



INTERNATIONAL ATOMIC ENERGY AGENCY
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INTERNATIONAL CENTRE FOR THEORETICAL PHYSICS
34100 TRIESTE (ITALY) - P.O.B. 586 - MIRAMARE - STRADA COSTIERA 11 - TELEPHONES: 224281/23456
CABLE: CENTRATOM - TELEX 460392-I

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SPRING COLLEGE ON AMORPHOUS SOLIDS
AND THE LIQUID STATE

14 April - 18 June 1982

SMALL POLARON CONDUCTION

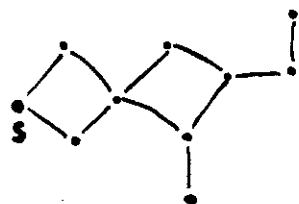
(Lectures I & II)

H. REIK

Fakultät für Physik
Albert-Ludwigs-Universität
Hermann-Herder-Strasse 3
7800 Freiburg
Federal Republic of Germany

These are preliminary lecture notes, intended only for distribution to participants.
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$$H_c = \sum_s \epsilon_s c_s^\dagger c_s + \sum_{s,i} t_{s,i} (c_s^\dagger c_{s+i} + c_{s+i}^\dagger c_s)$$



$$H_{ph} = \sum_{q,\lambda} \hbar \omega_{q\lambda} b_{q\lambda}^\dagger b_{q\lambda}$$

$$H_{cp} = \sum_{q,\lambda,s} \hbar \omega_{q\lambda} \alpha_{\lambda q}^s i (b_{q\lambda} e^{i\vec{q} \cdot \vec{R}_s} - b_{q\lambda}^\dagger e^{-i\vec{q} \cdot \vec{R}}) c_s^\dagger c_s$$

$$\alpha_{ac}(q) = \frac{\Xi_a^2 |q|^2}{2\bar{\gamma} V \omega_{ac}(q)}$$

$\bar{\gamma}$ = mass density

$$\alpha_{opt}(q) = \frac{A^2}{2\bar{\gamma} V \omega_{op}^2(q)}$$

$\bar{\gamma}$ = reduced mass density.

1. Complete squares in Hamiltonian:

$$H = \sum_s \epsilon_s c_s^\dagger c_s + \sum_{s,i} t_{s,i} (c_s^\dagger c_{s+i} + c_{s+i}^\dagger c_s)$$

$$+ \sum_{q\lambda} \hbar \omega_{q\lambda} [b_{q\lambda}^\dagger + i \alpha_{\lambda q}^s \sum_s e^{i\vec{q} \cdot \vec{R}_s} c_s^\dagger c_s] [b_{q\lambda} - i \alpha_{\lambda q}^s \sum_{s'} e^{-i\vec{q} \cdot \vec{R}_{s'}} c_{s'}^\dagger c_{s'}]$$

$$- \sum_{q\lambda} \hbar \omega_{q\lambda} \alpha_{\lambda q} \sum_{s,s'} e^{i\vec{q} \cdot (\vec{R}_s - \vec{R}_{s'})} c_s^\dagger c_s c_{s'}^\dagger c_{s'}$$

$$B_{q\lambda}^\dagger = b_{q\lambda}^\dagger + i \alpha_{\lambda q}^s \sum_s e^{i\vec{q} \cdot \vec{R}_s} c_s^\dagger c_s$$

$$b_s^\dagger = c_s^\dagger \prod_{q\lambda} e^{i \alpha_{\lambda q}^s [e^{-i\vec{q} \cdot \vec{R}_s} b_{q\lambda}^\dagger + e^{i\vec{q} \cdot \vec{R}_s} b_{q\lambda}]}$$

$$b_s = c_s \prod_{q\lambda} e^{-i \alpha_{\lambda q}^s [e^{i\vec{q} \cdot \vec{R}_s} b_{q\lambda} + e^{-i\vec{q} \cdot \vec{R}_s} b_{q\lambda}^\dagger]}$$

Displaced phonons and polarons.

$$c_s^\dagger = b_s^\dagger \prod_{q\lambda} e^{-i \alpha_{\lambda q}^s [e^{-i\vec{q} \cdot \vec{R}_s} b_{q\lambda}^\dagger + e^{i\vec{q} \cdot \vec{R}_s} b_{q\lambda}]}$$

$$[b_s^\dagger, b_{s'}]_+ = \delta_{s,s'}$$

$$[b_s^\dagger, B_{q\lambda}]_- = 0$$

$$H = \varepsilon_s l_s^\dagger l_s + \sum_{s,i} t_{s,i} (l_s^\dagger l_{s+i} \chi_{s,s+i} + l_{s+i}^\dagger l_s \chi_{s+i,s})$$

$$+ \sum_{q,\lambda} \hbar \omega_{q,\lambda} B_{q,\lambda}^\dagger B_{q,\lambda}$$

$$- \sum_{\lambda,q} \hbar \omega_{q,\lambda} \alpha_{\lambda,q} \sum_s l_s^\dagger l_s$$

$$- \sum_{\lambda,q} \hbar \omega_{q,\lambda} \alpha_{\lambda,q} \sum_{s,s'} e^{i\vec{q}(\vec{R}_s - \vec{R}_{s'})} l_s^\dagger l_s l_{s'}^\dagger l_{s'}$$

$$\chi_{s,s+i} = \int_{q,\lambda} e^{-i\alpha_{\lambda,q}} [e^{-i\vec{q}\vec{R}_s} - e^{-i\vec{q}\vec{R}_{s+i}}] B_{q,\lambda}^\dagger + h.c.$$

Current operator

$$\vec{r} = \sum_s R_s c_s^\dagger c_s = \sum_s R_s l_s^\dagger l_s$$

$$i\hbar \dot{\vec{r}} = \sum_{s,i} t_{s,i} (\vec{R}_{s+i} - \vec{R}_s) (c_{s+i}^\dagger c_s - c_s^\dagger c_{s+i})$$

$$\begin{aligned} \vec{J} &= -\frac{e}{i\hbar} \sum_{s,i} t_{s,i} (\vec{R}_{s+i} - \vec{R}_s) (c_{s+i}^\dagger c_s - c_s^\dagger c_{s+i}) \\ &= -\frac{e}{i\hbar} \sum_{s,i} t_{s,i} (\vec{R}_{s+i} - \vec{R}_s) (l_{s+i}^\dagger l_s \chi_{s+i,s} - l_s^\dagger l_{s+i} \chi_{s,s+i}) \end{aligned}$$

Current correlation function (App)

$$\begin{aligned} \langle \vec{j}(t) \vec{j}(0) \rangle &= \sum_{s,i} \langle j(t) j(0) \rangle^{s,i} \\ &= \frac{e^2}{\hbar^2} \sum_{s,i} t_{s,i}^2 (\vec{R}_{s+i} - \vec{R}_s) (\vec{R}_{s+i} - \vec{R}_s) \\ &\quad \langle l_s^\dagger(t) l_{s+i}(t) \chi_{s,s+i}(t) l_{s+i}^\dagger l_s \chi_{s+i,s} \rangle \end{aligned}$$

$$= \frac{e^2}{\hbar^2} \sum_{s,i} t_{s,i}^2 (\vec{R}_{s+i} - \vec{R}_s) (\vec{R}_{s+i} - \vec{R}_s)$$

$$\langle l_s^\dagger(t) l_{s+i}(t) l_{s+i}^\dagger l_s \rangle$$

$$\underbrace{\langle \chi_{s+i,s}(t) \chi_{s,s+i} \rangle - \langle \chi_{s+i,s} \chi_{s,s+i} \rangle}_{P_{s,i}(t)} + \underbrace{\langle \chi_{s+i,s} \chi_{s,s+i} \rangle}_{Q_{s,i}}$$

$$\langle l_s^+(z) l_{s+i}(z) l_{s+i}^+ l_s \rangle$$

$$= e^{\frac{i}{\hbar} (z_s - z_{s+i}) z} \eta_e^{(s)} \eta_h^{(s+i)}$$

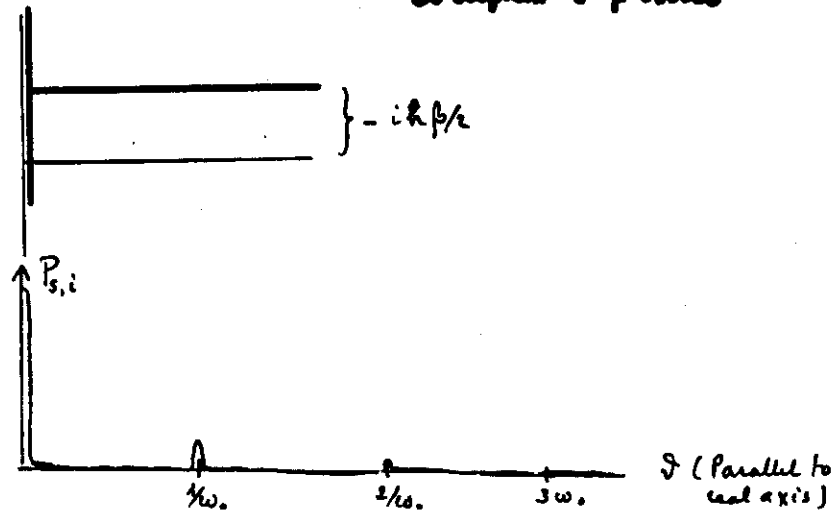
$$= e^{\frac{i}{\hbar} (z_s - z_{s+i}) z} \frac{e^{\beta(z_{s+i} - z_s)/2}}{4 \cosh \beta z_s/2 \cosh \beta z_{s+i}/2}$$

$$Q_{s,i} = e^{-2 \sum \lambda_q \sin^2(\frac{1}{2} \vec{q}(\vec{R}_{s+i} - \vec{R}_s)) \coth \hbar \omega_{\lambda q} \beta/2}$$

$$P_{s,i}(z) = Q_{s,i} \times$$

$$\left\{ e^{2 \sum \lambda_q \sin^2(\frac{1}{2} \vec{q}(\vec{R}_{s+i} - \vec{R}_s)) \frac{\cosh \omega_{\lambda q} (t + i\hbar \beta/2)}{\sinh \hbar \omega_{\lambda q} \beta/2} - 1} \right\}$$

complex z plane



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$$P^{(s,i)}(z) = Q^{(s,i)}$$

$$\exp \left\{ \sum_{q\lambda} \frac{2\lambda_q \sin^2 \frac{1}{2} \vec{q}(\vec{R}_{s+i} - \vec{R}_s)}{\sinh \hbar \omega_{\lambda q} \beta/2} \coth \{ \omega_{\lambda q} [Im z + \frac{1}{2} \hbar \beta] \} \right.$$

$$\times \exp \left\{ -i\hbar \sum \frac{2\lambda_q \sin^2 \frac{1}{2} \vec{q}(\vec{R}_{s+i} - \vec{R}_s)}{\sinh \hbar \omega_{\lambda q} \beta/2} \frac{\omega_{\lambda q}}{\sinh \{ \omega_{\lambda q} [Im z + \hbar \beta] \}} \right.$$

$$\times \exp \left\{ -\hbar^2 \sum \frac{\lambda_q \sin^2 \frac{1}{2} \vec{q}(\vec{R}_{s+i} - \vec{R}_s)}{\sinh \hbar \omega_{\lambda q} \beta/2} \omega_{\lambda q}^2 \coth \{ \omega_{\lambda q} [Im z + \hbar \beta] \} \right.$$

$$\text{Structure } \eta_{s,i} = 2 \sum_{q\lambda} \lambda_q \sin^2 \frac{1}{2} \vec{q}(\vec{R}_{s+i} - \vec{R}_s)$$

$$\omega_{\lambda q} = \omega_0 \quad (\text{no dispersion})$$

$$P^{s,i} = Q_{s,i} \exp \left\{ \eta^{(s,i)} \frac{1}{\sinh \hbar \omega_0 \beta/2} \coth \omega_0 [Im z + \frac{1}{2} \hbar \beta] \right\}$$

$$\times \exp \left\{ -i\omega_0 \hbar \eta^{(s,i)} \frac{\sinh \omega_0 [Im z + \hbar \beta]}{\sinh \hbar \omega_0 \beta/2} \right\}$$

$$\exp \left\{ -\omega_0^2 \hbar^2 \frac{1}{2} \eta^{(s,i)} \frac{\coth \omega_0 [Im z + \frac{1}{2} \hbar \beta]}{\sinh \hbar \omega_0 \beta/2} \right\}$$

$$\langle l_s^+(z) l_{s+i}(z) l_{s+i}^+ l_s \rangle = e^{\frac{i}{\hbar} (z_s - z_{s+i}) z}$$

$$\frac{e^{-(z_s - z_{s+i}) \frac{1}{\hbar} [Im z + \hbar \beta/2]}}{4 \cosh \beta z_s/2 \cosh \beta z_{s+i}/2}$$

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$$\sigma^{(s,i)}(\omega, T) = \int_{-\infty}^{\infty} dt e^{(i\omega + s)T} \int_{-\infty}^{\infty} d\lambda \langle \underline{j}(T - i\lambda) \underline{j}(0) \rangle^{s,i}$$

Complex τ plane

$$(1) \sigma^{(s,i)}(\omega, T) = 0$$

$$(2) \sigma^{(s,i)(2)}(\omega, T) = \iint_{-\infty}^{\infty} d\lambda d\tau e^{i\omega\lambda} \langle \underline{j}(\tau - i\lambda/2(1-\psi_{si})) \underline{j}(0) \rangle^{s,i}$$

Carry out the λ Integration

$$\sigma^{(s,i)(2)}(\omega, T) = \beta \frac{\sinh \omega \beta \hbar}{\omega \beta \hbar} e^{-\frac{1}{2} \omega \beta \hbar \psi_{si}} \int_{-\infty}^{\infty} d\tau e^{(i\omega - \frac{1}{2})\tau} \langle \underline{j}(\tau - i\beta \hbar (1-\psi_{si})) \underline{j}(0) \rangle^{s,i}$$

$$(3) \sigma^{(s,i)(2)}(\omega, T) = -i\beta \int_{-\infty}^{\infty} d\lambda e^{-\lambda \omega} \int_{-\infty}^{\infty} d\tau e^{i\omega \tau} \langle \underline{j}(-i\lambda \tau) \underline{j}(0) \rangle^{s,i}$$

$\frac{1}{2}\beta(1-\psi_{si})$ purely imaginary.

Insert in $\sigma^{(s,i)(2)}$

$$\sigma^{(s,i)(2)}(\omega, T) = \beta \frac{\sinh \omega \beta \hbar}{\omega \beta \hbar} e^{-\frac{1}{2} \omega \beta \hbar \psi_{si}} (\omega + [z_s - z_{s+i}] \frac{1}{\hbar})$$

$$\frac{1}{4 \cosh \beta z_s / 2 \cosh \beta z_{s+i} / 2}$$

$$-2 \sum \lambda_{sq} \sin^2 \frac{1}{2} \vec{q} (\vec{R}_{s+i} - \vec{R}_s) [\cot \hbar \omega \lambda_{sq} \beta \hbar - \frac{\cosh \hbar \omega \lambda_{sq} \beta \hbar \psi_{si}}{\sinh \hbar \omega \lambda_{sq} \beta \hbar}]$$

$$\int_{-\infty}^{\infty} d\tau e^{i\omega \tau + \frac{i}{\hbar} (z_s - z_{s+i})\tau} = i\hbar \sum \frac{2 \lambda_{sq} \sin^2 \frac{1}{2} \vec{q} (\vec{R}_{s+i} - \vec{R}_s)}{\sinh \hbar \omega \lambda_{sq} \beta \hbar} \omega \lambda_{sq} \sinh \hbar \omega \lambda_{sq} \beta \hbar \psi_{si} / 2$$

$$e^{-\frac{1}{2} \omega \beta \hbar \psi_{si}} \sum \frac{\lambda_{sq} \sin^2 \frac{1}{2} \vec{q} (\vec{R}_{s+i} - \vec{R}_s)}{\sinh \hbar \omega \lambda_{sq} \beta \hbar} \omega \lambda_{sq} \cosh \hbar \omega \lambda_{sq} \beta \hbar \psi_{si} / 2 \cdot \frac{e^{\frac{1}{2} \omega \beta \hbar \psi_{si}}}{\hbar^2} t_{s,i}^2 (\vec{R}_{s+i} - \vec{R}_s)$$

$$\omega + (z_s - z_{s+i}) \frac{1}{\hbar} = \sum \frac{2 \lambda_{sq} \sin^2 \frac{1}{2} \vec{q} (\vec{R}_{s+i} - \vec{R}_s)}{\sinh \hbar \omega \lambda_{sq} \beta \hbar} \omega \lambda_{sq} \sinh \hbar \omega \lambda_{sq} \beta \hbar \psi_{si} / 2$$

$$\int_{-\infty}^{\infty} e^{-\beta^2 A} d\beta = \frac{\sqrt{\pi}}{2 A^{1/2}}$$

$$\sigma^{(s,i)(2)}(\omega, T) = \beta \frac{\sinh \omega \beta \hbar}{\omega \beta \hbar} e^{-\frac{1}{2} \omega \beta \hbar \psi_{si}} (\omega + [z_s - z_{s+i}] \frac{1}{\hbar})$$

$$\frac{1}{4 \cosh \beta z_s / 2 \cosh \beta z_{s+i} / 2} \frac{\sqrt{\pi}}{2 \left(\sum \frac{2 \lambda_{sq} \sin^2 \frac{1}{2} \vec{q} (\vec{R}_{s+i} - \vec{R}_s)}{\sinh \hbar \omega \lambda_{sq} \beta \hbar} \omega \lambda_{sq} \cosh \hbar \omega \lambda_{sq} \beta \hbar \psi_{si} / 2 \right)}$$

$$e^{-2 \sum \lambda_{sq} \sin^2 \frac{1}{2} \vec{q} (\vec{R}_{s+i} - \vec{R}_s) [\cot \hbar \omega \lambda_{sq} \beta \hbar - \frac{\cosh \hbar \omega \lambda_{sq} \beta \hbar \psi_{si}}{\sinh \hbar \omega \lambda_{sq} \beta \hbar}]}$$

$$\frac{e^{\frac{1}{2} \omega \beta \hbar \psi_{si}}}{\hbar^2} t_{s,i}^2 (\vec{R}_{s+i} - \vec{R}_s) (\vec{R}_{s+i} - \vec{R}_s)$$

Solution of implicit eqns. High temperature

$$\omega + (\epsilon_s - \epsilon_{s+i}) \frac{1}{\kappa} - \sum 2\lambda q \sin^2 \frac{1}{2} \vec{q} (\vec{R}_{s+i} - \vec{R}_s) \omega q \lambda \psi_{s,i}/\hbar = 0$$

$$\psi_{s,i} = \frac{\omega + (\epsilon_s - \epsilon_{s+i})/\kappa}{\eta \cdot \omega_0}$$

$$\sigma^{(s,i)}(\omega, T) = \sigma^{(s,i)}(\omega, T) \frac{\sinh \hbar \kappa \omega \beta / 2}{\hbar \omega \beta / 2} e^{-\frac{1}{2} \hbar \beta \frac{[\omega + (\epsilon_s - \epsilon_{s+i})]}{(2) \eta \cdot \omega_0}}$$

$$\omega \sigma(\omega, T) = \frac{e^2}{\hbar^2} t_{s,i}^2 (\vec{R}_{s+i} - \vec{R}_s)(\vec{R}_{s+i} - \vec{R}_s) \frac{\sqrt{\pi}}{4 \cosh \beta \epsilon_s / \hbar \cosh \beta \epsilon_{s+i} / \hbar}$$

$$\frac{\beta}{2 \left(\sum \frac{2\lambda q \sin^2 \frac{1}{2} \vec{q} (\vec{R}_{s+i} - \vec{R}_s)}{\sinh \hbar \kappa \omega q \beta} \omega \lambda q^2 \cosh(1) \right)^{1/2}}$$

$$e^{-\eta \tanh \hbar \kappa \omega_0 \beta / 4}$$

$$= \frac{e^2}{\hbar^2} t_{s,i}^2 (\vec{R}_{s+i} - \vec{R}_s)(\vec{R}_{s+i} - \vec{R}_s) \frac{\sqrt{\pi}}{4 \cosh \beta \epsilon_s / \hbar \cosh \beta \epsilon_{s+i} / \hbar}$$

$$\frac{\beta}{2 \left(\frac{\eta \omega_0^2}{\frac{1}{2} \hbar \beta \omega_0} \right)^{1/2}} e^{-\eta \tanh \hbar \kappa \omega_0 \beta / 4}$$

$$U = \eta \hbar \omega_0 / 4$$

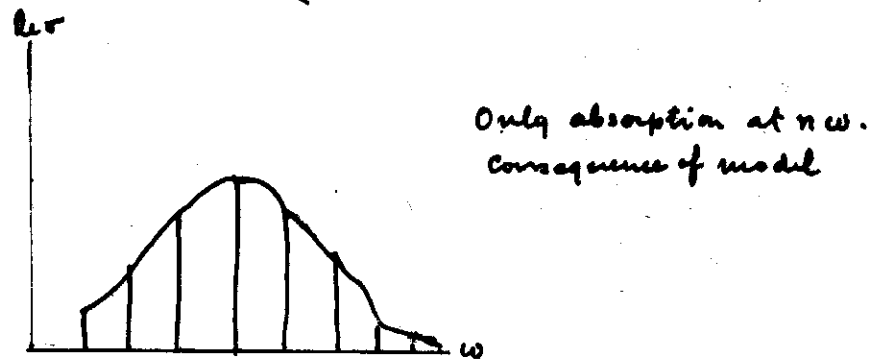
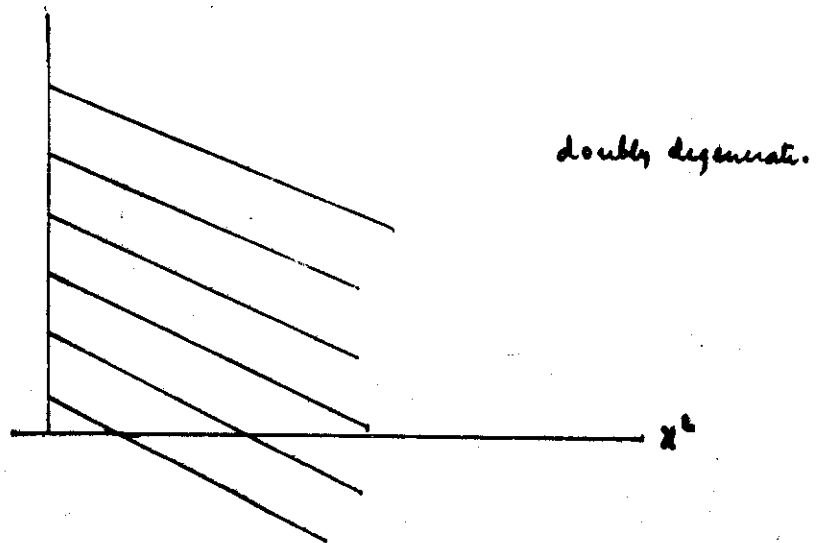
Ordered case

$$\sum_i (\vec{R}_{s+i} - \vec{R}_s)_z (\vec{R}_{s+i} - \vec{R}_s)_p = \delta_{zp} \frac{2a^2}{3}$$

Frequency of the real part of the conductivity

$$\sim \frac{\sinh \hbar \kappa \omega \beta / 2}{\hbar \omega \beta / 2} e^{-\frac{\hbar^2 (\omega + (\epsilon_s - \epsilon_{s+i}))^2}{16 U \hbar T}}$$

Suppose no dispersion no disorder, take two electronic levels $s, s+i$



What are the consequences if we include the effect of transfer integral in the calculation of the Polaron States?