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PDE'S Cellular Automata and Parallel Computing

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These are preliminary lecture notes, intended only for distribution to participants.

# PDE'S, CELLULAR AUTOMATA and PARALLEL COMPUTING

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# **Navier-Stokes Equations**

$$\partial_{t} \boldsymbol{u} + (\boldsymbol{u} \cdot \nabla) \boldsymbol{u} = -\frac{1}{\rho} \nabla \rho + \nu \nabla^{2} \boldsymbol{u} + \frac{1}{\rho} \boldsymbol{F}$$

$$\nabla \cdot \boldsymbol{u} = 0$$

- Generally difficult to solve
- Many established, highly efficient numerical techniques are available
- Many of these are not suitable to distributed memory parallel machines
- Treatment of flows in porous media or multiphase flows is difficult

## Micro- and Macro-Dynamics

- N-S equations are a macroscopic description of a fluid
- The large scale behaviour of a fluid is a 'limit' of the dynamics on the molecular scale
- An arbitrary microdynamics can be simulated, as long as N-S equations are recovered in the macroscopic limit
- It is possible to choose the microdynamics of the system to ease parallel implementations and to support the description of highly heterogeneous media

#### **Lattice Boolean Gases**

(Frisch, d'Humières, ...)

 Regular lattice with every lattice site x connected to b neighbouring sites

$$X + C_i$$
  $i = 1,...,b$ 

 The state of every site is encoded in b boolean variables n<sub>i</sub>, evolving according to the rule

$$n_i(\mathbf{X}+\mathbf{C}_i,t+1)-n_i(\mathbf{X},t)=\Omega_i(n_i(\mathbf{X},t))$$

- The lattice and  $\Omega_i$  must be chosen to recover N-S
- $n_i$  encodes the presence of a particle with velocity  $c_i$

#### From Boolean to Real

(McNamara, Zanetti)

- Conventional computers are oriented toward floating point performance
- Boolean gases suffer from statistical noise, i. e. huge quantities of sites are needed
- In 3D,  $\Omega_i$  is not expressible in a closed form
- The transformation from boolean to floating point ('mesoscopic' limit) solves the first two problems

$$n_i \rightarrow N_i = \langle n_i \rangle$$

# The Last Steps

 Chapman-Enskog limit (Higuera, Jimenez):

$$\Omega_{i} = \sum_{j=1}^{b} A_{ij} (N_{i} - N_{i}^{eq})$$

with  $A_{ij}$  constructed from the boolean microdynamics

 Enhanced Collisions (Higuera, Succi, Benzi):

 $A_{ij}$  can be derived from the conservation laws and the physical parameters, i.e. LBE is a model of hydrodynamics, unrelated to microdynamics of any type

#### The LBE Scheme

$$N_{j}(\boldsymbol{x}+\boldsymbol{c_{j}},t+1) = N_{j}(\boldsymbol{x},t) + \sum_{j=1}^{b} A_{jj}(N_{j}(\boldsymbol{x},t) - N_{j}^{eq}(\boldsymbol{x},t))$$

with

$$N_{i}^{eq} = \frac{\rho_{eq}}{b} \left( 1 + \frac{D}{c^{2}} u_{\alpha} c_{i\alpha} + \frac{D^{2}}{2c^{4}} \frac{b - 2\rho_{eq}}{b - \rho_{eq}} Q_{i\alpha\beta} u_{\alpha} u_{\beta} \right)$$

where

$$C = |\mathbf{c}_{i}|$$

$$Q_{i\alpha\beta} = c_{i\alpha}c_{i\beta} - \frac{c^{2}}{D}\delta_{\alpha\beta}$$

# Hydrodynamic Behaviour

- FCHC 4D (b = 24) lattice projected down in 1, 2 or 3 D
- The collision matrix has 4 different eigenvalues: 0 (conservation laws), λ (viscosity), σ and τ (spurious ghost fields)
- Adiabatic limit:

$$\rho(\mathbf{X},t) = \sum_{j=1}^{b} N_{j}(\mathbf{X},t)$$

$$J(\mathbf{X},t) = \rho(\mathbf{X},t) \mathbf{u}(\mathbf{X},t) = \sum_{j=1}^{b} \mathbf{c}_{j} \mathcal{N}_{j}(\mathbf{X},t)$$

obey the N-S equations under some constraints on the collision matrix eigenvalues

#### **LBE Pros**

- Maintenance of a main, general purpose and optimized code
- Simple, quick and efficient SIMD and MIMD parallel implementations
- Easy boundary conditions
- Easy simulation of flows in grossly irregular geometries

#### LBE Cons

- Low Mach number
- Moderate Reynolds number
- Stability implies that on a lattice step Δu/u<20%</li>

## LBE Applications

- Turbulent Thermal Convection
- Flows in Porous Media
- Oil Reservoir Modelling
- Multiphase Flows

# **Recent Developments**

- Unevenly Spaced Grids
- Compressible Flows
- Quantum Mechanics
- Lattice BGK Models

# **Future Developments**

 Finite Volume LBE on Embedded Adaptive Grids