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VI

SUMMER SCHOOL IN HIGH ENERGY PHYSICS AND COSMOLOGY

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**BASICS OF CFT, STRING COMPACTIFICATIONS, MIRROR SYMMETRY AND
SPACETIME TOPOLOGY CHANGE**

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Please note: These are preliminary notes intended for internal distribution only.

LECTURE V

• SUMMARY OF LECTURE IV:

- COMPLEX STRUCTURE PARAMETER SPACE OF M
DOES NOT APPEAR TO BE ISOMORPHIC TO
KÄHLER PARAMETER SPACE OF MIRROR W
 - RESOLUTION:
 - * KÄHLER SPACE $\rightarrow$ ENLARGED KÄHLER SPACE
 - * SOME CLASSICAL MATH. SINGS ARE
STILL PHYSICALLY SMOOTH IN CFT
- $\Rightarrow$ * "Mild" SPACETIME TOPOLOGY CHANGE

• LECTURE V

- CONSIDER SINGULARITIES WHICH ARE TRULY
BADLY BEHAVED AS CFT'S
 - USE STRUCTURE OF TYPE II STRINGS
TO SHOW PHYSICALLY SMOOTH
- $\Rightarrow$ * "DRASTIC" SPACETIME TOPOLOGY CHANGE

I) INTRODUCTION / REVIEW

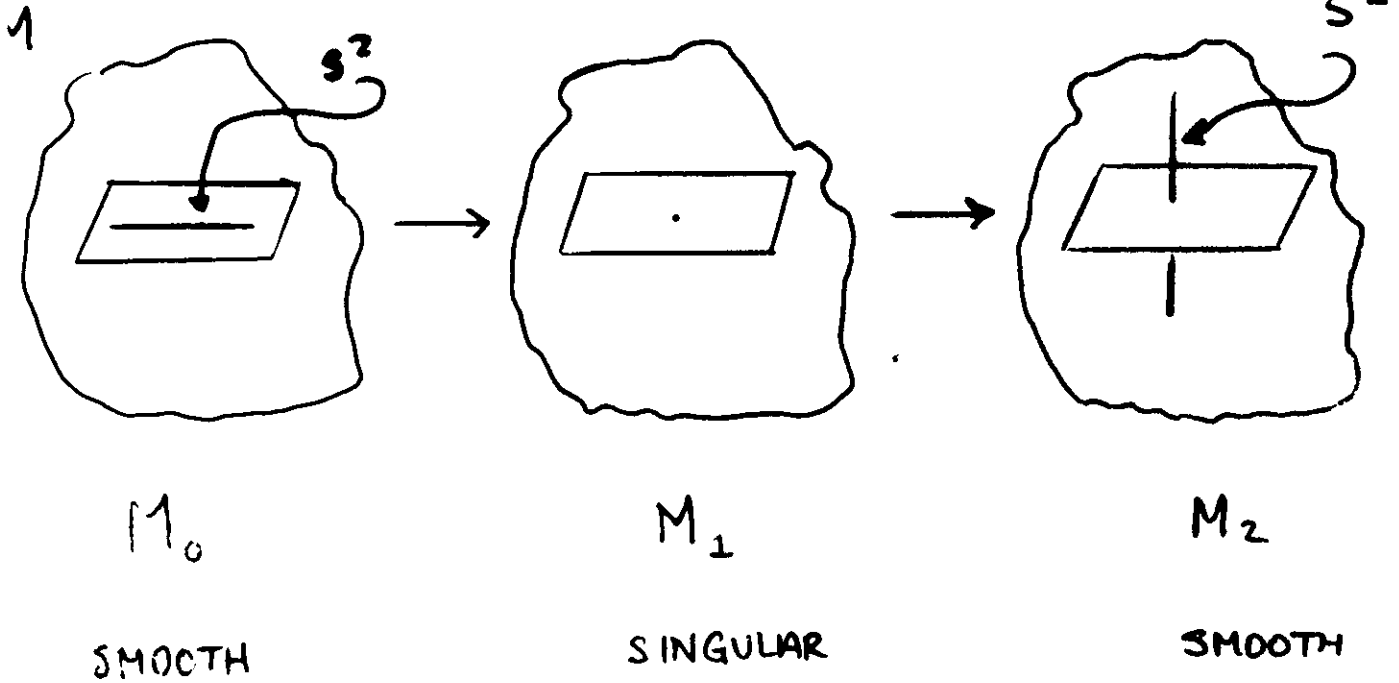
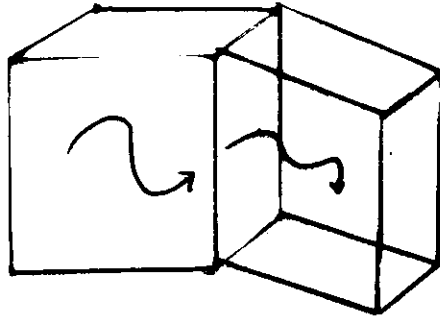
BASIC SET-UP:

- STRING VACUATION: $M_4 \times \{(2,2) C=9 \text{ SCF}\}$
- PARTICULAR HOMOGENEOUS: $M_4 \times \{C=9 \text{ SCF}\}$

PARAMETERS:

- MARGINAL OPERATORS: $\begin{cases} Q \cdot \bar{Q} = 1 \\ \bar{Q} \cdot Q = 1 \end{cases}$
- GEOMETRICAL DEFORMATIONS: $\begin{cases} \text{KAHLER STRUCTURE} \\ \text{COMPLEX STRUCTURE} \end{cases}$
- MIRROR SYMMETRY: $\begin{cases} KS_M \longleftrightarrow CS_W \\ CS_M \longleftrightarrow KS_W \end{cases}$

• UNDERSTANDING THE PHASES



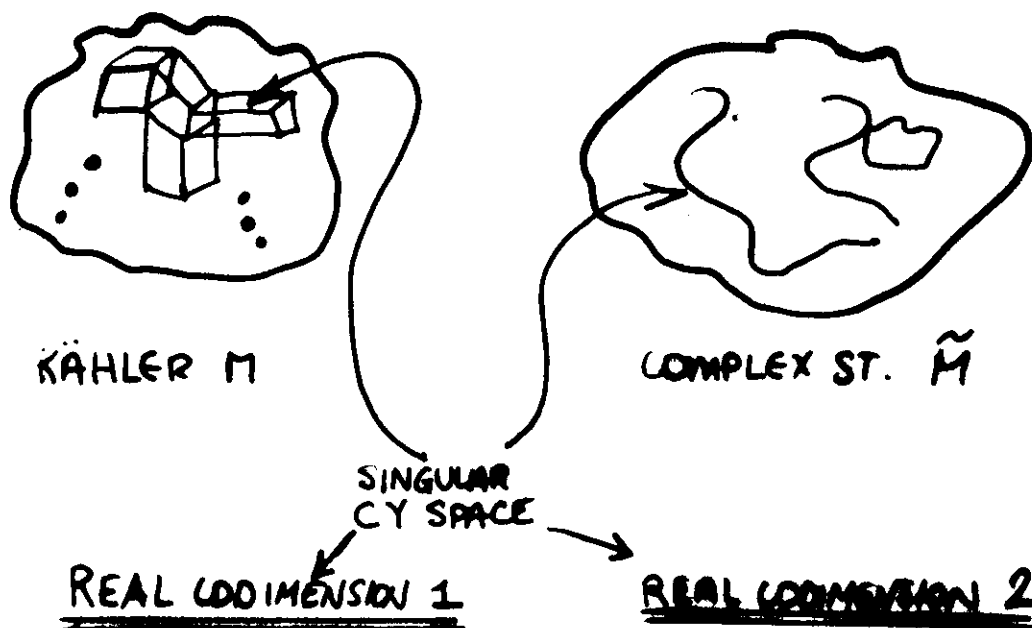
- MATHEMATICAL OPERATION: "FLOP"
- TOPOLOGY OF $M_2 \neq$ TOPOLOGY OF M_1
- CAN REPEAT $\rightarrow$ ADJOIN NUMEROUS REGIONS EACH BEING KÄHLER PARAMETER SPACE FOR "FLOPPED" MANIFOLDS

MATHEMATICALLY:

23.5
23.6
23.7

- ADJOINING ALL OF THESE KÄHLER CONES YIELDS "ENLARGED" KÄHLER MODULI SPACE
- TORIC VARIETY
- THIS OBJECT IS ISOMORPHIC MATHEMATICALLY TO COMPLEX STRUCTURE MODULI SPACE

PHYSICALLY: STILL A TROUBLING DISCREPANCY:

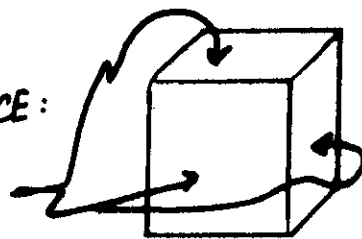


- RESOLUTION? DO WE TRUST THAT CFT'S ON LHS WALLS ARE SINGULAR? THESE ARE STRONGLY COUPLED THEORIES: α'/R^2 NOT SMALL.
- ON RHS: $\alpha'/\tilde{R}^2$ IS (CAN BE TAKEN) SMALL! STRONG COUPLING $\rightarrow$ WEAK COUPLING. $\rightarrow$ TRUST RHS - NONSINGULAR.

• THUS, BY MIRROR SYMMETRY HAVE LEARNED:

- MATHEMATICAL KÄHLER PARAMETER SPACE:

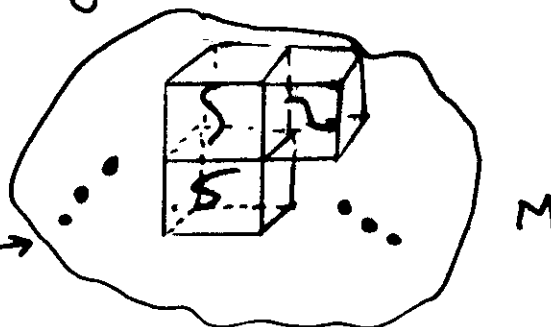
MATH. SINGULAR:



- PHYSICALLY: a) MUST ADJOIN ALONG COMMON WALLS:

• (AT MOST) PHYSICALLY SINGULAR

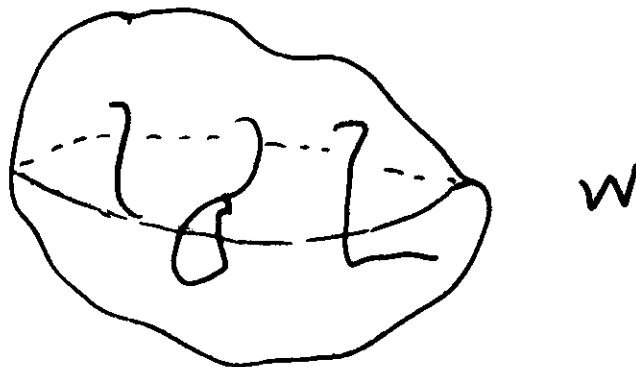
ON
CODIMENSION 1
SUBSET OF FACES →



b) MIRROR TO:

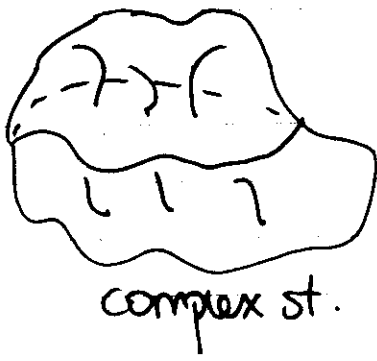
• (AT MOST) SINGULAR
ON REAL CODIMENSION 2

(TRUST PERT. TH*)



• $\begin{cases} \text{Is CFT REALLY SINGULAR AT THESE POINTS?} \\ \text{Is STRING THEORY SINGULAR " " " " } \end{cases}$

N=2 SUSY AND SPECIAL GEOMETRY

- TYPE II STRING ON $M_4 \times CY \Rightarrow N=2$ SUSY IN SPACETIME
- FOR IIB STRING: $h^{2,1} \leftrightarrow$ SPACETIME VECTOR MULTIPLTS
 $(A_\mu \quad \psi \quad \phi) \leftrightarrow U(1)^{h^{2,1}}$
- $h^{1,1} \leftrightarrow$ HYPERMULTIPLTS
 (FERMIONS; 2 complex scalars)
- Moduli Space: 

LOCAL GEOMETRY

$\updownarrow$

"SPECIAL GEOMETRY"

CONIFOLDS AND STROMINGERS RESOLUTION

• COORDINATES ON CS MODULI SPACE:

- (A_I, B^J) $I, J = 0, \dots, h^{2,1}$ SYMPLECTIC HOMOLOGY BASIS OF $H^3(M, \mathbb{Z})$.
 $(A_I \cap B^J = \delta_I^J; A_I \cap A_J = B_I \cap B_J = 0)$

- Let $z^J = \int_{B^J} \Omega$, $G_I = \int_{A_I} \Omega$; $\Omega = \text{HOL 3-FORM}$
 $\{z^J\} = \text{GOOD LOCAL (PROJECTIVE) COORDS ON } \mathcal{M}_{CS}$,
 $G_I = G_I(z)$.

• CONIFOLD POINT IN MODULI SPACE (ROUGHLY): SOME $z^J \rightarrow 0$

• MATHEMATICAL IMPLICATION: CURVATURE SINGULARITY:

- MONODROMY AROUND $z^J = 0$: $z^J \rightarrow z^J$
 $G_J \rightarrow G_J + z^J$
 $\Rightarrow$ NEAR $z^J = 0$, $G_J(z^J) = \frac{1}{2\pi i} z^J \ln z^J + \text{SINGULAR VALUE}$
- SPECIAL GEOMETRY : $K = -\ln(i \bar{z}^I G_I - i z^I \bar{G}_I)$
 METRIC : $g_{I\bar{J}}^{\mathcal{M}} = \partial_I \partial_{\bar{J}} K \Rightarrow g_{I\bar{J}}^{\mathcal{M}} \sim \ln(z^I \bar{z}^{\bar{J}})$

• PHYSICAL IMPLICATION:

- 4D LOW-ENERGY EFFECTIVE FIELD THEORY $\rightarrow$ SINGULAR
 NLTM ON $\mathcal{M}_{CS} \sim \int d^4x g_{I\bar{J}}^{\mathcal{M}} \partial \phi^I \partial \bar{\phi}^{\bar{J}}$

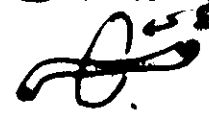
RESOLUTION:

• SINGULARITY IN $\mathcal{L}^{eff}$ $\leftrightarrow$ PREVIOUSLY MASSIVE STATES BECOME MASSLESS

• CURE SINGULARITY: INCLUDE THESE DEGREES OF FREEDOM IN $\mathcal{L}^{eff}$

• WHICH STATES?

• TYPE IIB: EXTREMAL CHARGED R-R BLACK HOLES

- 10-D: $\exists$ 3+1 BLACK-BRANE OF CHARGE $Q_{\Sigma_5} = \int_{\Sigma_5} F^{(5)}$
; ASSUME $Q = g_s \mathbb{Z}$.

- 4-D: WRAP 3-BRANE AROUND 3 CYCLES A_I, B^J ON CY .
 $\rightarrow$ 4-D CHARGED BLACK HOLES



• 4-D EM CHARGES: $\int_{A_I \times S^2} F^{(5)} = g_s \cdot n_I$; $\int_{B^J \times S^2} F^{(5)} = g_s \cdot m^J$

• 4-D EXTREMAL (BPS) MASS: (CRESOLE, D'ADAMO, FERARRA, INTRIGUEN)

$$M = g_s c^{K/2} |m^I G_I - n_I z^I|$$

• CONSIDER: $n_I = \delta_{IJ}$ $m^I = 0$ FOR J FIXED WITH $z^J \rightarrow 0$.

$\Rightarrow M \rightarrow 0$ AT CONIFOLD POINT.

• WILSONIAN $\mathcal{L}^{eff}$ CONTAINING MASSLESS $N=2$ B. HOLE Hyp.: SMOOTH

• INTEGRATE OUT B. HOLE $\rightarrow$ SHOULD RECOVER SINGULAR THEORY:

RUNNING COUPLING $\rightarrow \tau_{JJ} \sim \frac{1}{2\pi i} \ln z^J + \text{SING. VALUED}$

• $N=2 \Rightarrow \tau_{JJ} = 2G_J \Rightarrow G_J = \frac{1}{2\pi i} z^J \ln z^J + \text{SING. VALUED}$
 $\Rightarrow$ REPRODUCED SINGULARITY $\Rightarrow$ INCLUDING BH'S RESIDUES CONIFOLD

II) CONIFOLD TRANSITIONS & TOPOLOGY CHANGE

(BRG, D. MORRISON, A. STROMINGER)

• BASIC IDEA: CONSIDER A LESS "GENERIC" CONIFOLD POINT IN $\mathcal{M}^{c.s.}$ FOR WHICH

- MORE THAN 1, SAY P , 3-CYCLES DEGENERATE
- THESE P -3CYCLES ARE NOT HOMOLOGICALLY INDEPENDENT BUT SATISFY R HOMOLOGY RELATIONS.

• CLAIM: • BOSONIC POTENTIAL FOR SCALARS IN BLACK HOLE HYPERMULTIPLETS H^a $a=1, \dots, P$ HAS R FLAT DIRECTIONS.

- MOVING ALONG SUCH FLAT DIRECTIONS TAKES US TO ANOTHER BRANCH OF TYPE II STRING MODULI SPACE CORRESPONDING TO PROPAGATION ON A TOPOLOGICALLY DISTINCT CY SPACE.

• CLARIFYING REMARK:

- IN PREVIOUS RESULTS YIELDING TOPOLOGY CHANGE (WITTEN; ASPINWALL, BRG, MORRISON) HODGE #'S STAY FIXED - ONLY MORE SUBTLE TOP. INVARIANTS CHANGE
- IN PRESENT CONTEXT: HODGE #'S JUMP $(h^{2,1}, h^{1,1}) \mapsto (h^{2,1} - P + R, h^{1,1} + R)$ EVEN THOUGH PHYSICS IS SMOOTH.

• USEFUL BACKGROUND:

A) CONIFOLD POINTS IN $\mathcal{M}^{CS}$: (LEFSCHETZ, CANDELAS, de la Ossa, CANDELAS, GREEN, HUBSCH)

• DISCRIMINANT LOCUS ($P=0=dP$ FOR CY $P=0$)
 $\rightarrow$ SINGULAR CY'S

• CONIFOLD LOCUS: SINGULARITIES ON CY ARE NODES:

• LOCALLY $N=4$: $\sum_{i=1}^4 (w_i)^2 = 0$; SINGULAR AT APEX.
 AS $C=dC=0$.

• HOMOGENEOUS EQN $\Rightarrow (\lambda w_1, \dots, \lambda w_4)$ SATISFY IF $(w_1, w_2) \neq 0$
 $\Rightarrow$ CONE

• BASE OF CONE: INT. CONE WITH SPHERE IN $\mathbb{R}^8$
 $\sum_{i=1}^4 |w_i|^2 = r^2$

• WRITE: $\vec{w} = (w_1, w_2, w_3, w_4) = \vec{x} + i\vec{y}$

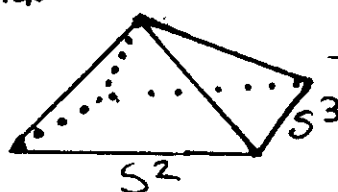
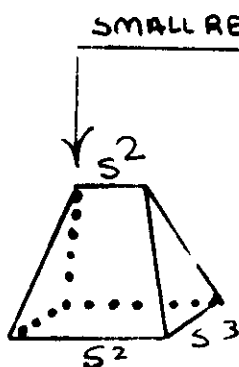
THEN EQNS: $\vec{x} \cdot \vec{x} = r^2/2$; $\vec{y} \cdot \vec{y} = r^2/2$; $\vec{x} \cdot \vec{y} = 0$

LOCUS: $S^3_{r/\sqrt{2}}$; $S^2_{r/\sqrt{2}}$ OVER S^3

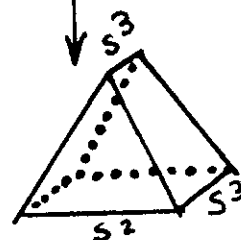
$\rightarrow$ BASE OF CONE $= S^2 \times S^3$

$$\epsilon=0 \rightarrow XY-UV=0$$

$$\begin{cases} XY-UV=0 \\ [X, \dots, V] \times [\lambda_1, \lambda_2] \\ \begin{pmatrix} X & U \\ V & Y \end{pmatrix} \begin{pmatrix} \lambda_1 \\ \lambda_2 \end{pmatrix} = 0 \end{cases}$$



$$\text{DEFORM: } \sum (w_i)^2 = \epsilon^2$$



• MONODROMY:

IF γ^a $a=1, \dots, k$ ARE k -VANISHING 3 CYCLES
 AT A CONIFOLD POINT, THEN A 3-cycle δ UNDERGOES
 MONODROMY $\delta \mapsto \delta + \sum_{a=1}^k (\delta \cap \gamma^a) \gamma^a$ ABOUT
 THIS PT IN $\mathcal{M}^3$.

B) TYPE IIB STRING THEORY ON A CY 3-FOLD

$\mathcal{N}=2$ SUSY: • $h^{2,1}$ U(1) VECTOR MULTIPLETS (+ GRAVIPHOTON)
 $h^{1,1} + 1$ HYPERMULTIPLETS (INCLUDING AXION/DILATON)

EXAMPLE:

- BEGIN WITH QUINTIC IN $\mathbb{P}^4$ ($h^{1,1}=1$
 $h^{2,1}=101$)
- DEFORM COMPLEX STRUCTURE TO 86 DIM
LOCUS

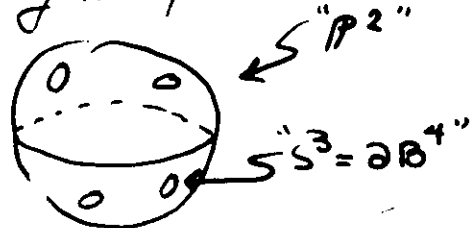
$$p(x) = x_1 g(x) + x_2 h(x) = 0$$

(g, h QUARTICS)

NOTE:

- THIS IS SINGULAR: $p \in dP$ VANISH AT
 $\{x_1 = x_2 = g(x) = h(x) = 0\}$ I.E. 16 POINTS
 (LYING ON $\mathbb{P}^2$ $x_1 = x_2 = 0$ IN $\mathbb{P}^4$)
- THESE ARE NODES (MATRIX OF 2nd DERIVATIVE NONDEG.)
- THE 16 S^3 's WHICH HAVE DEGENERATED, γ^a $a=1, \dots, 16$
 ARE NOT ALL HOMOLOGICALLY INDEPENDENT. RATHER
 $\sum_{a=1}^{16} \gamma^a = 0$.

$$\partial \left[\mathbb{P}^2 - \sum_{a=1}^{16} B^a \right] = \sum_{a=1}^{16} \gamma^a$$



BLACK HOLE RESOLUTION

SINGULARITY:

- BASIS $(A_I, B^J) : I, J = 1, \dots, 204$

- VANISHING CYCLES $\gamma^1, \dots, \gamma^{16}$ WITH $\sum \gamma^a = 0 \Rightarrow$ CAUTAKE

$$\gamma^1, \dots, \gamma^{15} = B^1, \dots, B^{15}; \quad \gamma^{16} = - \sum_{I=1}^{15} B^I$$

- AS USUAL:

$$z^I = \int_{B^I} \Omega \quad ; \quad G_J = \int_{A_J} \Omega$$

- MONODROMY AROUND $\{z^1 = z^2 = \dots = z^{15} = 0\}$

$$\delta \rightarrow \delta + (\delta \cap \gamma^a) \gamma^a \quad a = 1, \dots, 16$$

$$\int_{\delta} \Omega \sim \frac{1}{2\pi i} \sum_{a=1}^{16} (\delta \cap \gamma^a) \left(\int_{\gamma^a} \Omega \right) (\ln \int_{\gamma^a} \Omega) + \text{single valued}$$

$$\Rightarrow G_J = \int_{A_J} \Omega \sim \frac{1}{2\pi i} z^J \ln z^J + (1) \left(- \sum_{I=1}^{15} z^I \right) \left(\ln \left(\sum_{I=1}^{15} z^I \right) \right)$$

+ SINGLE VALUED

(AS BEFORE THIS DETERMINES SINGULARITY VIA SPECIFIC GEOMETRY)

- MASSLESS STATES: 16 DEGENERATING 3-cycles
 $\Rightarrow$ 16 MASSLESS BH HYPERMULTIPLETS
 $H^a \quad a=1, \dots, 16$

• WILSONIAN THEORY SMOOTH.

c. DO WE RECOVER KNOWN SINGULARITY IF WE INTEGRATE OUT BH MULTIPLETS?

• NEED BH CHARGES UNDER $U(1)^{15}$

- H^a ARISES FROM γ^a : $U(1)_J$ CHARGE = $Q_J^a = A_J \cap \gamma^a$
 AS TAKE $F^{(5)} = [\sum \alpha^I F_I^{(2)}]_{S.D.}$ w/ $\alpha^I = \text{dual of } A_I$

$$\Rightarrow Q_J^a = \delta_J^a \quad 1 \leq a \leq 15; \quad Q_J^{16} = -1 \quad \forall J.$$

• RUNNING COUPLES FROM INTEGRATING OUT THE H^a :

$$T_{IJ} \sim \frac{1}{2\pi i} \sum_{a=1}^{16} Q_I^a Q_J^a \ln(m^a)^{+i\epsilon}, \quad m^a \sim e^{K/2} |Q_I^a Z^I|$$

$$\Rightarrow T_{IJ} \sim \frac{1}{2\pi i} \delta_{IJ} \ln Z^J + \frac{1}{2\pi i} \ln \left(\sum_{K=1}^{15} Z^K \right) + \text{S. VALUES}$$

Now: $T_{IJ} = \partial_I G_J$ (N=2 SPECIAL GEOM) $\Rightarrow$

$$G_J \sim \frac{1}{2\pi i} Z^J \ln Z^J + \frac{1}{2\pi i} \left(\sum_{I=1}^{15} Z^I \right) \left(\ln \left(\sum_{I=1}^{15} Z^I \right) \right) + \text{S.V.}$$

LOW ENERGY EFFECTIVE ACTION WITH BLACK HOLES

MAIN QUESTION: WHAT IS PHYSICAL SIGNIFICANCE OF NONTRIVIAL HOMOLOGY RELATIONS BETWEEN VANISHING CYCLES?

• SCALAR BH POTENTIAL:

$$E_{\alpha\beta}^I = \left(\sum_{a=1}^{16} Q_a^I \epsilon_{\alpha\gamma} h_\gamma^{(a)} h_\beta^{(a)} - (\alpha \leftrightarrow \beta) \right); \quad I=1, \dots, 15$$

WITH $h_1^{(a)}, h_2^{(a)} \in H^{(a)} \quad a=1, \dots, 16.$

$$V \sim \sum E_{\alpha\beta}^I E_I^{\alpha\beta}$$

• FLAT DIRECTIONS?

$$\langle h_i^{(a)} \rangle = 0; \quad \langle \text{SCALARS IN 15 VECTORS} \rangle \neq 0$$

• BLACK HOLES GET MASS; MOVE BACK TO SMOOTH QUINTIC

$$Q_a^I \sim A_I \cap \gamma^a \quad \omega / \sum_{a=1}^{16} \gamma^a = 0 \Rightarrow \sum_a Q_a^I = 0 \quad \forall I$$

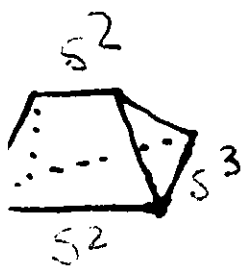
$$\Rightarrow \langle h_a^{(a)} \rangle = V^a \quad \forall a \text{ IS FLAT (UNIQUE UP TO GAUGE!)} \quad S^2$$



• 15 VECTORS / 15 HYPERS GET MASS; 1 HYPER MASSLESS

$$h_{0,0}^{2,1} = 101 \rightarrow 101 - 15 = 86; \quad h_{0,0}^{1,1} = 1 + 1 \rightarrow 2 + 1$$

PRECISELY HODGE NUMBERS GET FROM SMALL RESOLUTION!
(CANDELA, GREEN, HUBSCH)



• $\mathcal{H}$: Homology relation among γ^a

$\Rightarrow$ FLAT DIRECTION IN BH POTENTIAL

• MOVING ALONG SUCH FLAT DIRECTION TAKES US SMOOTHLY TO NEW BRANCH OF TYPE II STRING MODULI SPACE.

• APPARENTLY PHYSICALLY REALIZED CY TRANSITIONS DISCUSSED SOME YEARS AGO IN INSIGHTFUL PAPERS OF CANDELAS, GREEN, HUBSCH.

• GENERAL CASE:

• P VANISHING CYCLES

• R HOMOMOLOGY RELATIONS $\Rightarrow P-R$ U(1)'s

• P BLACK HOLE HYPERMULTIPLETS

• R FLAT DIRECTIONS IN BH POTENTIAL

• (GENERIC) DEFORMATION:

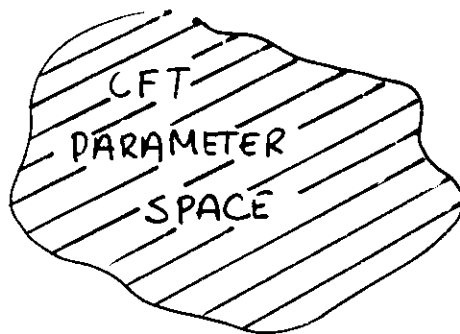
• LOSE $P-R$ VECTORS

• GAIN R HYPERMULTIPLETS

$$\Rightarrow (h^{2'}, h'') \mapsto (h^{2'} - (P-R), h'' + R) \quad (\Delta\chi = 2P)$$

MODULI SPACE OF $N=2$ SCFT'S

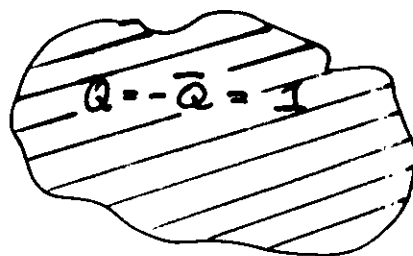
CFT



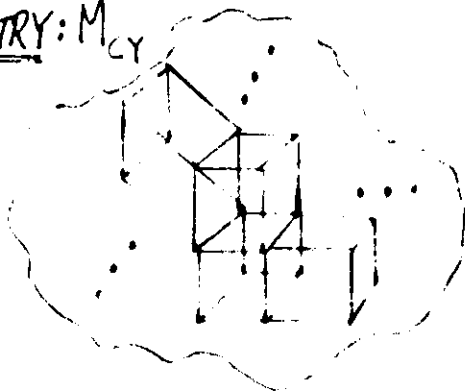
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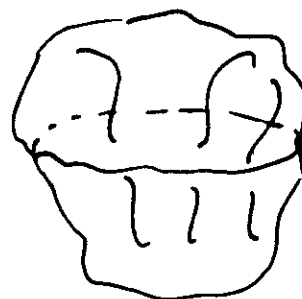
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GEOMETRY: M_{CY}



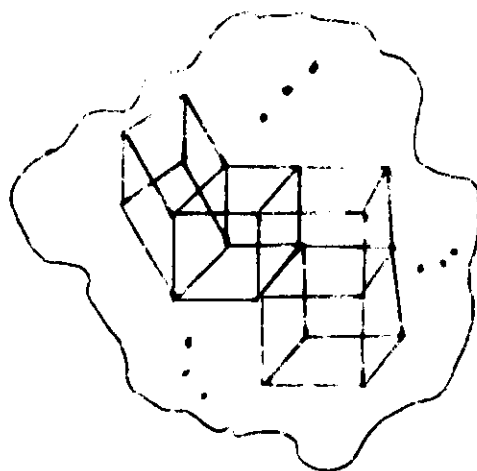
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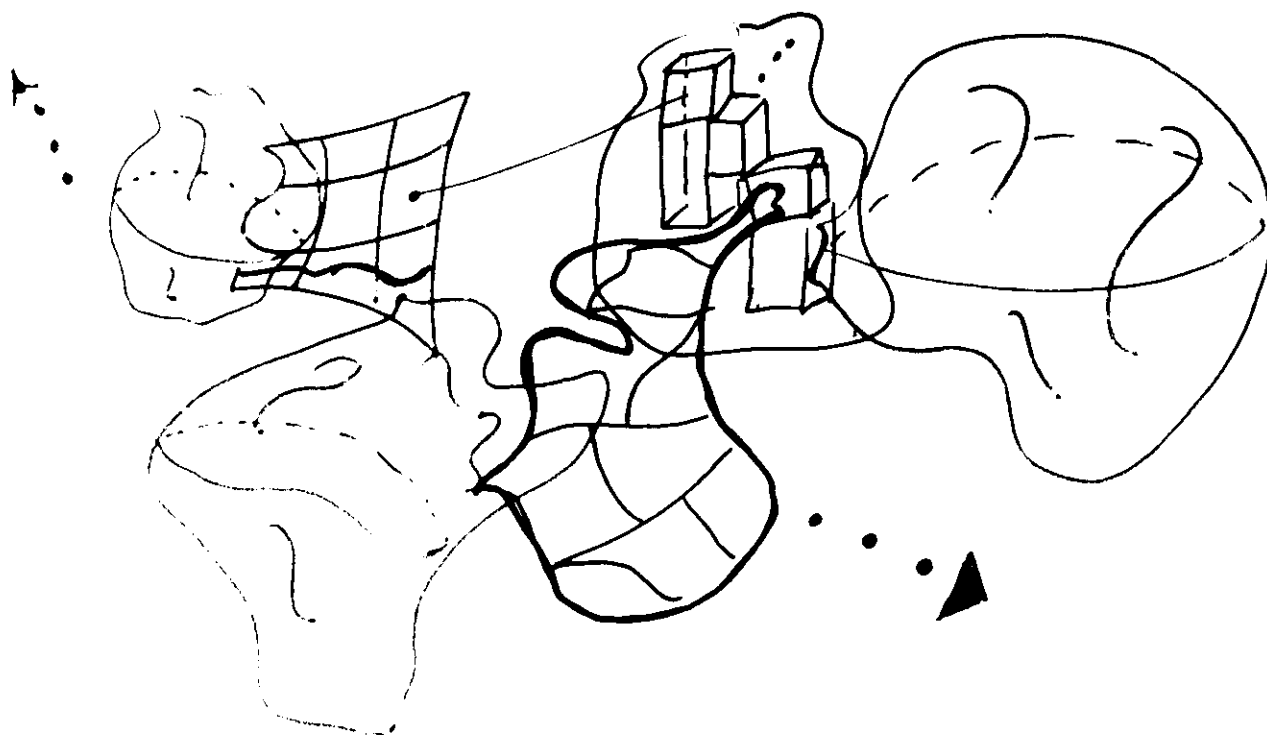
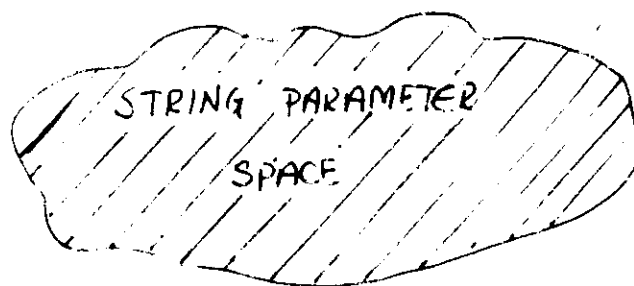
GEOMETRY: $\tilde{M}_{CY}$



$\times$



MODULI SPACE OF TYPE II STRING



III) IMPLICATIONS - CONCLUSIONS

- *1) TOWARDS UNIFICATION OF STRING VACUA
- *2) PHYSICALLY SMOOTH TOPOLOGY CHANGE
OF MORE "IDEALIC" SCET.
- *3) MIRROR SYMMETRY
- *4) BLACK HOLES AS ELEMENTARY PARTICLES