



INTERNATIONAL ATOMIC ENERGY AGENCY
UNITED NATIONS EDUCATIONAL, SCIENTIFIC AND CULTURAL ORGANIZATION
INTERNATIONAL CENTRE FOR THEORETICAL PHYSICS
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**WINTER COLLEGE ON 'MULTILEVEL TECHNIQUES IN
COMPUTATIONAL PHYSICS'**

**Physics and Computations with Multiple Scales of Lengths
(21 January - 1 February 1991)**

H4.SMR 539/17

Computer Simulation of the Liquid-Vapour Phase Transition

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DIPARTIMENTO FISICA TEORICA
TRIESTE

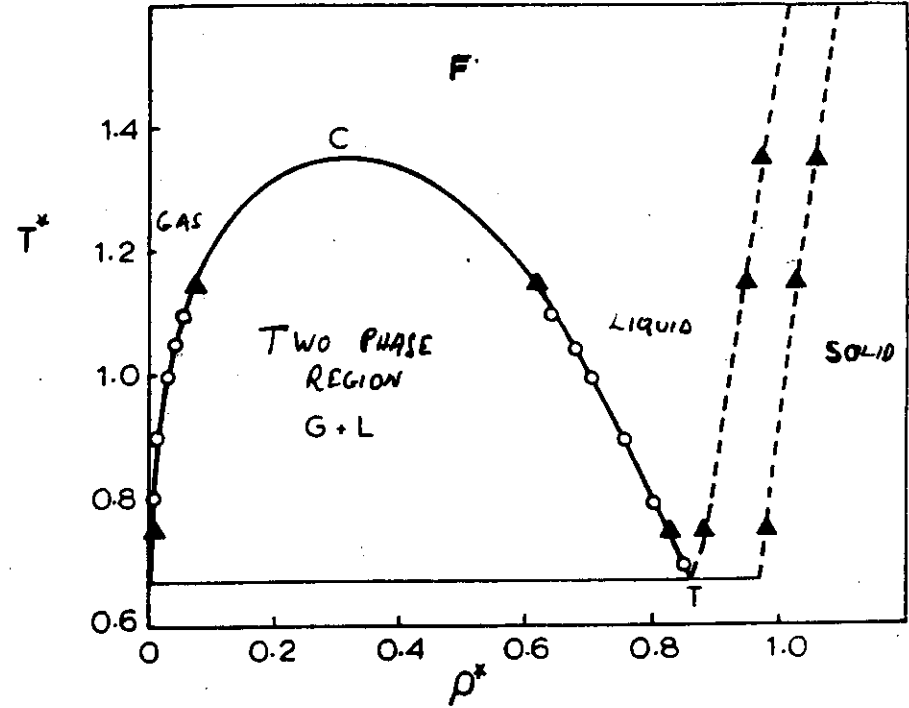
COMPUTER SIMULATION OF
THE LIQUID-VAPOUR PHASE TRANSITION

LIQUID-VAPOUR PHASE TRANSITION

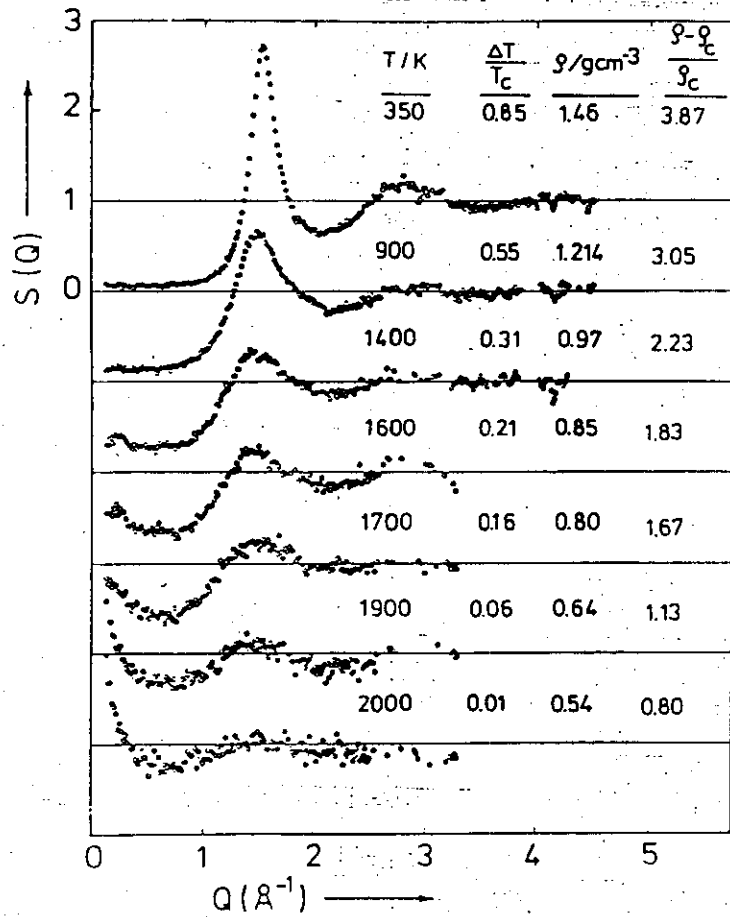
MODEL FOR FLUID Argon

OPEN PROBLEMS:

- How critical parameters (P_c, T_c, ρ_c) are related to the microscopic forces?
- cross-over normal to critical
- dynamical aspects
- metastability in the two phase region
- phase separation in fluid mixtures



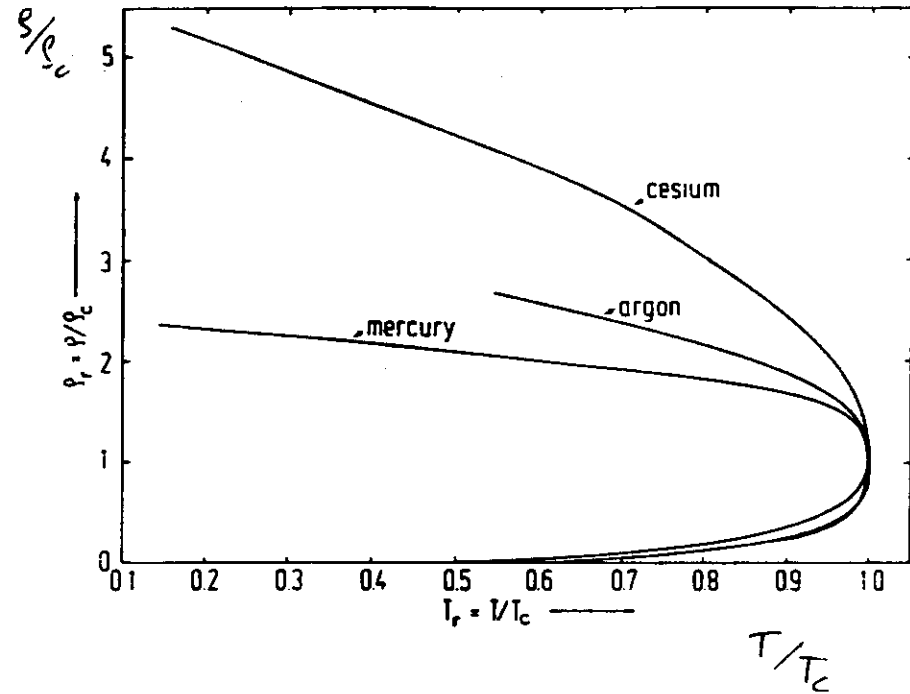
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Static structure factor of expanded liquid Rb determined by neutron diffraction (Franz *et al.* 1980).

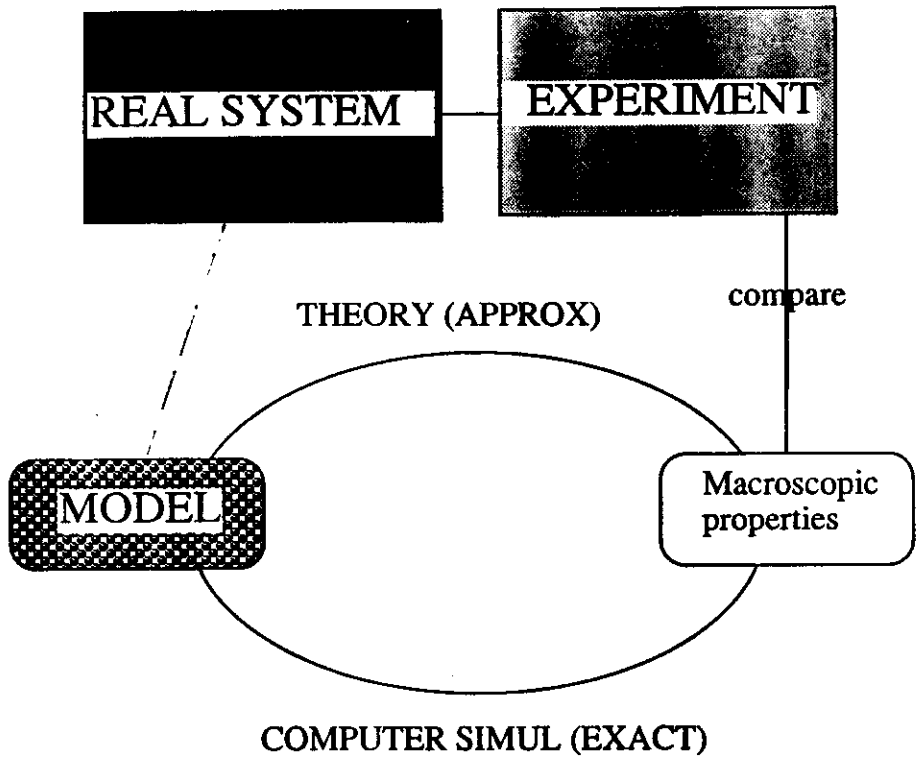
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COEXISTENCE CURVES



From: Freyland - Hensel

SIMULATION OF PHASE TRANSITIONS



PROBLEMS:

1. Finite size
2. Finite time

Critical properties are related to
 non analyticity of the free energy
 in the therm. limit

$$F(T, V, N) \quad \left. \begin{matrix} N \\ V \end{matrix} \right\} \longrightarrow \infty, \quad \rho = \text{const}$$

Finite size analysis for simulation
 of lattice models

(Ma Swendsen Binder)

Extension to
 "realistic fluids"

IN COLLABORATION
 K. BINDER , D. HEERMANN (MAINZ)

SIMULATION OF THE GAS LIQUID TRANSITION

- 1- Introduction and Motivations
- 2- MODEL: 2D-Lennard-Jones
- 3- BLOCK DISTRIBUTION FUNCTION (BDF)
- 4- RESULTS IN THE TWO PHASE REGION
- 5- MOMENTS OF THE BDF
- 6- DETERMINATION OF THE CRITICAL POINT
- 7- DISCUSSION AND FUTURE WORK

$$C_v \sim \langle \Delta E \Delta E \rangle$$

$$K_T \sim \langle \Delta p \Delta p \rangle$$

$$\chi \sim \langle \Delta M \Delta M \rangle$$

$$\langle \dots \rangle \sim \exp(-r/\xi)$$

$$\xi \quad T \gg T_c$$

$$\xi \rightarrow \infty \quad T \rightarrow T_c$$

$$\xi \sim |T - T_c|^{-\nu}$$

$$C_v \sim |T - T_c|^{-\alpha}$$

Liquid-Gas

$$K \sim |T - T_c|^{-\gamma}$$

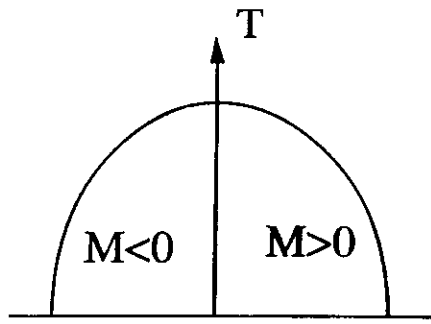
$$C_v \sim |T - T_c|^{-\alpha}$$

Magnet

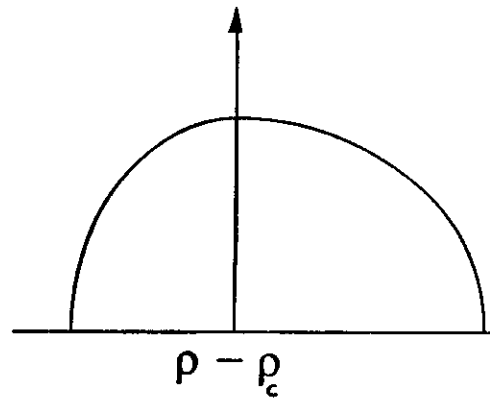
$$\chi \sim |T - T_c|^{-\gamma}$$

	α	γ
liquid-gas	0.08-0.10	1.2
magnet (uniax.)	0.08-0.10	1.15
3D Ising	0.12	1.25
2D Ising	(0)	1.75

ISING



LIQUID-GAS



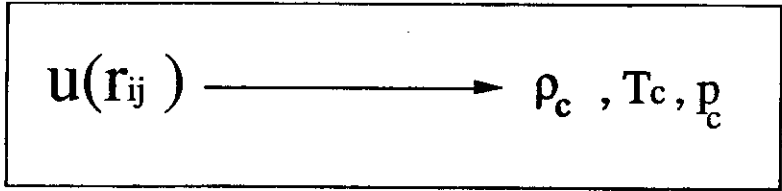
⇒ Traditional theories of liquids are not adequate

⇒ Attempts to combine with R.G.

$T=T_c$: universal properties like Ising model

no information about not univ. prop.

Hubbard and Schofield
Van Dieren and Van Leeuwen



Parola, Meroni and Reatto

$T \rightarrow T_c$

influence of the realistic potential

(13)

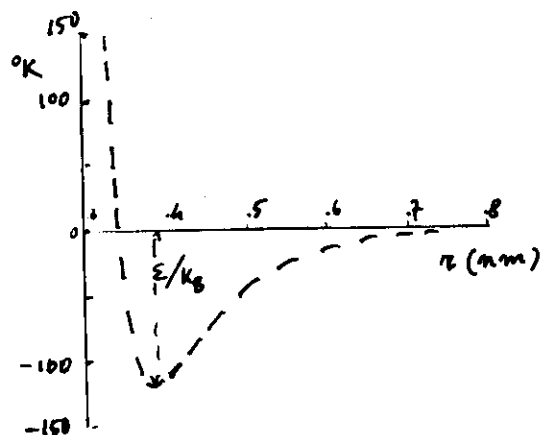
SIMULATION OF "REAL" CLASSICAL FLUIDS

$$H = \sum_i \frac{p_i^2}{2m} + \underline{V(r_1, \dots, r_N)}$$

$$V = \sum_{i>j} \phi(r_{ij}) + \dots \approx \sum_{i>j} \phi^{eff}(r_{ij})$$

EXAMPLE: LENNARD-JONES POTENTIAL

$$\phi^{LJ}(r) = 4\epsilon \left[\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^6 \right]$$



Argon: $\sigma = 0.34 \text{ nm}$
 $\epsilon/k_B \approx 120 \text{ K}$

(14)

COMPUTE AVERAGES IN THE ENSEMBLE

$$\langle A \rangle = \frac{\int dt A e^{-\beta H}}{\int dt e^{-\beta H}} = \quad (1)$$

$$= \lim_{t_0 \rightarrow \infty} \frac{1}{t_0} \int_0^{t_0} A(t) dt \quad (2)$$

(1) M.C. USUALLY CANONICAL ENP.

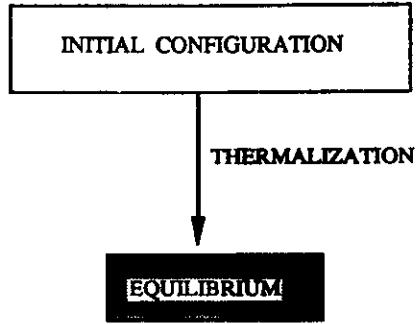
(2) M.D. USUALLY MICROCANONICAL

average temperature

$$\frac{f}{2} N k_B \langle T \rangle = \langle K \rangle$$

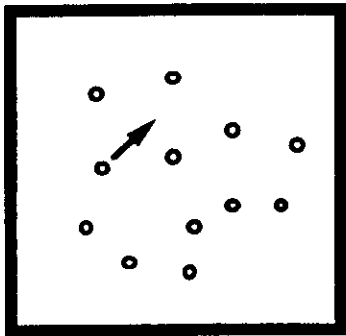
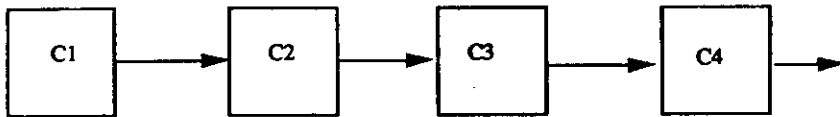
$$\begin{aligned} \dot{x}_i &= p_i \\ \dot{p}_i &= -\frac{\partial V}{\partial x_i} \end{aligned}$$

COMPUTER SIMULATION



M.C. SIMULATION

AVERAGES ON DIFFERENT CONFIGURATIONS



$$\mathbf{r}_i \rightarrow \mathbf{r}_i + \delta$$

$$\Delta E = E_{\text{new}} - E_{\text{old}}$$

Computer simulation near T_c

Indirect methods

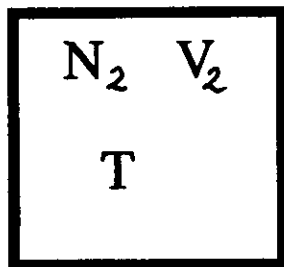
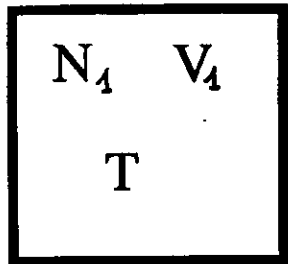
⇒ thermodynamic integration
(free energy)

⇒ umbrella sampling

⇒ particle insertion
(chemical potential)

GIBBS ENSEMBLE (M.C.)

M.C. simulation of 2D L.J. fluid



Potential

$$\varphi(r) = 4\epsilon \left\{ \left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^6 \right\}$$

$$\Delta V_1$$

$$\Delta V_2 = -\Delta V_1$$

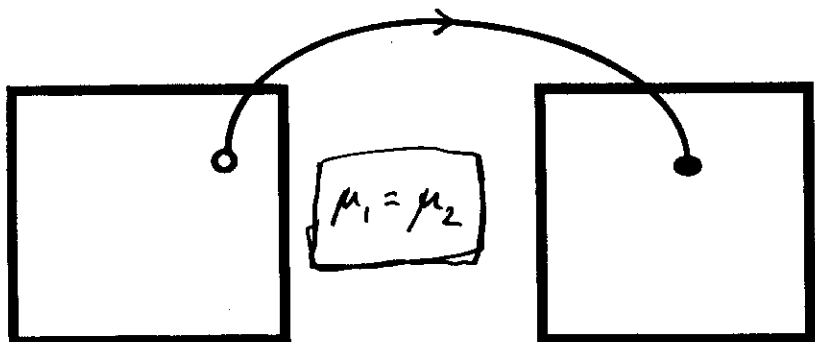
$$P_1 = P_2$$

canonical M.C. with P.B.C.

Units:

$$T^* = \frac{k_B T}{\epsilon}$$

$$\rho^* = \rho \sigma^2$$



EXCHANGE PARTICLES

Estimated critical point

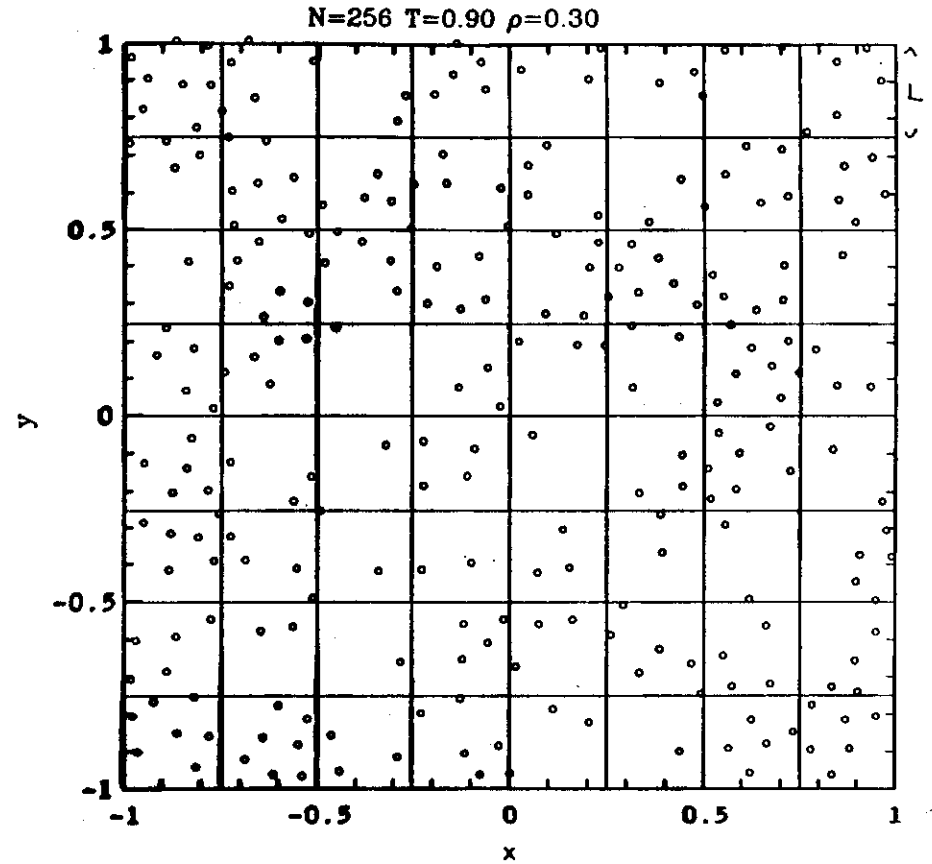
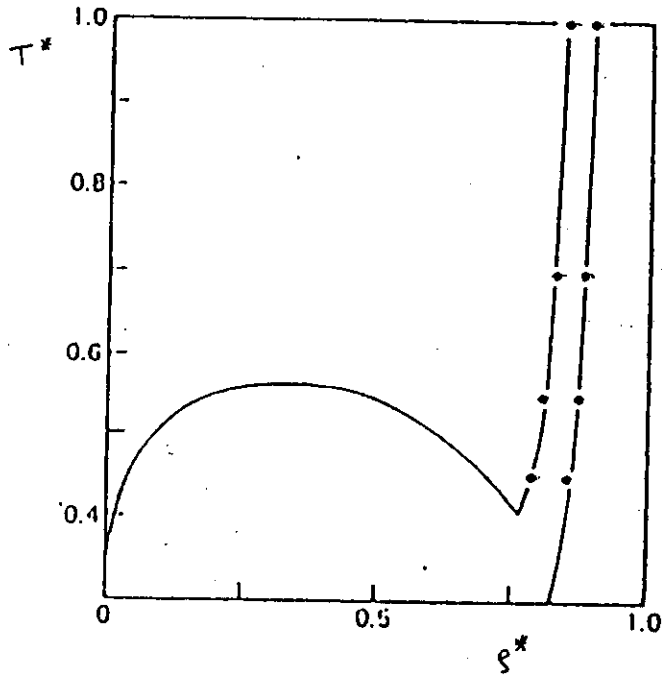
(2 D L.J.)

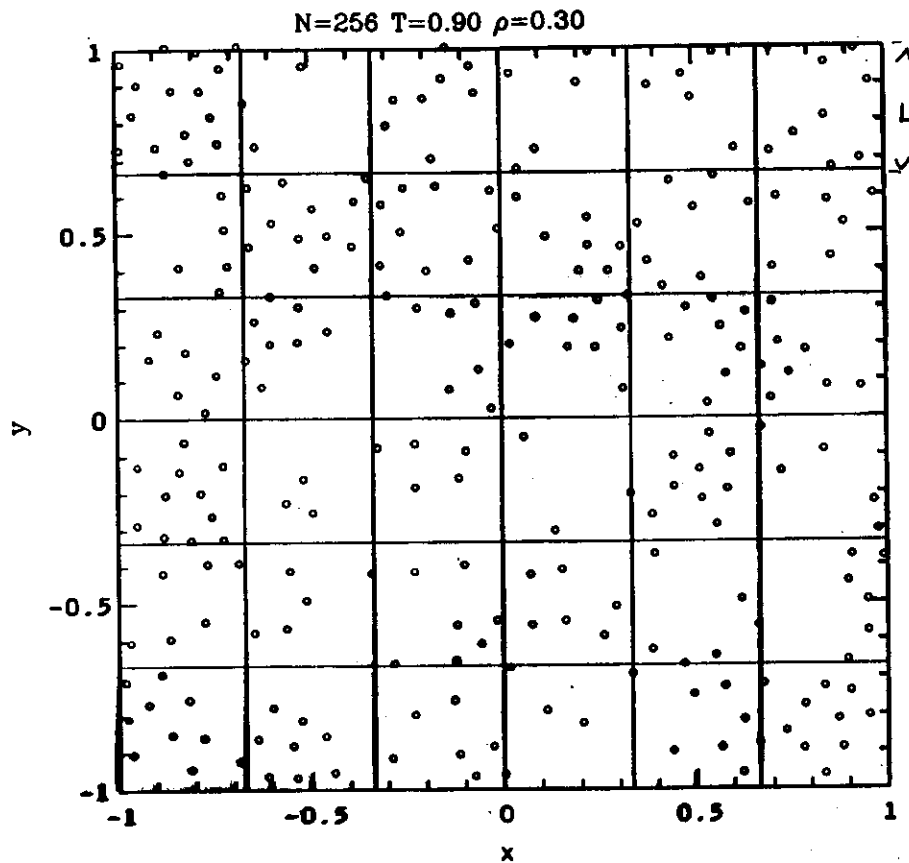
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(20)

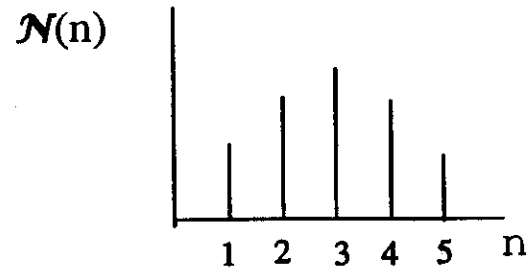
T_c^*	ρ_c^*
0.62 ÷ 0.7	0.38
0.56	0.32
0.53	0.33

Tsien & Valleau
Henderson (77)
Barker & al. (81)





Block distribution functions



$$\rho = n/L^2 \longrightarrow P_L(\rho)$$

$$T \gg T_c \quad \xi \ll L$$

$$P_L(\rho) \sim \exp \left\{ -(\rho - \langle \rho \rangle)^2 / \alpha^2 \right\}$$

$$\alpha^2 \sim k_B T \frac{K_T^L}{L^{-d}}$$

isothermal compress.

(in the therm. limit
 $K_T^L \rightarrow K_T$)

BLOCK DISTRIBUTION FUNCTIONS

length of the block $L = S/M$

= 2 M^2 BLOCKS

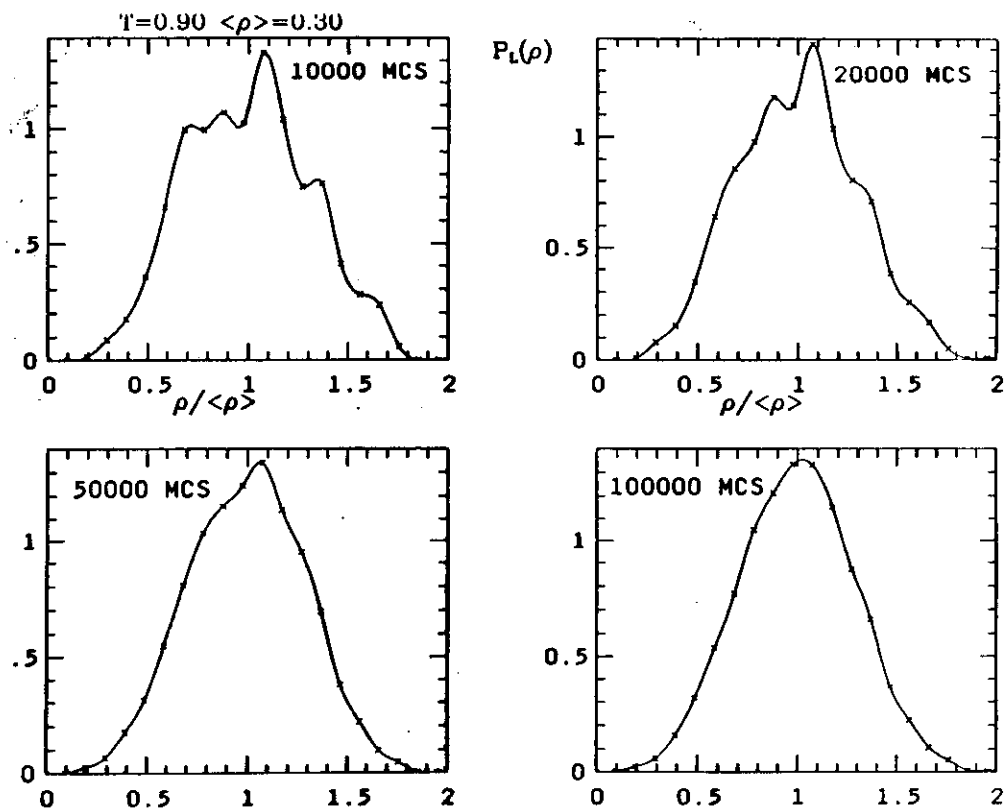
$\mathcal{N}_b(n_b)$ n. of blocks with n_b particles

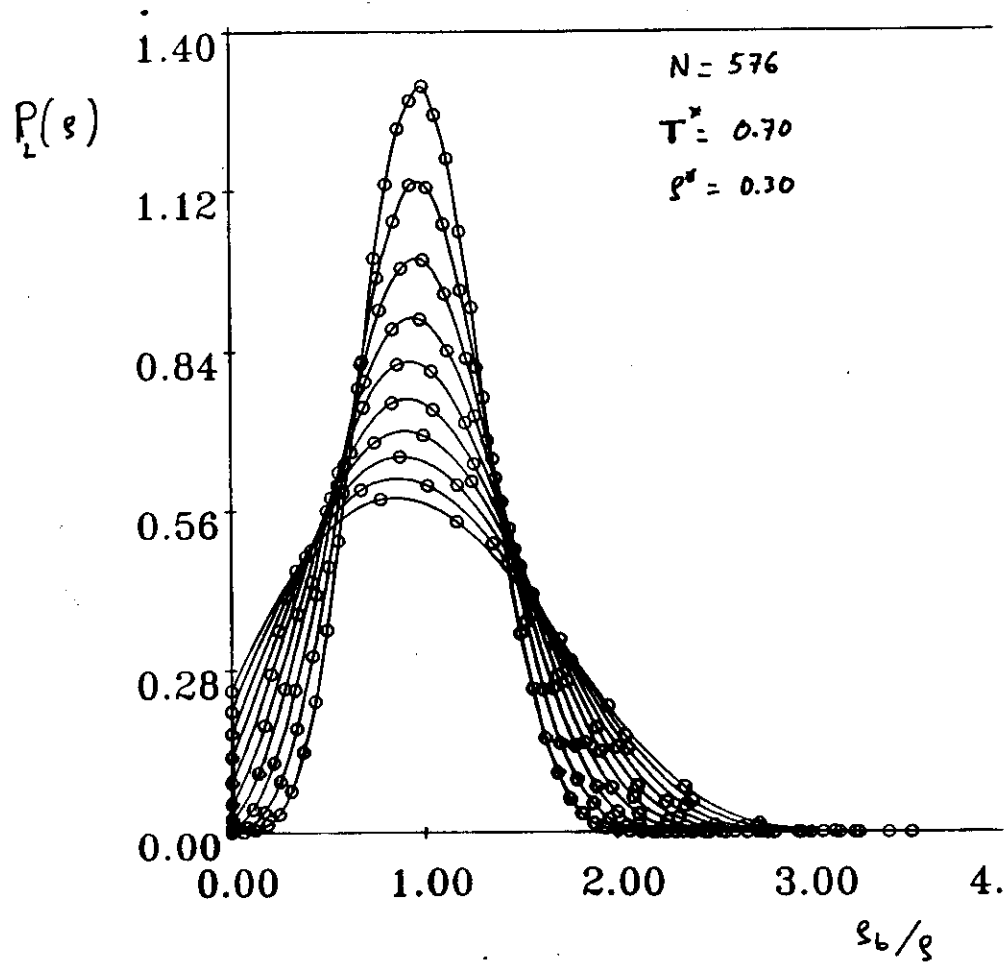
SUM RULES

$$\sum_{n_b} \mathcal{N}_b(n_b) = M^2$$

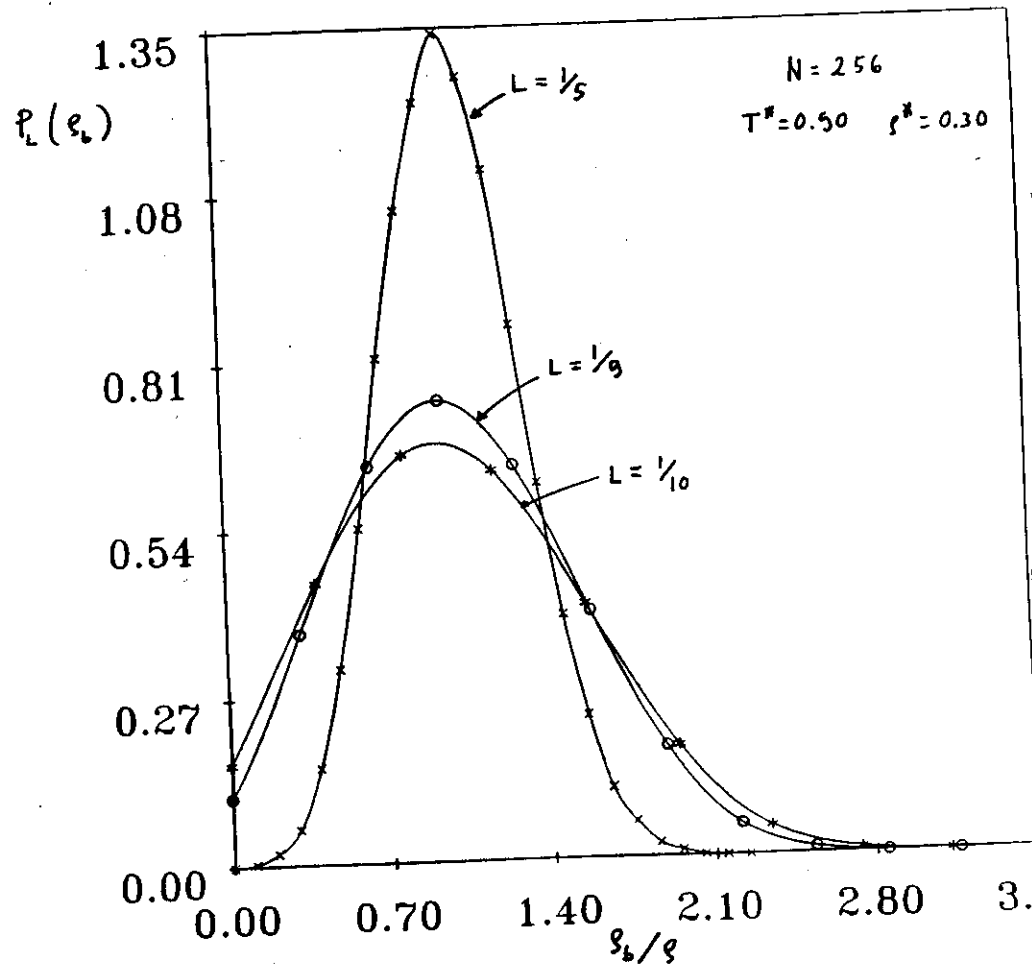
$$\sum_{n_b} \mathcal{N}_b(n_b) \cdot n_b = N$$

$$n_b \rightarrow \rho_b = n_b / L^2 : \rho_2(s)$$

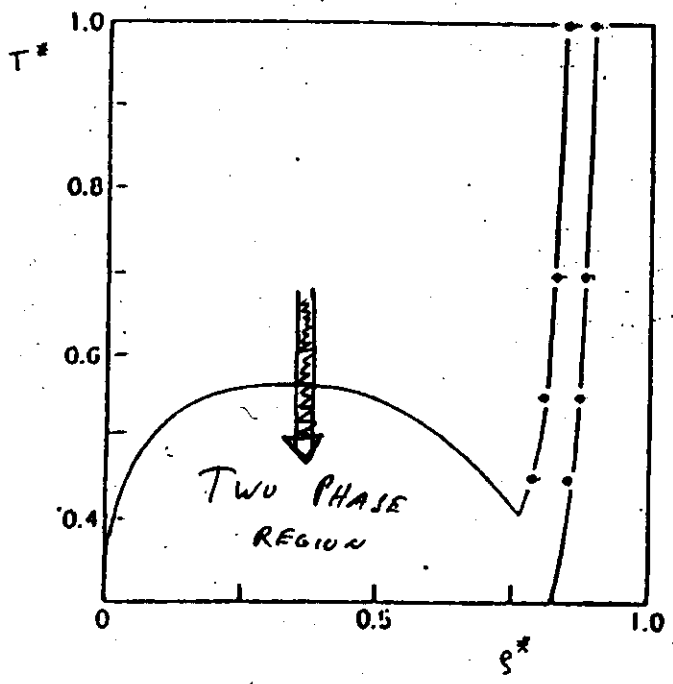




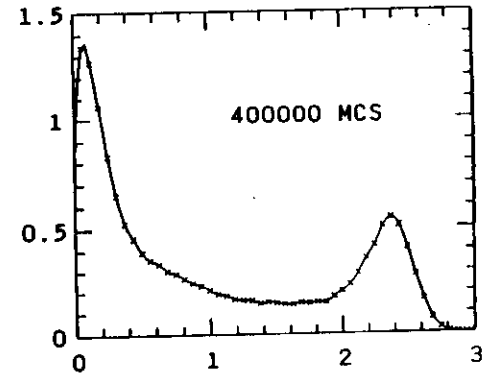
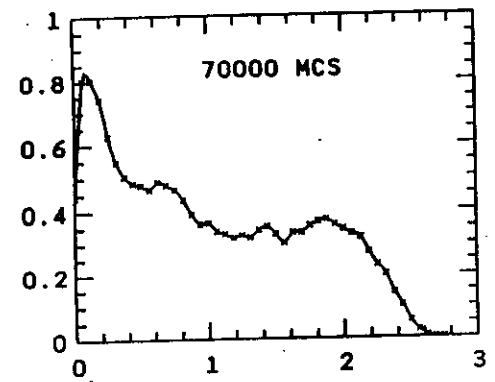
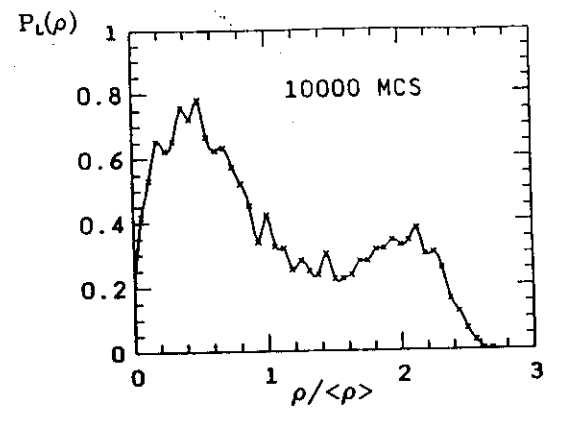
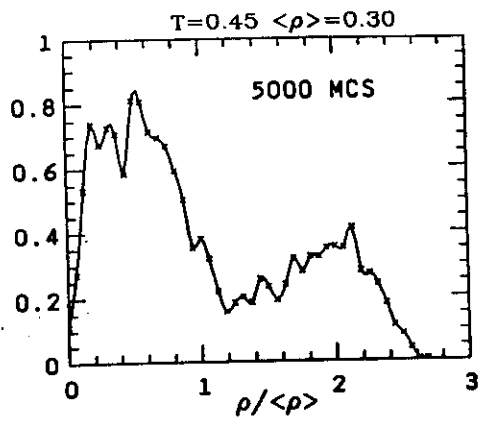
$T_c \approx 0.55 - 0.50$

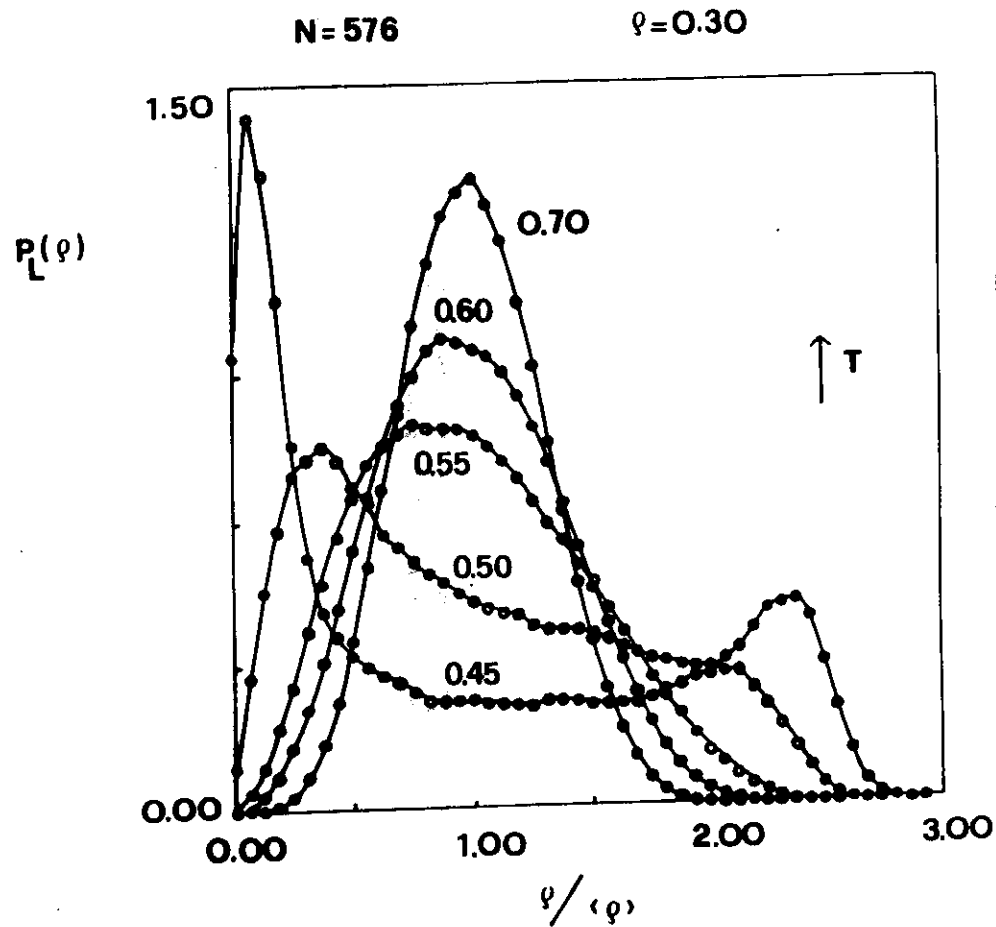
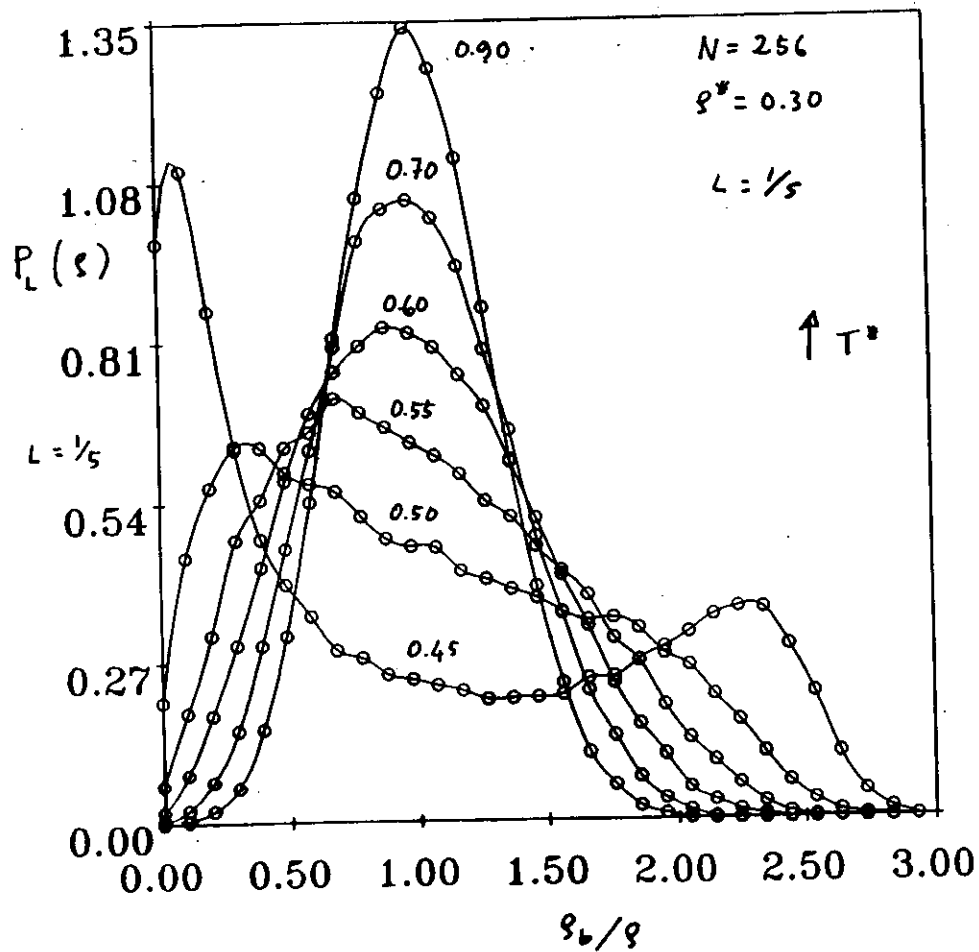


(27)

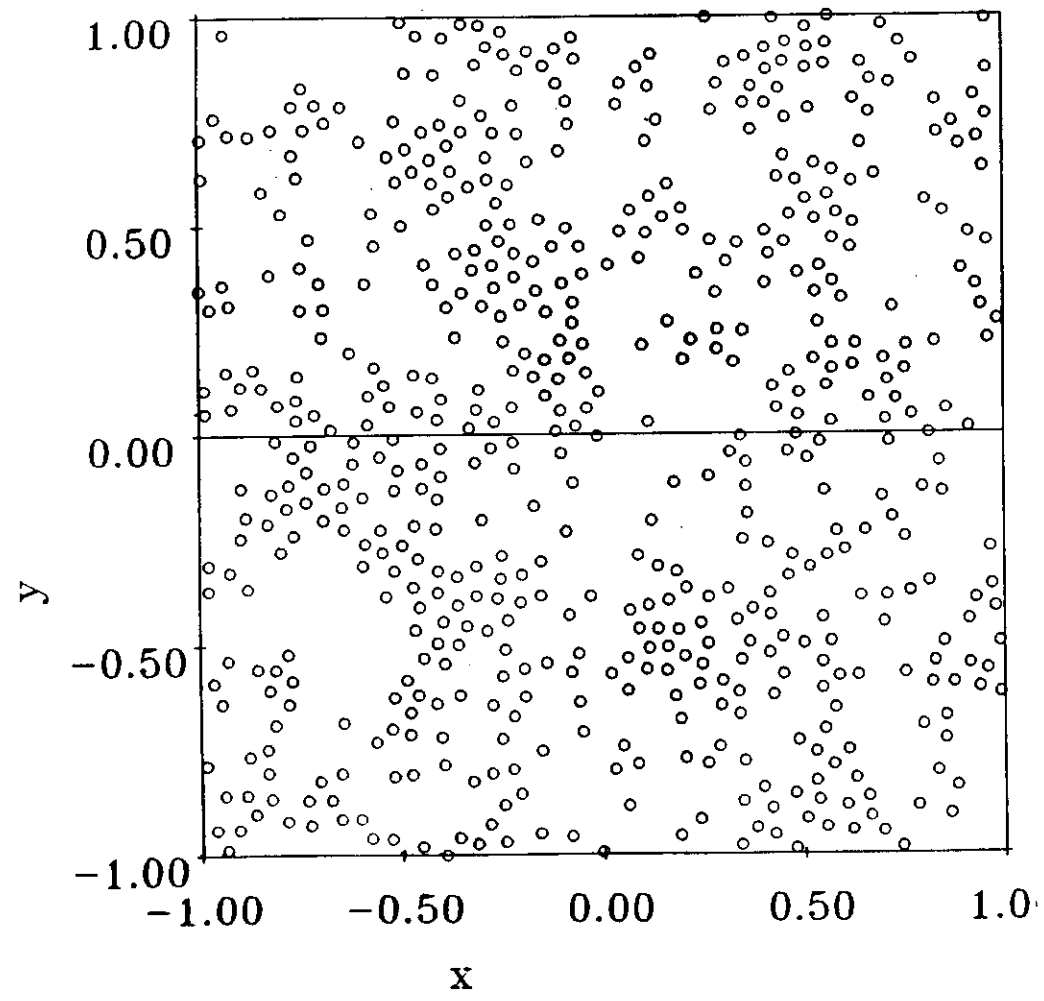


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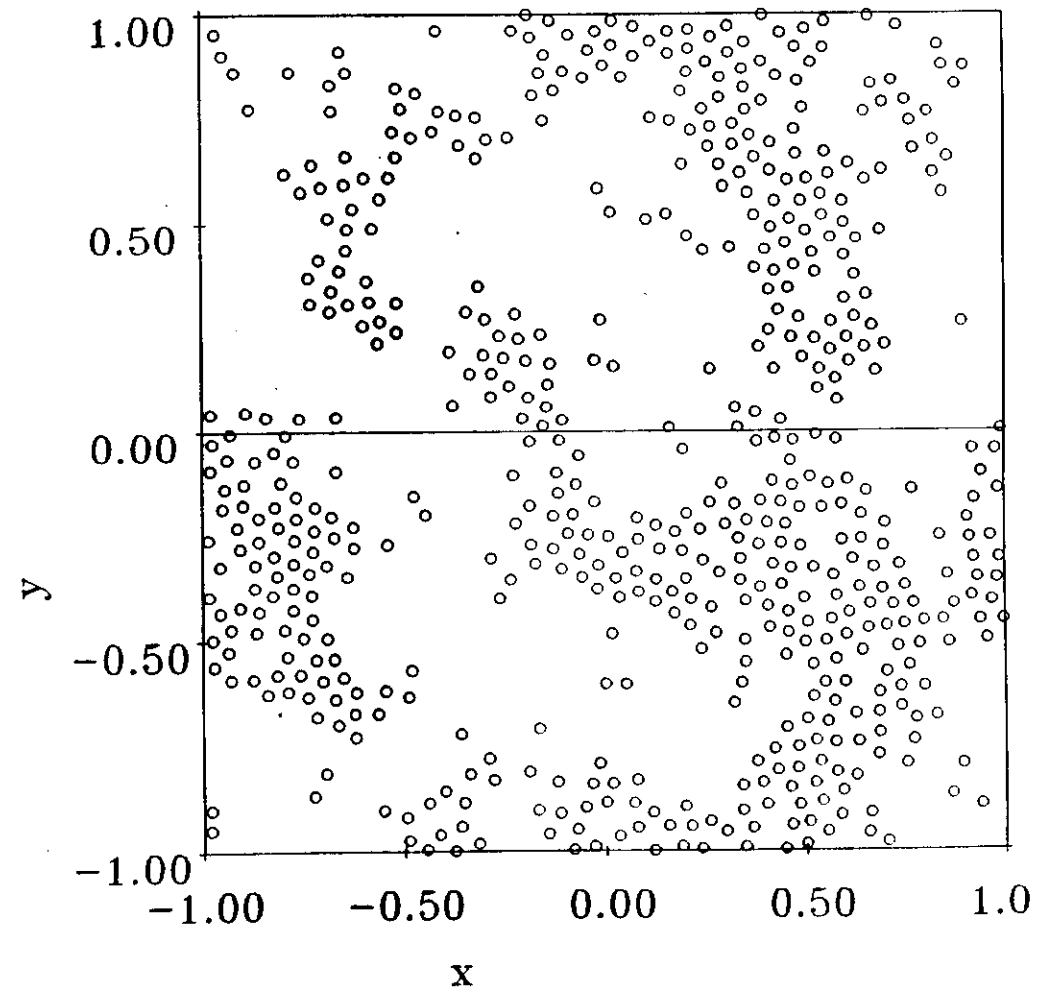




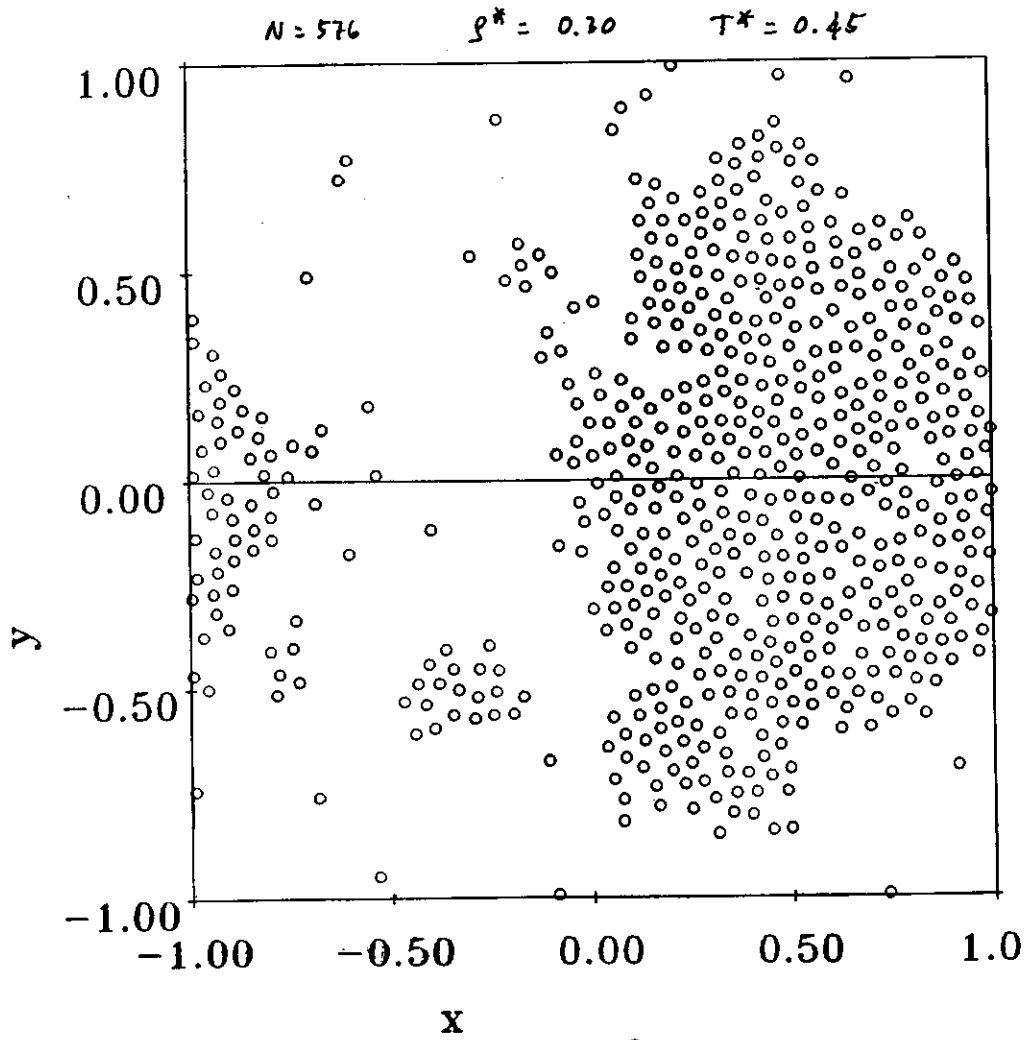
$N = 576$ $\beta^* = 0.30$ $T^* = 0.70$



$N = 576$ $\beta^* = 0.30$ $T^* = 0.50$

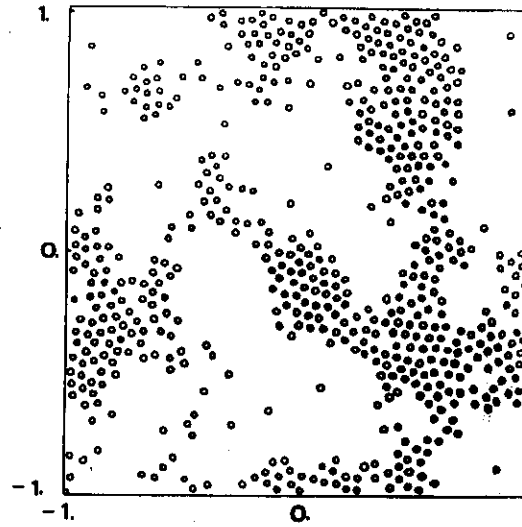


$\sim 10^6$ MCS

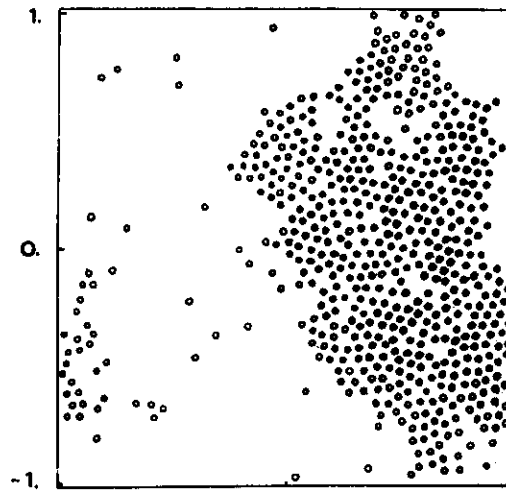


$\sim 10^6$ MCS

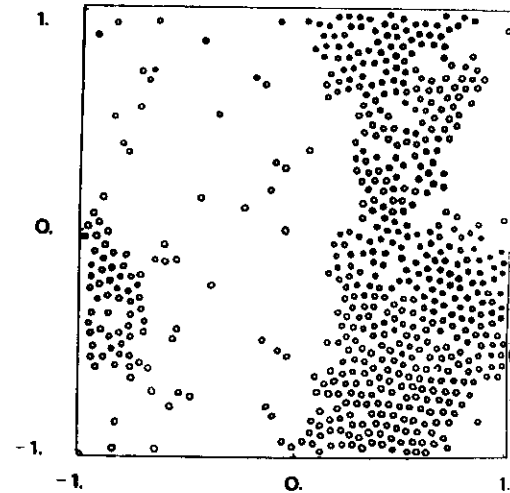
$N=576$ $T=0.45$ $\rho=0.30$ $MCS=10,000$



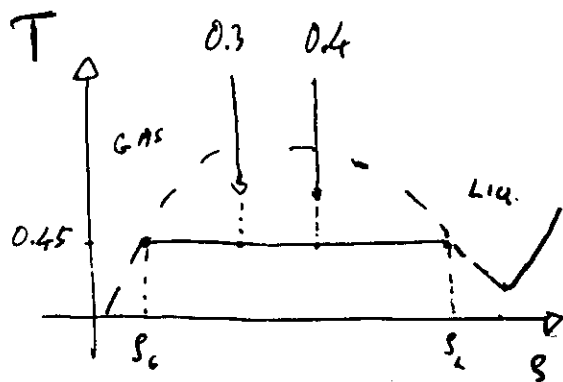
$N=576$ $T=0.45$ $\rho=0.30$ $MCS=900,000$



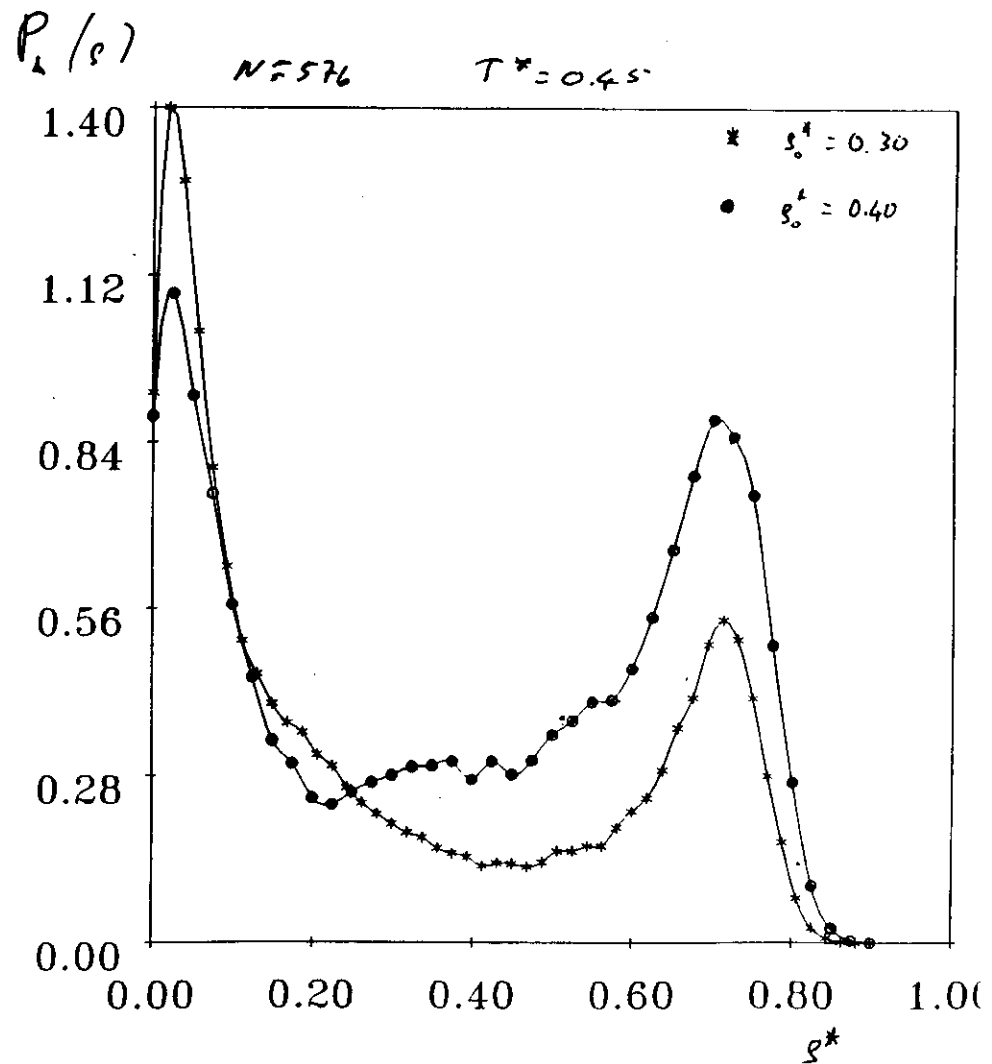
$N=576$ $T=0.45$ $\rho=0.30$ $MCS=400$

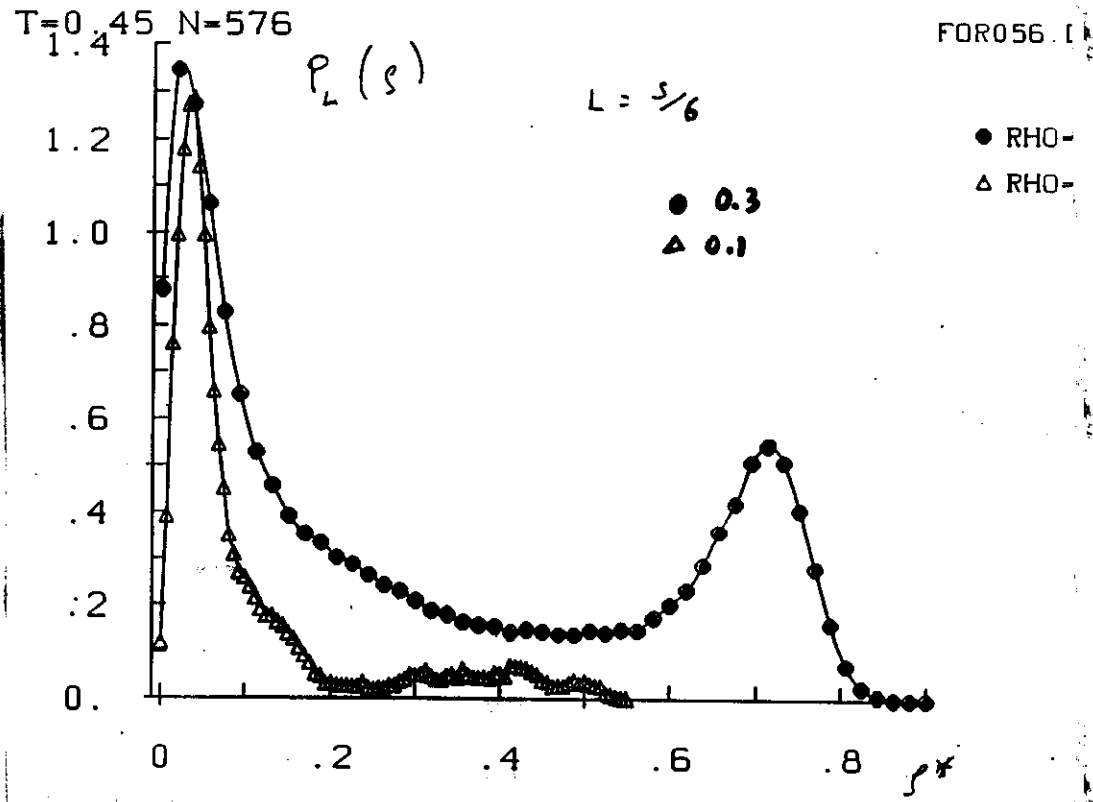
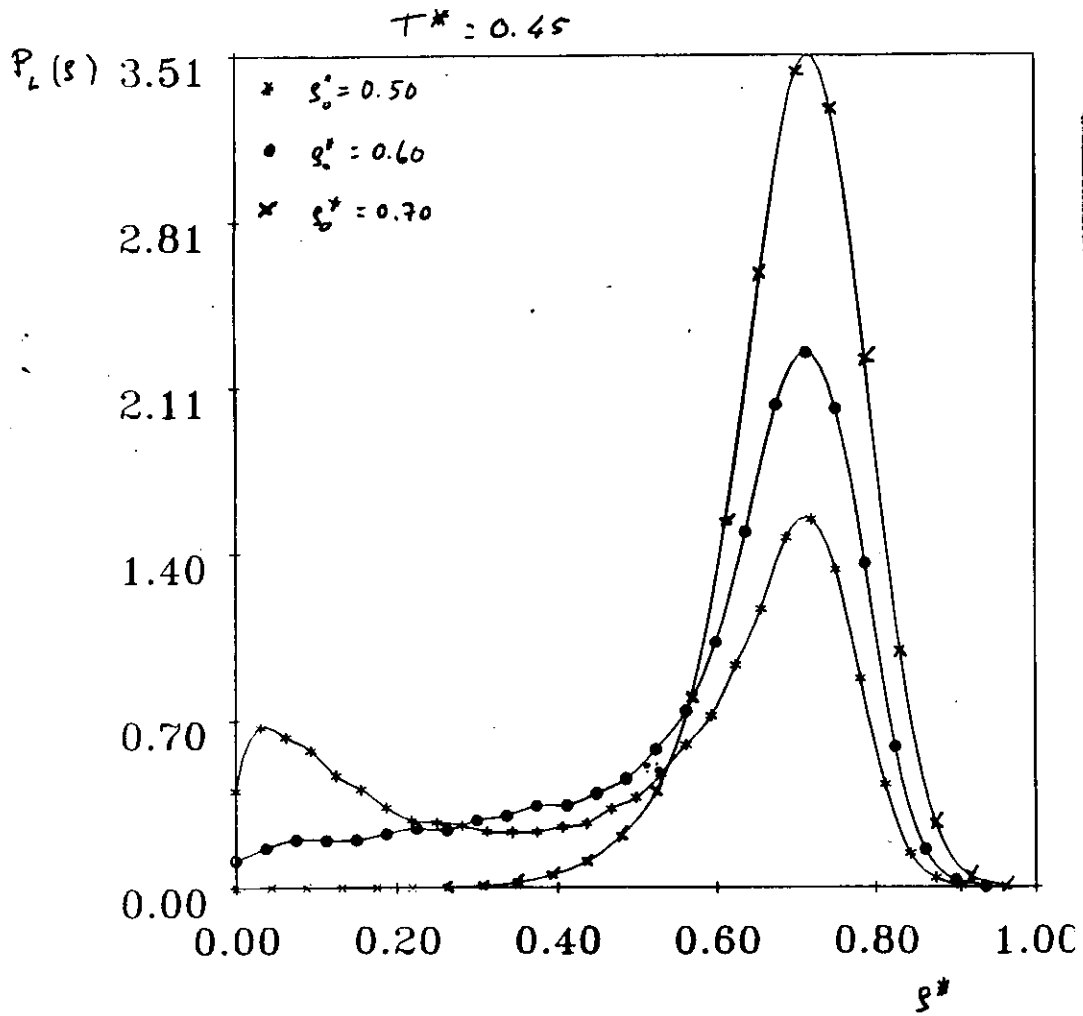


(35)



(36)





MOMENTS

$$\langle s^n \rangle_L = \int ds P_L(s) s^n$$

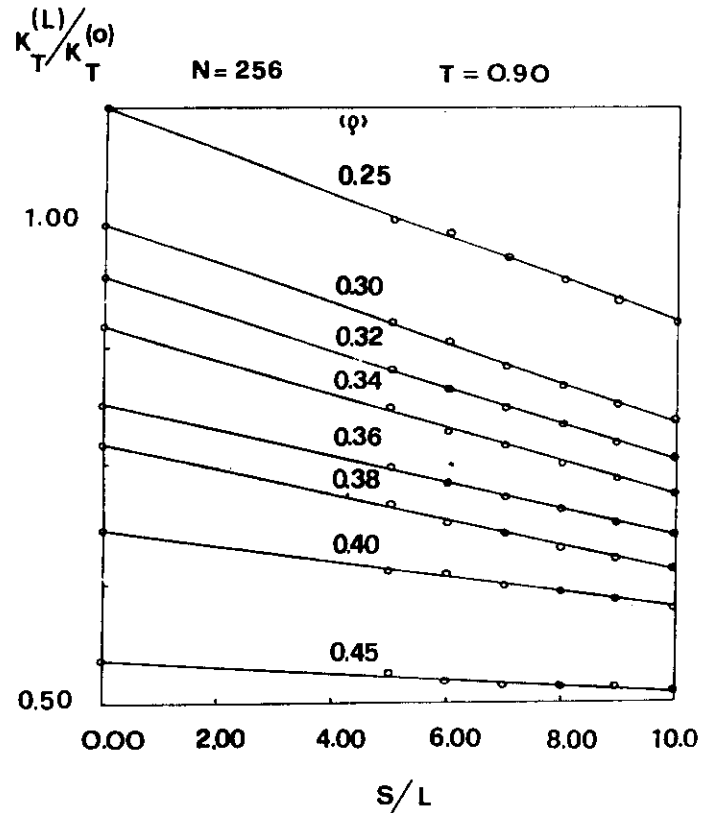
1-st $\langle s \rangle_L = \text{av. density of the block}$

2-nd is related to the isothermal compressibility

$$L^{-d} \langle s^2 \rangle_L - k_B T \frac{K_T^{(L)}}{K_T} = \langle (s - \langle s \rangle_L)^2 \rangle_L$$

$T \gg T_c$

$$K_T^{(L)} \sim K_T + c L^{-1}$$

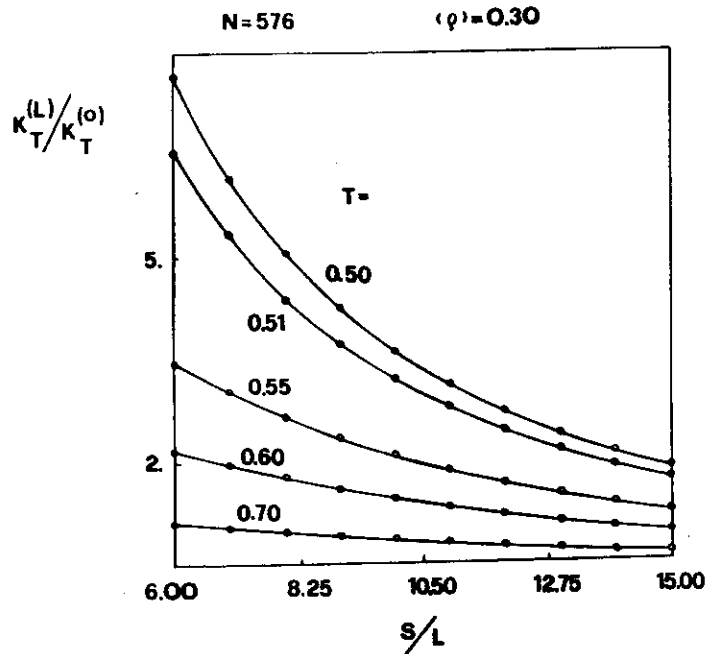


(4)

(4)

For a state in the two phase region

$$T \ll T_c \quad L \gg \xi$$



$$P_L(\rho) \propto A \frac{1}{\rho_{gas} (K_T^{gas})^{1/2}} \exp\left[-\frac{(\rho - \rho_{gas})^2 L^d}{2\rho_{gas}^2 k_B T K_T^{gas}}\right] +$$

$$B \frac{1}{\rho_{liq} (K_T^{liq})^{1/2}} \exp\left[-\frac{(\rho - \rho_{liq})^2 L^d}{2\rho_{liq}^2 k_B T K_T^{liq}}\right]$$

$$A(\langle \rho \rangle) \quad \text{and} \quad B(\langle \rho \rangle)$$

determined from $\mu_{gas} = \mu_{liq}$

(43)

(46)

$N = 576$ $T = 0.45$ $\langle \rho \rangle = 0.30$ COEX.

$\bar{\rho} = 0.45$ $RHO = 0.30$ $N = 576$

FOR056.DA

liquid branch

$P_L(\rho)$

$L = S/M$

○ M-15
△ M-10
+ M-6

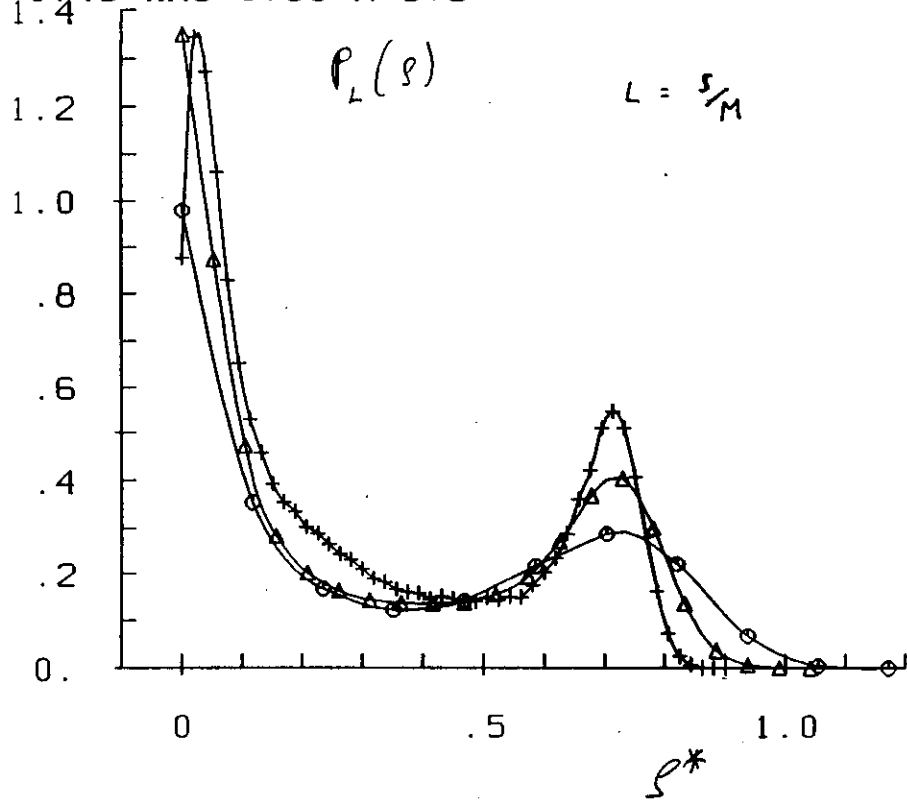


TABLE I

M_b	ρ_{liq}	$k_B T \rho_{liq} K_T^{liq}$
6	0.711	0.867
7	0.718	0.763
8	0.720	0.730
9	0.718	0.814
10	0.719	0.753
11	0.718	0.875
12	0.698	1.169
13	0.715	0.828
14	0.722	0.830
15	0.715	1.057

(45)

SOMETHING ABOUT BLOCKS

RESOLUTION IN DENSITY

$$\Delta p = \frac{1}{L^d} \quad L = \frac{S}{M}$$

LARGE $L \rightarrow$ GOOD RESOLUTION

BUT

$$\frac{\Delta m_b}{m_b} \text{ must be not too small}$$

OPT.

$$L_{\min} < L < L_{\max} \ll N^{1/d}$$

$$L_{\min} \text{ large enough} \quad \frac{\Delta p}{S} \approx 0.1$$

(46)

$$L = \frac{S}{M}$$

$$d=2$$

take M too LARGE $M^2 \gg N$

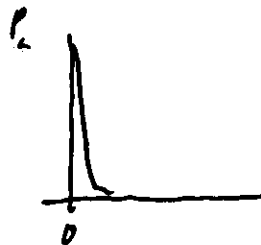
$$m_b = 0, 1$$

FROM SUM RULES

$$\mathcal{N}_b(0) \cdot 0 + \mathcal{N}_b(1) \cdot 1 \approx N$$

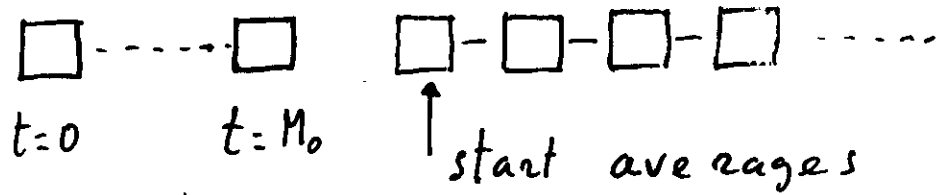
$$\mathcal{N}_b(0) + \mathcal{N}_b(1) \approx M^2$$

$$\mathcal{N}_b(0) \approx M^2$$



(47A)

Relaxation time to equilibrium



$$\tau_A^{rel} \ll M_0/N$$

→ metastable states



$$\tau_A^{rel} \approx \tau_0 \exp\left[\frac{\Delta U}{k_B T}\right]$$

(47B)

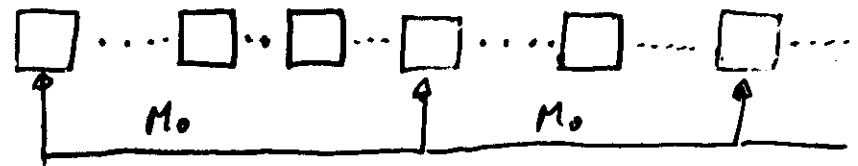
→ close to T_c

$$\Delta T = T - T_c$$

$$\tau_A^{rel} \sim |\Delta T|^{-\Delta_A(\Delta T)}$$

$$\chi_{AA} \sim \langle A^2 \rangle - \langle A \rangle^2 \rightarrow \tau_A^{rel} \sim \chi_{AA}$$

Average on independent configurations:



$$\tau_{AA} \ll (M - M_0)/N$$

4-th order reduced cumulant

$$\Delta \rho = \rho - \langle \rho \rangle$$

$$U_L(T, \rho) = 1 - \frac{\langle \Delta \rho^4 \rangle_L}{3 \langle \Delta \rho^2 \rangle_L^2}$$

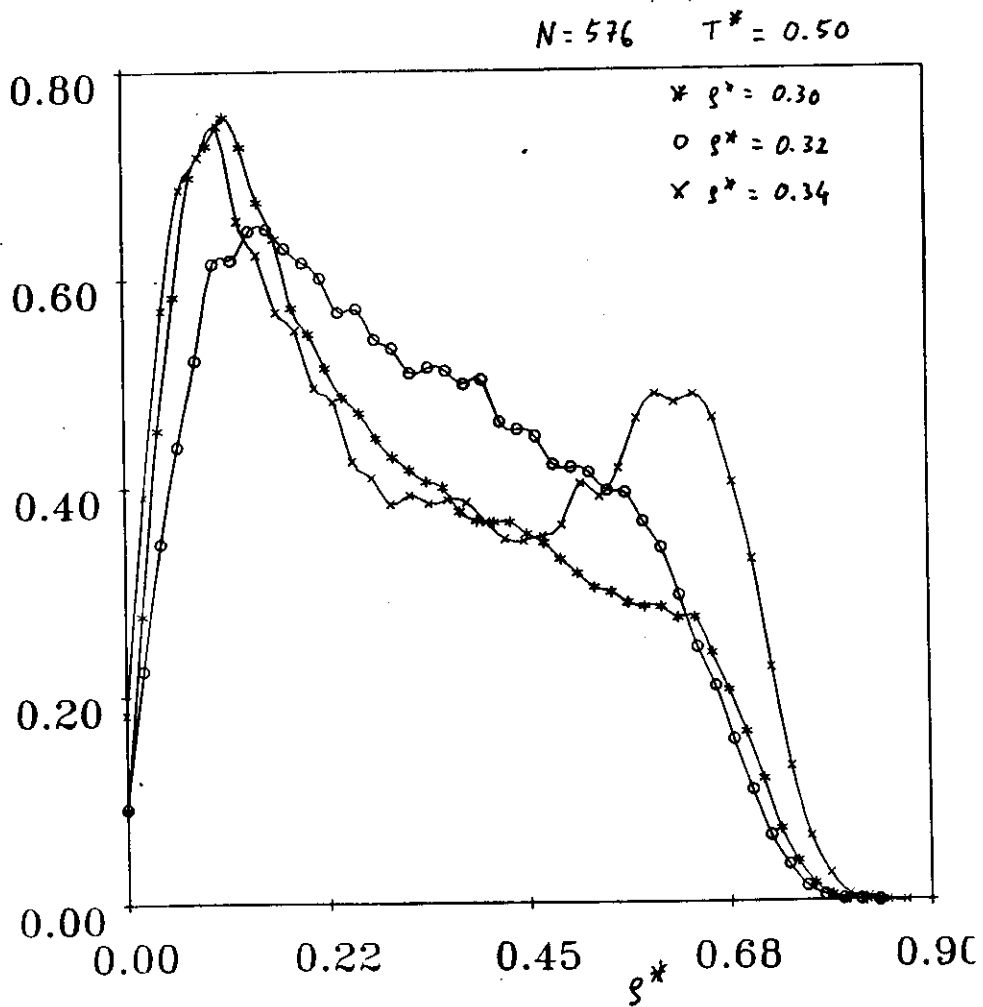
Using:

$$U_L(T) = 1 - \frac{\langle s^4 \rangle}{3 \langle s^2 \rangle^2}$$

$$T \gg T_c \quad U_L(T) \sim L^{-d} \rightarrow 0$$

$$T \ll T_c \quad U_L(T) \rightarrow 2/3 = 0.6666\dots$$

$$T \rightarrow T_c \quad U_L(T) \rightarrow U^*$$

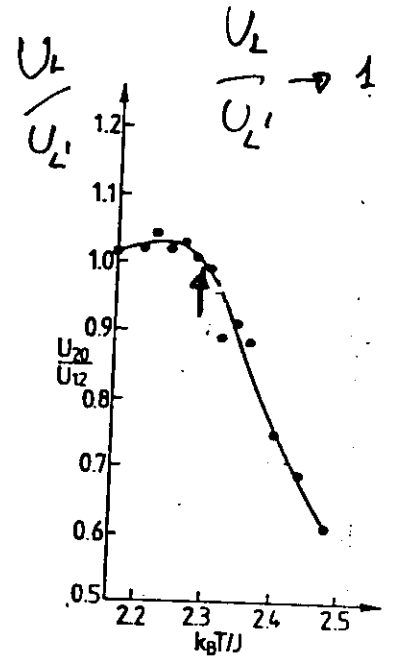
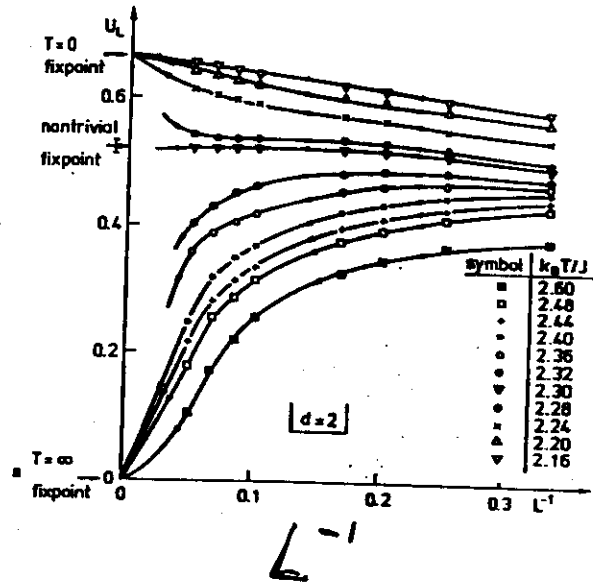
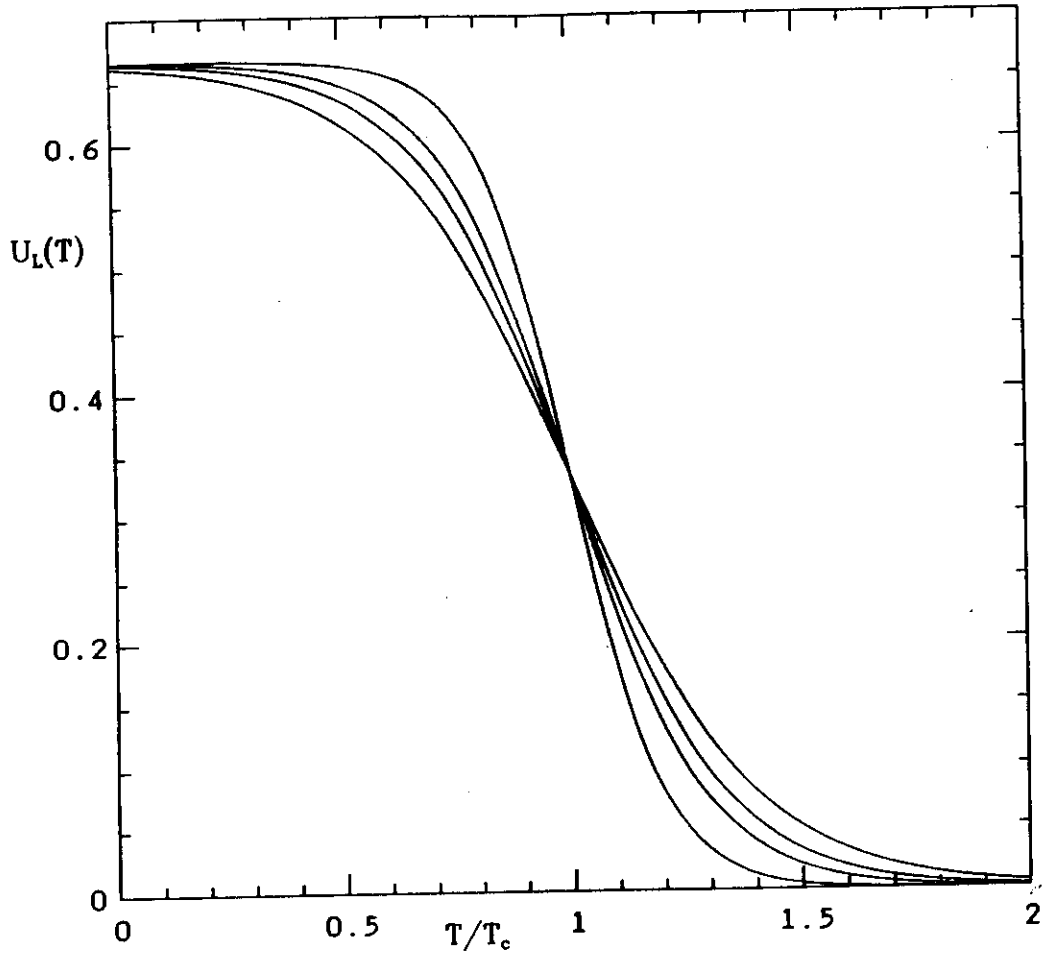


(40)

$$U(L,T) = 1 - \frac{\langle S^4 \rangle_L}{3 \langle S^2 \rangle_L^2}$$

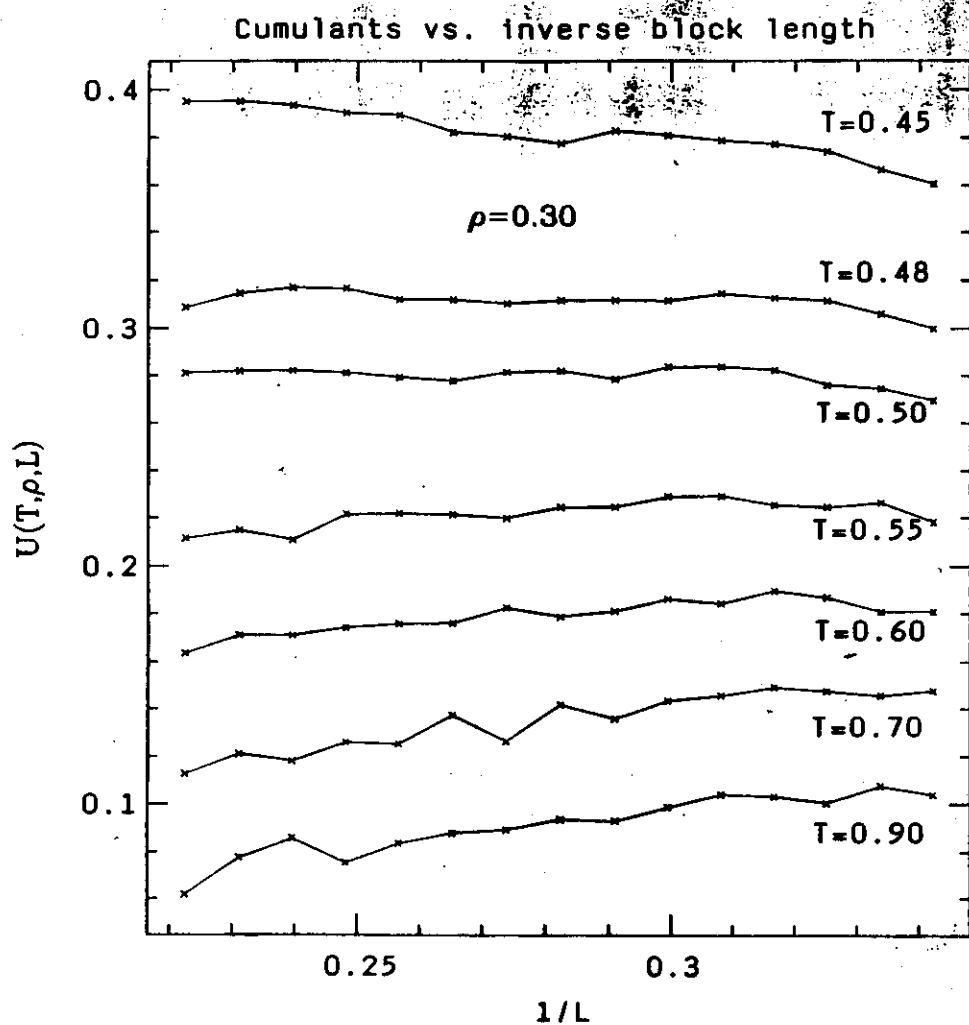
$$T \rightarrow \infty \quad U \rightarrow 0$$
$$T \rightarrow 0 \quad U = \frac{4}{3}$$

CUMULANTS FOR DIFFERENT L



$$\frac{k_B T_c}{J} = 2.269$$

$$P_L(\rho) = L^\nu \bar{P}\{(\rho - \rho_c)L^\nu, (p - p_c)L^\nu, L/\xi\}$$



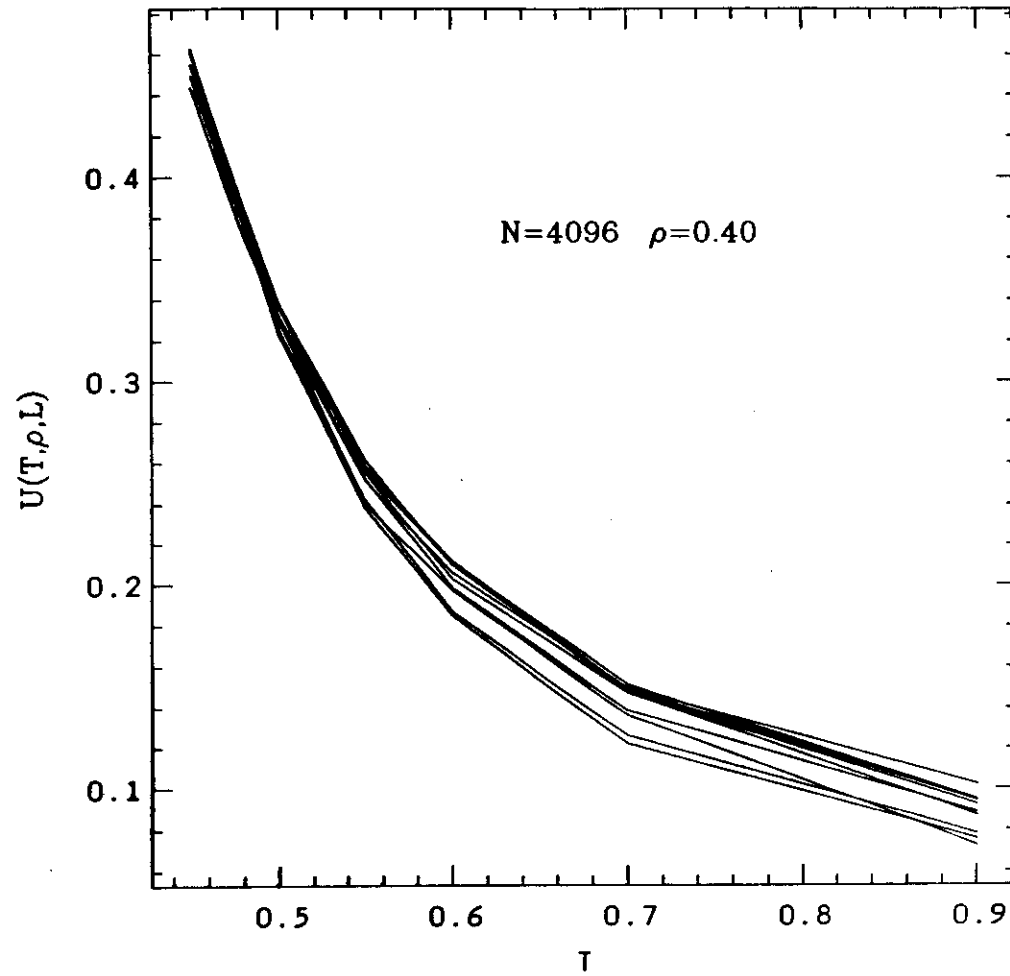
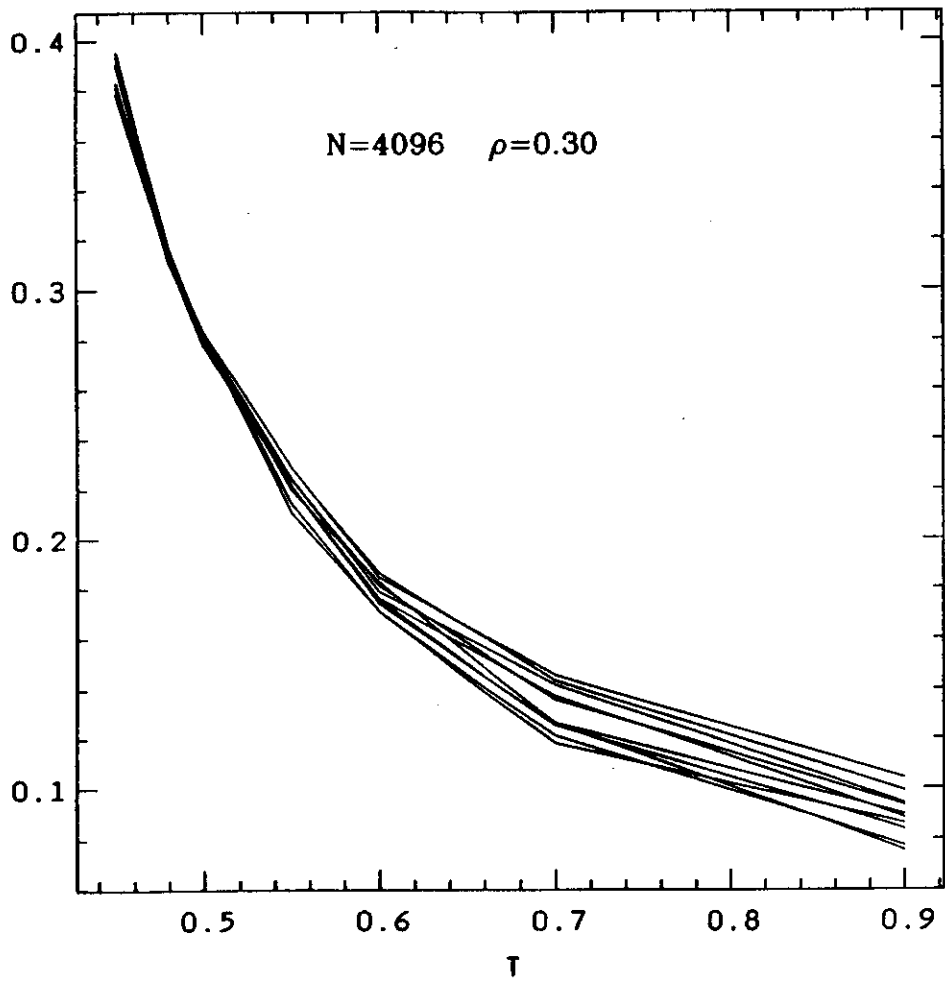
$$\rightarrow L^{\beta/\nu} \bar{P}\{(\rho - \rho_c)L^{\beta/\nu}, (p - p_c)|\epsilon|^{-\beta}, L|\epsilon|^\nu\}$$

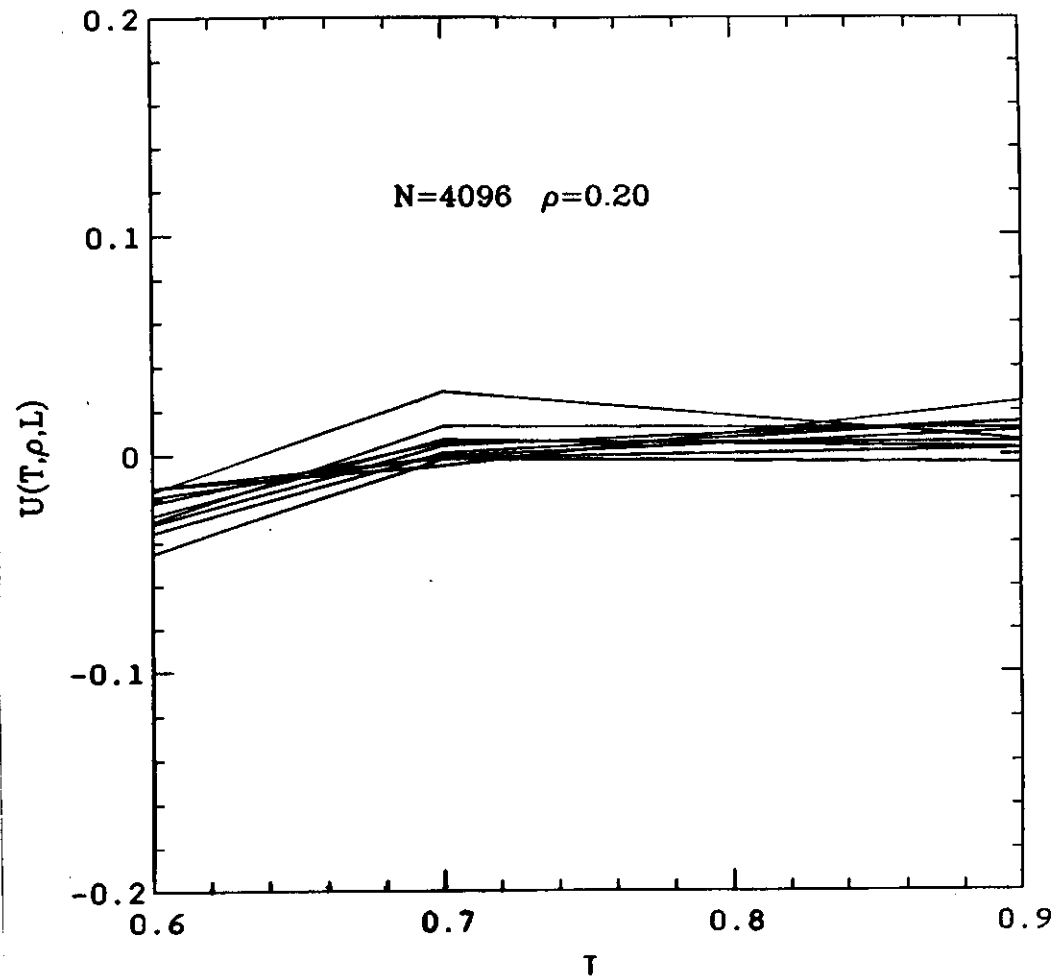
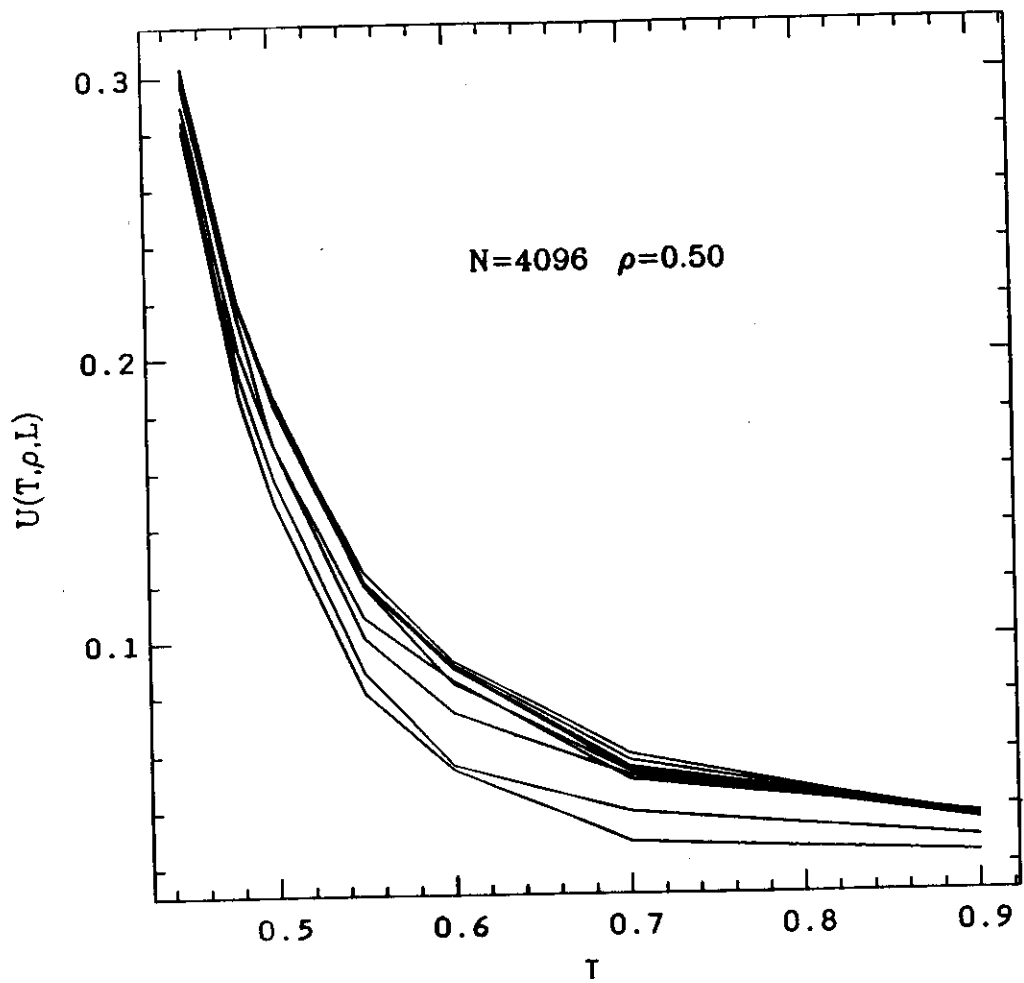
$$U_L \approx f_u\{(\langle \rho \rangle - \rho_c)\epsilon|^{-\beta}, L|\epsilon|^\nu\}$$

$$\langle \rho \rangle = \rho_c \quad T \rightarrow T_c$$

$$U_r \rightarrow U^* = f_u(0,0)$$

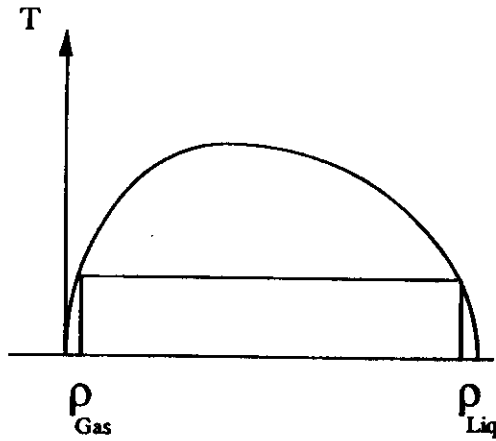
Cumulants for different block sizes



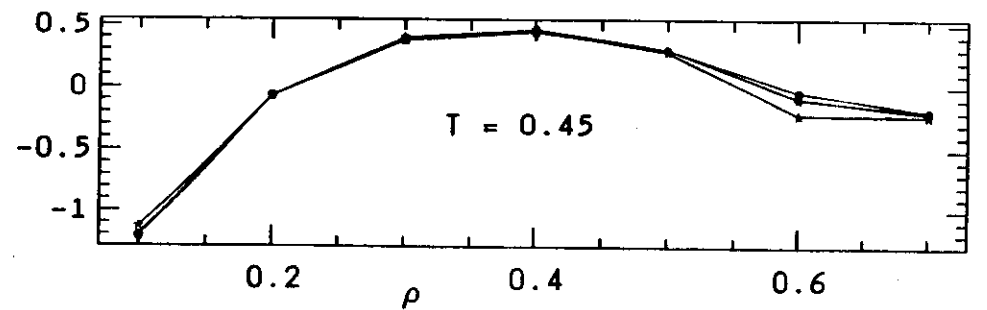
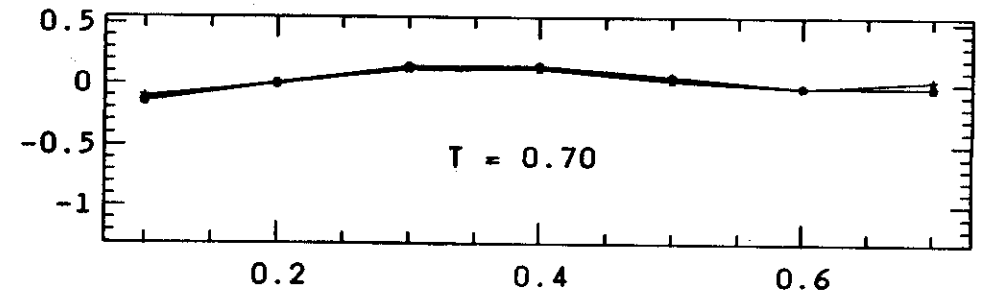
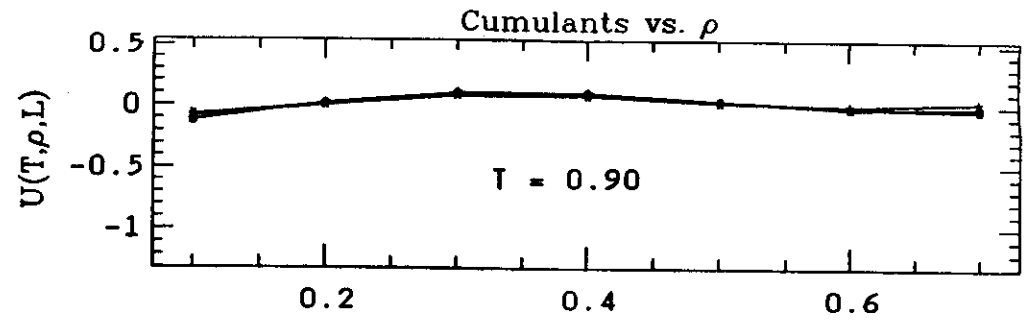
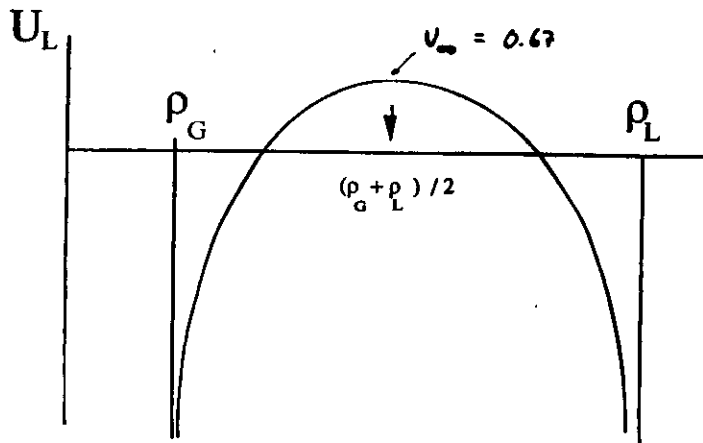


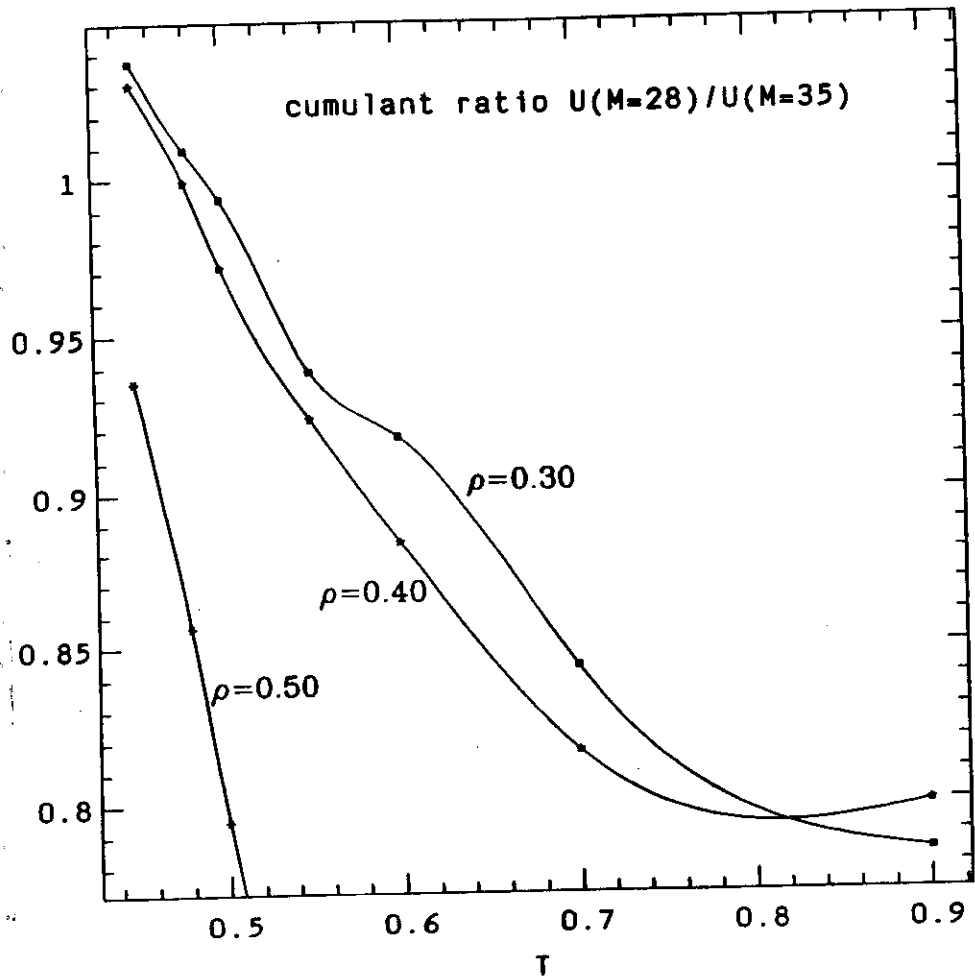
$T \ll T_c$

below the coexistence curve



$$P_L \sim a \delta(\rho - \rho_{Gas}) + b \delta(\rho - \rho_{Liq})$$





REDUCED CUMULANT

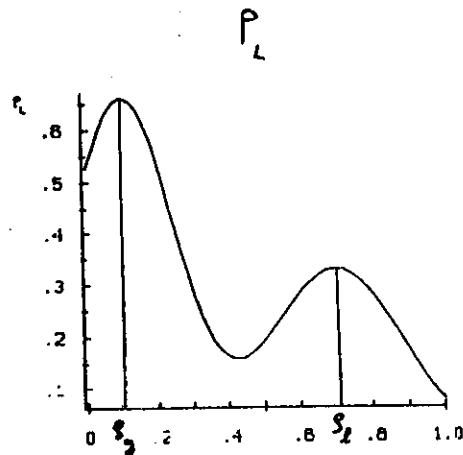
$$U_L = 1 - \frac{\langle (s - \langle s \rangle)^4 \rangle_L}{3 \langle (s - \langle s \rangle)^2 \rangle_L^2}$$

$T \gg T_c$

$\xi \ll L$

$U_L \sim L^{-d}$

$T \ll T_c$



$\xi \ll L$

$U_L \rightarrow 0.67, \quad s = \frac{1}{2} (s_1 + s_2)$

Conclusions

- 1- more detailed analysis of the results near T_c (scaling etc.)
- 2- Comparison with Gibbs ensemble
- 3- dynamics ?
- 4- critical slowing down

Future

- combine with other methods
Gibbs ensemble, G.C. MC.
- different systems
phase separation in fluid mixtures

