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**" ELECTRONIC STRUCTURE OF DISORDERED SURFACES
AND ITS APPLICATIONS "**

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These are preliminary lecture notes, intended only for distribution to participants.

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ELECTRONIC STRUCTURE OF DISORDERED SURFACES AND ITS APPLICATIONS

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MOTIVATION: surface disorder occurs frequently

- disordered / incomplete overlayers on substrates:
(Ag, Pd) on Ag(110); gases on TM surfaces (catalysis)
- surface disordering: $\text{Cu}_{50}\text{Au}_{50}(001)$
- surfaces of TM alloys with non-uniform composition:
(Ni, Pt) or (Cu, Ni) low-index surfaces
- semiconductor interfaces / superlattices: $\text{CdTe} / (\text{Cd}, \text{Hg})\text{Te}$

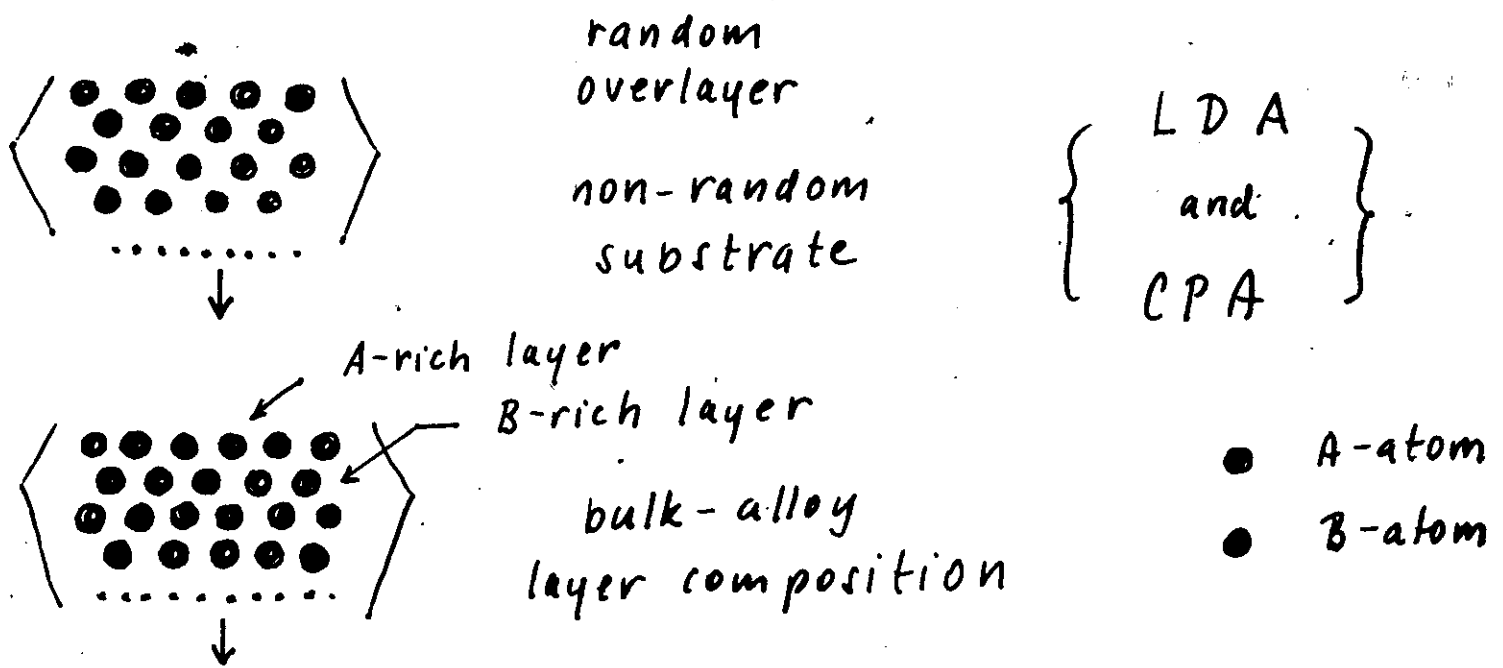
Electronic structure is key to microscopic understanding

AIM: develop electronic theory of systems with disordered surfaces / interfaces which

- starts from first principles ...
- semiinfinite sample geometry
- describes reliably concentration trends

To this end we combine

- TB-LMTO to describe electron states
- SGF method within the PL concept
- CPA to include disorder (inhomogeneous CPA)
- LSDA to account for electron correlations
- Empty spheres to describe vacuum



Basic approximation (in addition to LDA & CPA):
 Sample inhomogeneity in conjunction with LDA and CPA implies that each layer has different properties and layers are coupled each other $\rightarrow \infty$ set of equations

Surface-bulk approximation: only first N -top layers have different properties to be found selfconsistently, remaining layers have bulk properties

Typically concentrations / potentials differ from bulk ones in few top layers for metallic systems

To facilitate simultaneous treatment of disorder and surface we transform from original (z) to new (β) LMGD representation

$$\text{ORIGINAL} \rightarrow z_{RL}$$

$$\text{NEW} \rightarrow \beta_{RL} \equiv \beta_L$$

$$G(z) = (z - H)^{-1}$$

$$g^\beta(z) = (P^\beta(z) - S^\beta)^{-1}$$

$$H = C + \Delta^{1/2} S^0 \Delta^{1/2}$$

$$P^\beta(z) = (z - C) / \{\Delta + (z - \beta)(z - C)\}$$

$$S^0 = S^0 (1 - z S^0)^{-1}$$

$$S^\beta = S^0 (1 - \beta S^0)^{-1}$$

- off diagonal disorder

- site diagonal disorder

- large spatial extent of S^0

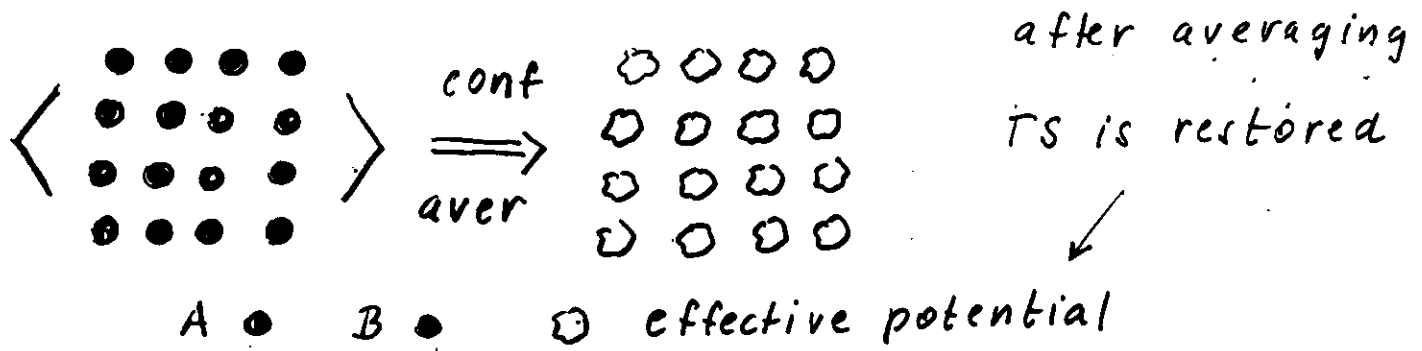
{KKR-CPA like}

- shortest possible spatial extent of S^β

Basic trick: we perform CPA averaging / introduce surface in β and final results transform back to physical z using scaling relation of $G(z)$ and $g^\beta(z)$

$$G_{RL,R'L'}(z) = \lambda_{RL}^\beta(z) \delta_{RR'} \delta_{LL'} + (\mu_{RL}^\beta(z) g_{RL,R'L'}^\beta(z) \mu_{R'L'}^\beta(z))$$

COHERENT POTENTIAL APPROXIMATION



$$c_A \begin{pmatrix} \odot & \odot & \odot \\ \odot & \bullet & \odot \\ \odot & \odot & \odot \end{pmatrix} + c_B \begin{pmatrix} \odot & \odot & \odot \\ \odot & \bullet & \odot \\ \odot & \odot & \odot \end{pmatrix} = \begin{pmatrix} \odot & \odot & \odot \\ \odot & \odot & \odot \\ \odot & \odot & \odot \end{pmatrix} \quad \begin{array}{l} \text{CPA} \\ \text{equation} \end{array}$$

CPA concept applied to each layer in question ($i=1,2,\dots,N$)

$$\mathcal{P}_i^\beta(z) = \langle \mathcal{P}_i^\beta(z) \rangle + [\mathcal{P}_i^{\beta,A}(z) - \mathcal{P}_i^\beta(z)] \phi_{ii}^\beta(z) [\mathcal{P}_i^{\beta,B}(z) - \mathcal{P}_i^\beta(z)]$$

$$\phi_{ii}^\beta(z) = \frac{1}{N_{ii}} \sum_{k_{ii}}^{SBZ} f_{ii}^\beta(k_{ii}, z | \{ \mathcal{P}_j^\beta(z) \}) \quad \begin{array}{l} \text{layer } iF \\ \text{central quantity} \end{array}$$

$\mathcal{P}_i^\beta(z)$ - coherent potential function (non-random)
determined from above coupled set of Eqs.
(coupled via $\phi_{ii}^\beta(z)$)

Properties of $\mathcal{P}_i^\beta(z)$:

- site diagonal, complex function
- non diagonal with respect to orbital index $L=lm$ due to lowering symmetry at the surface (bulk!)

PRINCIPAL LAYERS (PL)

Localized nature of S^β allows to define naturally PL which is very useful tool to determine layer GF (TB-LMTO !)

Def.: solid is divided into PL (p) such that only nn PL are coupled by S^β (i.e. off-diagonal part of problem, P^β -diagonal one)

$$S_{pq}^\beta = S_{01}^\beta \delta_{p,q+1} + S_{10}^\beta \delta_{p,q-1}, \quad S_{pp}^\beta = S_{00}^\beta$$

(bulk S^β are assumed)

Thus only $S_{00}^\beta(k_\parallel)$, $S_{01}^\beta(k_\parallel) = [S_{10}^\beta(k_\parallel)]^+$ fully describe our system [and $S^{\beta,\lambda\lambda}$, $S^{\beta,\lambda 0} = [S^{\beta,0\lambda}]^+$ for overlayer]

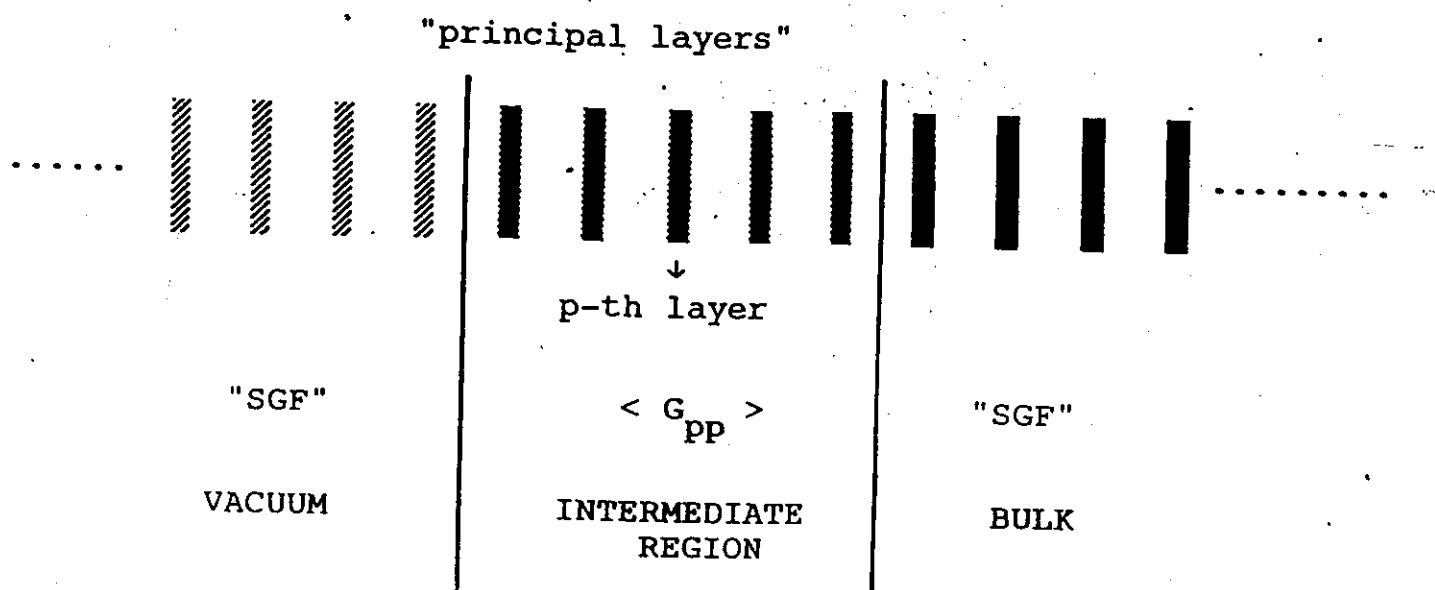
PL includes 1 to few atomic layers depending on lattice type, sample face & spatial extent of S^β

Example: fcc(001), fcc(111), bcc(110) $\rightarrow$ PL $\equiv$ 1 at. layer
 fcc(110), bcc(100) $\rightarrow$ PL $\equiv$ 2 at. layers

fcc - 1st nn, bcc - 2nd nn : most-localized MTM repres.

Remark: if p-th PL includes more atomic layers; ϕ_{ii}^β in the CPA set is a superdiagonal element of ϕ_{00}^β

Description of the semi-infinite geometry



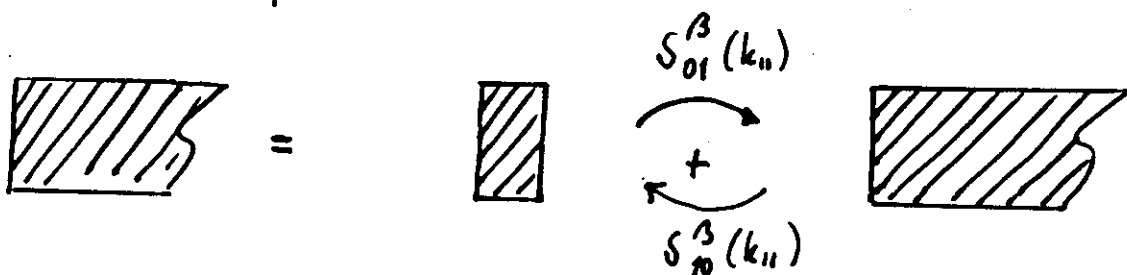
$P_{-1} - S_{00}$	$-S_{01}$								
$-S_{10}$	$P_0 - S_{00}$	$-S_{01}$							
	$-S_{10}$	$P_1 - S_{00}$	$-S_{01}$						
		$-S_{10}$	$P_2 - S_{00}$	$-S_{01}$					
			$-S_{10}$	$P_3 - S_{00}$	$-S_{01}$				
				$-S_{10}$	$P_4 - S_{00}$	$-S_{01}$			
					$-S_{10}$	$P_5 - S_{00}$	$-S_{01}$		
						$-S_{10}$	$P_6 - S_{00}$		
semi-infinite vacuum		intermediate region				semi-infinite bulk alloy			

Remark: $P_5 = P_6 = \dots P^{\text{bulk}}$, $P_0 = P_{-1} = \dots P^{\text{vac}}$

SURFACE GREEN'S FUNCTION (SGF)

We determine SGF employing removal invariance principle (RIP) avoiding knowledge of bulk GF and necessity of $k_{\perp}$ -integration of corresponding Dyson equation (cleavage perturbation method)

SGF by definition relates to ideal substrate, i.e. all layers have identical properties



$$\left| \mathcal{G}^{\beta, \sigma}(k_{||}, z) = \left\{ \mathcal{P}^{\beta, b}(z) - S_{00}^{\beta}(k_{||}) - S_{01}^{\beta}(k_{||}) \mathcal{G}^{\beta, \sigma}(k_{||}, z) S_{10}^{\beta}(k_{||}) \right\}^{-1} \right|$$

Remarks:

- (i) selfconsistent equation for $\mathcal{G}^{\beta, \sigma}(k_{||}, z)$ is solved iteratively in complex energy plane followed by deconvolution to real axis at the end
- (ii) simple and effective alternative to layer-doubling (RDT) and RS-MST techniques
- (iii) serves as an exact identity to simplify expressions for interfaces, superlattices, etc.

Summary :

Originally infinite problem in PL-indices is reduced to a finite one using surface GF to describe the semiinfinite bulk / vacuum

$P_1 - S_{00}$ $- S_{10} \mathcal{G}^v S_{01}$	$-S_{01}$	0	0
$-S_{10}$	$P_2 - S_{00}$	$-S_{01}$	0
0	$-S_{10}$	$P_3 - S_{00}$	$-S_{01}$
0	0	$-S_{10}$	$P_4 - S_{00}$ $- S_{01} \mathcal{G}^b S_{10}$

surface Green's functions

$\mathcal{G}^v$ and $\mathcal{G}^b$ are found from
removal invariance principle

$$\mathcal{G}^b(k_{||}, z) = [P^b(z) - S_{00}(k_{||}) - S_{01}(k_{||}) \cdot \mathcal{G}^b(k_{||}, z) S_{10}(k_{||})]^{-1}$$

$$\mathcal{G}^v(k_{||}, z) = [P^v(z) - S_{00}(k_{||}) - S_{10}(k_{||}) \cdot \mathcal{G}^v(k_{||}, z) S_{01}(k_{||})]^{-1}$$

LAYER-RESOLVED $k_{||}$ -INTEGRATED GF:

$$\Phi_p(z) = \frac{1}{N_{||}} \sum_{k_{||}} \langle g(k_{||}, z) \rangle_{pp} = f(\{\rho_p(z)\})$$

- (i) depends on all ρ_p relating them via a set of coupled layer-CPA equations
- (ii) enters also LDA part thus making necessary link of both CPA and LDA

LAYER-RESOLVED LDA-POTENTIAL $V_p^\alpha(r)$: $\alpha = A, B$

$$V_p^\alpha(r) = -\frac{Z^\alpha}{r} + V_p^{\alpha, H}(\tilde{\rho}(r)) + V_p^{\alpha, ex}(\tilde{\rho}(r)) + \sum_L \sum_q M_{pq}^{sL} \bar{Q}_q^L$$

— spherical part of $\rho^\alpha(r)$

$\bar{Q}_p^L$ — net charge in a given layer s^α

monopole & dipole terms of multipole expansion of $\rho^\alpha(r)$

$$\bar{Q}_p^L \propto \sum_\alpha c_p^\alpha \left\{ \int \gamma_L(r') \left(\frac{r}{s^\alpha}\right)^L \rho_p^\alpha(r) dr - Z^\alpha \delta_{L,0} \right\} \quad \begin{array}{ll} \ell=0 & \text{mono} \\ \ell=1 & \text{dipole} \end{array}$$

$$\bar{Q}_p^{\ell=0} \text{ from } \phi_p^{LL}(z) \quad \bar{Q}_p^{\ell=1} \text{ from } \phi_p^{LL'}(z)$$

M_{pq}^{sL} are intra- and inter-layer Madelung constants describing elst interaction of mono/dipo charges

ITERATION LOOP:

$$V_p^{\alpha, \text{inp}}(r) \xrightarrow{\boxed{\text{CPA}}} \bar{Q}_p^L \xrightarrow{\text{LDA}} V_p^{\alpha, \text{out}}(r)$$

WORK FUNCTION:

i) Clean Ag (001)

theory 4.78 eV (present)

4.74 eV (FLAPW: Surf. Sci. 243, 317 (1991))

4.95 eV (SEGF: PRB 37, 66 P2 (1988))

experiment: 4.2 - 4.64 eV (surface roughness)

Model: 'sandpaper'-like $\Rightarrow \sim 20\%$ vacancies
4.45 eV

Dipole barrier: $8 \text{ Ag} + 10 \text{ ES} (9 \text{ Ag} + 9 \text{ ES})$ } 7.1 - 7.2 eV
 $M_{Pq}^{s, l=1} = 0$

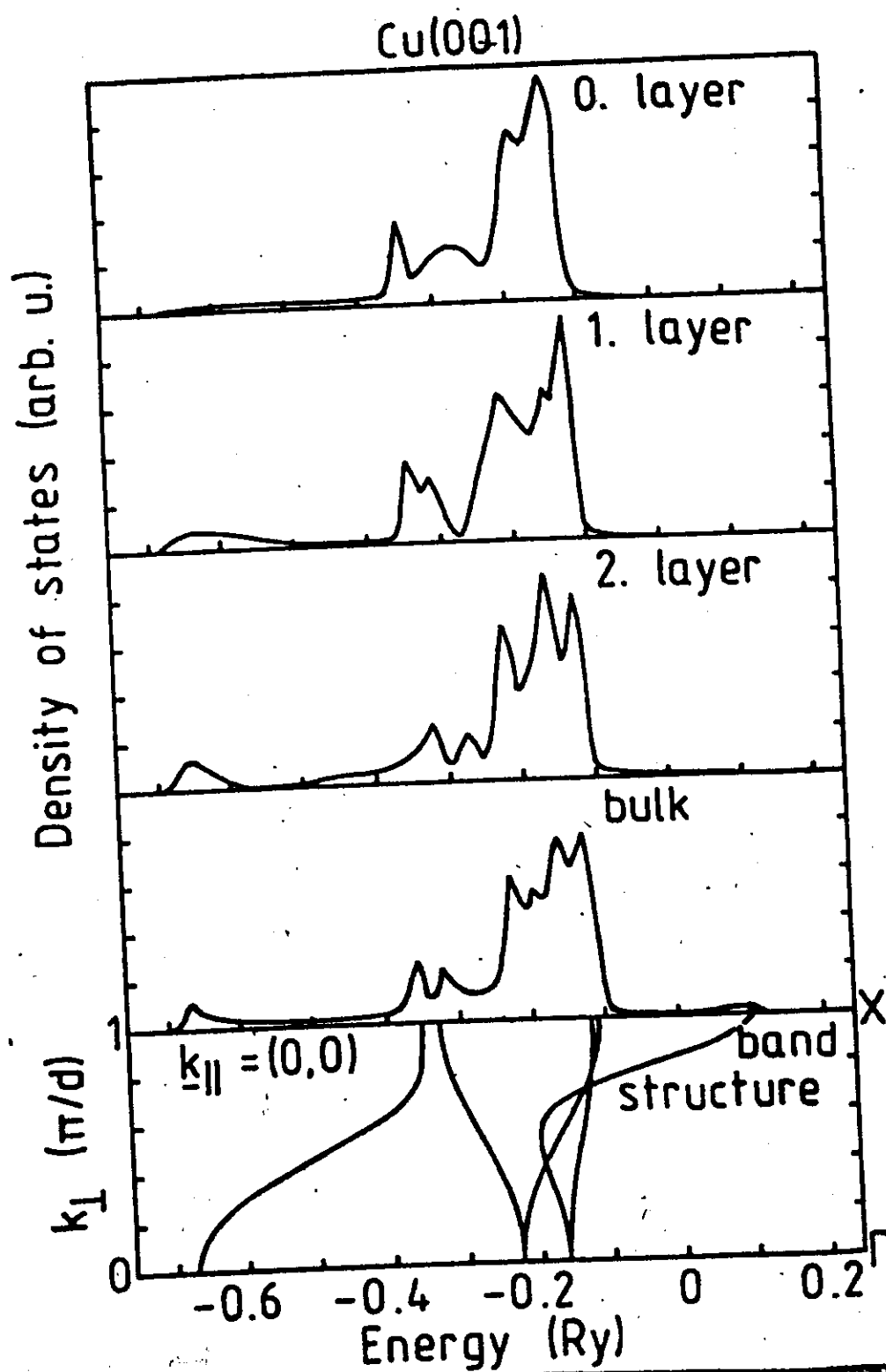
(ii) Clean Pd (001)

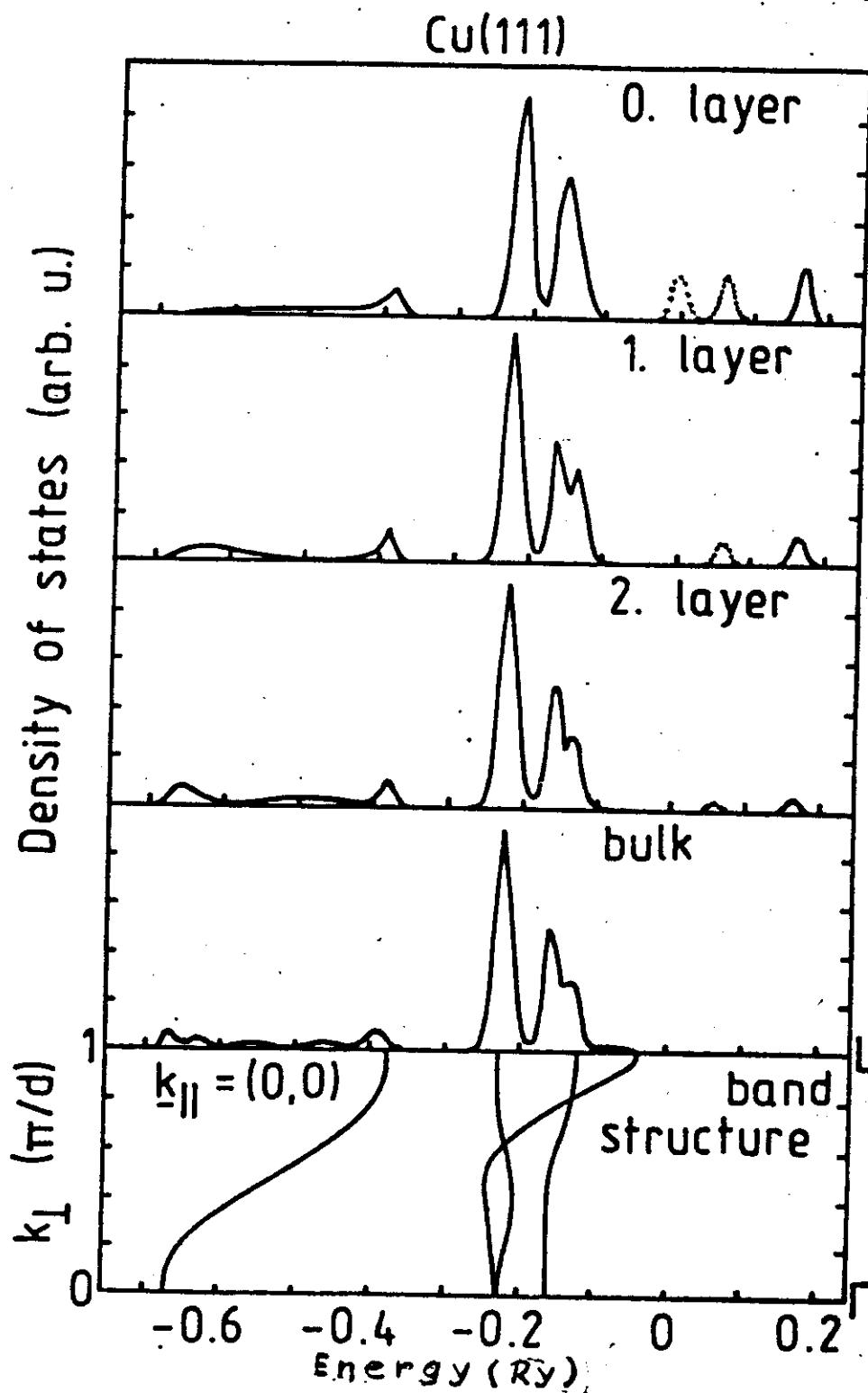
theory equilibrium: 5.83 eV expanded: 5.61 eV
(Ag)

experiment: 5.8 eV

(iii) (Ag, Pd) overlayer on Ag (001)

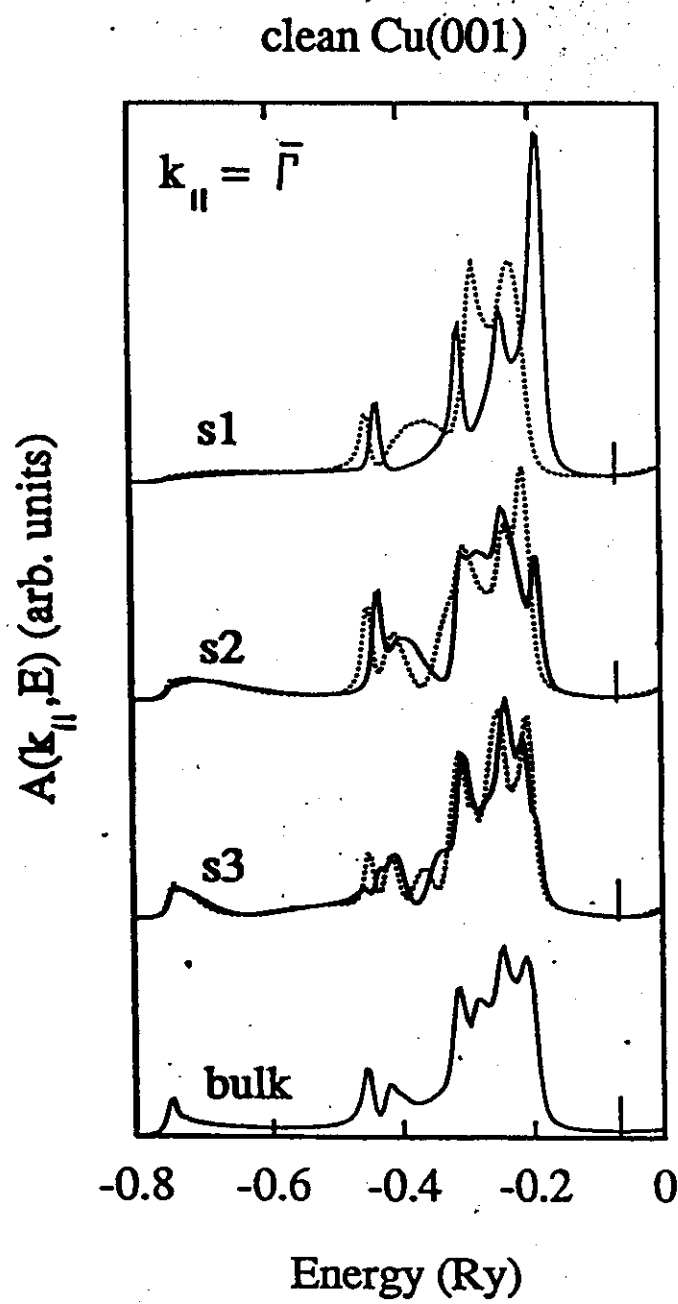
monotonic increase from 4.78 eV (clean Ag (001)) to
5.67 eV (Pd monolayer) $\rightarrow$ expanded Pd (001)





— scf calculations

• nscf calculations: bulk potentials up to the surface

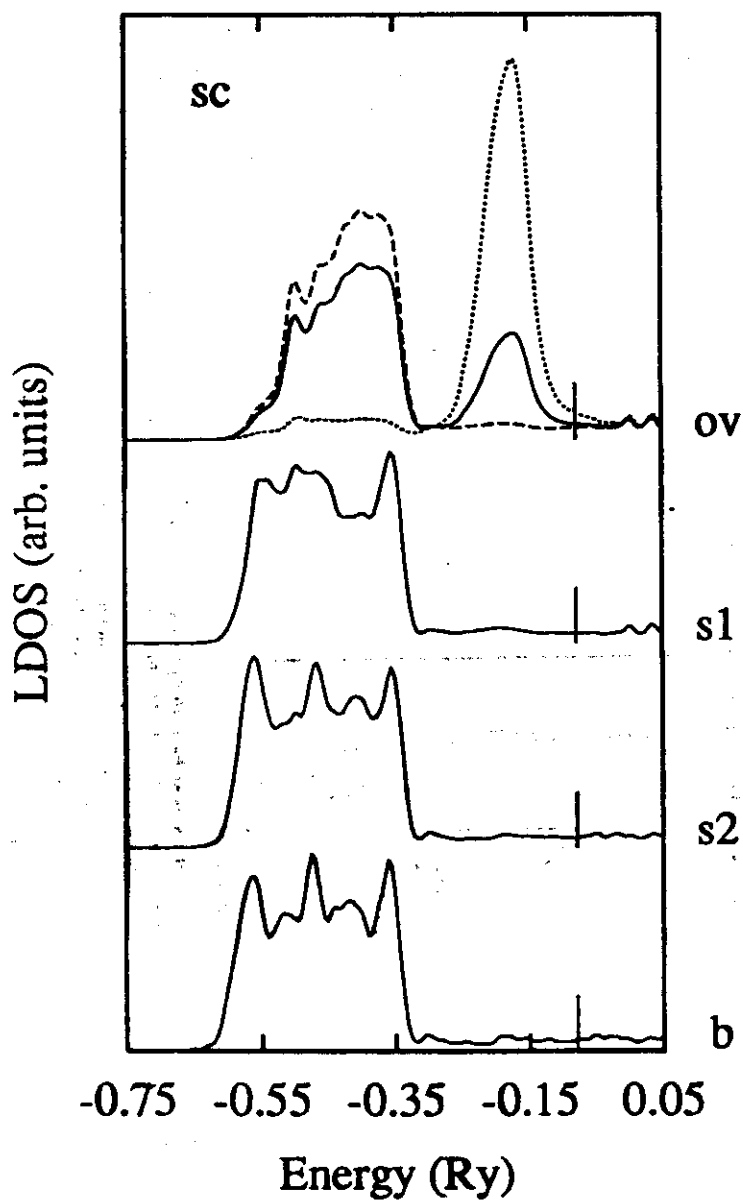


Intermediate region: 2 ES + overlayer + 2 substrate layers
 SBZ integration: 21 special $k_{||}$ -points

Layer-resolved DOS

— total DOS
 --- Ag LDOS
 Pd LDOS

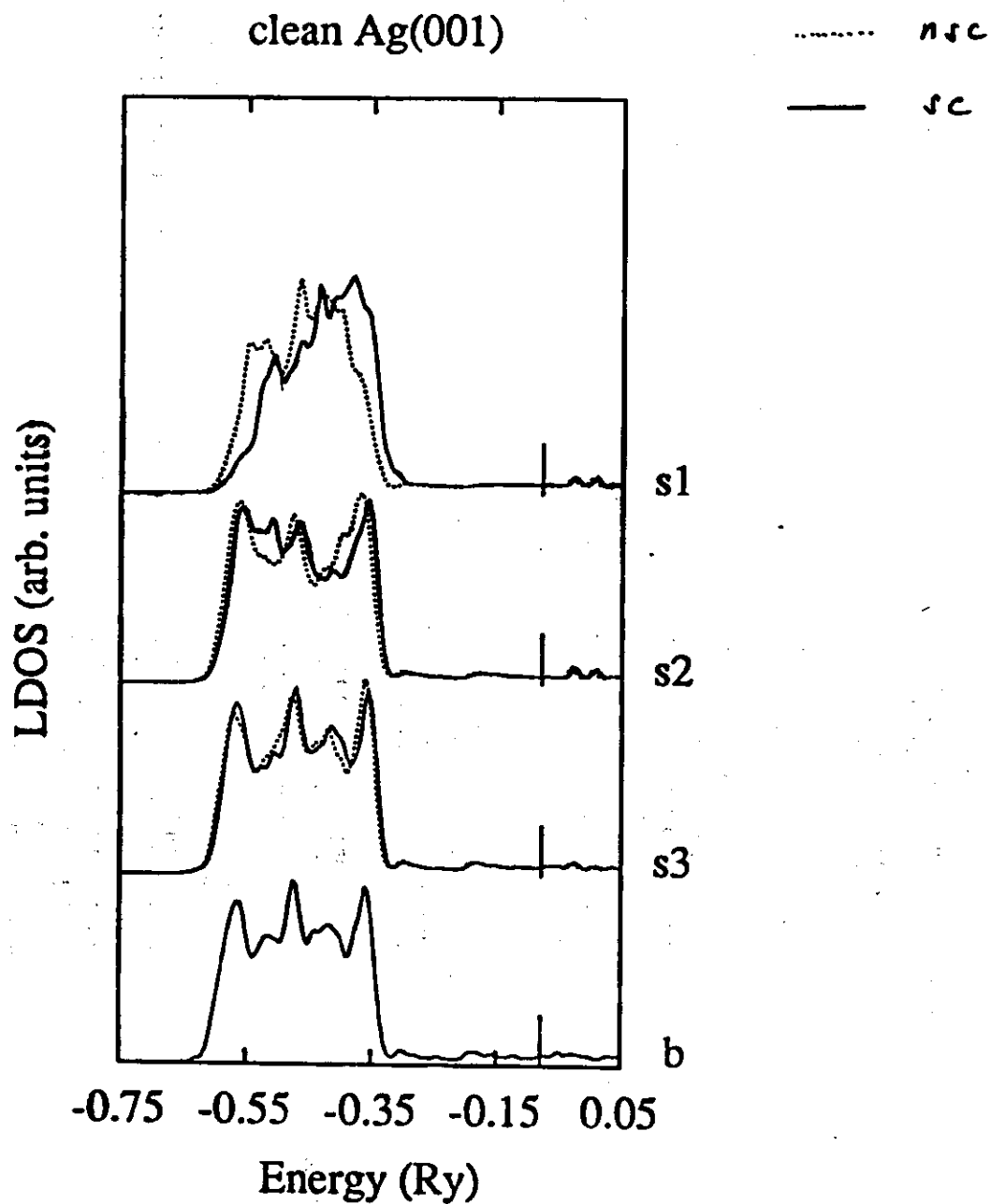
$\text{Ag}_{75}\text{Pd}_{25}$ on Ag(001)



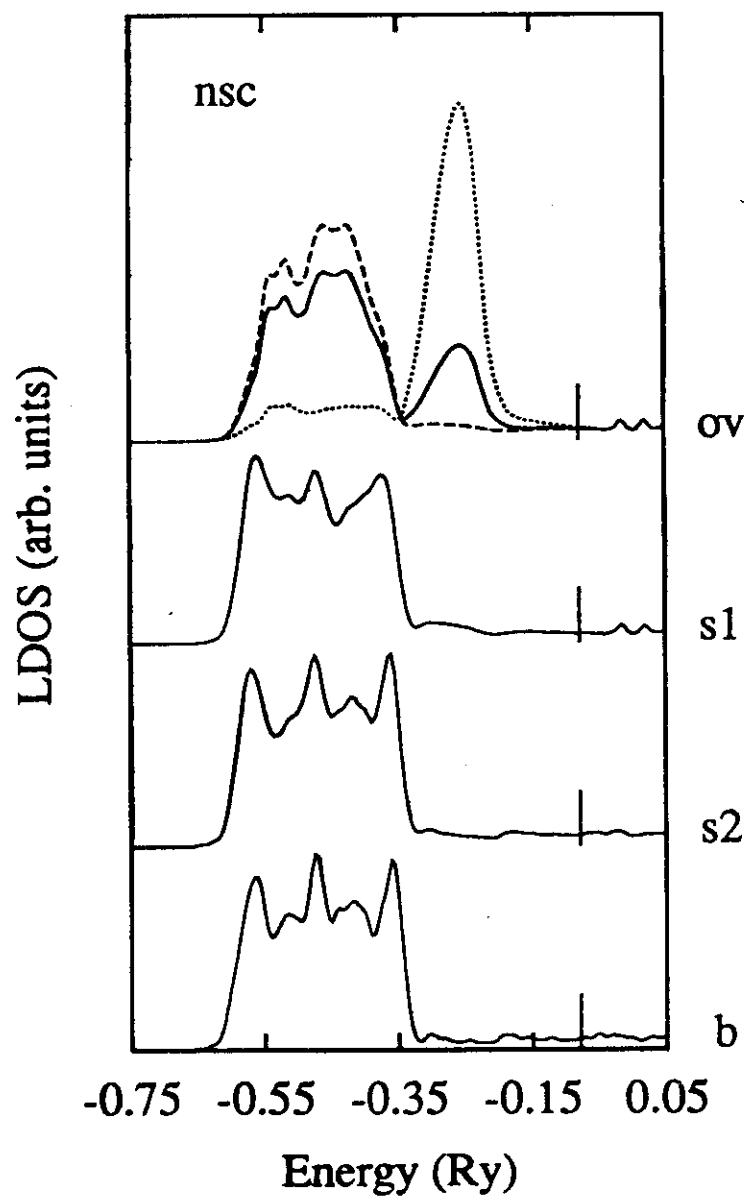
Intermediate region: 2 FS + 3 substrate layers

SBZ integration: 21 special k_n -points

Layer-resolved DOS



$\text{Ag}_{75}\text{Pd}_{25}$ on Ag(001)



ORDERING ON SURFACES: Cu(001) c(2x2)-Pd

LEED exp: $\approx 50\%$ Pd coverage, c(2x2) pattern found

Two possibilities:

- (i) CuPd surface alloy (ii) Pd overlayer
(limited to top sample layer)

Remark: similar case is Au coverage on Cu(001)

Dynamical LEED calculation $\rightarrow$ surface alloy

ARPES $\rightarrow$ surface alloy (Pd-induced features
similar as bulk CuPd)

New approach : PRL 69 (1992) 308

- we study the stability of disordered phases - $\text{Cu}_{50}\text{Pd}_{50}$
and $\text{Pd}_{50}\text{Vac}_{50}$ overlayers on Cu(001) towards the
formation of ordered phase

(analogy to magnetic instabilities of $\chi(k)$ - e.g. bcc
paramagnetic Cr $\rightarrow$ antiferromagnetic Cr).

- ordering is studied in the frame of effective 2D Ising
model constructed from corresponding ab-initio
electronic structure within the GPM

$$H = \sum_{R \neq R'} V_{RR'} \eta_R \eta_{R'}, \quad V_{RR'} = V_{RR'}^{AA} + V_{RR'}^{BB} - V_{RR'}^{AB} - V_{RR'}^{BA}$$

$$V_{RR'}^{\alpha\alpha'} = -\frac{1}{\pi} \text{Im} \text{tr} \int_{-\infty}^{E_F} \{ t_R^{\alpha}(z) \bar{g}_{RR'}(z) t_{R'}^{\alpha'}(z) \bar{g}_{R'R}(z) \} dz \quad z = E + i0 \quad \alpha, \alpha' = A, B$$

$t_R^{\alpha}(z)$ - overlayer CPA t -matrix

$$\bar{g}_{RR'}(z) = \frac{1}{N_{||}} \sum_{k_{||}} \bar{g}(k_{||}, z) e^{ik_{||}(R-R')} \quad (R, R' \in \text{overlayer})$$

- we look for extrema of $V(k_{||}) = \sum_R V_{0R} e^{ik_{||}R}$

(i) $\max \{-V(k_{||})\} = 0 \rightarrow \text{no ordering}$

(ii) $\max \{-V(k_{||})\}$ at $k_{||}^0 = (0,1) \rightarrow c(2 \times 2)$ ordering

($k_{||}^0$ - 2D analogy of Lifshitz special points)

Remark: V_{0R} up to 11-th shell were determined

Values of V_{0R} for first few shells:

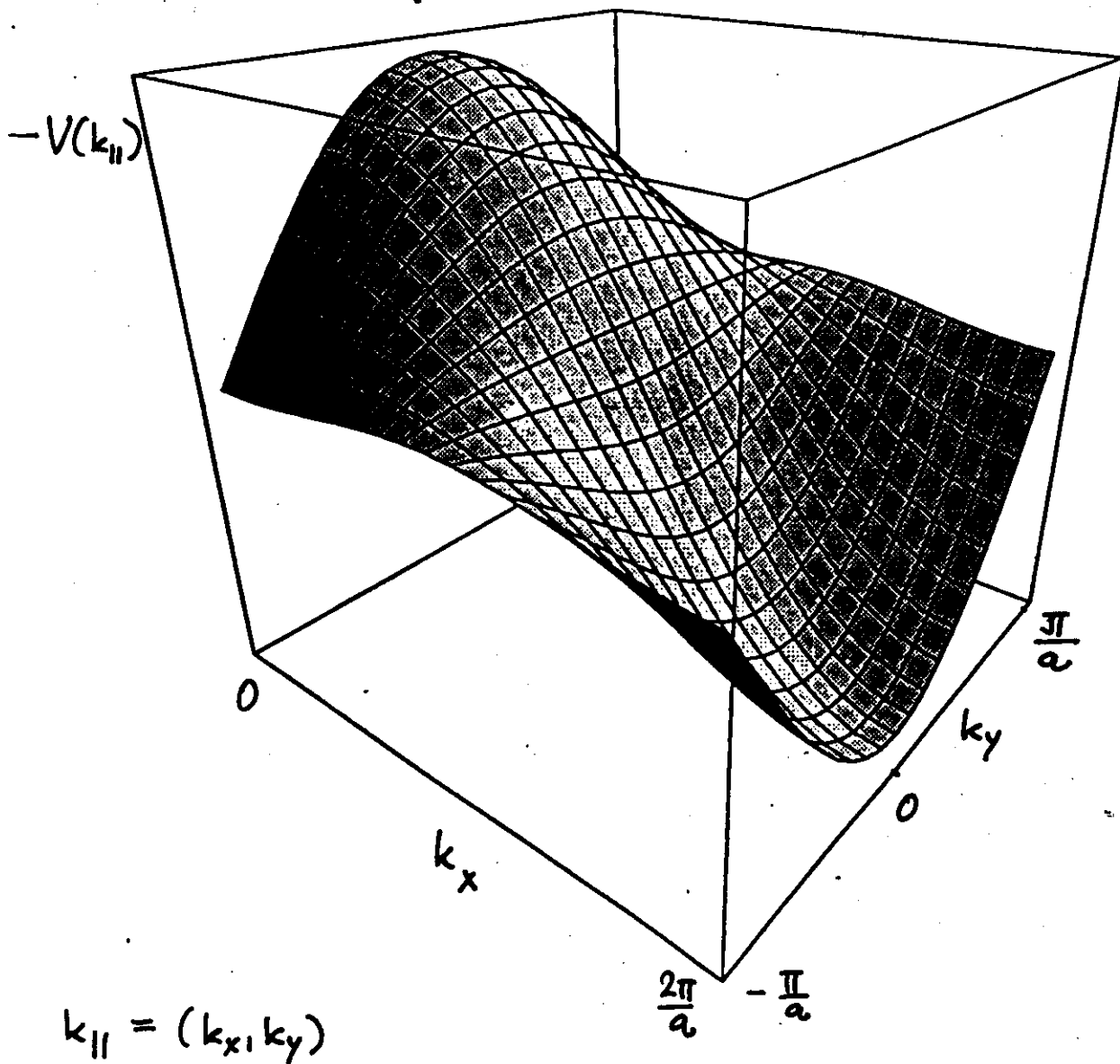
$$V_{0R} = 6.51, -0.59, -0.36, -0.18 \text{ [mRy]} \quad \dots \quad \text{CuPd}$$

$$V_{0R} = -19.79, 1.38, -0.39, -0.03 \text{ [mRy]} \quad \dots \quad \text{PdVAc}$$

Remark: the same tendency to ordering (1st NN dominating); not the type of superstructure

• Pd₅₀ Vac₅₀ overlayer on fcc (001) - Cu

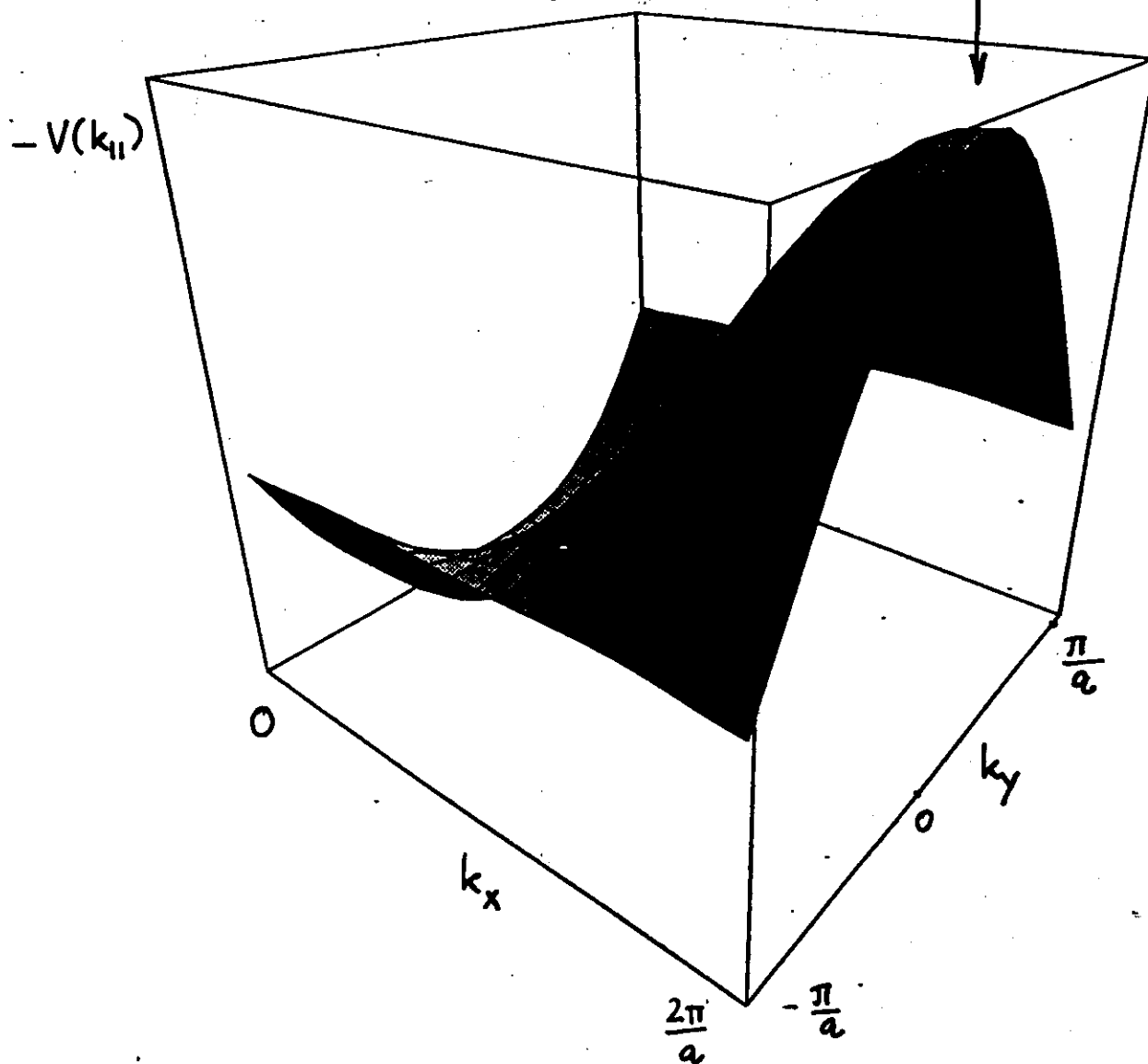
maximum at $k_{||} = (0,0)$: ORDERED STRUCTURE
UNSTABLE



$\text{Cu}_{50}\text{Pd}_{50}$ overlayer on fcc (001) Cu

maximum at $k_{||} = \frac{2\pi}{a}(1,0)$

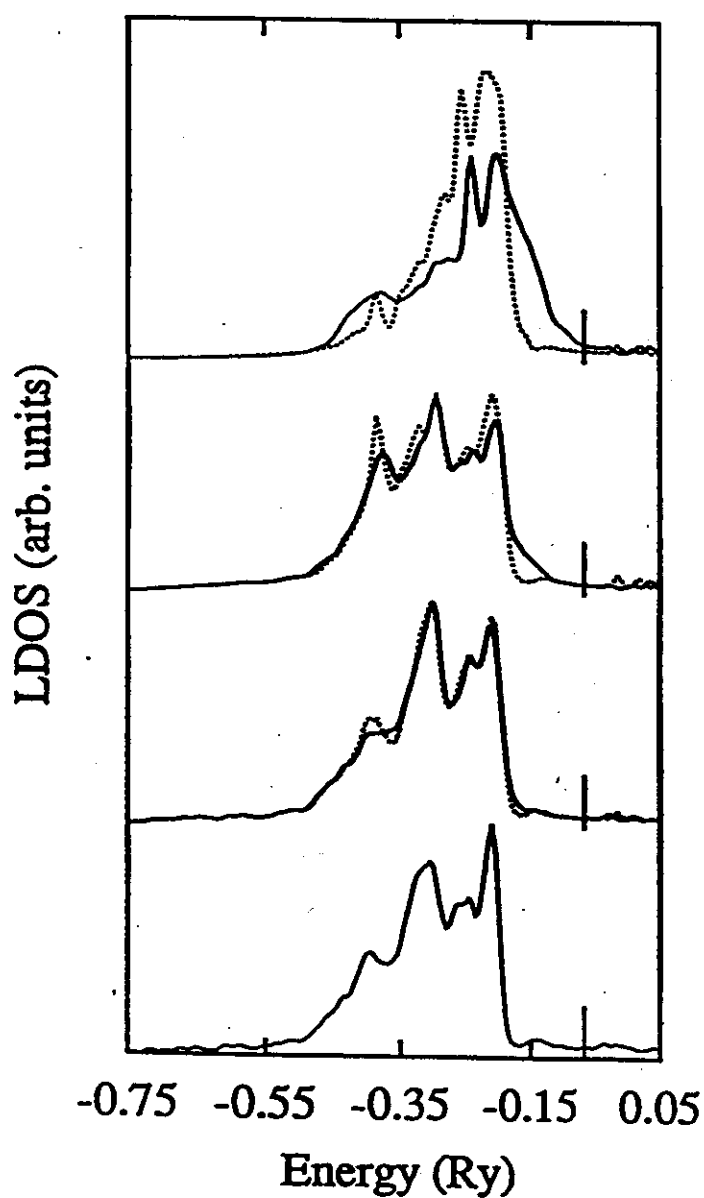
STABLE ORDERED
STRUCTURE



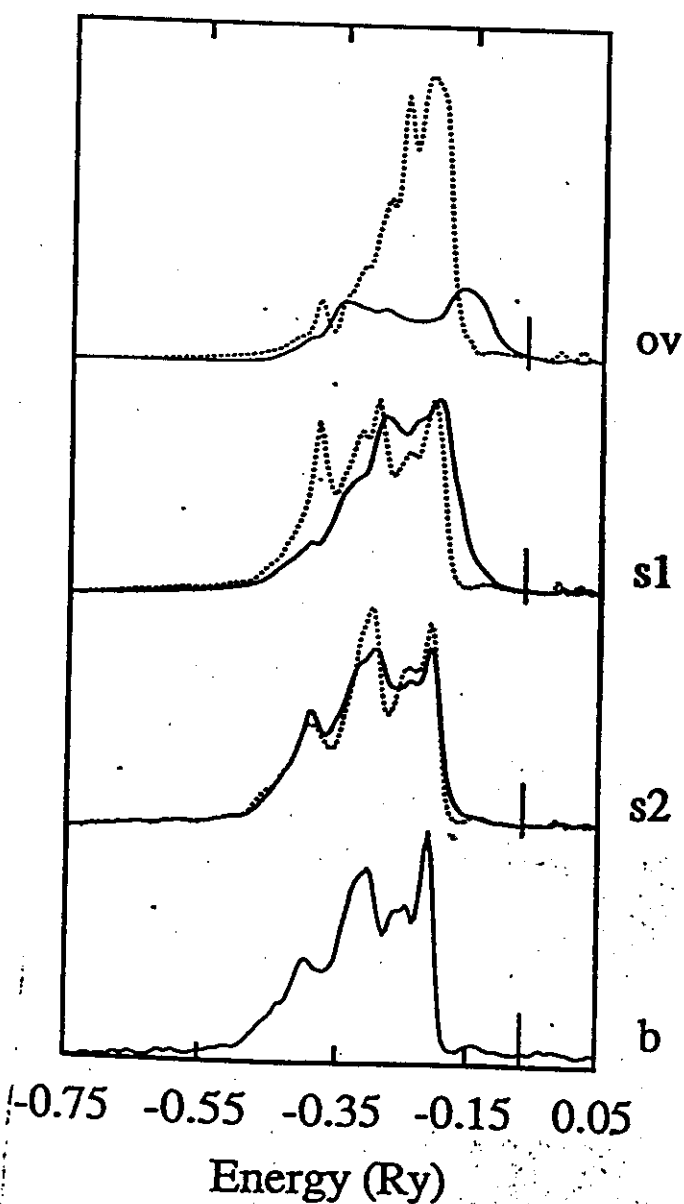
$$k_{||} = (k_x, k_y)$$

— Total DOS for Cu(001) coveredaged by overlayer
 Total DOS for clean Cu(001)

$\text{Cu}_{50}\text{Pd}_{50}$ on Cu(001)



$\text{Pd}_{50}\text{Vac}_{50}$ on Cu(001)



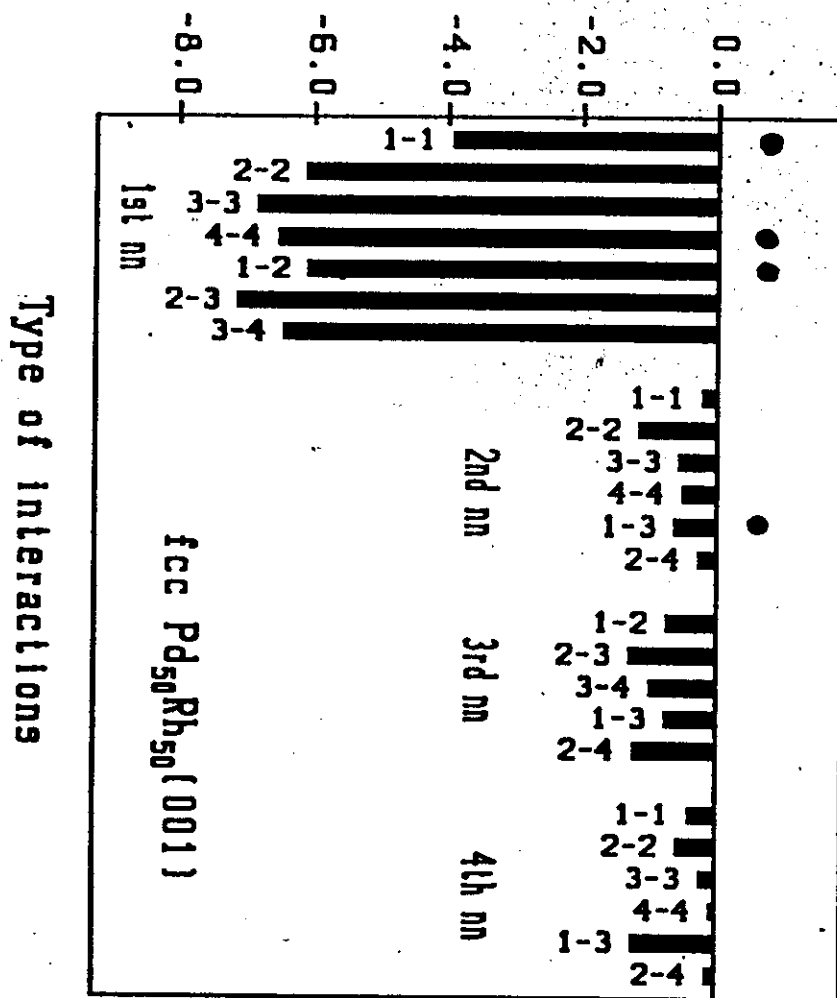
MONTÉ-CARLO SIMULATIONS OF SEGREGATION

We study segregation in $\text{Cu}_{1-x}\text{Ni}_x(001)$ alloys based on ab-initio Ising model derived for a homogeneous reference medium — scf bulk alloy

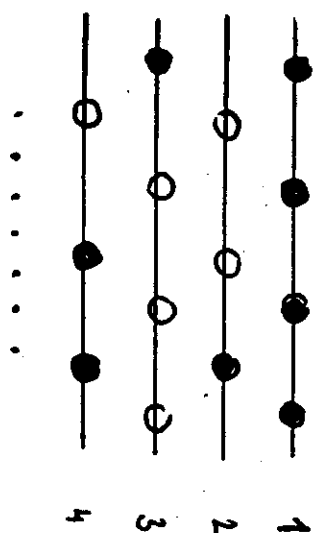
Details:

- computational cell consists of 8 fcc(001) alloy planes each containing 15×17 atoms with periodic boundary condition in layers || surface
- we start from a random configuration and interchange randomly two atoms. Their energy difference ΔE :
 - (i) $\Delta E < 0$ — favorable, represents a new configuration
 - (ii) $\Delta E > 0$ — represents a new configuration only if $\exp\{-\Delta E/kT\} > \tilde{a}$, $\tilde{a}$ — random number $\in (0,1)$
- this step is repeated and the system approaches the thermodynamical equilibrium if the number of interchanges is sufficiently large
- during MC steps, two bottom layers of slab have fixed bulk concentration

Pair interactions (mRy/atom)

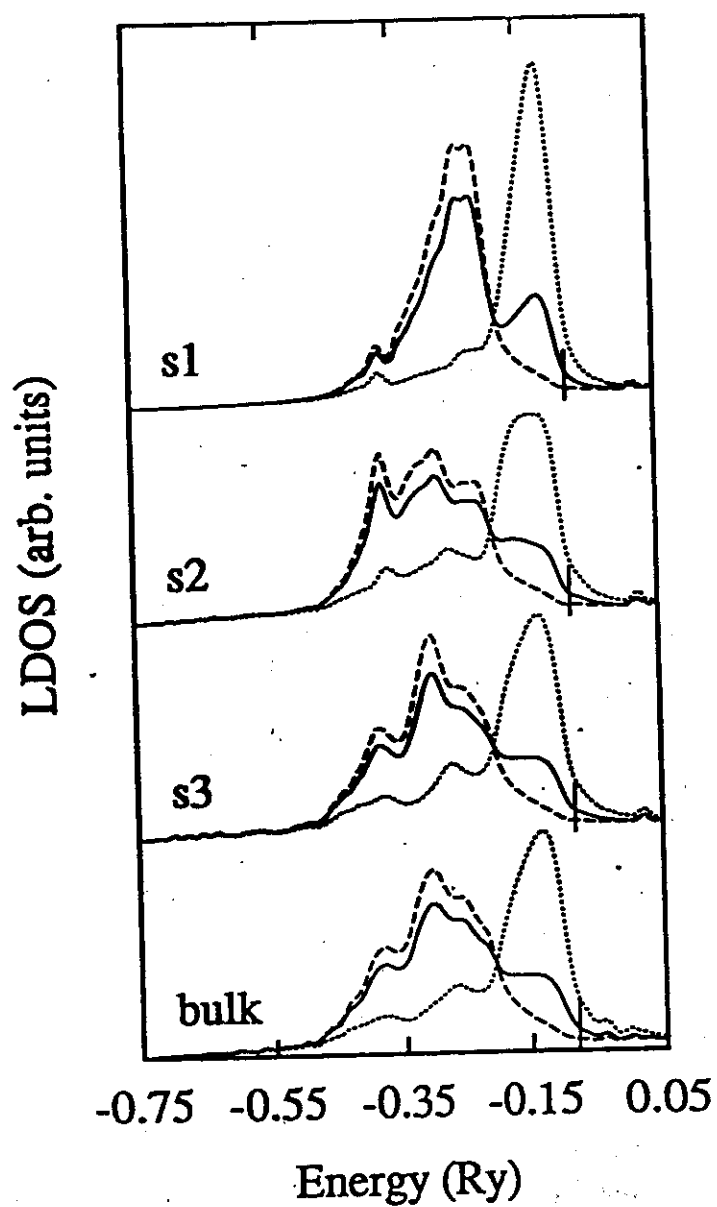


homogeneous profile



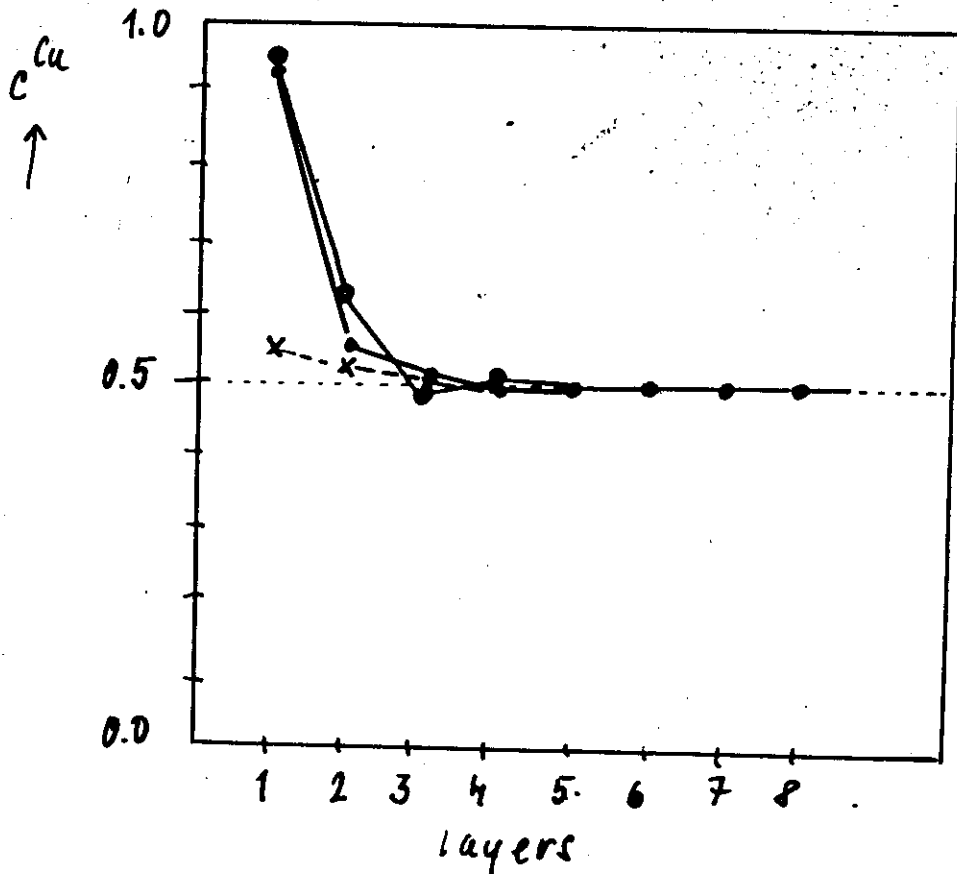
— total DOS
- - - Cu - LDOS
..... Ni - LDOS

fcc - Cu₇₅Ni₂₅(001)



Results of MC simulations: fcc-Cu₅₀Ni₅₀ (001)

$T = 800 \text{ K}$



—
V_{OR} and
D_R incl

—
V_{OR} = 0

only D_R incl

- - -
D_R = 0

only V_{OR} incl

Remarks:

- (i) Values of on-site terms of Ising Hamiltonian are decisive for segregation
- (ii) Good agreement with experiment over all concentration range

SUMMARY

We have developed self theory of electronic states of disordered surfaces within the LDA

- use of screened structure constants allows to include semiinfinite geometry simply and effectively
- use of inhomogeneous CPA allows a unified treatment of clean surfaces, various stages of adsorption (low coverage $\rightarrow$ films) and surface segregation
- use of multipole expansion of charge density at the surface gives proper description of dipole barrier
- in conjunction with GPM forms a theoretical background for study of ordering and segregation phenomena at surfaces from first principles

OUTLOOK

- surface magnetism: magnetic overlays, surfaces of mag. allo.
- relativistic effects: 4d and 5d metals, Dirac equation
- magnetic GPM: interplay of chemical & magnetic interaction

