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SMR.998d - 27

Research Workshop on Condensed Matter Physics
30 June - 22 August 1997
MINIWORKSHOP ON
QUANTUM WELLS, DOTS, WIRES
AND SELF-ORGANIZING NANOSTRUCTURES
11 - 22 AUGUST 1997

"Electronic Structure of Few-Particle Quantum Dots"

PART I

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I

These are preliminary lecture notes, intended only for distribution to participants.

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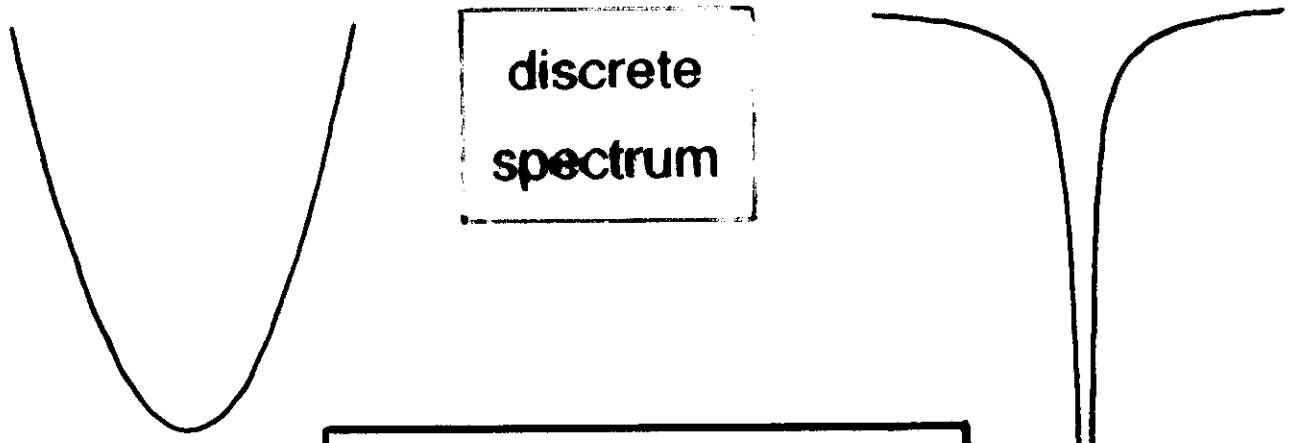
Outline

- Few Electron Quantum Dots:
Artificial Atoms
 - QD-Hydrogen: a 2D harmonic oscillator
 - QD-Helium: an illustrative example
- Correlations in Few-Particle Systems
- Chaos induced by Coulomb interaction?

Artificial Atoms

artificial: $\sim r^2$

natural: $\sim 1/r$

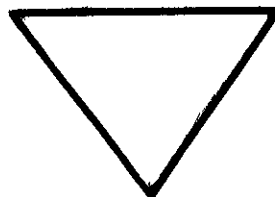


number of electrons



"quantum dot element"

Coulomb
interaction



kinetic
energy

confinement energy

Hamiltonian

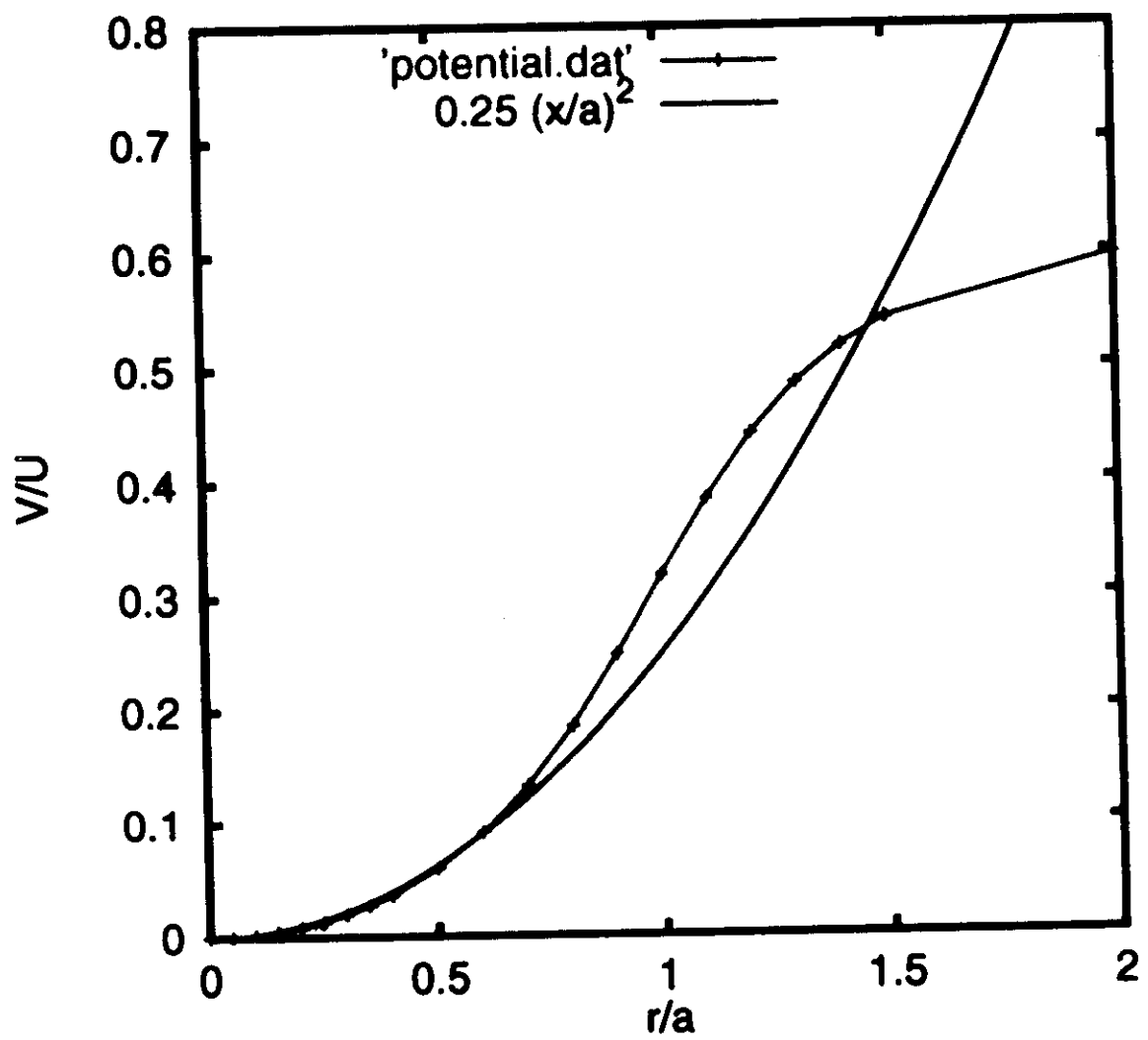
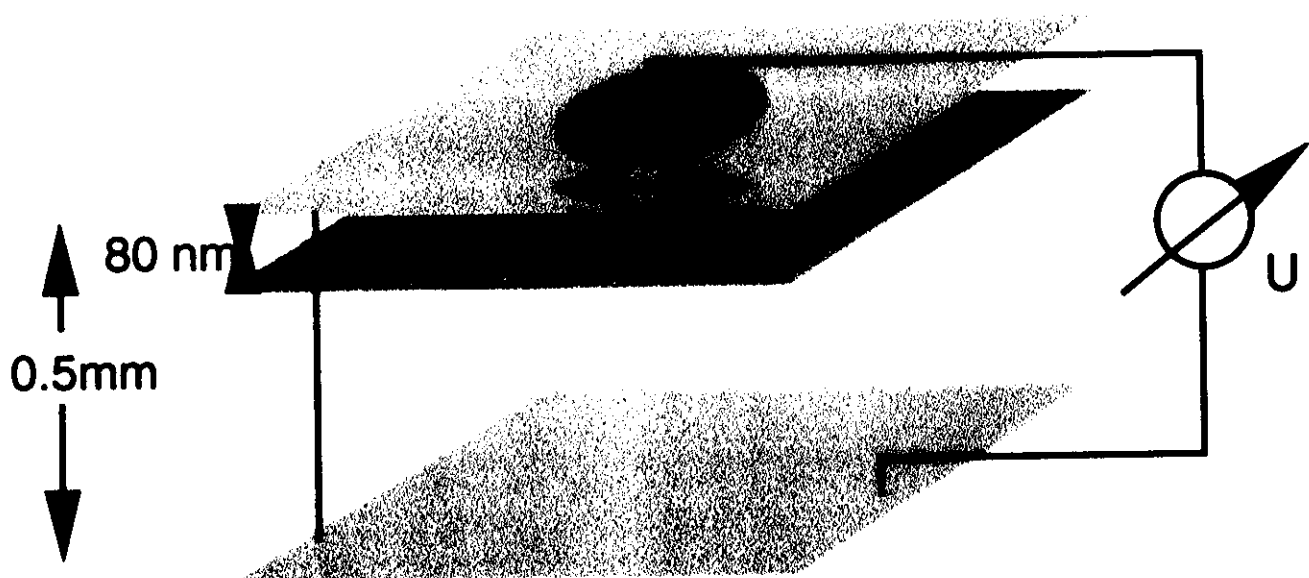
2D-electron		
in perpendicular magnetic field	Zee-man-energy	Coulomb-interaction

$$\mathcal{H} = \sum_{i=1}^{N_s} \left[\frac{1}{2m^*} \left(\vec{p}_i + \frac{e}{c} \vec{A}(\vec{r}_i) \right)^2 + \frac{g^* \mu_B}{\hbar} \vec{B} \cdot \vec{S}_i + V_{conf}(\vec{r}_i) \right] + \sum_{i < j}^{N_s} \frac{e^2}{\epsilon |\vec{r}_i - \vec{r}_j|}$$

confining potential

$$V_{conf}(r) = \frac{1}{2} m^* \Omega_0^2 r^2$$

Elektrostatistisches Potential



1-Particle Quantum Dot: QD-Hydrogen

$$H_0 = \frac{1}{2m^*} \left(\vec{p} + \frac{e}{c} \vec{A} \right)^2 + \frac{1}{2} m^* \Omega_0^2 r^2$$

symmetric gauge

$$= \hbar \Omega_+ (a_+^\dagger a_+ + \frac{1}{2}) + \hbar \Omega_- (a_-^\dagger a_- + \frac{1}{2})$$

effective confinement frequency

$$\Omega_{\pm} = \sqrt{\Omega_0^2 + \omega_c^2/4} \quad \omega_c = \frac{eB}{m^*c}$$

angular momentum

$$L_z = \hbar (a_+^\dagger a_+ - a_-^\dagger a_-)$$

eigenstates $|N_+ N_- \rangle$:

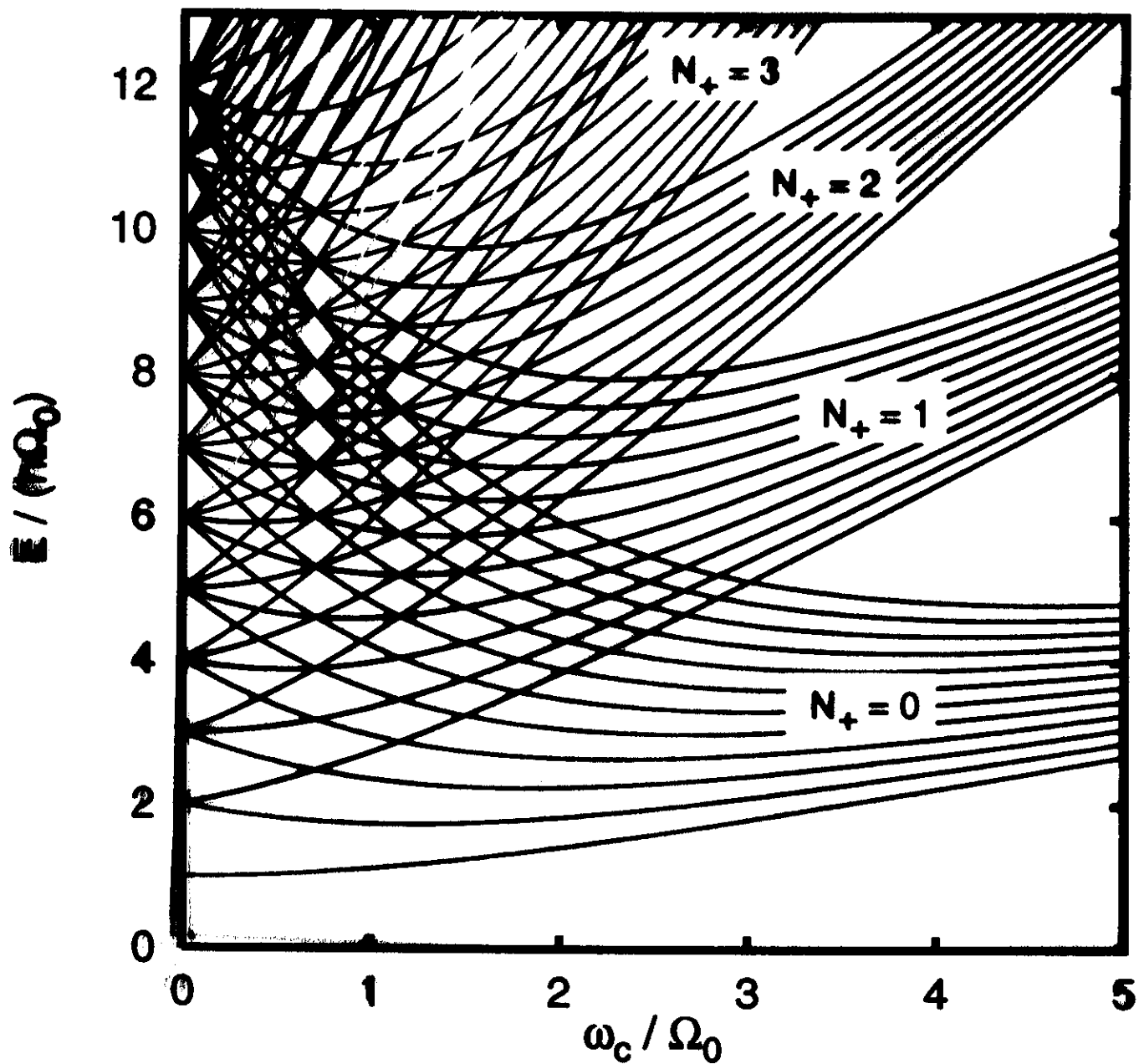
right(+) and left(-) circular polarized modes

s-states: $N_+ - N_- = 0$

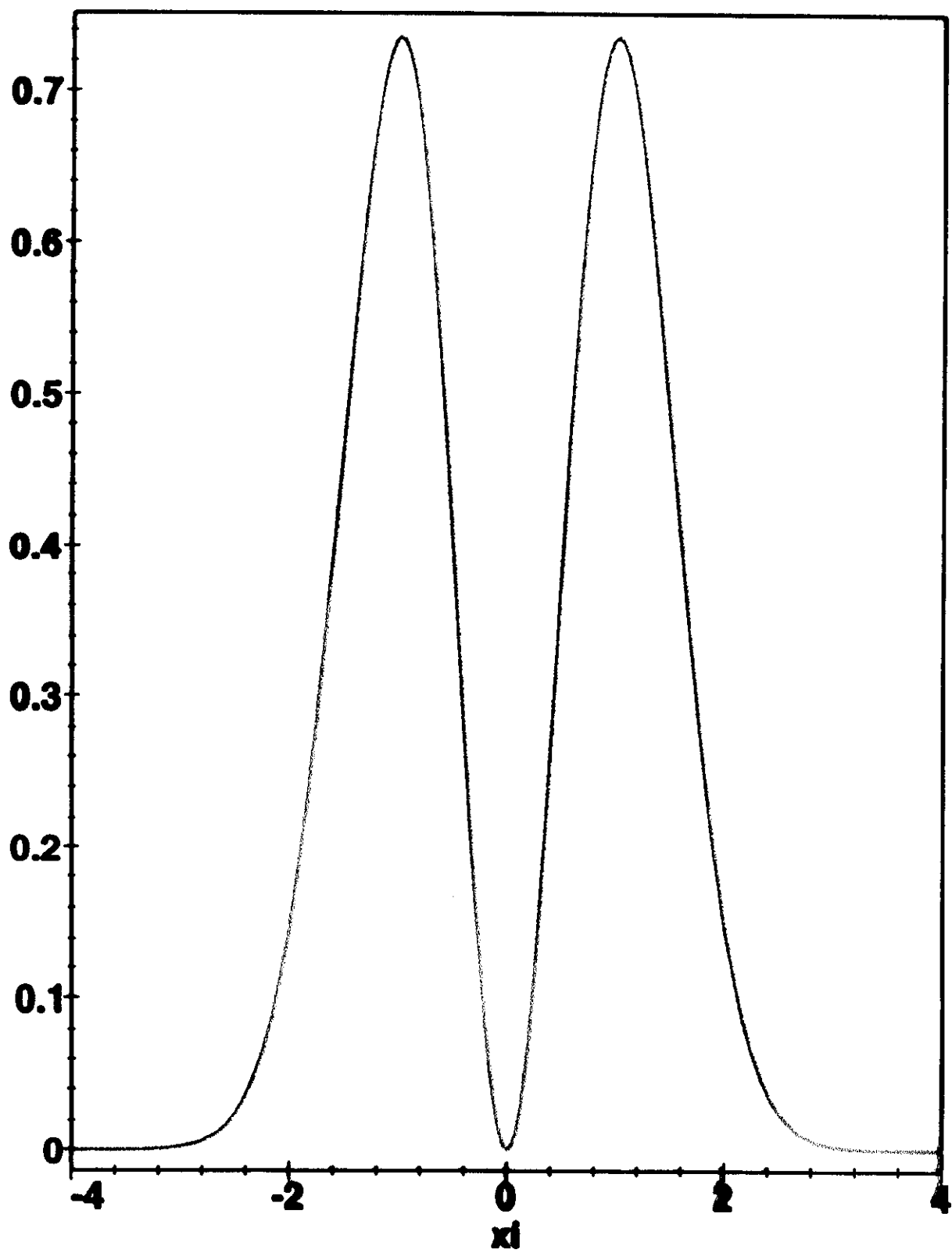
p-states: $N_+ - N_- = \pm 1$

1-Teilchen-Spektrum

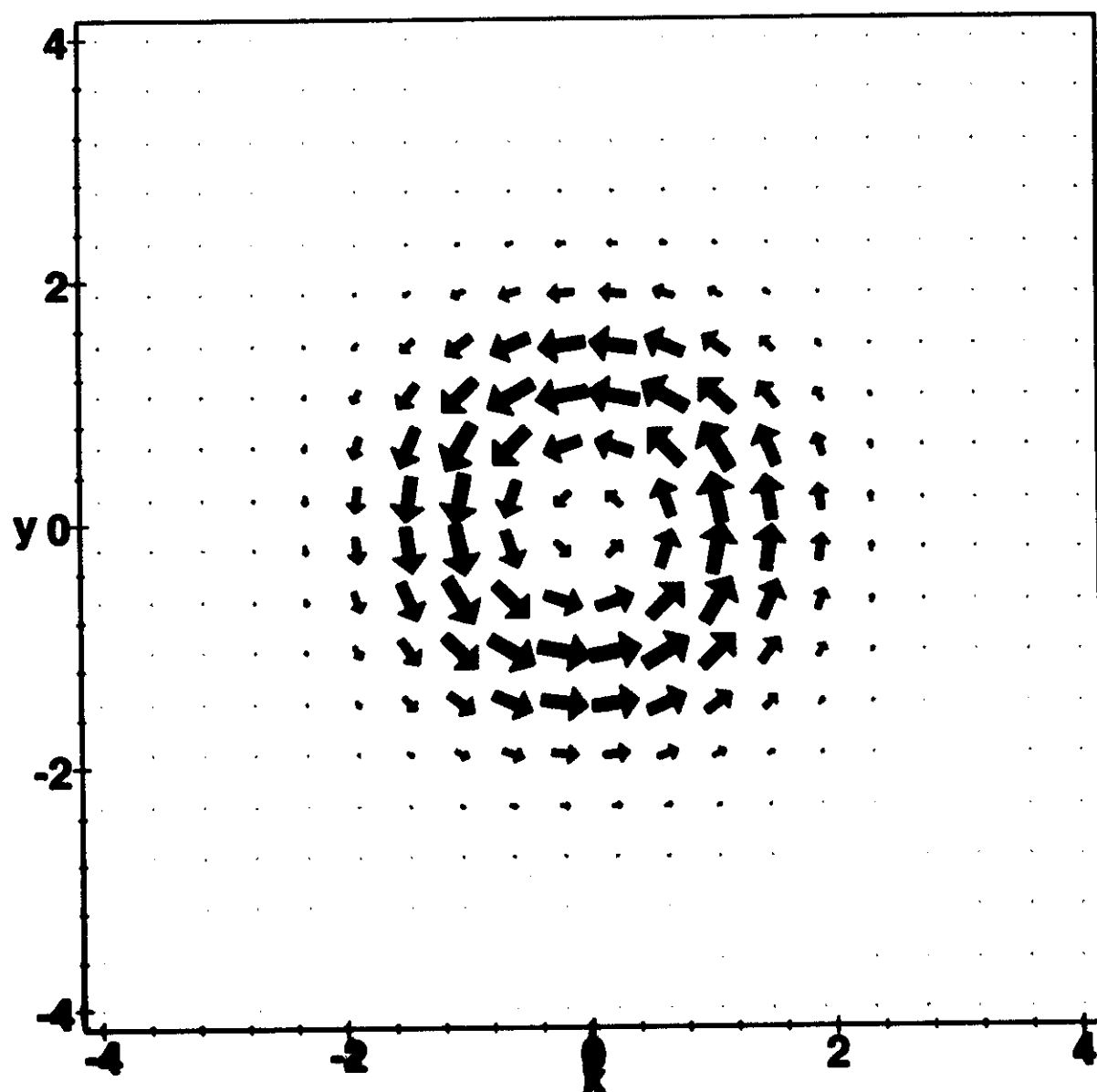
Quantendot-Wasserstoff



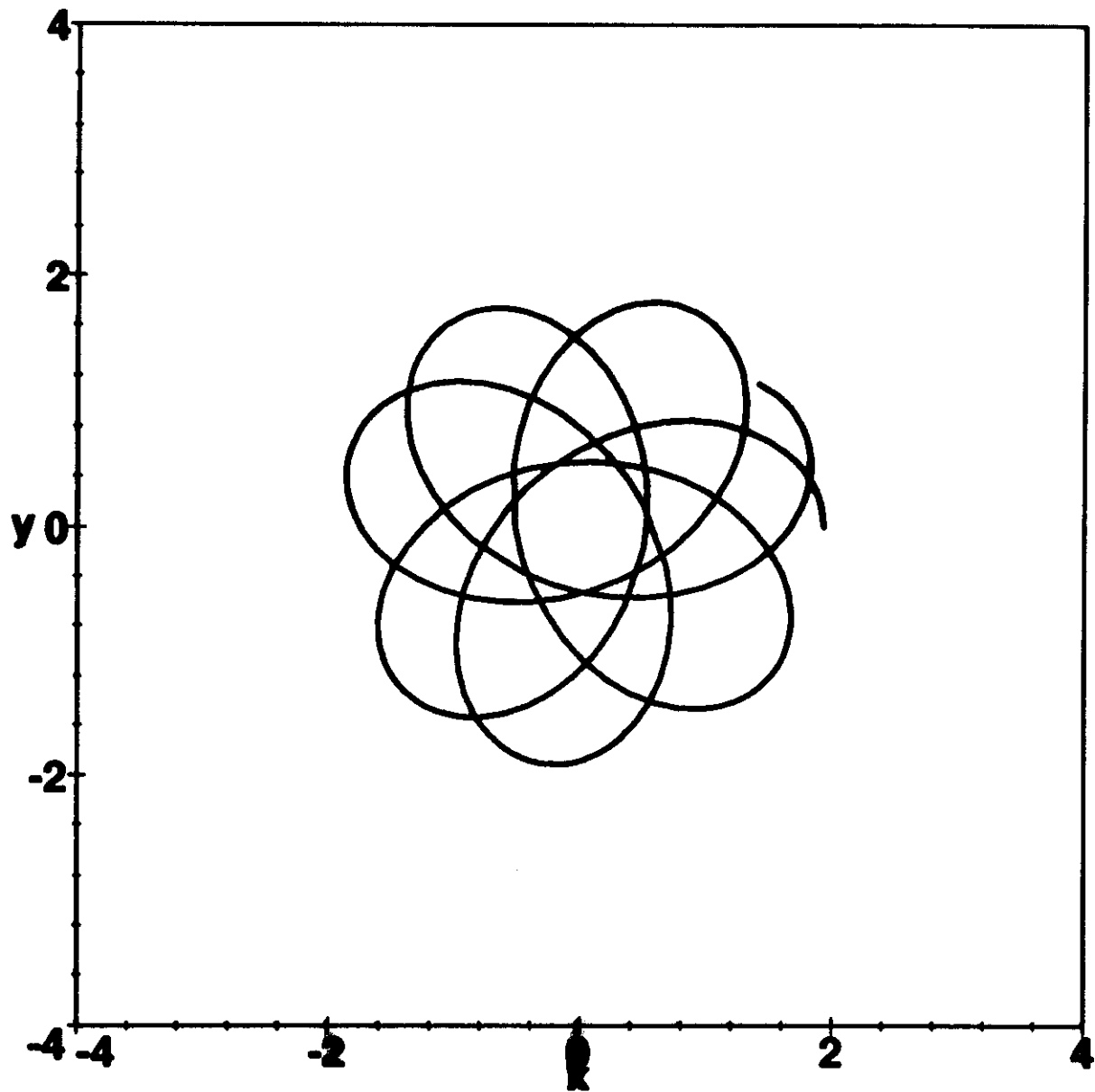
density, $N_p := 0$; $N_m := 1$;



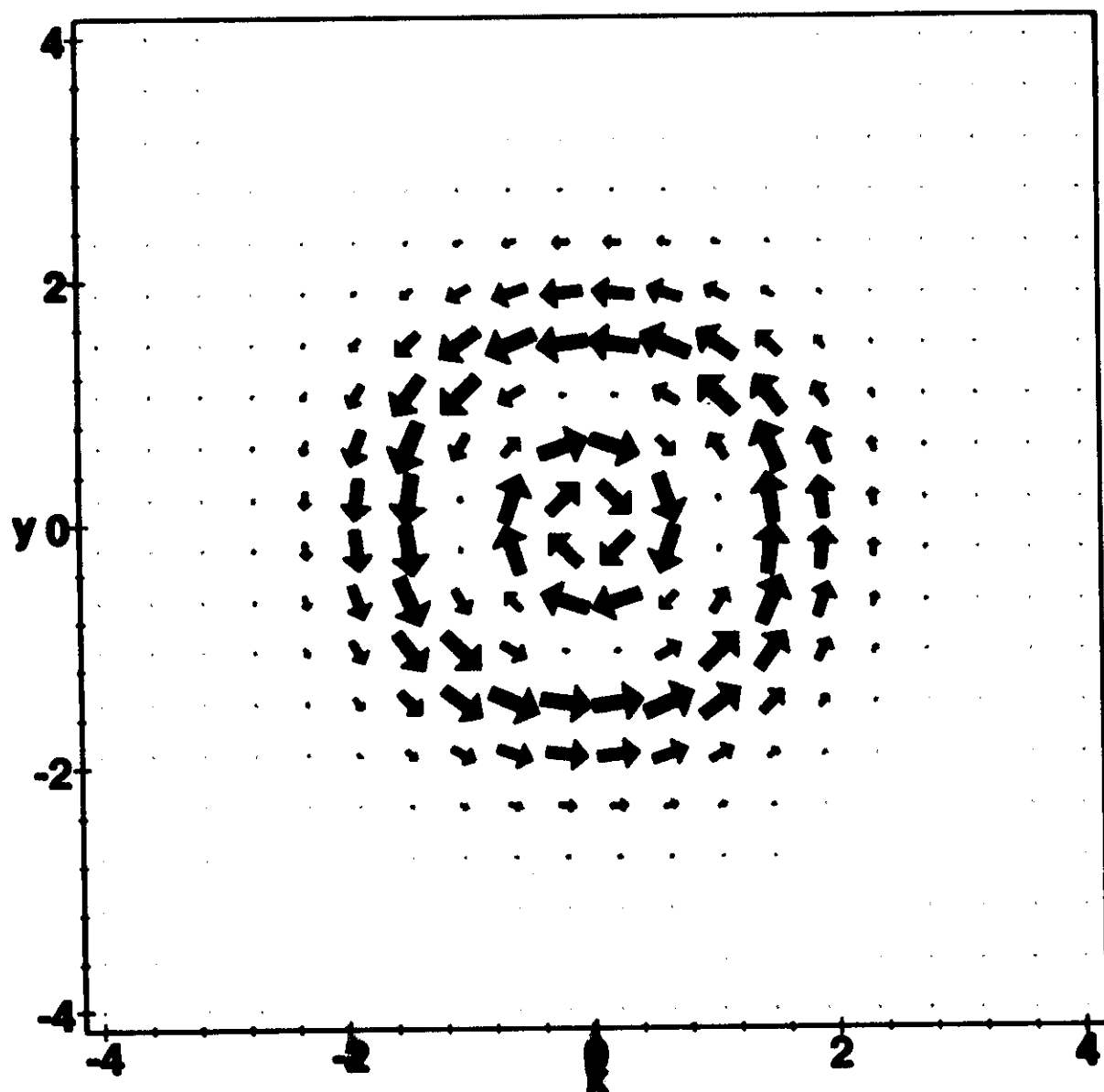
current density, $N_p := 1$; $N_m := 0$; $ratw := 1$;



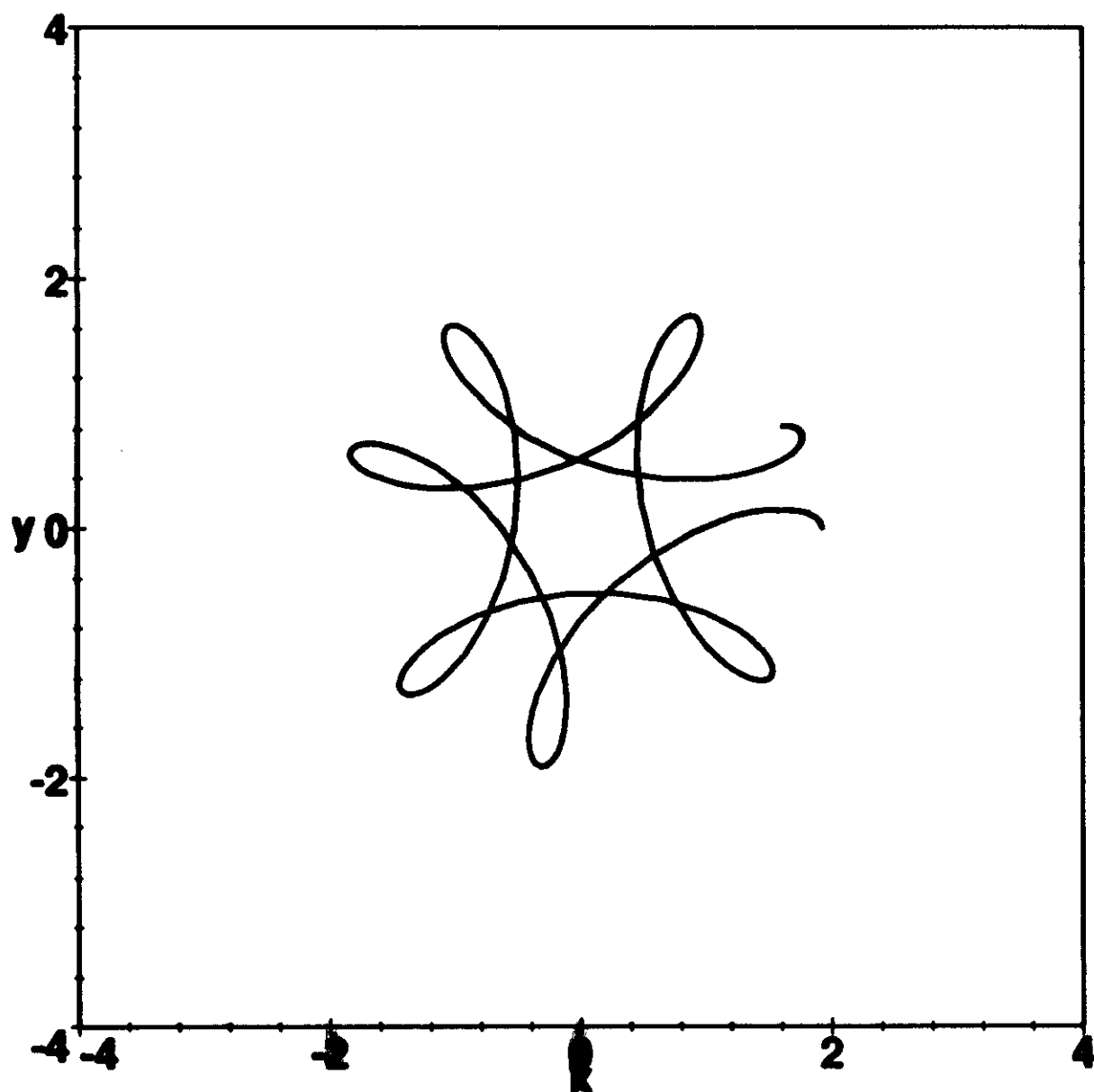
orbit, Np:= 1; Nm:= 0; ratw:= 1;



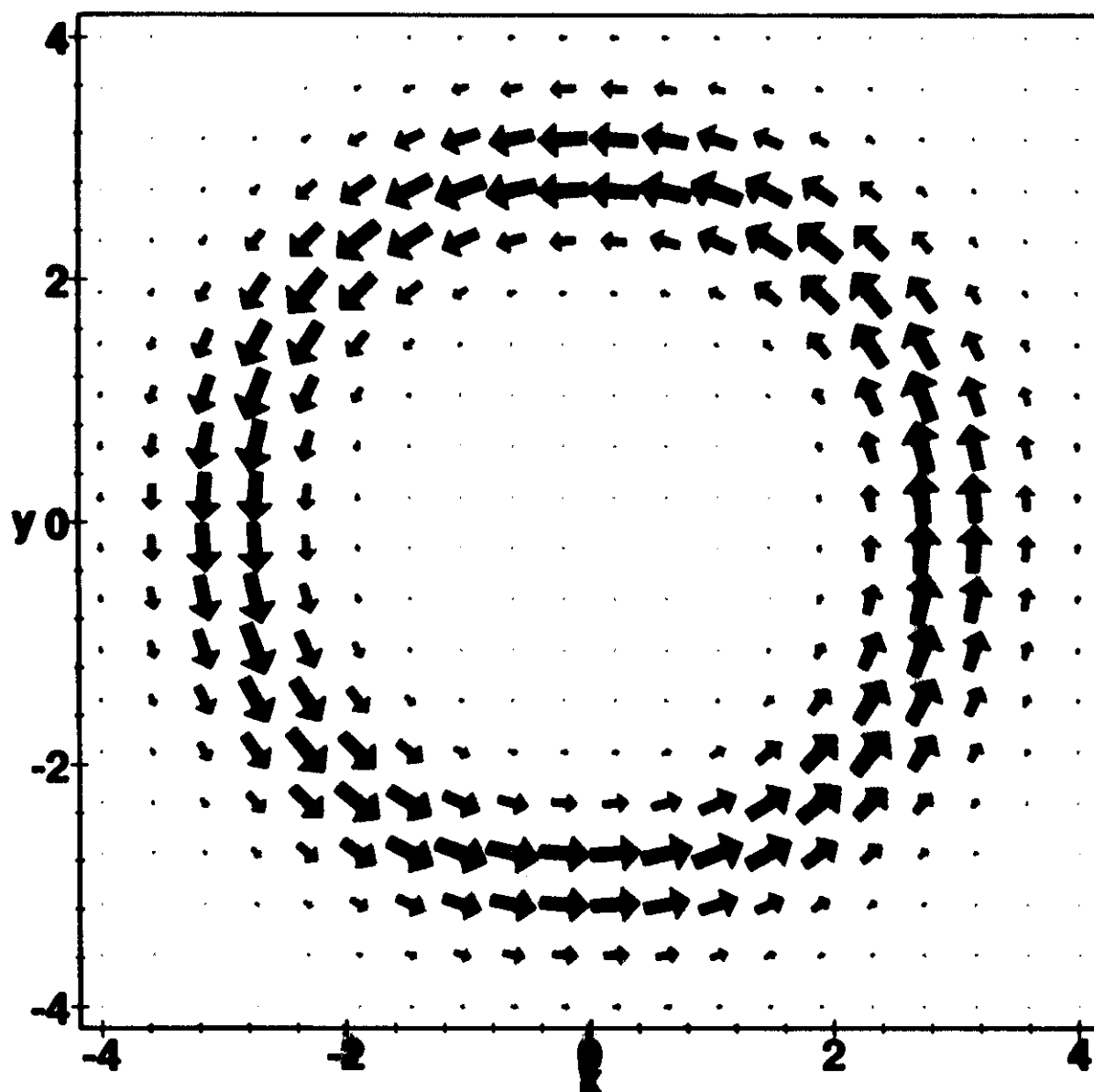
current density, $N_p := 0$; $N_m := 1$; $ratw := 1$;



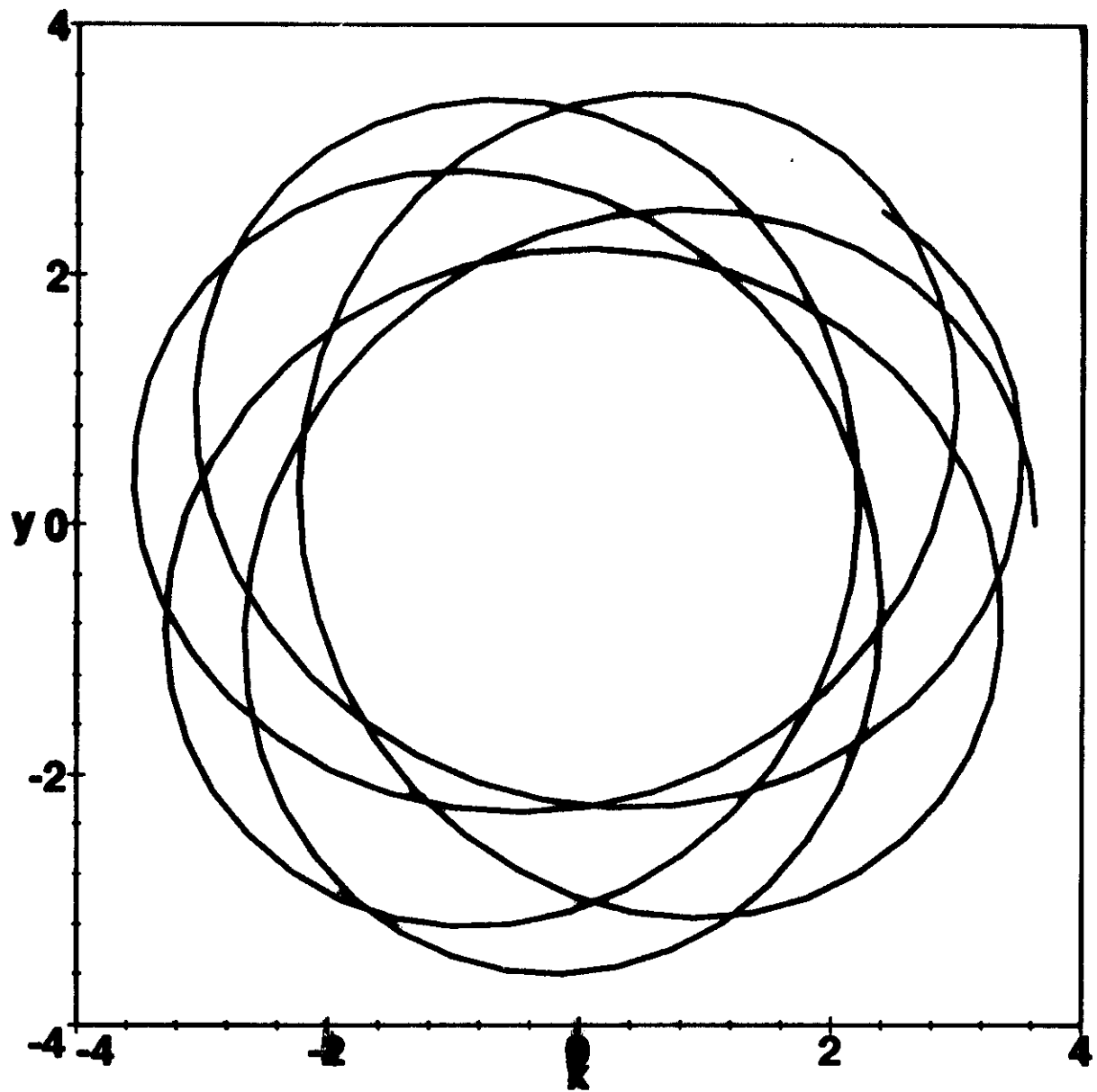
orbit, Np:= 0; Nm:= 1; ratw:= 1;



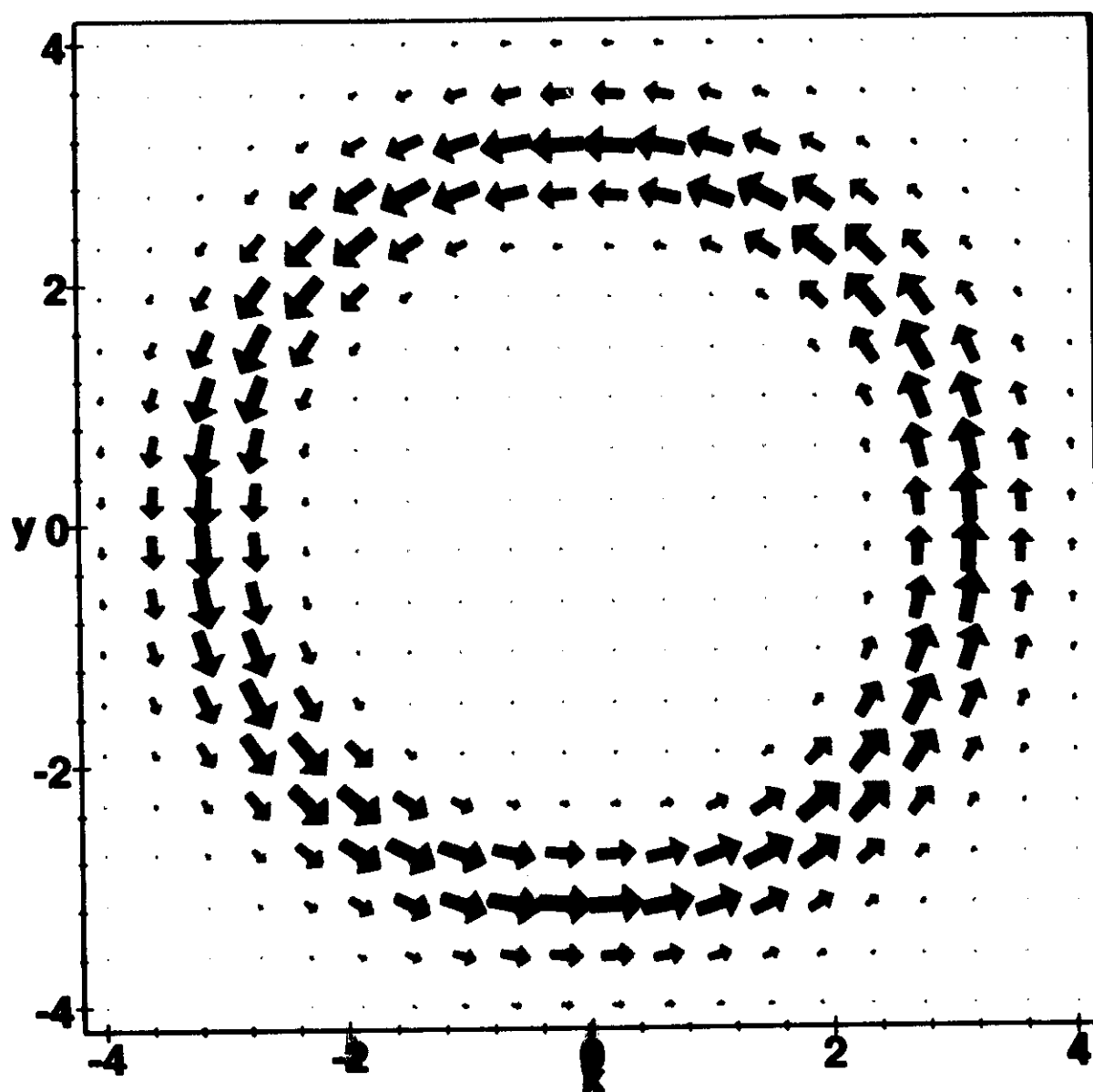
current density, $N_p := 8$; $N_m := 0$; $ratw := 1$;



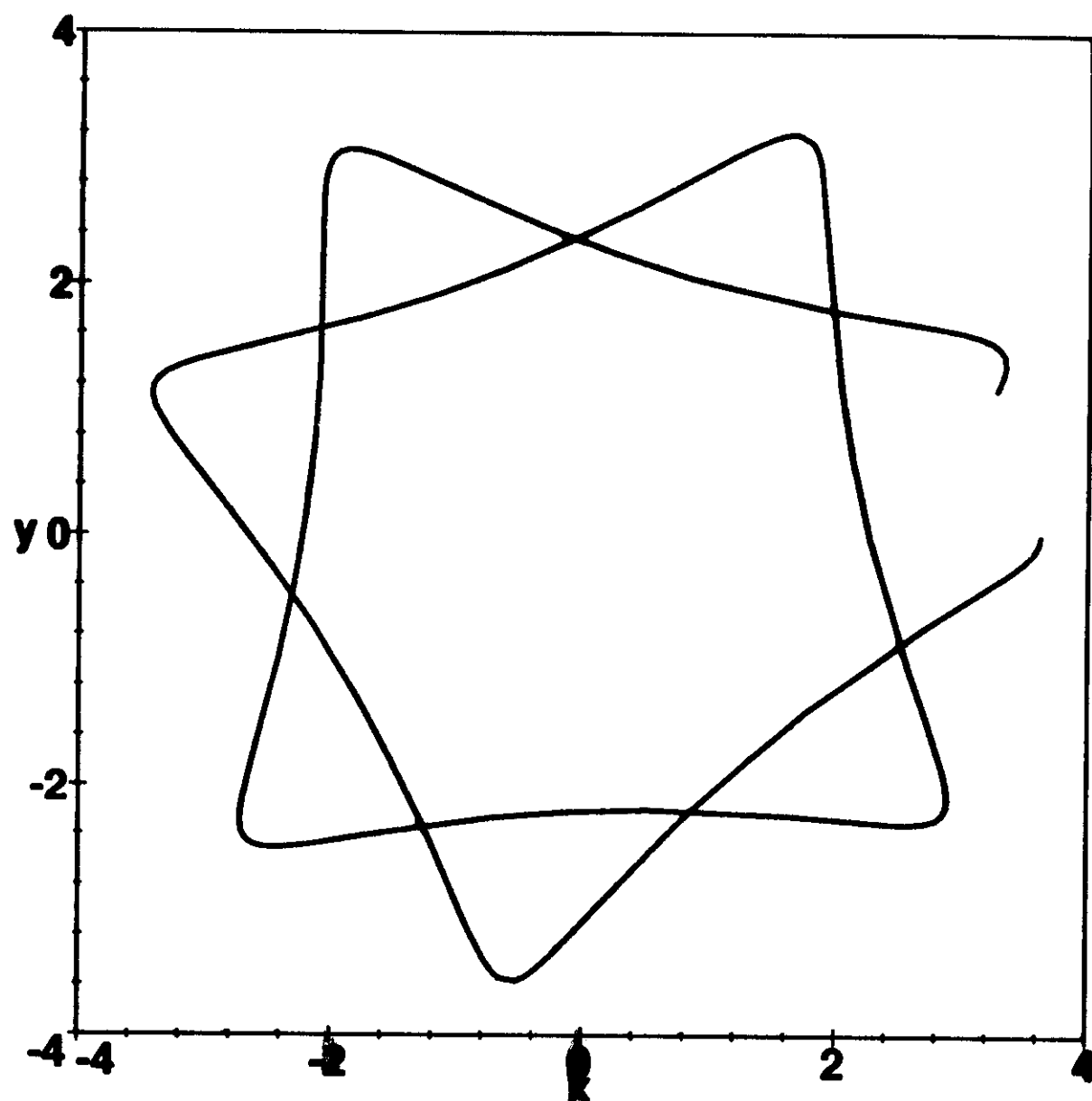
orbit, Np:= 8; Nm:= 0; ratw:= 1;



current density, $N_p := 0$; $N_m := 8$; $ratw := 1$;



orbit, Np:= 0; Nm:= 8; ratw:= 1;



QD-Helium

$$\mathcal{H} = H_0(1) + H_0(2) + \frac{e^2}{\epsilon |\vec{r}_1 - \vec{r}_2|}$$

couples motion of both particles

separation into center-of-mass and relative motion

$$\begin{aligned}\vec{R} &= \frac{1}{2}(\vec{r}_1 + \vec{r}_2) & \vec{r} &= \vec{r}_1 - \vec{r}_2 \\ \vec{P} &= \vec{p}_1 + \vec{p}_2 & \vec{p} &= \frac{1}{2}(\vec{p}_1 - \vec{p}_2)\end{aligned}$$

$$\mathcal{H} = H_0(\vec{R}, \vec{P}) + H_0(\vec{r}, \vec{p}) + \frac{e^2}{\epsilon r}$$

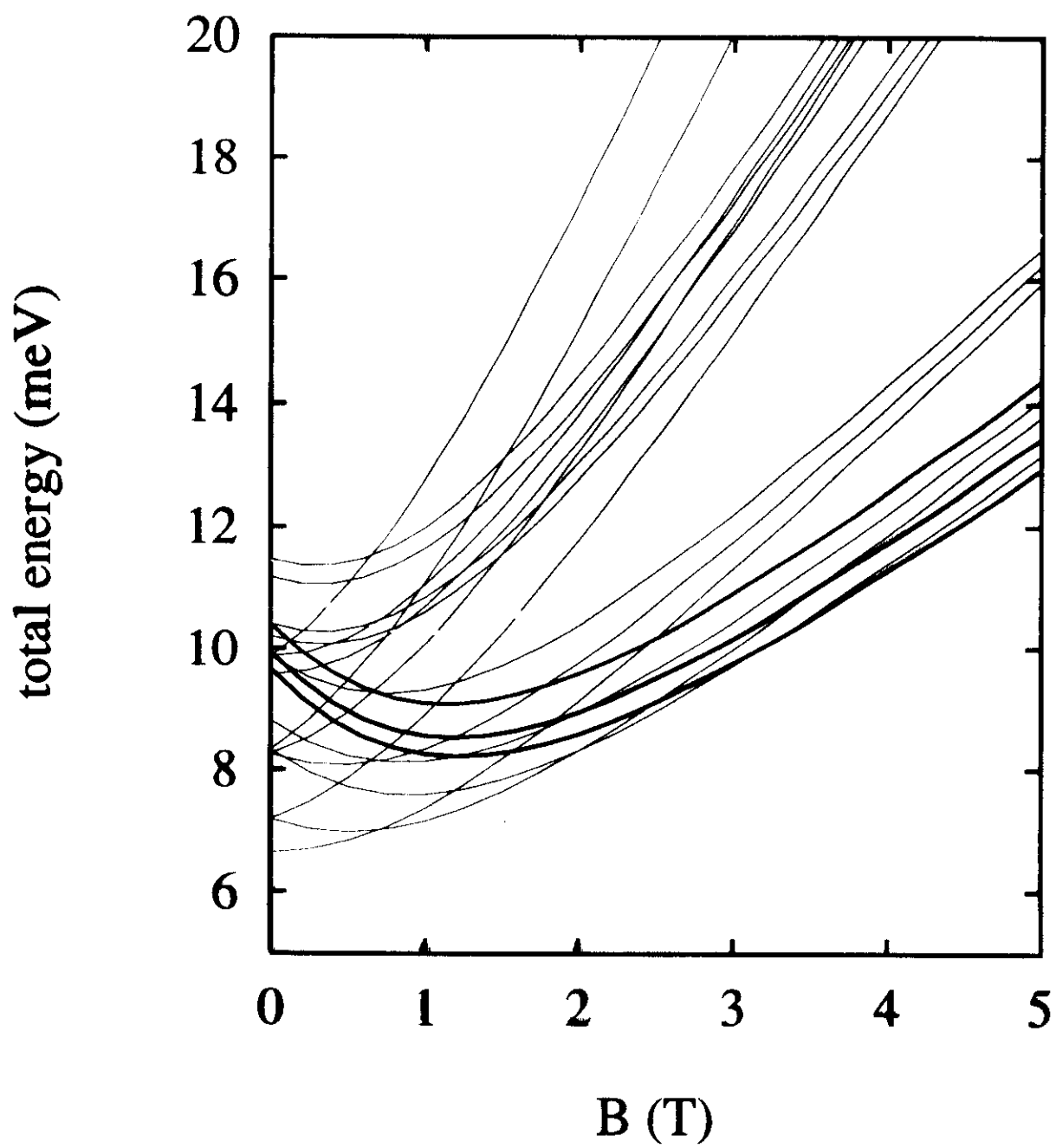
generalization for N particles:

$$\mathcal{H}_N = H(R_N, P_N) + \mathcal{H}_{N-1}(\vec{r}_i, \vec{p}_i)$$

Huser & Merkt,
Maksym & Chakraborty, ...

⇒ **generalized Kohn's Theorem
for FIR spectroscopy**

Interacting Spectrum



Spin and Angular Momentum

$$|\Psi(1, 2)\rangle = |\Phi_{CM}(1, 2)\rangle |\phi_{rel}(1, 2)\rangle \chi(S, S_z)$$

$$L_z^{rel} |\phi\rangle = \hbar M |\phi\rangle, \quad M = n_+ - n_-$$

Pauli-Principle: $P_{12} |\phi, \chi\rangle = -|\phi, \chi\rangle$

$$\begin{aligned} \langle \vec{r} | P_{12} |\phi, \chi\rangle &= \langle -\vec{r} | \phi, \chi\rangle \\ &= \langle \vec{r} | e^{iL_z \pi / \hbar} |\phi, \chi\rangle \end{aligned} \quad (1)$$

$\Rightarrow$

M odd	$\rightarrow \chi = \chi(1, S_z)$	(Triplet)
M even	$\rightarrow \chi = \chi(0, 0)$	(Singlet)

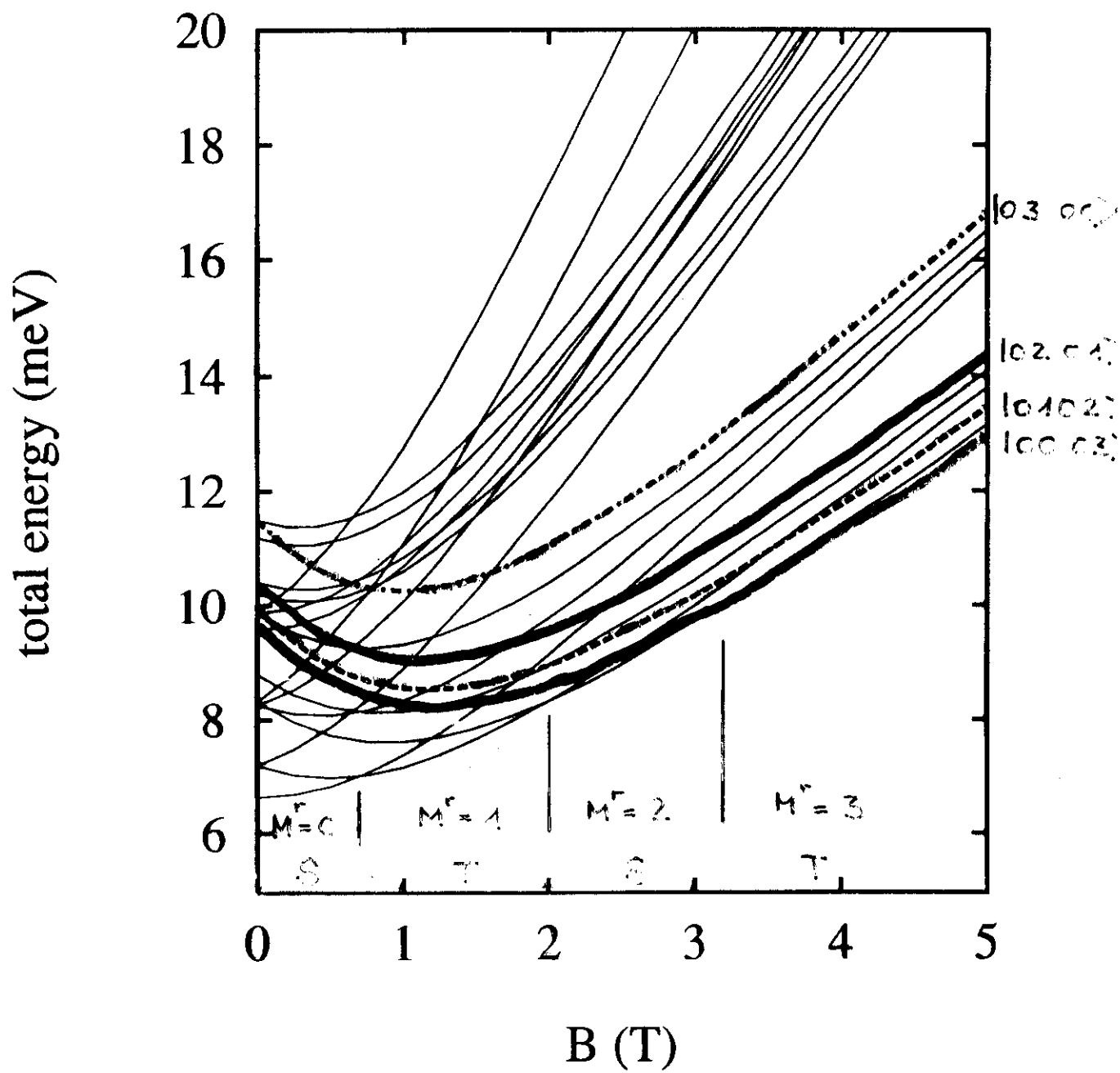
generalization for N electrons:

$$P_{ij} \vec{r}_m = e^{iL_z \alpha_{ijm} / \hbar} \vec{r}_n$$

for suitable relative coordinates

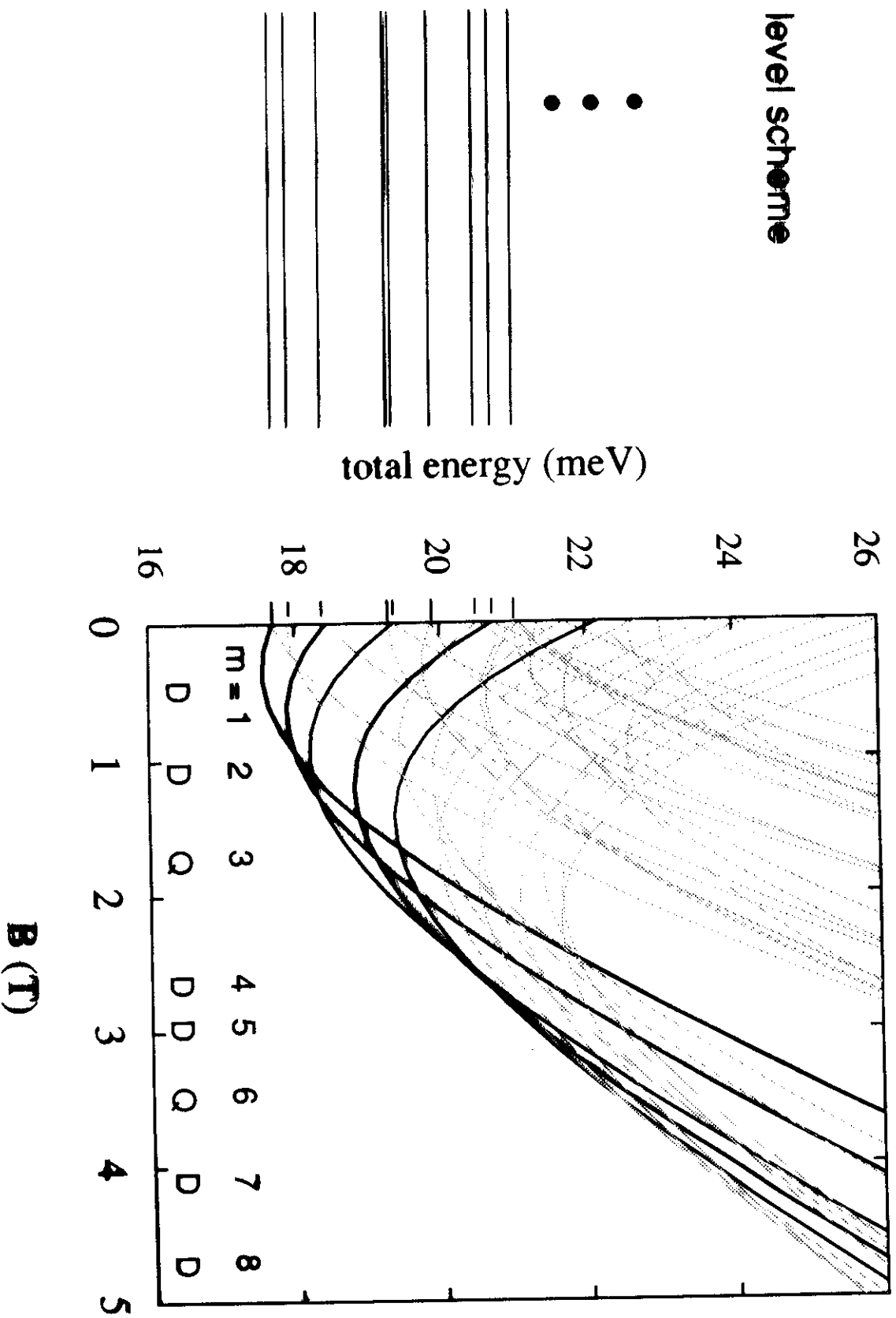
magic angular momentum quantum numbers
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Interacting Spectrum



Spectrum of QD-Lithium

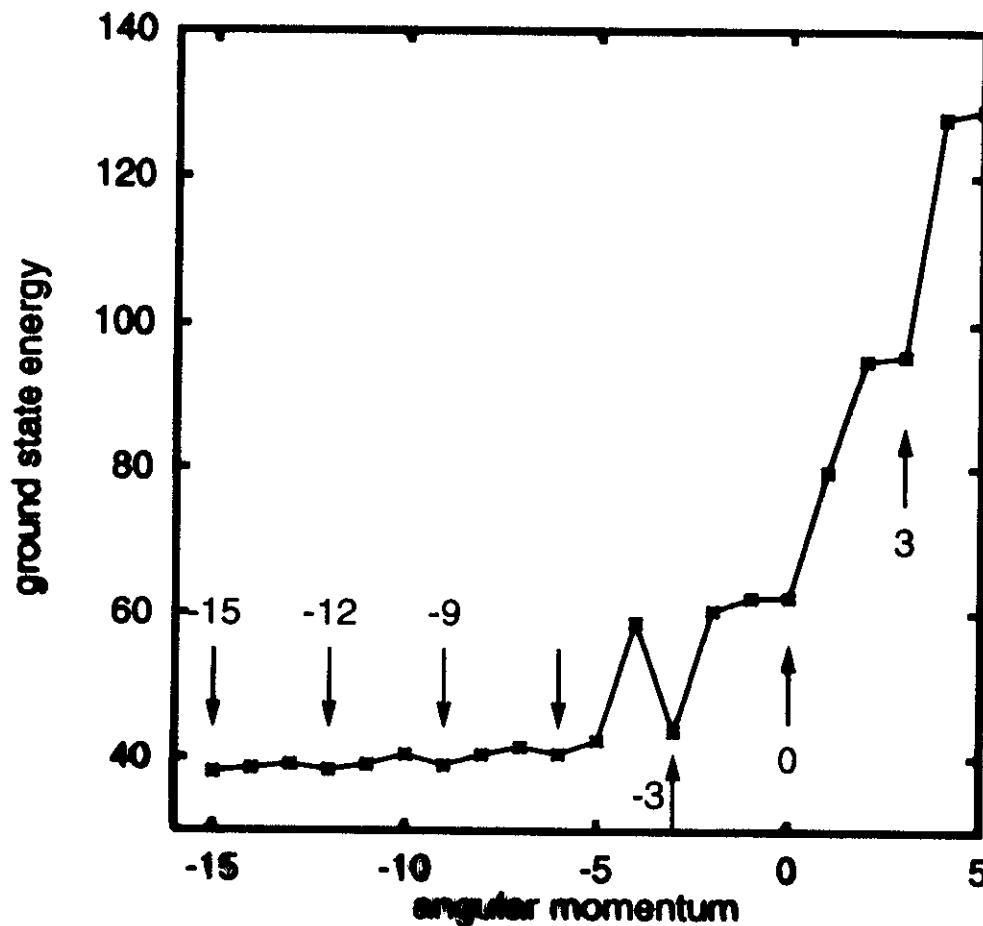
(3 electrons)



Magic Angular Momenta

Spin-polarized states occur

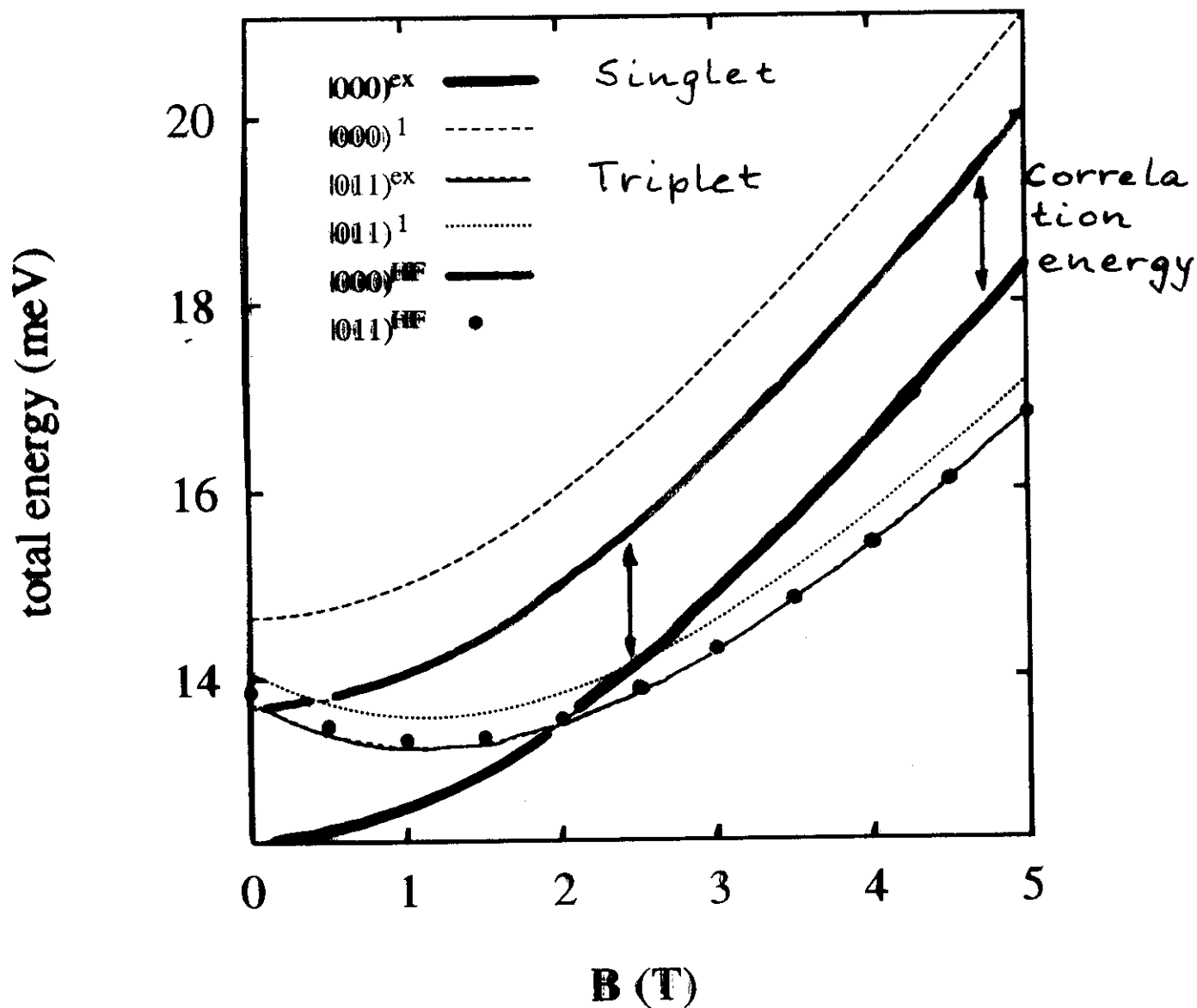
$$\text{for } \left. \begin{array}{l} \text{QD-Helium} \\ \text{QD-Lithium} \\ \text{QD-Beryllium} \end{array} \right\} \text{ at } L_z = \begin{cases} 1, 3, 5, \dots \\ 3, 6, 9, \dots \\ 2, 6, 10, \dots \end{cases}$$



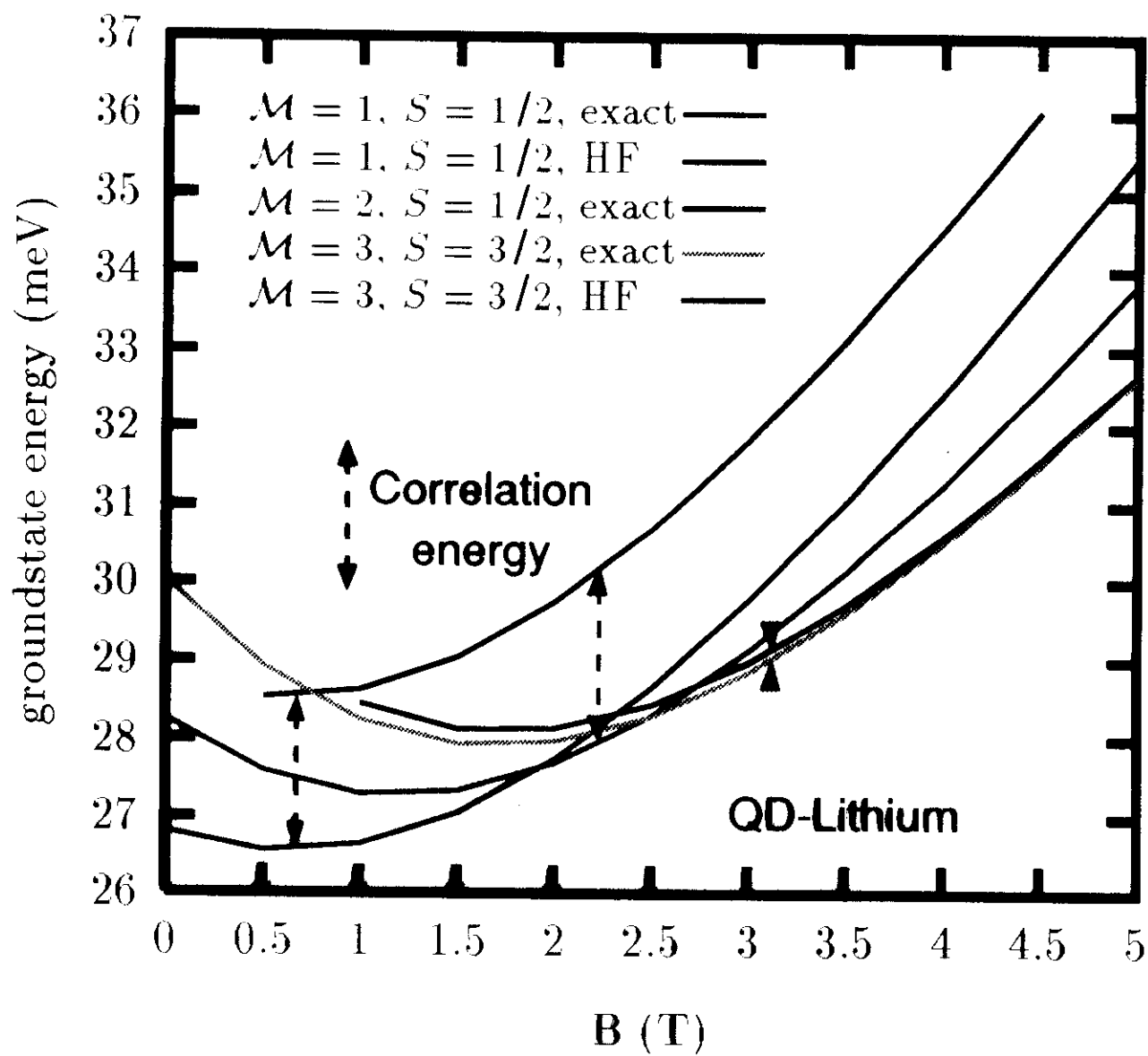
result from **Pauli-Principle** and
minimum energy vibrational modes

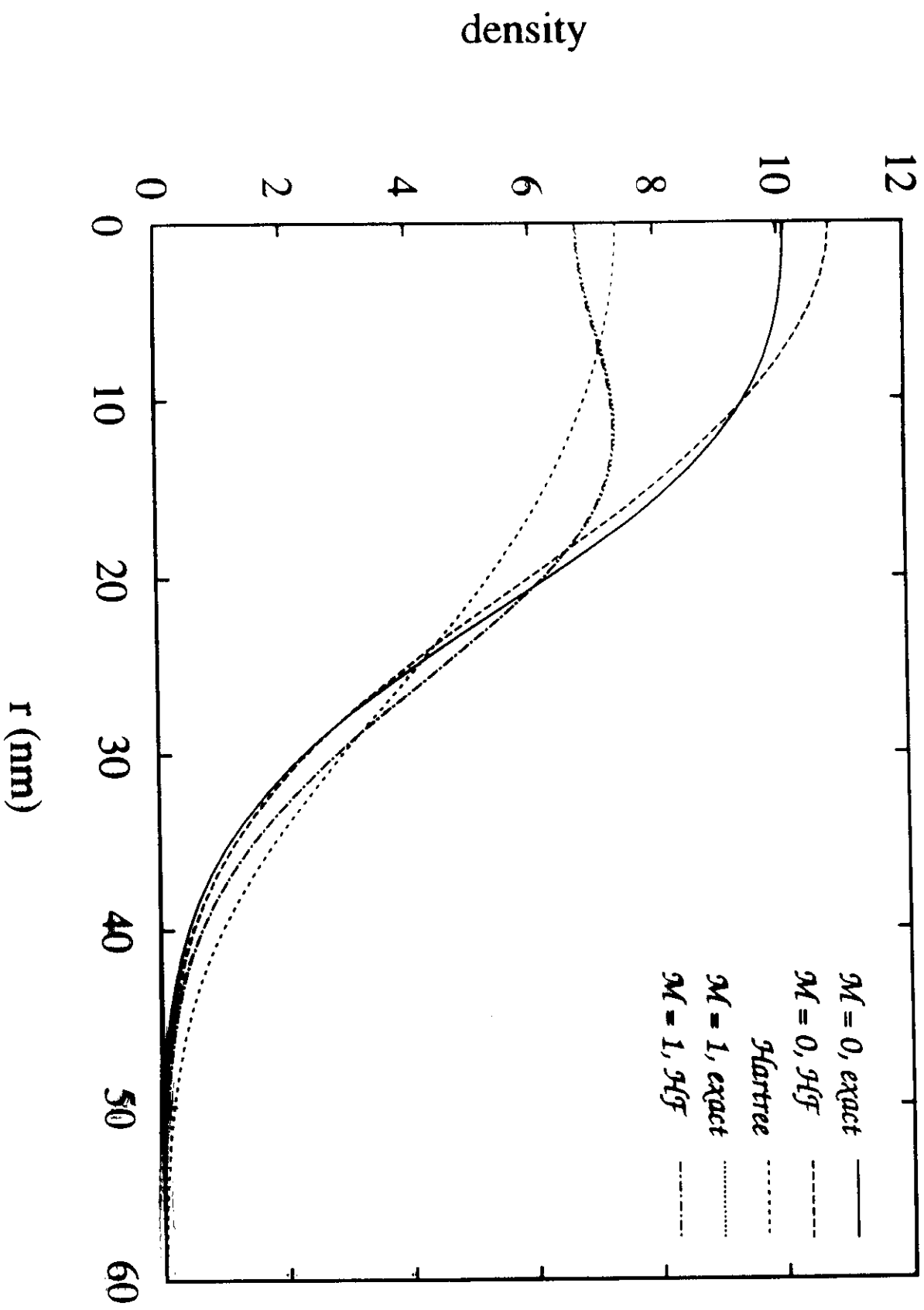
Maksym

Correlation Energy



Correlation Energy





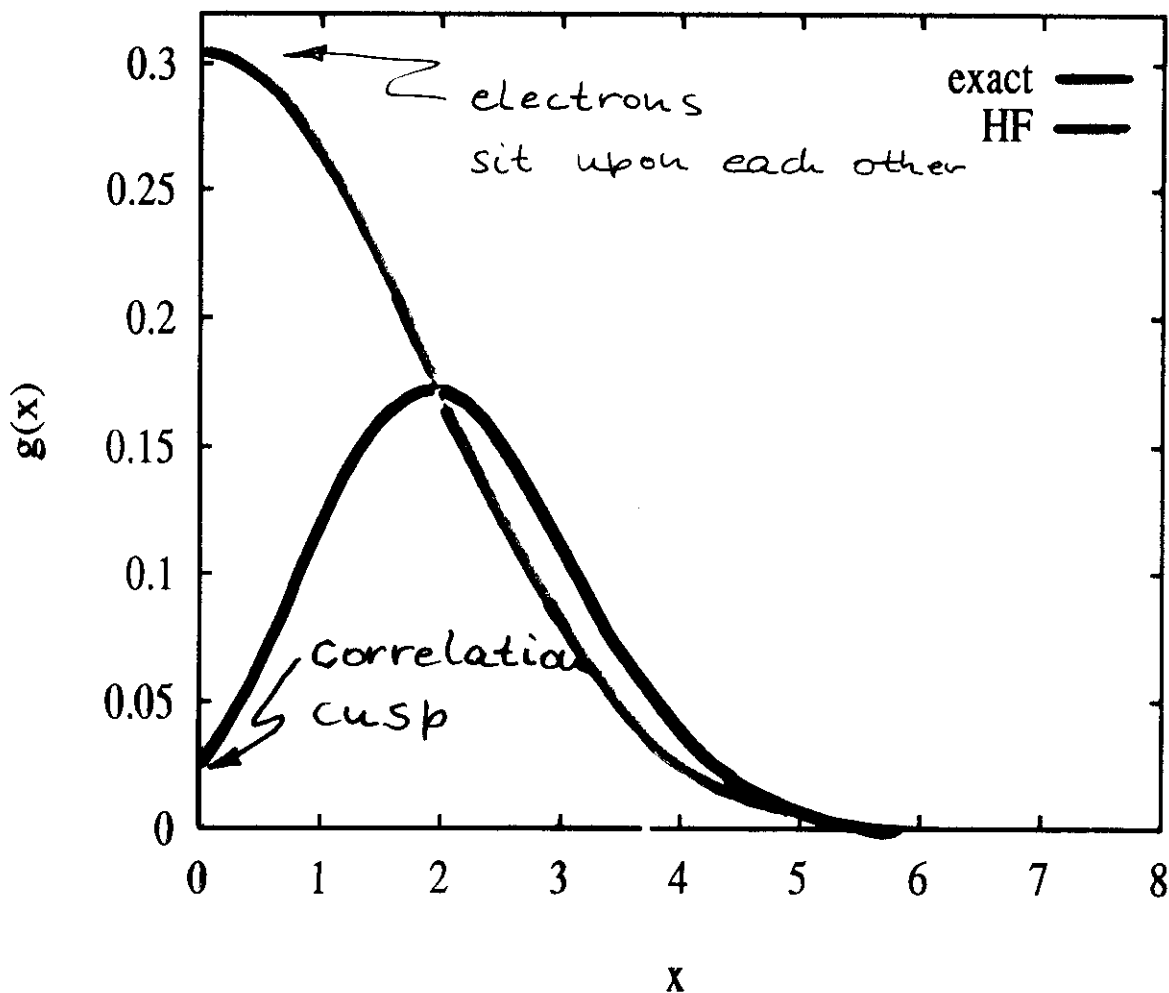
Pair-Correlation Function

why is the HF singlet energy that high?

Probability, for the electrons to be
a distance r apart from each other

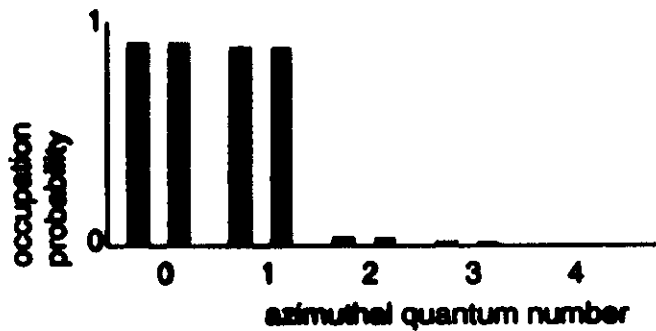
$$g(r) \propto \langle \delta(\vec{r} - (\hat{r}_1 - \hat{r}_2)) \rangle$$

$$\propto |\langle r | \tilde{O}^{rel} O^{rel} \rangle|^2$$

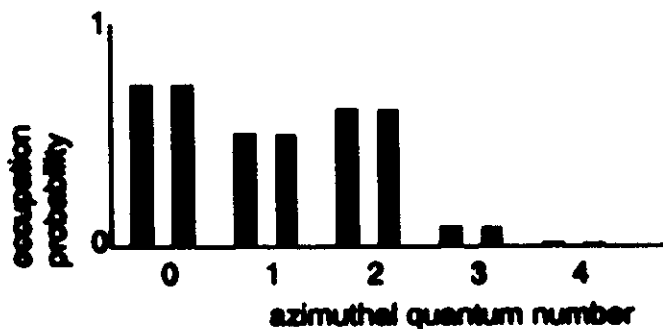


Wave Functions

uncorrelated 4-electron-state:



strongly correlated 4-electron-state:



Figures from J. Palacios *et al.*,
Superlattices and Microstructures, 15, 91, (1994)

strongly correlated states:
many single-particle states participate

Level Statistics

distribution of level spacing indicative of
chaos in quantum systems

generic probability

for finding 'nearest neighbor' level spacing s :

$$P(s) \propto \begin{cases} e^{-s/\langle s \rangle} & \text{Poisson} & \text{regular} \\ s e^{-s^2/\langle s \rangle^2} & \text{GOE} & \text{chaotic, } B=0 \\ s^2 e^{-s^2/\langle s \rangle^2} & \text{GUE} & \text{chaotic, } B \neq 0 \end{cases}$$

consequence of reduced symmetries

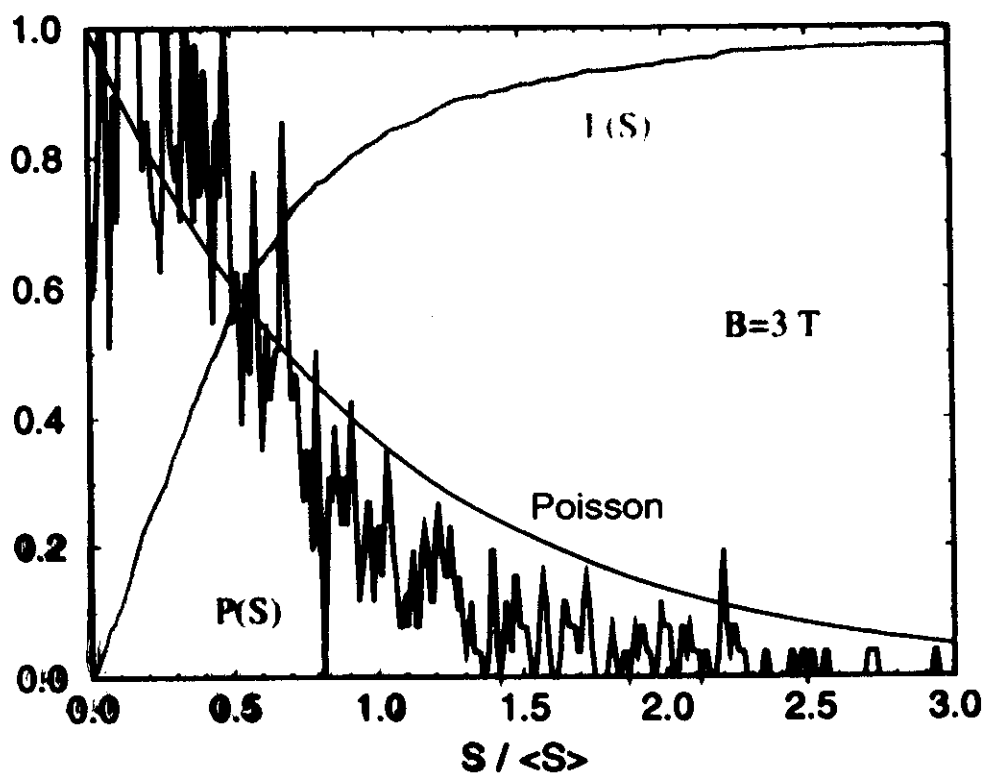
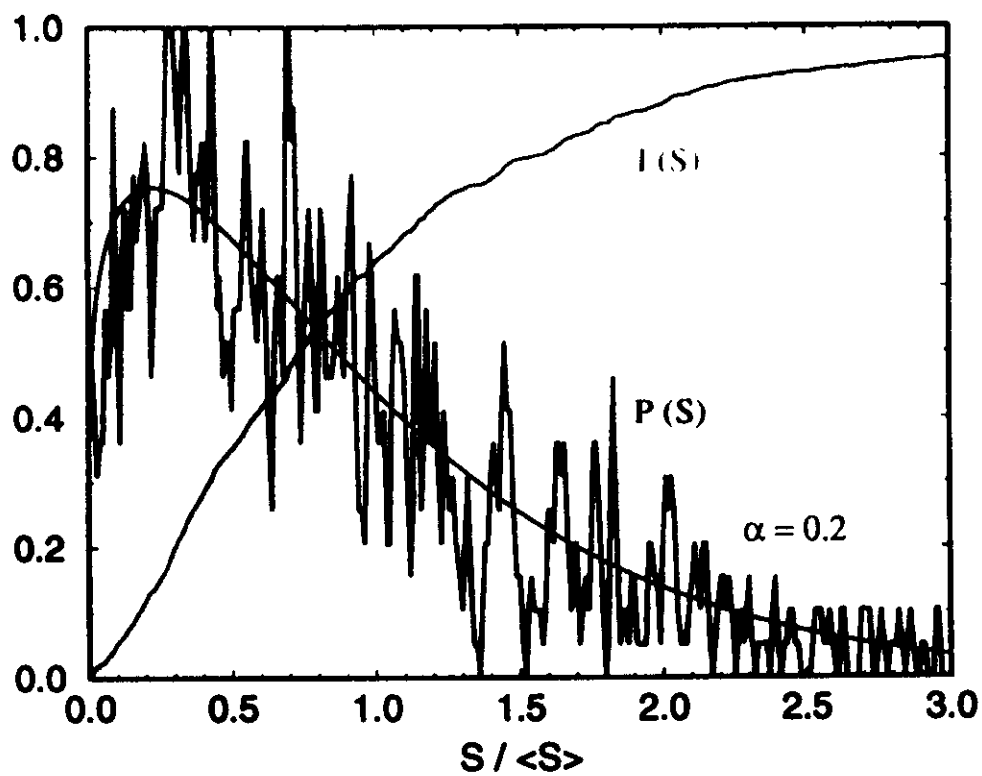
phenomenological interpolation formula:

Brody function $P_\alpha(s) \propto s^\alpha e^{-\beta_\alpha s^{\alpha+1}}$

$\alpha = 0 \rightarrow \text{Poisson}$

$\alpha = 1 \rightarrow \text{GOE}$

From Chaos to Regular behavior



Summary

- Coulomb interaction
 - fine structure
 - magic numbers
 - correlations
- Correlations
 - strong in partially spin-polarized states
 - weak in fully spin-polarized states
- chaos by interaction
oscillator localization vs Coulomb repulsion

