



H4.SMR/645-14

**SCHOOL ON PHYSICAL METHODS FOR THE
STUDY OF THE UPPER AND LOWER
ATMOSPHERE SYSTEM**

26 October - 6 November 1992
Miramare - Trieste, Italy

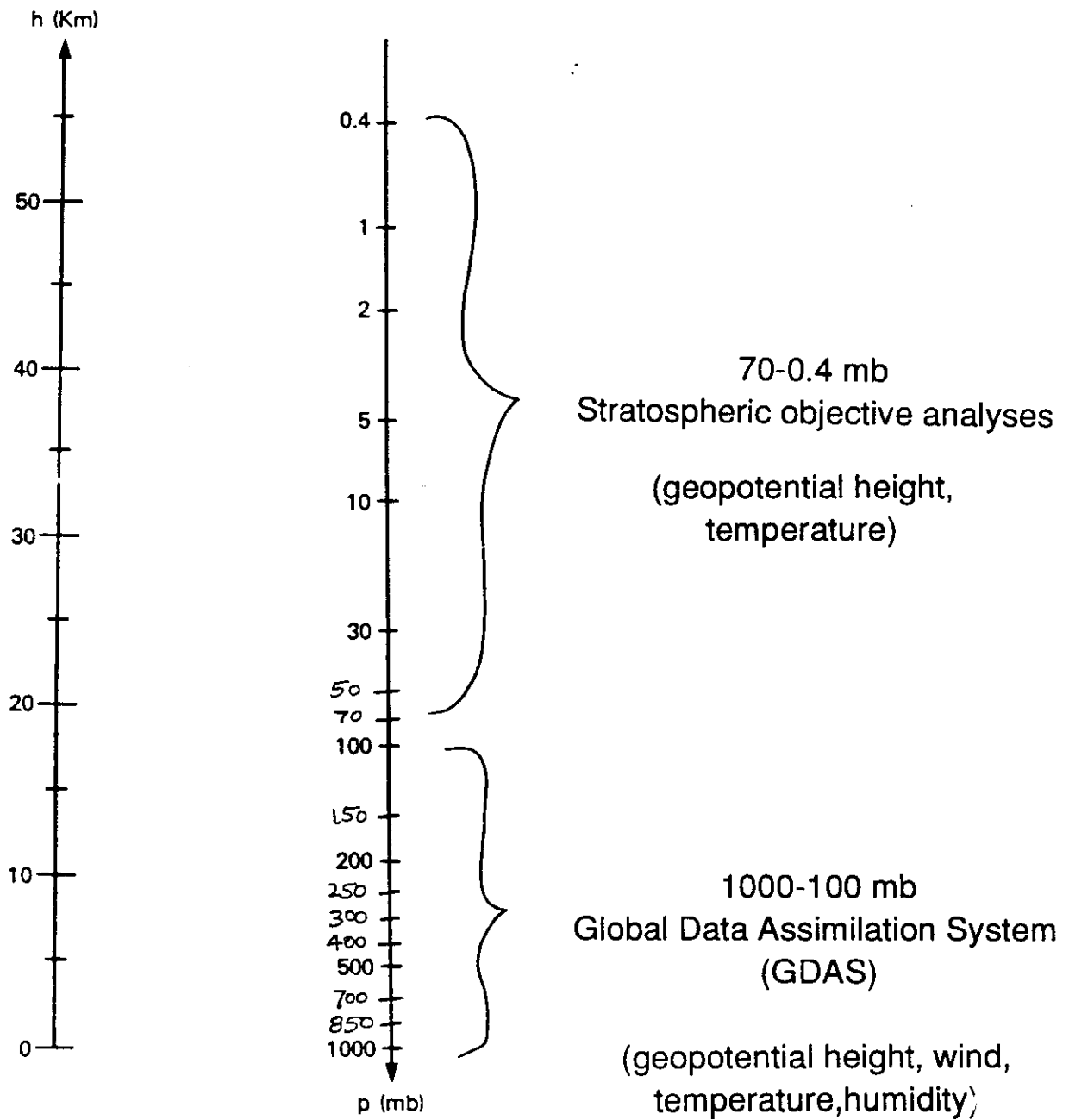
National Meteorological Center (NMC)
daily operational analyses
1000-1 mb (0-50km)

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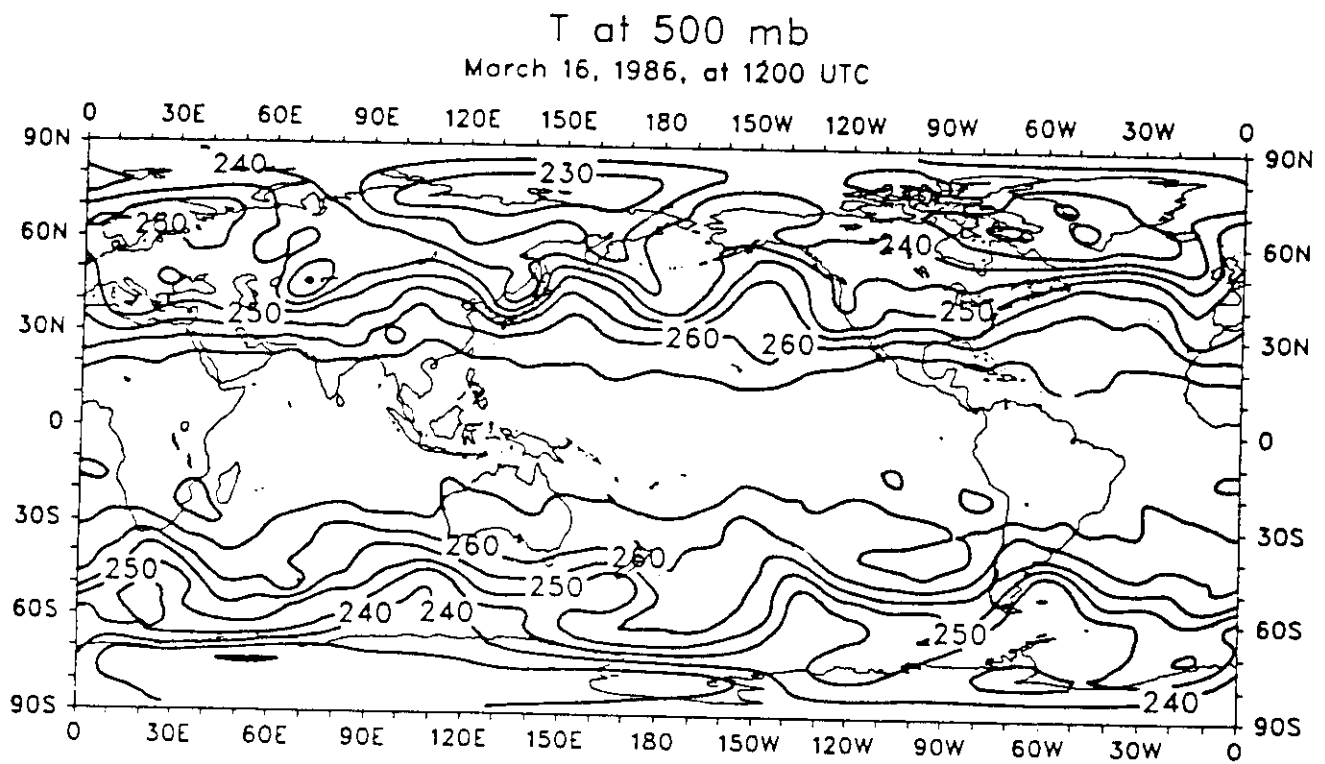
National Meteorological Center (NMC)
daily operational analyses
1000-1 mb (0-50 km)

- combination of NMC Global Data Assimilation System
tropospheric analyses (1000-100 mb) and NMC operational
stratospheric analyses (70-1 mb)
- data origin, characteristics, biases and changes with time
- example of calculations with geopotential height data:
 - temperatures
 - winds
 - eddy flux quantities
 - vertical propagation of planetary waves

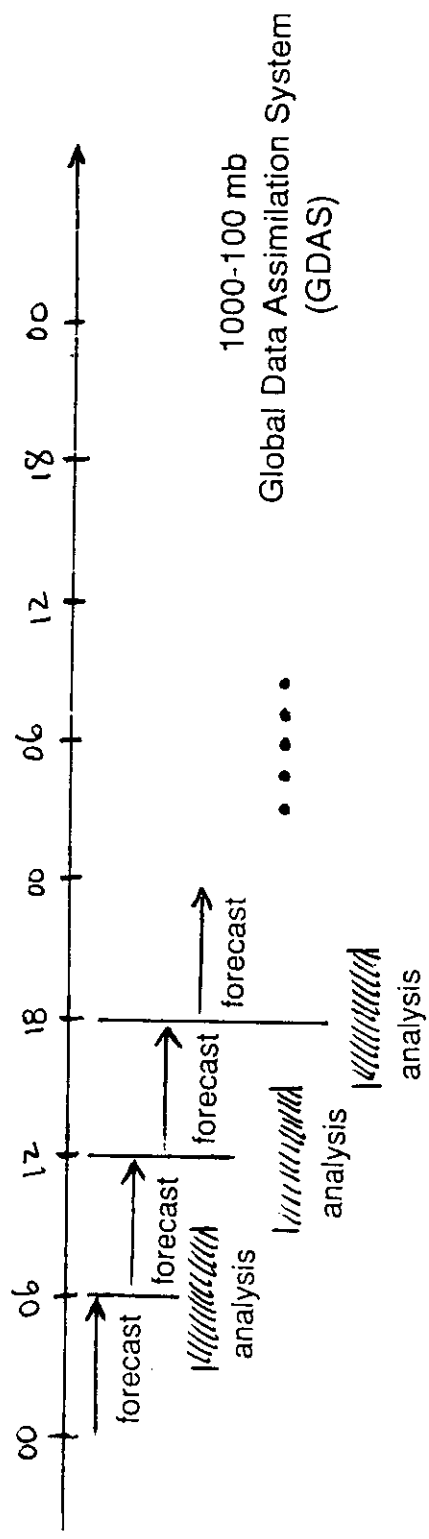
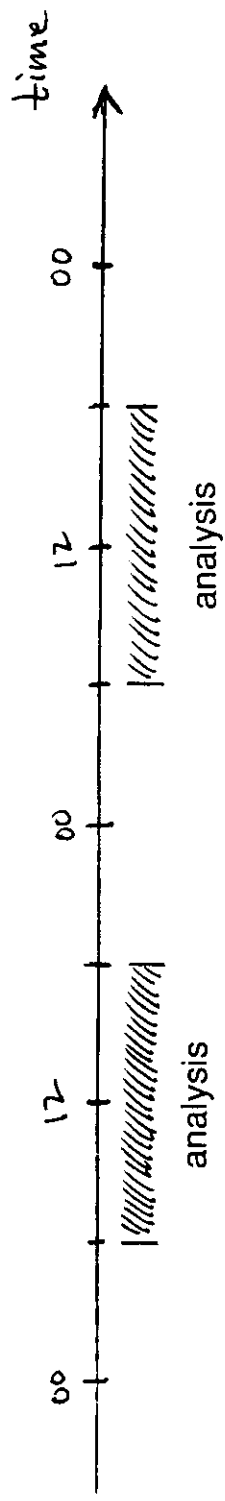
National Meteorological Center (NMC) daily operational analyses

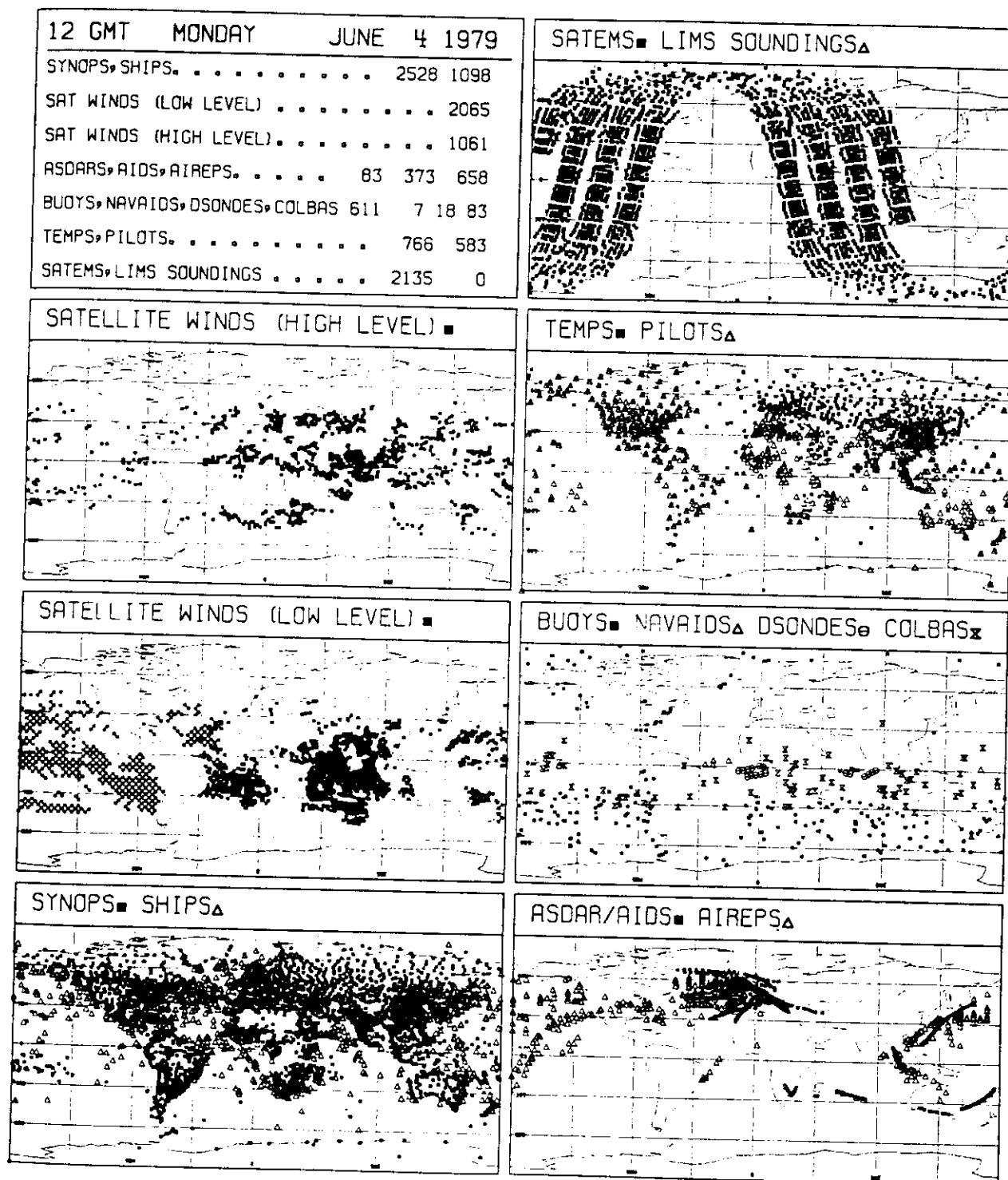


Example: global temperature analysis at 500 mb



70-0.4 mb Stratospheric objective analyses



FIG. 4. FGGE level II-5 data distribution. 4 June 1979 12 GMT ± 3 h.

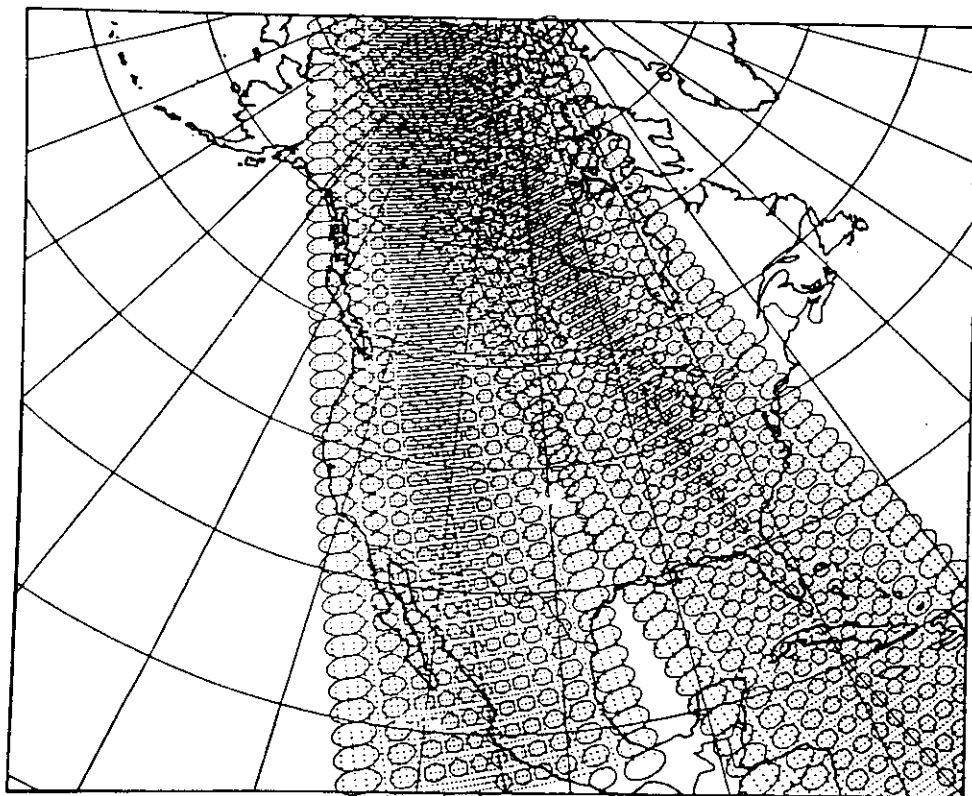


FIG. 2. HIRS and MSU scan patterns for two consecutive orbits.

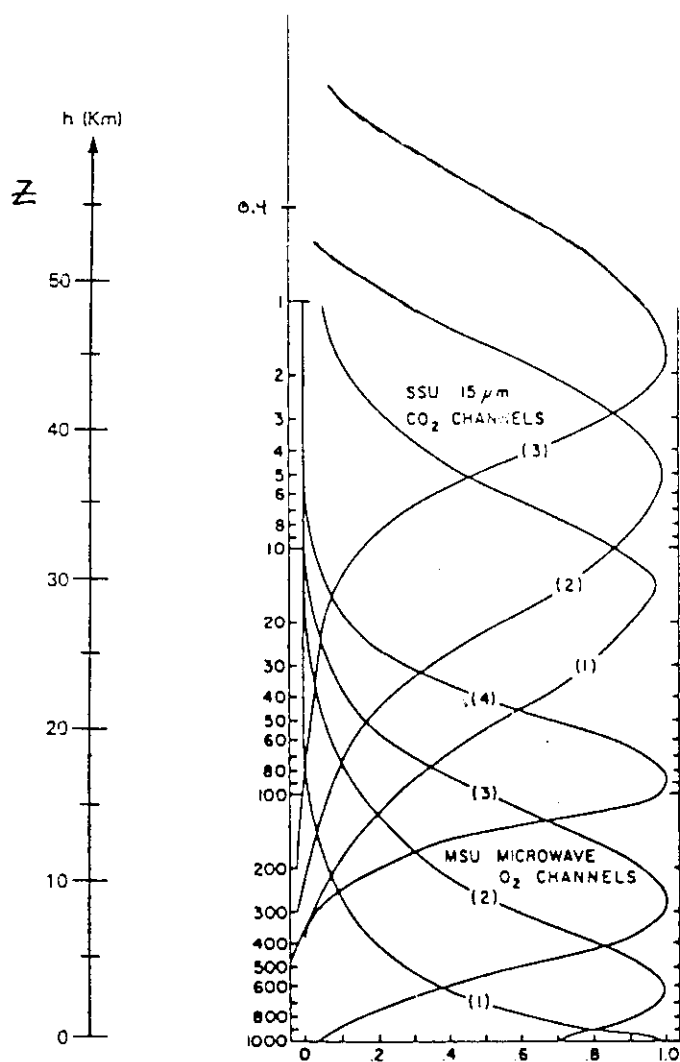
Horizontal scan patterns for stratospheric
temperature sounders

Some cautionary notes regarding use of NMC operational data

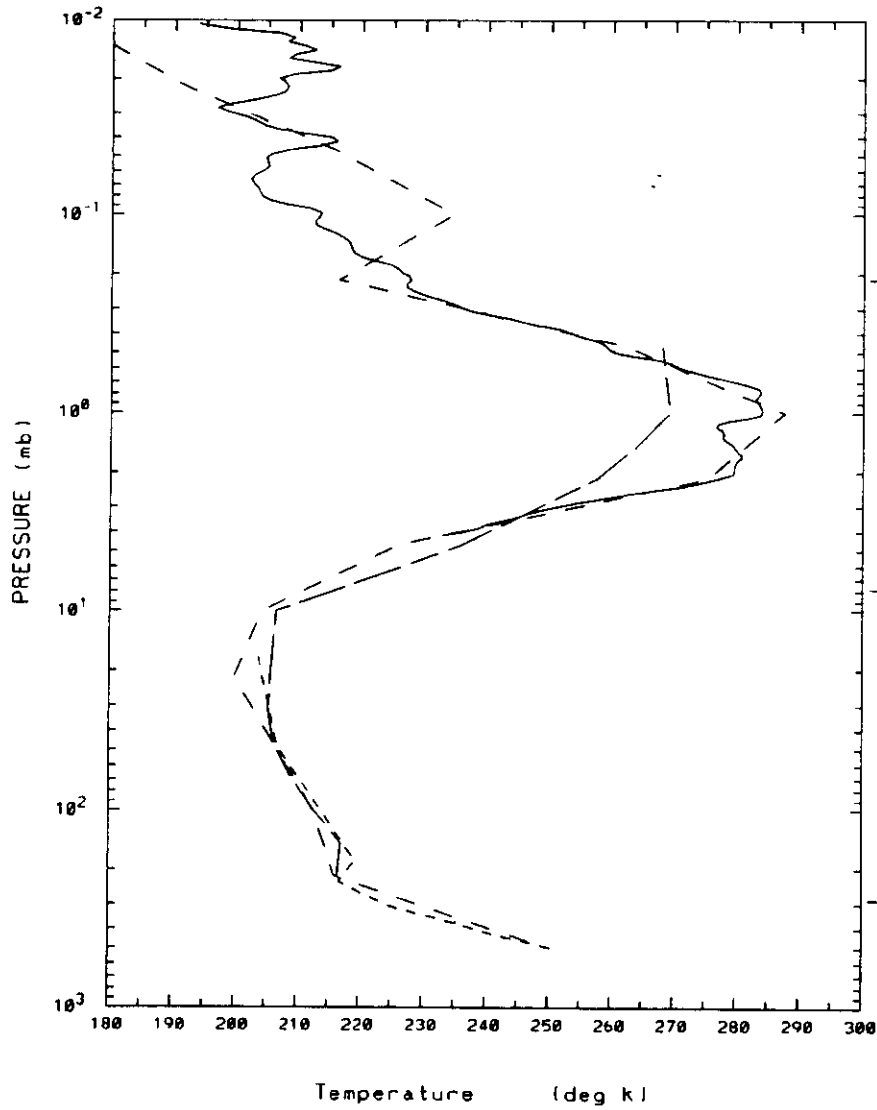
- 1) stratospheric temperature biases, due to low vertical resolution of satellite measurements
- 2) changes in products occur when forecast model changes (troposphere) or satellites change (stratosphere)

Satellite weighting functions $W(z)$ used for
stratospheric temperature retrievals

$$\text{SATELLITE MEASURED TEMPERATURE} = \int_0^{\infty} T(z) \cdot W(z) dz$$



UARS Temperature



(LIDAR)

OHP
44.00 6.00
12/14/91
20.43.0. 0

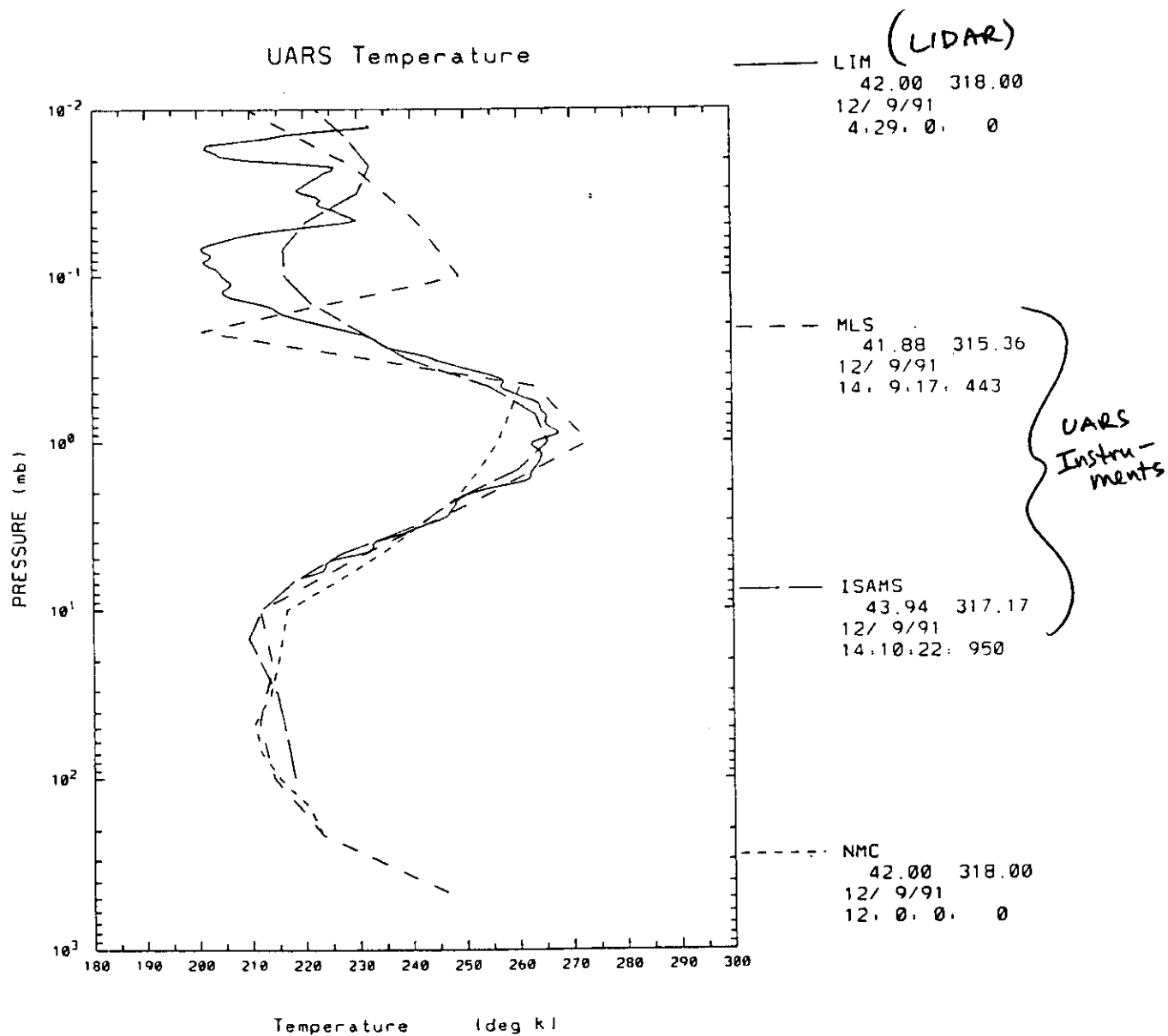
(Microwave
Limb
Sounder on UARS)

MLS
44.32 6.06
12/14/91
21.15.30. 1

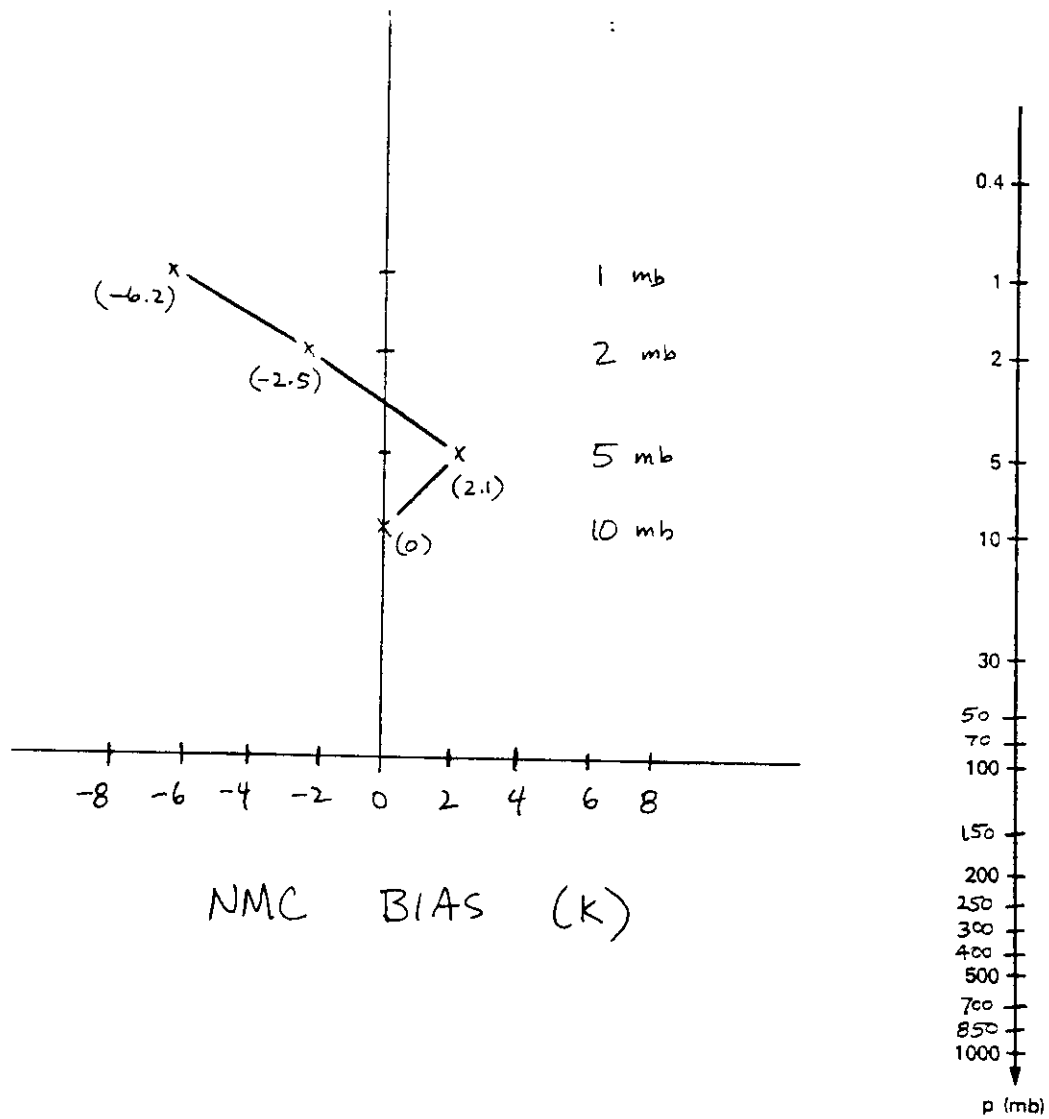
NMC
43.90 5.70
12/14/91
12.0.0. 0

(RADIOSONDE)

NIMES/COURBE
43.87 4.39
12/14/91
11.0.0. 0



NMC stratospheric temperature bias



Example of change in tropospheric analyses
resulting from change in forecast model

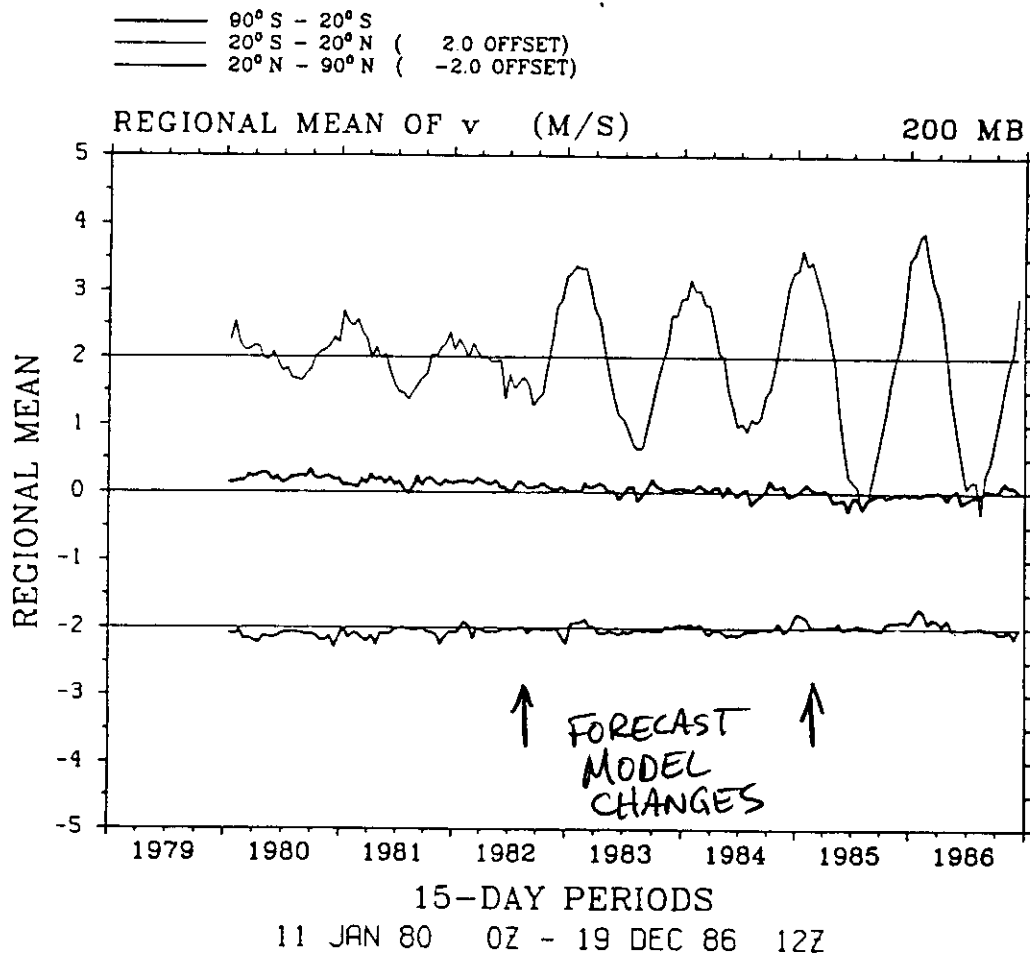
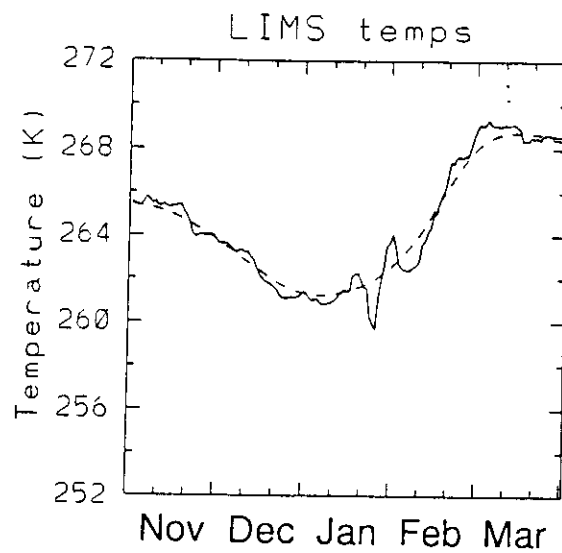


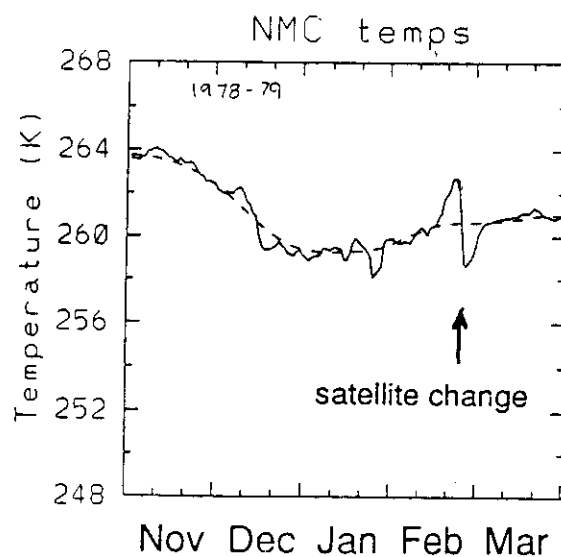
Fig. 14. 15-day averages of areal averages of v at 200 mb, in $m s^{-1}$.

— note change in wind velocity coincident with
changes in forecast
model,
from Trenberth (1992)

Discontinuities introduced by satellite changes



Limb
Infrared
Monitor of the
Stratosphere
data



NMC
data

CHANGES DETERMINED BY STATISTICAL EVALUATION

Another method of determining these changes is by statistical evaluation. Determinations based on this approach for 2 mb are also shown in Figure 1 (denoted by Step Regression). An equation of the form

$$T = X + Wt + b_1d_1 + \dots + b_8d_8$$

was fit to the monthly mean temperatures (T) for a particular latitude and level. The period of record is defined for each of the eight periods. An offset or step function (d) is then determined for each of the 8 periods in the whole data set. The X is a relative bias for the whole data set (reference is arbitrary) and Wt is the slope (W) with time (t) for the 91 months of record, from October 1978 to April 1986.

The relative pattern of agreement among the different methods depicted in Figure 1 is reasonably good. However, differences of 2 to 4 C degrees are common so that uncertainty in adjustments of this magnitude should be expected.

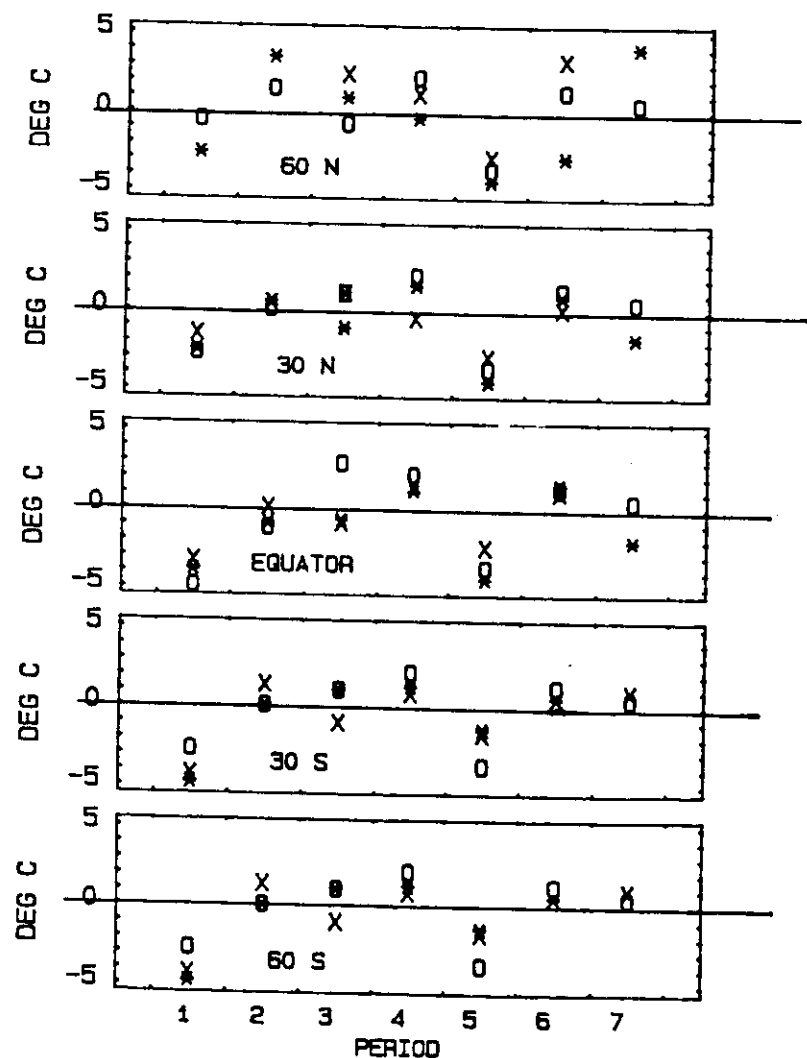
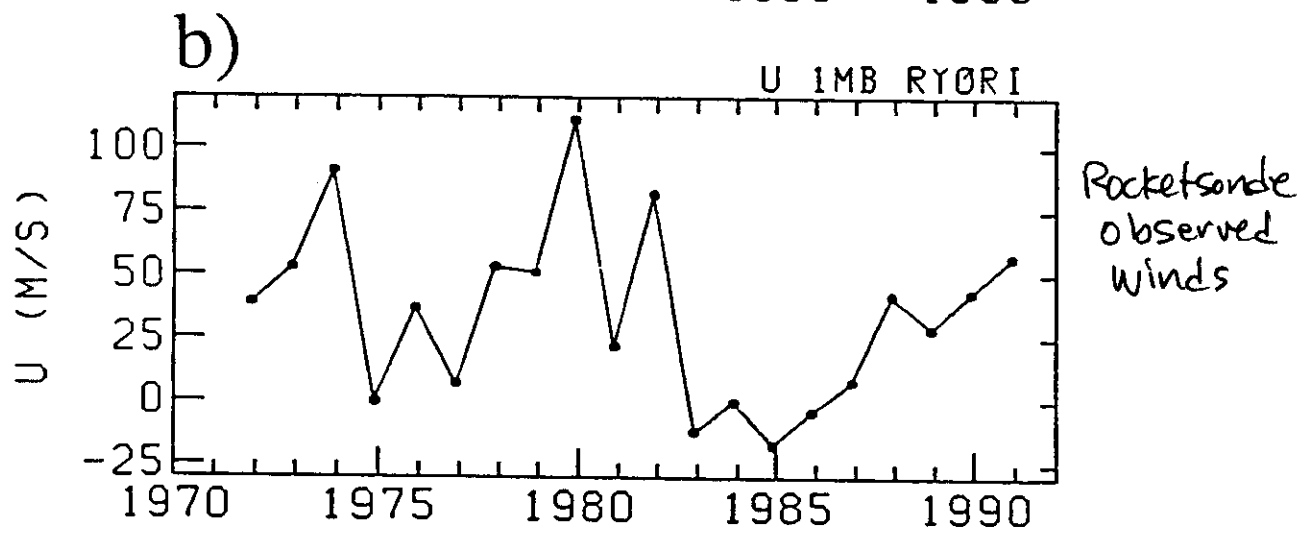
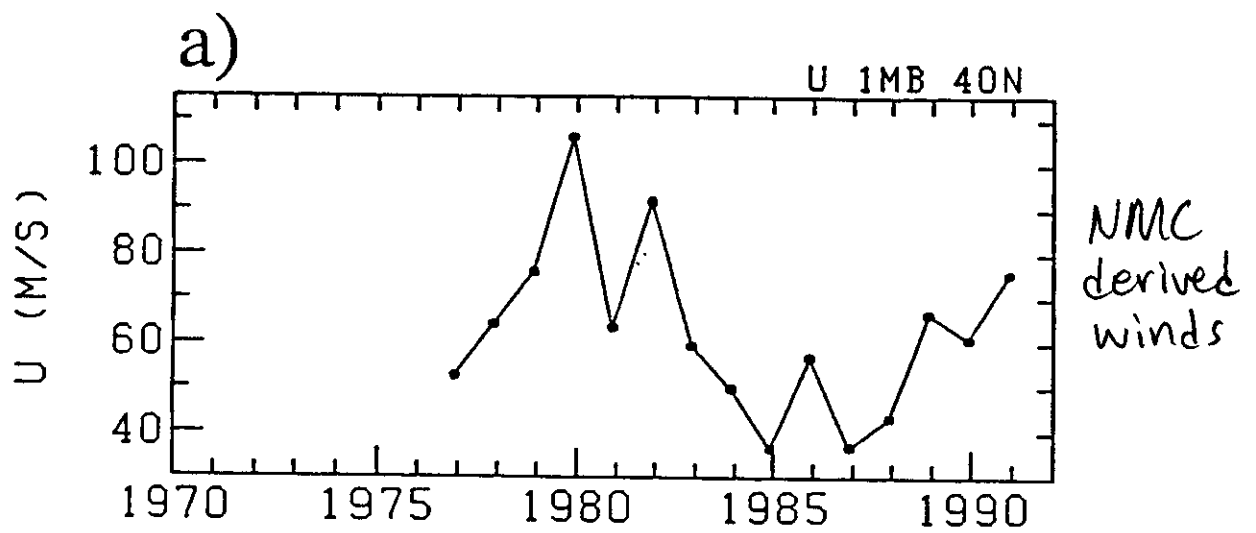


Fig. 1. Temperature changes in the 2 mb NMC analyses, as inferred from rocket comparisons (o,Rocket), comparison of analyses around change dates (x,Joint), and regression (*,Step Regression), as explained in text.



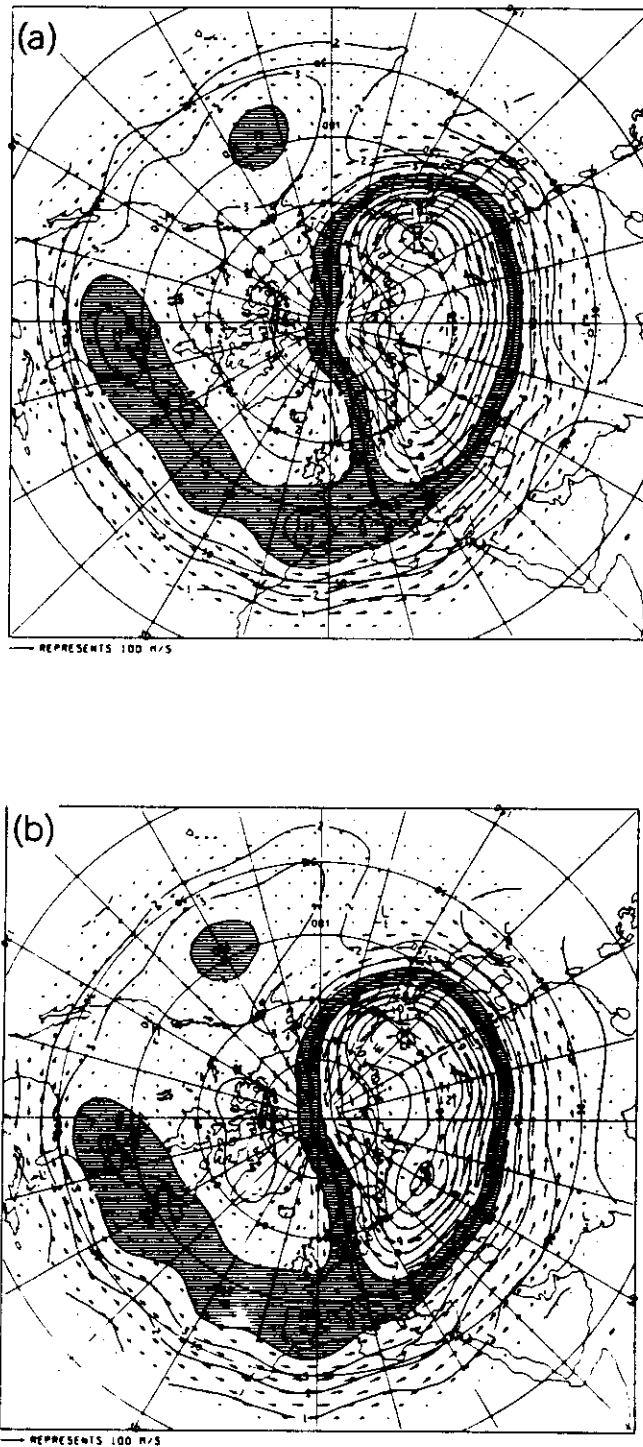


Figure 1. Ertel's potential vorticity, Q , and geostrophic winds evaluated on the 850 K isentropic surface for 7 December 1981. (a) Using data from NOAA-6; (b) using data from NOAA-7. Units: $\text{K m}^2 \text{ kg}^{-1} \text{ s}^{-1} \times 10^{-4}$. In these units, areas with values of Q between 4 and 6 are shaded.

NMC data available at Trieste

Global analyses of geopotential height
covering 1000-1 mb (17 pressure levels)

Once daily analyses covering 1979-1990

Stored as zonal Fourier coefficients
(up to wavenumber 8) on a regularly spaced
latitude grid (spacing near 4.4 degrees)

$$\begin{array}{c} \text{longitude} \searrow \\ \text{latitude} \downarrow \\ \text{Pressure} \swarrow \end{array} Z(\lambda, \phi, p) = \sum_{k=1}^8 ZS(\phi, p) \cdot \sin(k\lambda) + ZC(\phi, p) \cdot \cos(k\lambda)$$

↑
zonal wavenumber

TEMPERATURE, WIND AND WAVE FLUX DERIVATIONS

Temperatures are derived from the geopotential grids via the hydrostatic relation:

$$T = -R \frac{\partial \Phi}{\partial \ln p} \quad (1)$$

Here Φ is the geopotential height, $\ln p$ is the logarithm of pressure, and R the gas constant ($287 \text{ m}^2 \text{ s}^{-2} \text{ K}^{-1}$). Although simply called temperature here, this expression actually results in the virtual temperature (the temperature of dry air having the same density and pressure as the true, moist air).

Zonal mean winds are evaluated from the so-called gradient wind expression:

$$\frac{\bar{u}^2}{a} \tan \phi + 2\Omega \sin \phi \cdot \bar{u} + \frac{1}{a} \frac{\partial \bar{\Phi}}{\partial \phi} = 0 \quad (2)$$

Values at the equator are evaluated from the curvature expression:

$$\bar{u} + \frac{1}{2\Omega a} \cdot \frac{\partial^2 \bar{\Phi}}{\partial \phi^2} = 0 \quad (3)$$

Fleming and Chandra (1989) have compared equatorial winds derived from these data via Eq. 3 with monthly mean rawinsonde and rocketsonde measurements. They show that the derived winds reproduce the general characteristics of the quasi-biennial oscillation (QBO) and semi-annual oscillation (SAO), although the amplitude of the QBO is substantially underestimated in the derived winds.

Zonal Fourier components of horizontal winds are derived from the zonal and meridional momentum equations linearized about the zonal mean zonal wind \bar{u} :

$$\frac{\bar{u}}{a \cos \phi} \frac{\partial u'}{\partial \lambda} + \hat{f} v' + \frac{1}{a \cos \phi} \frac{\partial \Phi'}{\partial \lambda} = 0 \quad (4a)$$

$$\frac{\bar{u}}{a \cos \phi} \frac{\partial v'}{\partial \lambda} + \tilde{f} u' + \frac{1}{a} \frac{\partial \Phi'}{\partial \phi} = 0 \quad (4b)$$

In Eqns. (2)-(4), λ is longitude, ϕ latitude, u and v the zonal and meridional velocities, respectively, and a is earth radius (6.37×10^6 m). Overbars denote zonal means and primes deviations therefrom, and

$$\hat{f} = \left[2\Omega \sin \phi - \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\bar{u} \cos \phi) \right] \quad (4c)$$

$$\tilde{f} = \left[2\Omega \sin \phi + \frac{2\bar{u}}{a} \tan \phi \right] \quad (4d)$$

where Ω is earth rotation rate ($7.3 \times 10^{-5} \text{ sec}^{-1}$). The coupled equations (4a-b) are solved for each zonal wave number to evaluate the spectral coefficients of u' and v' ; note that these equations become singular for zonal wave k as

$$\delta \equiv \frac{(\bar{u} \cdot k / a \cos \phi)^2}{\hat{f} \tilde{f}} \rightarrow 1 \quad (4e)$$

For the data presented here, geostrophic winds were substituted if $\delta > 0.5$ (the results are not sensitive to this exact value). The use of these balanced winds and their improved treatment of stratospheric data (as compared to geostrophic winds) is discussed extensively in Randel (1987).

Statistics of several derived quantities are also presented here. A measure of the zonal mean static stability is given by the Brunt-Vaisala (or bouyancy) frequency squared:

$$N^2 = \frac{R}{H} \left(\frac{\partial \bar{T}}{\partial z} + \frac{\kappa \bar{T}}{H} \right) \quad (5)$$

where H is the atmospheric scale height, chosen here to be 7000 m, $z = H \ln(1000 \text{ mb}/p)$ is the vertical coordinate, and $\kappa = 2/7$. Relatively large values of N^2 correspond to strong stability with respect to vertical motions. A central quantity in the quasi-geostrophic theory of Rossby wave propagation and stability of the zonal mean flow is the latitudinal gradient of zonal mean potential vorticity:

$$\begin{aligned} \bar{q}_y = & \frac{2\Omega}{a} \cos \phi - \frac{1}{a^2} \frac{\partial}{\partial \phi} \left[\frac{1}{\cos \phi} \frac{\partial}{\partial \phi} (\bar{u} \cos \phi) \right] \\ & - (2\Omega \sin \phi)^2 \cdot e^{z/H} \frac{\partial}{\partial z} \left(e^{-z/H} \cdot \frac{1}{N^2} \cdot \frac{\partial \bar{u}}{\partial z} \right) \end{aligned} \quad (6)$$

A zonal cross section of potential temperature (Θ) and Rossby-Ertel potential vorticity (PV) is included for each month. Potential temperature is calculated as:

$$\Theta = T \left(\frac{1000 \text{ mb}}{p} \right)^\kappa \quad (7)$$

and potential vorticity from:

$$PV = -g(2\Omega \sin \phi + \zeta) \cdot \frac{\partial \Theta}{\partial p} \quad (8a)$$

Here ζ is the relative vorticity, and $g = 9.81 \text{ m} \cdot \text{sec}^{-2}$. For the zonal mean, this reduces to

$$\overline{PV} = -g \left[2\Omega \sin \phi - \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\bar{u} \cos \phi) \right] \cdot \frac{\partial \bar{\Theta}}{\partial p} \quad (8b)$$

PV is scaled in units of $10^{-6} \text{ K m}^2 \text{ kg}^{-1} \text{ s}^{-1}$, termed a PV unit. Statistics of the zonally-averaged poleward eddy heat flux ($\overline{v'T'}$) and poleward eddy momentum flux ($\overline{u'v'}$) are calculated and displayed. Also calculated are estimates of the zonal mean Eliassen-Palm (EP) flux vectors and their divergence. The orientation of the EP flux vectors indicates the wave group velocity or direction of wave activity flux (where such quantities are well defined), and the relative importance of heat versus momentum fluxes. The EP flux

divergence (also called the wave driving) gives a concise measure of the net effect of wave-induced heat and momentum fluxes in terms of a force per unit mass on the zonal mean flow; it is also equivalent to the poleward flux of quasi-geostrophic potential vorticity. A review of the motivation and theory behind EP flux diagnostics can be found in Edmon *et al.* (1980). Here the EP flux vectors and divergences are estimated using the primitive equation expressions, neglecting terms involving the vertical velocity. The zonal mean wave driving (DF) is given by:

$$\begin{aligned} DF &= \frac{1}{\rho(z) \cdot a \cos \phi} (\nabla \cdot \vec{F}) \\ &= \frac{1}{\rho(z) \cdot a \cos \phi} \left[\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (F_\phi \cos \phi) + \frac{\partial F_z}{\partial z} \right] \end{aligned} \quad (9)$$

with

$$F_\phi = \rho(z) \cdot a \cos \phi \left[-\overline{u'v'} - \frac{R}{H} \cdot \frac{\partial \bar{u}}{\partial z} \cdot \frac{\overline{v'T'}}{N^2} \right] \quad (10a)$$

$$F_z = \rho(z) \cdot a \cos \phi \cdot \left[\hat{f} \cdot \frac{R}{H} \cdot \frac{\overline{v'T'}}{N^2} \right] \quad (10b)$$

Here $\rho(z)$ is the basic state density (proportional to $e^{-z/H}$). EP flux diagrams here follow the convention in Randel *et al.* (1987); notably, the vectors are multiplied by $e^{z/2H}$ so that they are visible throughout the stratosphere. Because the use of balanced winds in the tropics is problematic, wave fluxes and EP diagnostics are displayed only polewards of 20° in each hemisphere.

Vertical derivatives are approximated by finite differences between adjacent pressure levels, followed by linear interpolation in log pressure back to the original pressure levels. Because of the separate analysis schemes at 100 and 70 mb, and the high number of derivatives used, eddy heat fluxes occasionally show discontinuous behavior between the 100 and 70 mb levels. Further vertical differentiation of these fluxes can

result in obviously bad EP flux divergence estimates; this problem is alleviated here by smoothing the eddy heat flux values across 100 and 70 mb prior to the EP flux calculations. Meridional derivatives are approximated by centered finite differences on the Gaussian latitude grid, and longitudinal derivatives are calculated spectrally.

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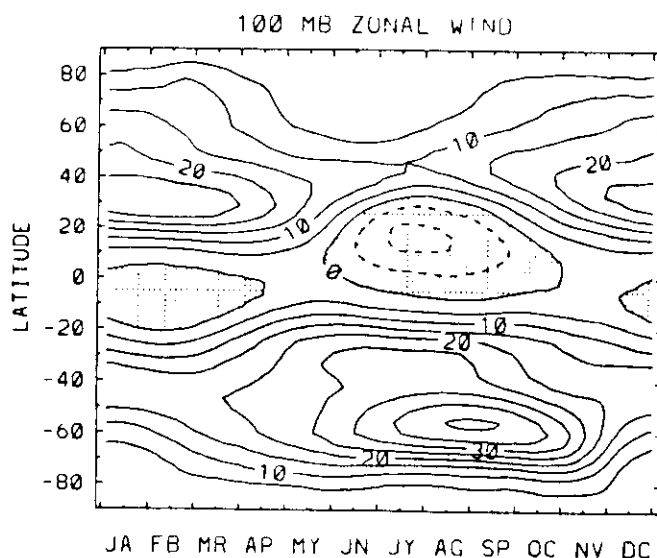
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P.O. Box 3000
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February 1992

Global Atmospheric Circulation Statistics, 1000-1 mb

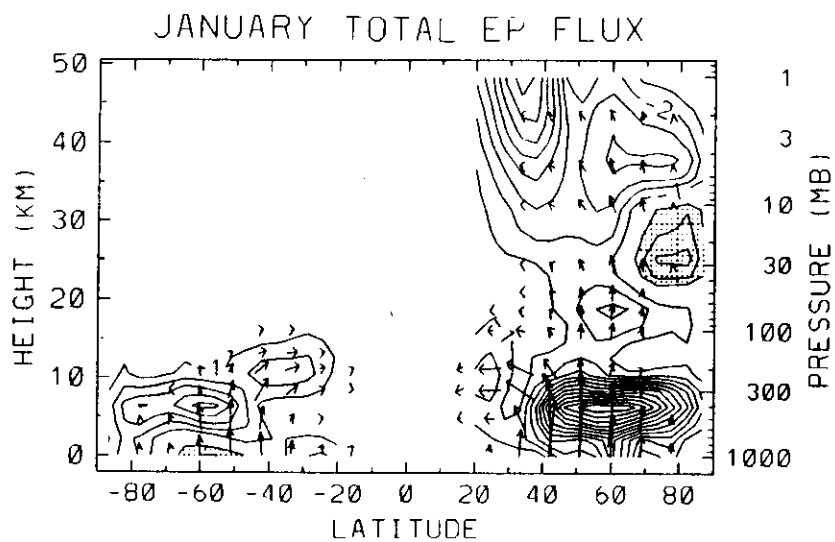
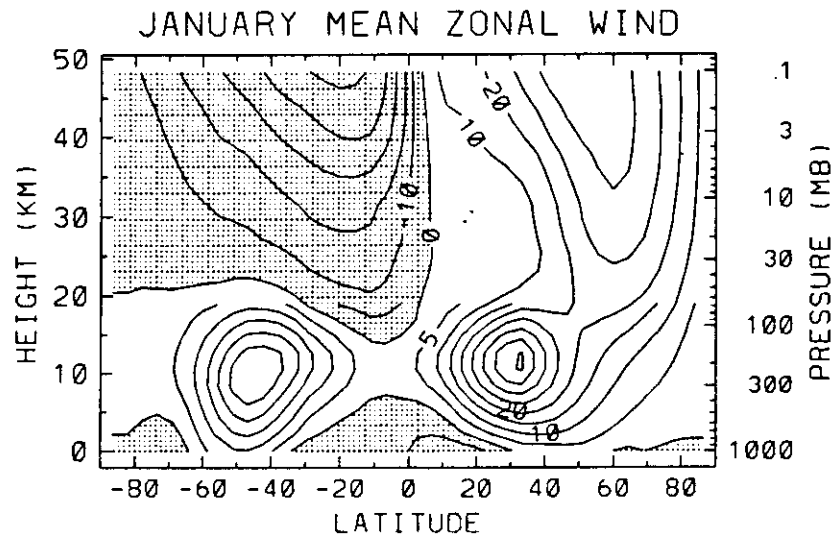
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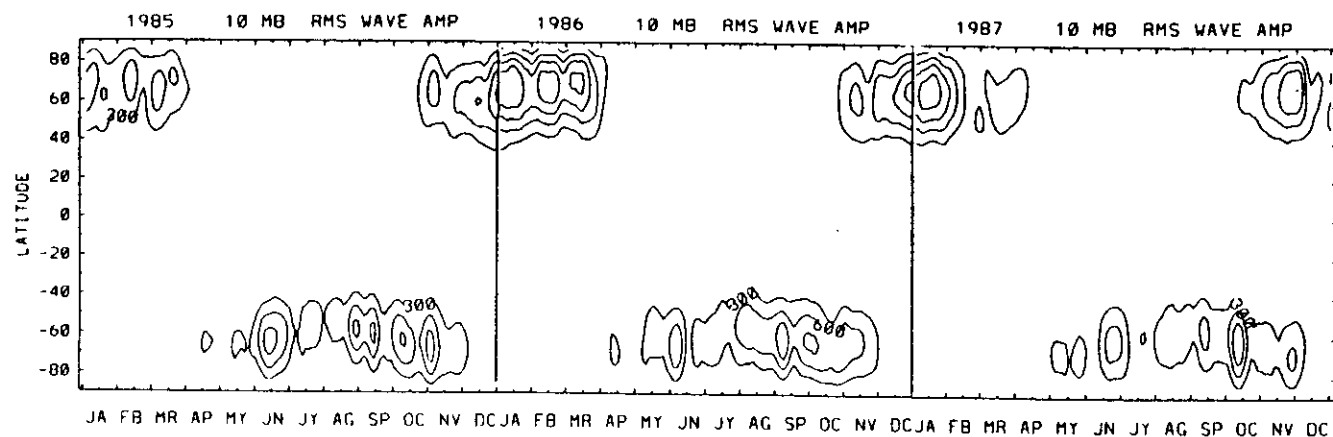
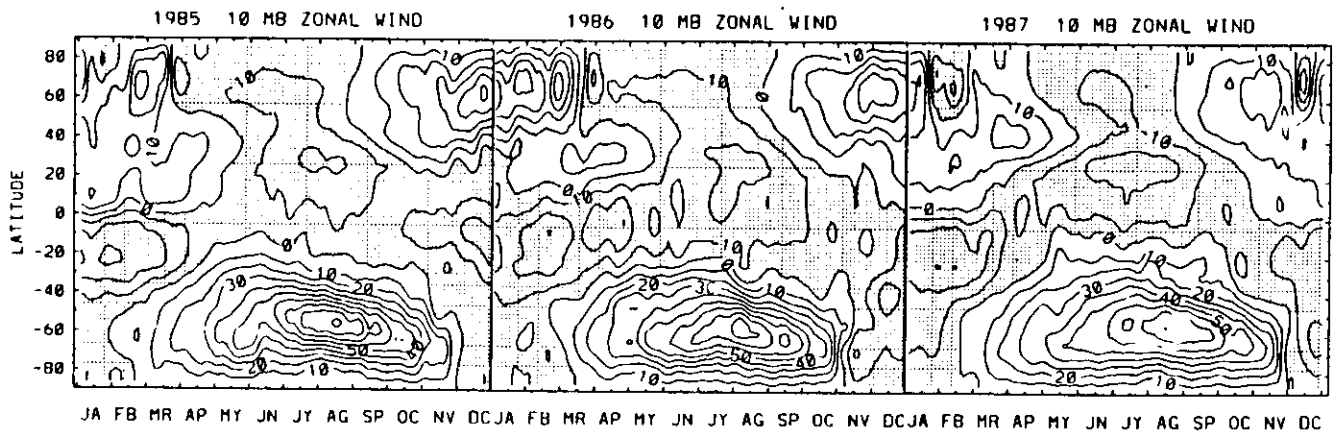
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BOULDER, COLORADO

Climatological means

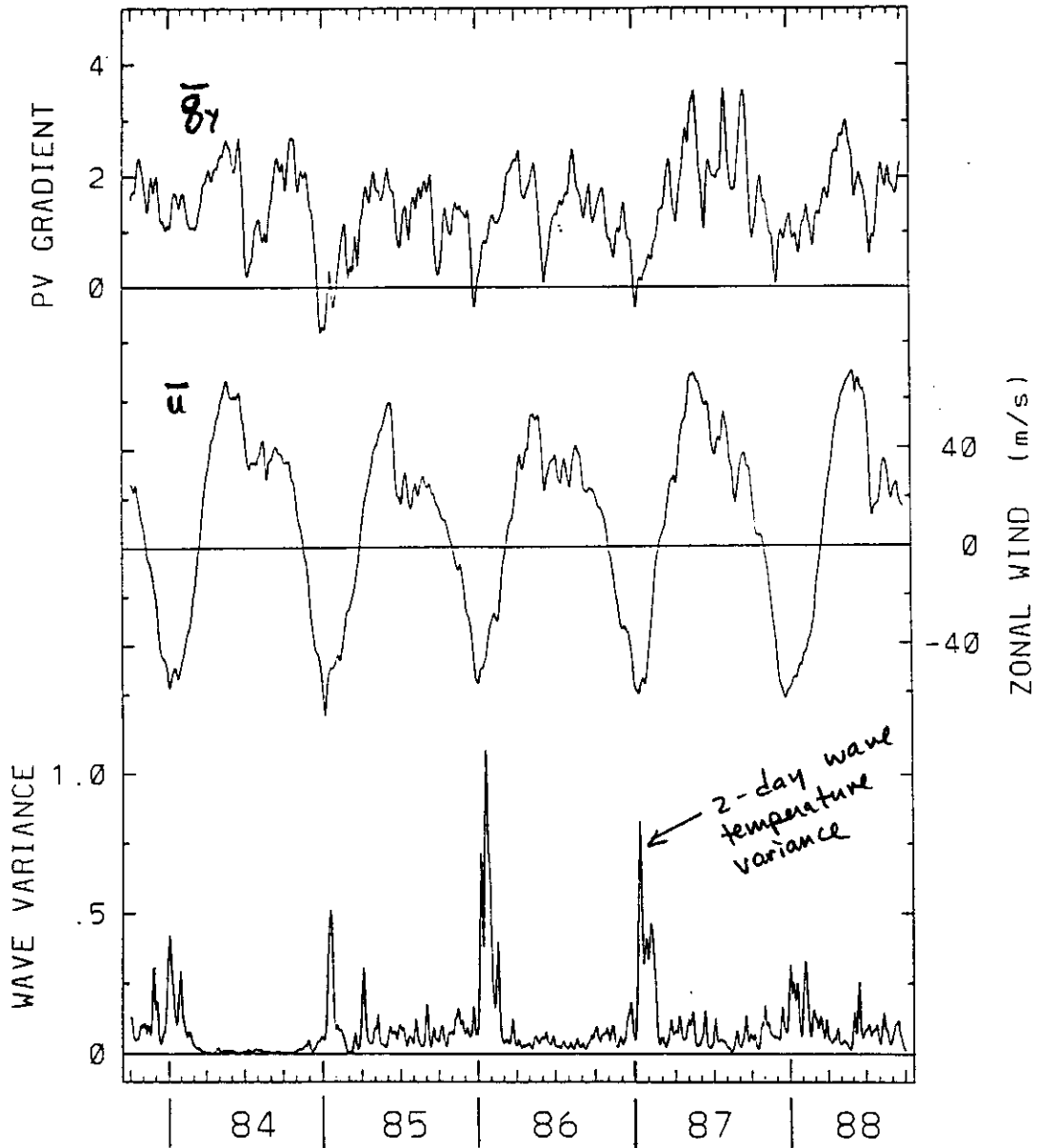


Reference: Randel, 1992: Global Atmospheric Circulation Statistics
NCAR Technical Note #366, 256 pp.

Stratospheric variability



Time series at 1 mb, 20° S



Reference: Randel (1992) in preparation: "The 2-day wave in NMC operational stratospheric analyses"

Fig. 2

Vertically propagating planetary waves

Reference: Randel, 1988, Tellus 40 257-271

