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LECTURE ON SUSY GUTS AND INFLATION

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Please note: These are preliminary notes intended for internal distribution only.

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Inflation & Particle Physics

Inflationary Cosmology contains a large number of parameters. They include

$$H_0, \Omega_b, \Omega_{\text{CDM}}, \Omega_{\text{HDM}}, n_s, n_g, \Lambda, \dots$$

Any 'Sensible' theory hopefully predicts at least some of them, consistent with the observations.

Why should particle physics be relevant?

If we think that GUTs/Susy GUTs are relevant, then have to worry about monopoles

Plus if we can do inflation with GUTs we can 'better trust' the calculation.

R-symmetric Superpotential

$$W = \kappa S \bar{\phi} \phi - \mu^2 S$$

$\uparrow \uparrow$
Gauge non-singlets

$\nwarrow$ gauge singlet

Under R:

$$\left. \begin{array}{l} S \rightarrow e^{i\gamma} S \\ \bar{\phi} \phi \rightarrow \text{invariant} \end{array} \right\} W \rightarrow e^{i\gamma} W$$

Note: S^2, S^3 excluded by R-sym.

From $W \rightarrow V$ (potential)

$$V(S, \phi, \bar{\phi}) = \kappa^2 |S|^2 [|\phi|^2 + |\bar{\phi}|^2] \\ + |\kappa \phi \bar{\phi} - \mu^2|^2 + \mathbf{D}\text{-terms}$$

Look for supersymmetric minimum

Get

$$|\langle \phi \rangle| = |\langle \bar{\phi} \rangle| = \mu / \sqrt{\kappa} \quad (\kappa > 0) (\equiv M_x)$$

$$\langle S \rangle = 0$$

Consider Early Universe With Chaotic Initial Conditions

In particular, for $S > S_c (\equiv M_x)$, it is

favorable to have $\phi = \bar{\phi} = 0$

$$(V = 2(\kappa^2 |S|^2 - \kappa \mu^2) |\phi|^2 + \mu^4 + \kappa^2 |\phi|^4)$$

Energy density $\sim \mu^4 \Rightarrow$ inflation

Q/ What 'drives' S to zero?

A/ Look at Radiative Corrections
(κ not small!)

Note: S non-vanishing $\Rightarrow$ SUSY broken
 $F_S = \mu^2$

One Loop:

$$\Delta V(S) = \sum_i \frac{(-)^F}{64\pi^2} M_i^4(S) \ln\left(\frac{M_i^2(S)}{\Lambda^2}\right)$$

So

$$V_{\text{eff}}(S) = \mu^4 + \frac{\kappa^2}{32\pi^2} \left[2\mu^4 \ln\left(\frac{\kappa^2 S^2}{\Lambda^2}\right) \right.$$

$$+ (\kappa S^2 - \mu^2)^2 \ln\left(1 - \frac{\mu^2}{\kappa S^2}\right)$$

$$\left. + (\kappa S^2 + \mu^2)^2 \ln\left(1 + \frac{\mu^2}{\kappa S^2}\right) \right]$$

For $S \gg S_c$, it reduces to

$$V_{\text{eff}}(S) \sim \mu^4 \left[1 + \frac{\kappa^2}{32\pi^2} \left(\ln \frac{\kappa^2 S^2}{\Lambda^2} + \frac{3}{2} \right) \right]$$

$S' > S_c$, μ^4 dominates $\rightarrow$ inflation

$S \approx S_c$, $\phi, \bar{\phi} \rightarrow$ driven to their non-zero values

$\Rightarrow$ gauge symmetry broken

Below S_c , the term $K^2 S^2 \phi^2$ is increasingly more effective in driving S to its supersymmetric value (of zero)

Remark :

Inflation ends before S reaches $S_c \Rightarrow \phi, \bar{\phi}$ cannot break the GUT symmetry but may have GUT vevs (e.g. $SO(10), (SU(3))^3$)

'Slow roll' conditions

$$\ddot{\phi} + 3H \dot{\phi} \approx -\frac{\partial V}{\partial \phi}$$

$$\epsilon \ll 1, \quad |\eta| \ll 1$$

$$\epsilon = \frac{M_p^2}{16\pi} \left(\frac{V'}{V} \right)^2, \quad \eta = \frac{M_p^2}{8\pi} \frac{V''}{V}$$

$$\eta = \left(K M_p / 4\pi M_x \right)^2 \frac{1}{8\pi} \left[(3x^2 - 1) \ln \left(1 - \frac{1}{x^2} \right) + (3x^2 + 1) \ln \left(1 + \frac{1}{x^2} \right) \right]$$

blows up as $x \rightarrow 1$ ($x \equiv S/S_c$)

Note: SUSY breaking effects are characterized by a TeV scale mass

$\Rightarrow$ small corrections to $V_{\text{eff}}(S)$

$\Rightarrow$ inflation does not depend on details of SUSY breaking

Let us look at $\delta\rho/\rho$ (or better still
 $(\Delta T/T)_Q \leftarrow \text{measured by COBE} \right) \sim 7 \times 10^{-6}$

$$(\Delta T/T)_Q \approx \sqrt{\frac{32\pi}{45}} \frac{V^{3/2}}{M_P^3 V'} = \underbrace{\left(8\pi N_Q\right)^{1/2} \left(M_x/M_P\right)^2}_{\substack{\uparrow \\ \text{Reminiscent of} \\ \text{cosmic strings!!}}}$$

$\therefore \Delta T/T$ is small $\because M_{\text{GUT}}$

is about 3 orders of magnitude
 smaller than M_P

With $N_Q \sim 50$ say, $M_x \sim 10^{15.5} \text{ GeV}$

Spectral Index n of density fluctuations is readily estimated

$$n \approx 0.98 \text{ (close to H-Z)}$$

What about K ?

$$\frac{K}{x_Q} \sim \frac{8\pi}{\sqrt{N_Q}} \frac{M_x}{M_P}$$

With $x_Q \sim 10$ (so $S \sim 10 S_c \Rightarrow S \ll M_P$
 $\Rightarrow$ planck scale corrections are small)

you get $K \sim 10^{-2}$

('tiny' coeffs.
absent)

Does this work in ^{realistic} $GUTS$??

Need $SO(10)$ or higher rank

Simplest example I know of

is based on $G \equiv SU(3)_c \times SU(3)_L \times SU(3)_R$
(‘trinification’)

This model, at least in the

‘higgs’ sector, is far simpler
than $SO(10)$.

