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SMR.959 - 41

MINIWORKSHOP ON STRONG ELECTRON CORRELATIONS
"Disorder and Interaction in Quantum Systems
and Their Classical Analogs"

(1 - 19 July 1996)

"Non-Fermi-liquid behaviour
in correlated electron systems"

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These are preliminary lecture notes, intended only for distribution to participants.

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Non-Fermi-liquid behavior

in correlated electron systems

H.v. Löhneysen

Introduction:

- heavy fermions and Kondo effect

Non-Fermi liquid at a magnetic instability:

- $\text{CeCu}_{6-x}\text{Au}_x$: $T_N(x) \geq 0$ for $x \geq x_c$
- pressure tuning
- effect of magnetic field

Non-Fermi liquid arising from single-ion effects

- two-channel Kondo effect
- distribution of T_K at metal-insul. transition

Summary and outlook

Benedict Bogenberger

Steffen Mock

Thomas Pietrus

Günther Portisch

Hans-Georg Schlager

Almut Schröder

Michael Sieck

Oliver Stockert

Thomas Trappmann

Matthias Waffenschmidt

Frank Huster

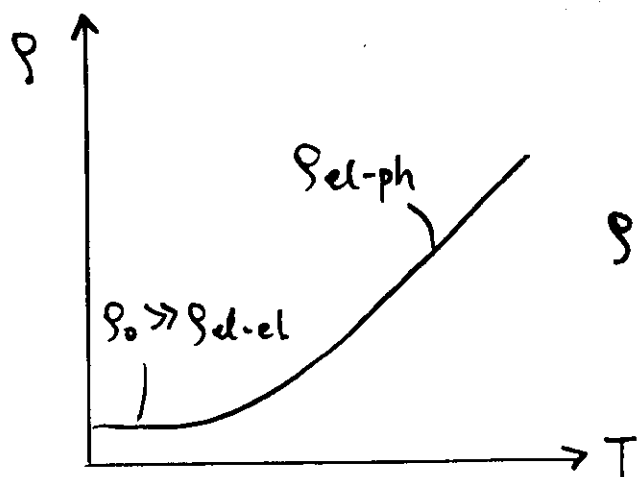
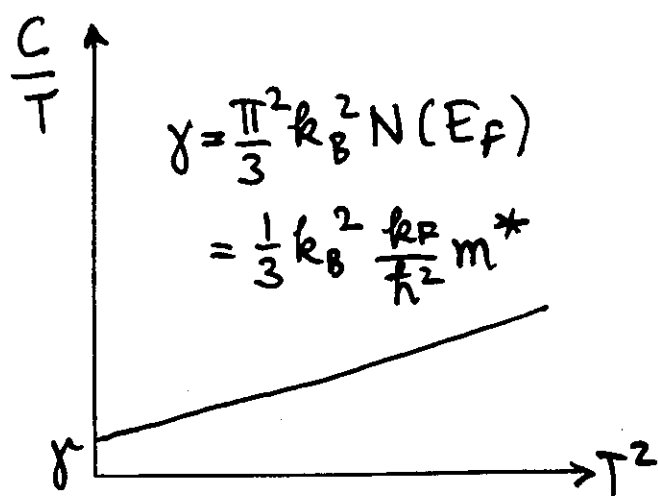
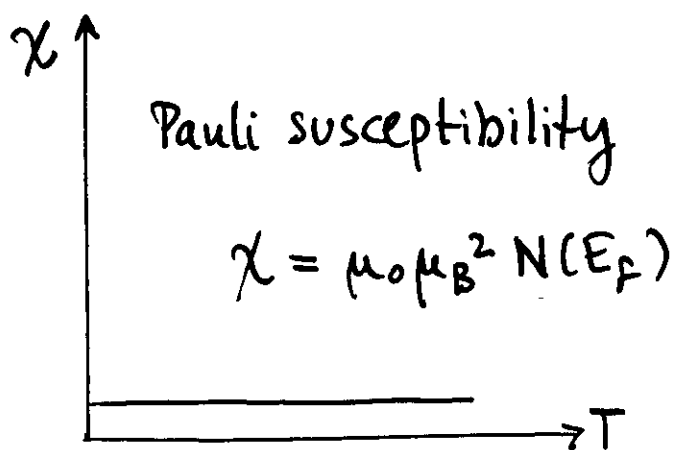
Anja Neubert

M. Loewenhaupt, J. Lynn, N. Pyka

B. Lüthi, F. Bruls, ...

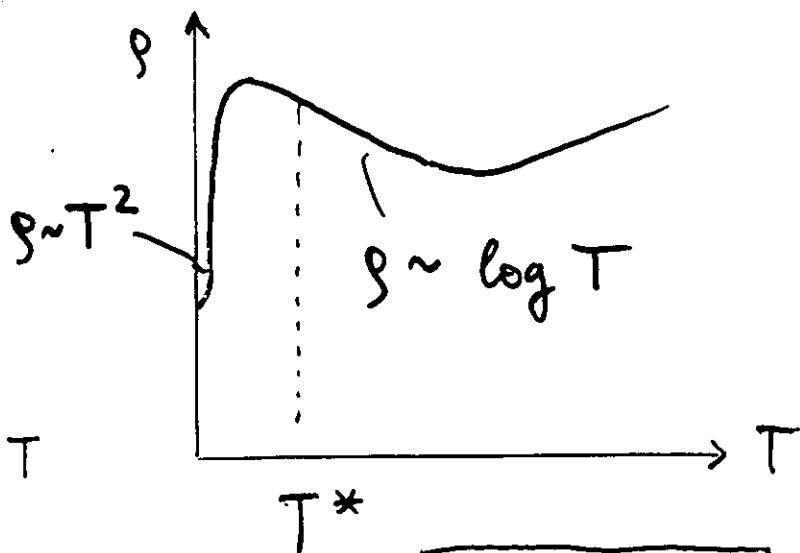
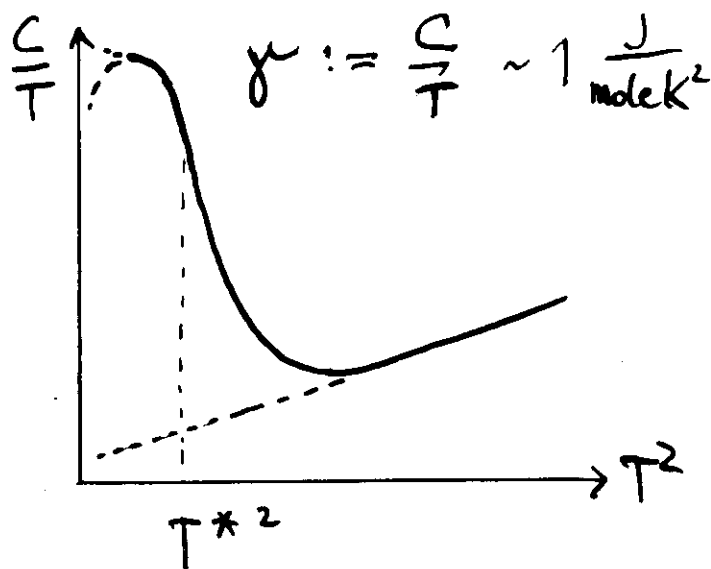
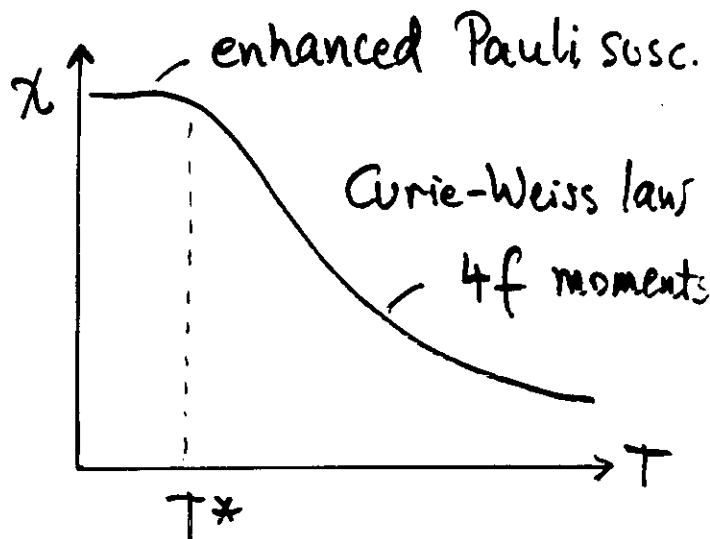
Introduction to heavy-fermion metals

Normal metal



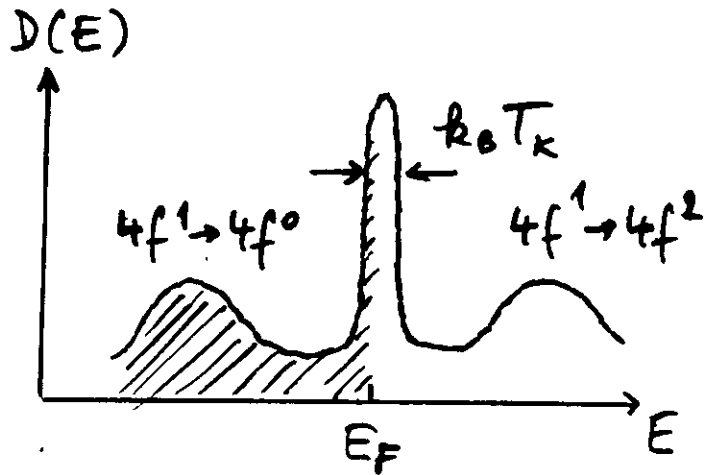
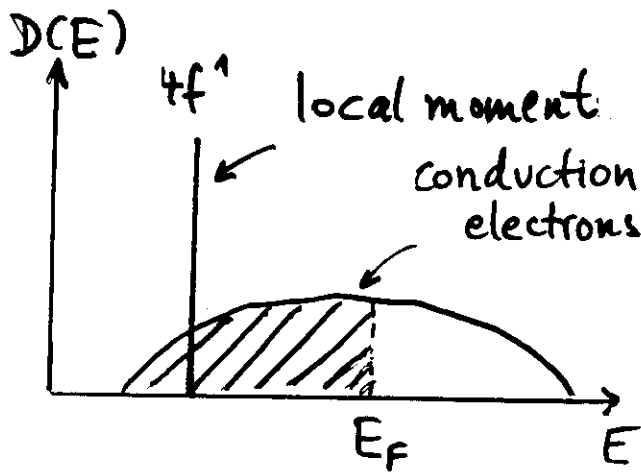
$$m^* \approx m_0$$

heavy-fermion metal



$m^* \approx 100 m_0$

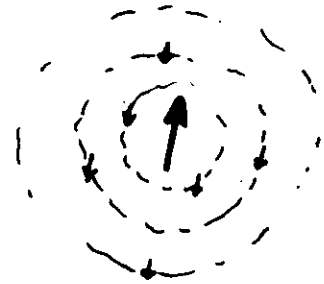
Kondo effect



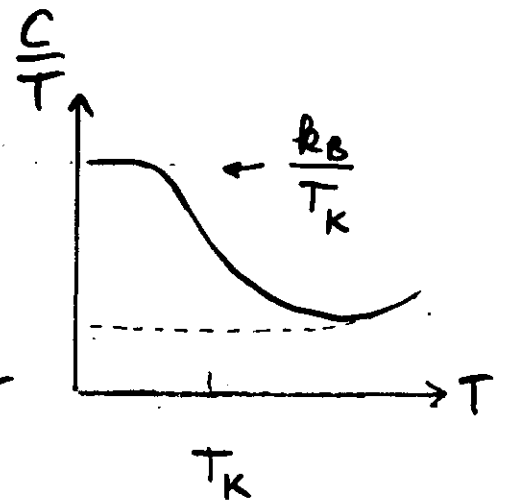
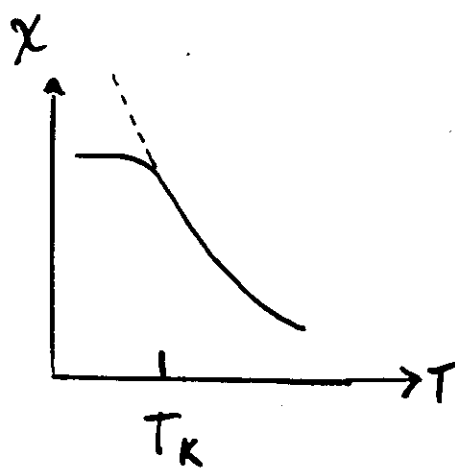
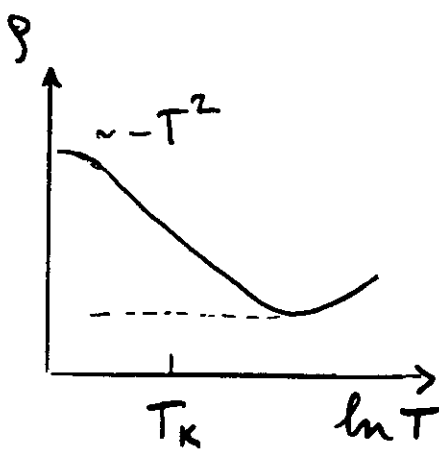
Hybridization between $4f$ and conduction electrons
 + $4f$ on-site correlations

↓
 "Resonance" near E_F builds up at low T

Screening of $4f$ moment
 by conduction electrons



Kondo anomalies at low T :



local Fermi liquid Nozieres

$T_K = 10\text{K} : \frac{C}{T} \sim 0.8 \frac{\text{J}}{\text{mol} \cdot \text{K}^2}$

Fermi liquid description of heavy-fermion metals

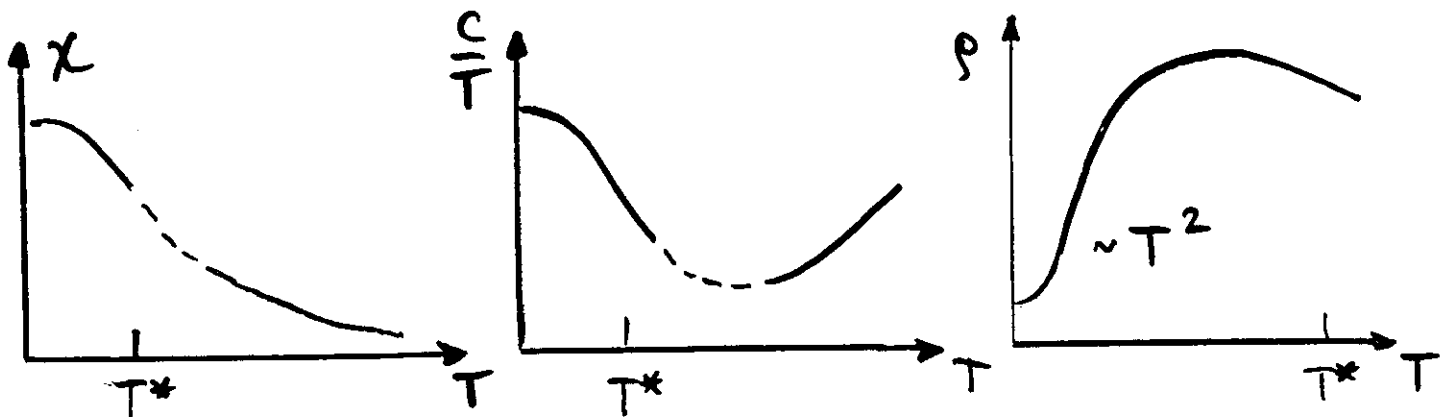
Fermi liquid (FL)
for interacting fermions

- one-to-one correspondence between excitations of FL and noninteracting Fermi gas
- transport equation for QP

Heavy-fermion metals:

$T \gg T^* \approx T_K$: conduction electrons, local moments

$T \ll T^*$: strongly interacting quasiparticles



Fermi liquid (Landau)

$$\gamma = \frac{C}{T} \sim m^* \sim \frac{1}{T^*}$$

$$\chi \sim \frac{m^*}{1 + F_0^a}$$

$$\Delta \rho = AT^2 \sim \left(\frac{T}{T^*}\right)^2 \sim \gamma^2 T^2$$

Non-Fermi liquid at a magnetic instability

heavy-fermion systems

competition between nonmagnetic
and magnetic groundstates

local singlet formation $T_K \sim \exp\left(-\frac{1}{N(E_F)J}\right)$

RKKY interaction

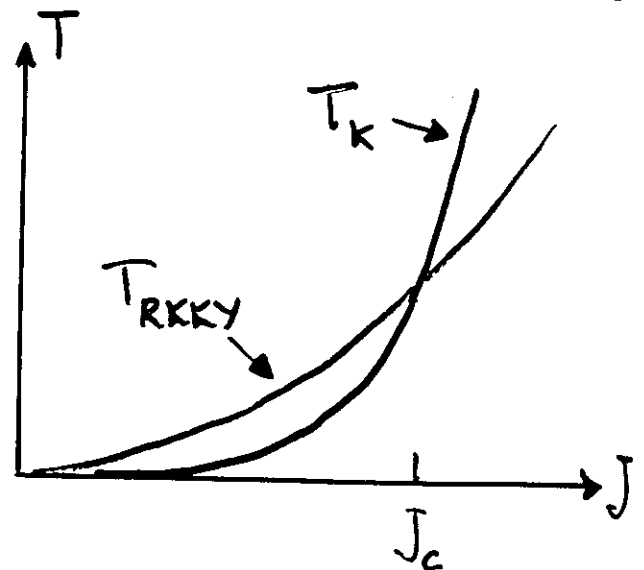
$$T_{\text{RKKY}} \sim J^2 N(E_F)$$

Doniach

$$0 \leq J \ll J_c$$

local moments
magnetic order

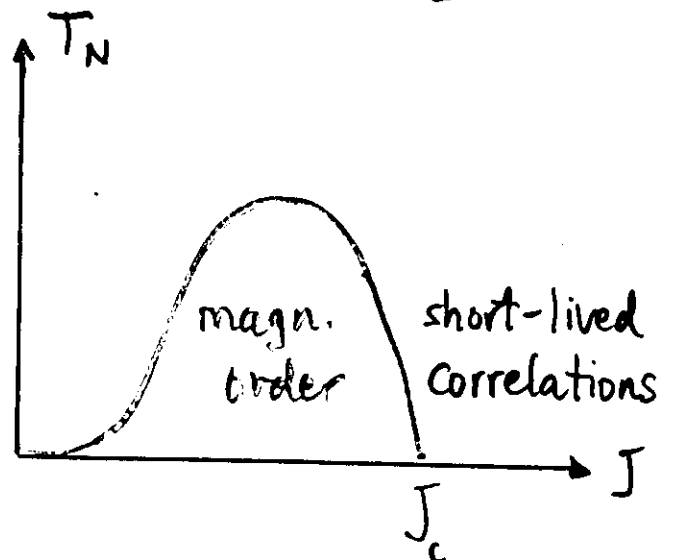
"small" Fermi volume



$$J > J_c$$

local singlets

"large" Fermi volume

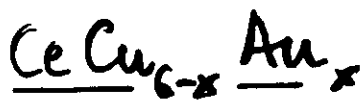


what happens at J_c ?

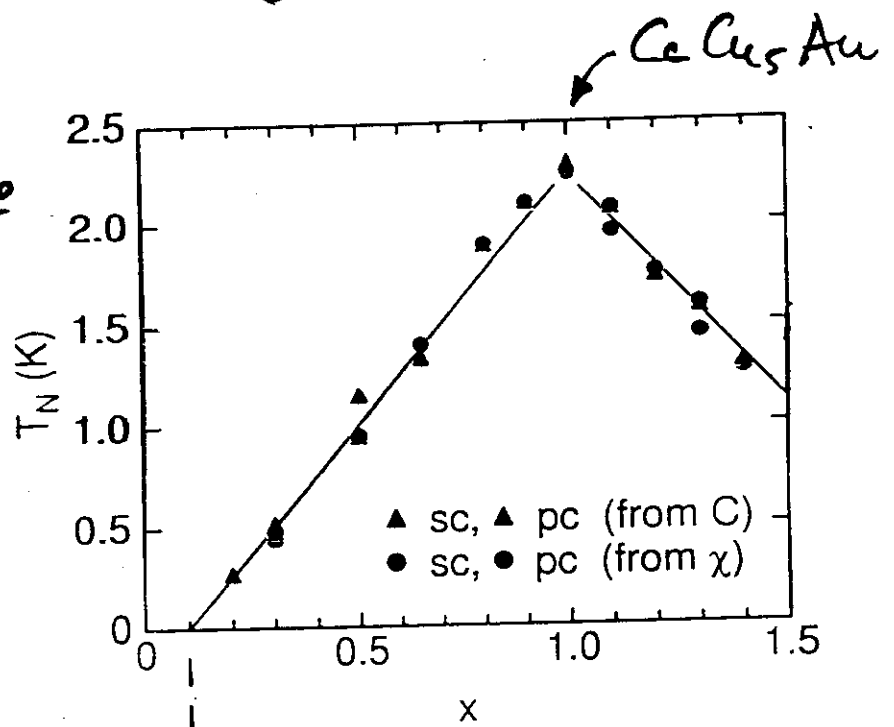
Magnetic instability in $CeCu_{6-x}Au_x$

$CeCu_6$: heavy fermions with $\gamma = 1.6 \text{ J/mol K}^2$
nonmagnetic groundstate Onuki, Amato,
 short-lived short-range AF correlations
 Aeppli, Rossat-Mignod

Alloying with Au: long-range AF order



neg. lattice pressure
 explains $T_N(x)$
 for $x < 1$



Fermi liquid
no magn. order

magn. order

Fermi liquid
 +
 local moments

$T_K \approx 6K$ $x_c \approx 0.1$

magnetic instability $J = J_c$

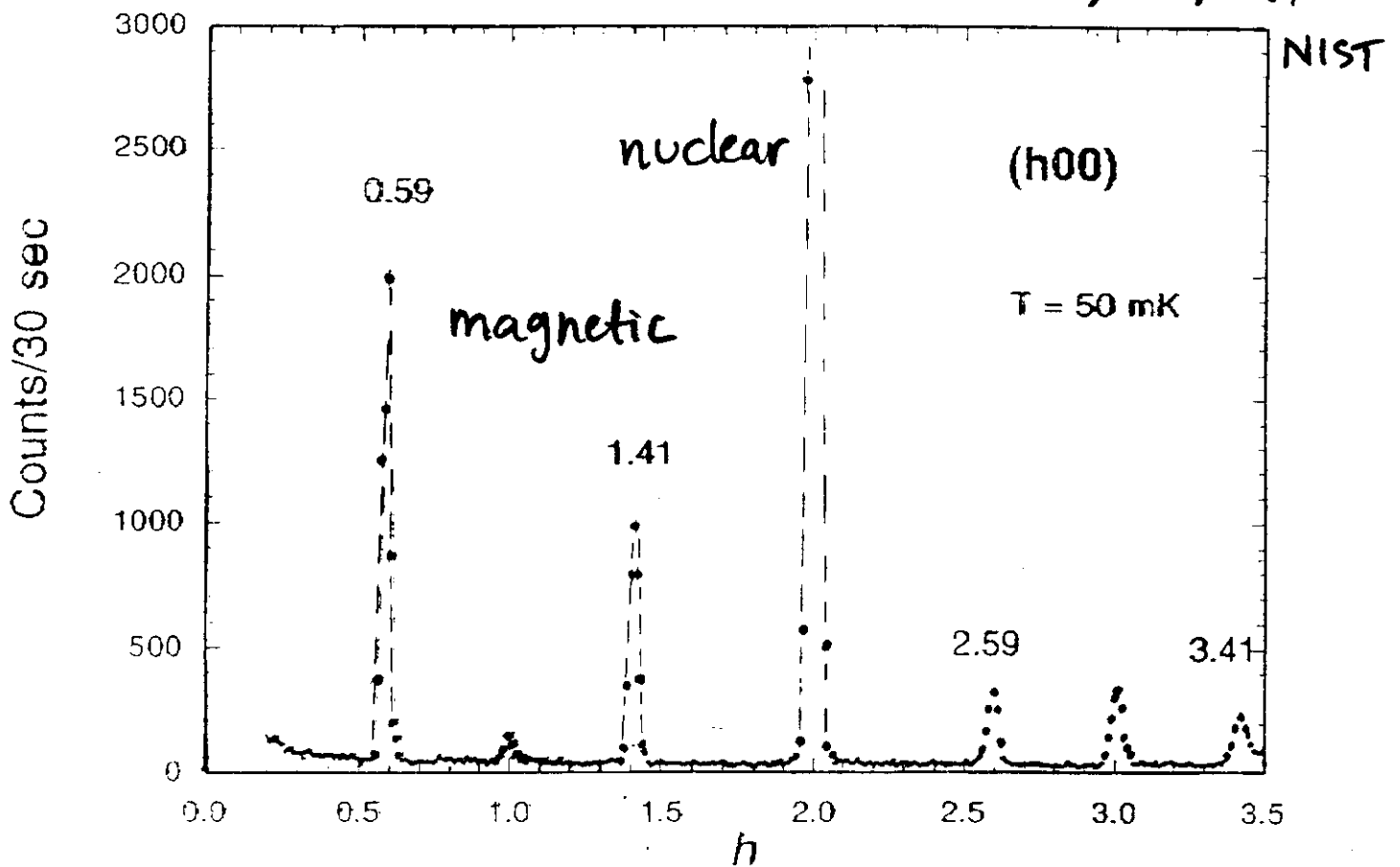
phase transition at $T = 0$?

Elastic neutron scattering on a CeCu_{5.5}Au_{0.5}
single crystal

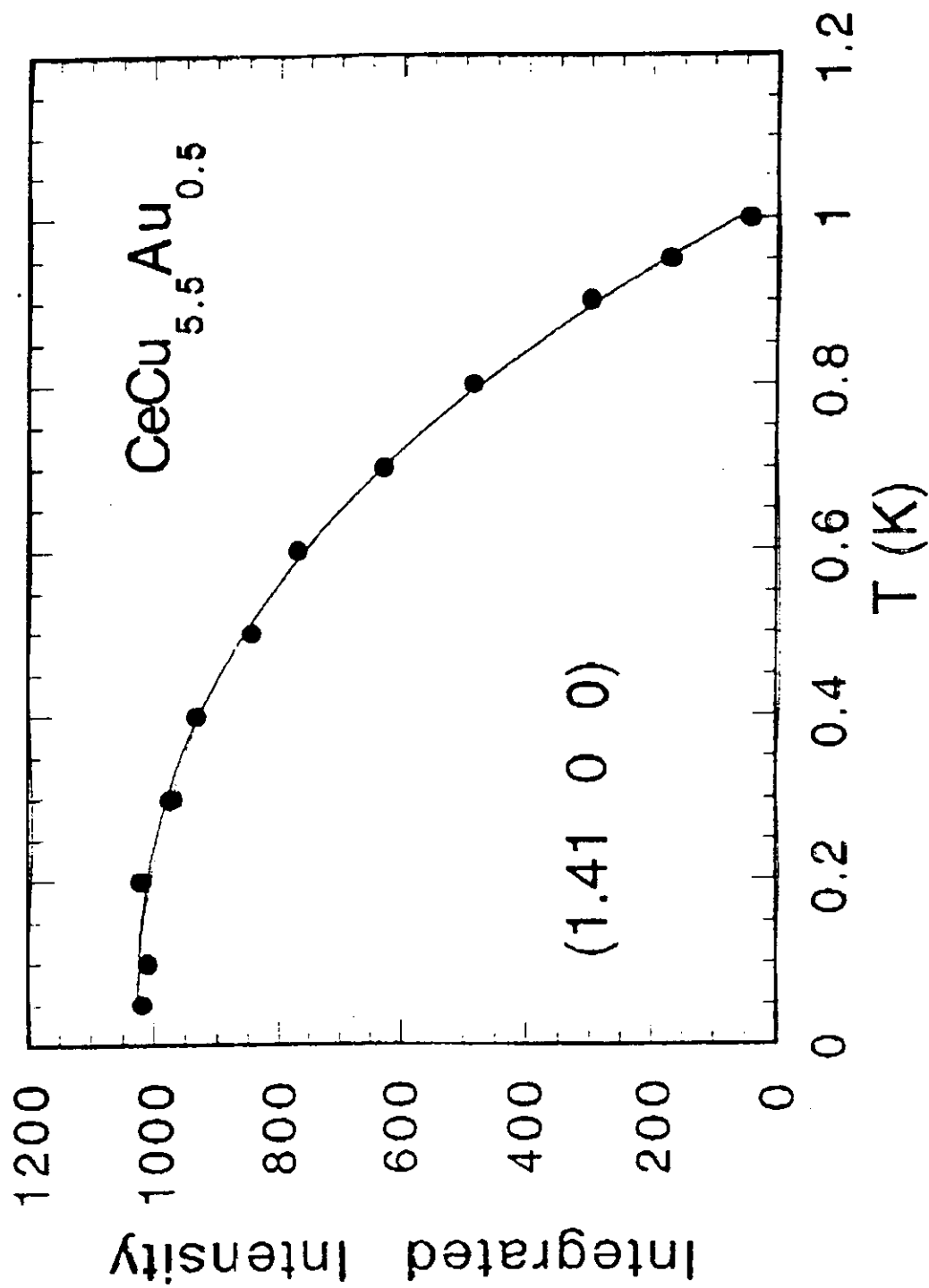
incommensurate magnetic reflections

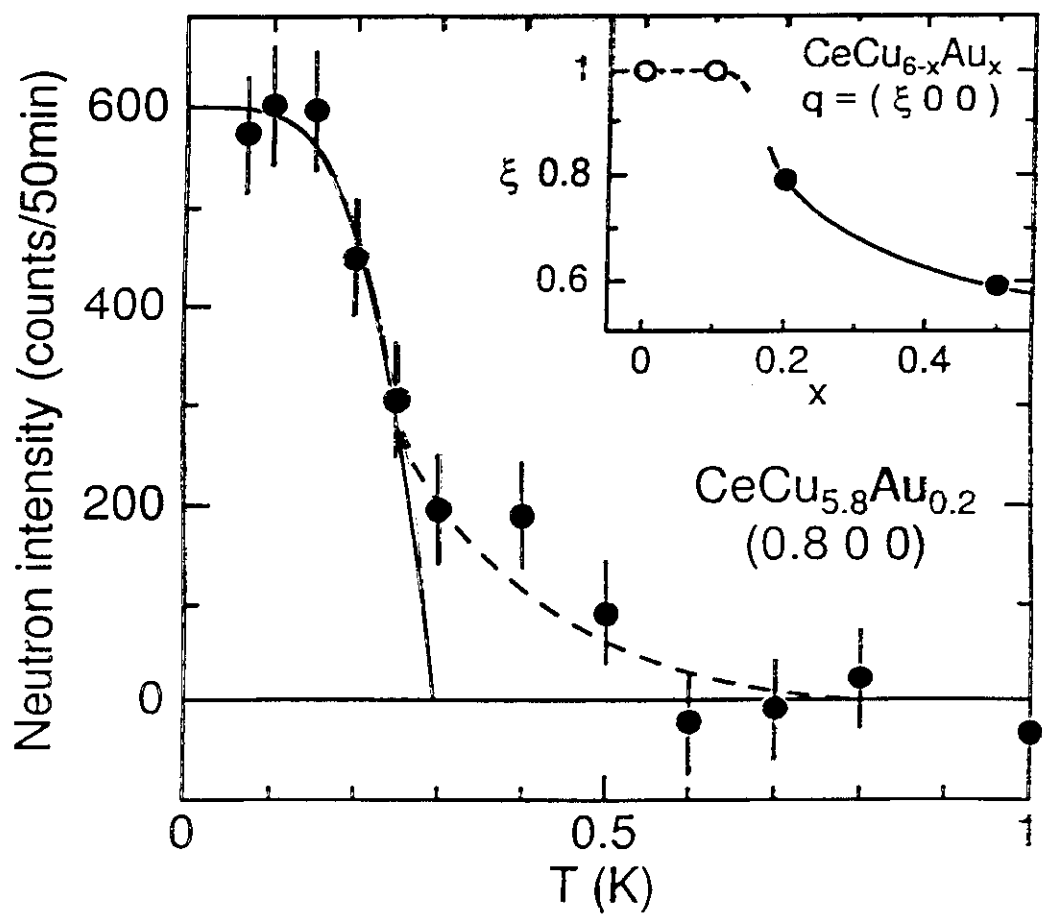
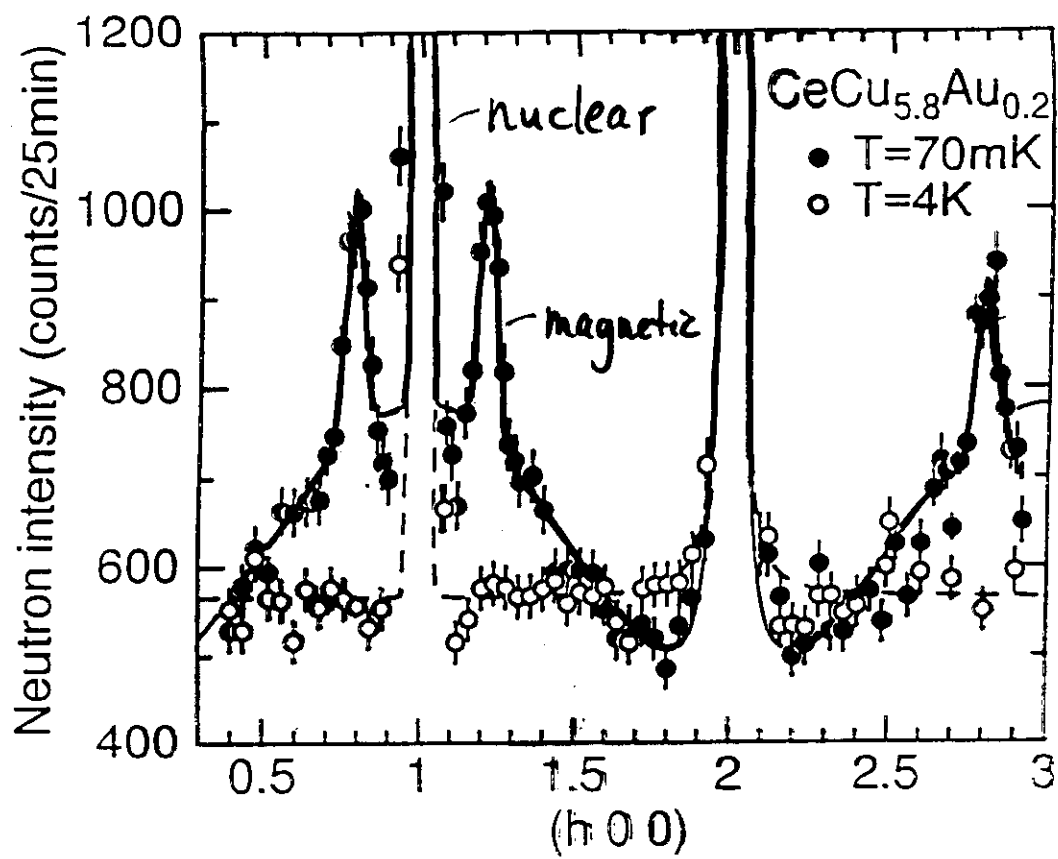
$$(\pm 0.59 \ 0 \ 0)$$

A. Schröder, J. Lynn, ...

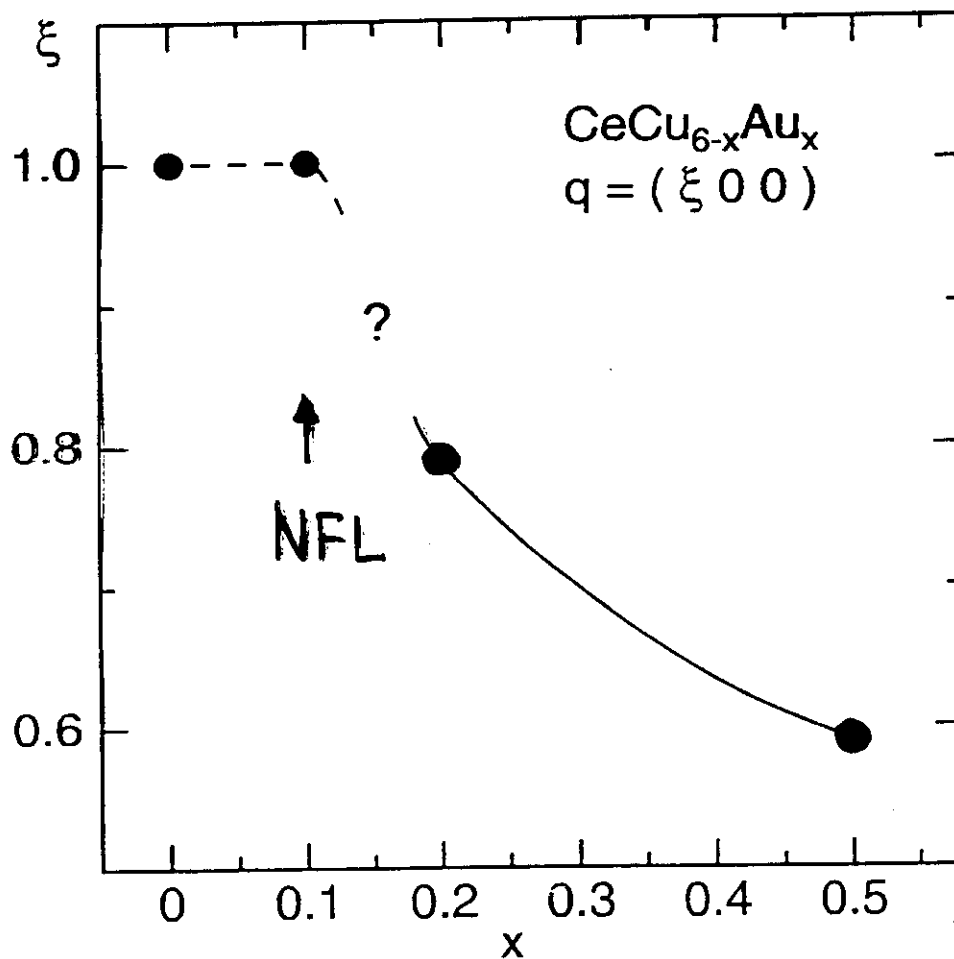


different from (100) and $(0 \ 1 \pm 0.15 \ 0)$
inelastic correlations in CeCu₆



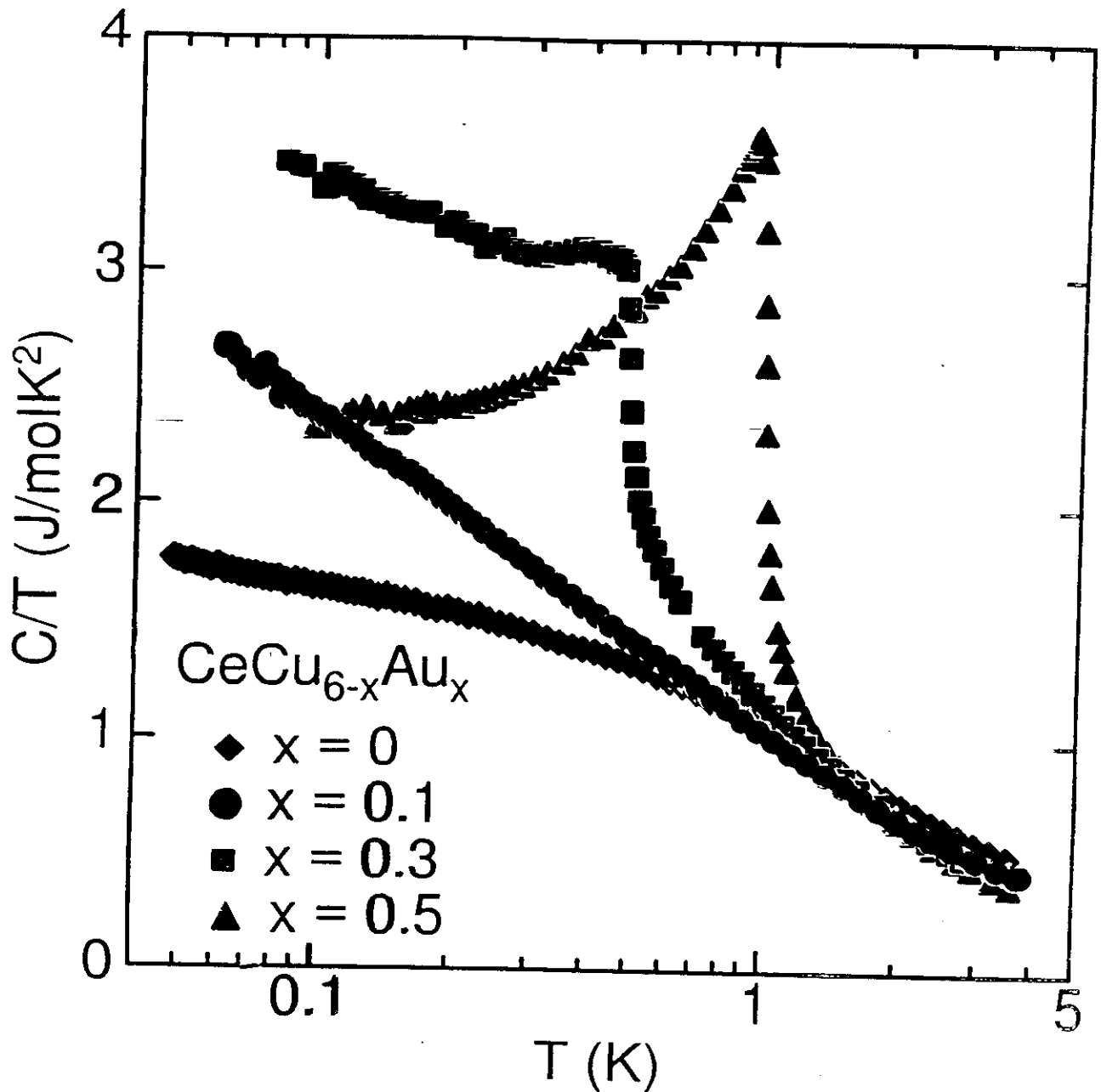


Evolution of magnetic order
out of short-range correlations



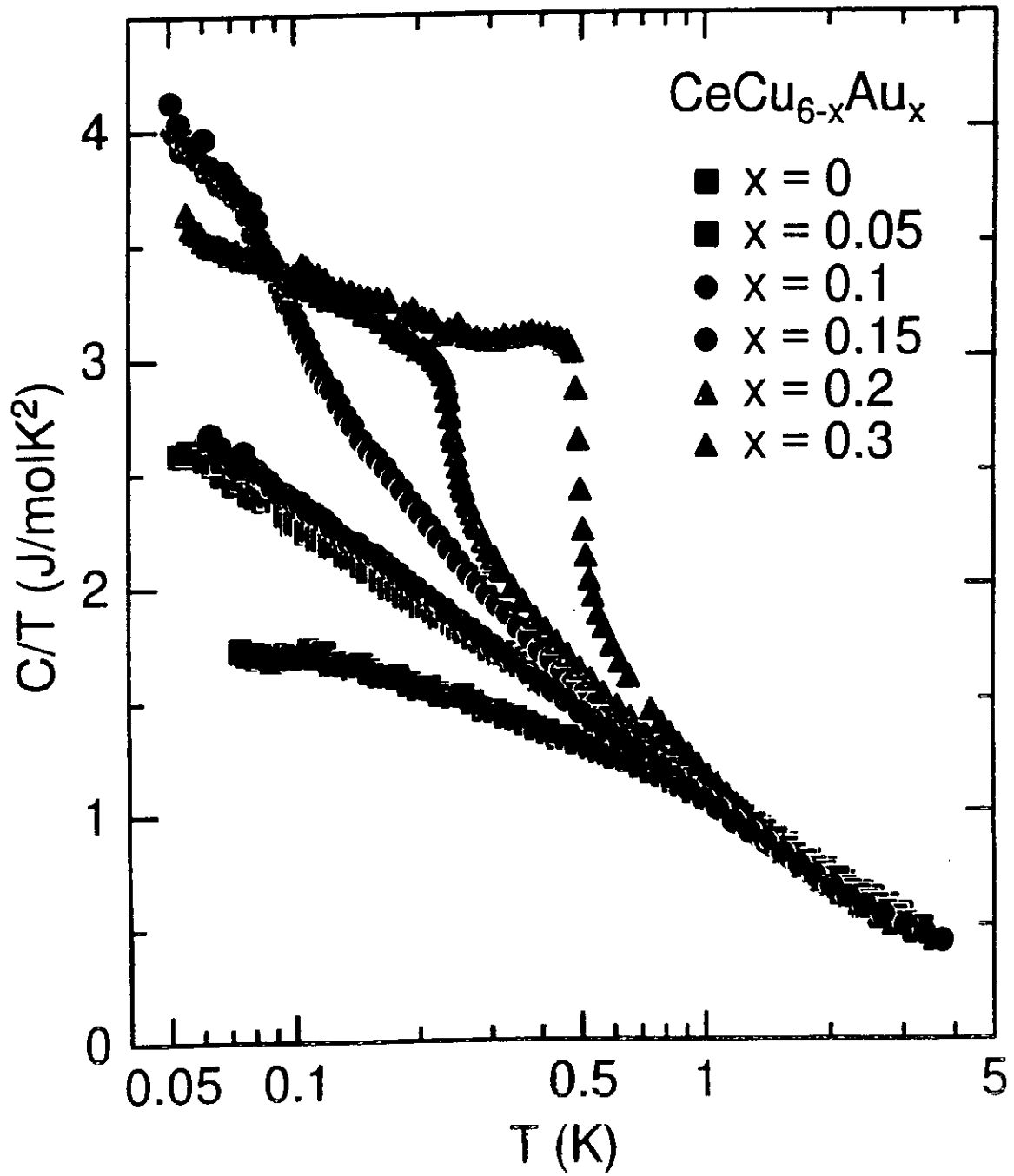
cf. CeRu₂Si₂ → La_xCe_{1-x}Ru₂Si₂
Q = (0.3 0 0)

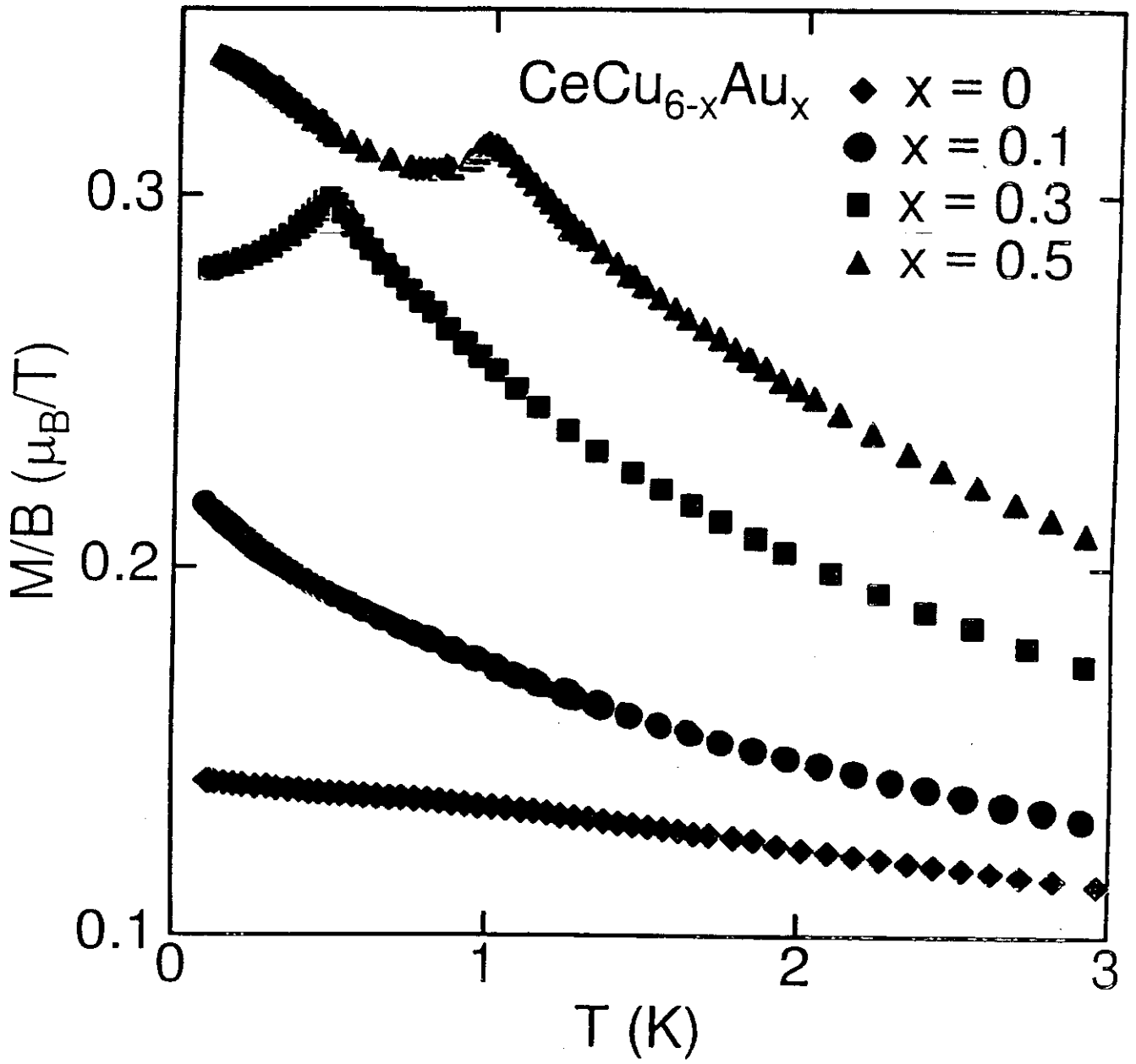
Specific heat of $\text{CeCu}_{6-x}\text{Au}_x$ single crystals

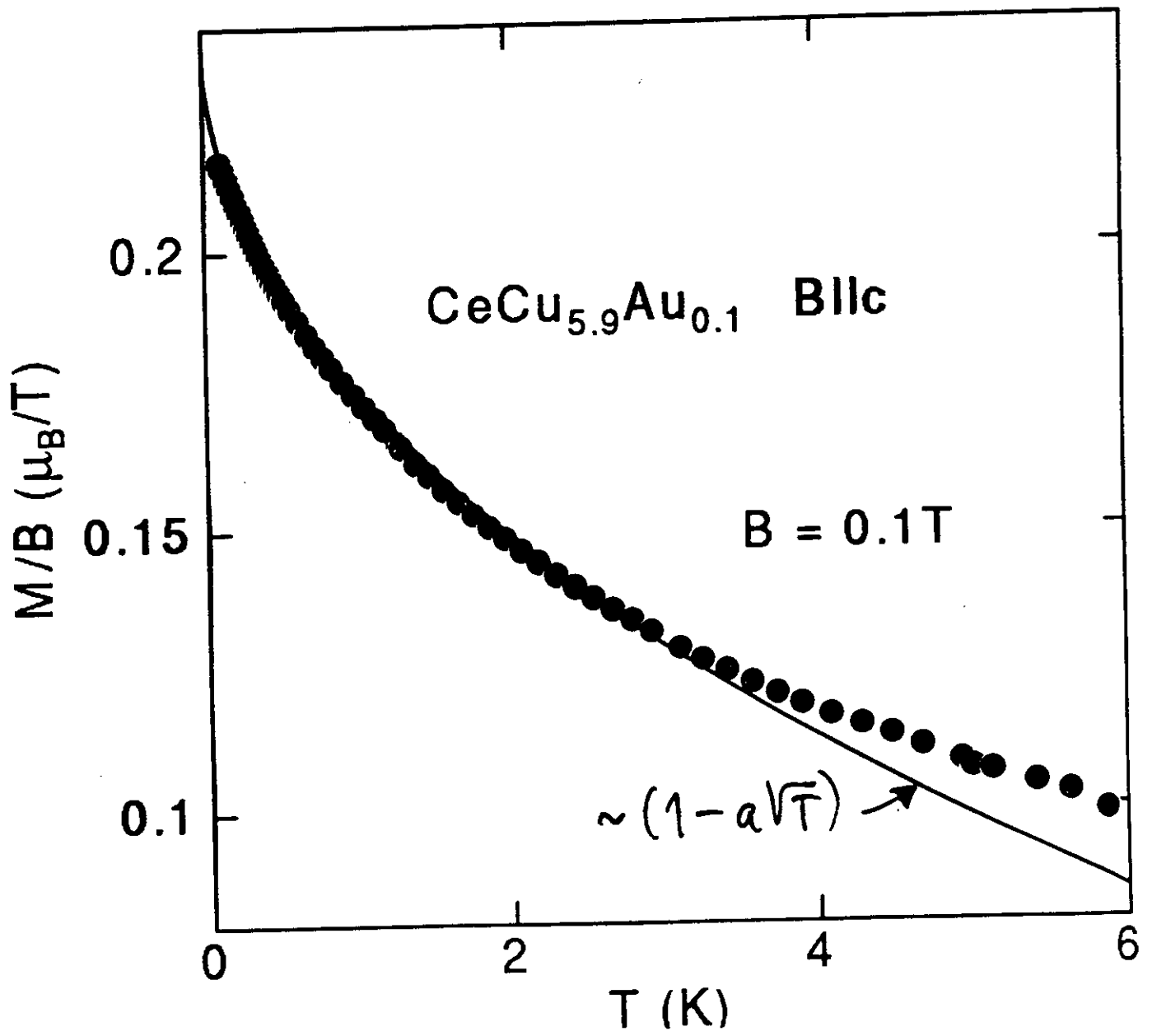


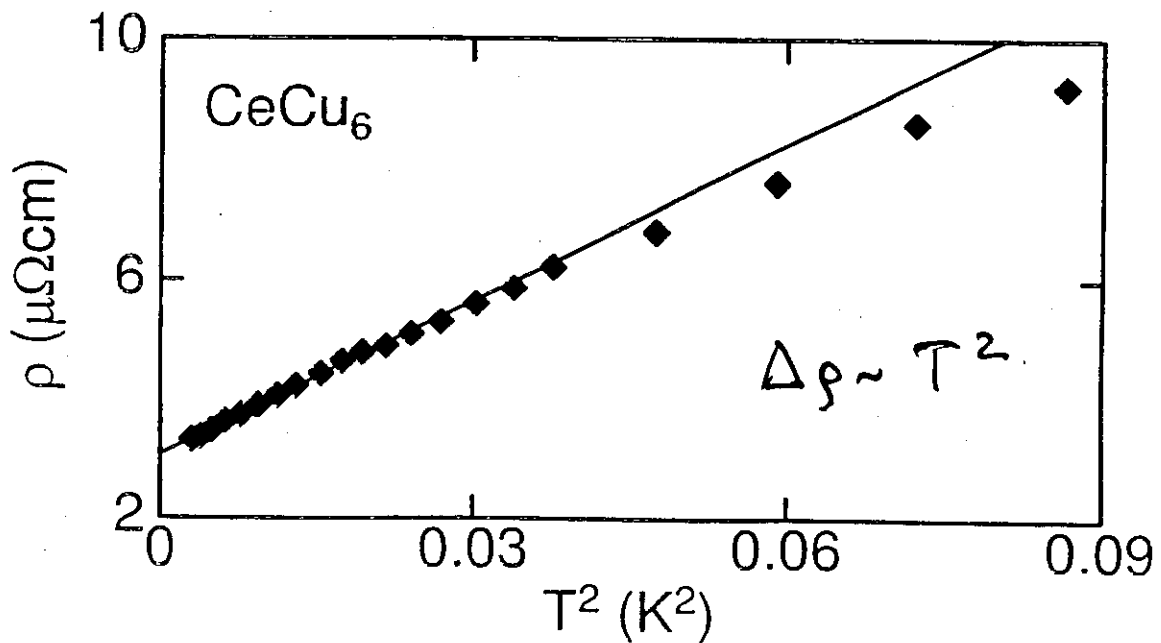
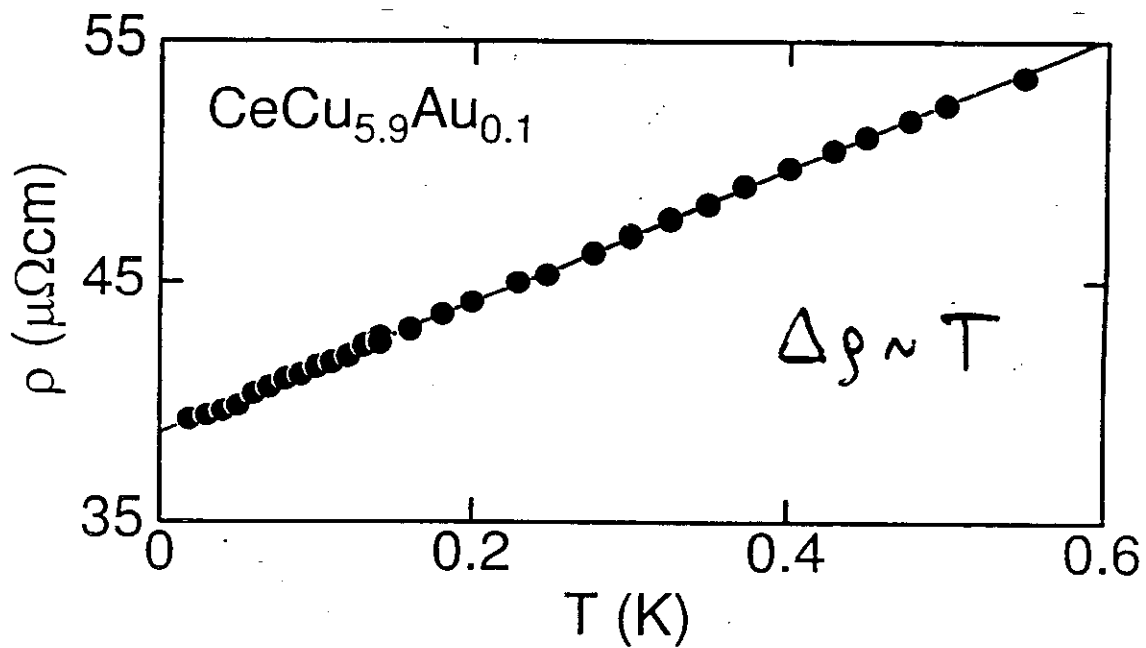
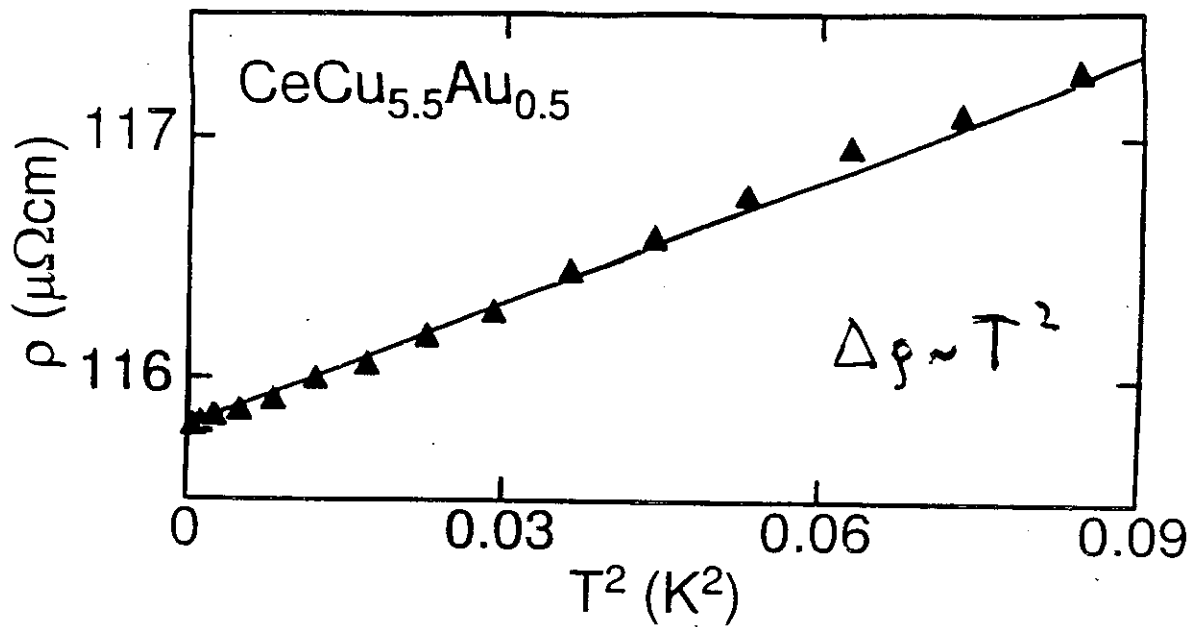
μSR for $x=0.1$: $\mu < 0.003 \mu_B$
(same as for CeCu_6)

A. Amato et al.





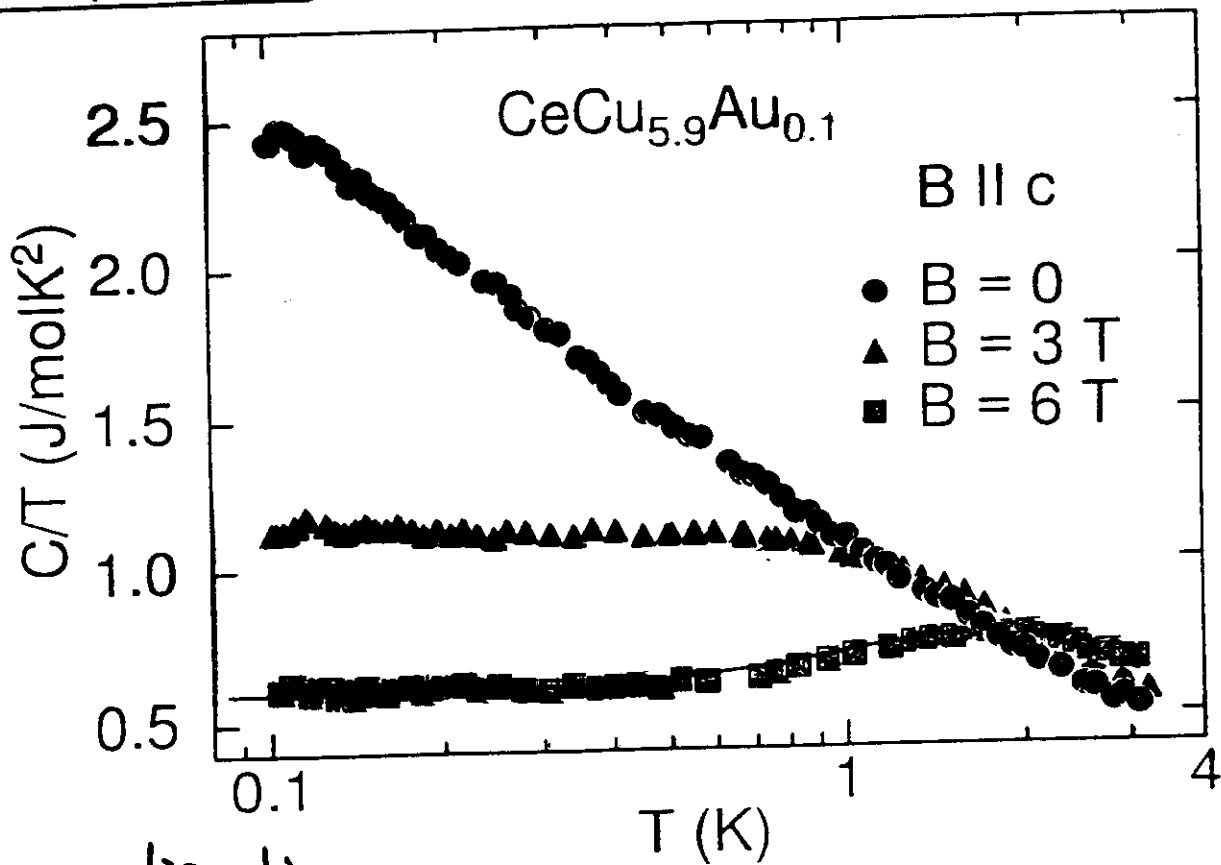




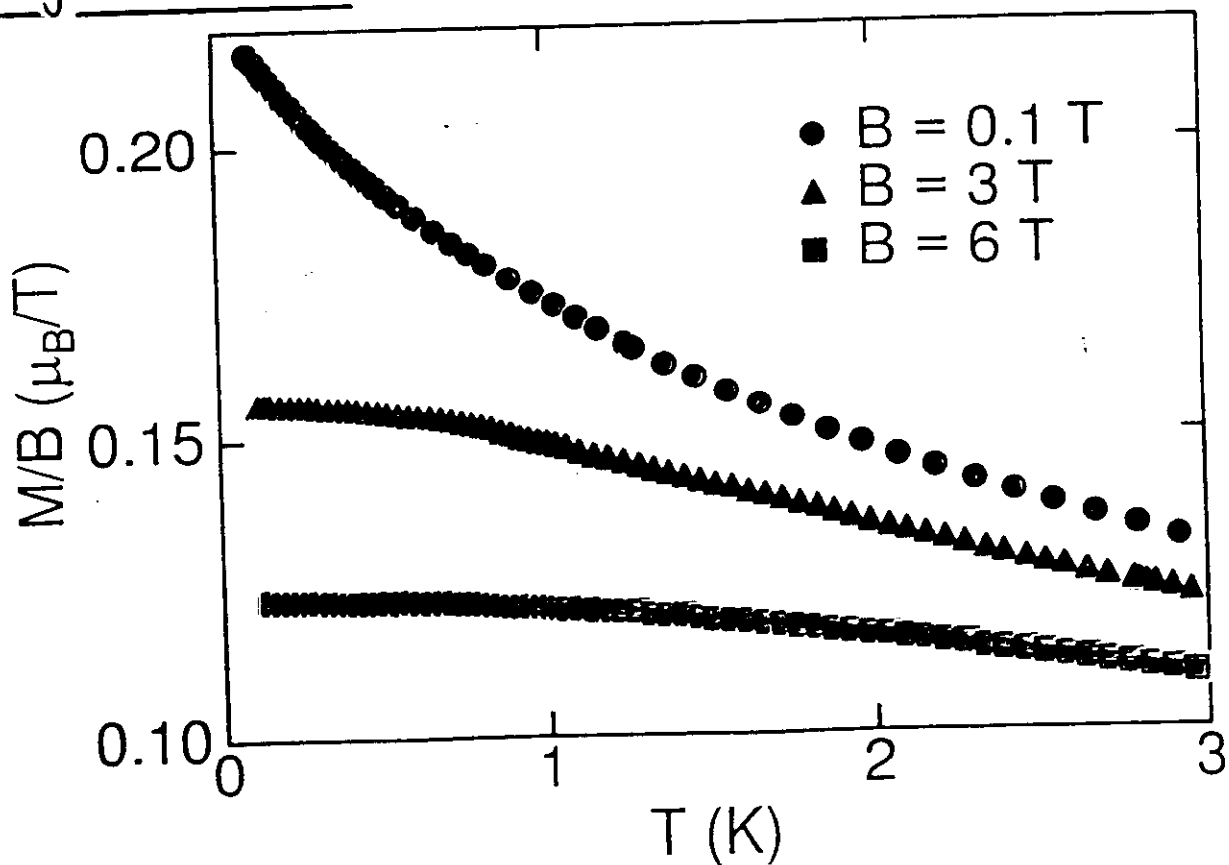
Recovery of Fermi-liquid in high magnetic fields

Specific heat

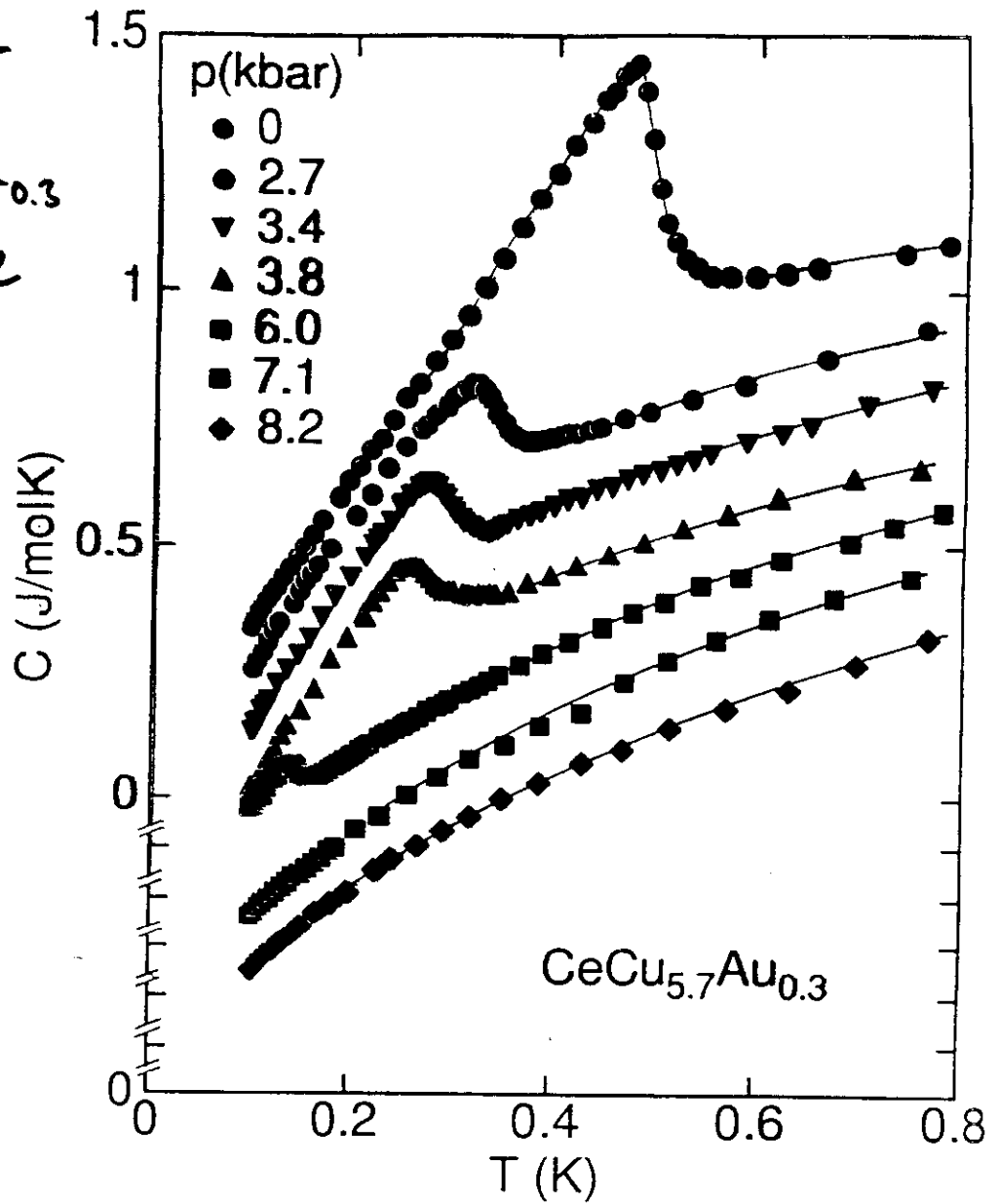
(hf contribution subtracted)



magnetization



Specific heat
of $\text{CeCu}_{5.7}\text{Au}_{0.3}$
under pressure



Néel temperature

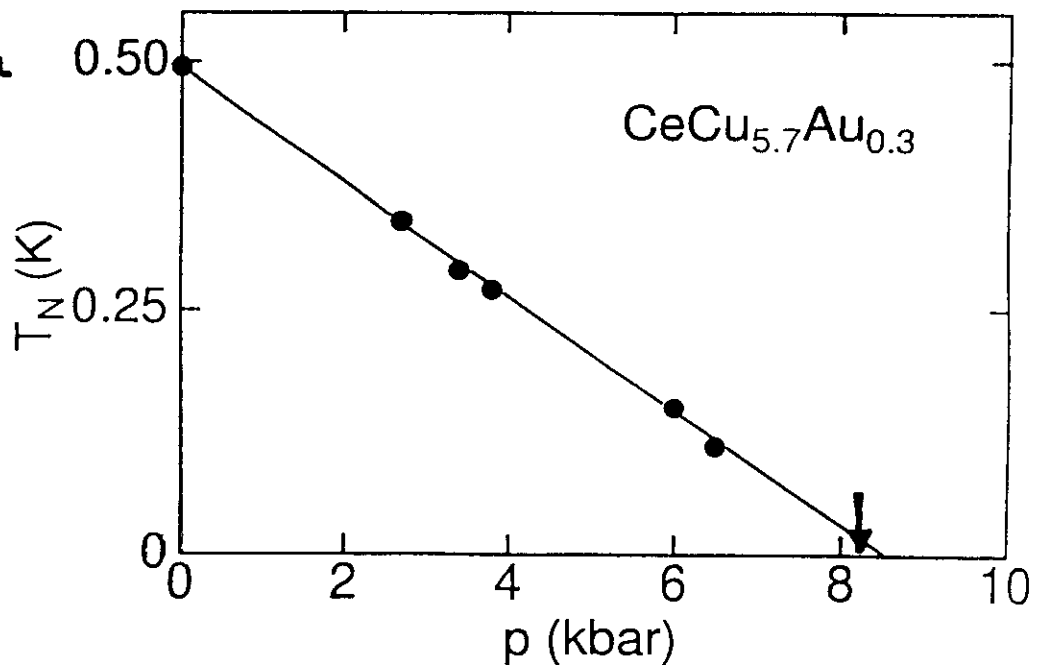
$$T_N \sim |p - p_c|^\mu$$

$$\mu = 1$$

cf.

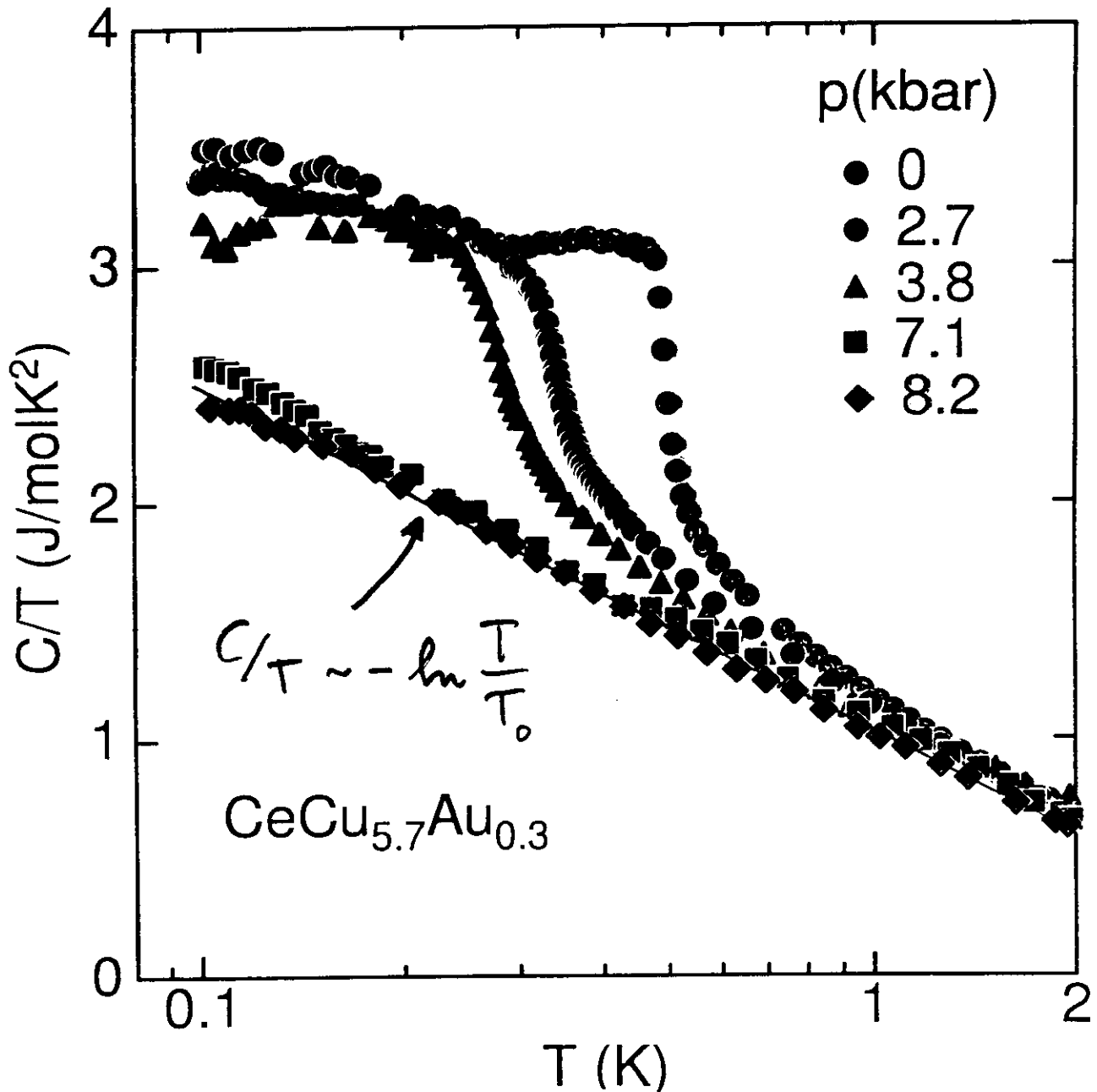
$$T_N \sim |x - x_c|^\mu$$

$$\mu = 1$$



Tuning of non-Fermi-liquid behavior

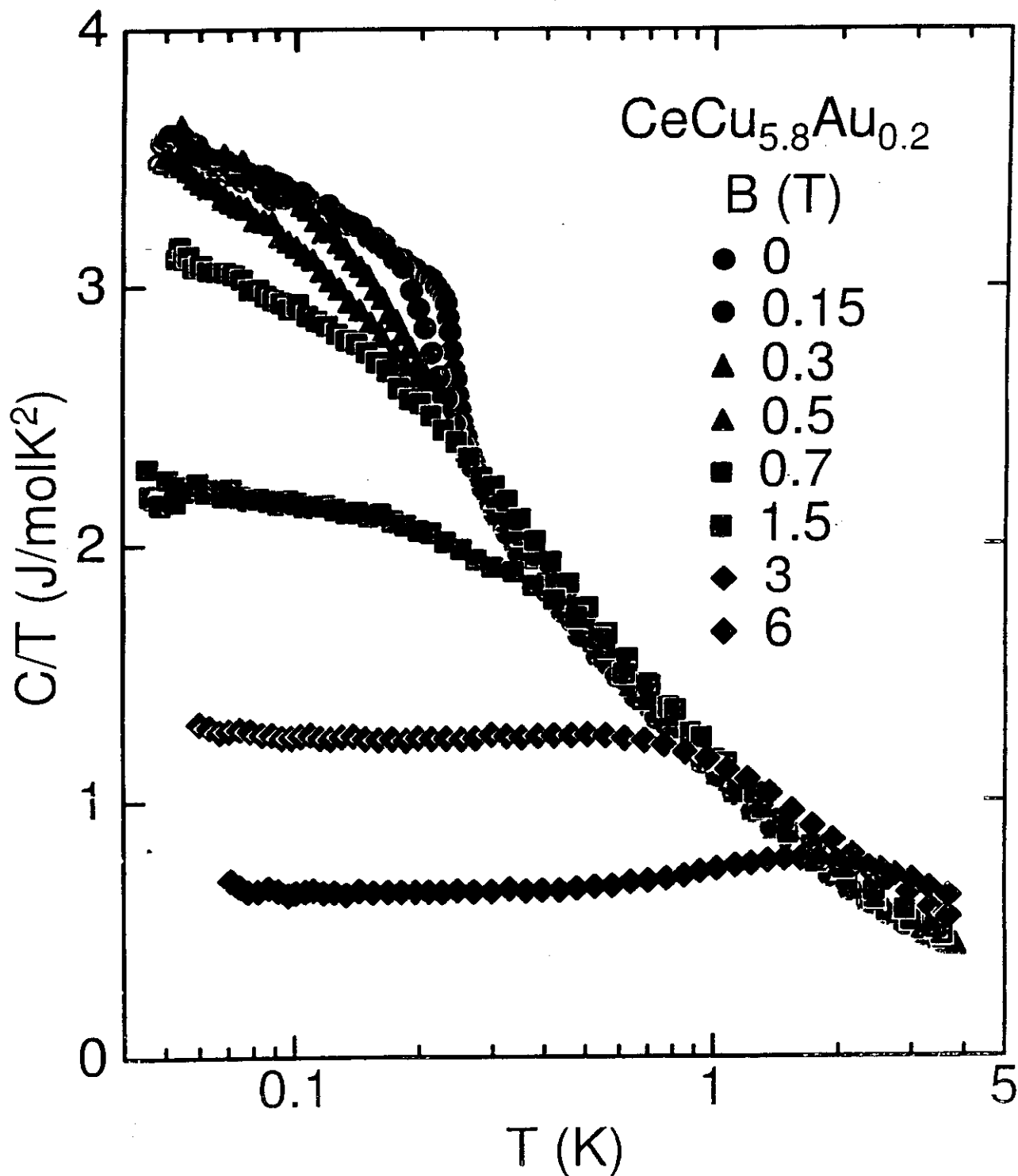
with hydrostatic pressure



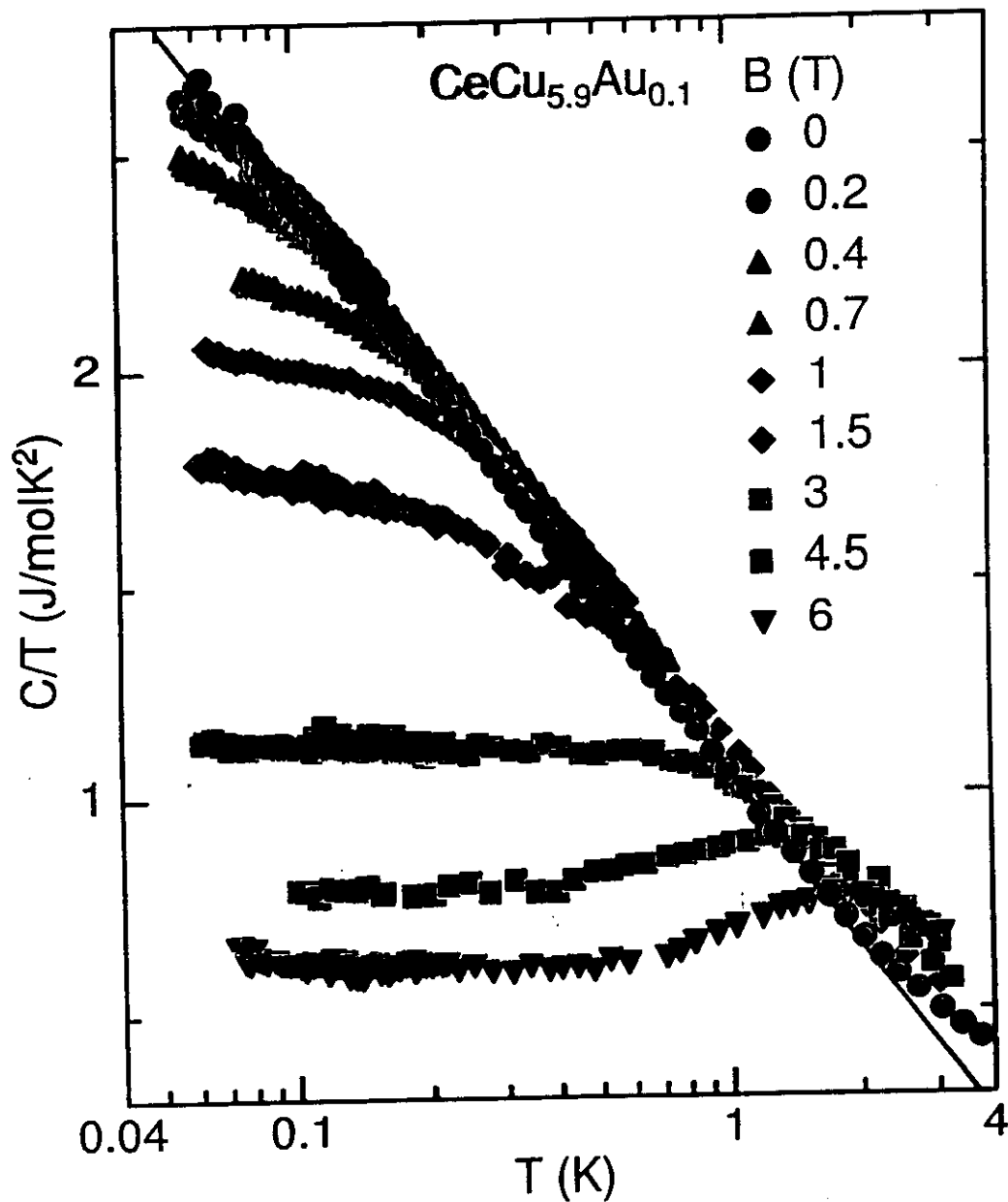
Specific heat of $\text{CeCu}_{5.8}\text{Au}_{0.2}$

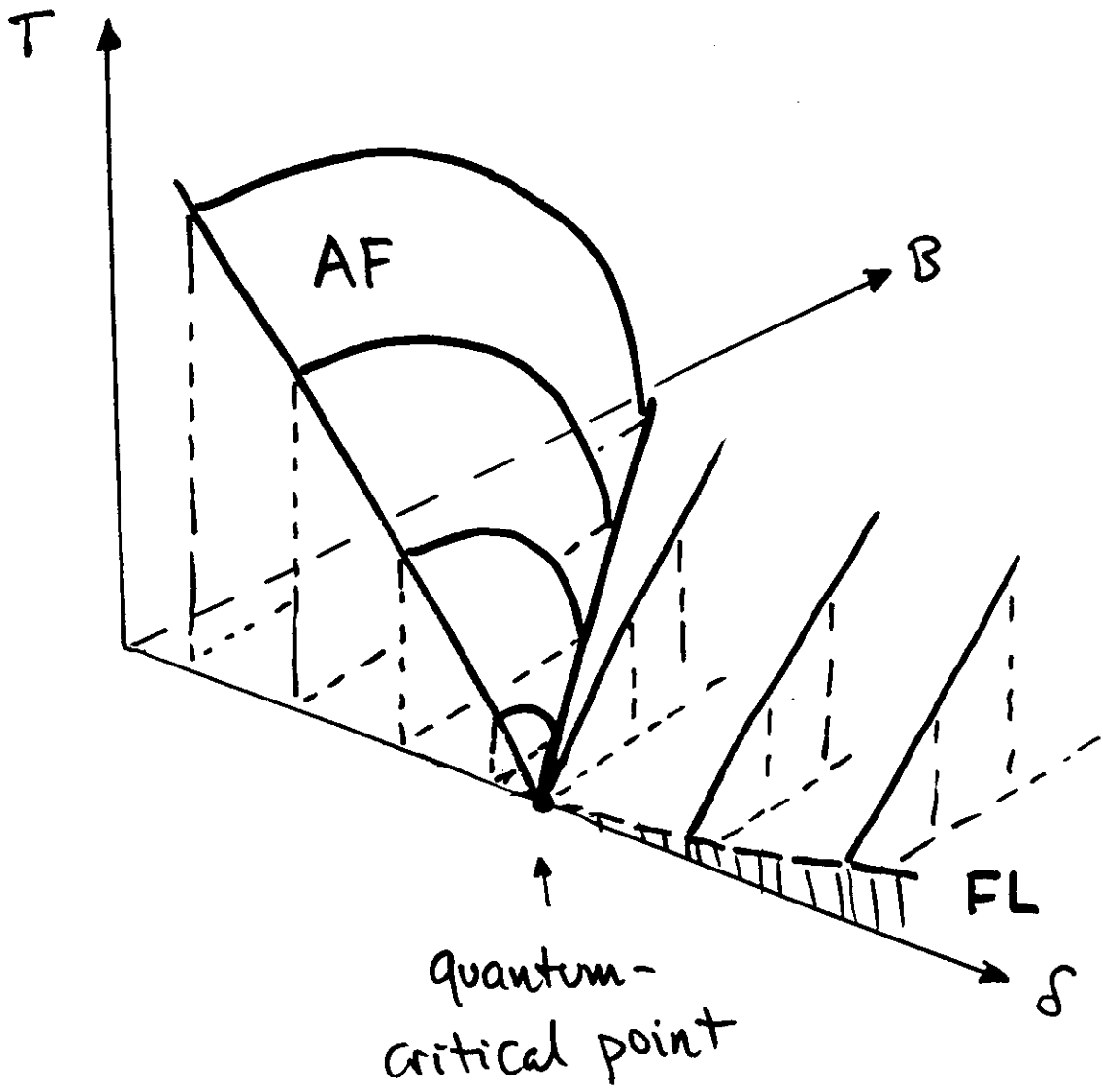
with $T_N = 0.25\text{K}$

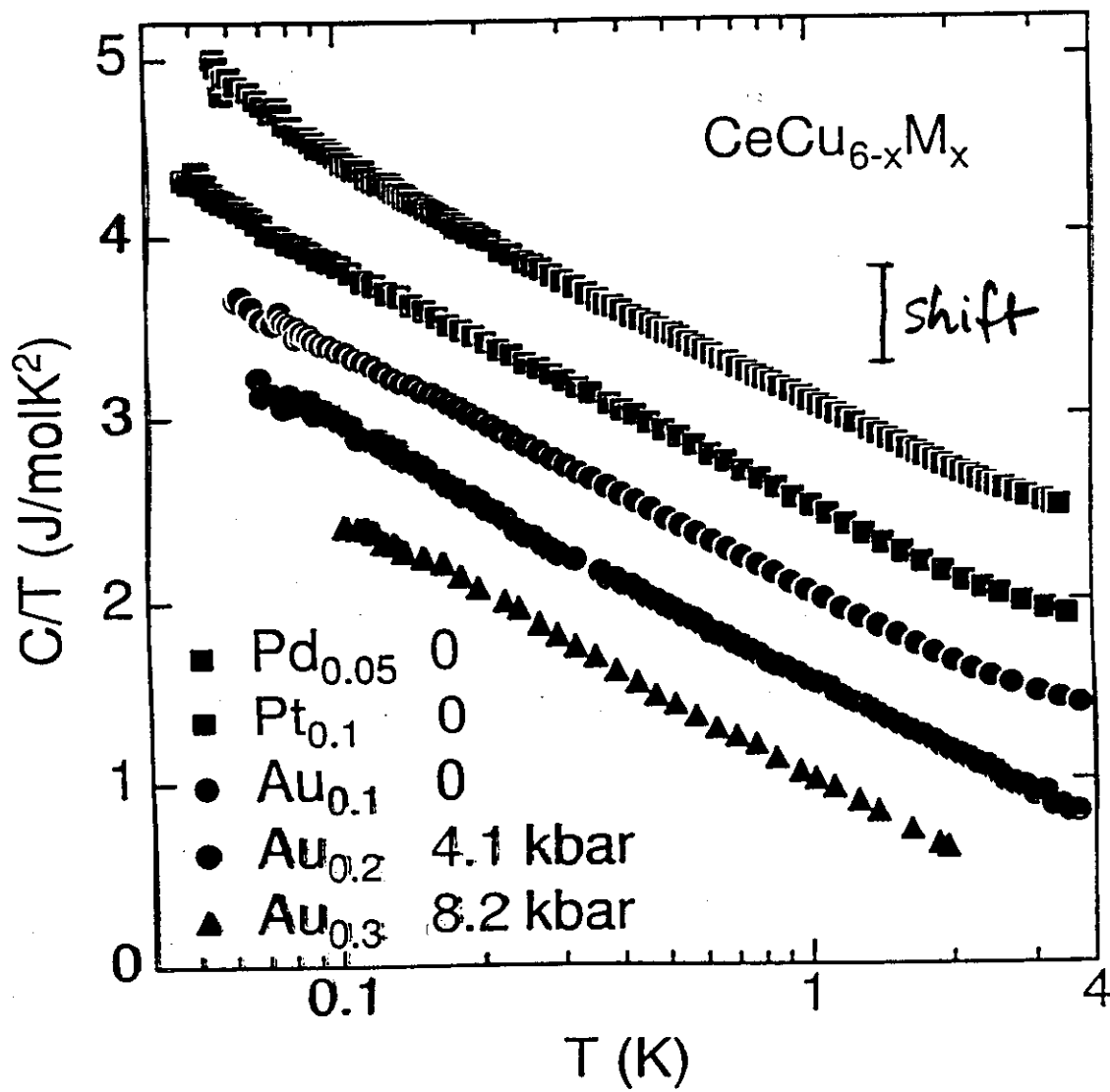
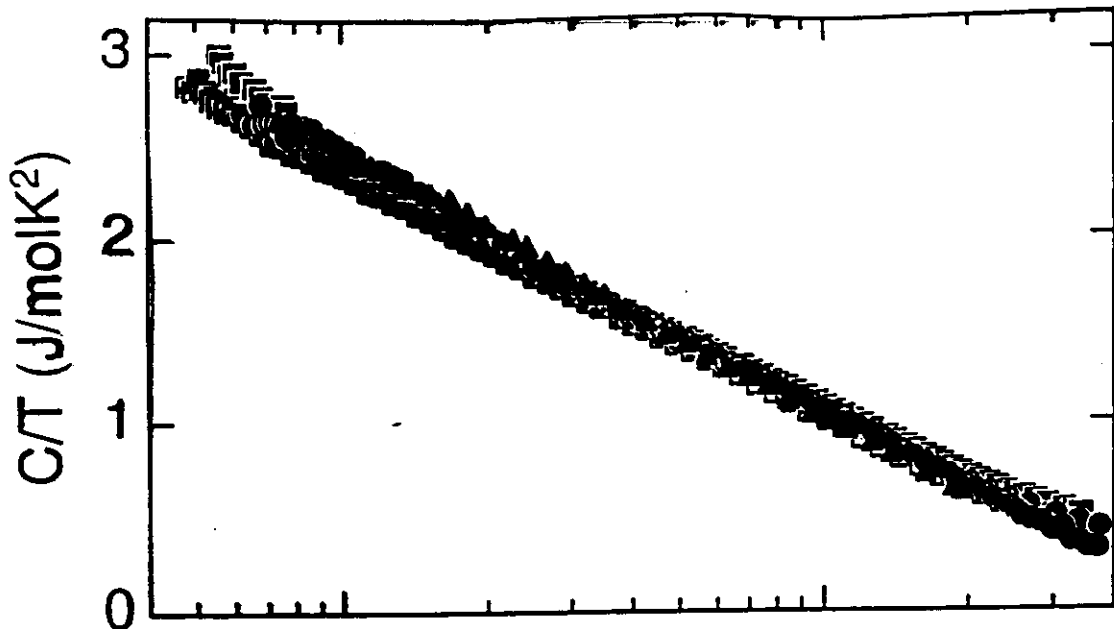
(hf contribution subtracted)



Crossover to Fermi-liquid behavior in magnetic fields







Universality in $CeCu_{6-x}M_x$: $C/T = a \ln(T/T_0)$
 ↑
 (1)

Other systems showing NFL behavior

Magnetic instability

$CeCu_{6-x}Au_x$, $-Pd_x$, $-Pt_x$	
$CePt(Si_{1-x}Ge_x)$	Steglich
$Ce_{0.1}La_{0.9}Cu_{2.2}Si_2$?	Andraka
$UCu_{5-x}Pd_x$?	Andraka, Stewart
$UCu_{4+x}Al_{8-x}$	Steglich
$U(Pt_{1-x}Pd_x)_3$	Kim, Stewart
$U_xY_{1-x}Pd_3$?	Andraka, Tsevik <u>1991</u>
$MnSi$, $ZrZn_2$	Lonzarich

Ce_7Ni_3 Takabatake
 $CeRh_2Si_2$
 $CePd_2Si_2$
Lonzarich, Thompson

Multichannel Kondo effect

$U_xY_{1-x}Pd_3, \dots$?	Maple <u>1991</u>	$Ce_{0.1}La_{0.9}Cu_{2.2}Si_2$?
$U_xTh_{1-x}Ru_2Si_2$	Amatsuka	Andraka
$U_xTh_{1-x}Be_{13}$	Aliev	

disordered metals (two-level tunneling states)

Distribution of Kondo temperatures

$UCu_{5-x}Pd_x$?	Mc Laughlin
Si:P (above metal-insulator transition)	

Other mechanisms leading to NFL behavior: single-ion effects

- Multichannel Kondo effect $\mathcal{H} = - \sum_{\alpha} J S s_{\alpha}$
conduction-electron channels $\alpha = 1, \dots, N$

if $N > 2S$ "overscreening"

$N=2, S=\frac{1}{2}$: \uparrow $(\downarrow \uparrow \downarrow)$ \uparrow

local non-Fermi liquid

$$\frac{C}{T} \sim \chi \sim - \ln T/T_K$$

generalization: internal degree of freedom
of impurity \rightarrow "pseudospin"

e.g. orbital coupling between conduction electrons
and f electrons with quadrupole moment:

pseudospin : electric quadrupole moment,

\Rightarrow two conduction-electron channels

quadrupolar Kondo effect D. Cox

- Distribution of T_K in disordered metals

more precisely: incoherent superposition of
local Fermi liquids

Distribution of Kondo temperatures

single impurity: $T_K \sim \exp\left(-\frac{1}{N(E_F)J}\right)$

usual metal hosts: $N(E_F), J$ well-defined

Metal-insulator transition with disorder

e.g. heavily doped semiconductors:

Si:P : $N_c \approx 3.5 \cdot 10^{18} \text{ cm}^{-3}$

formation of intrinsic magnetiz moments
already on the metallic side due to disorder

distribution of J and $N(E_F) \Rightarrow P(T_K)$

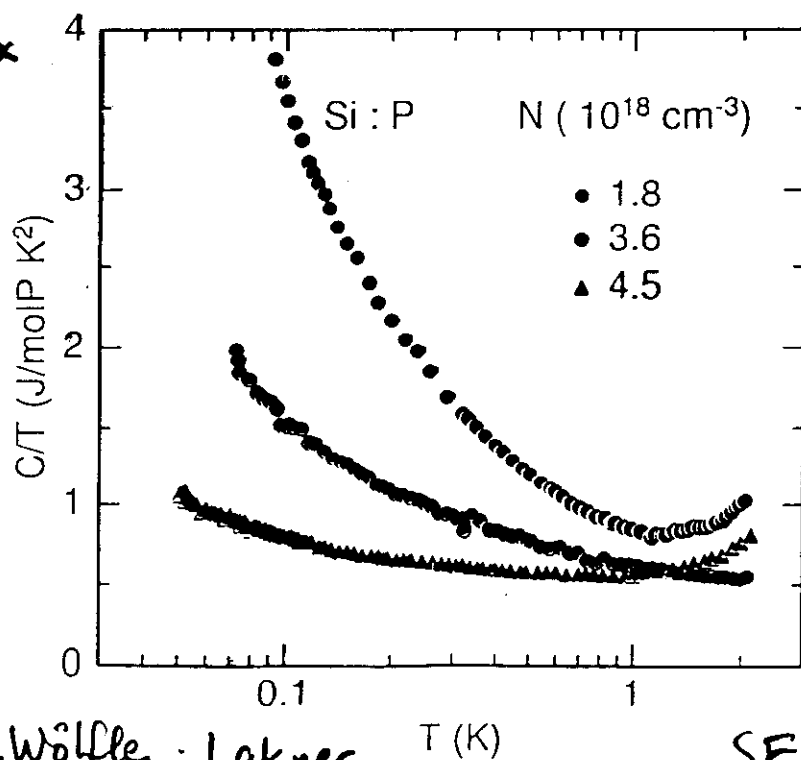
$$P(T_K) \sim T_K^{-\alpha_K}$$

↓

$$\frac{C}{T} \sim T^{-\alpha_K}$$

exptl. $\alpha_K \approx 0.8$

theor. $\alpha_K \approx 0.9$



Langenfeld, Wölfle ; Lakner

T (K)

SFB 195

Conclusions

- NFL behavior can have different microscopic origin: single-ion vs. $T=0$ phase transition
- $\text{CeCu}_{6-x}\text{Au}_x$: concentration and pressure tuning of magnetic instability

Open problems:

- $\Delta\chi \sim \sqrt{T}$, $C/T \sim -\ln(T/T_0)$, $\Delta\rho \sim T$
more generally: scaling properties?
- role of disorder?
- what type(s) of magnetic instability lead(s) to NFL behavior?
- microscopic model? Spin fluctuations as $T_N \rightarrow 0$
(FM vs AF)
- NFL: crossover or critical behavior for $T \rightarrow 0$
- one or two components: 3d vs. 4f heavy fermions