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INTERNATIONAL CENTRE FOR THEORETICAL PHYSICS

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Lecture III

SUMMER SCHOOL IN HIGH ENERGY PHYSICS AND COSMOLOGY

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CONFORMAL FIELD THEORY

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Please note: These are preliminary notes intended for internal distribution only.

LECTURE III

correction to notes LECTURE II:

$$G_1(x) = (1-x)^{1/2} x^{-1/2}$$

$$G_2(x) = (1-x)^{-1/2} x^{1/2}$$

Multi-spinon formulation of $Pu(z)$, $wz\omega$

OPE: $\phi^\alpha(z) \phi^\beta(w) = (-1)^q (z-w)^{-1/2} \varepsilon^{\alpha\beta} \left(1 + \frac{1}{2} (z-w)^2 T(w) + \dots \right)$
 $+ (-1)^q (z-w)^{1/2} (t_a)^{\alpha\beta} \left(j^a(w) + \frac{1}{2} (z-w) \partial j^a(w) + \dots \right)$

Modes:

$$\phi^\alpha(z) \chi_q(0) = \sum_{m \in \mathbb{Z}} z^{m + \frac{q}{2}} \phi_{-m - \frac{q}{2} - \frac{1}{4}}^\alpha \chi_q(0)$$

$$q = 0(1) \Leftrightarrow j^3 = \text{integer (half-integer)}$$

Generalized commutation relation:

$$\sum_{l \geq 0} c_l^{(-\frac{1}{2})} \left(\phi_{-n - \frac{q+1}{2} - l + \frac{3}{4}}^\alpha \phi_{-n - \frac{q}{2} + l + \frac{5}{4}}^\beta - \left(\alpha \leftrightarrow \beta \right) \right)$$

$$= (-1)^q \varepsilon^{\alpha\beta} \int_{m+n+q-1}$$

where $(1-x)^\alpha = \sum_{l \geq 0} c_l^{(\alpha)} x^l$

From this, properties of N -spinon states follow.

Example: $q=0, n=m=1$

$$(\phi_{-\frac{3}{4}}^+ \phi_{-\frac{1}{4}}^- - \phi_{-\frac{3}{4}}^- \phi_{-\frac{1}{4}}^+) |0\rangle = \overset{\text{singlet}}{\phi_{-\frac{3}{4}}^+ \phi_{-\frac{1}{4}}^-} |0\rangle = 0,$$

(agrees with result from 4-point function)

Systematic approach:

look for symmetry structure compatible
with multi-spinon structure of the
Hilbert space



Yangian $Y(\mathfrak{sl}_2)$

Origin: mathematical structure of
Maldacene-Shenker spin chains
with inverse-square exchange.

Definition of $\Upsilon(\mathcal{G})$ ($\mathcal{G}$ is finite dim. Lie algebra)

↓

Hopf algebra generated by Q_0^a, Q_1^a :

RELATIONS

$$(Y1) \quad [Q_0^a, Q_0^b] = f^{abc} Q_0^c$$

$$(Y2) \quad [Q_0^a, Q_1^b] = f^{abc} Q_1^c$$

$$(Y3) \quad [Q_1^a, [Q_1^b, Q_0^c]] + (\text{cyclic in } a, b, c) \\ = A^{abc, def} \{Q_0^d, Q_0^e, Q_0^f\}$$

$$(Y4) \quad [[Q_1^a, Q_1^b], [Q_0^c, Q_1^d]] + [Q_1^c, Q_1^d], [Q_0^a, Q_1^b]] \\ = (A^{abp, qrs} f^{cdp} + A^{cdp, qrs} f^{abp}) \{Q_0^q, Q_0^r, Q_1^s\}$$

with $f^{abc} f^{abc} = -2N \delta^{ab} \quad (\mathcal{G} = \mathfrak{su}(N))$

$$A^{abp, def} = \frac{\hbar^2}{4} f^{edp} f^{bez} f^{cfz} f^{pqr}$$

MULTIPLICATION

$$\Delta_{\pm}(Q_0^a) = Q_0^a \otimes 1 + 1 \otimes Q_0^a$$

$$\Delta_{\pm}(Q_1^a) = Q_1^a \otimes 1 + 1 \otimes Q_1^a \pm \frac{\hbar}{2} f^{abc} Q_0^b \otimes Q_0^c$$

REMARKS

1) this is non-commutative
non-co-commutative
Hopf algebra, i.e. a quantum group

2) Structure related to Yang's R-matrix

$$R = (u_1 - u_2) + \hbar P^{12}$$

hence the name 'Yangian'

3) for $\hbar \rightarrow 0$ Yangian goes into half
to loop algebra on $\mathfrak{g}$: $\hat{\mathfrak{g}}_+$

back to $\mathfrak{su}(2)$, WZW:

define
$$\begin{cases} Q_0^a = J_0^a \\ Q_1^a = \frac{1}{2} f^a_{bc} \sum_{n \geq 0} J_{-n}^b J_n^c \end{cases}$$

* these satisfy (Y1) ... (Y4),
hence generate Yangian $Y(\mathfrak{sl}_2)$

* commuting with the Q 's are:

$$H_1 = L_0, \quad H_2 = \sum_{n \geq 0} n J_{-n}^a J_n^a, \quad \text{etc.}$$

$\Rightarrow$ can decompose each L_0 level in
irreps of $Y(sl_2)$, each characterized
by values of $H_2, H_3, \dots$

precise content of these multiplets
fixed by representation theory of $Y(sl_2)$
(Chari - Pressley)

multiplets can be described in terms
of multi-spinon states

EXAMPLE 2-spinon states: $(n_2 \geq n_1 \geq 0)$

$$\Phi_{n_2, n_1}^{t, a} = (t^a)_{\alpha\gamma} \phi_{-n_2 - \frac{3}{2}}^{\alpha} \phi_{-n_1 - \frac{1}{2}}^{\gamma} |0\rangle$$

$$\Phi_{n_2, n_1}^s = \varepsilon_{\alpha\gamma} \phi_{-n_2 - \frac{3}{2}}^{\alpha} \phi_{-n_1 - \frac{1}{2}}^{\gamma} |0\rangle$$

At level $L_0 = 5$, $n_2 + n_1 = 4$, have the following
multiplets of $Y(sl_2)$

$$(4,0), \quad H_2 = 45 : \quad \Phi_{4,0}^{t,a} \quad \Phi_{4,0}^s$$

$$(3,1), \quad H_2 = 31 : \quad \Phi_{3,1}^{t,a} - \frac{1}{15} \Phi_{4,0}^{t,a} \quad \Phi_{3,1}^s + \frac{3}{15} \Phi_{4,0}^s$$

$$(2,2), \quad H_2 = 25 : \quad \Phi_{2,2}^{t,a} - \frac{1}{6} \Phi_{3,1}^{t,a} - \frac{11}{120} \Phi_{4,0}^{t,a}$$

GENERAL RESULT:

Yangian Highest Weight States are linear combinations of "fully polarized N -Spinon states"

$$\phi_{-\frac{2N-1}{4}-n_N}^+ \cdots \phi_{-\frac{3}{4}-n_2}^+ \phi_{-\frac{1}{4}-n_1}^+ |0\rangle$$

$$n_N \geq \dots \geq n_2 \geq n_1 \geq 0$$

They can be labelled by $\{n_N, \dots, n_2, n_1\}$

The L_0, H_2 eigenvalues are:

$$\begin{cases} L_0 = \frac{D^2}{4} + \sum_{i=1}^N n_i \\ H_2 = \sum_{i=1}^N 2(n_i + \frac{i-1}{2})(n_i + \frac{i}{2}) \end{cases}$$

The $SU(2)$ content for $\{n_N, \dots, n_2, n_1\}$ is

$$(\frac{1}{2}) \otimes (\frac{1}{2}) \otimes \dots \otimes (\frac{1}{2})$$

where $c_{i+1} \oplus c_i$ Symmetrized iff $n_{i+1} = n_i$.

The union of all these multiplets is the full Hilbert space of the theory.

Application: new expressions for the Virasoro characters in this theory

$$\chi_{g^2}^{\text{Vir}} = q^{-\frac{j}{2}} \sum_{m_1, m_2, \dots}^{\infty} q^{\frac{1}{2}(m_1^2 + m_2^2 + \dots - m_1 m_2 - m_2 m_3 - \dots)} \\ \times \frac{1}{(q)_{m_1}} \prod_{a \geq 2} \left[\frac{1}{2} (m_{a-1} + m_{a+1} + \int_a z_{j+1}) \right]_{m_a} q$$

where $\left\{ \begin{matrix} m_2, m_4, \dots, m_{2j} \\ m_1, m_3, \dots, m_{2j} \end{matrix} \right\}$ odd, rest even.

$$\left[\begin{matrix} a \\ b \end{matrix} \right]_q = \frac{(q)_a}{(q)_b (q)_{a-b}} \quad (q)_a = \prod_{n=1}^a (1 - q^n)$$

In the sum, m_1 is the number of spinors.

This formula first conjectured by Kedem et al.

About correlation functions: one may write down Ward Identity based on Yangian symmetry.

This takes over the role of the KZ equation!

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SPINON PICTURE FOR $SU(2)$, $k=1$, WZW

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YANGIAN $Y(\mathfrak{g})$

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