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## **Workshop on Fluid Mechanics**

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### **Similarity of steady stratified flows**

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These are preliminary lecture notes, intended only for distribution to participants

# (1)

## Similarity of Steady Stratified Flows

Incompressibility:

$$u_d \frac{\partial p}{\partial z_d} = 0$$

Continuity:

$$\frac{\partial(\rho u_d)}{\partial z_d} = 0, \quad \text{or} \quad \frac{\partial u_d}{\partial z_d} = 0.$$

Equations of motion:

$$\boxed{\rho u_d \frac{\partial u_i}{\partial z_d} = -\frac{\partial p}{\partial z_i} - g \rho \delta_{i3}. \quad (i=1,2,3)}$$

$$\delta_{i3} = \begin{cases} 0 & \text{if } i \neq 3 \\ 1 & \text{if } i = 3 \end{cases} \quad g \downarrow$$


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Let \* denote "laboratory", ^ denote "nature".

and let

$$y_i^* = \frac{x_i^*}{L^*}, \quad \hat{y}_i = \frac{\hat{x}_i}{\hat{L}}.$$

At corresponding points,

$$y_i^* = \hat{y}_i.$$

$$\hat{\rho} = \hat{\rho}_0 + \hat{\rho}_1 (\hat{y}_1 \cdot \hat{y}_2 \cdot \hat{y}_3).$$

$$\hat{u}_\alpha \frac{\partial \hat{\rho}}{\partial \hat{y}_\alpha} = 0. \quad (1)$$

$$\frac{\partial \hat{u}_\alpha}{\partial \hat{y}_\alpha} = 0. \quad (2)$$

$$\hat{\rho} \hat{u}_\alpha \frac{\partial \hat{u}_i}{\partial \hat{y}_\alpha} = - \frac{\partial \hat{\pi}}{\partial \hat{y}_i} - \hat{L} g \hat{\rho} \delta_{i3}, \quad (3)$$

where

$$\hat{\pi} = p + \hat{\rho}_0 g \hat{x}_3.$$

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In the laboratory, arrange to have

$$\rho^* = \rho_0^* + \rho_1^* (y_1^* \cdot y_2^* \cdot y_3^*).$$

$$\hat{\rho}_1 / \rho_1^* = r.$$

$$m = \hat{L} / L^*$$

$$u_i^* = (\hat{\rho} / r m \rho^*)^{1/2} \hat{u}_i. \quad (4)$$

and let

$$\pi^* = p^* + \rho_0^* g \hat{x}_3^*.$$

then

$$\textcircled{4} \text{ and } \textcircled{1} \rightarrow u_\alpha^* \frac{\partial \rho^*}{\partial y_\alpha^*} = 0 \quad (5)$$

$$\boxed{\frac{\partial u_{\alpha}^*}{\partial y_{\alpha}^*} = 0.}$$

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④ gives

$$\rho^* u_{\alpha}^* \frac{\partial u_i^*}{\partial y_{\alpha}^*} = \frac{\hat{p}}{r_m} \hat{u}_{\alpha} \frac{\partial \hat{u}_i}{\partial \hat{y}_{\alpha}}$$

$$\boxed{\frac{\partial}{\partial y_{\alpha}^*} \parallel \frac{\partial}{\partial \hat{y}_{\alpha}}}$$

so that (3) becomes

$$r_m \rho^* u_{\alpha}^* \frac{\partial u_i^*}{\partial y_{\alpha}^*} = - \partial \hat{\pi} / \partial y_i^* - r \hat{L} g \rho^* \delta_{i3}$$

$$\boxed{\hat{p}^* u_{\alpha}^* \frac{\partial u_i^*}{\partial y_{\alpha}^*} = - \partial \pi^* / \partial y_i^* - L^* g \rho^* \delta_{i3}.} \quad (7)$$

if we take

$$\pi^* = \hat{\pi} / r_m$$

and identify  $\pi^*$  with (i.e.  $\pi^* \equiv$ )

$$\hat{p}^* + \rho_0^* g \hat{x}_3^* .$$

For compressible fluids,

$$\rho/p^{1/\gamma} = \text{const. } e^{-S/c_p}$$

(  $S \equiv \text{entropy}, \gamma = c_p/c_v.$  )

Let

$$\lambda = (\rho/\rho_c)/(\rho/\rho_c)^{1/\gamma}.$$

Then

$$u_\alpha \frac{\partial \lambda}{\partial x_\alpha} = 0$$

$$\hat{\lambda} = \hat{\lambda}_0 + \hat{\lambda}_1, \quad \lambda^* = \lambda_0^* + \lambda_1^*,$$

$$r = \hat{\lambda}_1/\lambda_1^*.$$

$$u_i^* = (\hat{\lambda}/r m \lambda^*)^{1/2} \hat{u}_i, \quad \pi^* = \hat{\pi}/r m$$

where

$$\pi^* = \int \frac{dp^*}{\rho^*} + \lambda_0^* g z_3^*, \quad \hat{\pi} = \int \frac{d\hat{p}}{\hat{\rho}} + \hat{\lambda}_0 g \hat{z}_3.$$

$$\rho^* = \frac{\hat{\rho}}{\lambda^*}.$$

$$\hat{\rho}' = \frac{\hat{\rho}}{\hat{\lambda}}.$$

$$\frac{\rho^{*'} \cdot \hat{\rho}'}{(\rho^*)^{1/\gamma}} = \frac{\hat{\rho}'}{(\hat{\rho})^{1/\gamma}} = \text{constant}$$

similar  
inclusions  
as for  
comp. fluids

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Significance of

$$\hat{u}_i^* = (\hat{\rho}/r m \rho^*)^{1/2} \hat{u}_i .$$

$$r = \hat{\rho}_i / \rho_i^* ,$$

$$m = \hat{L} / L^* .$$

At corresponding points, then,

$$\frac{\rho^* u_i^{*2}}{\rho_i^* L^*} = \frac{\hat{\rho} \hat{u}_i^2}{\hat{\rho}_i \hat{L}} .$$

If

$$q^{*2} = u_i^{*2} u_i^* . \quad \hat{q}^2 = \hat{u}_i^* \hat{u}_i .$$

then, since (flow patterns being the same)

$$r = \frac{\hat{\rho}_i}{\rho_i^*} = \frac{\text{grad } \hat{\rho}_i}{\text{grad } \rho_i^*} \frac{\hat{L}}{L^*} . \quad \text{grad } \hat{\rho}_i = |\text{grad } \hat{\rho}_i|$$

we have, at corresponding points,

$$\frac{\rho^* q^{*2}}{g(\text{grad } \rho_i^*) L^*} = \frac{\hat{\rho} \hat{q}^2}{g(\text{grad } \hat{\rho}_i) \hat{L}^*} , \quad \text{or,}$$

With

$$g^* = g \frac{\text{grad } \hat{P}_i}{\rho_0} L^* \quad \hat{g} = g \frac{\text{grad } \hat{P}_i}{\hat{\rho}_i} \hat{L},$$

$$u_i^* = \left(\frac{\rho^*}{\rho_0}\right)^{1/2} \hat{u}_i. \quad \hat{u}_i = \left(\frac{\hat{\rho}_i}{\rho_0}\right)^{1/2} \hat{u}_i.$$

$$s^* = |u_i^*|, \quad \hat{s} = |\hat{u}_i|,$$

we have

$$\frac{s^{*2}}{g^* L^*} = \frac{\hat{s}^2}{\hat{g} \hat{L}} \quad (8)$$

$\sim$

$$F^{*2} = \hat{F}^2. \quad (\text{Equality of densimetric Froude numbers everywhere.})$$

Note that (4) implies

$$s^{*2} = \hat{s}^2 / rm$$

or

$$(\text{grad } s^*) L^* = (\text{grad } \hat{s}) \hat{L} \left(\frac{1}{rm}\right)^{1/2}$$

or

$$\frac{s^{*2}}{(\text{grad } s^*)^2 L^{*2}} = \frac{\hat{s}^2}{(\text{grad } \hat{s})^2 \hat{L}^2} \quad (9)$$

Eqs. (8) & (9) give

$$\frac{g^*}{(\text{grad } s^*)^2 L^{*2}} = \frac{\hat{g}}{(\text{grad } \hat{s})^2 \hat{L}^2}. \quad (10)$$

or

$$R_i^* = \hat{R}_i \quad (\text{Richardson numbers equal everywhere.})$$

