



INTERNATIONAL ATOMIC ENERGY AGENCY
UNITED NATIONS EDUCATIONAL, SCIENTIFIC AND CULTURAL ORGANIZATION



INTERNATIONAL CENTRE FOR THEORETICAL PHYSICS

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**EXPERIMENTAL WORKSHOP ON HIGH TEMPERATURE
SUPERCONDUCTORS & RELATED MATERIALS
(BASIC ACTIVITIES)**

12 - 30 MARCH 1990

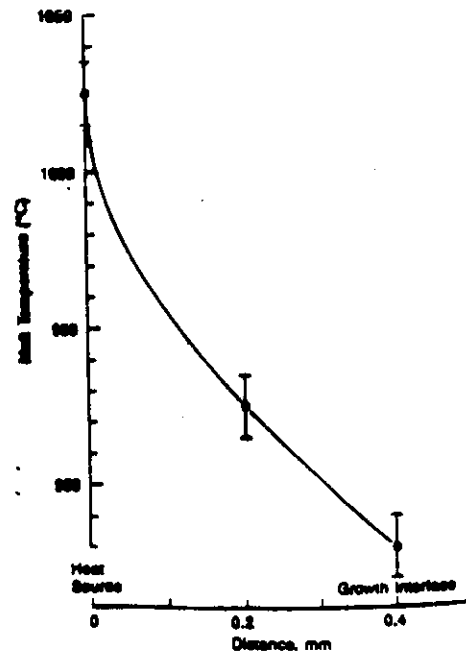
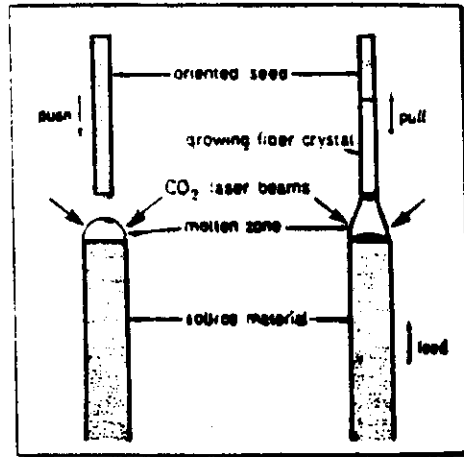
**CRITICAL CURRENTS
Part III**

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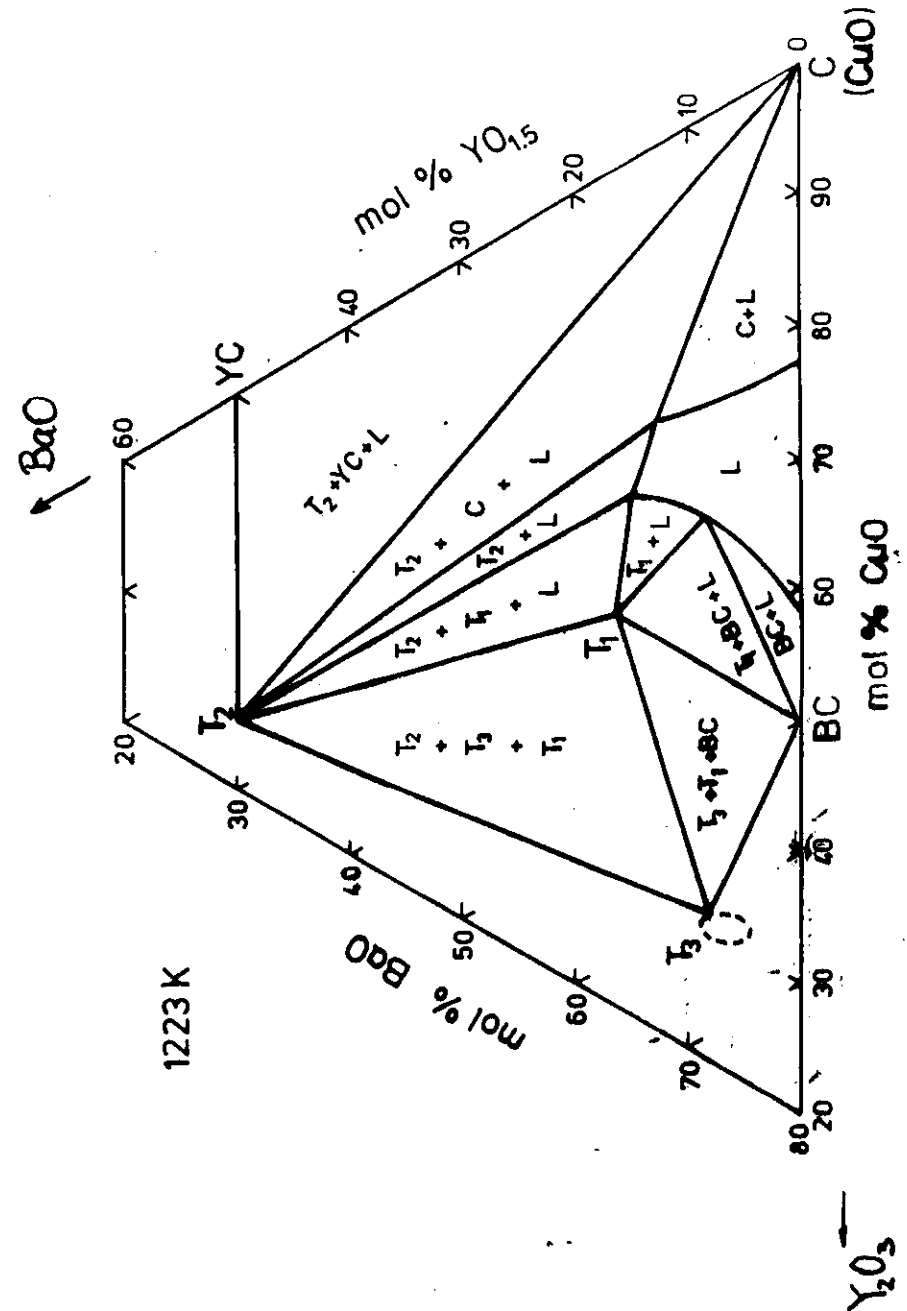
Bi-Sr-Ca-Cu-C
[Kapitulnik et al., Science]

(2212)



At 68K:

$$J_c(0) = 6 \cdot 10^4 \text{ A/cm!}$$



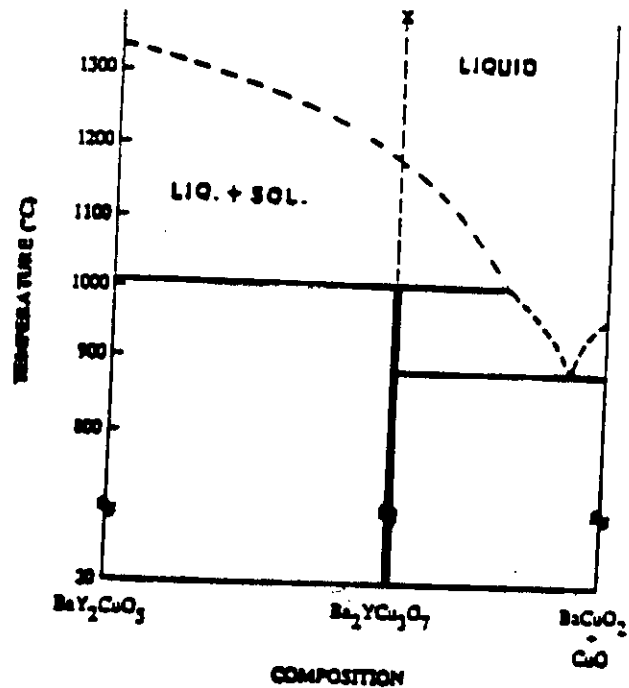
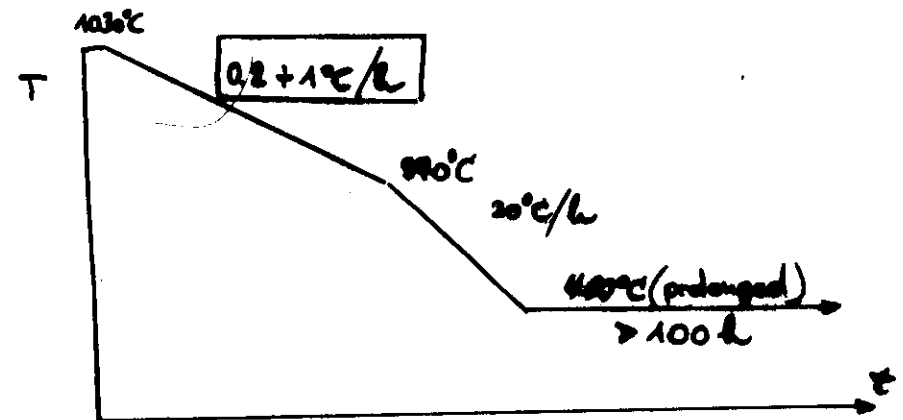


FIG. 1. Qualitative sectional phase diagram near the $\text{YBa}_2\text{Cu}_3\text{O}_{7-x}$ superconducting phase.

Single crystals of $\text{YBa}_2\text{Cu}_3\text{O}_{7-x}$

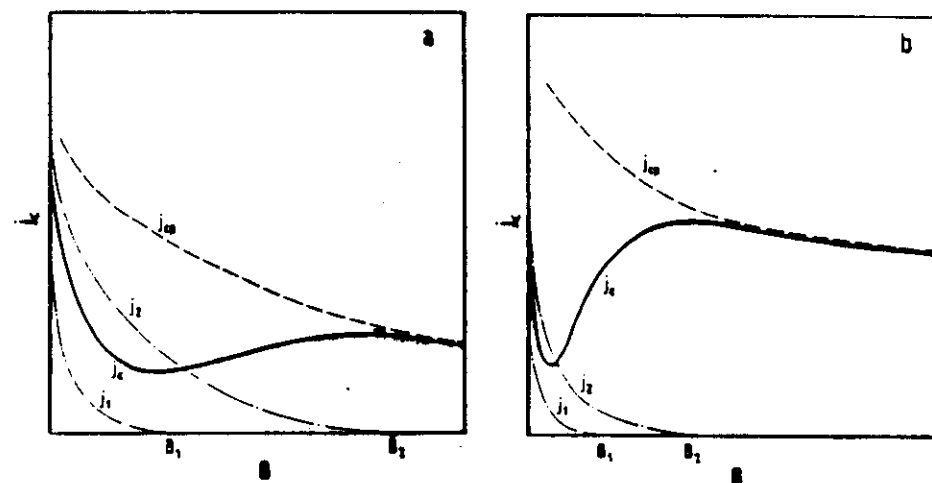
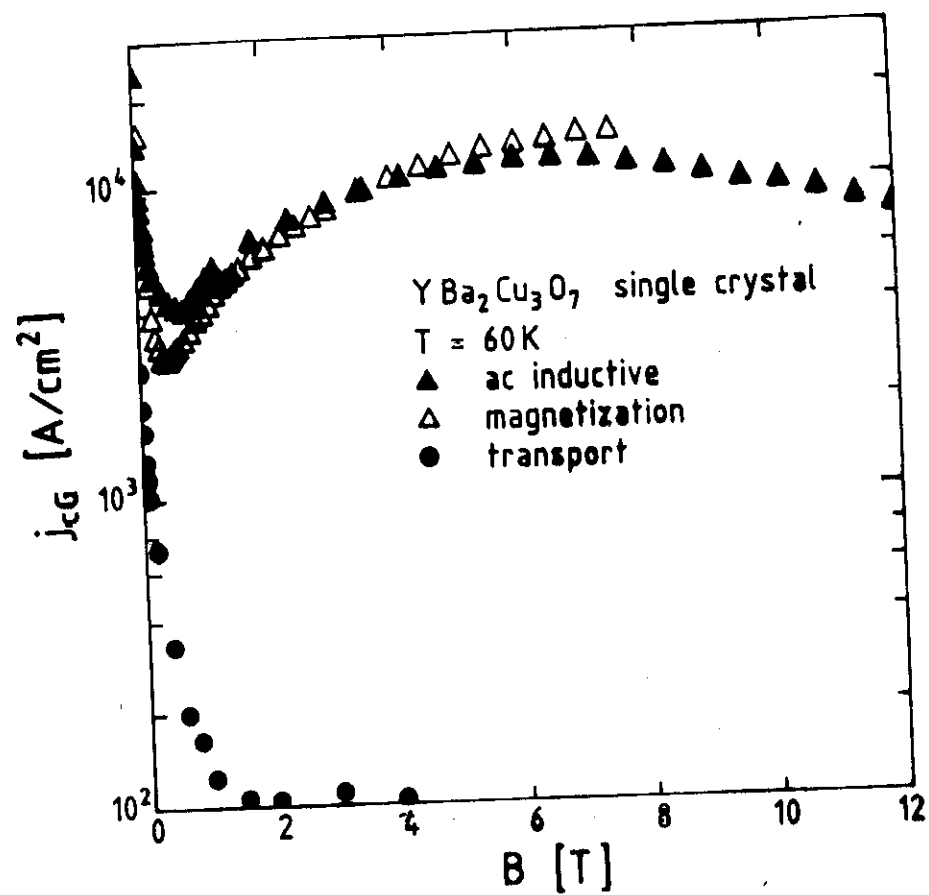
Flux method, CuO flux

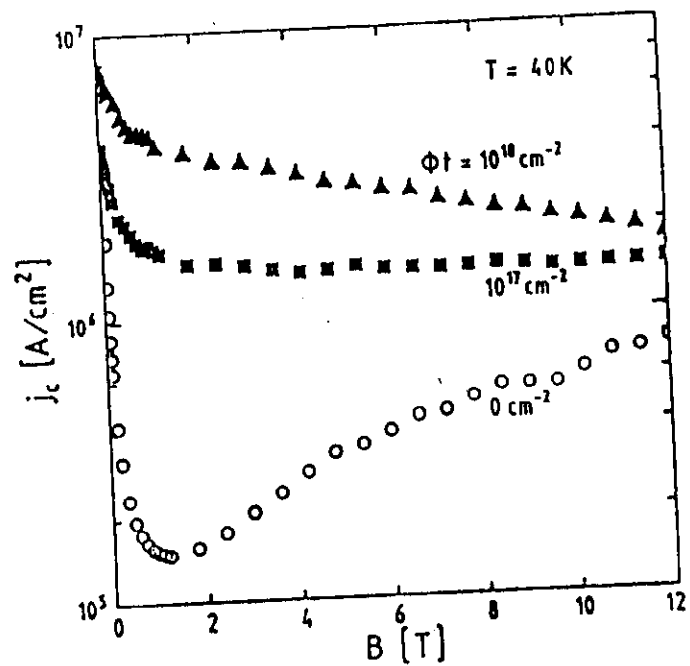
Crucible materials, Al_2O_3 , Au



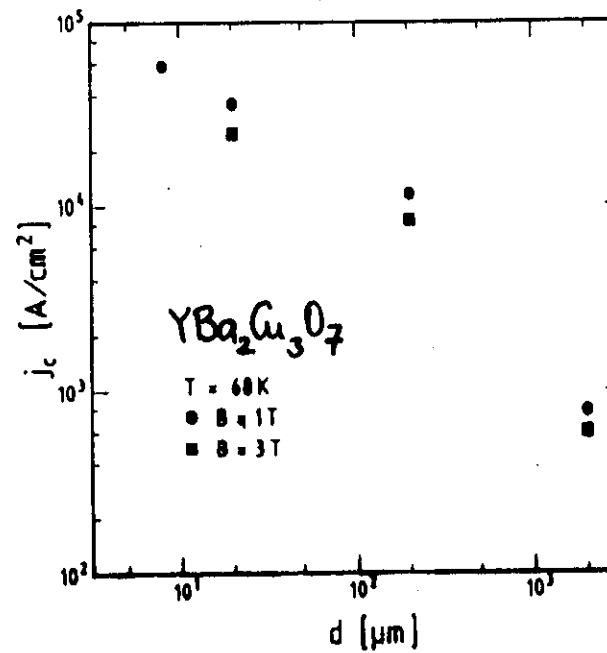
Results : $> 5^\circ/\text{h}$: $d_c \leq 0.1 \text{ mm}$
 $T_c \approx 90 \pm 0.1 \text{ K}$
 $\rho_{\text{lock}} < 80 \mu\Omega \text{ cm}$

$\leq 0.5^\circ/\text{h}$: $d \leq 2 \text{ mm}$
 $T_c \approx 87 \text{ K (in } \text{Al}_2\text{O}_3)$
 $\approx 90 \pm 0.1 \text{ K (in Au)}$
 $\rho_{\text{lock}} \approx 80 \pm 100 \mu\Omega \text{ cm}$
 weak links due to 1% AP_T substitution.



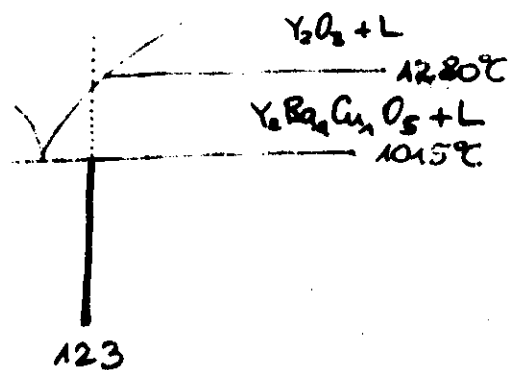


"WEAK LINK" INTRAGRAINS (MONOCRYSTAL)



Melt Textured Growth:

— Jin et al.



— Murakami et al. : $J_c \approx 10^4 A/cm^2$ bei 1T } $77K$
 $= 10^3 A/cm^2$ bei 10T

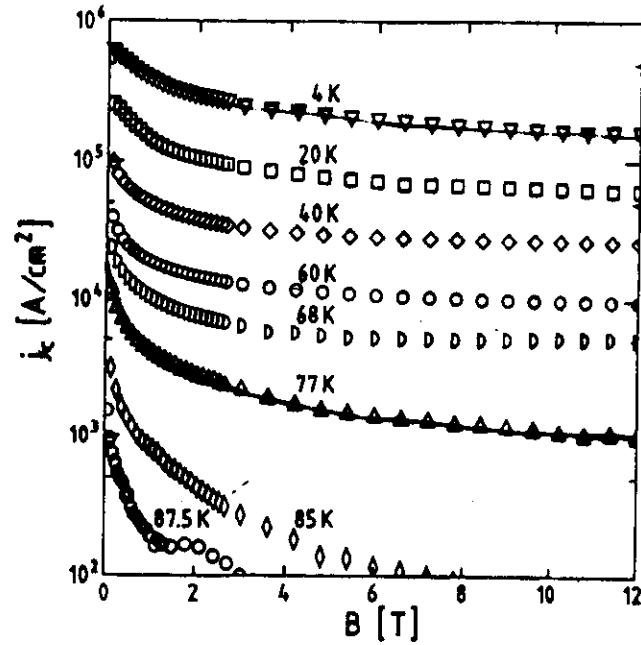
— Salama et al. $J_c = 7.5 \times 10^4 A/cm^2$
 (0T, 77K, 1ms)

$J_c = 1.8 \times 10^4 A/cm^2$
 (0T, 77K, $> 1s$)

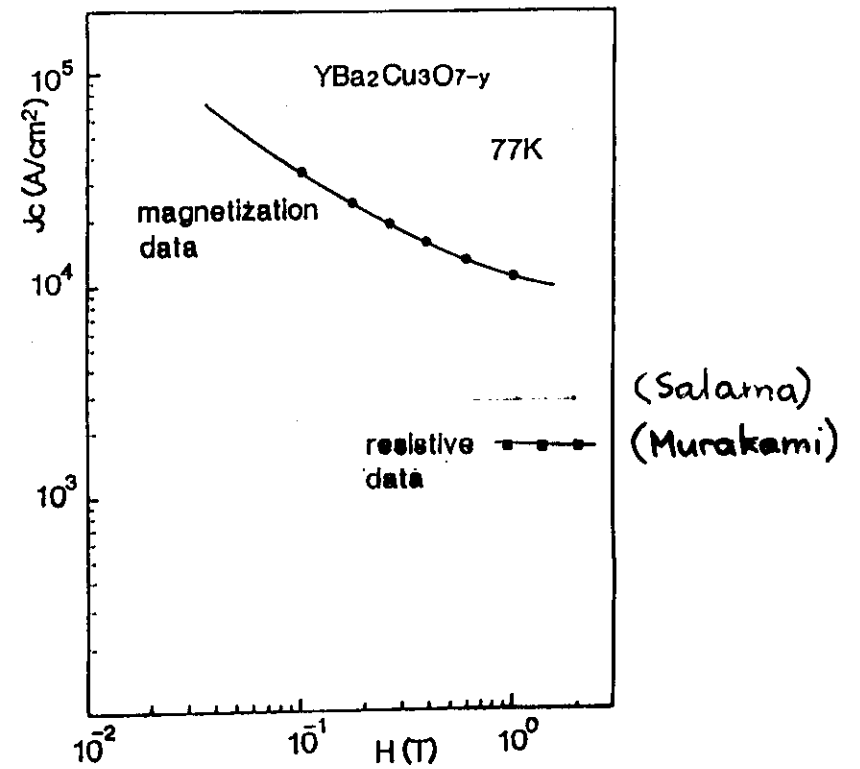


Murakami et al.

$\text{YBa}_2\text{Cu}_3\text{O}_7$: melt textured growth



77K: not longer weak link character



Melt Textured Samples

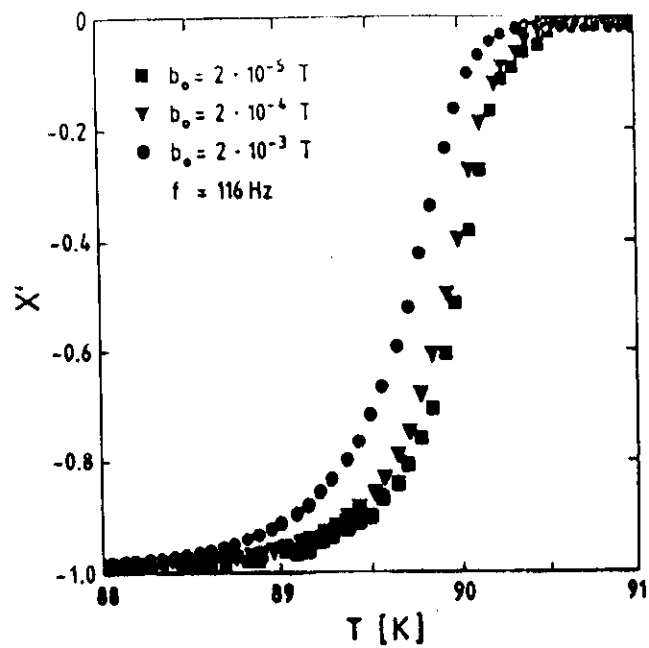


Fig. 2

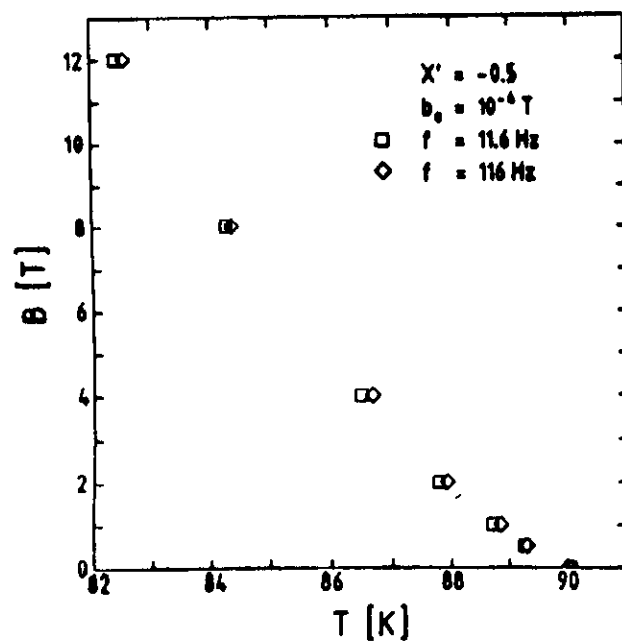


Fig. 3

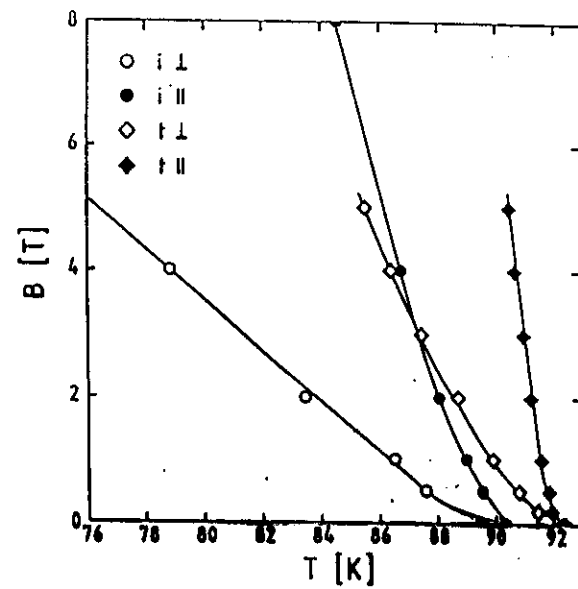


Fig. 1

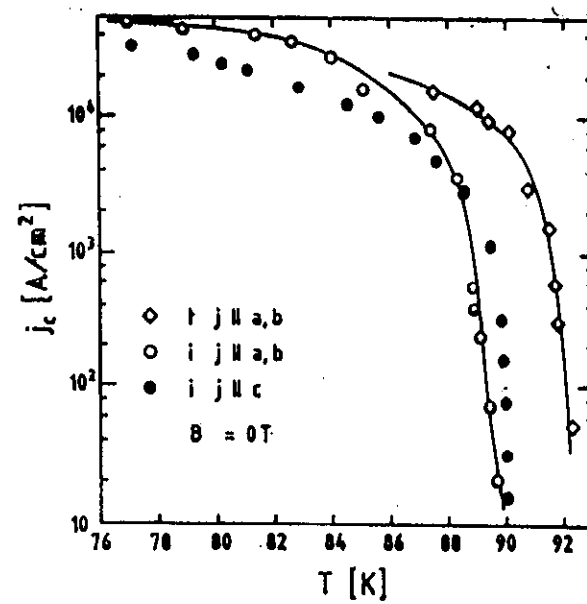


Fig. 2

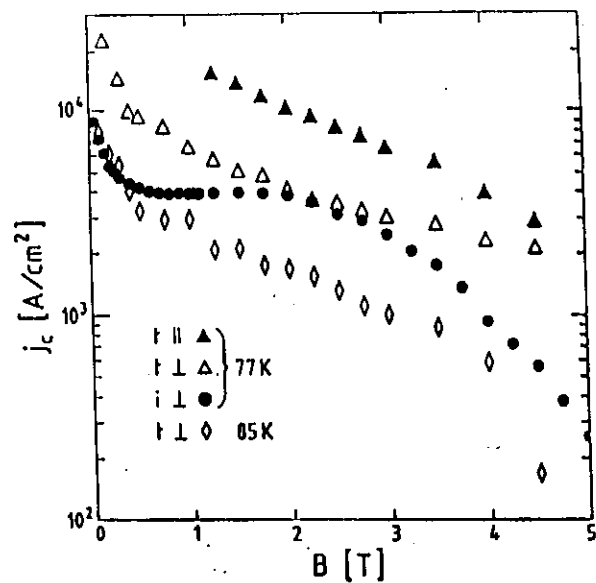


Fig. 3

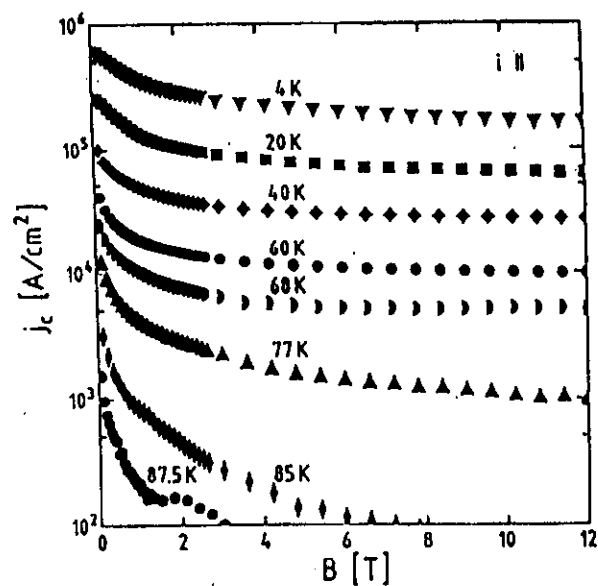


Fig. 4a

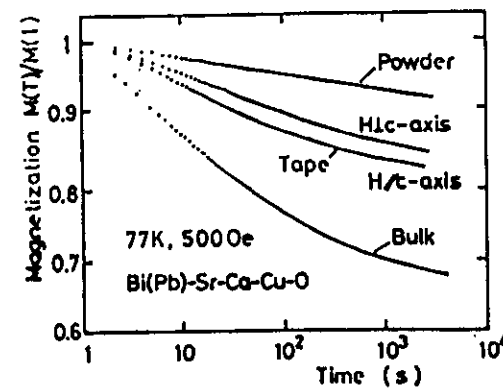
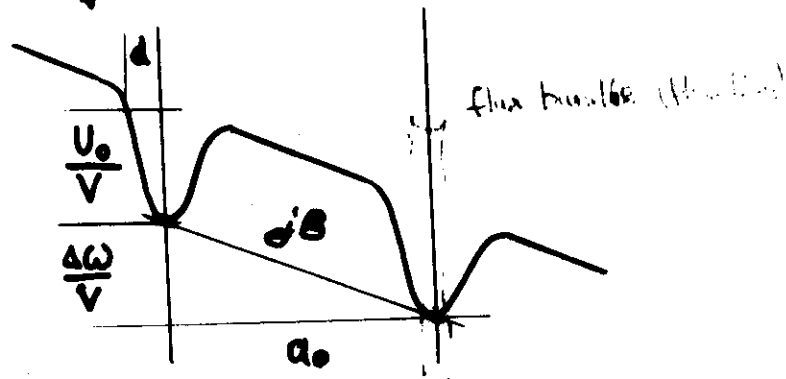


FIGURE 2
Magnetic relaxation at 77K and 5000e of tape,
bulk and powder samples

Thermally activated flux creep

$U_0, \Delta W$: absolute (flux bundle)

$\frac{U_0}{V}, \frac{\Delta W}{V}$: normed (flux line)



$$U_0 = j_{c0} B V d$$

$$\Delta W = j B V a_0$$

$$\Delta W = U_0 \frac{j}{j_{c0}} \left(\frac{a_0}{d} \right)$$

j_{c0} : non-relaxed critical current density

j : instantaneous, relaxed cr. c. density

V : Bundle volume

a_0 : jumping distance

d : Width of potential well
~ reversible displacement

U_0 : Pinning potential

At temperature T , 'particle' jumps

with a rate $R = R_0 \cdot e^{-\frac{U_0}{k_B T}}$ (non directional)

(R_0 : jumping frequency)

Voltage leads to directional jumping:

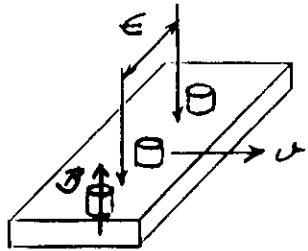
$$R_f = R_0 \cdot e^{-\frac{U_0 - \Delta W}{k_B T}}$$

$$R_b = R_0 \cdot e^{-\frac{U_0 + \Delta W}{k_B T}}$$

ΔW = Energy picked up by 'particle' when introducing a voltage

→ jumping rate : $-\frac{U_0}{k_B T}$

$$R_{net} = 2 R_0 e^{-\frac{U_0}{k_B T}} \sinh\left(\frac{\Delta W}{k_B T}\right)$$



Mit Bewegung der
Flusslinien verbunden
new E -Feld

$$E = vB$$

$$E = R_{\text{neu}} \cdot j \cdot B$$

$$E = 2\alpha \cdot B \cdot \Omega \cdot e^{-\frac{U_0}{k_B T}} \sinh\left(\frac{U_0}{k_B T} \frac{j}{j_0}\right)$$

If j_c has a variation
of $j_{c0} \frac{k_B T}{U_0}$, E varies by e

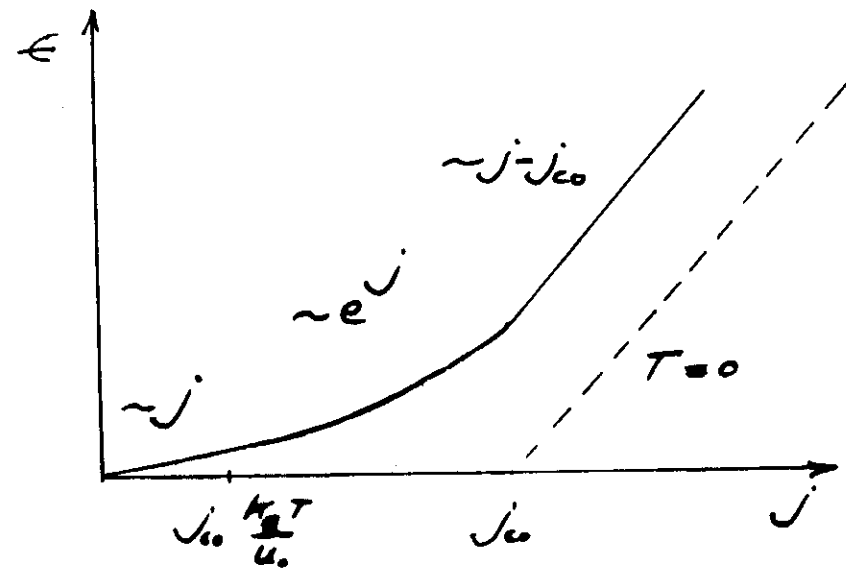
Different cases:

$$\sinh x \rightarrow \frac{e^x}{2}, x \gg 1$$

$$\sinh x \rightarrow x, x \ll 1$$

$U_0 < k_B T$	$U_0 > k_B T$
-	$\Delta W > k_B T$
$\Delta W < k_B T$	$\Delta W < k_B T$

$$: U_0 > k_B T$$

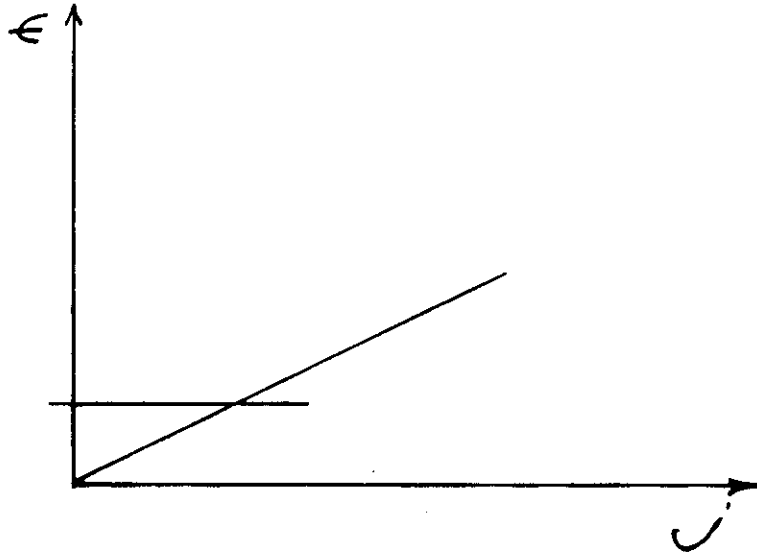


$$E = f_j(j - j_{c0}) + E_{c0}, \quad j > j_{c0}$$

$$= \alpha \cdot B \cdot \Omega \cdot e^{-\frac{U_0}{k_B T}} \left(1 - \frac{j}{j_0}\right), \quad j_0 \frac{k_B T}{U_0} < j < j_{c0}$$

$$= 2\alpha \cdot B \cdot \Omega \cdot \frac{U_0}{k_B T j_0} \cdot j \cdot e^{-\frac{U_0}{k_B T}}, \quad j < j_0 \frac{k_B T}{U_0}$$

$$U_0 \lesssim \mu_0 T$$

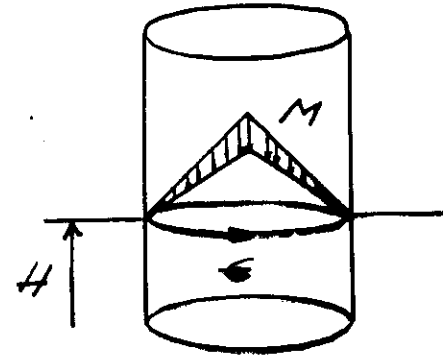


$$U = R_s j,$$

$$R_s = 29.8 \Omega \cdot \frac{U_0}{\mu_0 T} \frac{1}{j_c} e^{-\frac{U}{U_0}}$$

- Finite voltmeter resolution
 - Short sample length
- lead to finite j_c !

Electrodynamics of Flux Creep



Cylinder with radius
in an external field

Faraday'sches Gesetz

$$\text{rot } \vec{E} = - \frac{\partial \vec{B}}{\partial t}$$

$$E = - \frac{1}{2\pi R} \frac{\partial \Phi}{\partial t}$$

Bean'sches Modell

$$j = J \frac{M}{H}$$

$$\Phi = \frac{2}{3} \pi \mu_0 R^3 j$$

Φ eliminieren

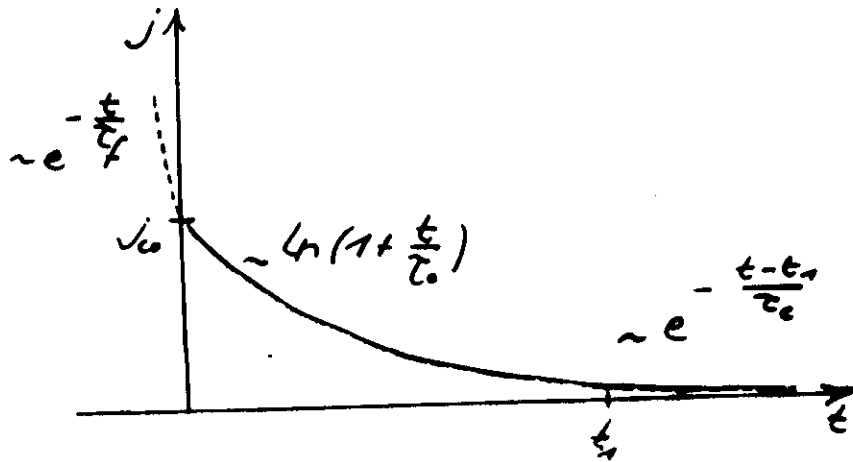
$$E = - \frac{1}{6} \mu_0 R^3 \frac{dj}{dt}$$

E eliminieren

$$\text{DGL für } j = j(t)$$

Time dependence of j_c

$$u_0 > k_B T$$



$$----- j(t) = j_0 + 4j_0 e^{-\frac{t}{\tau_f}}$$

$$\tau_f = \mu_0 R^2 / 6 \rho_f \dots \sim 10^{-6} s$$

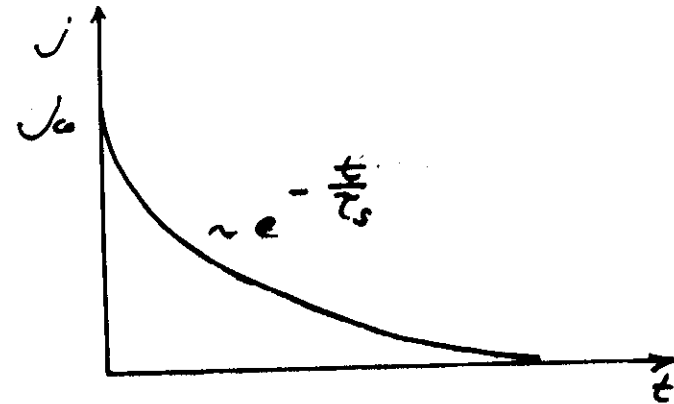
$$----- j(t) = j_0 [1 - \frac{k_B T}{u_0} \ln(1 + \frac{t}{\tau_0})]$$

$$\tau_0 = \frac{1}{6} \mu_0 \frac{R^2}{\rho_s \Omega} \frac{k_B T}{u_0} \frac{j_0}{B}, \sim 10^{-6} \dots 10^{-7} s$$

$$j(t) = j_0 \frac{k_B T}{u_0} e^{-\frac{t-t_1}{\tau_c}}, \quad t > t_1 \sim e^{-\frac{u_0}{k_B T}}$$

$$\tau_c = \frac{1}{2} \tau_0 e^{\frac{u_0}{k_B T}}$$

$$u_0 \lesssim k_B T$$

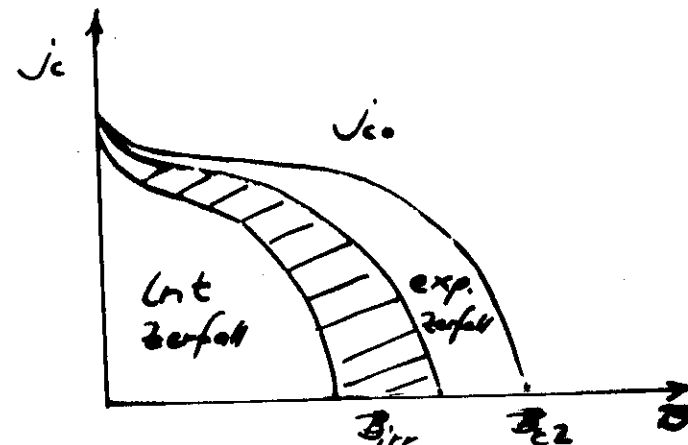


$$j(t) = j_0 e^{-\frac{t}{\tau_c}}, \quad t \geq 0$$

$$\tau_s = \mu_0 R^2 / 6 \rho_s$$

$$\tau_f \ll \tau_s \ll \tau_c$$

Irreversibility Line



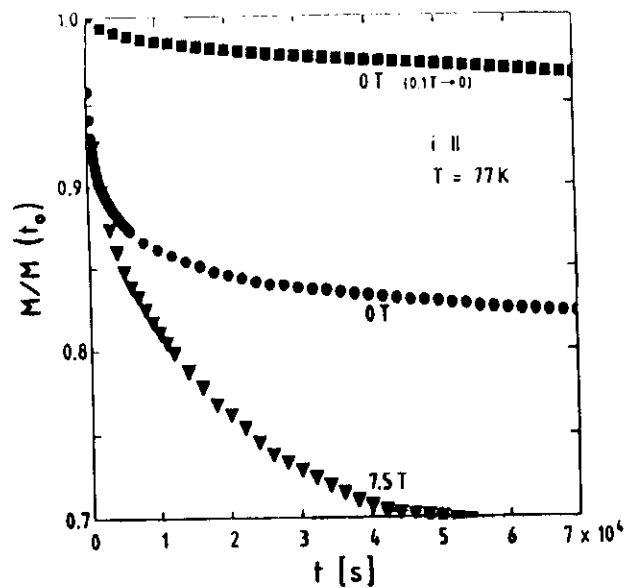


Fig. 1

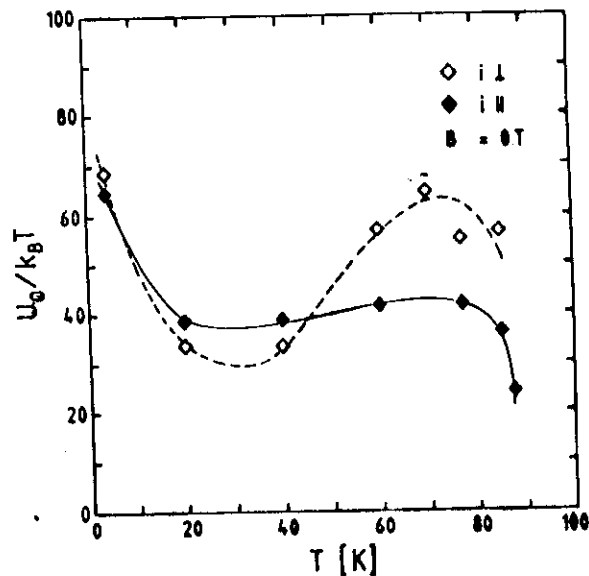
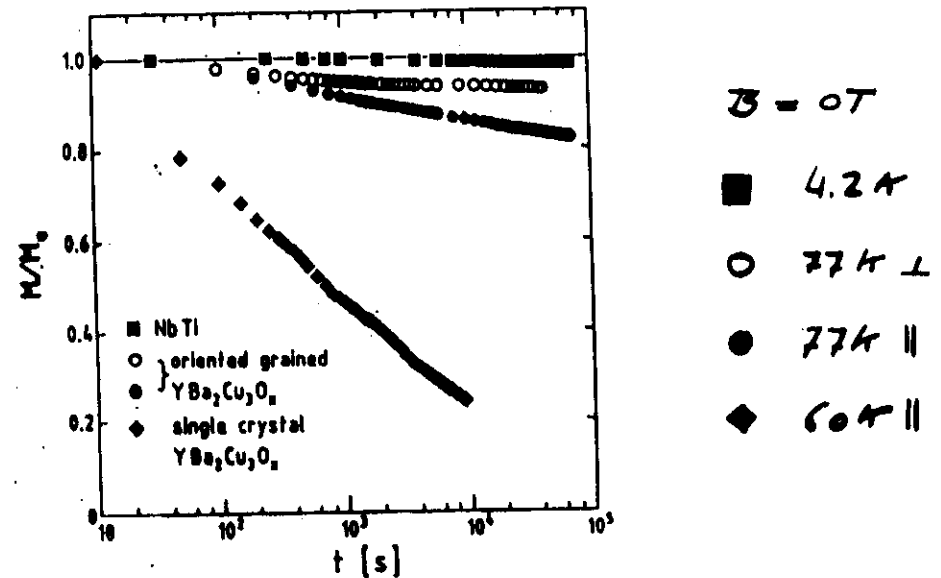


Fig. 2

Comparison with Nb 49 wt.% Ti

$$T_c = 9 K$$

$$j_c(4.2 K, 0 T) = 3.5 \cdot 10^5 \frac{A}{cm^2}$$



		$U_0/k_B T$	$U_0 [meV]$
oriented grained $YBa_2Cu_3O_x$ 77 K	//	47	272
	⊥	55	365
NbTi	4.2 K	487	176
	7.5 K	182	118
single crystal $YBa_2Cu_3O_x$	60 K	8	41

$B = 0 T$ $B \perp a, b$ $t_s = 75 s$
 $B_s = 3 T$ $\dot{B} = 12 mT/s$

