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College on Computational Physics

15 May - 9 June 1995

Computational Studies of High Energy Physics

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Computational Studies
of
High Energy Physics

R. S.

w/ Claudio Rebbi

at Boston University

using TM's CM-5

Outline

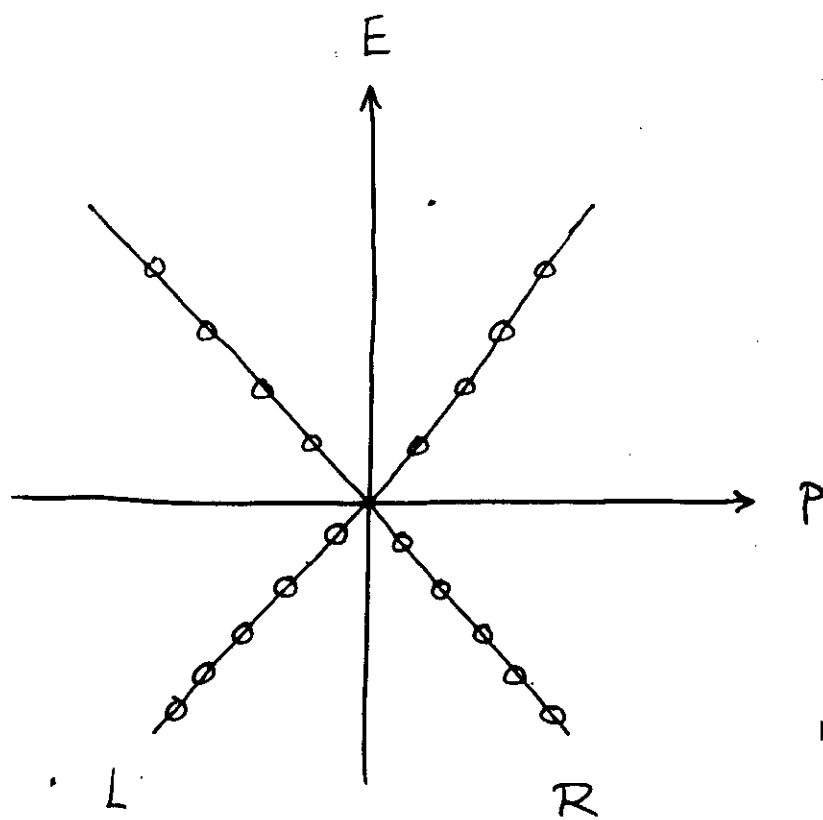
- Symmetry (anomalous)
- Vacuum Structure of gauge theory
- The sphaleron and topological transitions
- The computational Problem (anomalous baryon number violation)
- The Computational Strategy
- Computer animation of results
- More on the computation

Symmetry

	global	local
interactions	chiral symm.	weak
fields	electric charge	electromag.

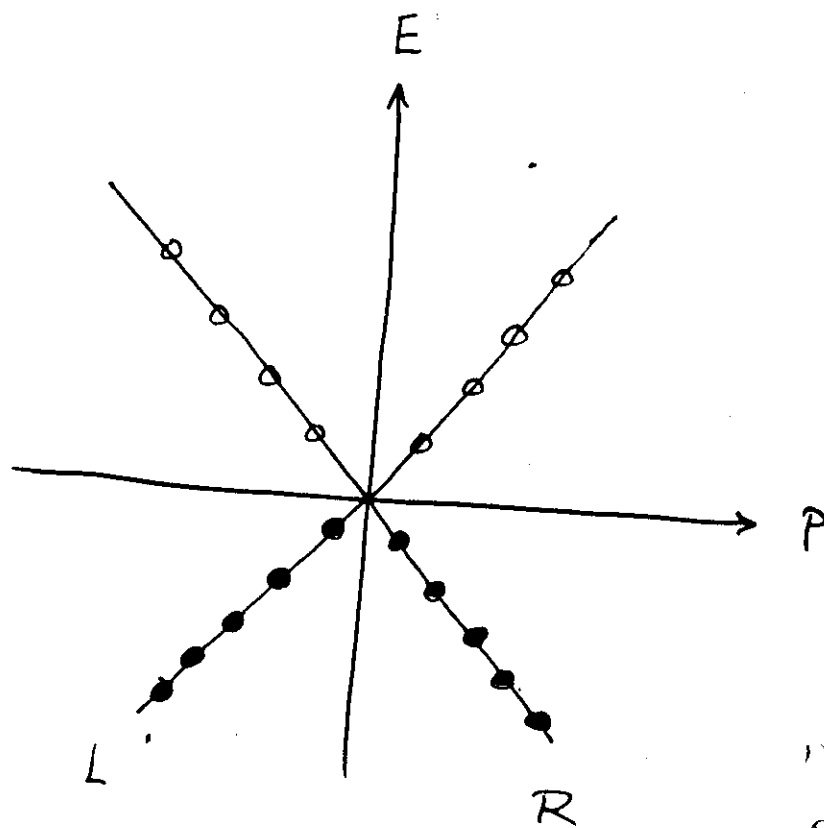
There is yet more: Anomalous Symm.
(purely quantum)

- a symmetry that is broken by quantum effects.
- Chiral gauge symmetries (weak interactions) suffer from this.



↑ Positive
energy
states

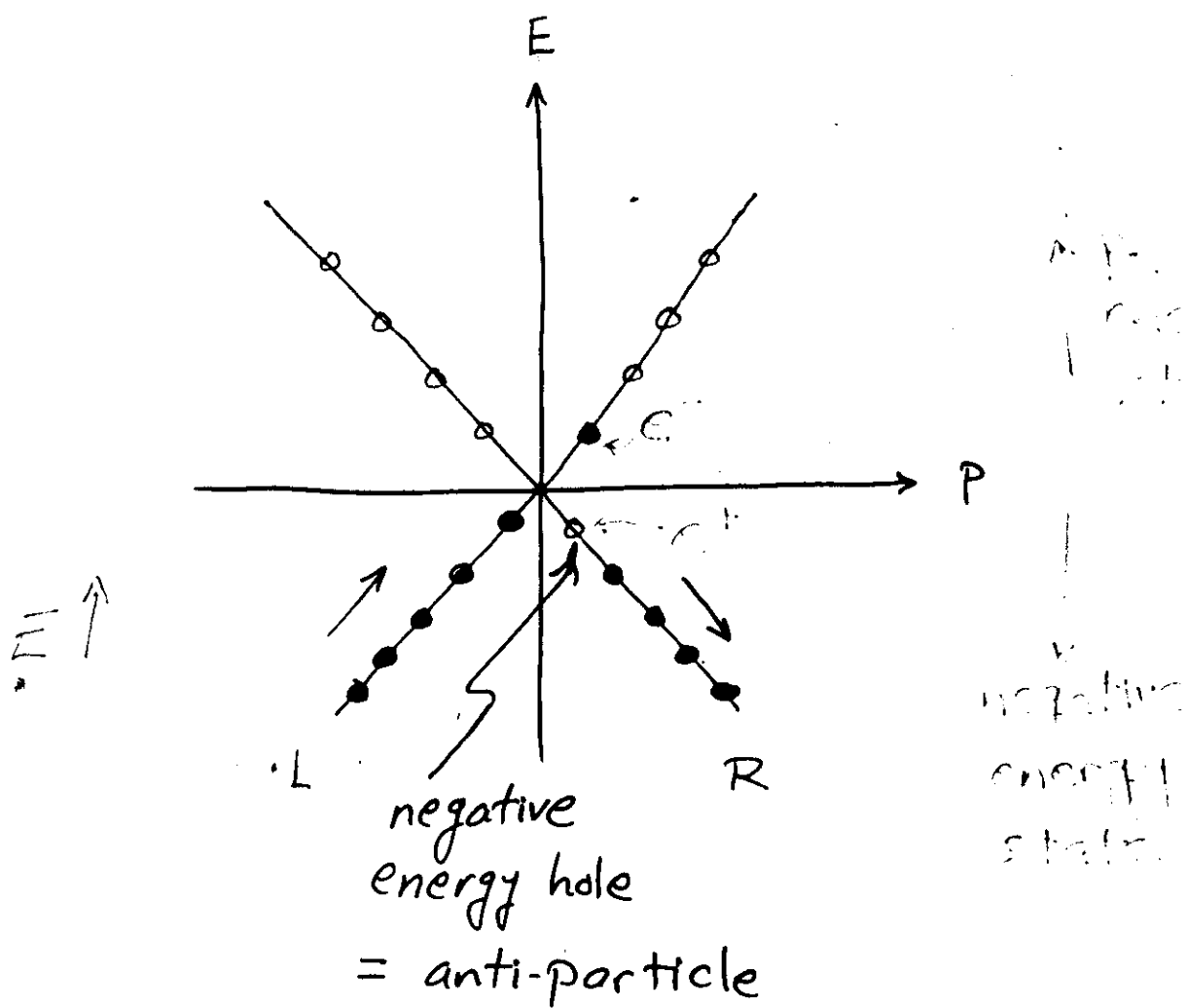
↓
negative
energy
states



↑ Positive
energy
states

↓
negative
energy
states

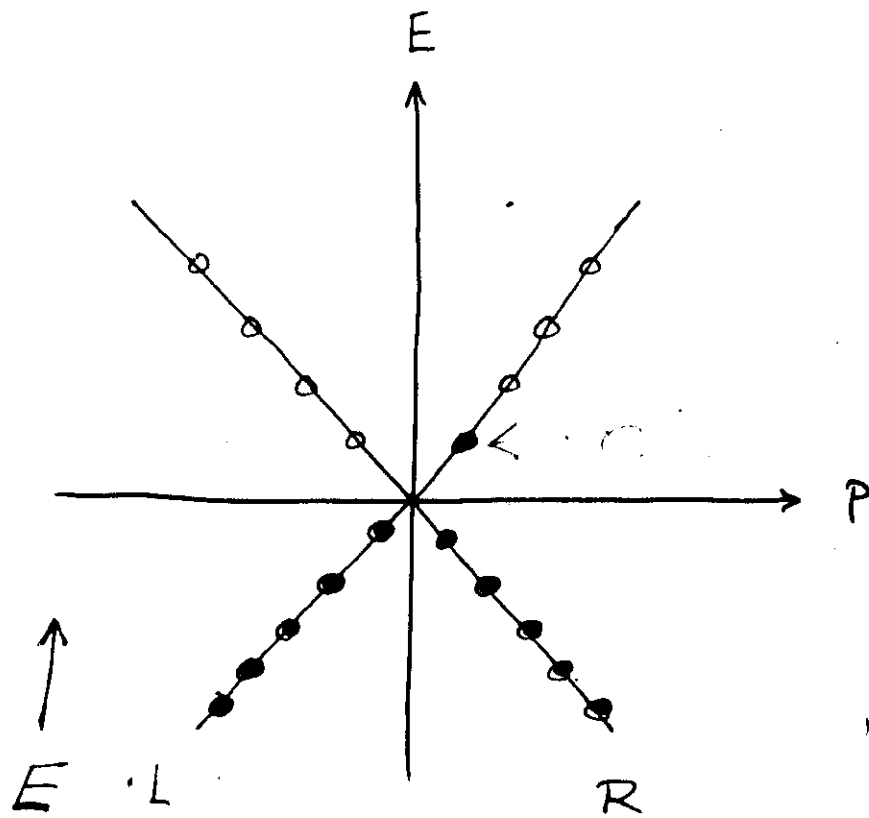
The vacuum
(all negative energy states
are filled)



$$V_{\text{vac}} = \langle \psi | \hat{H} | \psi \rangle$$

$$\Delta E = 0$$

Anomalous production



Any other
state

negative
energy
state

Chiral theory in which L handed particles are charged and R handed particles are neutral.

Thus, even though $Q = \text{electric charge}$ is classically conserved, it is not quantum mechanically.

Vacuum Structure

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu}^2 + |D_\mu \phi|^2 - \frac{\lambda}{2} (|\phi|^2 - v^2)^2$$

$$\underbrace{D_\mu \phi}_{\text{covariant derivative}} \equiv \partial_\mu \phi - i e A_\mu \phi$$

(one space + one time dim): $1+1$ ($\mu=0,1$)

This is called the Abelian Higgs model.

The vacuum has $|\phi| = v (\neq 0)$

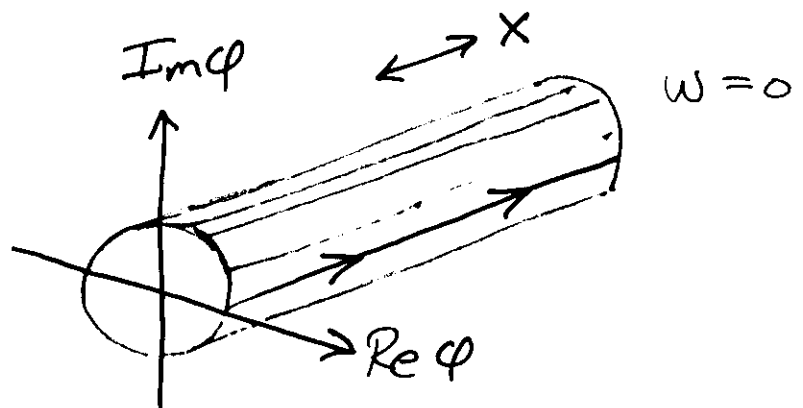
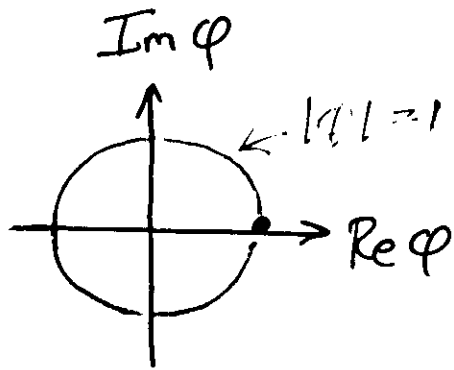
$\phi(x=+\infty) = \phi(x=-\infty)$: Periodic B.C.

$$\phi_{\text{vac}} : S^1 \rightarrow U(1) \sim S^1 \Rightarrow$$

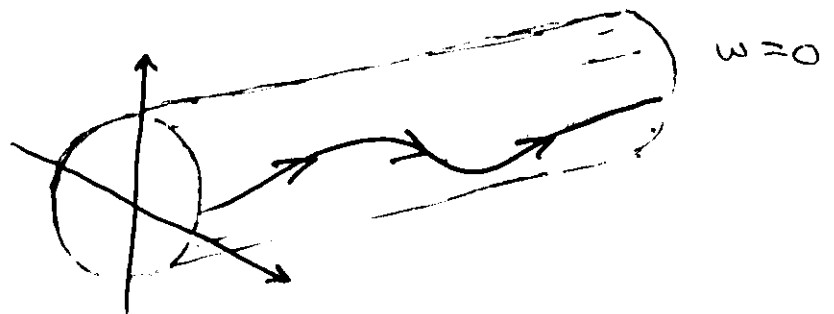
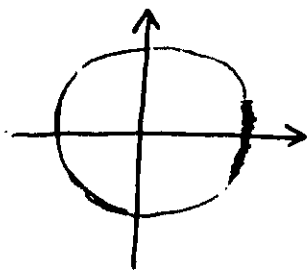
$$\underbrace{W[\text{vac}]}_{\text{winding \#}} \in \mathbb{Z}$$

winding #.

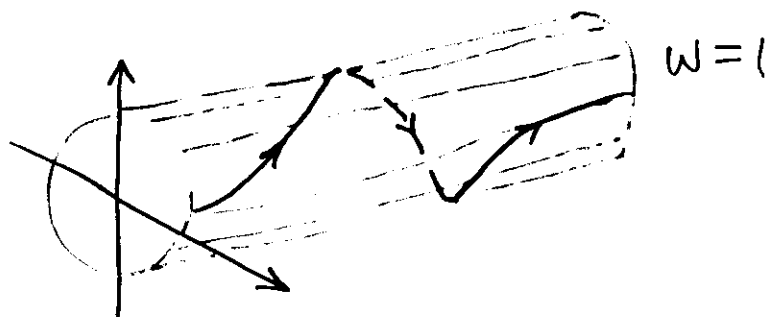
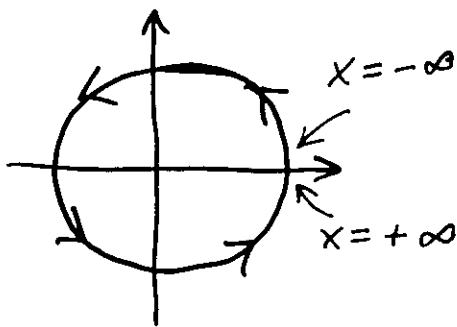
Vacuum Structure of 1+1 Abelian Higgs (Jackiw + Rebbi)



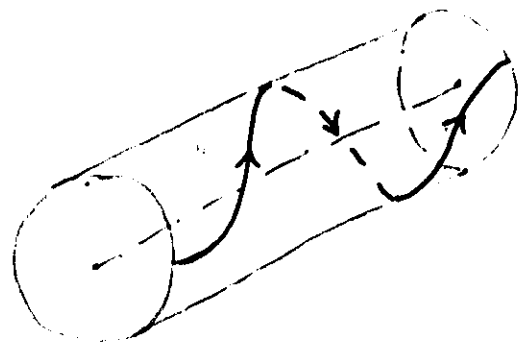
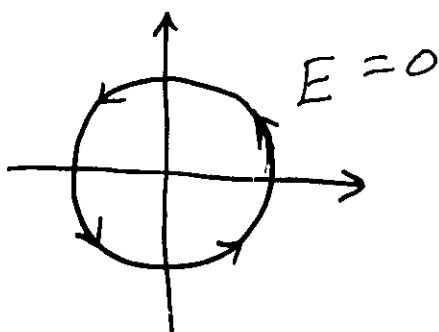
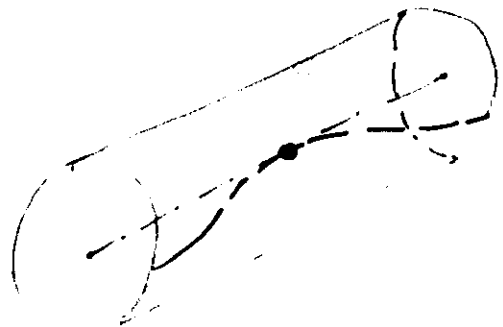
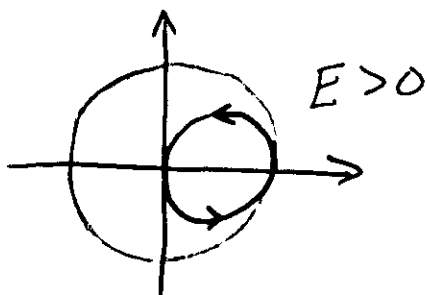
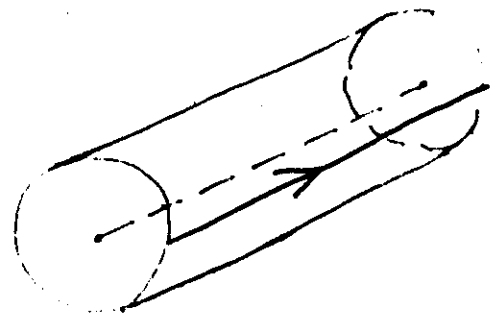
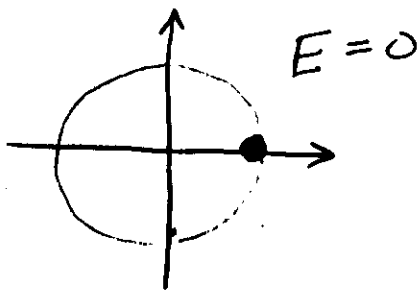
↓ topologically equivalent



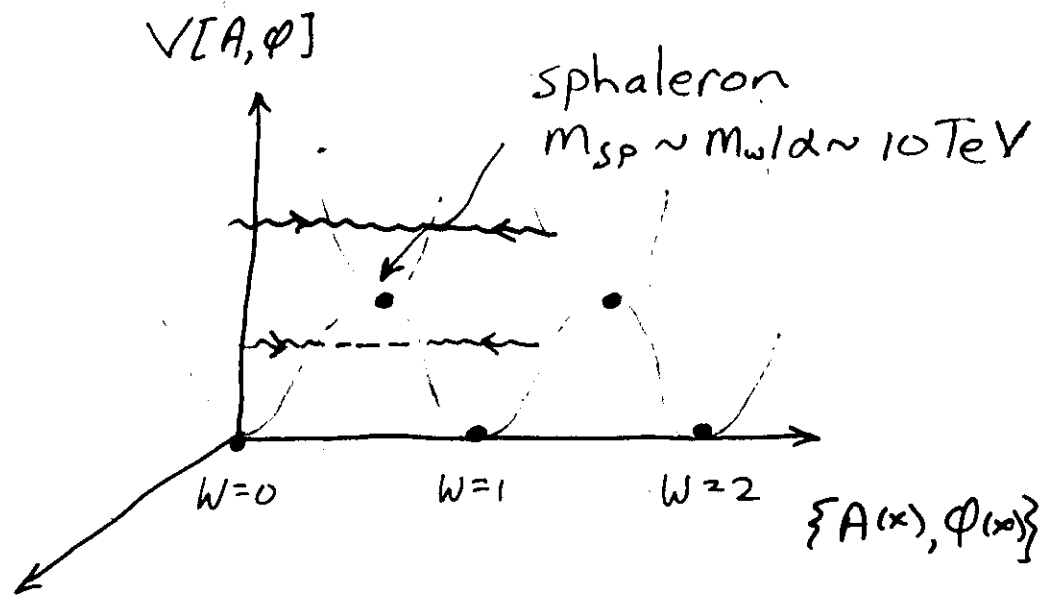
↓ topologically non-equiv.



In transitions between inequiv.
vacua, you must leave the manifold
of vacuum configurations (go over
the sphaleron)



The sphaleron

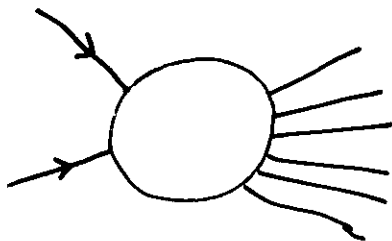


Field space: $A(x) = \int d^3P \tilde{A}(P) \underbrace{e^{iP \cdot x}}_{e_p(x) = \text{basis vectors}}$

Axial anomaly $\Rightarrow \Delta B = \Delta W$ (non perturbative!!)

low energy: $\sigma \sim e^{-4\pi/\alpha}$ (tunneling)

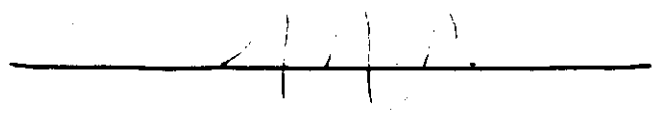
High Energy: $\sigma ?$ High T : $\sigma \sim \mathcal{O}(1)$

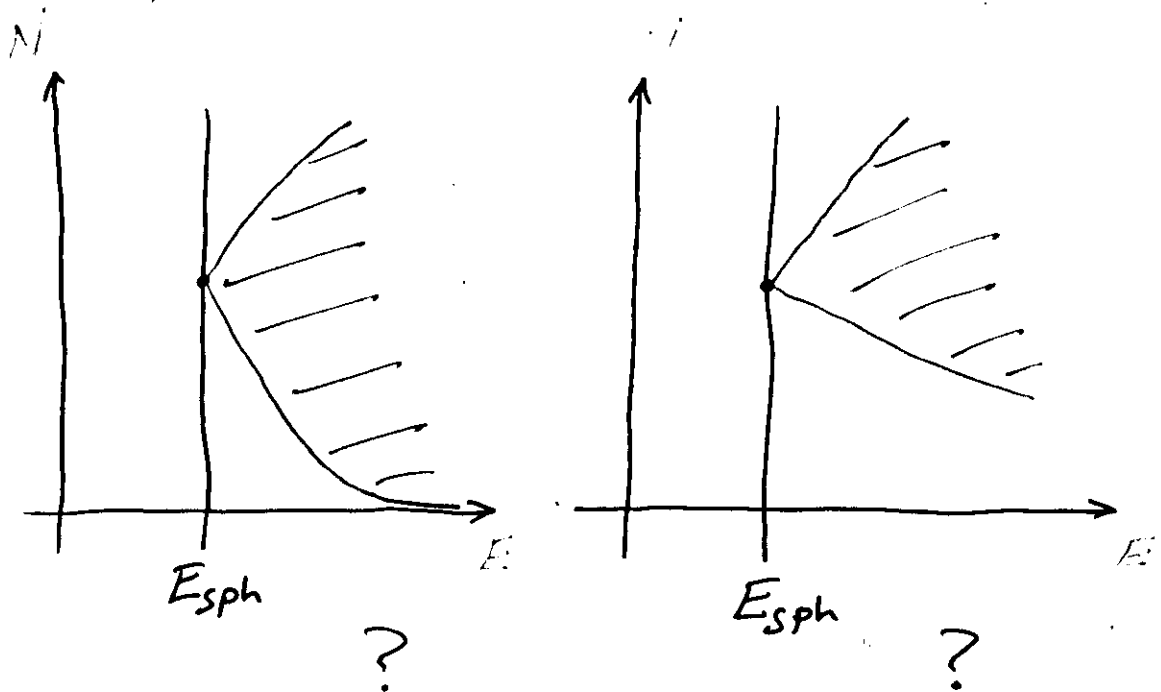


The Computational Problem

$\Rightarrow \quad \Leftarrow$
Linear: Identify N

passes over
the sphaleron.

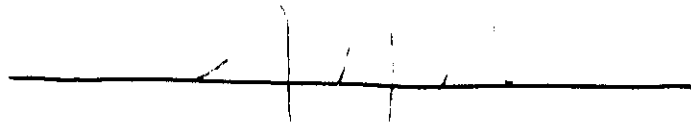
$\Rightarrow$ 
Topology Change ($\Delta W \neq 0$)



Find the boundary in the E - N plane for solutions with $\Delta W \neq 0$ (pass over the sphaleron).

Computational Strategy

Implement a time reversed evolution:



start w/ init. config on top of
the barrier (i.e. $\exists x \ni \phi(x) = 0$)
parameterized by $\{c_i\}$.



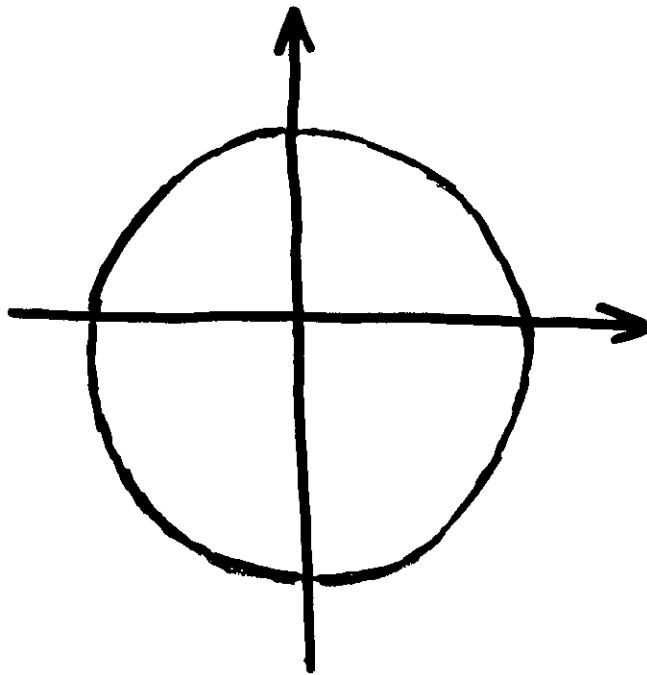
Evolve until linear regime and pull out N : $N = \sum_n |a_n|^2$.

where $\phi_{\text{sym}} = \sum_n a_n \phi_n$
 $\uparrow$
 normal modes

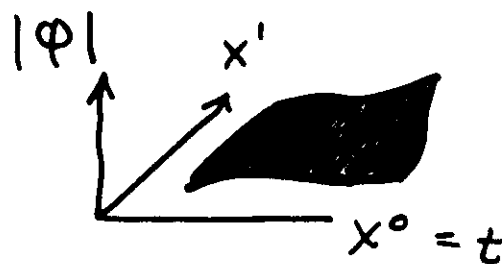
Computational Details

- gauge invariant (space) discretization of equations of motion.
- time evolution implemented with the leap-frog algorithm
 - 2nd order accurate
 - Exact time reversal
 - Conserves E exactly in the linear regime
- gauge invariant subtlety.
When pulling out the $\tilde{a}_n$, must gauge transform to "small" perturbations about trivial vacuum.

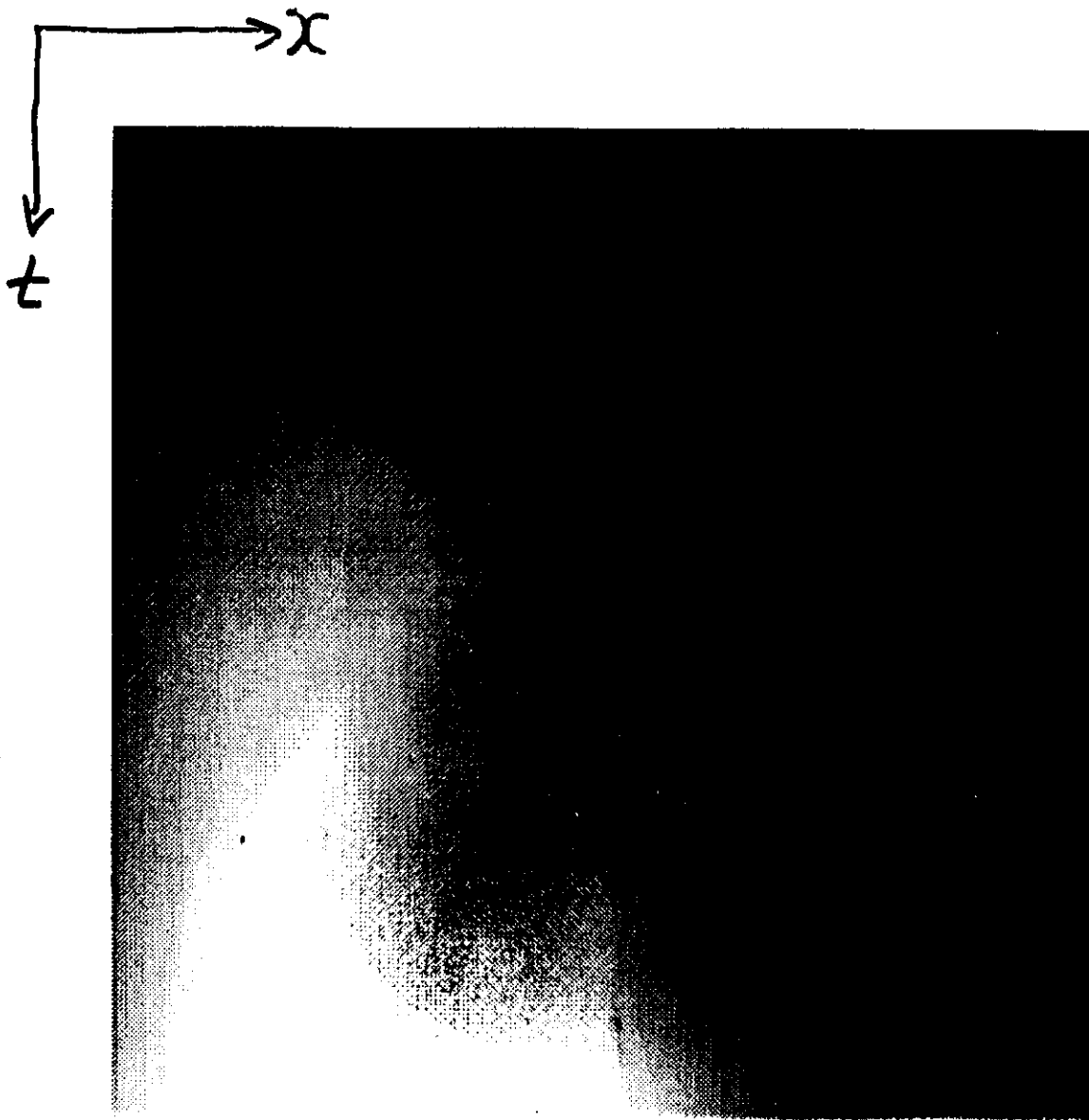
Phase is color coded :



$\varphi(x', x^0) \leftarrow$ 3-D colored surface

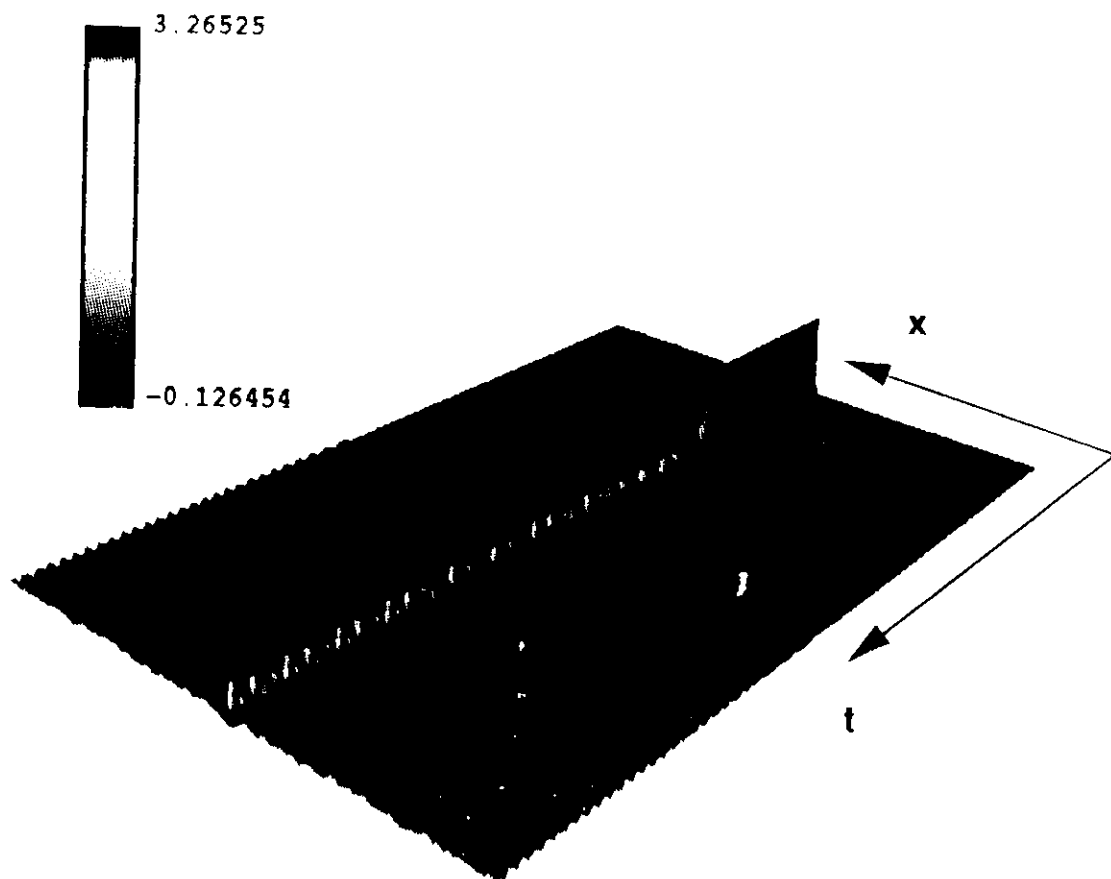


color code may not be unique.
Map will be chosen to give
pretty pictures.



Instanton transition in the 2-d Abelian Higgs model

Journal of Mathematical Physics, Vol. 34, No. 1, 1993



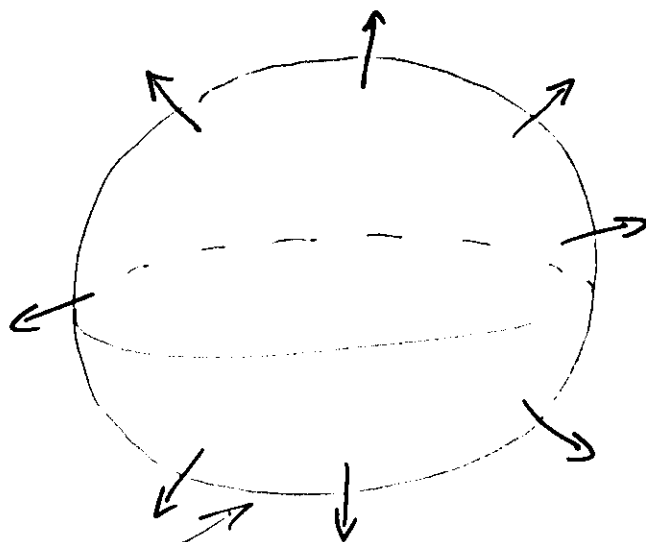
Sphaleron decay in 2D Abelian
Higgs model: the Higgs field

Figure 1: Sphaleron decay in 2D Abelian Higgs model

The Standard model in 3+1 dim

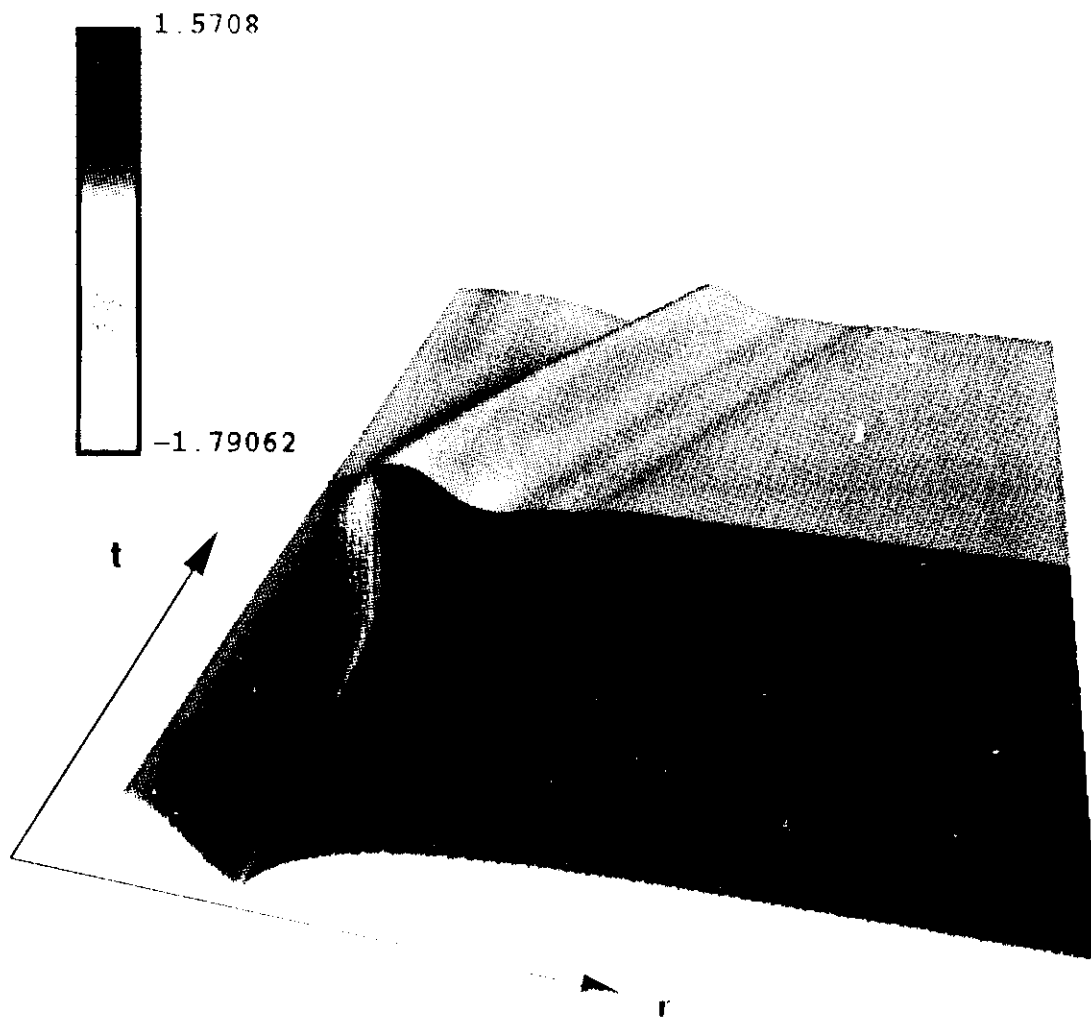
$$\begin{array}{ccc}
 A_\mu(\underline{x}, t) & \left. \vphantom{\begin{array}{c} A_\mu(\underline{x}, t) \\ \Phi(\underline{x}, t) \end{array}} \right\} \begin{array}{c} \text{radial} \\ \text{reduction} \end{array} \longrightarrow & a_\mu(r, t) \\
 \Phi(\underline{x}, t) & & \chi(r, t) \\
 \mu = 0, 1, 2, 3 & & \phi(r, t) \\
 & & \mu = 0, 1
 \end{array}$$

Radial symm. gives an effective theory in 1+1 dim (r, t) .

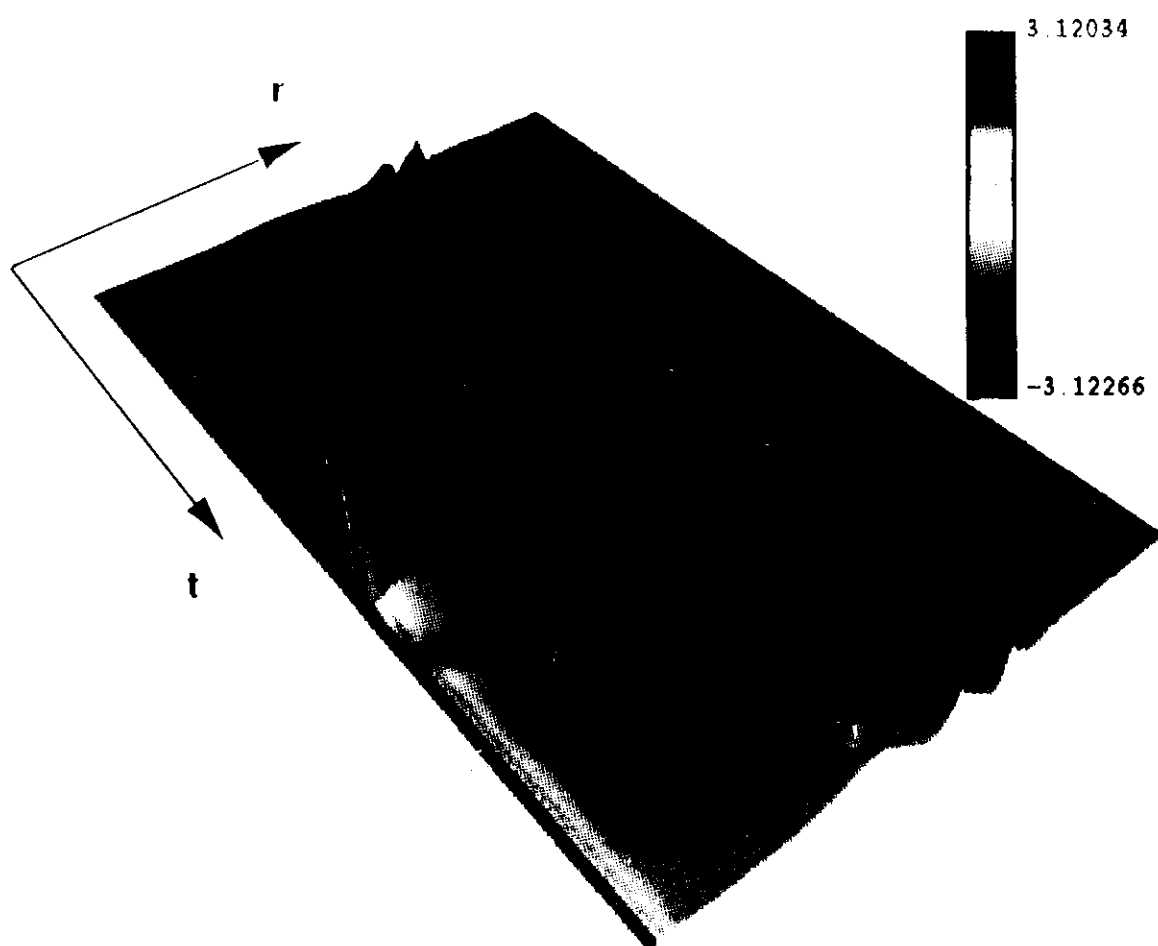


Energy density $\rho \sim \frac{1}{r^2} \sim \frac{1}{t^2}$

Much better linearization properties.



Sphaleron decay in the 4D SU(2) Higgs model:
the gauge field



Topology changing process in the
4D SU(2) Higgs model:
the gauge field

More on the Computation.

1. $\chi = \chi_{sph} + \delta\chi$
 $\varphi = \varphi_{sph} + \delta\varphi$
 $\pi_\chi = \delta\pi_\chi$

} parametrize by c_i

note: $\delta\chi(r=\infty) = \delta\chi(r=-\infty) = 0$
etc.

$\Rightarrow \exists c > 0 \ni \chi(r) = c.$

2. Evolve for a short time

3. Gauge transform to $\text{Arg } \chi = -\frac{\pi}{2}$

4. Evolve until system linearizes $\epsilon \cong \epsilon_{lin} \cong \sum_n \omega_n |a_n|$

5. Extract normal mode amplitudes

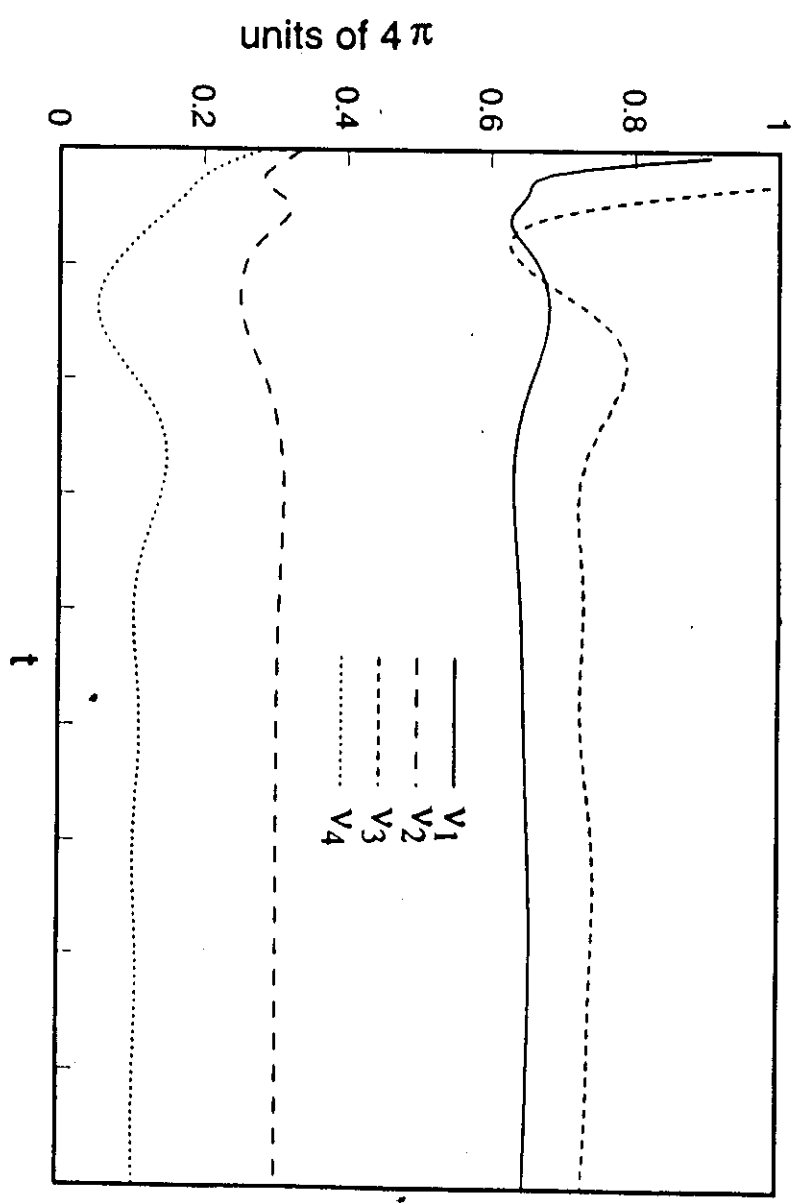
a_n and calculate $\nu = \sum_n |a_n|^2 \equiv \nu(c_i)$

- discretized normal modes numerically
- continuous normal modes analytically

4 types of modes.

1 of mass m_h (Higgs)

3 of mass m_w (gauge bosons)



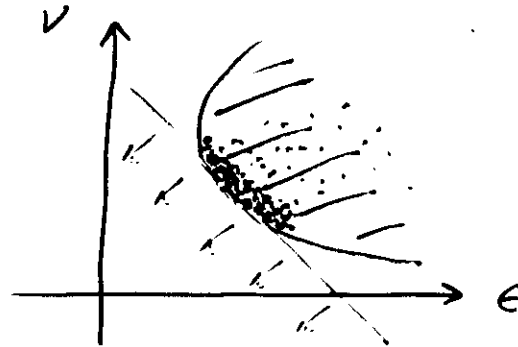
50 modes

$N = 511$

$\Delta t = 0.1$

$\lambda = 0.1$

Sample $\{C_i\}$ stochastically with a weight $e^{-\beta\epsilon - \mu\nu}$ favoring small ν .



$$\lambda = 0.1$$

$$\dim \text{ of } \{C_i\} = 20 \times 8$$

$$N = 256$$

$$\# \text{ normal modes} = 4 \times 50$$

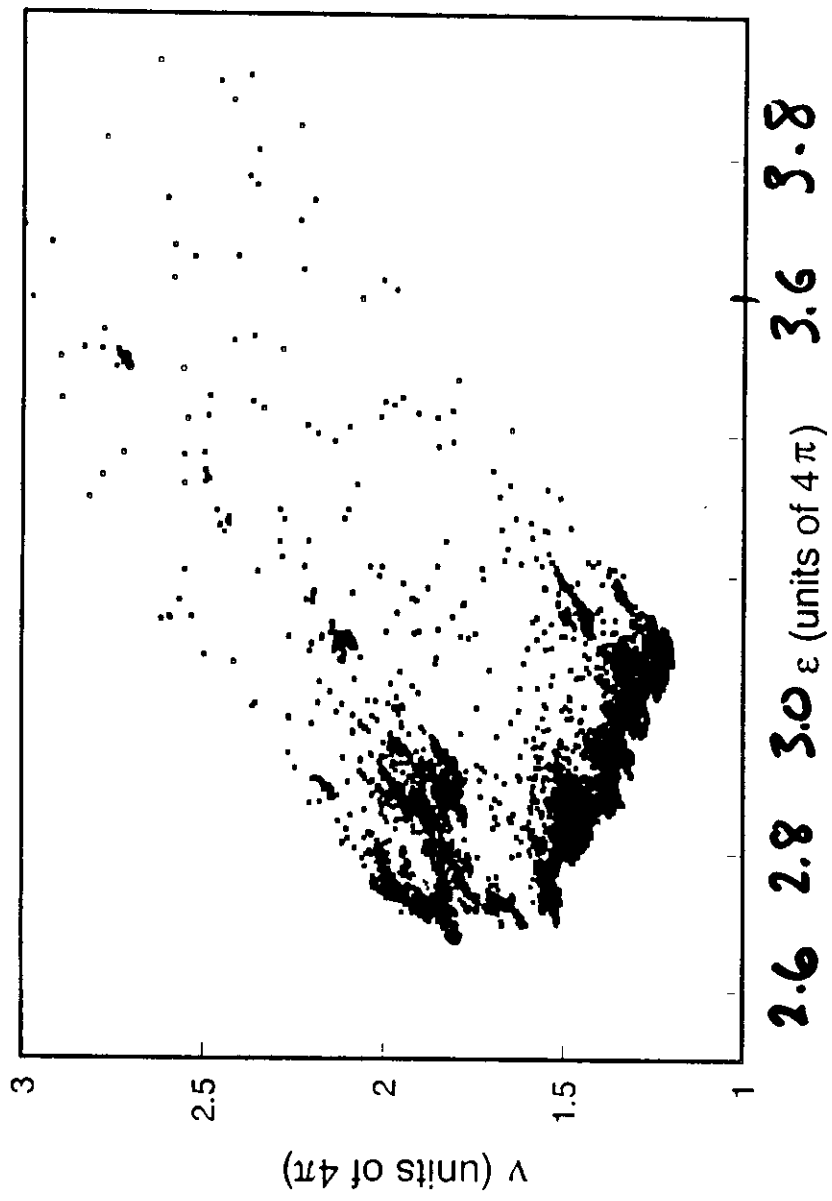
$$\Delta r = 0.2$$

Approximately 50,000 data points
(60 hrs CPU)

note: U.V. cutoff gives $\omega_{\max} \sim \frac{1}{\Delta r}$

In practice $\omega_{\max} = \frac{c}{\Delta r}$ $c < 1$

$$\therefore \min \frac{\nu}{\epsilon} \sim \frac{1}{\omega_{\max}} \sim \frac{\Delta r}{c}$$



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III. Plans for the future

- Our MC results are at their limit for our grid size
 - Larger lattice + bigger comp.
- Implement the full complex RST evolution. (below barrier)
 - initial testing w/ $1+1$ Dim.
Abelian Higgs complete.
(formidable numerical problem)
 - adapt to $3+1$ $SU(2)$ Higgs and see what happens.