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SPRING COLLEGE ON AMORPHOUS SOLIDS
AND THE LIQUID STATE

14 April - 18 June 1982

PAIR POTENTIALS IN AMORPHOUS METALLIC ALLOYS:
ORDER, STABILITY AND DYNAMICS
Lecture III

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These are preliminary lecture notes, intended only for distribution to participants.
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Summary of 2nd lecture

Construction of a pseudopotential:

- Operator approach
 - optimized OPW-pseudopotentials
"hard core"
- Scattering approach
 - HAA-type potentials
filled to non linear charge density
"soft core"

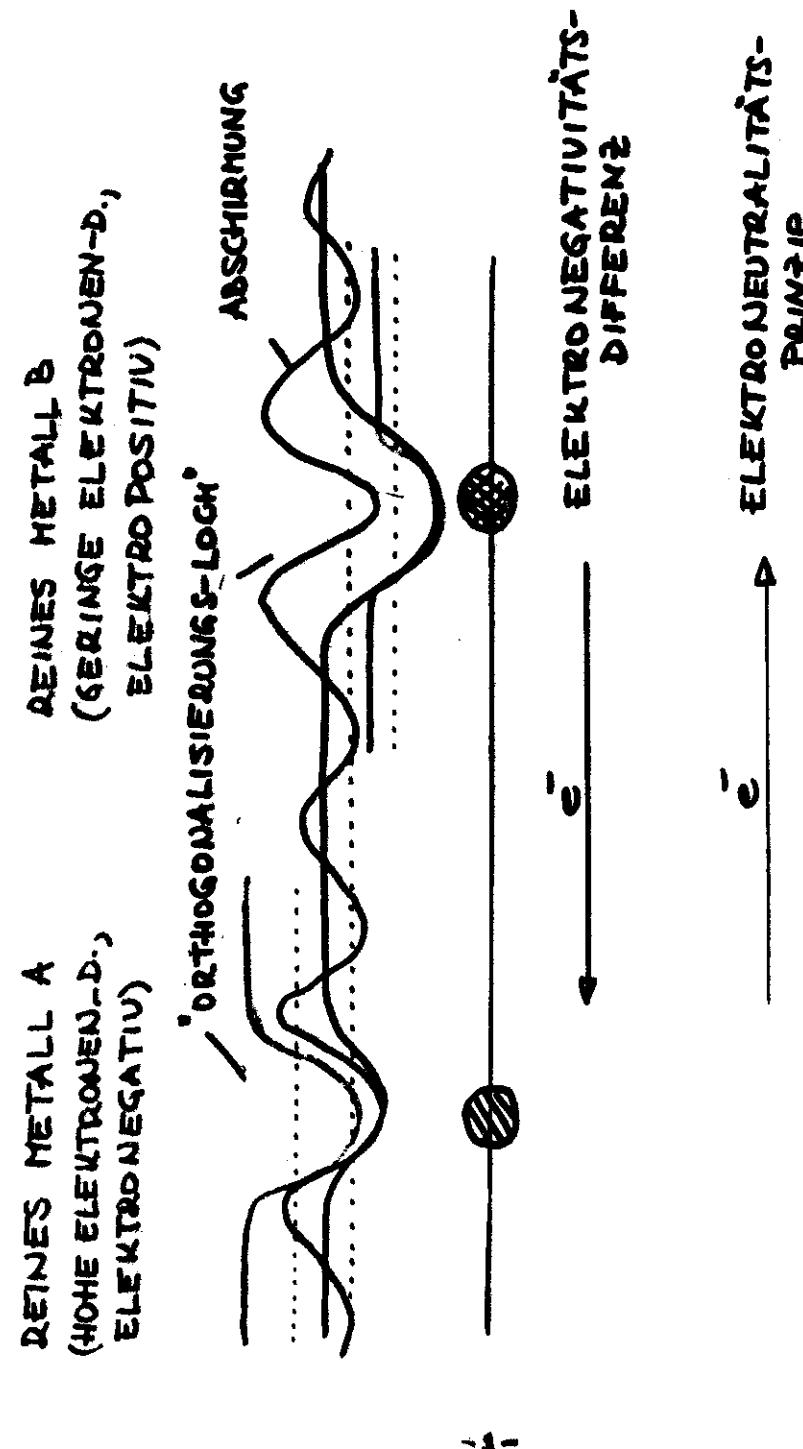
Exchange and correlation in the screening function

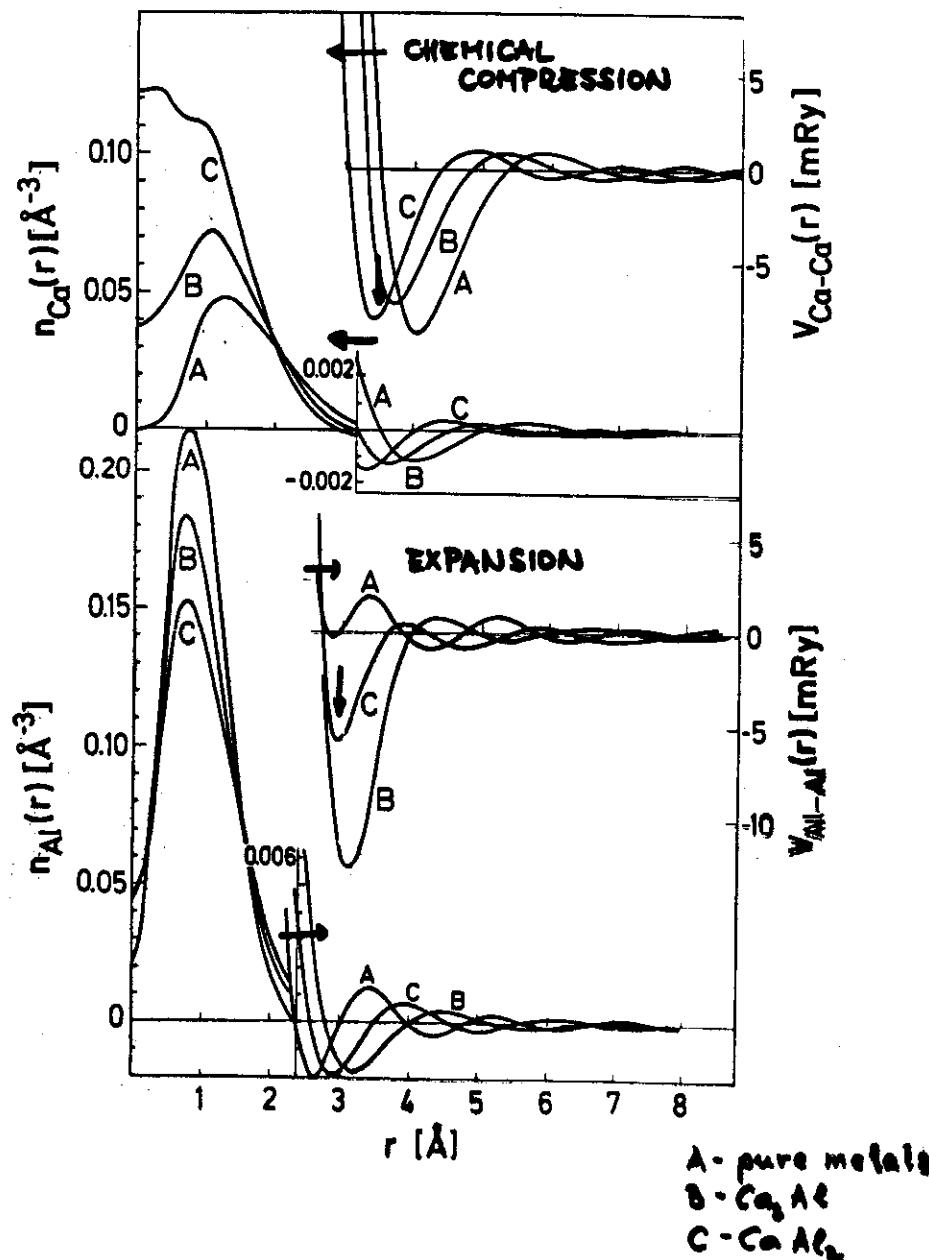
- Very important.
- require: - ratio's fraction of core pres. sum rule
- realistic e-e short range corr.

Pseudopotentials in alloys

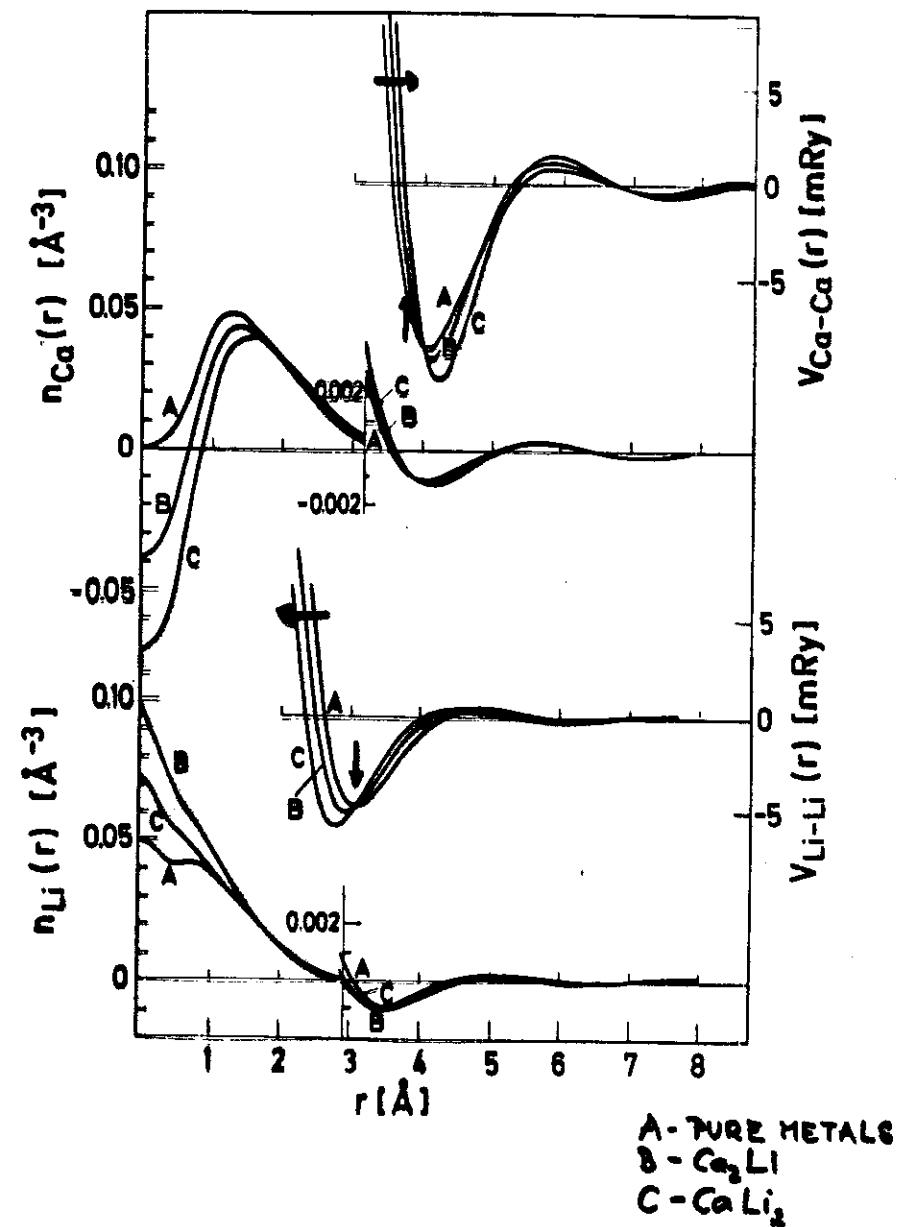
PP concentration dependent due to

- change in core-shift
- change in orthogonalization charges
- different energy-dependence
- shift of Fermi-level (different screening)





-3-



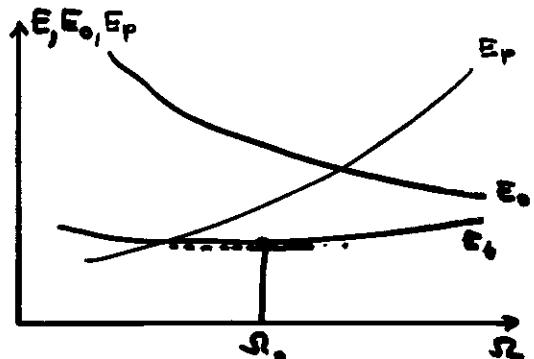
-4-

PHASE STABILITY IN PURE METALS
AND CRYSTALLINE ALLOYS

-27-

- PURE METALS

EQUILIBRIUM VOLUME DETERMINED BY
COMPETING VOLUME AND PAIR FORCES



$$E = E_{\Omega} + E_p$$

97-99% 1-3%

$$\rho = \rho_{\Omega} + \rho_p$$

>0 <0

$$\delta = \delta_{\Omega} + \delta_p$$

>0 >0

EQUILIBRIUM N.N. DISTANCE

≈ POSITION OF ATTR. MIN. IN $V(R)$

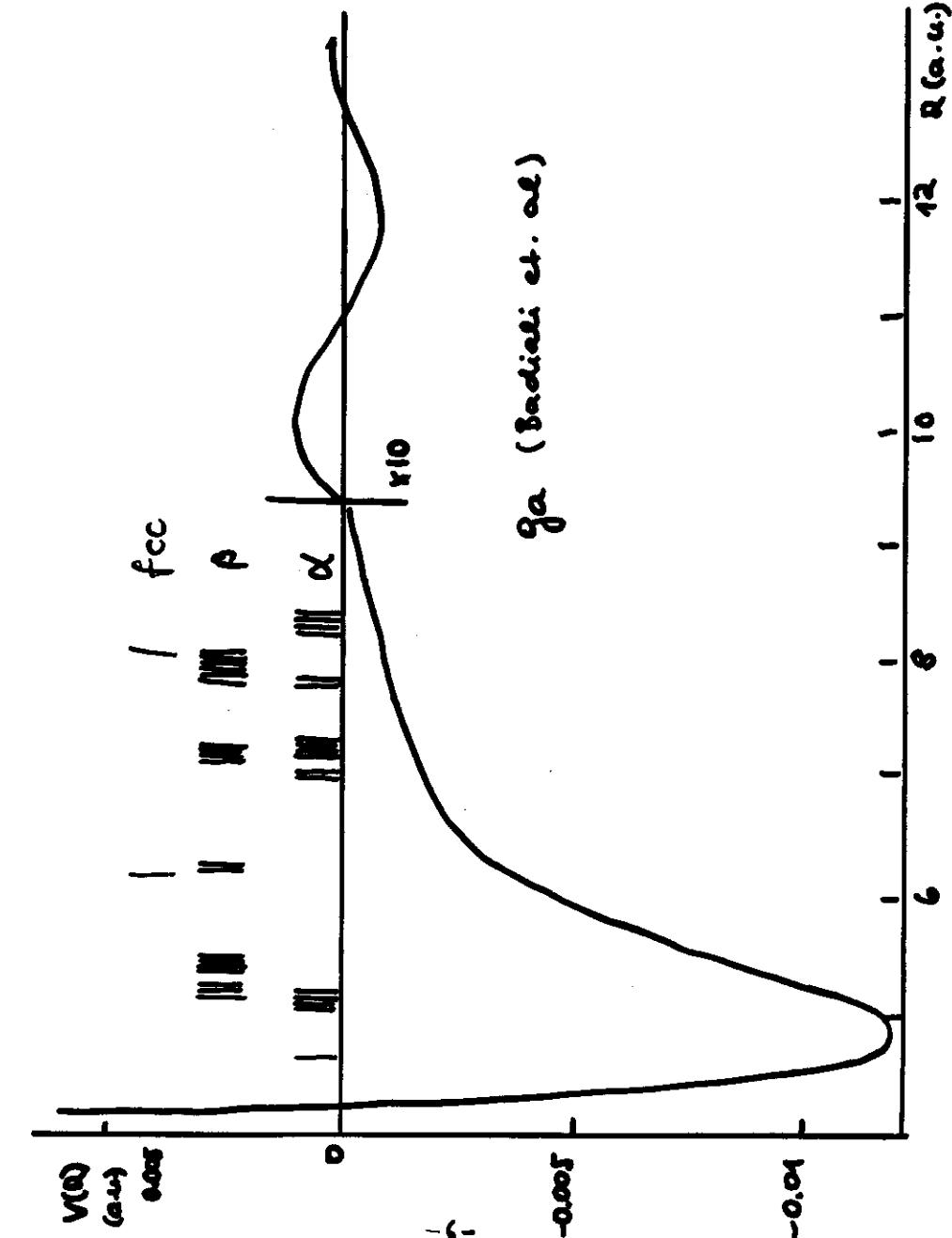
↔ CLOSE-PACKED STRUCTURE (Nb, Ti, Hf)

≠ POSITION OF ATTR. MIN. IN $V(R)$

↔ DISTORTED, OPEN STRUCTURE

e.g. g_a, s_h

-5-



-6-

-28-

Li-Mg

PHASENDIAGRAMM

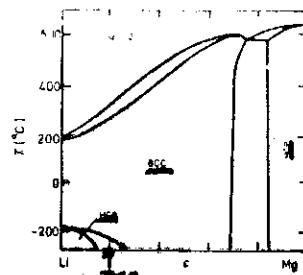
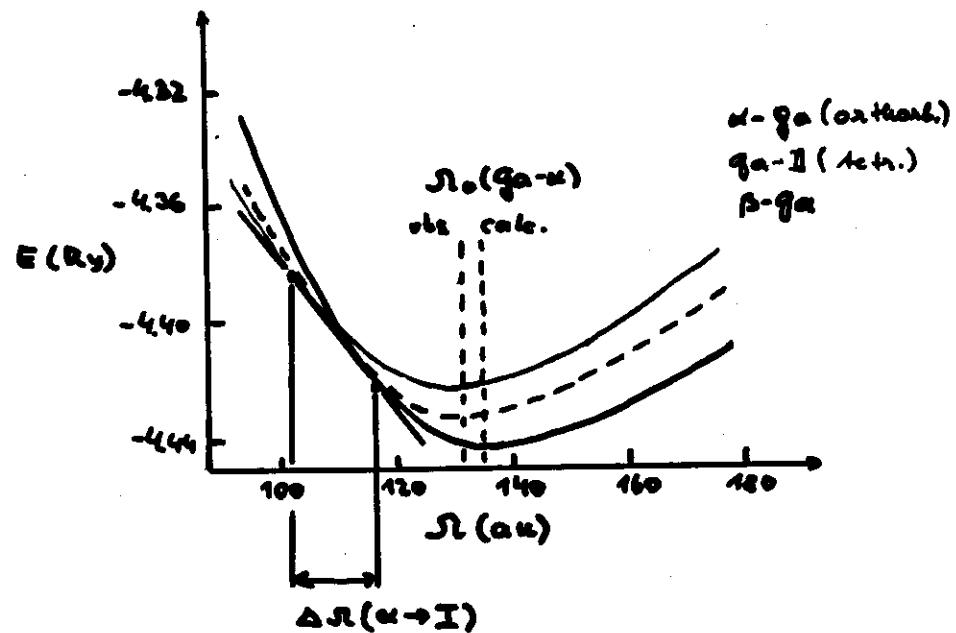


Figure 1. Phase diagram of the Li-Mg system (schematic).



Pressure-induced structural phase transformations
in solid γa

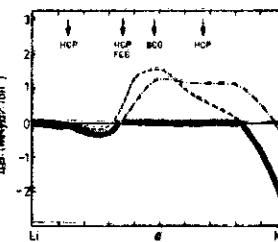


Figure 2. Structural energy differences ΔE relative to the FCC phase. The vertical arrows mark the concentrations at which the Fermi sphere touches a Jones-zone plane.

STRUKTURELLE
ENERGIEDIFFERENZEN
 $\downarrow \sim Q = 2k_F$

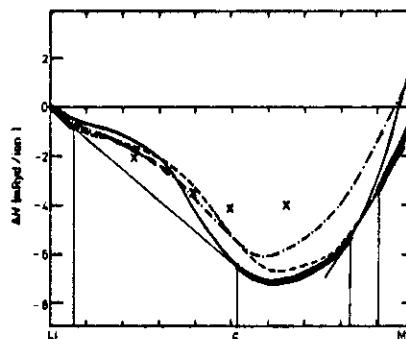


Figure 3. Heat of formation ΔH (— BCC, — HCP, — FCC) against concentration. The common tangent construction for the stability limits of the different phases is indicated. The crosses refer to the heat of formation measured by Mashkovetz and Puchkov (1965) for the liquid alloys.

BILDUNGSENTHALPIE
 — BCC
 — HCP
 — FCC

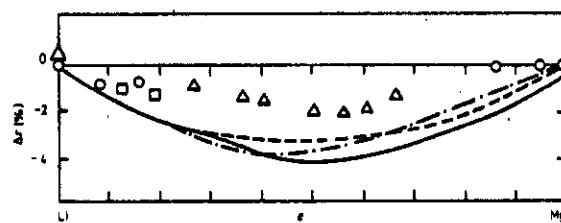


Figure 4. Deviations of the mean atomic radius from Vegard's rule (in percent). Full curve BCC phase; broken HCP; chain FCC (theoretical); Δ BCC, \square HCP, \square FCC (experimental, as referenced in the text).

BINARY ALLOYS

Finnis '74:

For local potentials (pure metal)

$$E_{\text{tot}} = \underbrace{2E_{\text{eg}} - 2J_0 B_{\text{eg}}}_{\downarrow} + \frac{1}{2} V_{\text{ind}}^{(2)}(0) \frac{\pi^2}{N} \sum_{\mathbf{k}} F(\mathbf{k})$$

"free electron" contribution to the heat of formation
(Pettifor '81)

$$\Delta H_{\text{eg}} = -\frac{5}{9} \frac{2}{\pi^2} (2.21 - 0.641 \bar{r}_s - 0.023 \bar{r}_s^2) [\Delta(\frac{1}{r_s})]^2$$

(1) loss of exchange and correlation
energy on alloying destabilizes the alloy

(2) relaxation of kinetic energy on alloying
stabilizes the alloy

(1) dominates at low electron densities
— — — high — — —

\bar{x} ... average valence

\bar{r}_s ... average electron dens. parameter

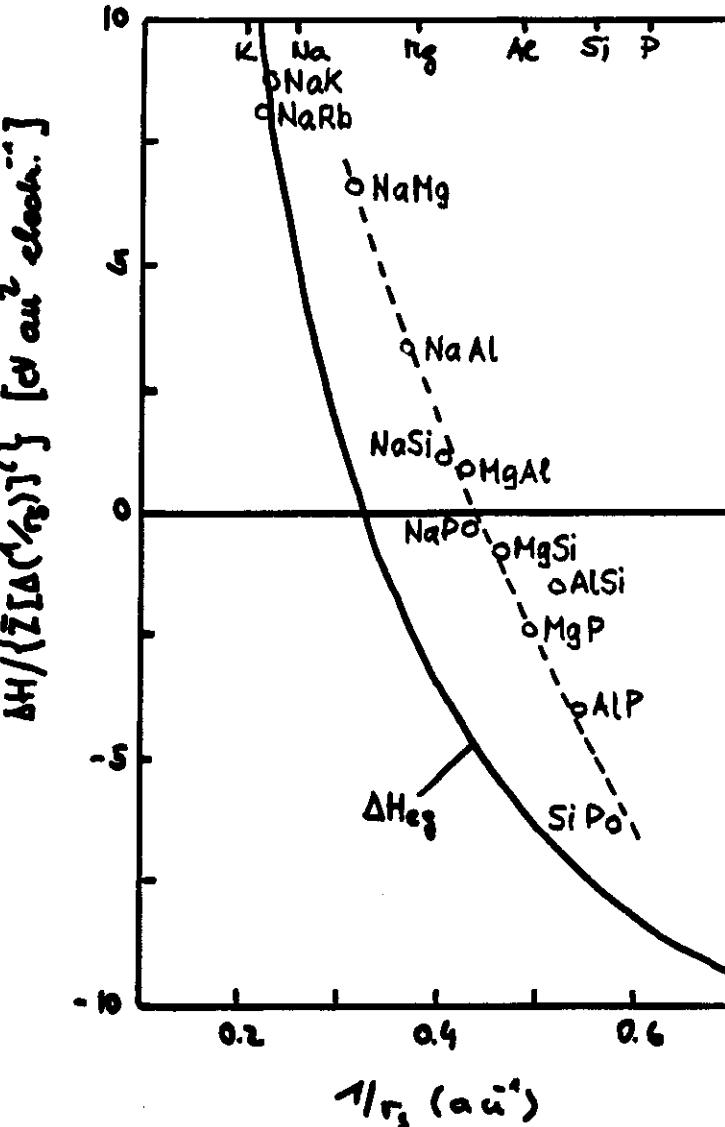
$\Delta(\frac{1}{r_s})$... difference in the two electron density parameter

$$\Delta H > \Delta H_{\text{eg}}$$

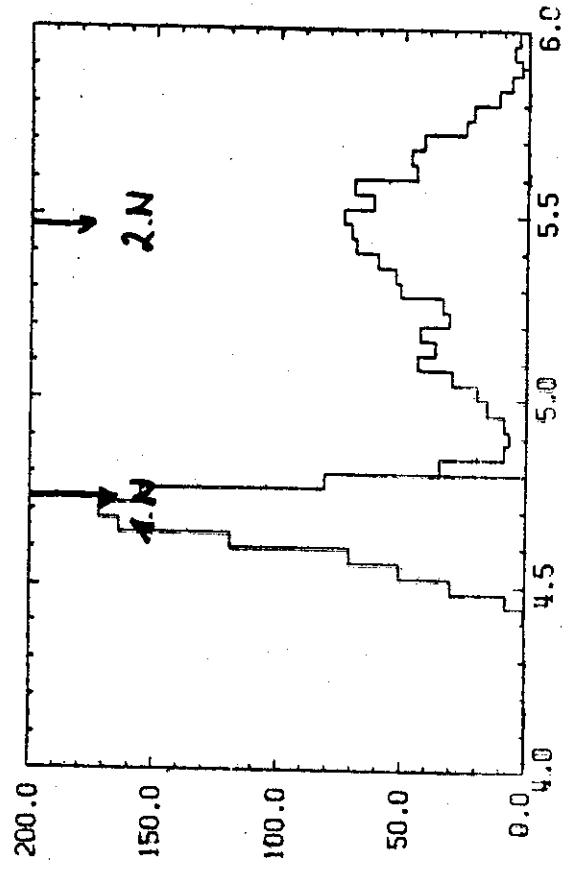
$$\Rightarrow \Delta H_p + \Delta(\frac{1}{2} V_{\text{ind}}(0)) > 0$$

pair interactions destabilize mixed crystals

Pettifor and Gelatt



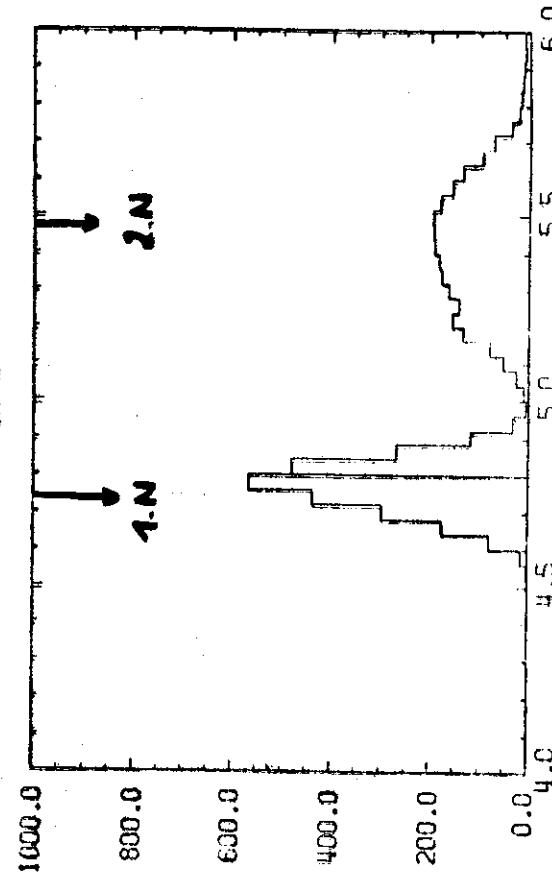
○ LDF-band structure calculation
(assuming CsCl-structure)



RDF(r), K29-RB71, PARTIELL K, HAI(100)

-1-

**STATISCHE GITTERVERTEILUNGEN IN K-RB
V. HEIMDORF UND HAFFNER (1978)
STATIC LATTICE DISTANCES**

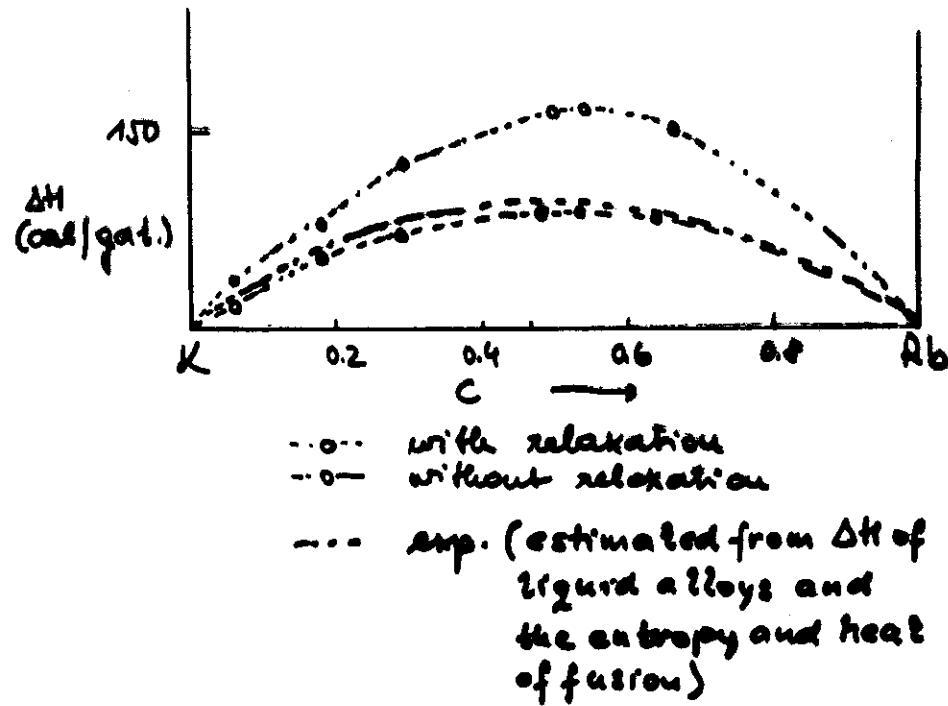


RDF(r), K29-RB71, PARTIELL K, HAI(100)

RADIALE VERTEILUNGSFUNKTION

-12-

K-Rb



ORDER-DISORDER TRANSITIONS

COMPLETE DISORDER

$$E = E_{\text{JL}} + \frac{1}{N} \sum_{i,j} V(R_{ij}) + \Delta E_{\text{bs}}$$

$$\Delta E_{\text{bs}} = c(1-c) \sum_{\vec{R}} F(\vec{R})$$

LONG-RANGE ORDER

BRAGG-WILLIAMS ORDER PARAMETERS η

$$\eta = (f_A - c)/(1-c) \quad f_A \dots \text{fraction of sites of type "A" occupied by A-atoms}$$

$$W(\vec{R}) = S(\vec{R}) \{ \bar{w}(\vec{R}) - c\eta \Delta w(\vec{R}) \} + \eta S_A(\vec{R}) \Delta w(\vec{R})$$

Young '79

SHORT-RANGE ORDER

WARREN-COWLEY ORDER PARAMETERS α_i

$$|D(\vec{R})|^2 = \frac{c(1-c)}{N} \sum_i \alpha_i(R_i) e^{-\frac{R_i^2}{2}}$$

$$E = E_{\text{JL}} + \frac{1}{N} \sum_{i,j} V(R_{ij}) + c(1-c) \sum_i N_i \alpha_i(R_i) \Delta V(R_i)$$

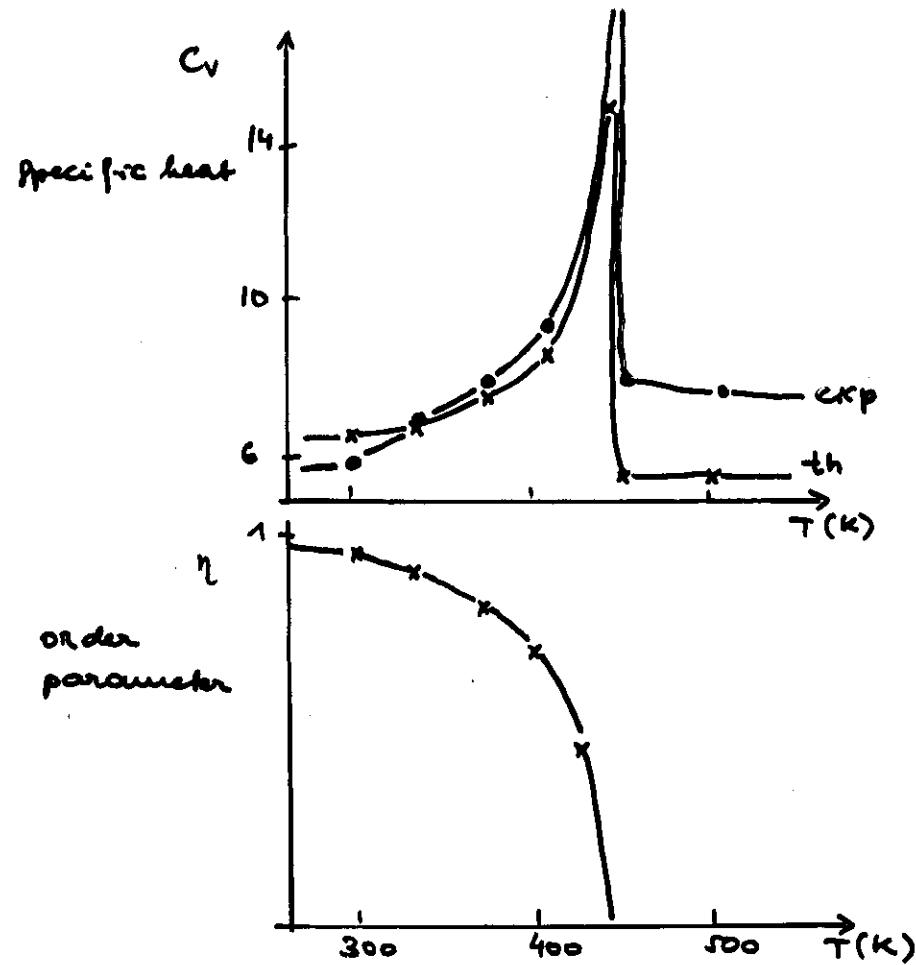
num of coordination shells
of radius R_i , coord. no. N_i

ORDER-PARAMETERS DETERMINED BY VARIATIONAL CONDITION

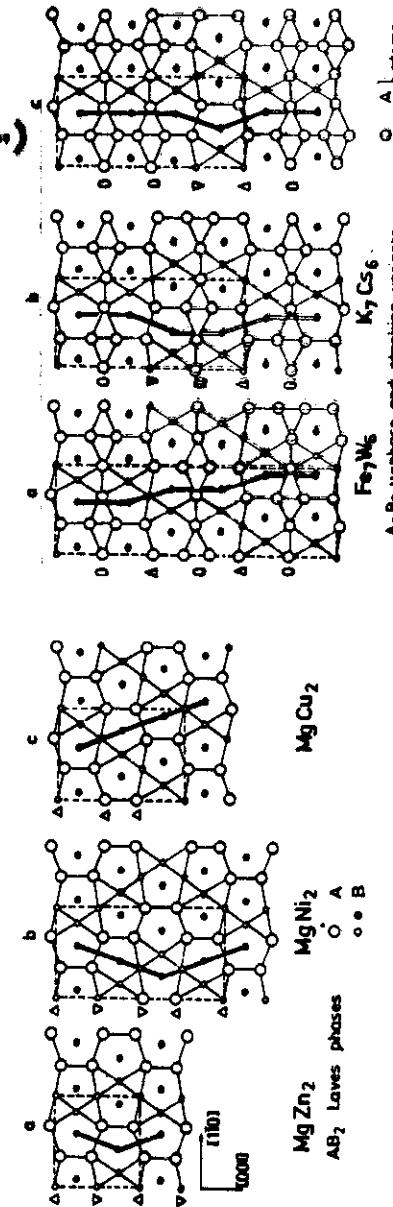
$$\left(\frac{\partial F}{\partial \eta} \right)_{S,T} = 0$$

$$\left(\frac{\partial F}{\partial \alpha_i} \right)_{S,T} = 0$$

LANGREICHWEITIGE ORDNUNG IN Mg_3Cd (LEUNG ET AL.)
LONG RANGE ORDER



-15-



-16-

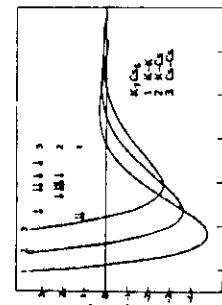


FIG. 5. Effective interatomic pair potentials in kJ/mol . The arrows indicate interatomic distances occurring in the observed h.c.p. structure.

FIG. 6. Enthalpies of formation of intermetallic compounds against the ratio of the atomic radii of the elements. ▲ Laves phase C14 (calculated structural parameters), ○ K₃Cd₂ structure (observed structural parameters), ● K₂Cd₃ structure (calculated structural parameters).

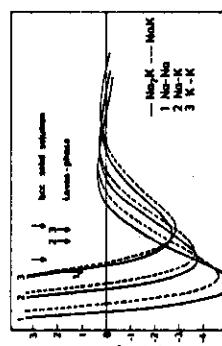
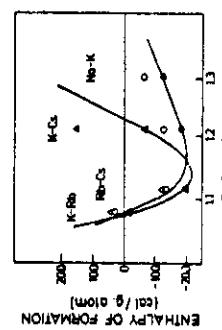
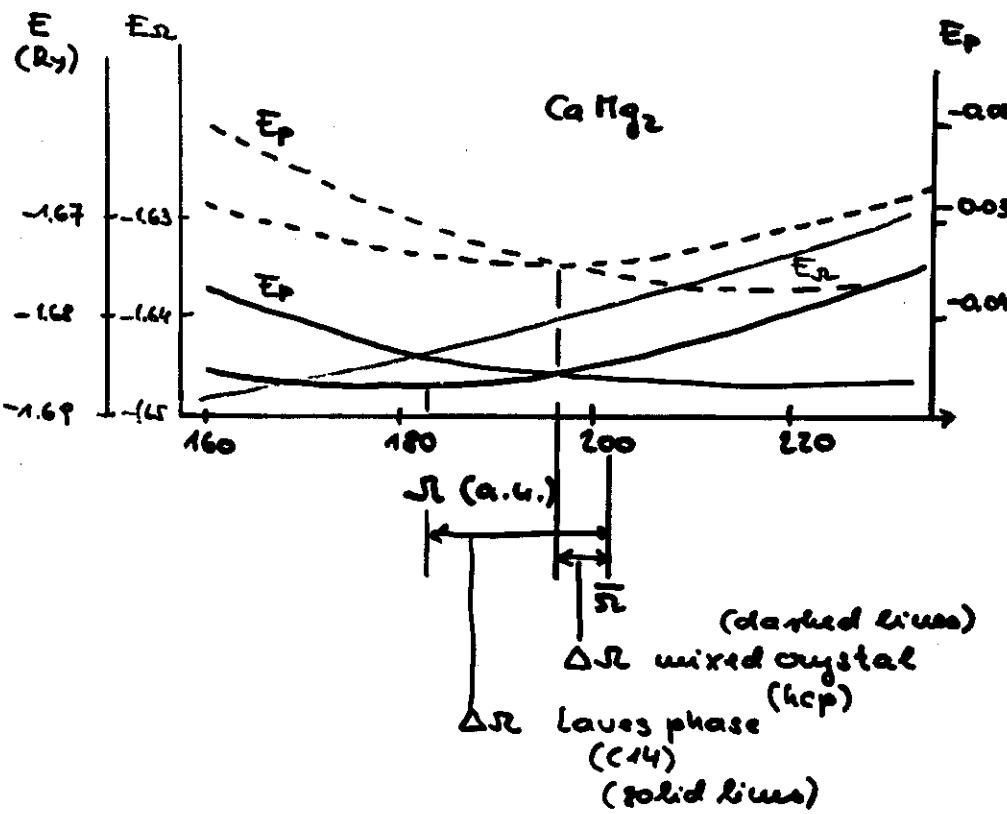


FIG. 5. Effective interatomic pair potentials for Na-K and Rb-K point and broken lines, respectively) at the equilibrium volume. The arrows indicate the nearest-neighbor distances in the C14 and in the h.c.p. structures.

Volume of formation of mixed crystals and intermetallic compounds



-17-

STRUCTURE AND THERMODYNAMICS OF LIQUID METALS AND ALLOYS

Calculation of the structure from the interatomic forces

- by solving one of the integral equations of the theory of liquid and liquid mixtures; difficult, successful for some cases, but not even attempted for alloys.
- by computer simulation: hard numerical work, very successful, confirms validity of DFT and OPW-pair potentials. But is it a theory?
- by thermodynamic perturbation theory

THERMODYNAMIC VARIATIONAL METHOD

Gibbs-Bogoliubov inequality

$$F \leq F_0 + \langle V \rangle_0$$

free energy of a reference system

expectation value of the perturbation

upper bound to exact free energy

Reference systems:

Hard-spheres: analytic solution of PY-equ., quite realistic $S(\mathbf{q})$

or semi-empirical Verlet-Ward $S(\mathbf{q})$

-18-

One-component plasma:

analytic thermodynamic functions,
approx. analytic $S(q)$

HS: H. Jones '73 OCP: Ross + De Nijs '60
 Groud + Ashcroft Parrinello et al '81
 Young et al
 Kafri
 i

$\text{Na}, \text{K}, \text{Rb}, \text{Cs}$: $F_{\text{OCP}} < F_{\text{HS}}$ S_{HS} slightly better

Li, polyvalent metals : $F_{\text{HS}} < F_{\text{OCP}}$ S_{HS} better

HS: corrections for soft repulsive potentials
and attractive potentials possible

OCP: how to go beyond G-B method?

$$\text{HS: } F \leq E_0 + E_{\text{bs}} + E_{\text{cs}} + \frac{3}{2} k_B T - TS_{\text{hs}} - \frac{1}{2} \gamma^2 T^2$$

kin. energy ~~some~~ electronic entropy

$$= E_R + E_p + \dots$$

$$\begin{aligned} \text{2-space} \quad & \left\{ \begin{array}{l} E_0 = 2E_{\text{cg}} + 2e/R_0 \\ E_{\text{bs}} = E_2 = \frac{N_0}{(2\pi)^2} \int F(r) S_{\text{hs}}(r) d^2 r \\ E_{\text{cs}} = \frac{\pi^2 e^2}{\pi} \int (S_{\text{hs}}(r) - 1) d^2 r \end{array} \right. \end{aligned}$$

$$\begin{aligned} \text{r-space} \quad & \left\{ \begin{array}{l} E_R = E_0 + \frac{1}{N} \sum_r F(r) + \lim_{R \rightarrow \infty} \left(\frac{4\pi r^2 e^2}{3R^2} + F(r) \right) \\ E_p = \frac{n_0}{2} \int g_{\text{hs}}(r) V(r) d^3 r \end{array} \right. \end{aligned}$$

VARIATIONAL CONDITION: $\left(\frac{\partial F}{\partial \sigma} \right)_{R,T} = 0$

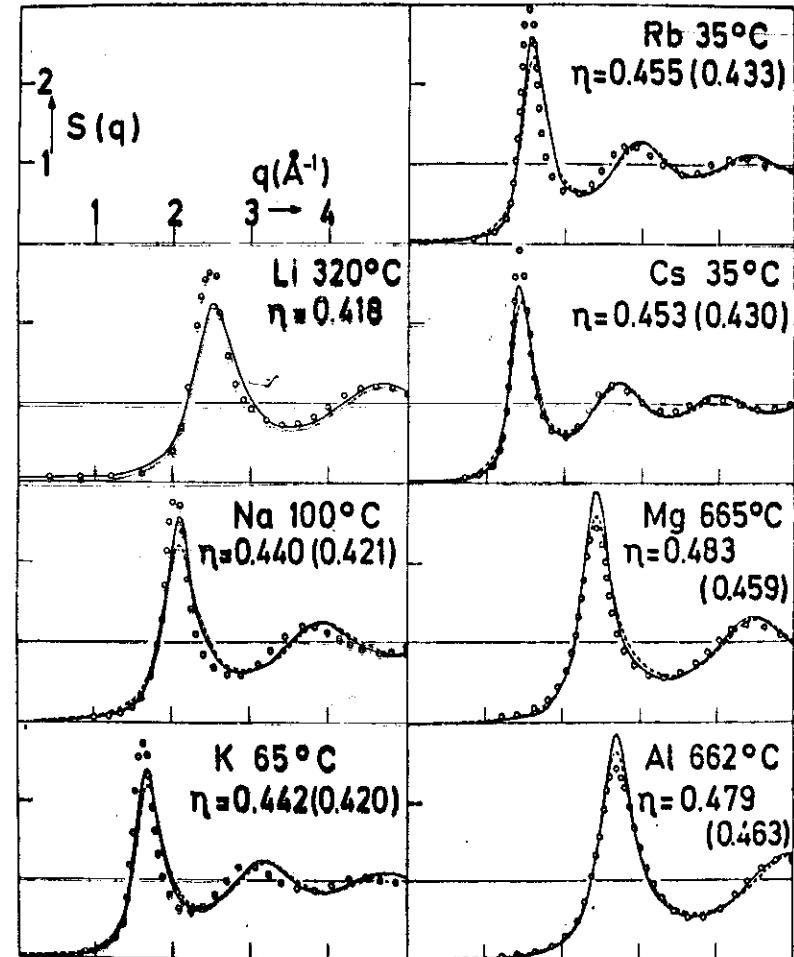


Fig. 1

Thermodynamic properties of liquid metals close to their melting point, as calculated using the GB-HS variational technique (Hoover '77)

a) Excess entropy $-SE/k_B$

	te.	exp
Li	3.75	3.70
Na	3.47	3.55
K	3.41	3.51
Rb	3.65	3.63
Cs	3.60	3.56
Hg	4.10	3.45
Ac	4.18	3.60

b) Bulk modulus $B_T (10^{10} \text{ dyne cm}^{-2})$

	te.	exp
Li	8.1	9.3
Na	3.7	3.8
K	1.7	1.8
Hg	21.6	19.7
Ac	37.4	30.7

c) Thermal expansion coeff. $\alpha \times (10^4 \text{ K}^{-1})$

	te.	exp
Li	3.9	3.2
Na	5.2	2.8
K	6.4	2.9
Hg	1.71	1.66
Ac	1.19	1.22

d) Specific heat

	te.	C_p/k_B	$\kappa \cdot C_p/C_v$
Li	4.3	3.7	1.3
Na	4.0	3.8	1.4
K	4.6	3.9	1.4
Rb	5.0	3.8	1.5
Cs	3.4	3.9	1.7
Hg	4.2	4.0	1.3
Ac	3.9	3.5	1.2

Gibbs-Dogo-Jacobson method for liquid mixtures

$$E_0 = 2E_{eg} + \frac{1}{\pi} (c \sigma_A + (1-c) \sigma_B) / R_0$$

$$E_{eg} = \frac{g_0}{(2\pi)^3} \int_0^\infty \sum_{i,j=1}^n (c_i c_j)^{1/2} F_{ij}(2) S_{ij}(2) d^3 r$$

↑
Machnicki-Langreth
pairwise structure f.

$$\rightarrow \frac{g_0}{(2\pi)^3} \int_0^\infty (F_{NN}(2) S_{NN}(2) + 2F_{Nc}(2) S_{Nc}(2) + f_{cc}(2) f_{cc}(2)) d^3 r$$

↑
Blatina-Thornton

$$E_{eg} = \frac{c^2}{\pi} \int_0^\infty \sum_{i,j=1}^n (c_i c_j)^{1/2} z_i z_j (S_{ij}(2) - 1) d^3 r$$

$$= \frac{c^2}{\pi} \int_0^\infty \left\{ \frac{1}{2} (S_{NN}(2) - 1) - 2 \bar{z} \Delta z S_{Nc}(2) - \bar{z}^2 (f_{cc}(2) - c(1-c)) \right\} d^3 r$$

$$E_R = E_0 + \frac{1}{2} [c V_{NN}(0) + (1-c) V_{BB}(0)]$$

$$- \lim_{R \rightarrow \infty} \sum_{i,j=1}^n c_i c_j \left(\frac{4 \pi^2 i^2 j^2 e^2}{3 R_0 2^5} + F_{ij}(2) \right)$$

$$E_p = \frac{n_0}{2} \sum_{i,j=1}^n c_i c_j \int_{A,B} g_{ij}(R) V_{ij}(R) d^3 R$$

$$= \frac{n_0}{2} \int [g_{NN}(R) \tilde{V}(R) + 2g_{Nc}(R) V_{Nc}(R) + f_{cc}(R) \Delta V(R)] d^3 R$$

Variational conditions:

$$\left(\frac{\partial F}{\partial \sigma_A} \right)_{\sigma_B, R, T} = 0 \quad \left(\frac{\partial F}{\partial \sigma_B} \right)_{\sigma_A, R, T} = 0$$

σ_A, σ_B ... hard sphere diameters

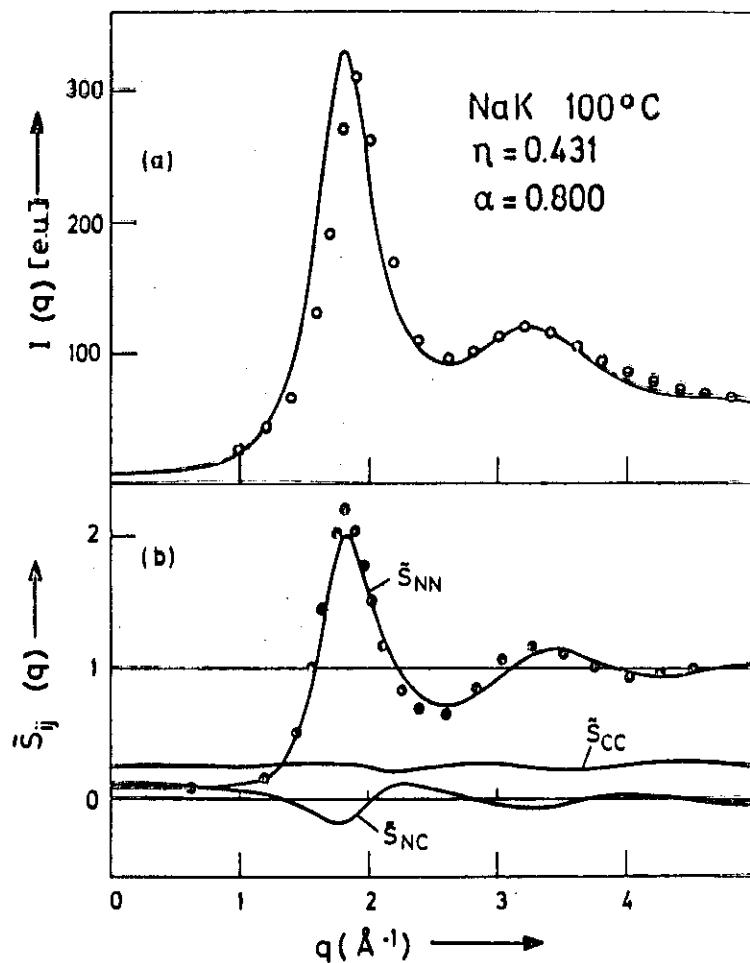
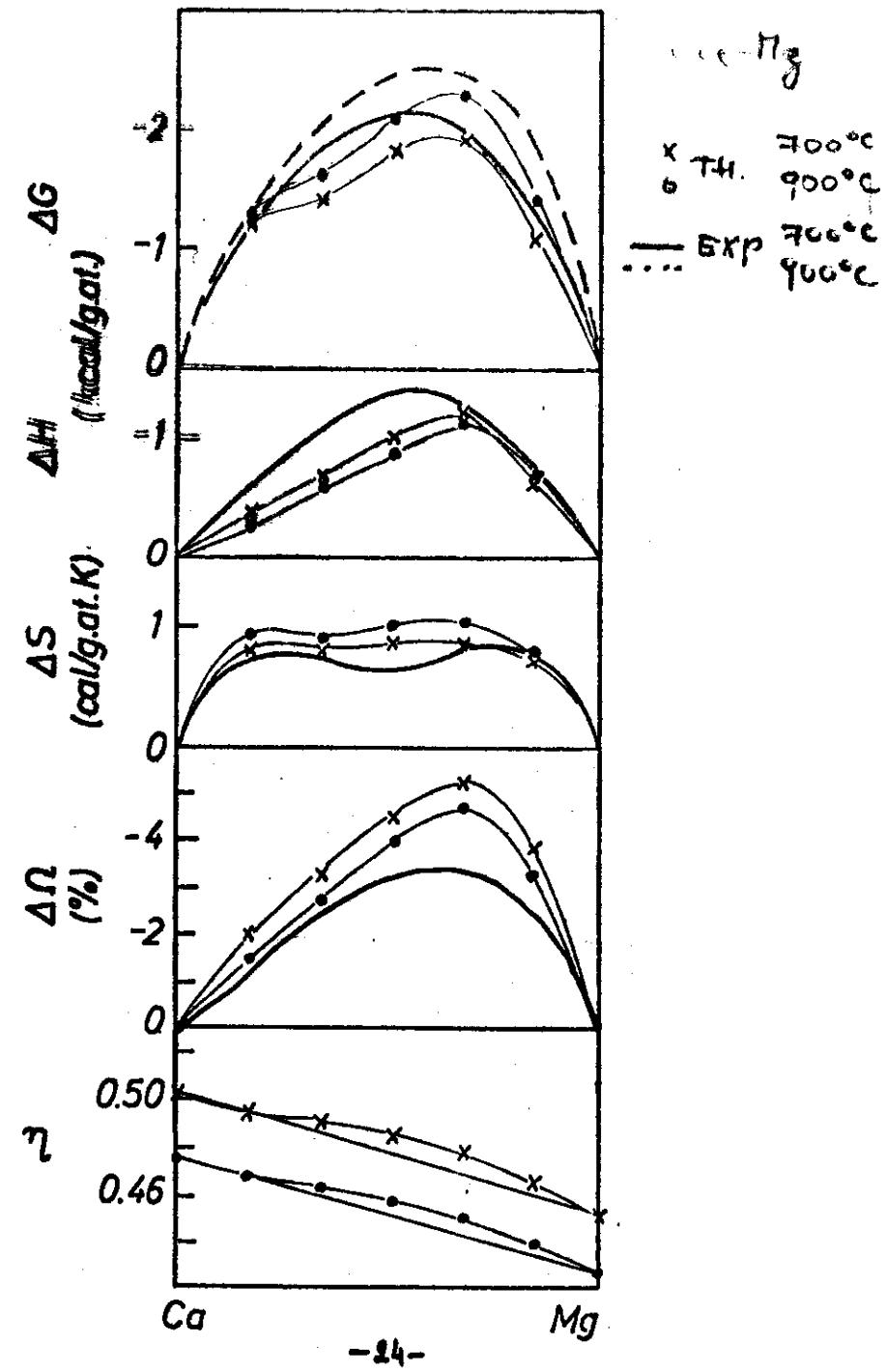
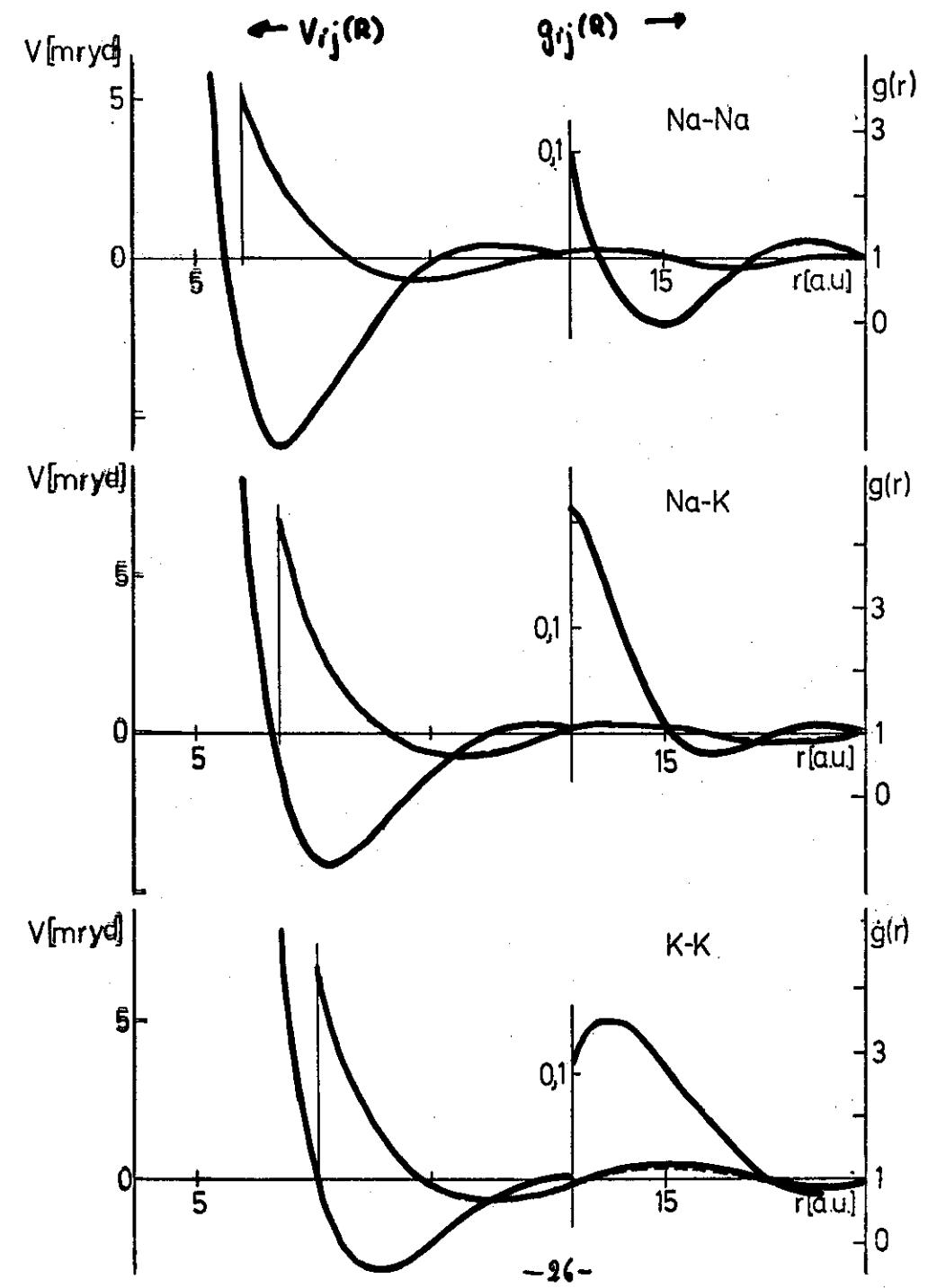
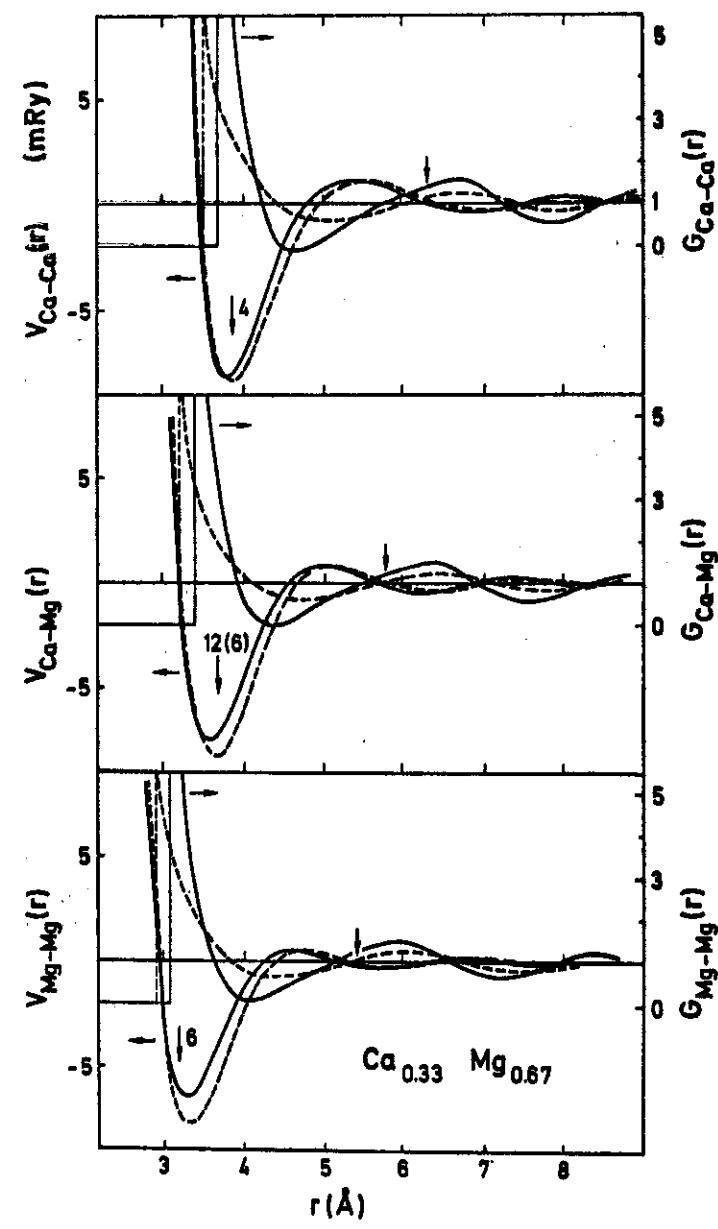


Fig. 4

-13-





References to 3rd lecture

Applications of pseudopotentials to structural stability

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Structure and thermodynamics of liquid metals and alloys

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Gibbs-Hogdgin's reference system

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