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The Hodge Decomposition Theorems (II)

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These are preliminary lecture notes, intended only for distribution to participants

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The Hodge decomposition theorems II

Christopher Deninger

In this writeup we give a proof of the identity $\Delta = 2\bar{\Delta}$ where $\Delta = dd^* + d^*d$ is the ordinary Laplacian and $\bar{\Delta} = \bar{\partial}\bar{\partial}^* + \bar{\partial}^*\bar{\partial}$ is the $\bar{\partial}$ -Laplacian on a Kähler manifold M . From the harmonic forms representation of cohomology classes it follows from this that if M is compact:

$$H_{dR}^n(M, \mathbb{C}) = \bigoplus_{p+q=n} H^{p,q}, \quad \overline{H^{p,q}} = H^{q,p}$$

where $H^{p,q} \cong H^q(M, \mathbb{R}^p)$. This is the Hodge decomposition of cohomology.

The proof of $\Delta = 2\bar{\Delta}$ on a Kähler manifold is surprisingly involved. We will largely follow [We]. Other approaches can be found in [G-H] and [Weil]. Let us begin with some representation theory of the Lie algebra $\mathfrak{sl}_2 = \mathfrak{sl}_2(\mathbb{C})$ of 2×2 complex matrices with trace zero and bracket $[A, B] = AB - BA$. Recall that a representation of a Lie algebra $\mathfrak{g}$ in a $\mathbb{C}$ -vector space V is a homomorphism of Lie algebras

of $\mathfrak{L} \rightarrow \text{End } V$ where $\text{End } V$ is the Lie algebra of endomorphisms of V with bracket $[X, Y] = XY - YX$.
 Note that $\mathfrak{sl}_2$ is 3-dimensional with basis

$$X = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}, \quad Y = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}, \quad H = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

We have $[X, Y] = H$, $[H, X] = 2X$, $[H, Y] = -2Y$.

Recall that since $SL_2 = SL_2(\mathbb{C}) = \{g \in GL_2(\mathbb{C}) \mid \det g = 1\}$ is simply connected we have a 1-1 correspondence

$$(1.1) \quad \left\{ \begin{array}{l} \text{Lie alg. repr.} \\ \mathfrak{sl}_2 \xrightarrow{\rho} \text{End } V \end{array} \right\} \longleftrightarrow \left\{ \begin{array}{l} \text{Lie group repr.} \\ SL_2 \xrightarrow{\pi} GL(V) \end{array} \right\} \quad \text{if } \dim V < \infty$$

$$\text{given by: } \begin{array}{ccc} \rho & \longmapsto & \pi(\exp A) \cdot v := (\exp \rho(A)) \cdot v \\ d\rho & \longleftarrow & \pi \end{array}$$

In particular for every Lie algebra representation ρ we can consider the action on V of the element

$$(1.2) \quad w = \exp\left(\frac{i\pi}{2}(X+Y)\right) = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} \in SL_2.$$

It will be very important for us to have a formula for this action in a particular instance. To deduce this formula let us prove it first for a very concrete representation: The group SL_2 acts on polynomials $P \in \mathbb{C}[T_1, T_2]$ by the formula

$$(g \cdot P)(T_1, T_2) := P((T_1, T_2)g)$$

(3)

This defines a representation π of SL_2 on $\mathbb{C}[T_1, T_2]$.

Note that for $g = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$ we have

$$(1.3) \quad (\pi(g) P)(T_1, T_2) = P(aT_1 + cT_2, bT_1 + dT_2).$$

The representation π decomposes into the representations π_m of SL_2 on homogeneous polynomials of degree m .

$$\text{Note that } \pi_m(g) (T_1^{m-k} T_2^k) = (\pi_1(g) T_1)^{m-k} (\pi_1(g) T_2)^k$$

because π acts by $\mathbb{C}$ -algebra automorphisms.

If $\rho_m = d\pi_m$ is the corresponding Lie algebra representation of $\mathfrak{sl}_2$ on homogeneous polynomials of degree m then

$$(1.4) \quad (\rho_1 P)(T_1, T_2) = P(aT_1 + cT_2, bT_1 + dT_2)$$

for P homogeneous of degree 1 and ρ_m is obtained by extending ρ_1 as a derivation. In particular we have:

$$\rho_1(H) T_1 = T_1, \quad \rho_1(X) T_1 = 0, \quad \rho_1(Y) T_1 = T_2$$

$$\rho_1(H) T_2 = T_2, \quad \rho_1(X) T_2 = T_1, \quad \rho_1(Y) T_2 = 0.$$

By induction it follows from this that if we set

$$u_{-1} = 0, \quad u_0 = T_1^m, \quad u_k = \binom{m}{k} T_1^{m-k} T_2^k$$

then we have the following relations:

$$u_k = \frac{1}{k!} \rho_m(Y)^k u_0 \quad \text{for } k=0, 1, \dots, m$$

$$(1.5) \quad \rho_m(H) u_k = (m-2k) u_k$$

$$\rho_m(Y) u_k = (k+1) u_{k+1}$$

$$\rho_m(X) u_k = (m-k+1) u_{k-1}.$$

It is clear that the representation ρ_m is irreducible of dimension $m+1$. It is known that every irreducible representation of dimension $m+1$ of $\mathfrak{sl}_2$ is equivalent to ρ_m but we will not need this fact. Now let us compute the action of the element w in SL_2 on homogeneous polynomials of degree m i.e. the action of $\pi_m(w)$. Since $w = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}$ we have

by (1.3)

$$\begin{aligned} \pi_m(w) T_1^{m-k} T_2^k &= (iT_2)^{m-k} (iT_1)^k \\ &= i^m T_1^k T_2^{m-k}. \end{aligned}$$

Hence

$$(1.6) \quad \pi_m(w) (\rho_m(Y)^k u_0) = i^m \frac{k!}{(m-k)!} (\rho_m(Y)^{m-k} u_0).$$

This is the formula we have been referring to above. We will require it in a more abstract setting:

(5)

Let $\rho : \mathfrak{sl}_2 \rightarrow \text{End } V$ be any finite dimensional representation of $\mathfrak{sl}_2$. A vector $v_0 \in V$ is said to be primitive of weight λ if

$$v_0 \neq 0, \quad \rho(H) v_0 = \lambda v_0 \quad \text{and} \quad \rho(X) v_0 = 0.$$

Then letting $v_{-1} = 0$ and setting $v_k = \frac{1}{k!} \rho(Y)^k v_0$

for $k = 0, 1, 2, \dots$ one obtains:

Lemma: (a) $\rho(H) v_k = (\lambda - 2k) v_k$

(b) $\rho(Y) v_k = (k+1) v_{k+1}$

(c) $\rho(X) v_k = (\lambda - k + 1) v_{k-1}$.

Moreover $\lambda = m$ where $m+1 = \dim \langle v_0, v_1, \dots \rangle$ and the subrepresentation of V generated by the v_i is irreducible.

The proof follows essentially from the commutation relations in $\mathfrak{sl}_2$, e.g. for (a): Let $\rho(H)v = \lambda v$, then:

$$\begin{aligned} \rho(H)(\rho(Y)v) &= (\rho(H)\rho(Y) - \rho(Y)\rho(H))v + \rho(Y)\rho(H)v \\ &= \rho([H, Y])v + \lambda \rho(Y)v \\ &= -2\rho(Y)v + \lambda \rho(Y)v \\ &= (\lambda - 2)\rho(Y)v. \end{aligned}$$

This implies (a) by induction. The rest of the proof

is left to the reader as an exercise c.f. [46] V 3.7. #. (5)

It is clear from the lemma and from (1.5) that mapping u_k to v_k induces an equivalence of ^{the} $\mathfrak{sl}_2$ -representations ρ_u and $\rho|_{\langle v_0, v_1, \dots \rangle}$. Hence we obtain from (1.6) the formula:

$$(1.7) \quad \pi(w) (\rho(Y)^k v_0) = i^m \frac{k!}{(m-k)!} (\rho(Y)^{m-k} v_0)$$

where $d\pi = \rho$, v_0 as above. ($\lambda = m \in \mathbb{N}$ automatic by the lemma).

After these preparations let us do some hermitian exterior algebra which will eventually lead to a particular $\mathfrak{sl}_2$ -representation.

Let E be a $\mathbb{C}$ -vector space of dimension n . Let $E' = \text{Hom}_{\mathbb{R}}(E, \mathbb{R})$ i.e. the real dual space of the underlying real vector space of E and set

$$F = E' \otimes_{\mathbb{R}} \mathbb{C} = \text{Hom}_{\mathbb{R}}(E, \mathbb{C}).$$

Clearly $\dim_{\mathbb{C}} F = 2n$ and F is equipped with a complex conjugation. One checks that

$$F = F^{1,0} \oplus F^{0,1}$$

where $F^{1,0} = \text{Hom}_{\mathbb{C}}(E, \mathbb{C})$ ($\mathbb{C}$ -lin. Homs)

$F^{0,1} = \text{Hom}_{\mathbb{C}\text{-anti}}(E, \mathbb{C})$ ($\mathbb{C}$ -antilin. Homs).

Clearly $\overline{F^{1,0}} = F^{0,1}$. On the complex exterior algebra $\wedge F$ of F we get an induced decomposition:

$$\wedge F = \bigoplus_{p=0}^{2n} \wedge^p F$$

where $\wedge^p F = \bigoplus_{r+s=p} \wedge^{r,s} F$, $\wedge^{r,s} F = \wedge^r F^{1,0} \otimes_{\mathbb{C}} \wedge^s F^{0,1}$.

Clearly $\overline{\wedge^{r,s} F} = \wedge^{s,r} F$ with respect to the natural complex conjugation on $\wedge^p F = (\wedge_{\mathbb{R}}^p E') \otimes_{\mathbb{R}} \mathbb{C}$. The elements $\omega \in \wedge^p F$ are called p -forms. We say that ω is real if $\bar{\omega} = \omega$. Let $(\wedge^p F)_{\mathbb{R}} \cong \wedge_{\mathbb{R}}^p E'$ denote the subspace of real p -forms.

Now suppose that E carries a Hermitian inner product $\langle, \rangle$. We may view $\langle, \rangle$ as an element of $\wedge^2 F$. As such it lies actually in $\wedge^{1,1} F$. There is a basis

$z_1, \dots, z_n$ of $\wedge^{1,0} F = F^{1,0} = \text{Hom}_{\mathbb{C}}(E, \mathbb{C})$ such that

$$\langle, \rangle = \sum_{\mu=1}^n z_{\mu} \otimes \bar{z}_{\mu}. \quad (\text{Exercise!})$$

(8)

The fundamental 2-form associated to $\langle, \rangle$ is by definition

$$(2.1) \quad \Omega := -\frac{1}{2} \operatorname{Im} \langle, \rangle = \frac{i}{2} \sum_{\mu=1}^p z_{\mu} \wedge \bar{z}_{\mu} = \sum_{\mu=1}^p x_{\mu} \wedge y_{\mu}$$

if we set $x_{\mu} = \frac{1}{2}(z_{\mu} + \bar{z}_{\mu})$, $y_{\mu} = \frac{1}{2i}(z_{\mu} - \bar{z}_{\mu})$. Note here

that $z_{\mu} \wedge \bar{z}_{\mu} = \frac{1}{2}(z_{\mu} \otimes \bar{z}_{\mu} - \bar{z}_{\mu} \otimes z_{\mu})$. Obviously Ω

is a real form of type $(1,1)$. By (2.1) we have:

$$\Omega^n = n! \, x_1 \wedge y_1 \wedge \dots \wedge x_n \wedge y_n \quad \text{in } \Lambda^{2n} E' = (\Lambda^{2n} F)_{\mathbb{R}}.$$

The bilinear form $S = \operatorname{Re} \langle, \rangle$ defines a real scalar product on the $\mathbb{R}$ -vector space underlying E . We have

$$S = \sum_{\mu} (x_{\mu} \otimes x_{\mu} + y_{\mu} \otimes y_{\mu}).$$

Let S' be the scalar product induced on the real dual E' . By definition the x_{μ}, y_{μ} form an orthonormal basis of E' with respect to S' and hence

$$(2.2) \quad \operatorname{vol} = \frac{1}{n!} \Omega^n$$

is an orientation of E' i.e. an ON basis of $\Lambda^{2n} E'$ with respect to the scalar product induced by S' .

Note that vol depends only on the Hermitian inner product $\langle, \rangle$ on E .

(9)

With respect to S' on E' and vol there is a Hodge star operator

$$* : \Lambda^p E' \longrightarrow \Lambda^{2n-p} E' \quad \text{for every } 0 \leq p \leq 2n.$$

We will extend it $\mathbb{C}$ -linearly to a $*$ -operator

$$* : \Lambda^p F \longrightarrow \Lambda^{2n-p} F.$$

Let $\pi_p : \Lambda F \rightarrow \Lambda^p F$, $\pi_{p,q} : \Lambda F \rightarrow \Lambda^{p,q} F$ be the natural projections and set

$$W = \sum_p (-1)^p \pi_p : \Lambda F \rightarrow \Lambda F$$

$$J = \sum_{p,q} i^{p-q} \pi_{p,q} : \Lambda F \rightarrow \Lambda F.$$

Note that since $J = i$ on $\Lambda^{1,0} F = F^{1,0}$, $J = -i$ on $\Lambda^{0,1} F = F^{0,1}$ the operator J is the $\mathbb{C}$ -algebra extension to ΛF of the operator J on F given by $(J\varphi)(e) = \varphi(ie)$.

Let

$$L : \Lambda F \rightarrow \Lambda F \quad \text{be defined by } L(v) = \Omega \wedge v.$$

Clearly $L(\Lambda^{p,q} F) \subset \Lambda^{p+1, q+1} F$ and L is a real operator i.e. commutes with complex conj. on ΛF .

Recall that $\Lambda^p F$ has a natural Hermitian inner

product given by:

$$\langle \alpha, \beta \rangle_{\text{vol}} = \alpha \wedge * \bar{\beta} \quad \text{for } \alpha, \beta \in \wedge^p F.$$

Let L^* be the adjoint of L with respect to this inner product. We have the following relations:

(2.3) Thm.: 1) $* * = w$, $J^2 = w$

2) $L^* = w * L *$ in particular L^* is real, $L^*(\wedge^p F) \subset \wedge^{p-2} F$.

3) $* \Pi_{pq} = \Pi_{u-q, u-p} *$

4) $[L, w] = [L, J] = [L^*, w] = [L^*, J] = 0$

5) $[L^*, L] = \sum_{p=0}^{2u} (u-p) \Pi_p$.

Proof: 1), 2) follow by an immediate computation

3) implies that $L^*(\wedge^{r,s} F) \subset \wedge^{r-1, s-1} F$ and hence 4) follows from 3). Finally 3) and 5) are consequences of the next lemma which is proved by tedious but straightforward calculation:

(2.4) Lemma: Suppose that A, B, M are mutually disjoint increasing multiindices (e.g. $A = \{\mu_1 < \mu_2 < \dots < \mu_a\}$ etc.) Let $z_A = z_{\mu_1} \wedge \dots \wedge z_{\mu_a}$ and similarly $\bar{z}_B$. Also set

$w_M = \prod_{\mu \in M} z_{\mu} \wedge \bar{z}_{\mu}$. Then we have:

$$*(z_A \wedge \bar{z}_B \wedge w_M) = i^{a-b} (-1)^{\frac{p(p+1)}{2} + m} (-2i)^{p-m} z_A \wedge \bar{z}_B \wedge w_{M'}$$

where $a = |A|$, $b = |B|$, $m = |M|$, $p = a+b+2m$ and $M' = \{1, \dots, n\} \setminus (A \cup B \cup M)$. (Our multiindices A, B, M are supposed to be subsets of $\{1, \dots, n\}$.)

Let us set $B = \sum_{p=0}^{2m} (n-p) \pi_p$ then we have

the commutation relations which follow from 4), 5):

$$[L^*, L] = B, \quad [B, L^*] = 2L^*, \quad [B, L] = -2L.$$

Hence we obtain a $\mathfrak{sl}_2$ -representation:

$$\rho : \mathfrak{sl}_2 \longrightarrow \text{End}(\wedge F)$$

by setting $\rho(X) = L^*$, $\rho(Y) = L$ and $\rho(H) = B$.

Let π_ρ be the associated representation of SL_2 .

Our main application of the $\mathfrak{sl}_2$ -theory will be to relate the $*$ -operator to the automorphism of $\wedge F$ $\# := \pi_\rho(w)$ for w as in (1.2). We will prepare this result by a lemma. If $\gamma \in \wedge^p F$ we set

$$e(\gamma) \in \text{Hom}(\wedge F, \wedge F), \quad e(\gamma)\varphi = \gamma \wedge \varphi$$

e.g. $L = e(\Omega)$. Note that if γ is a real 1-form

then we have

$$(2.5) \quad e^*(\gamma) := e(\gamma)^* = * e(\gamma) *.$$

Using this one checks that

$$(2.6) \quad [L^*, e(\gamma)] = -J e^*(\gamma) J^{-1}.$$

(2.7) Lemma: let γ be a real 1-form i.e. $\gamma \in (\Lambda^1 F)_{\mathbb{R}}$.

Then

$$\# e(\gamma) \#^{-1} = -i J e^*(\gamma) J^{-1}$$

$$\text{where } \# = \pi_p(w) = \exp\left(\frac{i\pi}{2} \rho(X+Y)\right) = \exp\left(\frac{i\pi}{2} (L^*+L)\right).$$

Proof: For $t \in \mathbb{C}$ let

$$\begin{aligned} e_t(\gamma) &= \exp(it(L^*+L)) \cdot e(\gamma) \cdot \exp(-it(L^*+L)) \\ &= \text{Ad}(\exp it(L^*+L)) \cdot e(\gamma) \\ &= \exp(it \text{ad}(L^*+L)) \cdot e(\gamma) \end{aligned}$$

Since $d\text{Ad} = \text{ad}$ and $\exp$ commutes with taking the derivative, Recall that $\text{ad}(u)v = [u, v]$.

Note that

$$e_0(\gamma) = e(\gamma)$$

$$e_{\frac{\pi}{2}}(\gamma) = \# e(\gamma) \#^{-1}$$

$$\frac{d}{dt} e_t(\gamma) = i(\text{ad}(L^*) + \text{ad} L) e_t(\gamma).$$

The relations

(15)

$$(\text{ad } L^*)^2 = 0, \quad (\text{ad } L) \cdot e(z) = 0 \quad (\text{clear since } \deg \Omega = 2)$$

$$\text{ad}(-B) \cdot e(z) = e(z) \quad (\text{by (1)}) \quad \text{and} \quad [L, L^*] = -B$$

show that

$$\tilde{e}_t(z) := \cos t \, e(z) + i \sin t \, \text{ad}(L^*) e(z)$$

is a solution of

$$\frac{d}{dt} \tilde{e}_t(z) = i (\text{ad}(L^*) + \text{ad}(L)) \tilde{e}_t(z).$$

Since $\tilde{e}_0(z) = e(z)$ we must have $e_t(z) = \tilde{e}_t(z)$ i.e.

$$e_t(z) = \cos t \, e(z) + i \sin t \, \text{ad}(L^*) e(z).$$

In particular

$$\begin{aligned} e_{\frac{\pi}{2}}(z) &= i \, \text{ad}(L^*) e(z) = i [L^*, e(z)] \\ &= -i \int e^*(z) \, J^{-1} \quad \text{by (2.6)}. \end{aligned}$$

This implies the lemma.

Now we come to the main technical result of this section.

(2.8) Main-Lemma: For $\varphi \in \Lambda^p F$ we have:

$$* \varphi = i^{p^2-n} J^{-1} \# \varphi.$$

Proof: The $*$ -operator on ΛF satisfies:

$$(1) \quad * 1 = \text{vol} = \frac{1}{n!} \Omega^n = \frac{1}{n!} L^n(1)$$

(2) For all real 1-forms γ we have on $\Lambda^p F$:

$$* e(\gamma) = (-1)^p e^*(\gamma) * \quad (\text{use (2.5)}) \quad , \quad p \geq 0.$$

Now $*$ is the only $\mathbb{C}$ -linear operator $\Lambda F \rightarrow \Lambda F$ satisfying (1) and (2) since the forms obtained from 1 by repeated application of $e(\gamma)$, γ as above span ΛF .

Let
$$\tilde{*} = i^{p^2-n} J^{-1} \# \quad : \quad \Lambda^p F \rightarrow \Lambda F$$

and extend $\tilde{*}$ to an endomorphism of ΛF .

In the notation of (1.7) the form $v_0 = 1$ is primitive of weight n (since $Bv_0 = nv_0$, $L^*v_0 = 0$) hence by (1.7) for $k=0$ we obtain:

$$\# 1 = \frac{i^n}{n!} L^n(1).$$

Thus

$$\tilde{*} 1 = i^{-n} \left(\frac{i^n}{n!} \right) L^n(1) = \text{vol}.$$

On the other hand for γ as above and $\varphi \in \Lambda^p F$ we have:

$$\begin{aligned} \tilde{*} e(\gamma) \varphi &= i^{(p+1)^2-n} J^{-1} \# e(\gamma) \varphi \\ &= i^{p^2-n} (-1)^p i J^{-1} \# e(\gamma) \#^{-1} \# \varphi \end{aligned}$$

$$\stackrel{(2.7)}{=} i^{p^2-n} (-1)^p e^*(\gamma) J^{-1} \# \varphi$$

$$= (-1)^p e^*(\gamma) \tilde{*} \varphi. \quad \text{Thus (1), (2) hold for } \tilde{*}. \text{ Hence } * = \tilde{*}.$$

We now apply these results to manifolds. Let M be a compact Hermitian manifold. Then by doing the constructions of the last section fibrewise starting from $E = T_{\xi}^{\mathbb{C}} M$, $\xi \in M$ the complex tangent space of M in ξ with its Hermitian metric, we get bundle maps $L : \Lambda^p T^*(M)_{\mathbb{C}} \rightarrow \Lambda^{p+2} T^*(M)_{\mathbb{C}}$, L^* , $B, *, \bar{\partial}, \bar{\partial}^*$ etc. Note that the pointwise fundamental forms Ω_{ξ} glue to a section $\Omega \in \Gamma(M, \Lambda^2 T^*(M)_{\mathbb{C}})$ the fundamental form of M . Also note that the above bundle maps induce maps of sections (differential forms)

$$L : \Gamma(M, \Lambda^p T^*(M)_{\mathbb{C}}) =: A^p(M) \rightarrow A^{p+2}(M) \text{ etc.}$$

It is easy to check that in this interpretation $L^* = \omega * L *$ becomes the adjoint of L with respect to the Hodge inner product.

On forms we also have exterior differentiation $d, \bar{\partial}$ together with their adjoints d^* with respect to the Hodge inner product. We now assume in addition that M is Kähler i.e. that $d\Omega = 0$ or in other words that $[L, d] = 0$. Then we have:

(3.1) Theorem: $\Delta = 2 \bar{\square}$ where $\Delta = dd^* + d^*d$ and

$$\bar{\square} = \bar{\partial}\bar{\partial}^* + \bar{\partial}^*\bar{\partial}.$$

The proof will be based on the relations (16) between d, d^*, L, L^* given by the next lemma.

$$\text{Set } d_c = J^{-1} d J = w J d J$$

$$d_c^* = J^{-1} d^* J = w J d^* J$$

then we have

(3.2) Crucial Lemma: a) $[L, d] = 0$, $[L^*, d^*] = 0$

$$b) [L, d^*] = d_c \quad , \quad [L^*, d] = -d_c^* .$$

Proof: a) is clear since M is Kähler. (Note that $[L^*, d^*] = 0$ is the adjoint assertion to $[L, d] = 0$.)

b) Since $[L, d^*]^* = [d, L]$ it is sufficient to prove $[L^*, d] = -d_c^*$ i.e. that

$$L^* d - d L^* = - J^{-1} d^* J .$$

We know that

$$(3.3) \quad d^* = - * d * = (-1)^{p+1} * d *^{-1} \text{ on } A^p(M) .$$

For $\varphi \in A^p(M)$ we have by (2.8)

$$\begin{aligned} \# d \#^{-1} \varphi &= i J (-1)^{p+1} * d *^{-1} J^{-1} \varphi \\ &\stackrel{(3.3)}{=} i J d^* J^{-1} \varphi = -i J^{-1} d^* J \varphi \end{aligned}$$

Now let

$$d_t = \exp(it(L^* + L)) \cdot d \cdot \exp(-it(L^* + L)) .$$

$$\begin{aligned} \text{Thus } d_t &= \text{Ad}(\exp it(L^* + L)) \circ d \\ &= \exp(it \text{ad}(L^* + L)) \circ d. \end{aligned}$$

Note that:

$$d_0 = d$$

$$(3.4) \quad d_{\frac{\pi}{2}} = \# d \#^{-1} = -i J^{-1} d^* J$$

$$\frac{d}{dt} d_t = i(\text{ad} L^* + \text{ad} L) d_t.$$

The relations

$$(\text{ad} L^*)^2 = 0, \quad \text{ad} L \circ d = [L, d] = 0, \quad \text{ad}(-B) \circ d = d$$

and $[L, L^*] = -B$ show that

$$\tilde{d}_t := \cos t \cdot d + i \sin t \text{ad}(L^*) d$$

is a solution of

$$\frac{d}{dt} \tilde{d}_t = i(\text{ad}(L^*) + \text{ad}(L)) \tilde{d}_t.$$

Since $\tilde{d}_0 = d$ we must have $d_t = \tilde{d}_t$ i.e.

$$d_t = \cos t \cdot d + i \sin t \text{ad}(L^*) d.$$

In particular

$$d_{\frac{\pi}{2}} = i \text{ad}(L^*) d = i [L^*, d]$$

which implies the assertion in view of (3.4).

(3.5) Corollary: 1) $[L, d_c] = 0$, $[L^*, d_c^*] = 0$, $[L, d_c^*] = -d$, $[L^*, d_c] = d^*$. (18)

$$2) \quad d_c d^* + d^* d_c = 0, \quad d d_c^* + d_c^* d = 0$$

Proof: 1) follows from (3.2) by noting that J commutes with the real operators L, L^* .

2) Since $d^* = [L^*, d_c]$ by 1) we have:

$$d^* d_c = L^* d_c d_c - d_c L^* d_c = -d_c L^* d_c$$

$$d_c d^* = d_c L^* d_c - d_c d_c L^* = d_c L^* d_c$$

so $d_c d^* + d^* d_c = 0$. The other equation follows by taking adjoints.

We can now prove Th. (3.1) : $\Delta = 2 \square$ as follows:

$$\begin{aligned} \Delta &= d d^* + d^* d = d [L^*, d_c] + [L^*, d_c] d \\ &= d L^* d_c - d d_c L^* + L^* d_c d - d_c L^* d. \end{aligned}$$

On the other hand

$$J^{-1} \Delta J = -d_c L^* d + d_c d L^* - L^* d d_c + d L^* d_c$$

so $\Delta = J^{-1} \Delta J$ since $d_c d = -d d_c$. On the other hand

$$J^{-1} \Delta J = J^{-1} d d^* J + J^{-1} d^* d J = d_c d_c^* + d_c^* d_c$$

and hence

$$2\Delta = \Delta + J^{-1} \Delta J = d d^* + d^* d + d_c d_c^* + d_c^* d_c.$$

On the other side the relations

$$d_c = i(\bar{\partial} - \partial) \quad , \quad d_c^* = -i(\bar{\partial}^* - \partial^*) \quad \text{imply}$$

$$\bar{\partial} = \frac{1}{2}(d - i d_c) \quad , \quad \bar{\partial}^* = \frac{1}{2}(d^* + i d_c^*)$$

and hence

$$\begin{aligned} 4\bar{\square} &= 4(\bar{\partial}\bar{\partial}^* + \bar{\partial}^*\bar{\partial}) = dd^* + i dd_c^* - i d_c d^* + d_c d_c^* \\ &\quad + d^* d - i d^* d_c + i d_c^* d + d_c^* d_c \end{aligned}$$

$$\stackrel{(3.5) 2)}{=} dd^* + d^* d + d_c d_c^* + d_c^* d_c$$

$$= 2\Delta \quad \text{q.e.d.}$$

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