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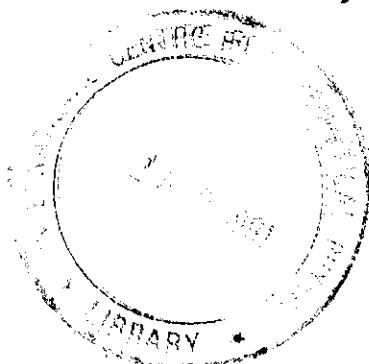
**WINTER COLLEGE ON "MULTILEVEL TECHNIQUES IN  
COMPUTATIONAL PHYSICS"**

***Physics and Computations with Multiple Scales of Lengths***  
**(21 January - 1 February 1991)**

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H4.SMR 539/11

***Non-Scalar Systems***



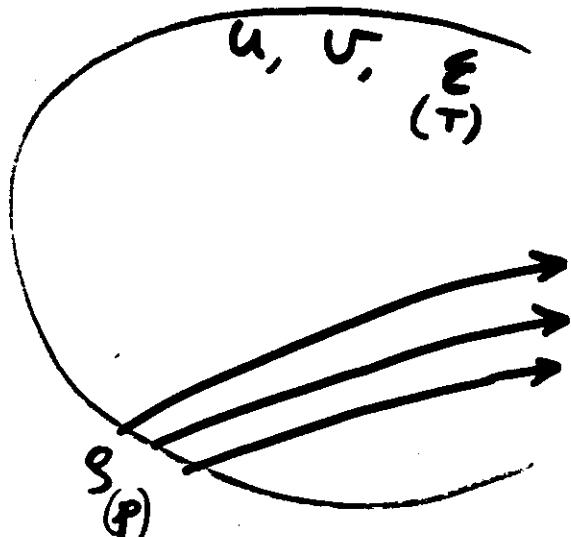
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# Compressible Navier-Stokes : 2D

- $-\mu \Delta u + \rho u u_x + \rho v u_y - (\lambda + \mu)(u_x + v_y)_x + p_x = 0$
- $-\mu \Delta v + \rho u v_x + \rho v v_y - (\lambda + \mu)(u_x + v_y)_y + p_y = 0$
- $-\kappa \Delta \epsilon + \rho u \epsilon_x + \rho v \epsilon_y + p(u_x + v_y)$   
 $-\mu(u_x + v_y)^2 - \lambda(u_x + v_y)^2 - 2\mu(u_x^2 + v_y^2) = 0$
- $(\rho u)_x + (\rho v)_y - \nu \Delta g = 0$
- $p = p(\epsilon, g)$

Non-elliptic BVP.  
unknown  $u, v, \epsilon, g, p$

B.C.:



Inviscid case:  $\lambda, \mu, \kappa \ll gl \max(|u|, |v|)$   
 $l$  = length at which  $(u, v, \epsilon)$  change.  $\rightarrow$  Euler eqn.  
 Usually there are viscous layers.

Incompressible:

$$Q = -\frac{1}{R} \Delta + u \partial_x + v \partial_y.$$

$$\begin{pmatrix} Q & 0 & \partial_x \\ 0 & Q & \partial_y \\ \partial_x & \partial_y & 0 \end{pmatrix} \begin{pmatrix} u \\ v \\ p \end{pmatrix} = 0$$

$$\det \uparrow = (-\Delta) Q \quad \text{elliptic.}$$

large Re: Sing. Pert.

# "True" MG Efficiency

- ~ .1 error reduction per cycle
- 1-FMG solver to  $O(h^2)$  accuracy  
~ 10 minimal work units
- Predicted theoretically

Simple elliptic  
mode analysis

1972-3

rigorous  $\square$

1977

General elliptic systems

1979-84

rigorous: exact ( $\Rightarrow$  opt.)

general  $\mathcal{G}$ , norms, cycles

1985-90

SD transonic

1976

etc.

SORTH

General non-elliptic  $O(h^2)$   
mode analysis

1990

YAVNEH

h-principal

$U = t$

die:

$$\text{Cauchy-Riemann} \quad \begin{pmatrix} \partial_x & \partial_y \\ \partial_y & -\partial_x \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix} = 0 \quad -\Delta$$

$$\text{2D Stokes} \quad \begin{pmatrix} -\Delta & 0 & \partial_x \\ 0 & -\Delta & \partial_y \\ \partial_x & \partial_y & 0 \end{pmatrix} \begin{pmatrix} u \\ v \\ p \end{pmatrix} \quad \Delta^2$$

$$\text{2D Incompressible Navier-Stokes} \quad \begin{pmatrix} Q & 0 & \partial_x \\ 0 & Q & \partial_y \\ \partial_x & \partial_y & 0 \end{pmatrix} \begin{pmatrix} u \\ v \\ p \end{pmatrix} \quad -Q\Delta$$

$Q = -\frac{1}{R}\Delta + \underline{u} \cdot \nabla$

$$\text{2D Euler} \quad \begin{pmatrix} g\underline{u} \cdot \nabla & 0 & 0 & 0 & \partial_x \\ 0 & g\underline{u} \cdot \nabla & 0 & 0 & \partial_y \\ g\partial_x & g\partial_y & \underline{u} \cdot \nabla & 0 & 0 \\ p\partial_x & p\partial_y & 0 & g\underline{u} \cdot \nabla & 0 \\ 0 & 0 & -p_g & -p_e & 1 \end{pmatrix} \begin{pmatrix} u \\ v \\ g \\ \epsilon \\ p \end{pmatrix} \quad \begin{aligned} g^3(\underline{u} \cdot \nabla)^2 \times \\ [(u \cdot \nabla)^2 - a^2 \Delta] \\ a^2 = p_g + \frac{p}{g^2} p_e \end{aligned}$$

$$\text{Compressible Navier-Stokes (on the viscous scale)} \quad k\mu(2\mu+\lambda) \Delta^3 (\underline{u} \cdot \nabla)$$

$$\text{Central Cauchy-Riemann} \quad \Delta^{2h}$$

$$\text{Central (Navier-) Stokes} \quad Q^h \Delta^{2h}$$

# Cauchy - Riemann Equations

$$① U_y - V_x = F$$

$$② U_x + V_y = G$$

One boundary condition

$$L_U = \begin{pmatrix} \partial_y & -\partial_x \\ \partial_x & \partial_y \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix} = \begin{pmatrix} F \\ G \end{pmatrix}$$

$$\det L = \partial_x^2 + \partial_y^2 = \Delta$$

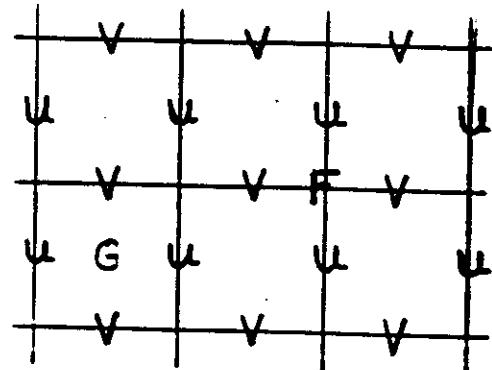
## Difference Equations :

Eq. 1 defined at grid nodes

Eq. 2 defined at cell centers

$$\partial_x^h = (T_x^{1/2} - T_x^{-1/2})/h,$$

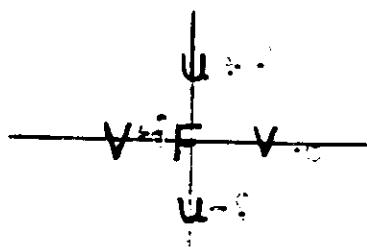
$\det L^h = \Delta^h$  5-point



## Relaxation on Equation 1 :

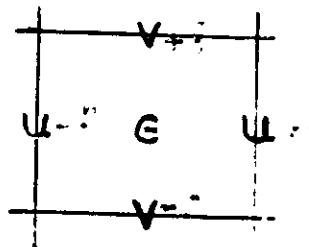
preserves  $R_2 = G - U_x - V_y$

smoothes  $R_1 = F - U_y + V_x$



## Relaxation on Equation 2 :

Smoothing Factor :  $\bar{\mu} = .5$



Red-Black :  $\bar{\mu} = .25$

# Central Differencing of Cauchy-Riemann Eq.

$$U_y - V_x = F$$

$$U_x + V_y = G$$

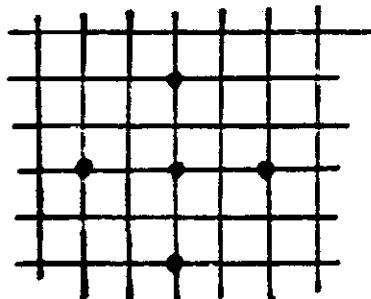
$$\frac{U^h(x, y+h) - U^h(x, y-h)}{2h} - \frac{V^h(x+h, y) - V^h(x-h, y)}{2h} = F^h(x, y)$$

$$\frac{U^h(x+h, y) - U^h(x-h, y)}{2h} + \frac{V^h(x, y+h) - V^h(x, y-h)}{2h} = G^h(x, y)$$

$$L^h \begin{pmatrix} U^h \\ V^h \end{pmatrix} = \begin{pmatrix} \frac{T_y - T_y^{-1}}{2h} & -\frac{T_x - T_x^{-1}}{2h} \\ \frac{T_x - T_x^{-1}}{2h} & \frac{T_y - T_y^{-1}}{2h} \end{pmatrix} \begin{pmatrix} U^h \\ V^h \end{pmatrix} = \begin{pmatrix} F^h \\ G^h \end{pmatrix}$$

$$\det L^h = \frac{1}{4h^2} (T_x^2 + T_x^{-2} + T_y^2 + T_y^{-2} - 4) = \Delta^{2h}$$

unstable modes:



$$\begin{matrix} & & 1 & & \\ & 1 & -4 & 1 & \\ & & 1 & & \end{matrix}$$

a	b	a	b	a	b
c	d	c	d	c	d
a	b	a	b	a	b
c	d	c	d	c	d
a	b	a	b	a	b
c	d	c	d	c	d

Positive type but "Quasi-elliptic"

# Distributive Relaxation

**Cauchy-Riemann:**  $L = \begin{pmatrix} \partial_y & -\partial_x \\ \partial_x & \partial_y \end{pmatrix}$

Relaxing Eq. 1 is equivalent to changing a function  $w_1$  at the relaxed vertex, with

$$u = \partial_y^h w_1, \quad v = -\partial_x^h w_1.$$

Relaxing Eq. 2 - changing  $w_2$  at the relaxed center

$$u = \partial_x^h w_2, \quad v = \partial_y^h w_2$$

$\Leftrightarrow$  GS relaxation for  $\underline{w} = \begin{pmatrix} w_1 \\ w_2 \end{pmatrix}$ , where

$$u = \partial_y w_1 + \partial_x w_2, \quad v = -\partial_x w_1 + \partial_y w_2$$

$$\text{or } \underline{u} = M \underline{w}, \quad M = \begin{pmatrix} \partial_y & \partial_x \\ -\partial_x & \partial_y \end{pmatrix}$$

$\Leftrightarrow$  GS for  $LM \underline{w} = \underline{f}$ ,  $LM = \begin{pmatrix} \Delta & 0 \\ 0 & \Delta \end{pmatrix}$

$$\therefore \bar{\mu}_L = \bar{\mu}_\Delta = \begin{cases} .5 & : \text{Lexicographic}^2 \\ .25 & : \text{RB}^2 \text{ (one or two sweeps/cycle)} \end{cases}$$

**Generally**  $\bar{\mu}_L = \bar{\mu}_{\det L}$  can be obtained for any choice of relaxation for  $\det L$

**Method:**  $M_{ij} = \text{cofactor of } L_{ji}$

$$\Rightarrow LM = (\det L) I$$

# REGULARIZATION OF SYSTEMS

Guiding Principle: Principal  $\det L = l_1 l_2 \dots$

$\Rightarrow \bar{\mu}_L = \max[\bar{\mu}_{l_1}, \bar{\mu}_{l_2}, \dots]$  can be obtained

Theorem (Brandt - Joseph).  $\det L = l_1 l_2$

$\Rightarrow \exists L_i$  such that  $L = L_1 L_2, \det L_i = l_i$ .

Relaxation of  $L, L_2 \underline{U} = \underline{f}$

by alternately relaxing  $L_1 \underline{U} = \underline{f}, L_2 \underline{U} = \underline{U}$

$$\begin{array}{c|c|c|c} \bar{\omega}_2 & \bar{U}_2 & \bar{\omega}_2 \\ \hline U_1 & \bar{\omega}_1 & U_1 \\ \hline \bar{\omega}_2 & \bar{U}_2 & \bar{\omega}_2 \end{array}$$

Example: Elasticity

$$\begin{pmatrix} M\Delta + \lambda \partial_{xx} & \lambda \partial_{xy} \\ \lambda \partial_{xy} & M\Delta + \lambda \partial_{yy} \end{pmatrix} = \begin{pmatrix} \partial_x & \partial_y \\ \partial_y & -\partial_x \end{pmatrix} \begin{pmatrix} M + \lambda & 0 \\ 0 & M \end{pmatrix} \begin{pmatrix} \partial_x & \partial_y \\ \partial_y & -\partial_x \end{pmatrix}$$

$\Rightarrow \bar{\mu} = \bar{\mu}_\Delta$  even for  $M \ll \lambda$ .

usually: find

$$LM = \begin{pmatrix} l_{11} & & & 0 \\ & l_{22} & & \\ & & \ddots & \\ & & & l_{qq} \end{pmatrix}$$

# Incompressible Navier Stokes: 2D

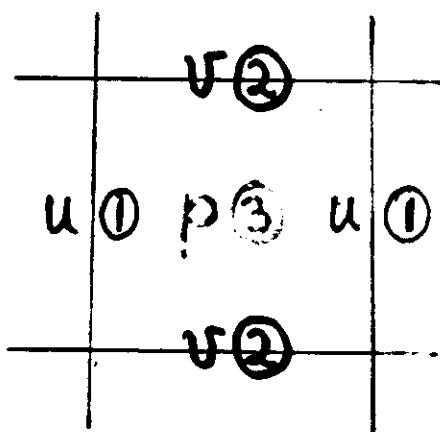
- ①  $QU + P_x = f_1$
- ②  $QV + P_y = f_2$
- ③  $U_x + V_y = f_3$

$$Q = -\frac{1}{h} \Delta + \tilde{U} \partial_x + \tilde{V} \partial_y$$

two boundary conditions

$$L \underline{U} = \begin{pmatrix} Q & 0 & \partial_x \\ 0 & Q & \partial_y \\ \partial_x & \partial_y & 0 \end{pmatrix} \begin{pmatrix} U \\ V \\ P \end{pmatrix} = \begin{pmatrix} f_1 \\ f_2 \\ f_3 \end{pmatrix}$$

$$\det L = -\Delta Q$$



## Differencing:

$$\partial_x^h = (T_x^{1/2} - T_x^{-1/2})/h. \quad \det L^h = -\Delta^h Q^h$$

$h$ -ellipticity of  $L^h$  depends on  $Q^h$ .

## Relaxation:

distribution  $M = \begin{pmatrix} 1 & 0 & -\partial_x \\ 0 & 1 & -\partial_y \\ 0 & 0 & \frac{Q}{\Delta} \end{pmatrix} \rightarrow LM = \begin{pmatrix} Q & 0 & 0 \\ 0 & Q & 0 \\ \partial_x & \partial_y & -\Delta \end{pmatrix}$

cofactors of last row in  $L$ , divided by their common divisor.

Relaxation of  $L$  is reduced to that of  $Q$ .

Box GS on coarsest grids

Higher-order smoothing boundary conditions

## FAS

FMG solution to  $O(h^2)$  in one  $W(2,0)$ .

Non staggered grids. Conservative schemes

## Relaxation:

① GS for  $Qu = f_1 - p_x$

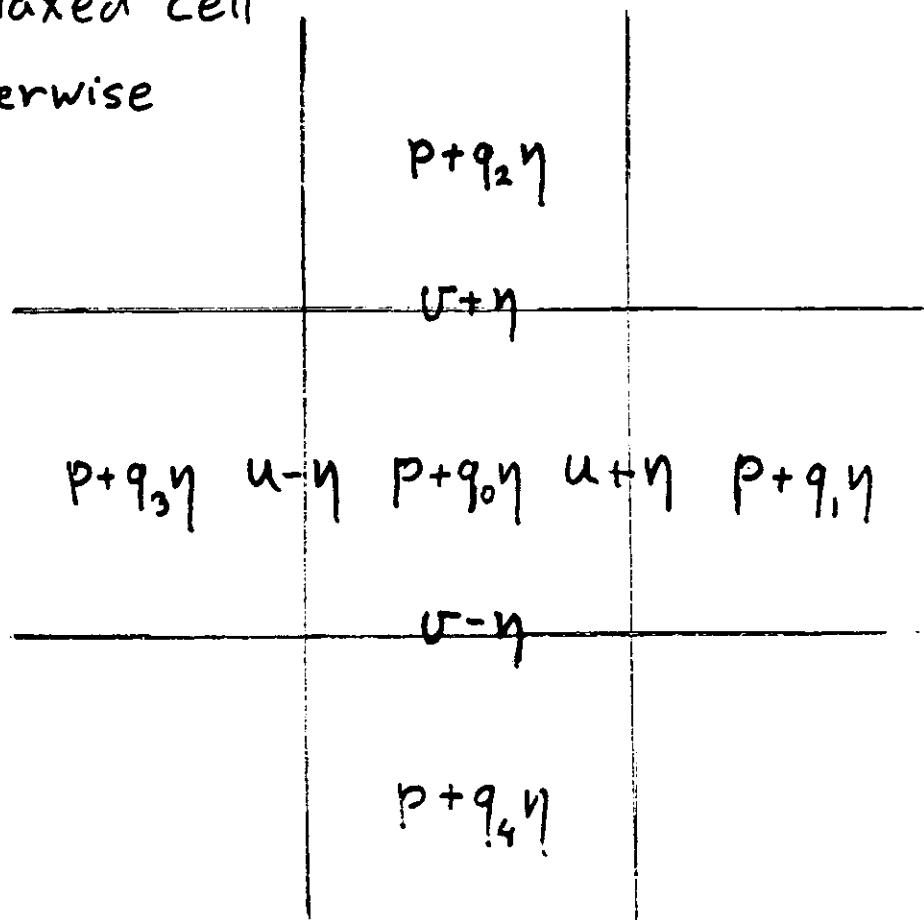
② GS for  $Qv = f_2 - p_y$

③ GS for  $U_x + V_y = f_3$

with  $\delta u = -\omega_x$ ,  $\delta v = -\omega_y$ ,  $\delta p = Q\omega$

$\Rightarrow \delta(Qu + p_x) \approx 0$ ,  $\delta(Qv + p_y) \approx 0$

$$\delta\omega = \begin{cases} h\eta & : \text{relaxed cell} \\ 0 & : \text{otherwise} \end{cases}$$



$\eta$  chosen to satisfy  $U_x + V_y = f_3$

$$\text{Convection Diffusion } Q = -\frac{1}{h} \Delta + \sum u_i \partial_i$$

Large  $hR$ :  $Q^h = -\sum \beta_i h \partial_{ii}^h + \sum u_i \partial_i^{2h}$   
AV (artificial viscosity)

For stability (i.e., in relaxation)  $\beta_i = O(|u_i|)$   
i.e., stream-wise AV is  $O(h|u|)$ . semi  $h$ -elliptic.

Intended alignment: for good approximation  
of cross-stream  $h$ -f or discontinuities:  
(1) grid aligned with the flow throughout  
(2)  $Q^h$  that avoid straddling the discontinuity

Relaxation: red-black Gauss-Seidel  $\frac{\beta_i}{|u_i|} > .5$

Block relaxation: only for the corresponding  
intended alignment, hence simple.

Slow smoothing in regions of accidental alignment  
is for "characteristic" components, which anyway  
are not approximated at all in other regions

Residuals transfer:  $\beta_i = 0 \Rightarrow O(h^2)$ .  
Downstream bias • Avoid straddling discontinuities

Interpolation: Avoid straddling discontinuities

