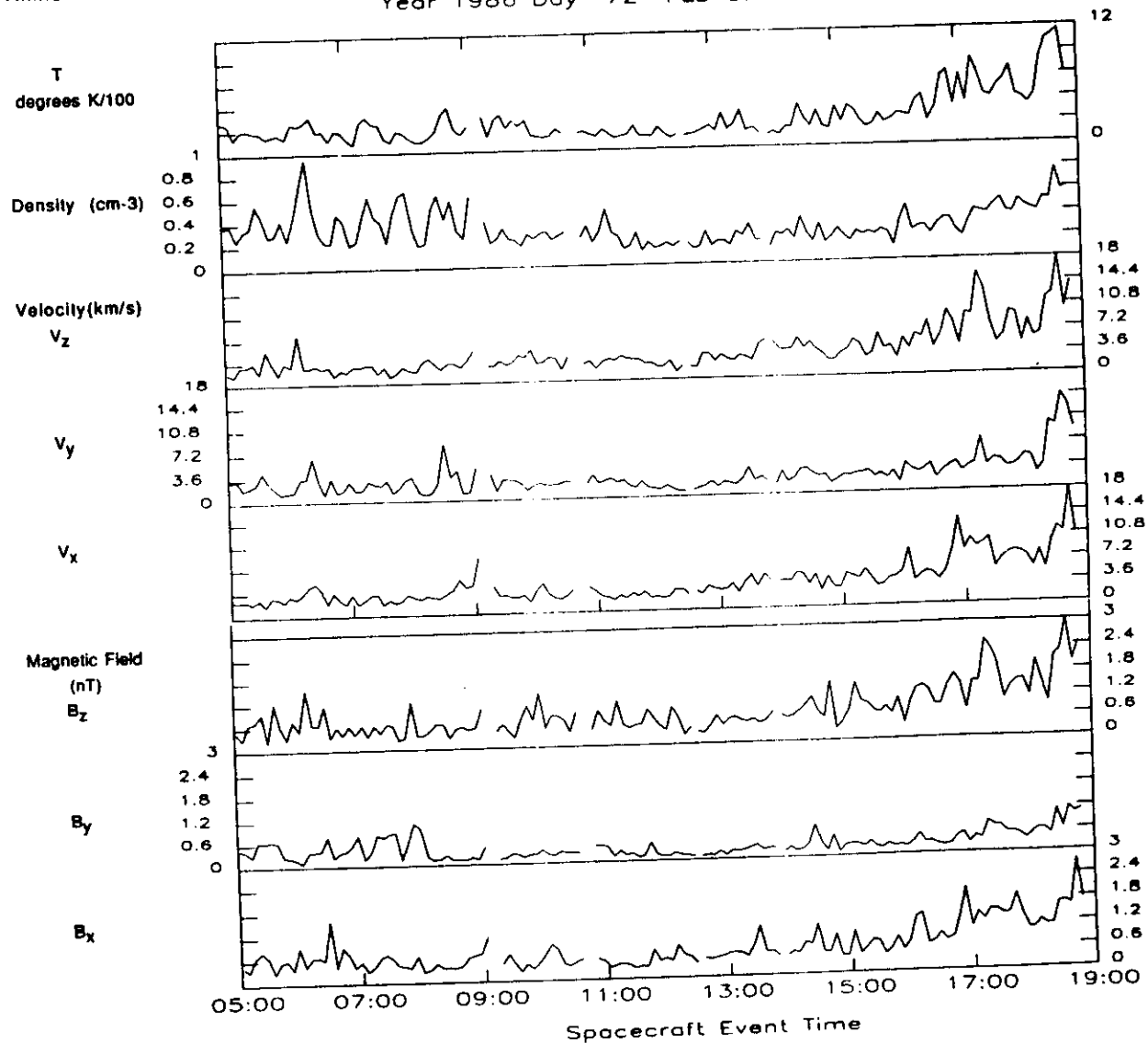
$$\langle \vec{B} \rangle =$$

$$(0.15, -1.25, 4.32)$$

$$|\vec{B}| = 4.5 \text{ nT}$$

R.M.S. Deviation

Year 1986 Day 72 FLD coordinates



WAVE ENERGY DENSITY

The wave energy density is obtained by adding the energy densities of the magnetic and electric components :-

$$W = \frac{1}{2} (\rho < v_1^2 > + \frac{< B_1^2 >}{\mu_0})$$

where $< v_1^2 >$ and $< B_1^2 > / \mu_0$ are the mean square amplitudes of the waves.

POYNTING VECTOR

The poynting vector, S, for the wave is defined as,

$$S = E \times H$$

where E_1 and H_1 are the wave fields and may be obtained from the data by putting

$$E_1 = v_1 \times B_0$$

where B_0 is the interplanetary magnetic field and v_1 is the velocity wave amplitude, and,

$$H_1 = \frac{B_1}{\mu_0}$$

giving,

$$S = (v_1 \times B_0) \times \frac{B_1}{\mu_0}$$

FINDING GROUP VELOCITY

It can be shown that S is equal to the wave energy density multiplied by the wave group velocity. i.e.

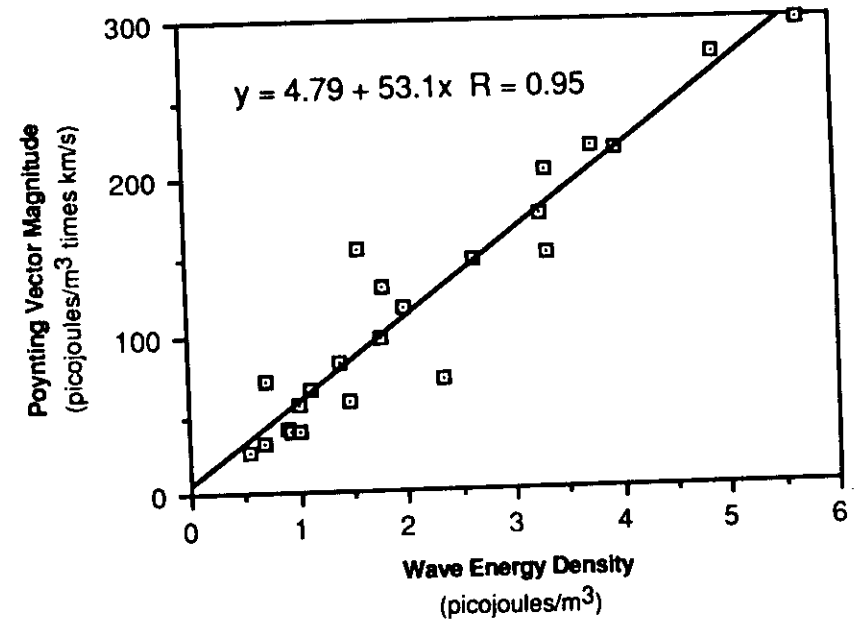
$$S = \frac{1}{2} \left(\epsilon_0 E_1^2 + \frac{B_1^2}{\mu_0} \right) V_g k$$

or, as measured,

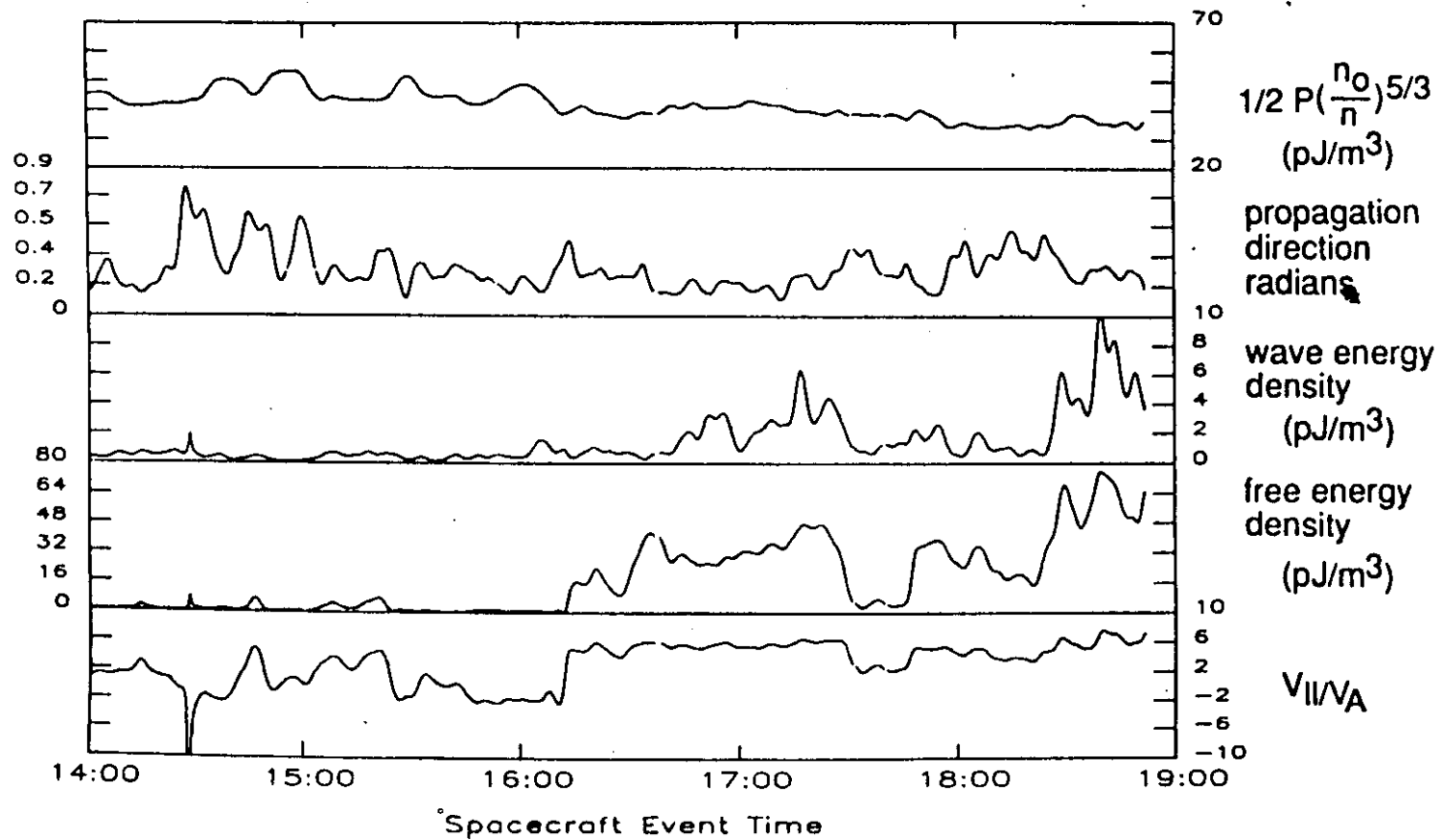
$$S = \frac{1}{2} \left(\rho \langle v_1^2 \rangle + \frac{\langle B_1^2 \rangle}{\mu_0} \right) V_g k$$

hence,

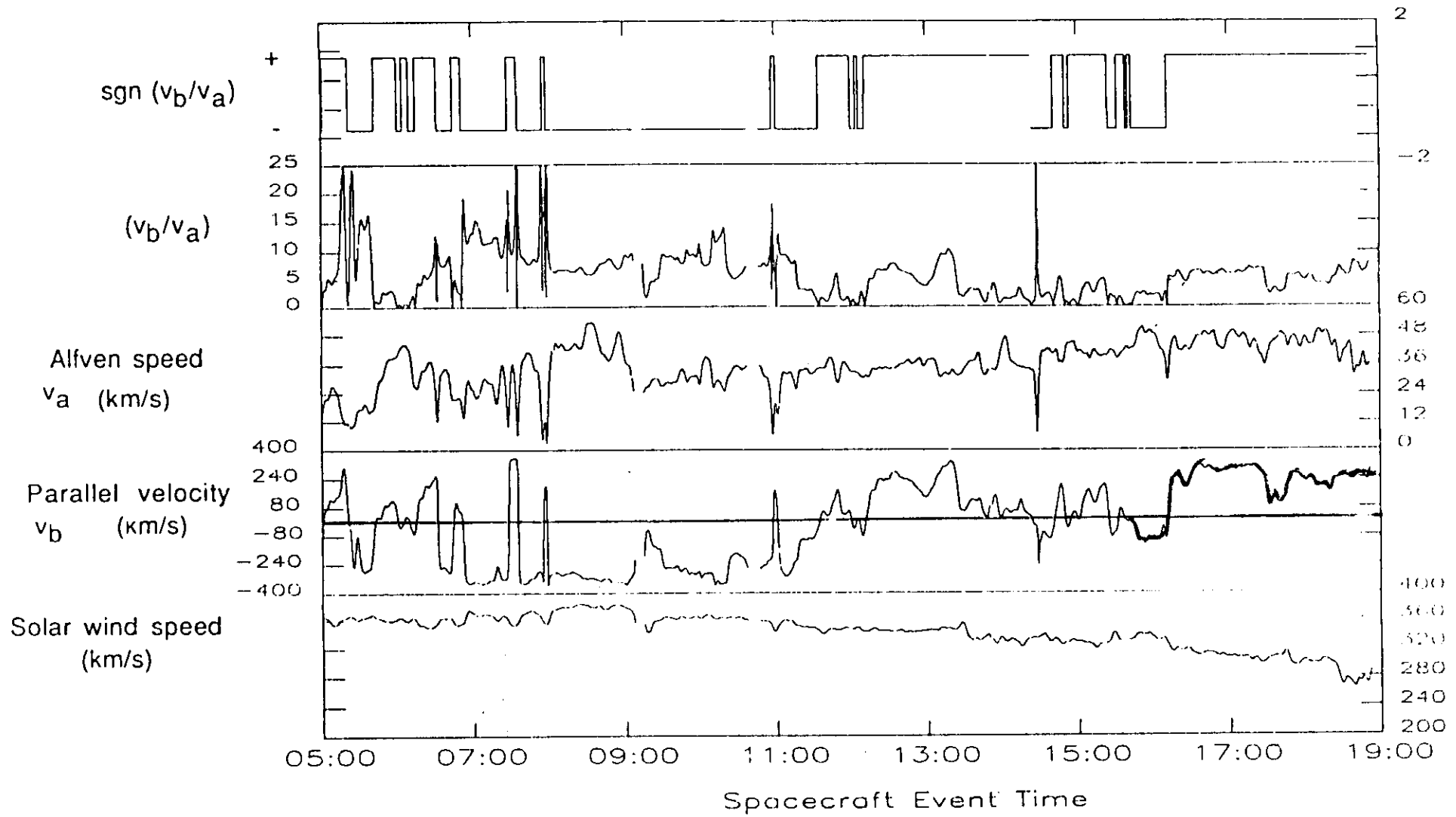
$$V_g = \frac{(\mathbf{v}_1 \times \mathbf{B}_0) \times \frac{\mathbf{B}_1}{\mu_0}}{\frac{1}{2} \left(\rho \langle v_1^2 \rangle + \frac{\langle B_1^2 \rangle}{\mu_0} \right)}$$



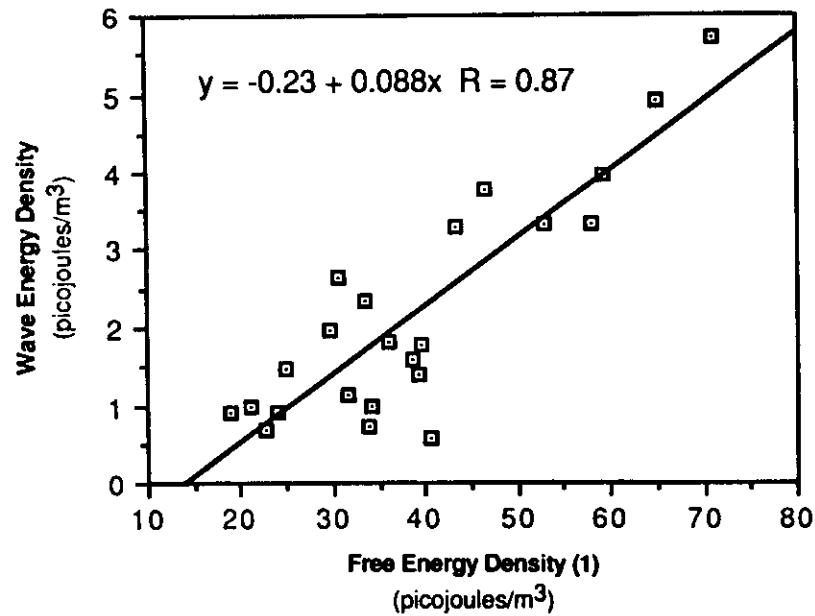
GIOTTO JPA FIS Solar Wind
Year 1986 Day 72



GIOTTO JPA FIS Solar Wind
Year 1986 Day 72



Plot Done! 2-AUG-87 15:35:23



FREE ENERGY

$$E_{\text{ring}} - E_{\text{soler}} = 2V_A [V_s \cos \phi - V_A] \left(\frac{1}{2} n_i m_i \right)$$

$$\text{if } V_s \cos \phi \gg V_A$$

$$\Delta E = \frac{2V_A}{V_s \cos \phi} \left(\frac{1}{2} \rho_i V_s^2 \cos^2 \phi \right)$$

$$V_A \sim 50 \text{ km/s} \quad V_s \cos \phi \sim 250 \text{ km/s}$$

$$\therefore \Delta E \approx 0.4 \text{ (kinetic energy of ion beam in solar wind frame)}$$

Detected energy in waves

$$\sim 0.09 \text{ (kinetic energy of ion beam)}$$

KINETIC EQUATION FOR WAVES

$$\textcircled{*} (V_{II} - V_A) \frac{\partial U}{\partial x} = \frac{F n_c}{\tau_1} - \frac{U}{\tau_2}$$

F free energy per ion

n_c neutral particle density

$1/\tau_1$ ionisation rate

$1/\tau_2$ damping rate.

Assumes time independent situation

$\textcircled{*}$ In our correlation

$$F \propto U$$

where F is a function of α

$\textcircled{*}$ What is the damping mechanism?

NON LINEAR DAMPING PROCESSES

Achternberg Ast. Astrophys 98, 161, 1981

$\textcircled{*}$ Resonant three wave interaction

$$\omega_1 + \omega_2 = \omega_3 \quad (\text{conservation of energy})$$

$$k_1 + k_2 = k_3 \quad (\text{conservation of momentum})$$

Requires waves travelling in opposite directions

$$k_1 \approx -k_2$$

then ω_3, k_3 is an ion acoustic wave

$\textcircled{*}$ Resonant interaction between thermal ions and the high + low frequency beat wave of two circularly polarised waves

$$(\omega_1 \pm \omega_2) - (k_1 \pm k_2) v_z = 0$$

(non linear Landau damping.)

$\textcircled{*}$ Transfers momentum + energy to solar wind ions $\textcircled{*}$

$$\text{Damping rate } \gamma_d = \frac{1}{\tau_d} = - \left(\frac{1}{32\pi} \right)^{1/2} \frac{e^2}{m^2} \frac{k}{v_A \tau_p^2} \frac{T_e}{mc^2} \int dk |H_k|^2$$

Sagdeev et al GRL 13, 85, 1986

