



INTERNATIONAL ATOMIC ENERGY AGENCY
UNITED NATIONS EDUCATIONAL, SCIENTIFIC AND CULTURAL ORGANIZATION



INTERNATIONAL CENTRE FOR THEORETICAL PHYSICS
34100 TRIESTE (ITALY) - P.O. B. 586 - MIRAMARE - STRADA COSTIERA 11 - TELEPHONES: 224281/3/3/4/5/6
CABLE: CENTRATOM - TELEX 460392-1

SMR/94- 29

SPRING COLLEGE ON AMORPHOUS SOLIDS
AND THE LIQUID STATE

14 April - 18 June 1982

IMPURITIES, TRAPS AND RECOMBINATION CENTERS
IN A SiHx
(Part I)

B. ABELES
Exxon Research and Engineering Co.
P.O. Box 45
Linden, N.J. 07036
U.S.A.

These are preliminary lecture notes, intended only for distribution to participants.
Missing or extra copies are available from Room 230.

AMORPHOUS SEMICONDUCTORS

Impurities Traps and Recombination Centers in a SiH_x

Application to Solar Cells

B. Abeles
Exxon Research and Engineering Co.
P. O. Box 45
Linden N. J. 07036

Collaborators
G. D. Cody
T. Tiedje
C. D. Wronski

1

IMPURITIES, TRAPS AND RECOMBINATION CENTERS

IN HYDROGENATED AMORPHOUS SILICON

B. Abeles
Exxon Research and Engineering Company
P. O. Box 45
Linden, NJ 07036

Abstract

The theme of the five lectures is the density of gap states (DGS) in hydrogenated amorphous silicon (a-SiH_x) and its significance to solar cells. The DGS in amorphous semiconductors is of fundamental importance because it provides information on the nature of the disorder and defects. The DGS is also of practical importance because it determines the ultimate performance of amorphous semiconductors in electronic devices. There are several reasons why a-SiH_x is enjoying prominence at this time over other amorphous semiconductors. First of all, of course, there is worldwide interest in this material because of its many novel technological applications (in particular for solar cells). From a scientific point of view a-SiH_x is so attractive because of its excellent photo-electric properties which allow one to probe the density of states with techniques known from crystalline semiconductors;

2

the ability to make reproducibly well characterized materials which can be doped n and p type and in which the density of states can be varied over many orders of magnitude. The central issue of amorphous semiconductor solar cells is the requirement of a low DGS combined with a low optical bandgap, ($E_G \sim 1.5$ eV). The low DGS is required to keep recombination at a low level so that most of the photogenerated carriers are extracted. A band gap of $E_G \sim 1.5$ eV is required so that a large portion of the solar spectrum is absorbed. In Lecture I we show that there exists a fundamental trade off between E_G and DGS--the lower the E_G is, the larger is the DGS. The optimum E_G a-SiH_x solar cells is near 1.7 eV. After discussing some of the basic concepts in photocarrier transport in Lecture I, we discuss in Lecture II transient photoconductivity techniques for studying the band tail DGS. The band tail states act as traps and determine the drift mobility of holes and electrons. In solar cells the band tails ultimately limit the open circuit voltage (Lecture V). In Lecture III we discuss the use of steady state photoconductivity techniques to determine the optical absorption and diffusion lengths of photocarriers. In Lecture IV we discuss deep level defect states and their capture cross section (e.g., dangling bonds), intrinsic to the material or induced by P or B doping, and observed in optical absorption ESR, photoluminescence and DLTS. An overview of the DGS of a-SiH_x is presented. In Lecture V we discuss the basic principles of operation of a thin film solar cell and the fundamental limits on the efficiency imposed by the DGS.

I. EFFECT OF DISORDER ON OPTICAL BAND GAP.

BASIC CONCEPTS IN PHOTO-CARRIER TRANSPORT

1. DENSITY OF STATES - FUNDAMENTAL CONCEPTS.
2. REQUIREMENTS ON DENSITY OF STATES FOR EFFICIENT SOLAR CELL.
3. RELATION BETWEEN EXPONENTIAL ABSORPTION EDGE AND BAND GAP.
4. FUNDAMENTAL CONCEPTS IN PHOTOCARRIER TRANSPORT - RECOMBINATION OF FREE CARRIERS - CAPTURE CROSS SECTION - QUASIFERMI LEVEL - DEMARCATION LEVEL - RESPONSE TIME, DRIFT MOBILITY AND FREE CARRIER MOBILITY.

II. BAND TAIL STATES

1. TIME RESOLVED TRANSPORT IN AMORPHOUS SEMICONDUCTORS - GENERAL
PHYSICAL CONCEPTS
2. TIME OF FLIGHT EXPERIMENTS
3. MULTIPLE TRAPPING MODEL OF DISPERSIVE TRANSPORT - CURRENT DECAY AT
SHORT TIMES ($t < t_T$) - TRANSIT TIME t_T - CURRENT DECAY AT LONG TIMES
($t > t_T$)
4. INTERPRETATION OF EXPERIMENTAL RESULTS OF TIME OF FLIGHT FOR $a\text{-SiH}_x$
- BAND TAIL PARAMETERS OF $a\text{-SiH}_x$
5. PHOTOCURRENT DECAY IN $a\text{-SiH}_x\text{:P}$
6. INDUCED ABSORPTION
7. MOVIE: "MULTIPLE TRAPPING-DISPERSIVE AND NONDISPERSIVE TRANSPORT."

5

III. STEADY STATE PHOTOCONDUCTIVITY

1. PHOTOCONDUCTIVITY WITH OHMIC CONTACTS (P.C.) BASIC PRINCIPLES
2. NON-LINEAR LIGHT DEPENDENCE OF P.C.
3. BIAS LIGHT ILLUMINATION
4. COMPARISON OF P.C. WITH PHOTO-THERMAL DEFLECTION SPECTROSCOPY
5. INFRARED QUENCHING
6. PHOTOCONDUCTIVITY WITH BLOCKING CONTACTS (SOLAR CELL COLLECTION
EFFICIENCY C.E.) BASIC PRINCIPLES
7. SAMPLE GEOMETRY AND TYPICAL EXPERIMENTAL RESULTS
8. INTERNAL PHOTOEMISSION
9. COMPARISON OF α DETERMINED FROM P.C. AND C.E.
10. C.E. WITH STRONGLY ABSORBED LIGHT
11. DIFFUSION LENGTH OF MINORITY CARRIERS FROM C.E.

6

IV. DEEP LEVEL DEFECT STATES IN a-SiH_x

1. HOW a-SiH_x FILMS ARE PREPARED.
2. DETERMINATION OF H, P AND B DOPING CONCENTRATIONS IN FILMS.
3. ELECTRICAL AND OPTICAL PROPERTIES OF P AND B DOPED FILMS.
4. DENSITY OF STATE MODEL FOR THE OPTICAL ABSORPTION AND EVIDENCE FOR DANGLING BOND STATES.
5. CAPTURE CROSS SECTION FOR HOLES AND ELECTRONS
6. SPIN RESONANCE OF DEFECT STATES
7. QUENCHING OF PHOTOLUMINESCENCE BY DEFECTS.
8. DLTS OF a-SiH_x: P
9. DEFLECTION LAYER MODULATED ESR.
10. SUMMARY OF DENSITY OF STATES OF a-SiH_x.

V. AMORPHOUS SEMICONDUCTOR SOLAR CELLS

1. HOW THE SOLAR CELL WORKS.
2. SOLAR SPECTRUM AND THE IDEAL SOLAR CELL.
3. SOLAR CELL PARAMETERS - OPEN CIRCUIT VOLTAGE V_{OC} , SHORT CIRCUIT CURRENT I_{SC} , FILL FACTOR F , AND EFFICIENCY $EF = FI_{SC}V_{OC}$.
4. RECIPE FOR EFFICIENT THIN FILM SOLAR CELL.
5. COMPARISON BETWEEN c-Si, c-GaAs and a-SiH_x solar cells.
6. FUNDAMENTAL LIMITATIONS ON V_{OC} , I_{SC} AND F IN AMORPHOUS SEMICONDUCTORS.
7. HOW TO FURTHER IMPROVE THE EFFICIENCY.

IMPURITIES, TRAPS AND RECOMBINATION CENTERS

IN HYDROGENATED AMORPHOUS SILICON

FIVE LECTURES

BASIC CONCEPTS OF DAS

I. EFFECT OF DISORDER ON OPTICAL BAND GAP.

BASIC CONCEPTS IN PHOTO-CARRIER TRANSPORT

II. BAND TAIL STATES

TIME RESOLVED PHOTOCONDUCTIVITY

III. STEADY STATE PHOTOCONDUCTIVITY

MAINLY TECHNIQUES TO DETERMINE
DEEP LEVEL STATES

IV. DEEP LEVEL DEFECT STATES IN a-SiH_x

DANGLING BONDS ETC

V. AMORPHOUS SEMICONDUCTOR SOLAR CELLS

RELEVANCE OF DAS

I. EFFECT OF DISORDER ON OPTICAL BAND GAP.

BASIC CONCEPTS IN PHOTO-CARRIER TRANSPORT

1. DENSITY OF STATES - FUNDAMENTAL CONCEPTS.

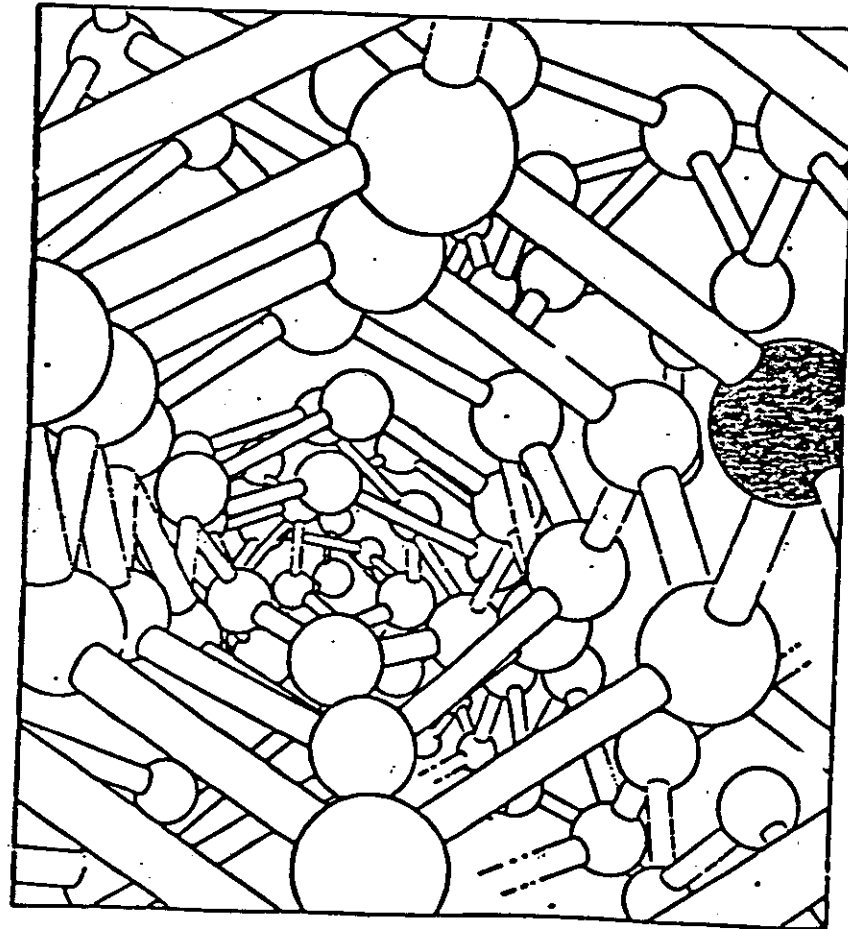
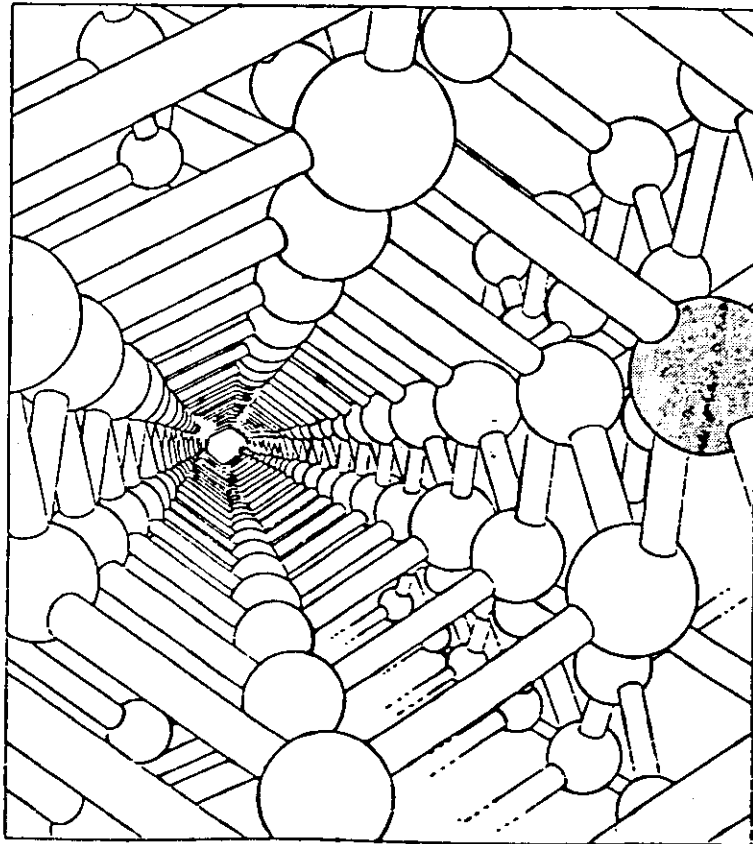
2. REQUIREMENTS ON DENSITY OF STATES FOR EFFICIENT SOLAR CELL.

3. RELATION BETWEEN EXPONENTIAL ABSORPTION EDGE AND BAND GAP.

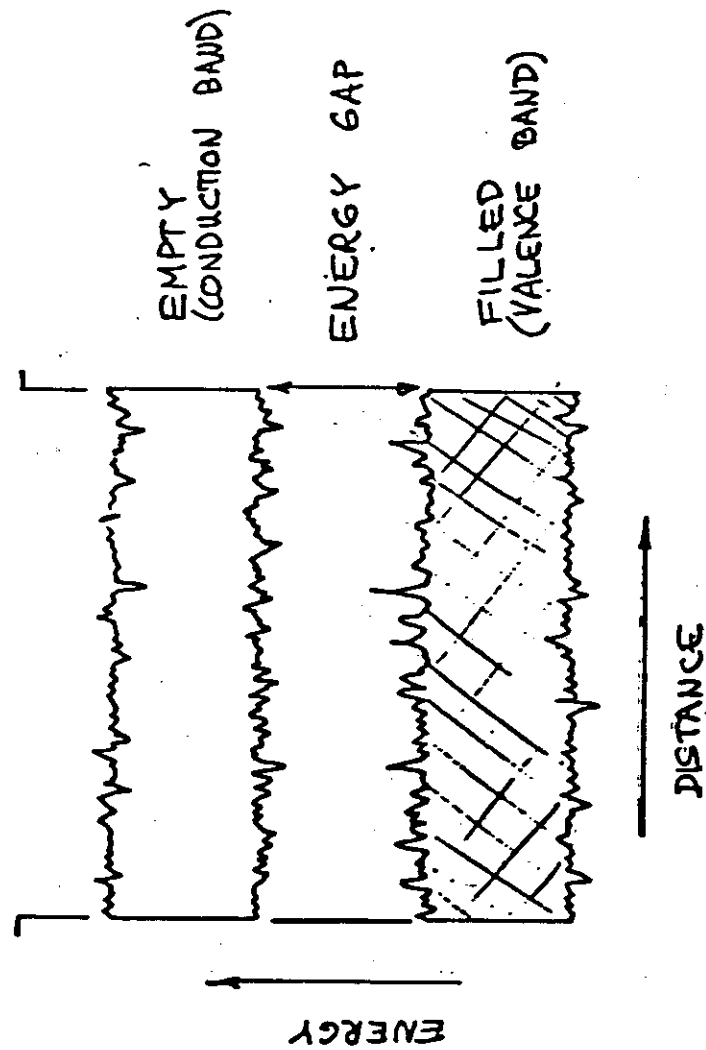
4. FUNDAMENTAL CONCEPTS IN PHOTOCARRIER TRANSPORT - RECOMBINATION OF FREE CARRIERS - CAPTURE CROSS SECTION - QUASIFERMI LEVEL - DEMARCATION LEVEL - RESPONSE TIME, DRIFT MOBILITY AND FREE CARRIER MOBILITY.

AMORPHOUS SILICON HAS NO LONG RANGE ORDER

THE STRUCTURE OF CRYSTALLINE SILICON IS HIGHLY
REGULAR AND SYMMETRIC



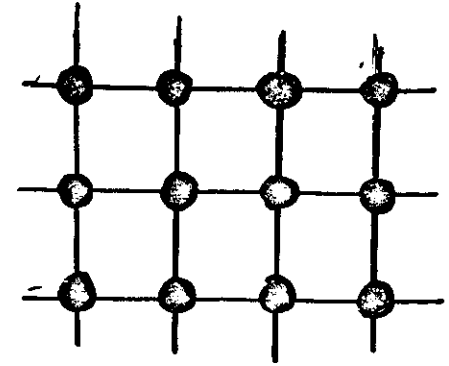
AN AMORPHOUS SEMICONDUCTOR



13

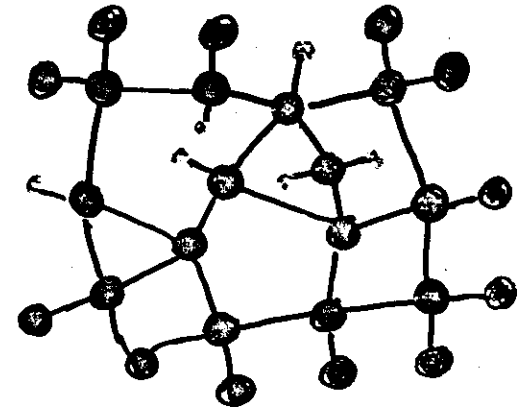
TETRAHEDRAL COORDINATION

$C-Si$

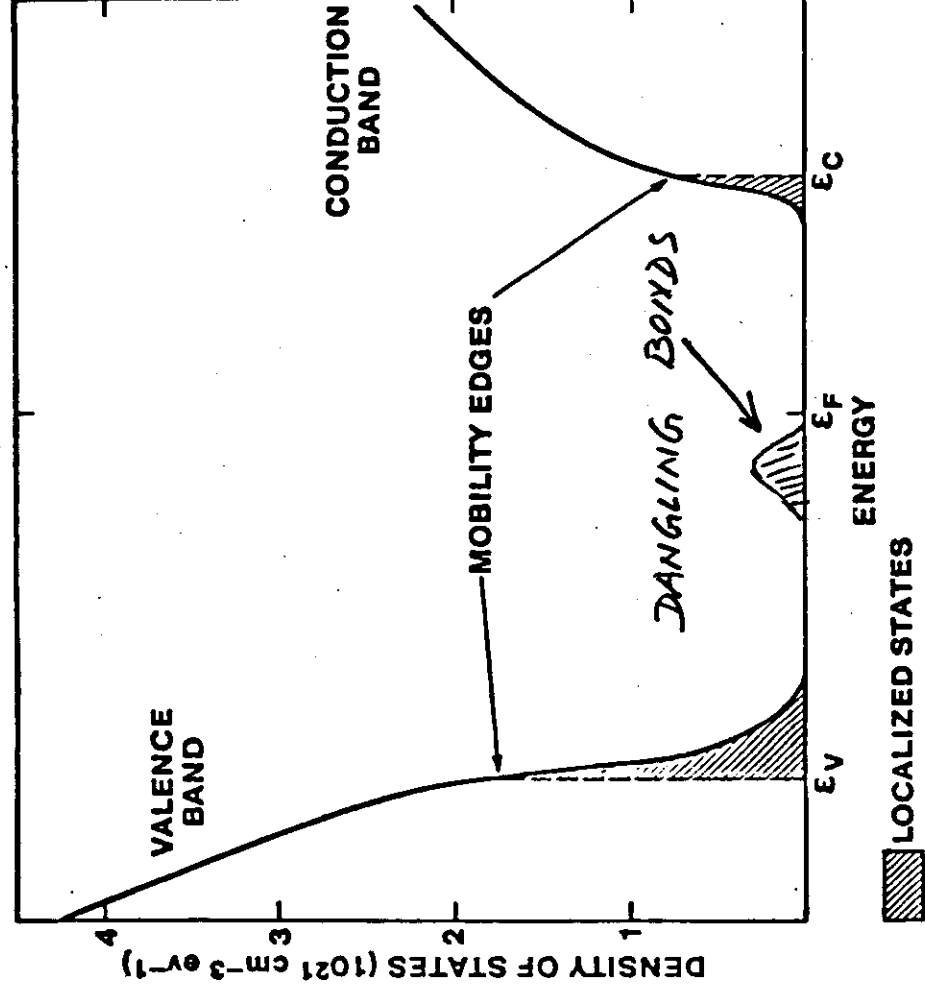
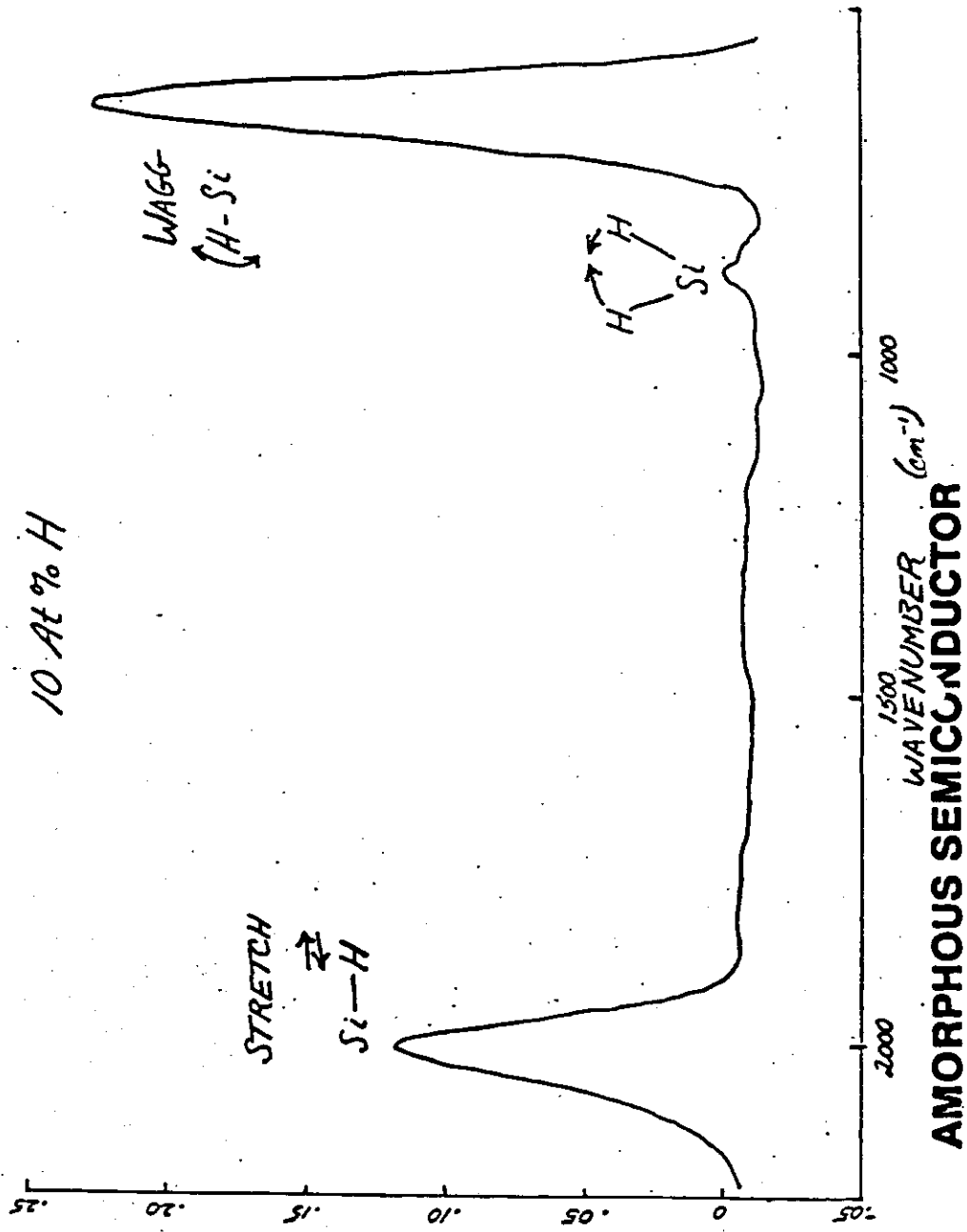


$\alpha-SiH_x$

DANGLING
BONDS ARE
HEALED BY H

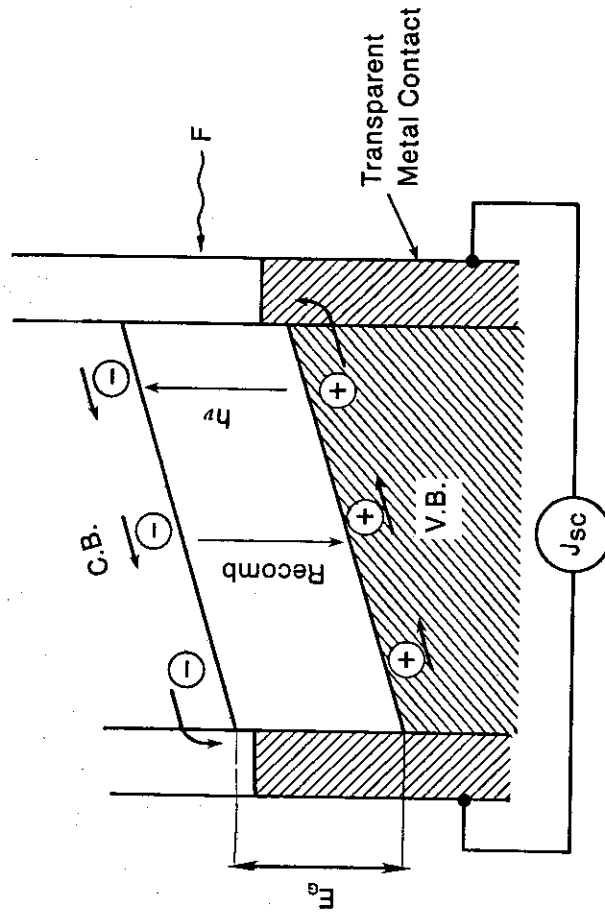


14



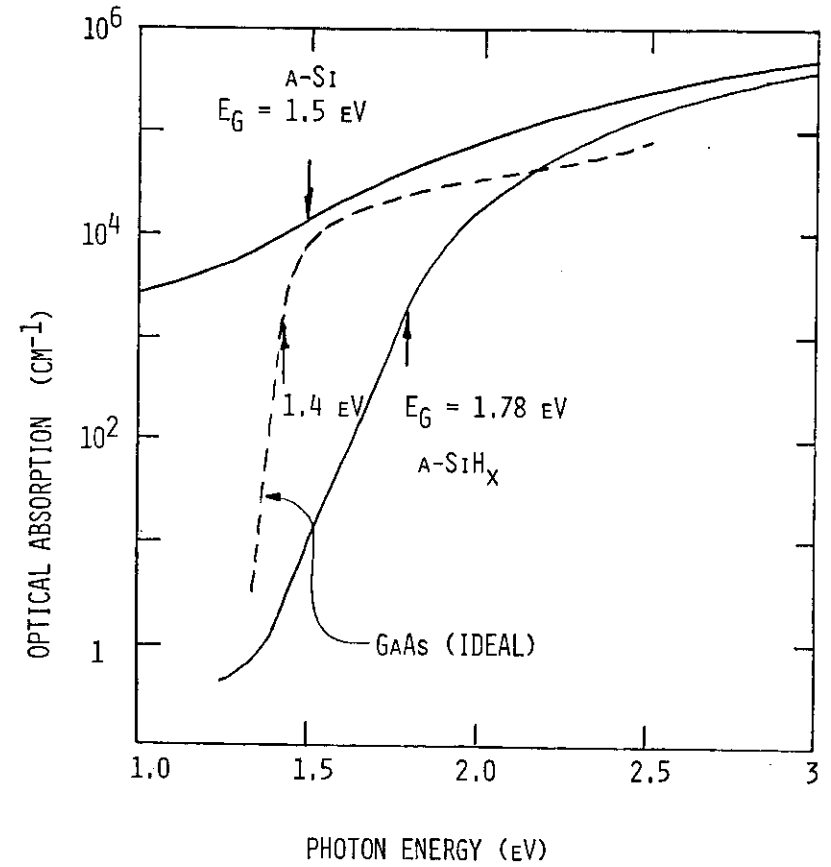
EFFICIENT THIN FILM SOLAR CELL

1. Band gap $E_g = 1.4 \text{ eV}$
2. Strong optical absorption ($d \sim 1 \mu\text{m}$)
3. Photogenerated carrier extraction efficiency $\sim 100\%$



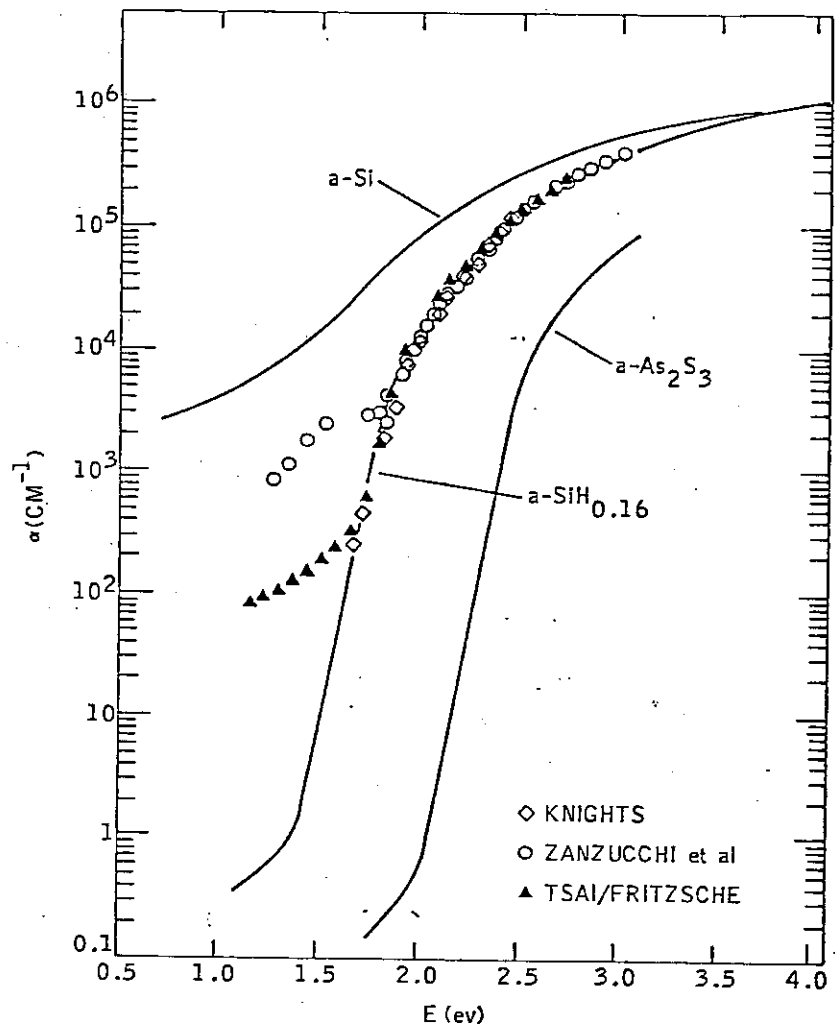
17

THE BANDGAP OF AMORPHOUS SILICON DEPENDS ON HOW THE MATERIAL IS MADE

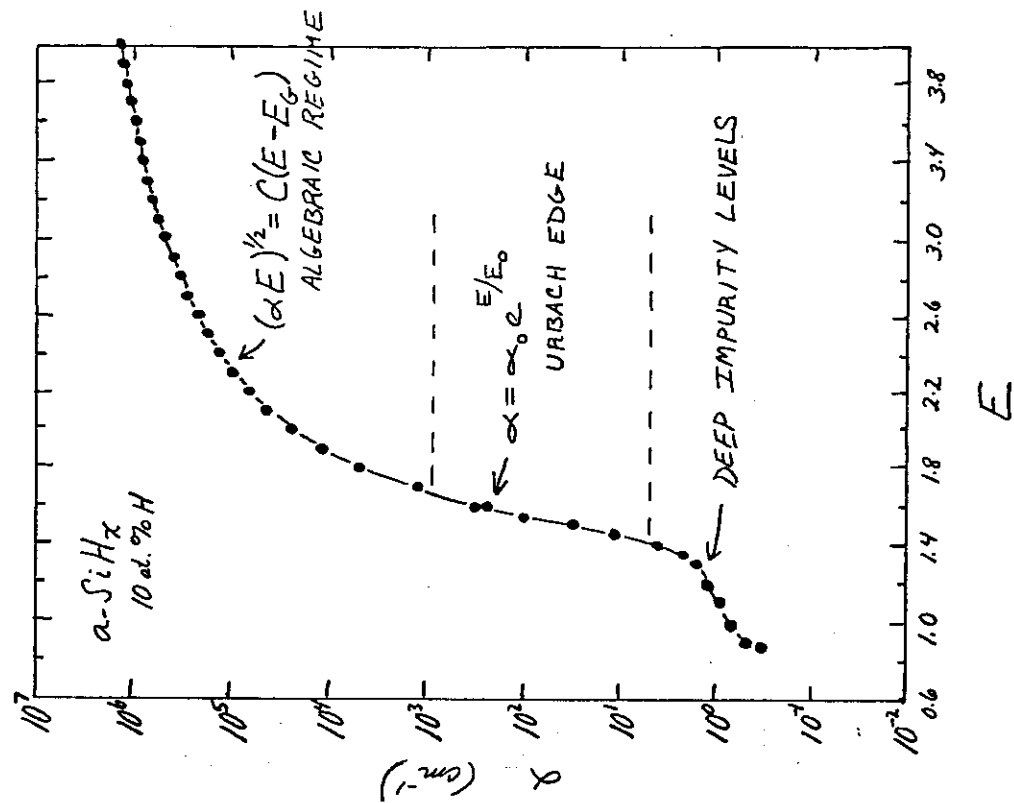


G. D. CODY
B. ABELES
C. R. WRONSKI

18



19



20

OPTICAL ABSORPTION

$$\alpha(E) \propto \int d\epsilon_1 d\epsilon_2 \delta(\epsilon_1 - \epsilon_2 - E)$$

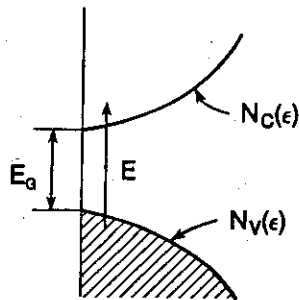
$$\propto N_C(\epsilon_1) N_V(\epsilon_2) M(\epsilon_1, \epsilon_2)$$

$$\text{For } N_C(\epsilon) \propto (\epsilon - E_C)^{1/2}$$

$$N_V(\epsilon) \propto (\epsilon - E_V)^{1/2}$$

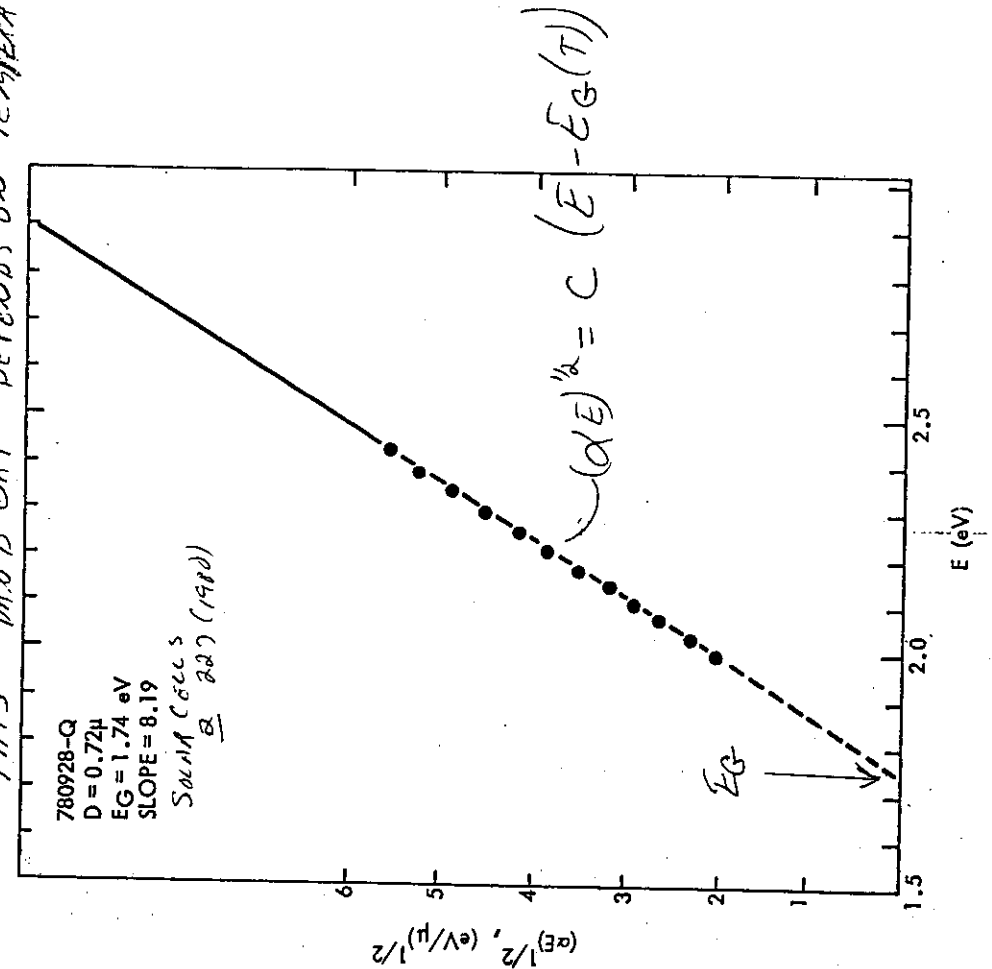
$$M(\epsilon_1, \epsilon_2) = \text{const.}$$

$$\alpha E \propto (E - E_G)^2$$



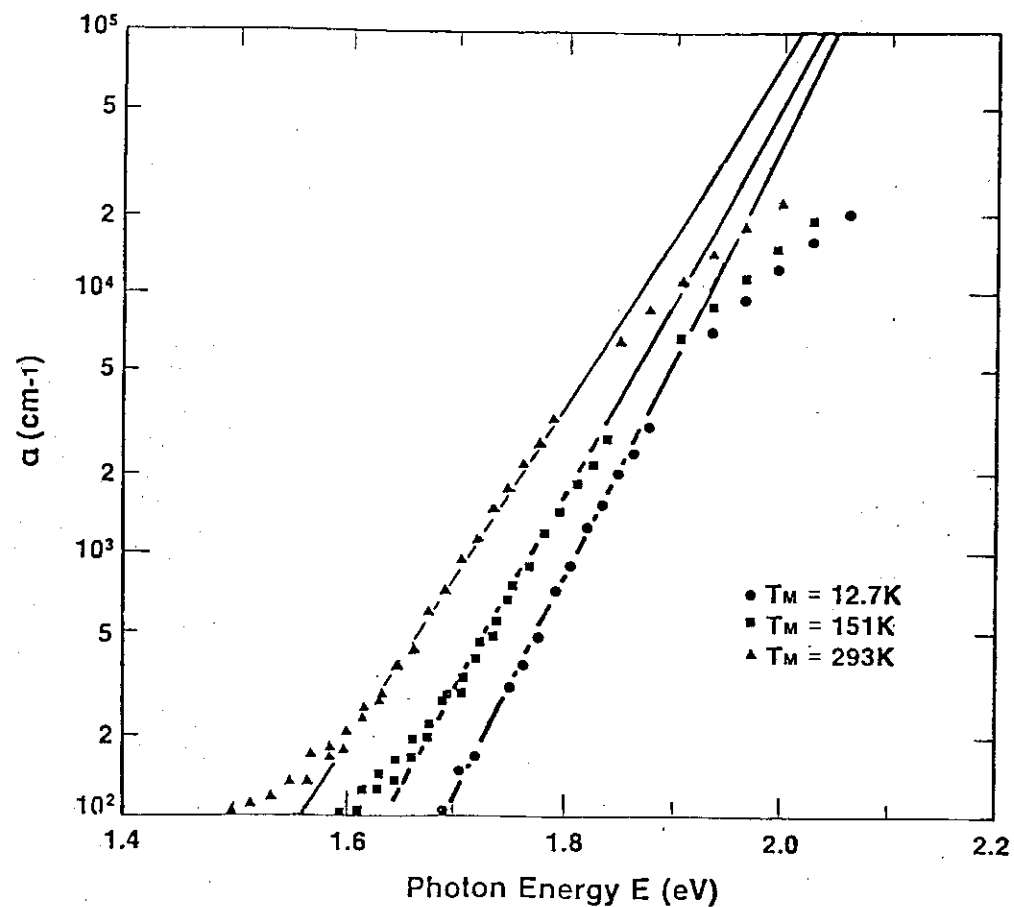
1. J. Tauc, In amorphous and liquid semiconductors, edited by J. Tauc (Plenum, London, 1976) Chapter 6.

WE CAN DEFINE A BAND GAP
THIS BAND GAP DEPENDS ON TEMPERATURE



THERMAL DISORDER BROADENS ABSORPTION EDGE

AND LOWERS BANDGAP



G. D. CODY
T. TIEDJE
B. ABELES

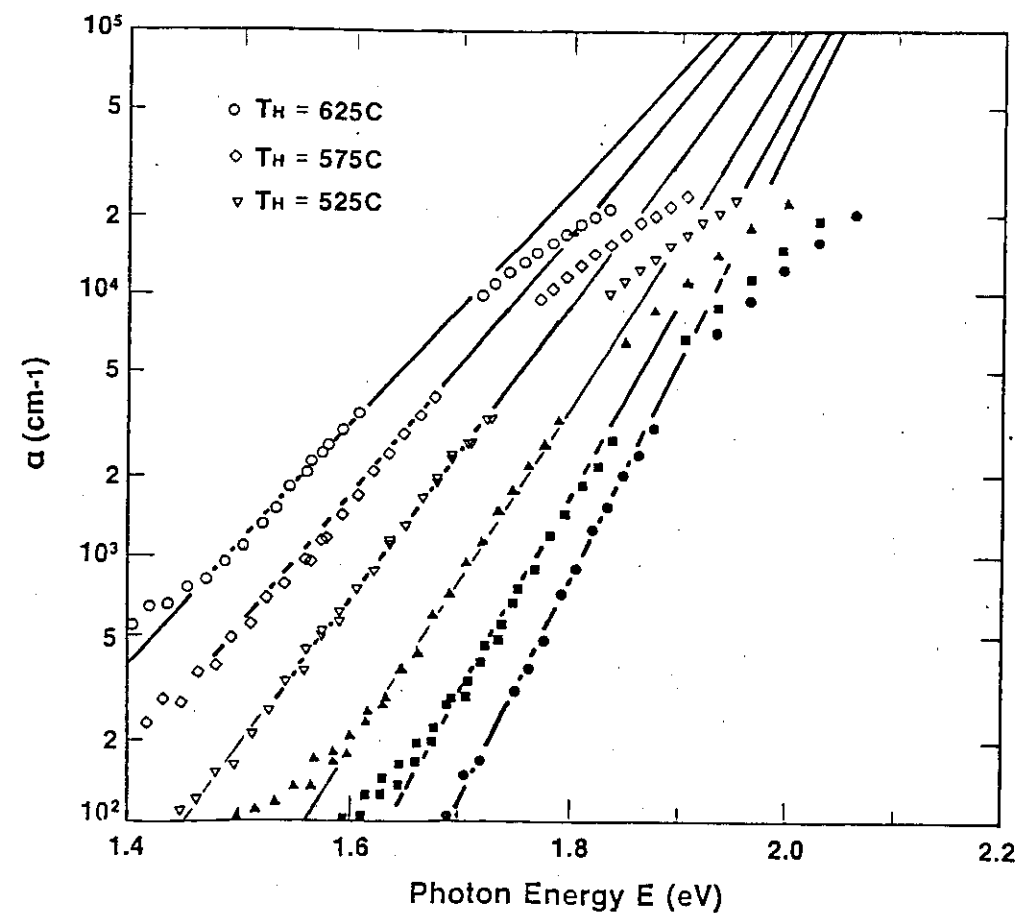
B. BROOKS
Y. GOLDSTEIN

23

STRUCTURAL DISORDER ALSO BROADENS ABSORPTION EDGE

AND LOWERS BAND GAP

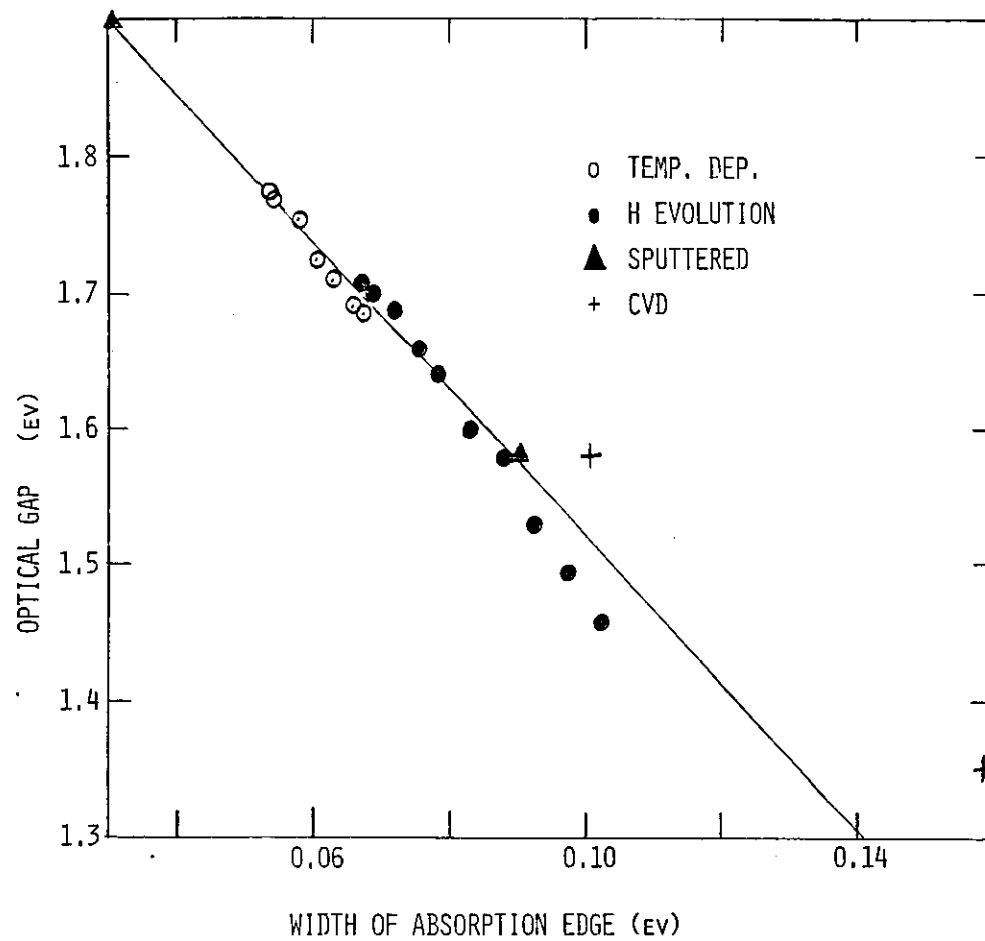
- LATTICE DISORDERED BY HEAT TREATMENTS THAT EVOLVE HYDROGEN



G. D. CODY, T. TIEDJE, B. ABELES, B. BROOKS, Y. GOLDSTEIN

24

IN AMORPHOUS SILICON OPTICAL GAP AND WIDTH OF
ABSORPTION EDGE ARE RELATED



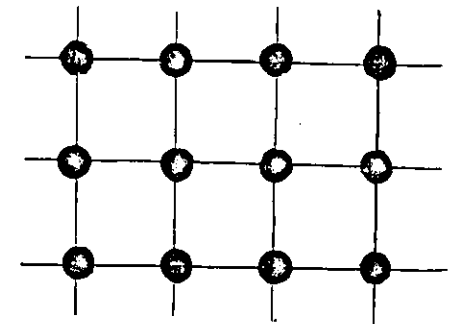
G. D. CODY T. MOUSTAKAS
T. TIEDJE B. BROOKS
B. ABELES Y. GOLDSTEIN

25

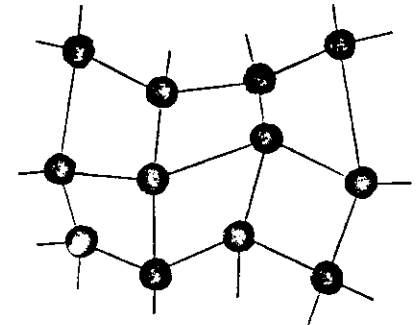
THERMAL AND STRUCTURAL DISORDER

ARE SIMILAR

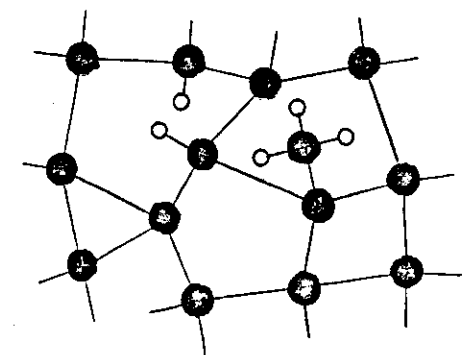
CRYSTAL
AT
0°K



CRYSTAL AT
ROOM
TEMPERATURE
(SNAP SHOT)



AMORPHOUS
MATERIAL



26

PHENOMENOLOGICAL MODEL ¹

EQUIVALENCE OF STRUCTURAL AND THERMAL DISORDER:

$$E_0(x, T) \propto \langle U^2 \rangle_T + \langle U^2 \rangle_x$$

$$E_G(T, x) = E_G(0, 0) - D \langle U_0^2 \rangle \left(\frac{E_0(T, x)}{E_0(0, 0)} - 1 \right)$$

$$x = \frac{\langle U^2 \rangle_x}{\langle U^2 \rangle_0} \quad \text{DISORDER PARAMETER}$$

$$\langle U^2 \rangle_x = \text{RMS STRUCTURAL DISORDER}$$

$$\langle U^2 \rangle_0 = .08 \text{ \AA RMS OF ZERO PT. MOTION}$$

$$D = 30 \text{ eV/\AA}^2 \text{ SECOND O DEF. POT.}$$

$$E_G(0, 0) = 2.0 \text{ eV}$$

¹ CODY ET AL., PHYS. REV. LETT. 47, 1480 (1981).

COHERENT POTENTIAL APPROXIMATION

(ABE AND TOYOZAWA, J. PHYS. SOC. JAP 50, 2185 (1981))

- DISORDER CHARACTERIZED BY GAUSSIAN DISTRIBUTION
WITH RMS OF RANDOM FLUCTUATIONS - W

- DENSITY OF STATES $\frac{1}{\pi B_u^2} \sqrt{(B_u^2 - \epsilon^2)}$
($B_{v,c}$ BAND WIDTH)

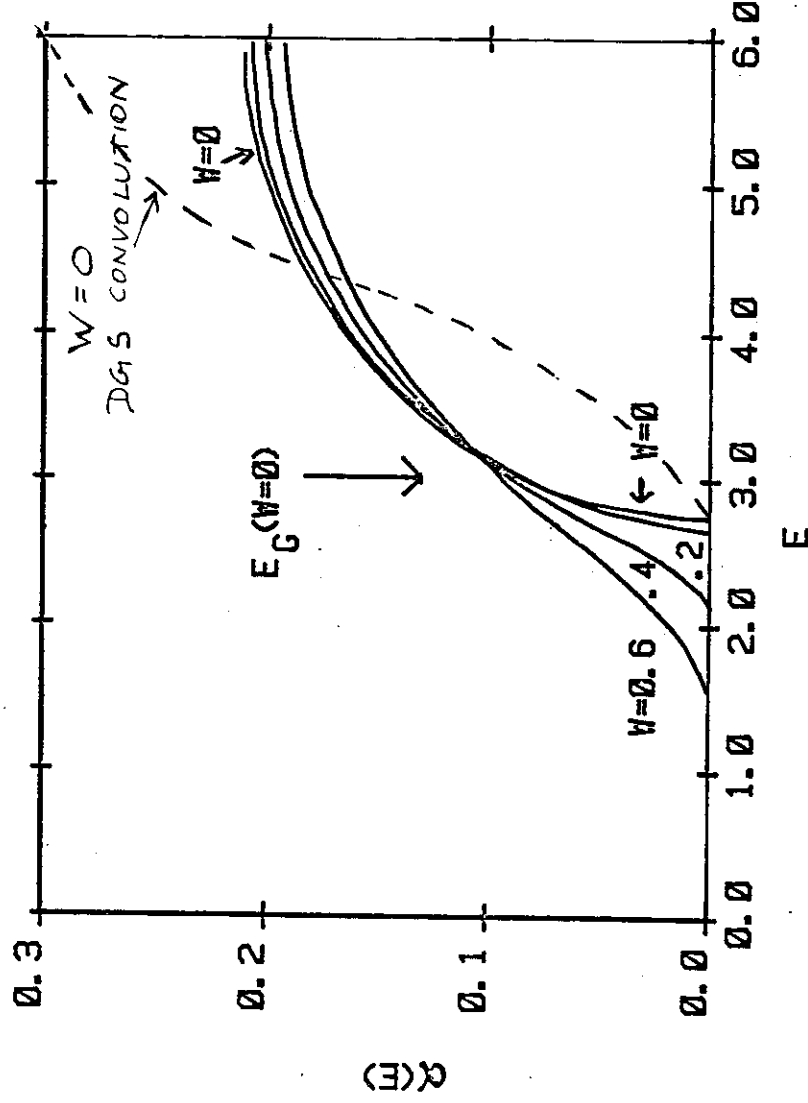
MATRIX ELEMENT - κ RELAXATION AND DISORDER INDUCED
SPATIAL CORRELATION

- THEORY: $E_0 = 0.5 \frac{W}{B}$
 $\alpha \tilde{E} \propto (E - E_G)^2$
 $E_G = 3 - 10 E_0$
- EXPERIMENT:
 $E_G = 2.1 - 6.2 E_0$

THE ONLY FEATURE IN THE THEORY IS
AT ZERO DISORDER ($W=0$)

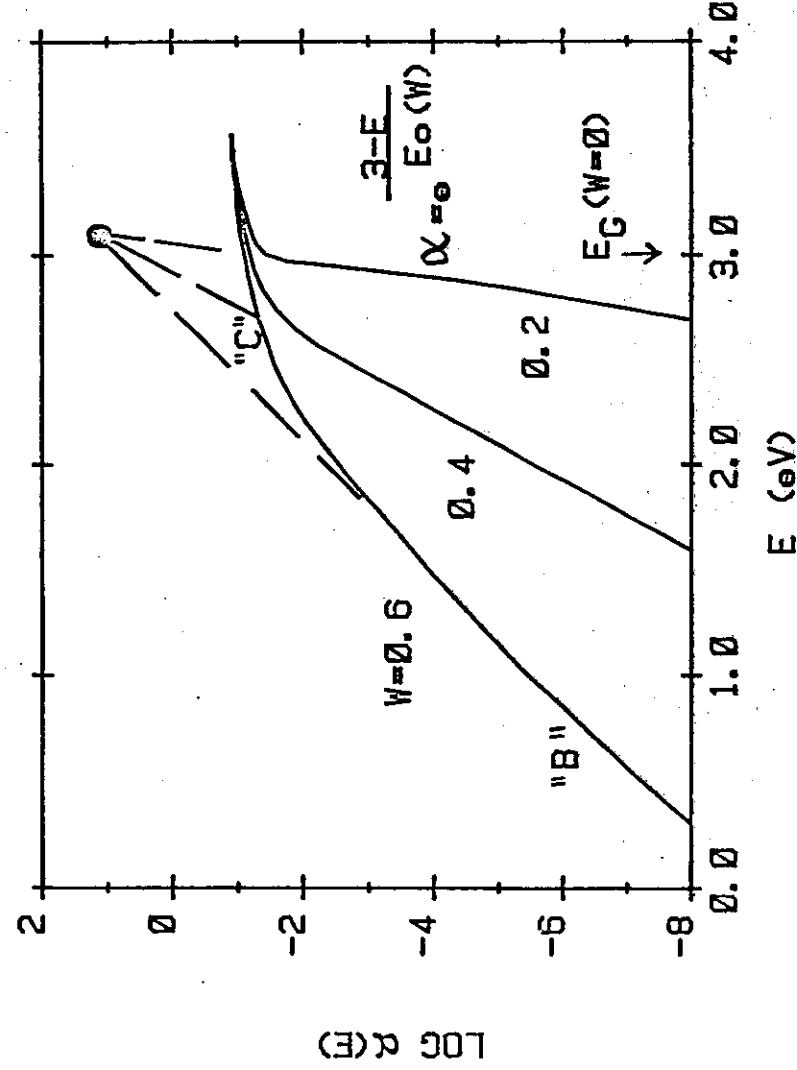
Dirac

ABE & TOYOZAWA
JPS JAPAN 50 2185 (1981)



29

THERE IS AN URBACH REGION AND $E_G(W=0)$ IS THE FOCUS

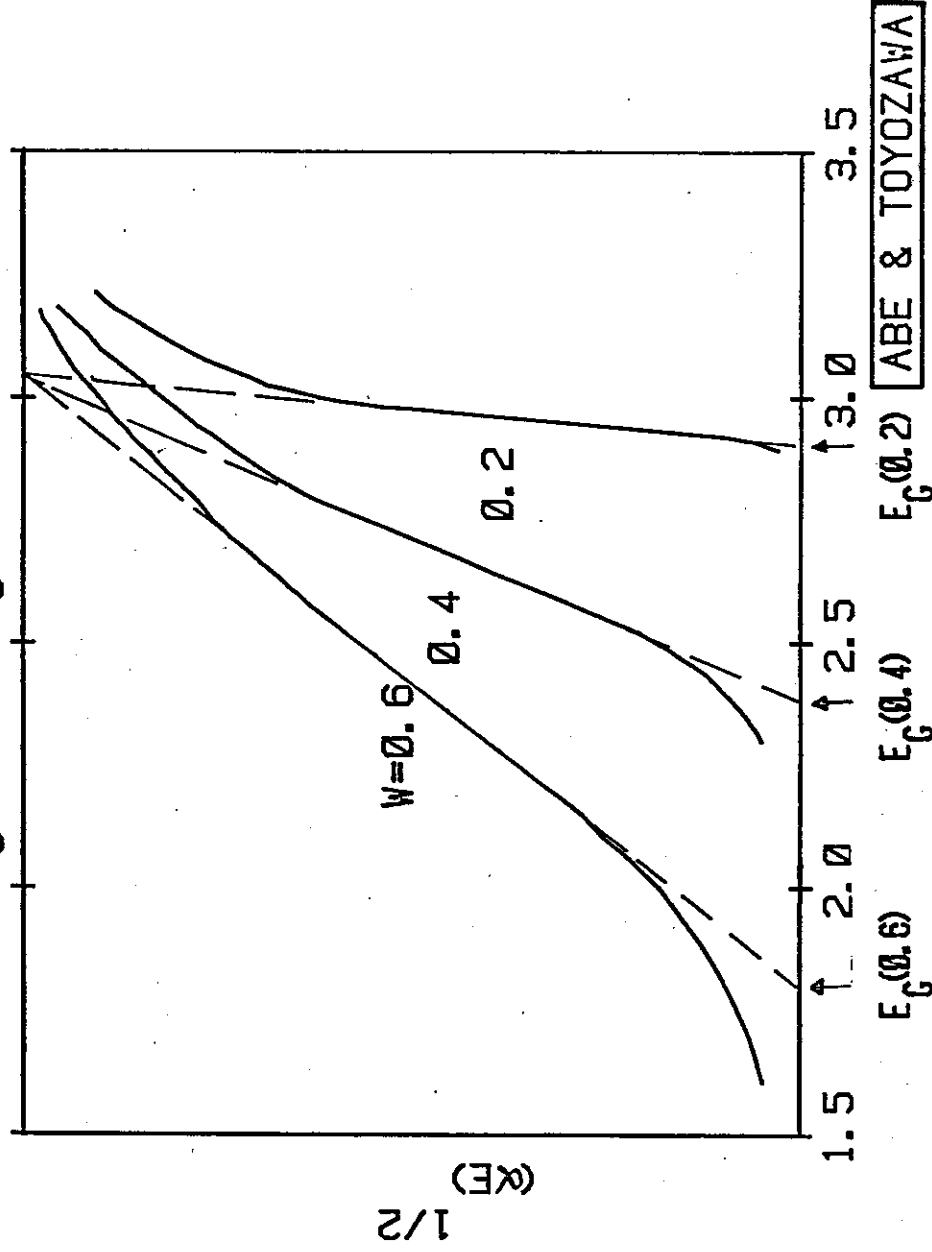


30

THERE IS A PARABOLIC REGION, AND IN THAT REGION

$E_G(w)$ CAN BE DEFINED AND

$$E_G(w) \approx 3 - 10E_0(w)$$



SUMMARY

- EDGE DEFINED BY E_G AND E_0 , ~~NOT THESE~~ ARE LINEARLY RELATED AND BOTH DOMINATED BY STRUCTURAL AND THERMAL DISORDER
- RECENT THEORY SUGGESTS THAT EDGE IS THE BROADENING OF A DIRECT GAP OF 2.1-2.2 EV FOR $\alpha\text{-Si:H}_x$ (ABE/TOYOZAWA)
- SHARPENING THE EDGE IMPROVES $\alpha\text{-Si:H}_x$ "ELECTRICALLY." IT ALSO REDUCES THE COUPLING TO THE SOLAR SPECTRUM

CONCEPTS IN PHOTOCONDUCTIVITY AND ALLIED PROBLEMS

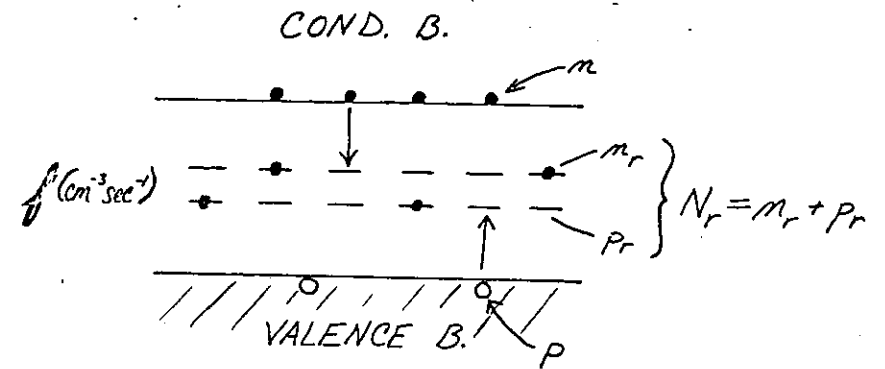
"ALL YOU EVER NEED TO KNOW
ABOUT PHOTOCONDUCTIVITY

ALBERT ROSE

RCA Laboratories
Princeton, New Jersey

1978
ROBERT E. KRIEGER PUBLISHING CO.
Huntington, New York

SINGLE SET OF RECOMB. C., LOW LIGHT $n, p \ll$



$$n = f \tau_n \quad p = f \tau_p$$

$n (p)$ = Free electron (hole) density

$n_r (p_r)$ = recombination c. occupied (unoccupied) by el.

$\tau_n (\tau_p)$ = Free electron (hole) lifetime

f = photogeneration rate

$$\tau_n^{-1} = K_n p_r$$

$$\tau_p^{-1} = K_p n_r$$

$K_n (K_p)$ capture probabilities

RECOMBINATION

EXAMPLE: HIGH ILLUMINATION

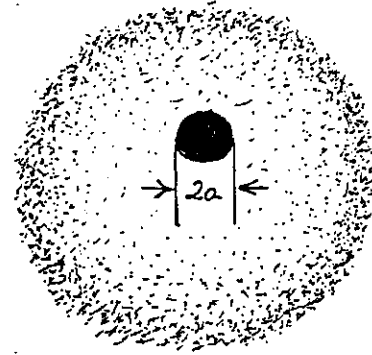
$$1 \text{ SUN} \sim 10^{17} \text{ photons/cm}^2 \text{ sec}$$

FILM $1 \mu\text{m}$

$$\phi = \frac{10^{17}}{10^{-4}} = 10^{21} \text{ photons/cm}^3 \text{ sec}^{-1}$$

$$\left. \begin{array}{l} \tau_n = 10^{-6} \text{ sec} \\ n = 10^{15} \text{ cm}^{-3} \\ \tau_p = 10^{-8} \text{ sec} \\ p = 10^{11} \text{ sec} \end{array} \right\} \begin{array}{l} \text{TYPICAL} \\ \text{a-SiH}_x \\ \text{SOLAR CELL} \end{array}$$

A. DIFFUSIVE: ($l < a$)



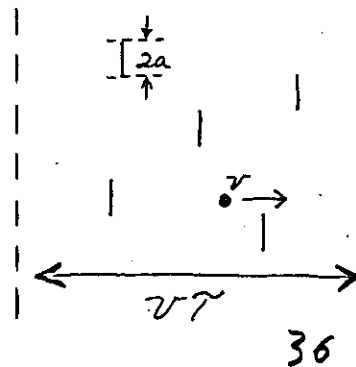
$$n(x) = \left(1 - \frac{a}{x}\right) n_0$$

$$\frac{dn_0}{dt} = -n_r 4\pi a^2 D \left. \frac{dn}{dx} \right|_{r=a}$$

$$\frac{dn_0}{dt} = -4\pi n_r a D n_0$$

$$\frac{1}{\tau} = 4\pi n_r a D \quad \text{or} \quad K = 4\pi a D$$

B. BALISTIC: ($l > a$)



$$n_r v \tau \pi a^2 = 1$$

$$\frac{1}{\tau} = n_r v \pi a^2$$

$$K = v \pi a^2$$

CAPTURE CROSS SECTION

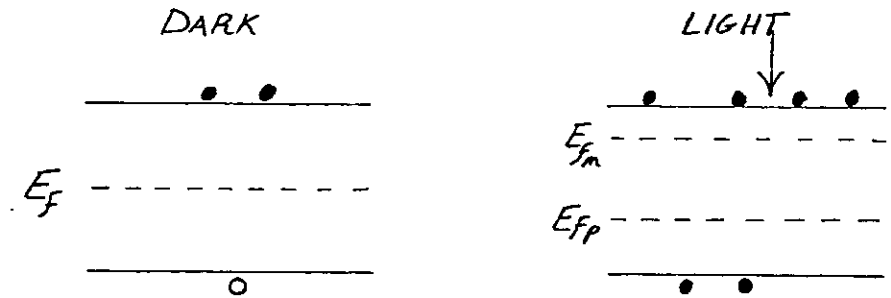
$$\frac{a_{bal}}{a_{diff}} = \sqrt{\left(\frac{4 \mu kT}{e a_{diff} v}\right)}$$

$$\begin{aligned} a_{diff} &= 1(\text{\AA}) \\ \mu &= 1 \text{ cm}^2/\text{V s} \\ v &= 10^7 \text{ cm/s} \\ T &= 300^\circ\text{K} \end{aligned}$$

Center	a_{diff} (Å)
Neutral (at dilms.)	1
Coulomb attractive ($kT = \frac{e^2}{\epsilon r}$)	50
Coulomb repulsive	10^{-3}

37

QUASI-FERMI LEVELS



$$n = N_c e^{-\frac{(E - E_{Fn})}{kT}}$$

$$p = N_v e^{-\frac{(E - E_{Fp})}{kT}}$$

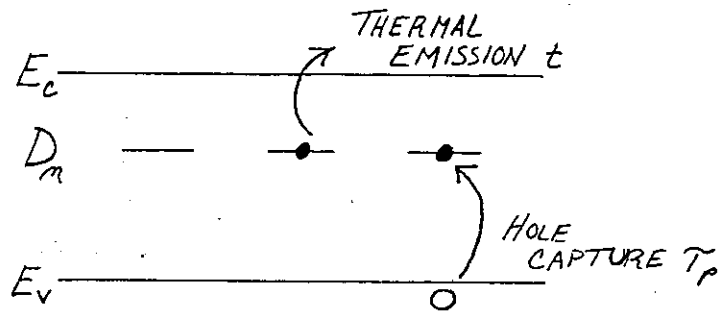
$$N_c = 2 \left(\frac{2\pi m^* kT}{h^2} \right)^{3/2} = 2 \times 10^{19} \text{ cm}^{-3}$$

$$(T = 300^\circ\text{K}, m^* = m)$$

majority carrier n if $n > p$

38

DEMARICATION LEVELS



AT DEMARICATION LEVEL $t = \tau_p$

$$t = \omega^{-1} e^{\frac{E_c - D_m}{kT}}$$

$$\tau_p = \frac{1}{\nu_s p}$$

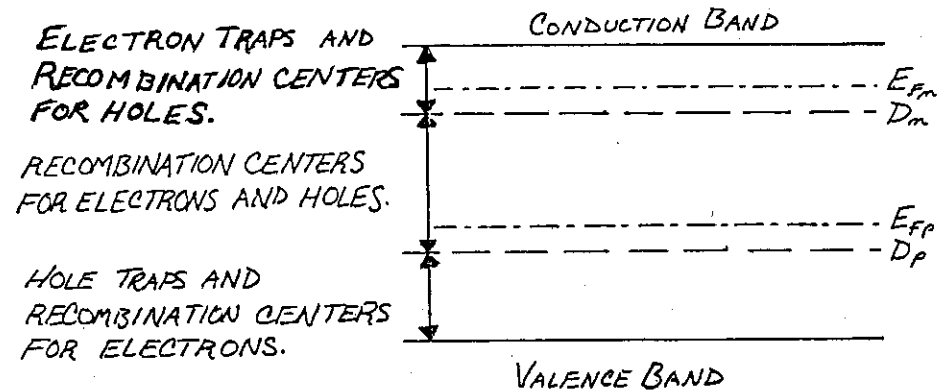
$\omega = N_c \nu_s$ attempt frequency
(by detailed balance)

$$\left. \begin{aligned} e^{\frac{-E_c - D_m}{kT}} &= \frac{\omega}{\nu_s p} \\ e^{\frac{-E_c - E_{fm}}{kT}} &= \frac{n}{N_c} \end{aligned} \right\}$$

$$E_c - D_m = E_c - E_{fm} + kT \ln \frac{n s_n}{p s_p}$$

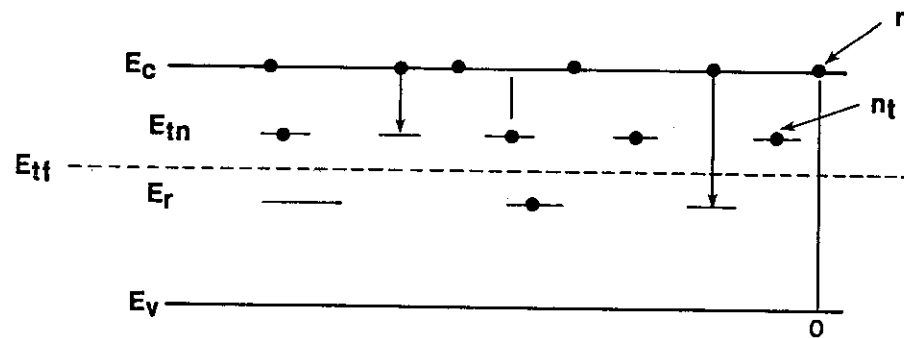
- BELOW D_m STATES ACT AS RECOMBINATION CENTERS FOR HOLES.

DEMARICATION LEVELS



$$E_{fm} - D_m = E_{fp} - D_p = kT \ln \frac{n s_n}{p s_p}$$

TRAPPING (Single Level)



$$\frac{n_t}{N_t} = \frac{n}{N_c} e^{\frac{-(E_c - E_{tn})}{kT}}$$

$$\text{response time } \tau_{on} = (1 + \frac{n_t}{n}) \tau_n$$

$$\text{drift mobility } \mu_{nD} = \mu_n \frac{n}{n_t + n}$$

$$\text{Free carrier mobility} = \mu_n$$

