



SMR.637/7

**ADVANCED WORKSHOP ON ARITHMETIC ALGEBRAIC
GEOMETRY**
(31 August - 11 September 1992)

The Hodge decomposition theorems (I)

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These are preliminary lecture notes, intended only for distribution to participants

The Hodge decomposition Theorem I

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References

- [G, H] P. Griffiths, J. Harris, Principles of Algebraic Geometry, John Wiley 1978
- [Wa] F. W. Warner, Foundations of diff. manifolds and Lie groups, Springer
- [We] R. O. Wells, Diff. analysis on complex manifolds, Springer GTM 65, 1980

In the above references general elliptic operators are treated. The text [We] is more elementary than [G, H] and [Wa]. In [We] the theory is based on pseudo-differential operators which are an interesting topic in itself. We will mostly follow [We].

We begin with generalities about the $*$ -operator first for Hermitian vector spaces and then for Riemannian or Hermitian Manifolds.

Let V be an n -dimensional real vector space with an inner product \langle, \rangle . It induces inner products on all exterior powers $\wedge^p V$ as follows: If e_1, \dots, e_n is an

$e_{i_1} \wedge \dots \wedge e_{i_p}$ for $1 \leq i_1 < \dots < i_p \leq m$ form an ON-basis of $\Lambda^p V$.

Since $\dim_{\mathbb{R}} \Lambda^m V = 1$ there are exactly two ON-bases of $\Lambda^m V$. These are called orientations. They correspond to an ordering of e_1, \dots, e_m up to an even permutation, since $e_{\sigma_1} \wedge \dots \wedge e_{\sigma_m} = \text{sgn}(\sigma) e_1 \wedge \dots \wedge e_m$ for all $\sigma \in \mathfrak{S}_m$.

Fix an orientation $\text{vol} \in \Lambda^m V$. Define the $*$ -operator as follows

$$* : \Lambda^p V \longrightarrow \Lambda^{m-p} V$$

Having fixed an ON-basis of V , e_1, \dots, e_m s.t. $\text{vol} = e_1 \wedge \dots \wedge e_m$ set

$$*(e_{i_1} \wedge \dots \wedge e_{i_p}) = \delta e_{j_1} \wedge \dots \wedge e_{j_{m-p}}$$

where $\{i_1, \dots, i_p, j_1, \dots, j_{m-p}\} = \{1, \dots, m\}$ and $\delta \in \{\pm 1\}$ is s.t.

$$(e_{i_1} \wedge \dots \wedge e_{i_p}) \wedge *(e_{i_1} \wedge \dots \wedge e_{i_p}) = \text{vol}.$$

For the \mathbb{R} -linear extension of $*$ to $\Lambda^p V$ we then have

$$(1.1) \quad \alpha \wedge * \beta = \langle \alpha, \beta \rangle \text{vol} \quad \text{for all } \alpha, \beta \in \Lambda^p V.$$

Setting $V_{\mathbb{C}} = V \otimes_{\mathbb{R}} \mathbb{C}$, \langle, \rangle on V extends to a Hermitian scalar product \langle, \rangle on $V_{\mathbb{C}}$. If $*$ is extended \mathbb{C} -linearly to $V_{\mathbb{C}}$ by $* \otimes \text{id}$ we have

$$(1.2) \quad \alpha \wedge * \bar{\beta} = \langle \alpha, \beta \rangle \text{vol} \quad \text{for all } \alpha, \beta \in \Lambda^p V_{\mathbb{C}} = \Lambda^p V \otimes_{\mathbb{R}} \mathbb{C}.$$

If $*$ is extended \mathbb{C} -antilinearly to $V_{\mathbb{C}}$ by

$\bar{*} \cong * \otimes c$, $c = \text{complex conjugation on } \mathbb{C}$ then

$$(1.3) \quad \alpha \wedge \bar{*} \beta = \langle \alpha, \beta \rangle \text{ vol.}$$

The $*$ -operator is almost an involution:

Set $\Lambda^* V = \bigoplus_{p=0}^m \Lambda^p V$ and denote by

$\pi_p: \Lambda^* V \rightarrow \Lambda^p V$ the p -th projection.

If we set

$$\omega: \Lambda^* V \rightarrow \Lambda^* V, \quad \omega = \sum_{p=0}^m (-1)^{m-p} \pi_p$$

then we have (exercise!)

$$(1.4) \quad * * = \omega.$$

Note that if m is even $\omega = \sum (-1)^p \pi_p$.

Now let M be a compact Riemannian manifold of dimension m . We assume that there exists a (C^∞) section vol of $\Lambda^m T^*(M)$ s.t. for all $\xi \in M$

$\text{vol}_\xi \in \Lambda^m T_\xi^*(M)$ is an orientation (with respect

to the inner product coming from the Riemannian structure).

If such a section vol exists the manifold is called orientable. The orientability of M depends in fact only on the manifold structure of M and not on the choice of a Riemannian metric.

Henceforth we assume M to be oriented and with a choice of orientation vol . We obtain smooth bundle maps

$$* : \Lambda^p T^*(M) \rightarrow \Lambda^{m-p} T^*(M) \quad \text{and} \quad * : \Lambda^p T^*(M)_{\mathbb{C}} \rightarrow \Lambda^{m-p} T^*(M)_{\mathbb{C}}$$

by defining $*$ fibrewise. They give isomorphisms of sections

$$A^p(M) \cong \Gamma(M, \Lambda^p T^*(M)_{\mathbb{C}})$$

$$* : A^p(M) \xrightarrow{\sim} A^{m-p}(M) \quad \text{s.t.} \quad \alpha \wedge * \bar{\beta} = \langle \alpha, \beta \rangle \text{vol}$$

for all $\alpha, \beta \in A^p(M)$. Note that this means

$$\alpha_{\xi} \wedge * \bar{\beta}_{\xi} = \langle \alpha_{\xi}, \beta_{\xi} \rangle_{\xi} \text{vol}_{\xi} \quad \text{for all } \xi \in M.$$

Before going on we have to recall how one integrates m -forms on an orientable m -dimensional ^{compact} manifold:

For $\varphi \in A^m(M)$ define $\int_M \varphi$ as follows:

Fix a finite atlas $\{U_{\alpha}, f_{\alpha} : V_{\alpha} \subset \mathbb{R}^m \xrightarrow{\sim} U_{\alpha} \subset M\}$ and

choose a subordinate partition of unity $\{\psi_{\alpha}\}$ i.e.

$\text{supp } \psi_{\alpha} \subset\subset U_{\alpha}$, $\sum_{\alpha} \psi_{\alpha} = 1$. Set:

$$\int_M \varphi := \sum_{\alpha} \int_{V_{\alpha}} f_{\alpha}^*(\psi_{\alpha} \varphi) = \sum_{\alpha} \int_{\mathbb{R}^m} g_{\alpha}(x) dx$$

where $f_{\alpha}^*(\psi_{\alpha} \varphi) = g_{\alpha}(x) dx_1 \wedge \dots \wedge dx_m$ for some $g_{\alpha} \in C_0^{\infty}(V_{\alpha})$.

The definition of $\int_M \varphi$ is independent of all choices (exercise)!

On $A^*(M) = \bigoplus_{p=0}^m A^p(M)$ we can now define

a Hermitian inner product by

$$(1.5) \quad (\alpha, \beta) := \int_M \alpha \wedge \bar{\beta} = \int_M \langle \alpha, \beta \rangle \text{ vol} \quad \text{for } \alpha, \beta \in A^p(M)$$

$$(\alpha, \beta) := 0 \quad \text{for } \alpha \in A^p(M), \beta \in A^q(M) \text{ with } p \neq q.$$

This is called the Hodge inner product on $A^*(M)$.

If M is a Hermitian complex manifold of real dimension $2m$ define the Hodge inner product with respect to the underlying Riemannian metric and the orientation coming from the complex structure. Set

$$A^{p,q}(M) = (p,q)\text{-forms in } A^{p+q}(M).$$

Prop. (1.6) $A^u(M) = \bigoplus_{p+q=u} A^{p,q}(M)$ is orthogonal with

respect to the Hodge inner product.

Proof: $\alpha \in A^{p,q}(M)$, $\beta \in A^{r,s}(M)$, $p+q = r+s$. Then

$\alpha \wedge \bar{\beta}$ has type $(m-r+p, m-s+q)$. It is a

$2m$ -form iff $\begin{matrix} m-r+p = m \\ m-s+q = m \end{matrix}$ i.e. $p=r, q=s$.

Prop. (1.7) M m -dimensional compact oriented Riemannian manifold. Then exterior derivation $d: A^*(M) \rightarrow A^*(M)$ has an adjoint d^* w.r.t. respect to the Hodge inner product and we have:

i) $d^* = \delta$ on $A^p(M)$ where $\delta = (-1)^{m+mp+1} * d * = (-1)^{m+mp+1} \bar{*} d \bar{*}$

ii) For the operator $\Delta = dd^* + d^*d$ on $A^*(M)$ or $A^p(M)$ we have $(\Delta\alpha, \beta) = (\alpha, \Delta\beta)$, $*\Delta = \Delta*$, $\bar{*}\Delta = \Delta\bar{*}$

Remark: The operator Δ is the Laplace operator on M .

Ex. Using the isomorphism $S^1 \cong \mathbb{R}/2\pi\mathbb{Z}$, $e^{2\pi i\theta} \leftrightarrow \theta$ the standard metric on \mathbb{R} induces a Riemannian structure on S^1 . Let $T^m = S^1 \times \dots \times S^1$ with the induced Riemannian structure.

Then we have for $\alpha \in A^0(T^m)$

$$\Delta\alpha = -\left(\frac{\partial^2}{\partial\theta_1^2} + \dots + \frac{\partial^2}{\partial\theta_m^2}\right)\alpha. \quad (\text{exercise!})$$

(The orientation chosen is the "obvious" one).

Proof of (1.7): i) For $\alpha \in A^{p-1}(M)$, $\beta \in A^p(M)$ we have

$$\begin{aligned} d(\alpha \wedge * \bar{\beta}) &= d\alpha \wedge * \bar{\beta} + (-1)^{p-1} \alpha \wedge d(* \bar{\beta}) \\ &\stackrel{(1.4)}{=} d\alpha \wedge * \bar{\beta} - \alpha \wedge * \delta \bar{\beta}. \end{aligned}$$

Hence

$$0 = \int_M d(\alpha \wedge * \bar{\beta}) = (d\alpha, \bar{\beta}) - (\alpha, \delta \bar{\beta}) \Rightarrow \text{i), rest: exercise.}$$

States theorem! $\partial M = \emptyset$

Notation: the kernel of Δ on $A^p(M)$

$$\mathcal{H}^p(M) := \text{Ker}(\Delta : A^p(M) \rightarrow A^p(M))$$

is called the space of "harmonic p -forms on M ".

Ex.: For $p=0$ we have $\Delta = d^*d$ on $A^0(M)$. Hence

$$\Delta \alpha = 0 \Rightarrow (\Delta \alpha, \alpha) = 0 \Rightarrow \|d\alpha\|^2 = 0 \Rightarrow d\alpha = 0 \Rightarrow \alpha \text{ constant}$$

(M being connected). It follows that

$$\mathcal{H}^0(M) = \mathbb{C}$$

For the $\bar{\partial}$ -operator there is the entirely analogous result

Prop. (8.8) M compact Hermitian complex manifold. Then

$$\bar{\partial} : A^{p,q}(M) \rightarrow A^{p,q+1}(M)$$

has an adjoint $\bar{\partial}^*$ with respect to the Hodge inner product and we have

$$\bar{\partial}^* = -\bar{*} \bar{\partial} \bar{*}$$

If $\bar{\Delta} := \bar{\partial} \bar{\partial}^* + \bar{\partial}^* \bar{\partial}$ denotes the "complex Laplacian"

we have: $\bar{\Delta} \bar{*} = \bar{*} \bar{\Delta}$ and $(\bar{\Delta} \alpha, \beta) = (\alpha, \bar{\Delta} \beta)$.

Notation: The kernel of $\bar{\Delta}$ on $A^{p,q}(M)$

$$\mathcal{H}^{p,q}(M) := \text{Ker}(\bar{\Delta} : A^{p,q}(M) \rightarrow A^{p,q}(M))$$

is called the space of "harmonic (p,q) -forms on M ".

We are now in the position to state the famous Hodge decomposition theorem for forms:

Th. (1.9) Let M be as in (1.7). For each p the space $\mathcal{H}^p(M)$ is finite dimensional and we have the following orthogonal sum decompositions of $A^p(M)$:

$$\begin{aligned} A^p(M) &\stackrel{(1)}{=} \mathcal{H}^p(M) \oplus \Delta A^p(M) \\ &\stackrel{(2)}{=} \mathcal{H}^p(M) \oplus d d^* A^p(M) \oplus d^* d A^p(M) \\ &\stackrel{(3)}{=} \mathcal{H}^p(M) \oplus d A^{p-1}(M) \oplus d^* A^{p+1}(M). \end{aligned}$$

In particular the equation $\Delta \omega = \alpha$ has a solution $\omega \in A^p(M)$ iff $\alpha \in A^p(M)$ is orthogonal to $\mathcal{H}^p(M)$.

Th. (1.10) Let M be as in (1.8). For all p, q the spaces $\mathcal{H}^{p,q}(M)$ are finite dimensional and we have orth. sum decompositions:

$$\begin{aligned} A^{p,q}(M) &= \mathcal{H}^{p,q}(M) \oplus \bar{\square} A^{p,q}(M) \\ &= \mathcal{H}^{p,q}(M) \oplus \bar{\partial} \bar{\partial}^* A^{p,q}(M) \oplus \bar{\partial}^* \bar{\partial} A^{p,q}(M) \\ &= \mathcal{H}^{p,q}(M) \oplus \bar{\partial} A^{p,q-1}(M) \oplus \bar{\partial}^* A^{p,q+1}(M). \end{aligned}$$

In particular the equation $\bar{\square} \omega = \alpha$ has a solution $\omega \in A^{p,q}(M)$ iff $\alpha \in A^{p,q}(M)$ is orthogonal to $\mathcal{H}^{p,q}(M)$.

Remark: There is in fact a general Hodge decomposition theorem for any "elliptic operator" on M : Δ and $\bar{\square}$ are examples of elliptic operators. The proof which we will sketch roughly below for Δ carries over in fact to elliptic operators.

Elementary remarks on (1.9): We have the important fact:

(1.11) Fact: $\Delta \alpha = 0 \iff d\alpha = 0$ and $d^* \alpha = 0$

Proof: " \Leftarrow " clear, " \Rightarrow " $0 = (\Delta \alpha, \alpha) = (dd^* \alpha, \alpha) + (d^* d \alpha, \alpha)$
 $= (d\alpha, d\alpha) + (d^* \alpha, d^* \alpha) = \|d\alpha\|^2 + \|d^* \alpha\|^2$
 $\implies d\alpha = 0, d^* \alpha = 0.$

Now it is clear that the decompositions in (1.9) are orthogonal (also use $d^2=0$) and hence (2) and (3) follow from (1).

(1.12) Def.: Let $\mathcal{H} : A^p(M) \rightarrow \mathcal{H}^p(M)$ denote the (orthogonal) projection with respect to the Hodge decomposition. The Green's operator

$$G : A^p(M) \rightarrow \mathcal{H}^p(M)^\perp$$

is defined by $G(\alpha) = \omega$ where ω is the unique (!) solution of $\Delta \omega = \alpha - \mathcal{H}(\alpha)$ in $\mathcal{H}^p(M)^\perp$.

Similarly for the $\bar{\square}$ -Hodge decomposition (1.10)

Theorems (1.9) and (1.10) have beautiful consequences for the study of de Rham and Dolbeault cohomology

Recall that

$$H_{dR}^p(M, \mathbb{C}) = \frac{\text{Ker}(d: A^p(M) \rightarrow A^{p+1}(M))}{\text{Im}(d: A^{p-1}(M) \rightarrow A^p(M))} = \frac{\{d\text{-closed } p\text{-forms}\}}{\{d\text{-exact } p\text{-forms}\}}$$

If a form $\beta \in d^* A^{p+1}(M)$ is d -closed $d\beta = 0$ then $\beta = 0$ (since $\beta = d^* \alpha \Rightarrow \|\beta\|^2 = (d^* \alpha, d^* \alpha) = (dd^* \alpha, \alpha) = (d\beta, \alpha) = 0$)

Hence the Hodge decomposition (3) of Th. (1.9) implies using (1.11) that

$$\text{Ker}(d: A^p(M) \rightarrow A^{p+1}(M)) = \mathcal{H}^p(M) \oplus dA^{p-1}(M)$$

Hence the map

$$(1.13) \quad \mathcal{H}^p(M) \xrightarrow{\sim} H_{dR}^p(M, \mathbb{C}), \quad \omega \mapsto [\omega]$$

is an isomorphism with inverse $[\alpha] \mapsto \mathcal{H}(\alpha)$.

Every dR -cohomology class then contains a unique harmonic representative. In particular (1.13) implies

$$\dim_{\mathbb{C}} H_{dR}^p(M, \mathbb{C}) < \infty !$$

Entirely analogous remarks apply to Dolbeault cohomology:

Setting

$$H_{\bar{\partial}}^{p,q}(M) := \frac{\text{Ker}(\bar{\partial} : A^{p,q}(M) \rightarrow A^{p,q+1}(M))}{\text{Im}(\bar{\partial} : A^{p,q-1}(M) \rightarrow A^{p,q}(M))}$$

we have canonically

$$H^p(M) \cong H_{\bar{\partial}}^{p,0}(M)$$

and in particular the $H_{\bar{\partial}}^{p,q}(M)$ are finite dimensional. Note that by Dolbeault's theorem we have

$$H^q(M, \Omega^p) \cong H_{\bar{\partial}}^{p,q}(M).$$

In our second section we give some indications about the proof of (1.9). We begin with Sobolev spaces and L^2 -spaces.

Let M be a compact oriented Riemannian manifold of dimension n and $E \rightarrow M$ a Hermitian vector bundle [in applications to Δ , $E = \wedge^p T^*(M)_\mathbb{C}$]

Set $(\alpha, \beta) = \int_M \langle \alpha, \beta \rangle_E \text{vol}$ for sections α, β of E

[in appl. to Δ this is the Hodge inner product]

Let $L^2(M, E)$ be the completion of the space of $(C^\infty-)$ sections $\Gamma(M, E)$ with respect to $(,)$.

This is a Hilbert space.

Define Sobolev norms $\| \cdot \|_{E, s}$ for any $s \in \mathbb{R}$ on $\Gamma(M, E)$ as follows:

Let $\{U_i, f_i, f_{E, i}\}$ be a finite trivializing cover for M, E :

$$\begin{array}{ccc} V_i \times \mathbb{C}^r & \xrightarrow{\sim} & E|U_i & ; \quad r = \text{rank } E \\ \downarrow & f_{E, i} & \downarrow & \\ \mathbb{R}^m \supset V_i & \xrightarrow{\sim} & U_i \subset M & \\ & f_i & & \end{array}$$

and let $f_i^* : \Gamma(U_i, E) \cong C^\infty(V_i)^r$, $\sigma \mapsto f_{E, i}^{-1} \circ \sigma \circ f_i$ be the induced map. Let $\{\psi_i\}$ be a partition of unity subordinate to $\{U_i\}$ and set for $\alpha \in \Gamma(M, E)$:

$$\| \alpha \|_{E, s}^2 = \sum_i \underbrace{\| f_i^*(\psi_i \alpha) \|_{\mathbb{R}^m, s}^2}_{\text{has compact support in } \mathbb{R}^m}$$

Here for $f: \mathbb{R}^m \rightarrow \mathbb{C}^r$ a C^∞ -function with compact support

$$\| f \|_{\mathbb{R}^m, s}^2 = \sum_{j=1}^m \| f_j \|_{\mathbb{R}^m, s}^2$$

where finally for $f \in C_0^\infty(\mathbb{R}^m)$ one sets:

$$(2.1) \quad \|f\|_{\mathbb{R}^m, s}^2 := \int_{\mathbb{R}^m} |\hat{f}(y)|^2 (1+|y|^2)^s dy \quad ; \quad |y|^2 = \sum_{j=1}^m y_j^2$$

$$\text{where } \hat{f}(y) = (2\pi)^{-m} \int_{\mathbb{R}^m} e^{-i\langle x, y \rangle} f(x) dx$$

is the Fourier transform.

Remark: On $C_0^\infty(\mathbb{R}^m)$ the $\|\cdot\|_{\mathbb{R}^m, s}$ -norm for $s=0, 1, \dots$ is equivalent to $\|\cdot\|'_{\mathbb{R}^m, s}$ defined by:

$$(2.1.1) \quad \|f\|'_{\mathbb{R}^m, s} := \left(\sum_{|\delta| \leq s} \int_{\mathbb{R}^m} |D^\delta f|^2 dx \right)^{1/2}$$

This follows from: $\widehat{D^\delta f}(y) = y^\delta \hat{f}(y)$, $\|f\|_2 = \|\hat{f}\|_2$

where $y^\delta = y_1^{\delta_1} \dots y_m^{\delta_m}$, $D^\delta = (-i)^{|\delta|} D_1^{\delta_1} \dots D_m^{\delta_m}$, $D_j = \frac{\partial}{\partial x_j}$, $|\delta| = \sum_{j=1}^m \delta_j$.

Proof is an exercise! Start with proving the inequality:

$$(1+|y|^2)^s \leq \sum_{|\delta| \leq s} y^{2\delta} \leq c (1+|y|^2)^s \text{ for suitable } c > 0.$$

The norms $\|\cdot\|_{E, s}$ actually depend on our choices. However due to the compactness of M different choices of trivializations $(U_i, f_{E, i}, f_i)$ and partitions of unity give rise to equivalent norms. We fix such choices in the following. Every norm comes from a scalar product (which?)
 Completing $\Gamma(M, E)$ with respect to $\|\cdot\|_{E, s}$ therefore defines a Hilbert space $W^s(M, E)$ for every $s \in \mathbb{R}$. There

is a natural isomorphism of topological vector spaces

$$(2.2) \quad L^2(M, E) = W^0(M, E)$$

such that the norms $\|\cdot\|_2 = \sqrt{(\cdot, \cdot)}$ and $\|\cdot\|_{E,0}$ are equivalent.

Note that since $\|\cdot\|_{E,s} \geq \|\cdot\|_{E,\epsilon}$ on $\Gamma(M, E)$ for $s \geq \epsilon$ there are natural inclusions

$$(2.3) \quad W^s(M, E) \hookrightarrow W^\epsilon(M, E) \quad \text{for } s \geq \epsilon$$

(2.4) Proposition: For every $s \in \mathbb{R}$ there is a natural argument continuous pairing \mathbb{C} -linear in the first \mathbb{C} -antilinear in the second

$$(\cdot, \cdot) : W^s(M, E) \times W^{-s}(M, E) \rightarrow \mathbb{C}$$

which identifies $W^{-s}(M, E)$ with the space $W^s(M, E)^*$ of \mathbb{C} -antilinear continuous linear forms on $W^s(M, E)$.

Proof: is reduced to pairing $f, g \in C_0^\infty(\mathbb{R}^m)$ as follows:

$$(f, g) = \int_{\mathbb{R}^m} f(x) \overline{g(x)} dx = \int_{\mathbb{R}^m} \hat{f}(y) \overline{\hat{g}(y)} dy$$

Hence

$$|(f, g)| \leq \int_{\mathbb{R}^m} |\hat{f}(y)| (1+|y|^2)^{s/2} |\hat{g}(y)| (1+|y|^2)^{-s/2} dy$$

$$\leq \|f\|_{\mathbb{R}^m, s} \|g\|_{\mathbb{R}^m, -s}$$

Thus the nat. \mathbb{C} -linear, \mathbb{C} -antilinear pairing

$$(\cdot, \cdot) : \Gamma(M, E) \times \Gamma(M, E) \rightarrow \mathbb{C}$$

induces a non-degenerate continuous pairing

$$(\cdot, \cdot) : W^s(M, E) \times W^s(M, E) \rightarrow \mathbb{C}.$$

Since the $W^s(M, E)$ are Hilbert spaces we conclude that $W^{-s}(M, E) = W^s(M, E)^*$.

The two basic results about Sobolev spaces are the following:

(2.5) Theorem (Sobolev) If $s > [\frac{m}{2}] + k + 1$ then

$$(2.6) \quad W^s(M, E) \subset \Gamma^k(M, E) := \begin{array}{l} \text{space of } k\text{-times cont.} \\ \text{diff. sections of } E \end{array}$$

and in particular

$$(2.7) \quad \bigcap_s W^s(M, E) = \Gamma(M, E) \quad \text{the } C^\infty\text{-sections of } E.$$

These assertions need some explanation: For any $s, t \in \mathbb{R}$ with $s \geq t$ there is the natural embedding (2.3)

$$W^s(M, E) \hookrightarrow W^t(M, E).$$

In particular for every $s \geq 0$

$$(2.8) \quad W^s(M, E) \subset W^0(M, E) = L^2(M, E).$$

Hence every element of $W^s(M, E)$ can be viewed as an equivalence class of measurable sections up to almost everywhere equality (with respect to the measure on Borel sets of M defined by $\mu(B) = \int_M \chi_B(x) \text{ vol}$). Assertion (2.6)

says that in every such equivalence class there is a unique representative in $\Gamma^k(M, E)$. Using (2.6) the meaning of (2.7) is clear as well.

For the proof see e.g. [We] IV (1.1)'.
:

The second important result is this:

(2.9) Theorem (Rellick): The natural inclusion

$$j: W^s(M, E) \hookrightarrow W^t(M, E) \quad \text{for } s > t$$

is a compact operator i.e. j maps bounded sets to relatively compact sets.

For the proof compare e.g. [Luk] IV (1.2)'.
:

The spaces $\Gamma(M, E)$ are by definition dense in $W^s(M, E)$.

Let

$$\text{OP}_k(E) = \text{"operators of order } k \text{ on } E\text{"}$$

be the space of \mathbb{C} -linear mappings

$$L: \Gamma(M, E) \rightarrow \Gamma(M, E)$$

which have a continuous extension

$$L_s: W^s(M, E) \rightarrow W^{s-k}(M, E)$$

for all s .

(2.10) Prop.: Any $L \in \text{OP}_k(E)$ has an adjoint $L^* \in \text{OP}_k(E)$

and $(L^*)_s = L_{k-s}^*$.

Proof: $L_s : W^s(M, E) \rightarrow W^{s-k}(M, E)$ induces

$$L_s^* : W^{s-k}(M, E)^* \rightarrow W^s(M, E)^*$$

$$\begin{array}{ccc} \parallel & (2.4) & \parallel \\ W^{k-s}(M, E) & \rightarrow & W^{-s}(M, E) \end{array}$$

Hence $Q_s := L_{k-s}^* : W^s(M, E) \rightarrow W^{s-k}(M, E)$.

Uniqueness of adjoints implies that

$$Q_s|_{W^t(M, E)} = Q_t \quad \text{for } t \geq s.$$

Hence $Q|_{\Gamma(M, E)}$ takes values in $\bigcap_s W^{s-k}(M, E) \stackrel{(2.5)}{=} \Gamma(M, E)$.

Clearly $Q = L^*$ on $\Gamma(M, E)$ and $L^* \in \text{OP}_k(E)$.

In the following we fix a $0 \leq p \leq m$ and set $E = \Lambda^p T^*(M)_\mathbb{C}$

and $W^s(M) := W^s(M, \Lambda^p T^*(M)_\mathbb{C})$. It is easy to see that

$\Delta : A^p(M) \rightarrow A^p(M)$ induces operators $\Delta_s : W^s(M) \rightarrow W^{s-2}(M)$

which are continuous. (For $s = 2, 3, \dots$ this is obvious by using (2.1.1); the general case is also not difficult)

Hence $\Delta \in \text{OP}_2 := \text{OP}_2(\Lambda^p T^*(M)_\mathbb{C})$.

According to (2.10) and the symmetry (1.7)ii) of Δ

we get

$$(2.11) \quad \Delta_{2-s}^* = \Delta_s$$

Using e.g. pseudodifferential operators as in [We] IV

one can show:

(2.12) Theorem: There exists a parametrix (or pseudo-inverse) for Δ i.e. an operator $P \in OP_{-2}$ s.t.

$$P \circ \Delta - \text{id} \in OP_{-1}$$

$$\Delta \circ P - \text{id} \in OP_{-1}.$$

In fact these techniques establish an analogue of (2.12) for general elliptic operators on M (including $\Delta, \bar{\Delta}$).

We need

(2.13) Prop.: Let $S \in OP_{-1}$. Then S is a compact operator, i.e. for all extensions $\tilde{S}_s : W^s(M) \rightarrow W^s(M)$ we have: \tilde{S}_s is compact.

Proof: This follows from the commutative diagram

$$\begin{array}{ccc} W^s(M) & \xrightarrow{\tilde{S}_s} & W^s(M) \\ S_s \searrow & & \nearrow j \\ & W^{s+1}(M) & \end{array}$$

using the compactness of j (Rellick's Lemma (2.9)).

Recall the notation

$$\mathcal{H}^P(M) = \text{Ker}(\Delta : A^P(M) \rightarrow A^P(M))$$

and let $\mathcal{H}^P(M)^\perp$ be the orthogonal complement in $L^2(M, E)$, $E = \wedge^P T^*M_\mathbb{C}$.

Before we can prove that $\mathcal{H}^P(M)$ is finite dimensional we need the following useful result from functional analysis.

(2.14) Prop.: Let \mathcal{B} be a Banach space, $S : \mathcal{B} \rightarrow \mathcal{B}$ a compact linear operator. Then setting $T = \text{id} - S$ one has

i) $\text{Ker } T$ is finite dimensional

ii) $T(\mathcal{B})$ is closed in \mathcal{B} and $\text{Coker } T = \mathcal{B}/T(\mathcal{B})$ is finite dimensional.

[An operator satisfying i), ii) is called a Fredholm operator]

For the proof of (2.14) in the case of Hilbert spaces we refer to [We] IV (4.6).

We have the following corollary of (2.12), (2.13) and (2.14):

(2.15) Thm.:

For any parametrix P as in (2.12)

the operators $\Delta \circ P$ and $P \circ \Delta$ have continuous extensions as Fredholm operators $W^s(M) \rightarrow W^s(M)$ for each $s \in \mathbb{R}$.

We can now prove the following basic fundamental

Theorem:

(2.16) Thm. 1 Setting

$$\mathcal{H}^p(M)_s = \text{Ker}(\Delta_s : W^s(M) \rightarrow W^{s-2}(M))$$

we have:

i) $\mathcal{H}^p(M)_s = \mathcal{H}^p(M)$ for all $s \in \mathbb{R}$

ii) Δ_s is a Fredholm operator, in particular

$\mathcal{H}^p(M)_s = \mathcal{H}^p(M)$ is finite dimensional.

Proof: First we show that $\dim \mathcal{H}^p(M)_s < \infty$ for all s .

Let P be a parametrix for Δ as in (2.12). By (2.15)

we know that the extension

$$(P \circ \Delta)_s : W^s(M) \rightarrow W^s(M)$$

has finite dimensional kernel. On the other hand

because of the commutative diagram

$$W^s(M) \xrightarrow{(P \circ \Delta)_s} W^s(M)$$

$$\begin{array}{ccc} & \searrow & \nearrow \\ \Delta_s & & P_{s-2} \\ & \searrow & \nearrow \\ & W^{s-2}(M) & \end{array}$$

we have $\text{Ker} \Delta_s \subset \text{Ker} (P \circ \Delta)_s$ hence $\dim \mathcal{H}^p(M)_s =$

$= \dim \text{Ker} \Delta_s < \infty$. Similarly one sees that Δ_s has

finite dimensional cokernel. This already implies

that Δ_s is Fredholm (the closedness of $\text{Im} \Delta_s$ being

automatic). [Hörmander - Schögl, Einführung in die P.D.-analysis

The proof of 2) will follow from the more general statement about regularity of the Δ -operator:

(2.17) Thm.: Suppose that for $\omega \in W^s(M)$ we have

$$\Delta_s \omega \in A^p(M)$$

then $\omega \in A^p(M)$. In particular $\mathcal{H}^p(M)_s = \text{Ker } \Delta_s \subset A^p(M)$.

Proof: For a parametrix P for Δ we have

$$P \circ \Delta - \text{id} = S \in \text{OP}_{-1}.$$

Now $\Delta \omega \in A^p(M)$ implies that $(P \circ \Delta)(\omega) \in A^p(M)$

(since $P(A^p(M)) \subset \bigcap_{s'} W^{s+2}(M) \stackrel{(2.5)}{=} A^p(M)$). Hence from

$$\omega = (P \circ \Delta)(\omega) - S(\omega)$$

and $S(\omega) \in W^{s+1}(M)$ (recall that $S \in \text{OP}_{-1}$) we

find: $\omega \in W^{s+1}(M)$. Arguing inductively we find:

$$\omega \in \bigcap_{t \geq s} W^t(M) = A^p(M) \text{ by (2.5).}$$

We can now prove the Hodge decomposition theorem (1.9).

We have to show that

$$A^p(M) = \mathcal{H}^p(M) \oplus \Delta A^p(M).$$

As $\Delta_2 : W^2(M) \rightarrow W^0(M) = L^2(M, E)$

is Fredholm by (2.16) ii) we know that $\text{Im } \Delta_2$

(3) closed. Hence

$$L^2(M, E) = (\text{Im } \Delta_2)^\perp \oplus \text{Im } \Delta_2 .$$

Now clearly $(\text{Im } \Delta_2)^\perp = \text{Ker } \Delta_2^*$

$$\stackrel{(2.11)}{=} \text{Ker } \Delta_0 = \mathcal{H}^p(M) .$$

$$\stackrel{(2.16) i)}{=} \mathcal{H}^p(M) .$$

Hence

$$L^2(M, E) = \mathcal{H}^p(M) \oplus \text{Im } \Delta_2 ,$$

and since $\mathcal{H}^p(M) \subset A^p(M)$ we get:

$$A^p(M) = \mathcal{H}^p(M) \oplus (\text{Im } \Delta_2 \cap A^p(M)) .$$

The regularity theorem (2.17) shows that

$$\text{Im } \Delta_2 \cap A^p(M) = \text{Im } \Delta$$

and hence we get

$$A^p(M) = \mathcal{H}^p(M) \oplus \Delta A^p(M)$$

i.e. the assertion.

