



the  
**abdus salam**  
international centre for theoretical physics

SMR.1508 - 14

## **SUMMER SCHOOL ON PARTICLE PHYSICS**

**16 June - 4 July 2003**

### **NEUTRINO PHYSICS**

#### **Lectures IV & V**

**S. PARKE  
Fermilab  
Batavia, IL  
U.S.A.**



## References:

Past ICTP Summer School Lectures:

Lisi - 2001 (ICTP web)

Akhmedov - 1999 hep-ph/0001264.

Pakvasa - 1997 hep-ph/9804426

+ references there in

## Recent Reviews

Pakvasa + Valle hep-ph/0301601

Gonzalez = Garcia + Nir hep-ph/0202058

## MSW + Solar $\nu$ 's

Smirnov - hep-ph/0305106

Parke - SLAC Summer School 1986

Fermilab - Conf - 86 - 131-T  
available on SPIRES

click Fermilab - Library - Server:

↓  
little known

Problem:

2

Kamiland:

Expected:  $87 \pm 6$

Observed: 54

Prob < 0.05%

$$\frac{\text{Diff.}}{\sqrt{\text{Exp}}} = \frac{87 - 54}{\sqrt{87}} = 3.5 = \frac{33}{9.3}$$

K2K:

Expected:  $80 \pm 6$

Observed: 56

Prob:  $\sim 1.3\%$

$$\frac{\text{Diff.}}{\sqrt{\text{Exp}}} = \frac{80 - 56}{\sqrt{80}} = \frac{24}{8.9} = 2.7$$

Redo with Poisson ?  
Statistics

Problem:

For 3 flavors:

SHOW

$$2 \sum_{i>j} g_m (U_{\alpha i} U_{\beta i}^* U_{\alpha j}^* U_{\beta j}) \sin^2 \Delta_{ij}$$

$$= \pm \sin \delta \sin^2 \theta_{12} \sin^2 \theta_{23} \sin^2 \theta_{13} \cos \theta_{33} \\ \sin \Delta_{21} \sin \Delta_{32} \sin \Delta_{31}$$

where  $\Delta_{ij} = \frac{8m_{ij}\omega}{4E}$

SAME for all  $\alpha \neq \beta$  (if  $\nu \neq \bar{\nu}$ )

- or + for  $\mu \rightarrow e, \tau \rightarrow \mu, e \rightarrow \tau$   
 + or - for  $e \rightarrow \mu, \mu \rightarrow \tau, \tau \rightarrow e$

Hint:  $\Delta_{31} - \Delta_{32} - \Delta_{21} = 0$

# $\theta_{13}$ and Beyond

4

Reactors:  $\bar{\nu}_e \rightarrow \bar{\nu}_e$  disappearance  
ala Chooz (KamLAND etc.)

$$P_{\bar{\nu}_e \rightarrow \bar{\nu}_e} \approx 1 - \sin^2 2\theta_{13} \sin^2 \Delta_{31}$$

$$\Delta_{31} = \frac{\delta m_{31}^2 L}{4E}$$

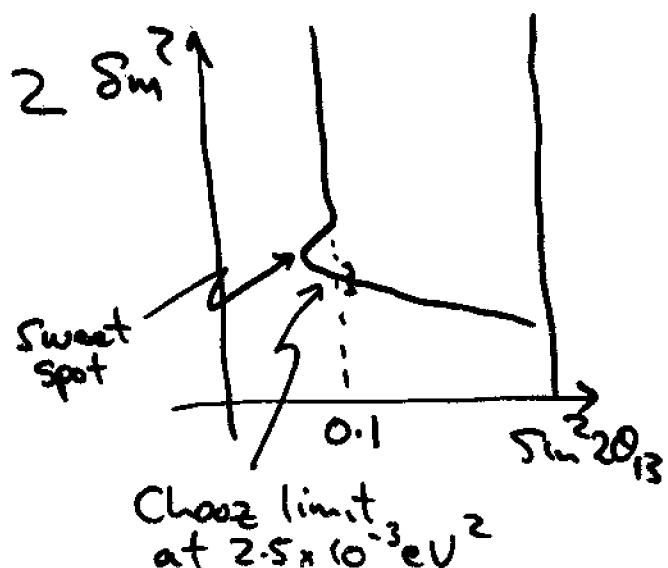
Currently from Chooz

$$\sin^2 2\theta_{13} < 0.05 \text{ at any } E.$$

$$\text{at } 2.5 \times 10^{-3} \text{ eV}^2: \sin^2 2\theta_{13} < 0.01$$

SK Smin

Could get a factor of 2  $\sin^2$   
by working at a better  
distance

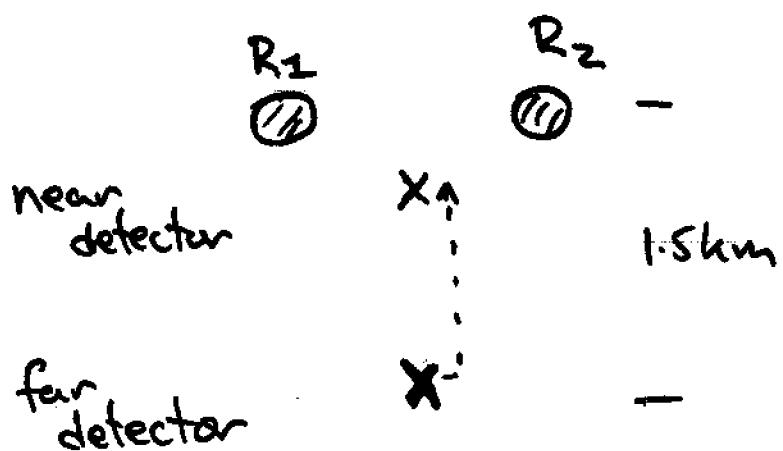


Oscillation Peak (or dip in disappearance)  $\xrightarrow{5}$   
occurs at

$$L = \frac{4\pi E}{28m_{atom}^2} = \frac{\pi E}{2(1.27)m_{atom}^2} = \frac{\pi \cdot 3 \times 10^{-3}}{2.54 \times 2.5 \times 10^{-3}} \text{ km}$$

$$\approx 1.5 \text{ km}$$

Also reduce systematic errors  
 $\rightarrow 1\%$



Differential Total Flux measurement

maybe able to get to

$$\sin^2 2\theta_{13} \sim 0.01$$

Shape  
Linder et al.  
'03

- To do better than Chooz, relatively easy
- to get to 0.01 HARD, DIFFICULT, ...

$$P_{\nu_e \rightarrow \nu_e} = \left| U_{e3} U_{e3}^* e^{-im_3 \omega / 2E} + U_{e2} U_{e2}^* e^{-im_2 \omega / 2E} + U_{e1} U_{e1}^* e^{-im_1 \omega / 2E} \right|^2$$

6

$$U_{e1} U_{e1}^* = 1 - U_{e2} U_{e2}^* - U_{e3} U_{e3}^* \quad \text{unitarity}$$

$$P_{\nu_e \rightarrow \nu_e} = \left| 1 - 2i S_{13}^2 e^{-2i \Delta_{31}} \sin \Delta_{31} - 2i C_{13} S_{12} e^{2i \Delta_{21}} \sin \Delta_{21} \right|^2$$

Square

$$= 1 - \sin^2 2\theta_{13} \left( \sin^2 \Delta_{31} + 2 S_{12}^2 \sin \Delta_{31} \cos \Delta_{32} \cdot \sin \Delta_{21} + S_{12}^2 \sin^2 \Delta_{21} \right)$$

$$- C_{13}^2 \sin^2 2\theta_{12} \sin^2 \Delta_{21}$$

At distances, intermediate between  
 (30km) Cleop and KamLAND, there are wiggles which, for  $\sin^2 2\theta_{13}$  not too small,  
 could be used to determine  $\sin^2 \theta_6, S_0^2$   
 $\sin^2 \theta_{\text{atm}}$ . see Petcov et al June 2003

clinton

Old UCI Page

## The President's View

Announcement (*excerpted from remarks at the MIT commencement, June 6, 1998*)

Paper

(submitted to  
Phys.Rev.Lett)

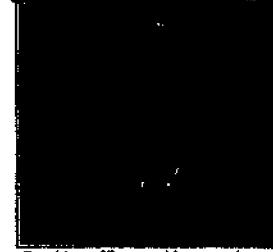
Clinton on  
Neutrinos

FAQ

Glossary

Links

[W]e must help you to ensure that America continues to lead the revolution in science and technology. Growth is a prerequisite for opportunity, and scientific research is a basic prerequisite for growth. Just yesterday in Japan, physicists announced a discovery that tiny neutrinos have mass. Now, that may not mean much to most Americans, but it may change our most fundamental theories -- from the nature of the smallest subatomic particles to how the universe itself works, and indeed how it expands.



President Clinton addresses the  
graduating class of MIT

This discovery was made, in Japan, yes, but it had the support of the investment of the U.S. Department of Energy. This discovery calls into question the decision made in Washington a couple of years ago to disband the Super-conducting Supercollider, and it reaffirms the importance of the work now being done at the Fermi National Acceleration Facility in Illinois.

The larger issue is that these kinds of findings have implications that are not limited to the laboratory. They affect the whole of society -- not only our economy, but our very view of life, our understanding of our relations with others, and our place in time.

The full text of the President's address is also available.

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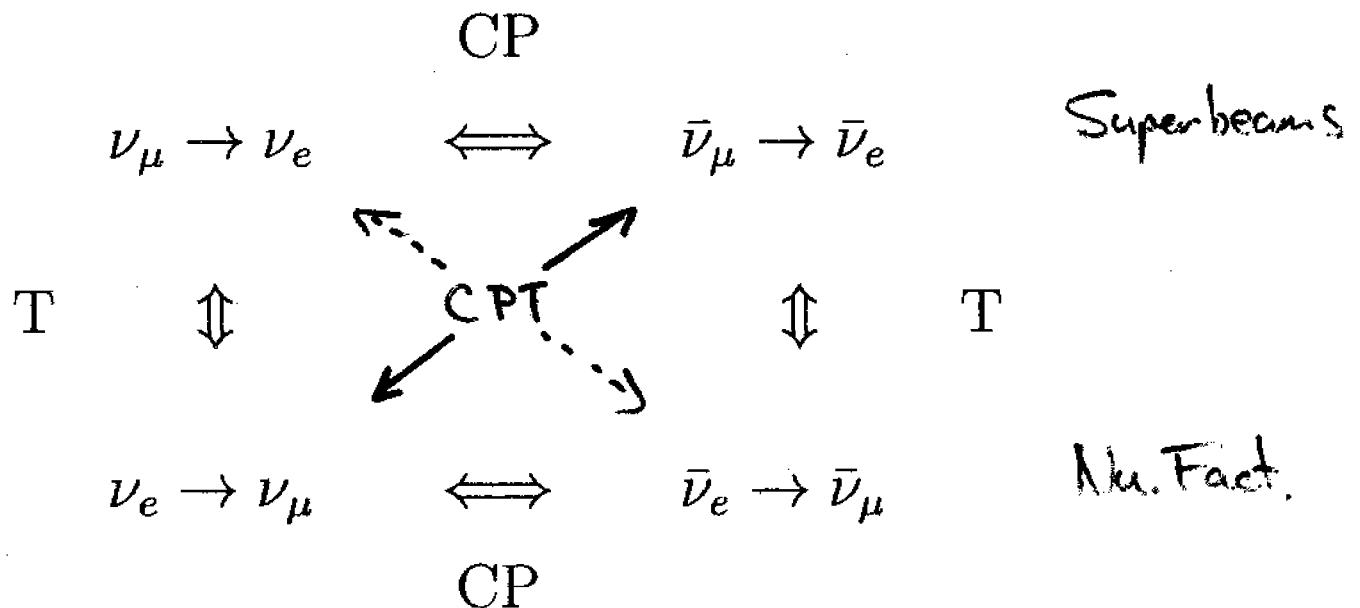
$\theta_{13}$  and

CP

Violation:

Fermion  
Mass

## Leptonic CP and T Violation in Oscillations



IN GENERAL (in vacuum):

CP Violation:

$$\alpha \neq \beta \quad P(\nu_\alpha \rightarrow \nu_\beta) \neq P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta)$$

T Violation:

$$\begin{aligned} \alpha \neq \beta \quad &P(\nu_\alpha \rightarrow \nu_\beta) \neq P(\nu_\beta \rightarrow \nu_\alpha) \\ \text{and} \quad &P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta) \neq P(\bar{\nu}_\beta \rightarrow \bar{\nu}_\alpha) \end{aligned}$$

CPT Violation:

$$\text{any } \alpha, \beta \quad P(\nu_\alpha \rightarrow \nu_\beta) \neq P(\bar{\nu}_\beta \rightarrow \bar{\nu}_\alpha)$$

$$\underline{\underline{P_{\nu_\mu \rightarrow \nu_e}}}$$

$$P_{\mu \rightarrow e} = | U_{e3} U_{\mu 3}^* e^{-im_3^2/2L} + U_{e2} U_{\mu 2}^* e^{-im_2^2/2L} + U_{e1} U_{\mu 1}^* e^{-im_1^2/2L} |$$

multiple by  $e^{+im_1^2/2L}$

$$\text{then } U_{e1} U_{\mu 1}^* = -U_{e3} U_{\mu 3}^* - U_{e2} U_{\mu 2}^* \quad \left. \right\} D_{ij} = \frac{8m_j L}{4E}$$

$$= | 2U_{e3} U_{\mu 3}^* e^{-i\Delta_{31}} \sin \Delta_{31} + 2U_{e2} U_{\mu 2}^* e^{-i\Delta_{21}} \sin \Delta_{21} |$$

multiple by  $e^{+i\Delta_{21}}$

$$= | e^{-i\Delta_{21}} 2U_{e3} U_{\mu 3}^* \sin \Delta_{31} + 2U_{e2} U_{\mu 2}^* \sin \Delta_{21} |$$

$$2 U_{e3} U_{\mu 3}^* = 2 S_{13} C_{13} \cdot S_{23} e^{-i\delta}$$

$$2 U_{e2} U_{\mu 2}^* \simeq 2 C_{13} S_{12} (C_{23} C_{12} - S_{13} \cancel{S_{12}} \cancel{S_{23}} e^{-i\delta})$$

$$P_{\mu \rightarrow e} = \left| e^{-i(\Delta_{32} + \delta)} \sqrt{P_{\text{atm}}} + \sqrt{P_\odot} \right|^2$$

with

$$P_{\text{atm}} = \left( S_{23} \sin^2 \theta_{13} \sin \delta_{31} \right)^2$$

$$P_\odot = \left( C_{13} C_{23} \sin 2\theta_{12} \sin \delta_{21} \right)^2$$

$$\text{Oscillation Max: } \Delta_{31} \approx \Delta_{32} = \frac{\pi}{2}$$

$$\text{Max: CP Violations: } \delta = \frac{\pi}{2}$$

$$P_{\mu \rightarrow e} = \left| \sqrt{P_{\text{atm}}} \pm \sqrt{P_\odot} \right|^2$$

$$(+\nu)(-\bar{\nu})$$

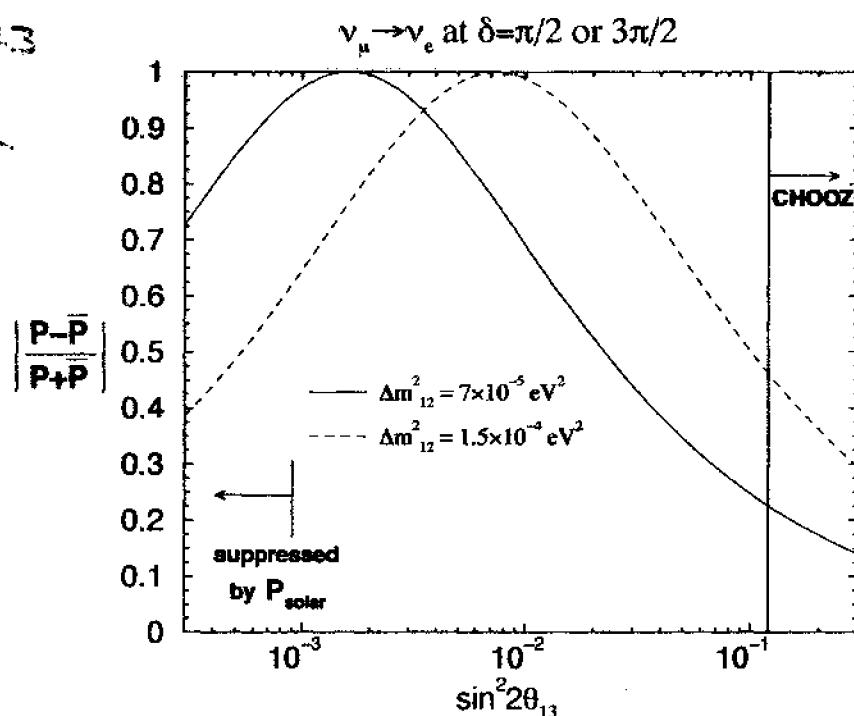
$$\text{Asym: } \frac{P - \bar{P}}{P + \bar{P}} = \frac{2 \sqrt{P_{\text{atm}}} \sqrt{P_\odot}}{(P_{\text{atm}} + P_\odot)}$$

# $P(\nu_\mu \rightarrow \nu_e)$

## Why Everybody is Excited!

- Maximum Allowed Asymmetry ( $\delta = \frac{\pi}{2}$  or  $\frac{3\pi}{2}$ ) for  $\nu_\mu \rightarrow \nu_e$  at first Oscillation Maximum in vac:
  - $P, \bar{P} = |a_{\mu \rightarrow e}^{atm} + a_{\mu \rightarrow e}^{\odot}|^2 \approx (\sin \theta_{23} \sin 2\theta_{13} \pm \sqrt{P_\odot})^2$
  - $|P - \bar{P}| \approx 4\sqrt{P_\odot} \sin \theta_{23} \sin 2\theta_{13}$
  - $P + \bar{P} \approx 2 \sin^2 \theta_{23} \sin^2 2\theta_{13} + 2P_\odot$
- easily generalized*
- 18<sup>II</sup>
  - off O.M.
  - matter

$$\begin{aligned} a_{\mu \rightarrow e}^{atm} &\propto U_{e3} \\ a_{\mu \rightarrow e}^{\odot} &\propto U_{e2} \end{aligned}$$



Asymmetry  
LARGE: ✓  
0

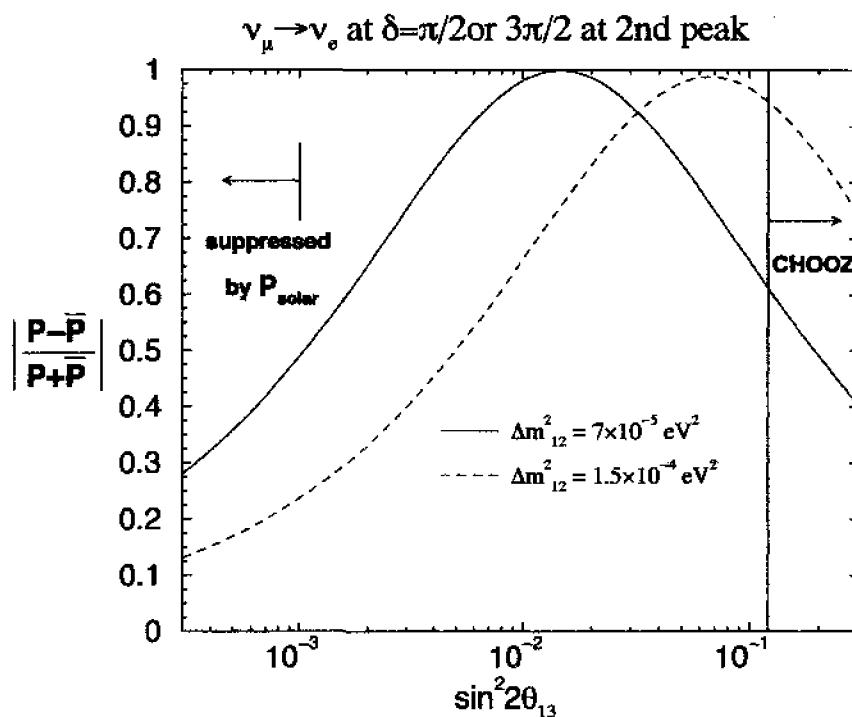
- Peak occurs at

$$\sin^2 2\theta_{13} \approx \frac{\sin^2 2\theta_{12}}{\tan^2 \theta_{23}} \left[ \frac{\pi}{2} \frac{\delta m_{12}^2}{\delta m_{13}^2} \right]^2$$

at OM  $\sqrt{P_\odot} = \cos \theta_{13} \cos \theta_{23} \sin 2\theta_{12} \sin \left( \frac{\pi}{2} \frac{\delta m_{12}^2}{\delta m_{13}^2} \right)$

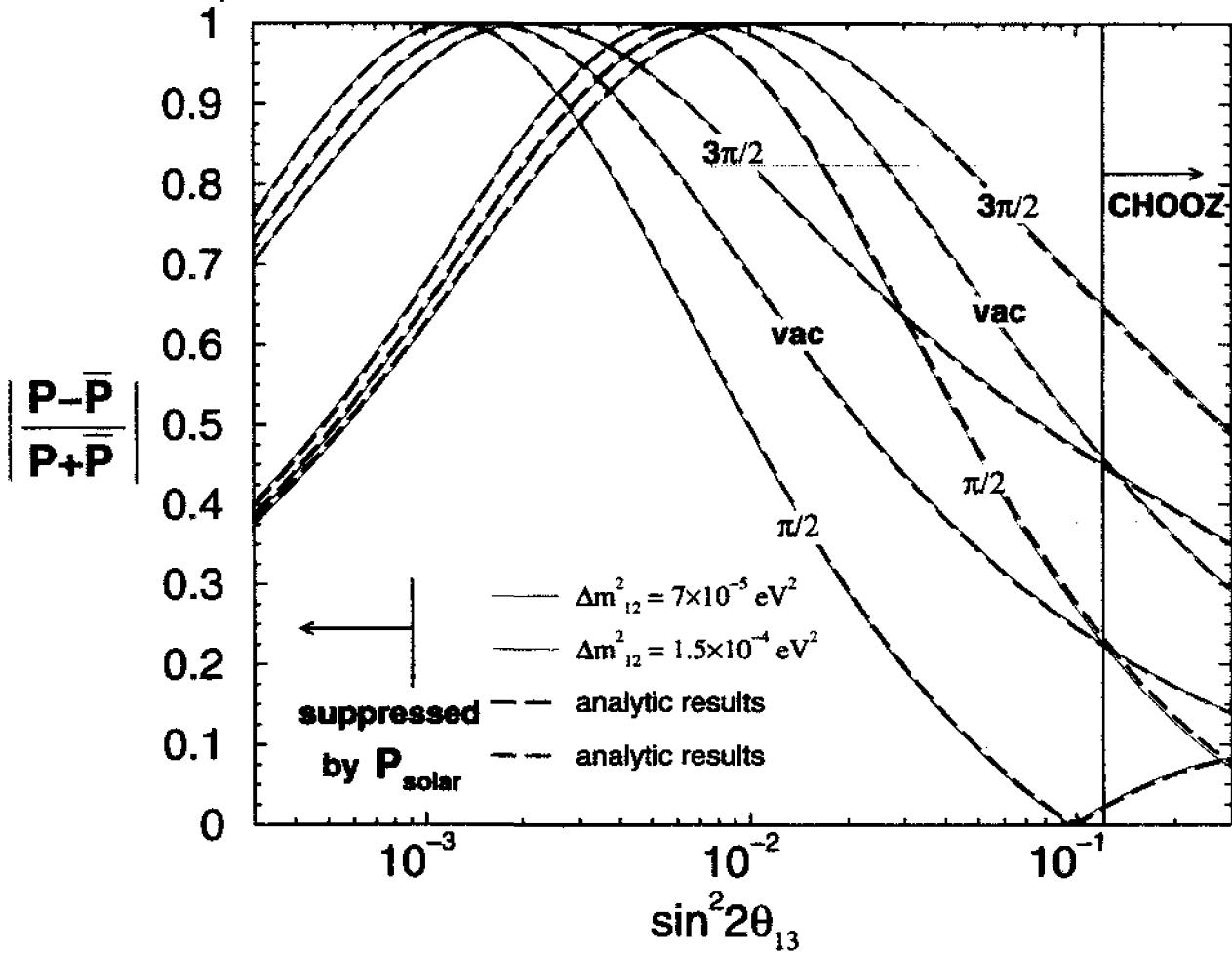
- For BK

## 2nd Peak



- Peak occurs at  $\sin^2 2\theta_{13} \approx \frac{\sin^2 2\theta_{12}}{\tan^2 \theta_{23}} \left[ \frac{3\pi}{2} \frac{\delta m_{12}^2}{\delta m_{13}^2} \right]^2$

$\nu_\mu \rightarrow \nu_e$  at  $\delta = \pi/2$  or  $3\pi/2$  in matter for  $L = 732$  km



$$P_{\mu \rightarrow e} = P_{\text{atm}} + 2\sqrt{P_{\text{atm}}} \sqrt{P_0} \cos(\Delta_{32} \pm \delta) + P_0$$

$$\cos(\Delta_{32} + \delta) = \cos \Delta_{32} \cos \delta + \underbrace{\sin \Delta_{32} \sin \delta}$$

CP violating term is

$$\sin \delta \cdot \sin 2\theta_{12} \sin 2\theta_{23} \sin 2\theta_{13} \cos \theta_{13} \\ * \sin \Delta_{21} \sin \Delta_{32} \sin \Delta_{31}$$

FULL EXPRESSION:

$$\begin{aligned}
 P_{\mu \rightarrow e} &= P_{\otimes} - 2 S_{23}^2 S_{2(13)}^2 S_{12}^2 \sin \Delta_{31} \cos \Delta_{32} \frac{\sin \Delta_{21}}{\text{small}} \\
 &+ P_{\odot} (1 - 2 S_{13} \tan \theta_{23} \tan \theta_{12} \cos \delta + S_{13}^2 \tan^2 \theta_{23} \tan^2 \theta_{12}) \\
 &+ 2 \sqrt{P_{\otimes} P_{\odot}} \cos (\Delta_{32} \mp \delta)
 \end{aligned}$$

$$P_{\otimes} = S_{23}^2 S_{2(13)}^2 \sin^2 \Delta_{31}$$

$$P_{\odot} = C_{13}^2 C_{23}^2 S_{2(12)}^2 \sin^2 \Delta_{21}$$

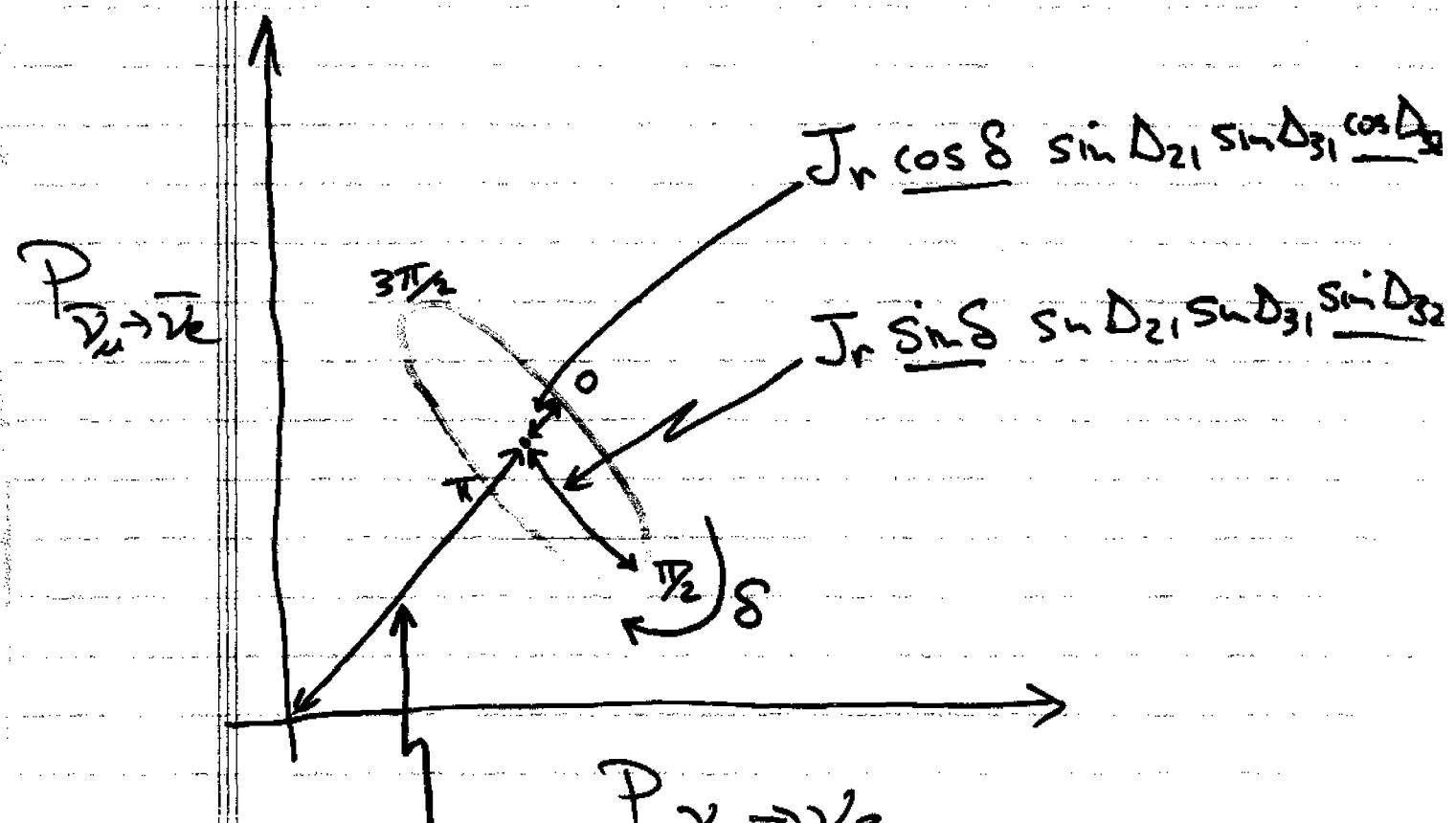
$P_{\odot}$  is irrelevant unless  $S_{13}$  is small

where  $(1 - 2 S_{13} \dots) \approx 1$

$$\frac{\Delta m}{\Delta_{31}} = \frac{S_{12}^2}{\sum m_i^2} \approx \frac{1}{30} \quad \text{also } \sin \Delta_{31} \approx 1 \\
 \cos \Delta_{32} \approx 0$$

## Bi-Probability Plot

$(P, \bar{P})$  varying only  $S_{CP}$



$$\sqrt{2} \sin^2 \theta_{23} \sin^2 \theta_{13} \sin^2 \Delta_{31}$$

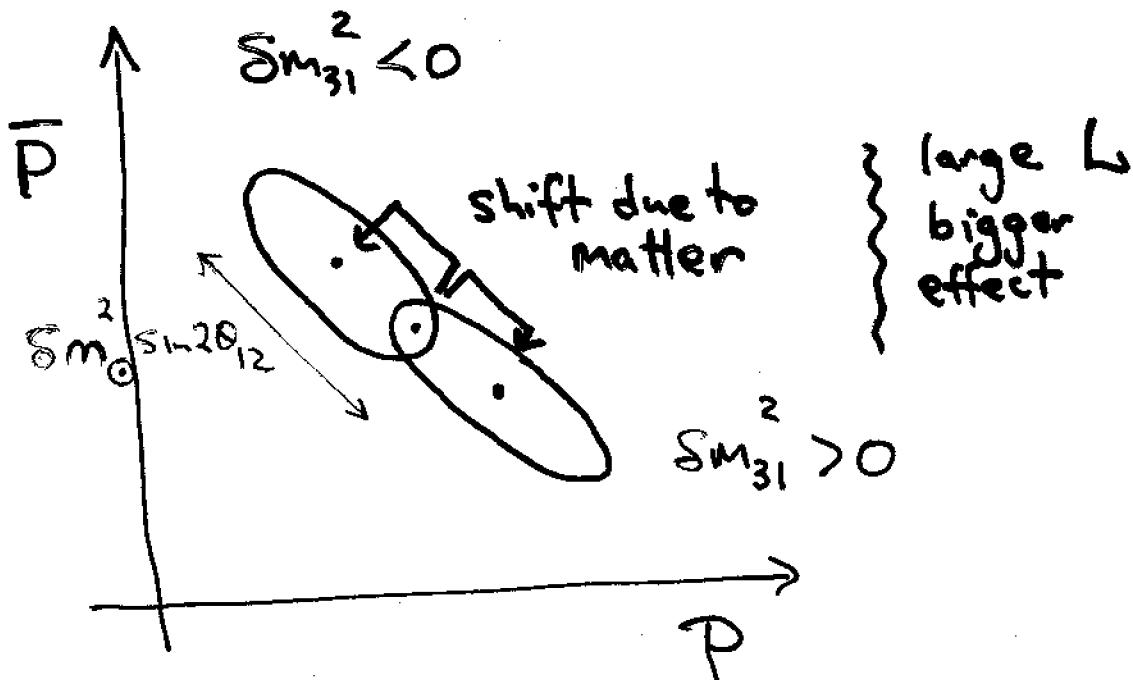
If 5MA or low solar sol

ellipse collapses to point.

LMA - Wow!!!

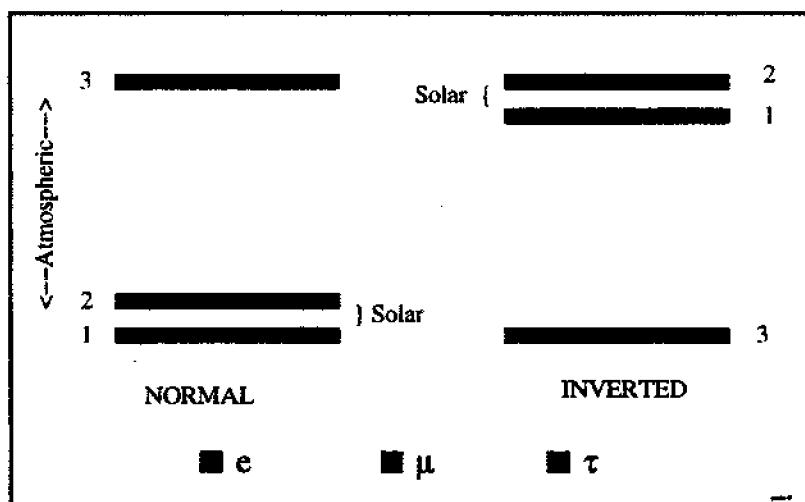
# Matter Effects:

18



$$P_{\text{mat}}^{\text{center}} \approx \left(1 \pm \frac{E_\nu}{6 \text{GeV}}\right) P_{\text{vac}}^{\text{center}}$$

O.M.  
for small  
 $E_\nu$



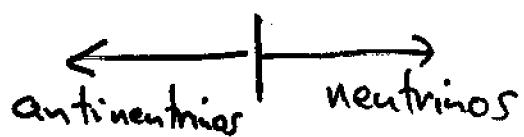
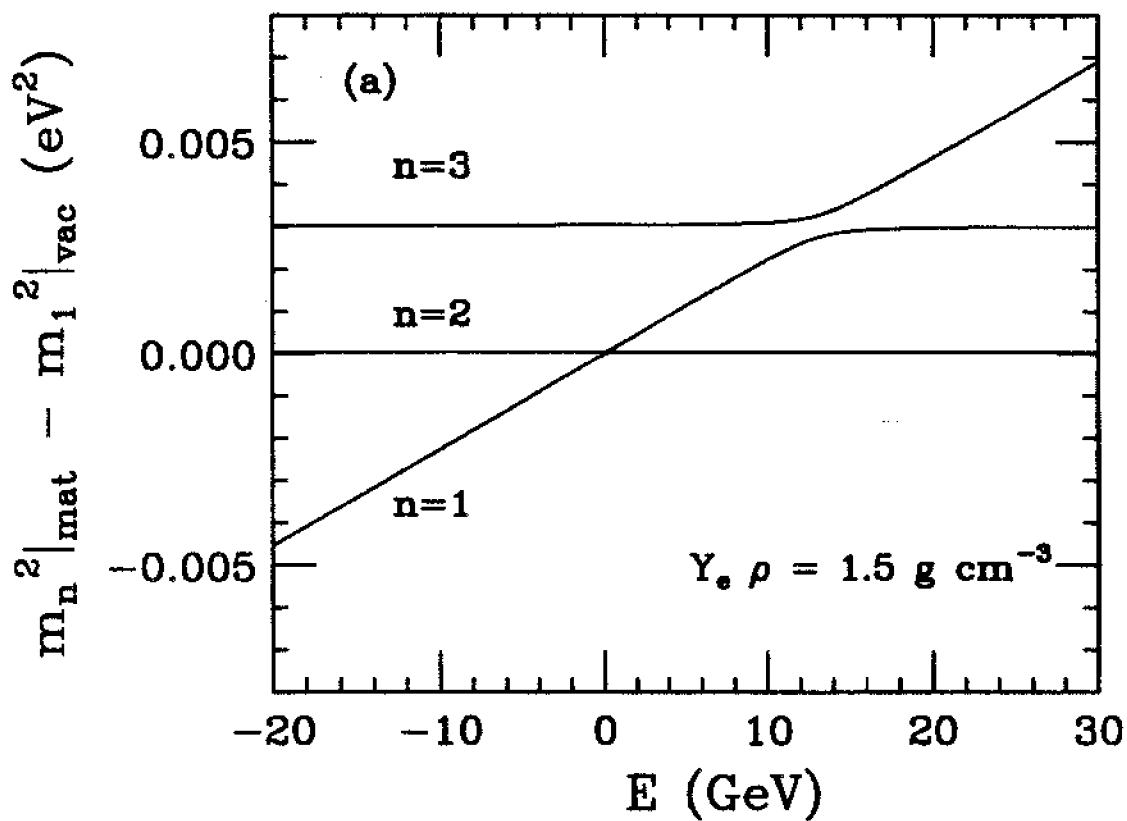
$\delta m_{31}^2 > 0$

$\delta m_{31}^2 < 0$

$$\left( \sin 2\theta_{13} S_{M_{31}}^2 \right)_{\text{matter}} = \left( \sin 2\theta_{13} S_{M_{31}}^2 \right)_{\text{vac}}^{19}$$

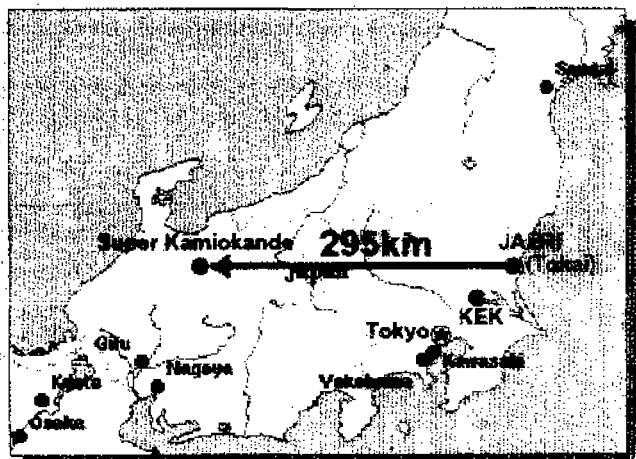
and

$$\left( \sin 2\theta_{12} S_{M_{21}}^2 \right)_{\text{MATTER}} = \left( \sin 2\theta_{12} S_{M_{21}}^2 \right)_{\text{vac}}$$



## JHF → Super-Kamiokande

- ✓ 295 km baseline
- ✓ Super-Kamiokande:
  - 22.5 kton fiducial
  - Excellent e/μ ID
  - Additional  $\pi^0/e$  ID
- ✓ Hyper-Kamiokande
  - 20× fiducial mass of SuperK
- ✓ Matter effects small
- ✓ Study using fully simulated and reconstructed data



Requires New Beamline:

~~<http://www.ornl.gov/>~~

LOI: hep-ex/0106019

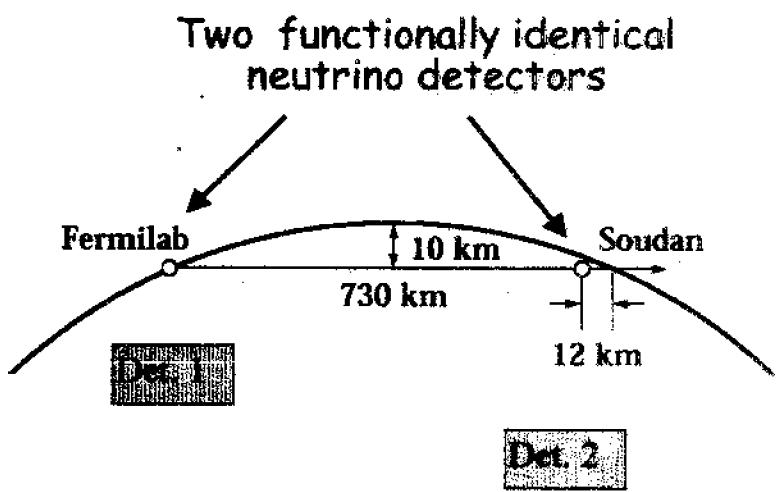
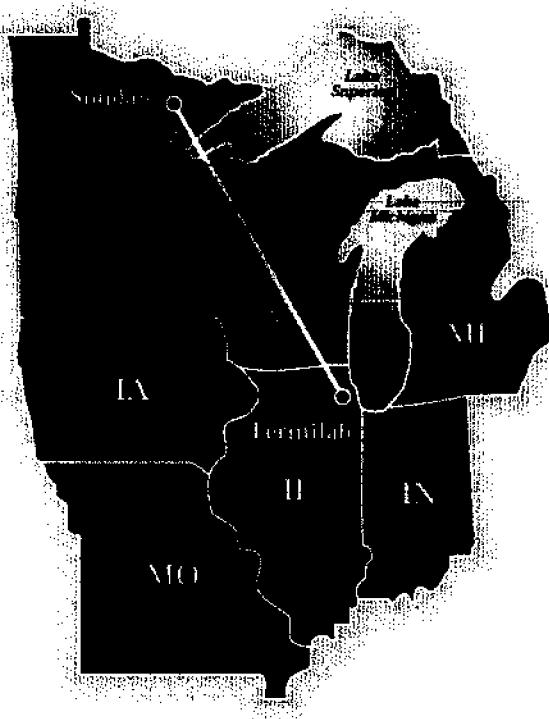
<http://www-nu.kek.jp/jhfnu/>

$$\bar{E}_\nu = 0.6 \text{ to } 1.0 \text{ GeV}$$

$\sim 20\%$  spread

$$L = 295 \text{ km}$$

## The NUMI Beamline



New Detector Required:

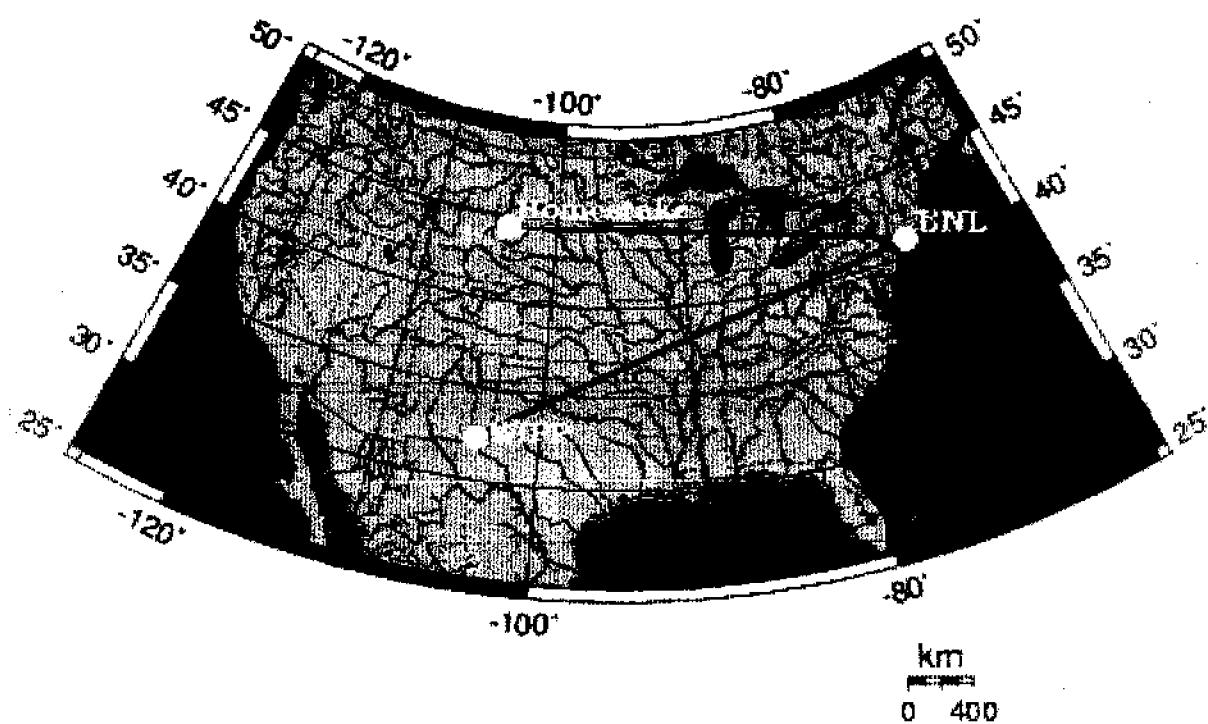
<http://www-off-axis.fnal.gov/>

LOI: hep-ex/0210005

$E \sim 2 \text{ GeV}$

$L \sim 732 \text{ km}$   
 $\pm 200 \text{ km}$

## Brookhaven to Homestake OR WIPP



$L = \underline{2540 \text{ km}}$  or  $2880 \text{ km}$

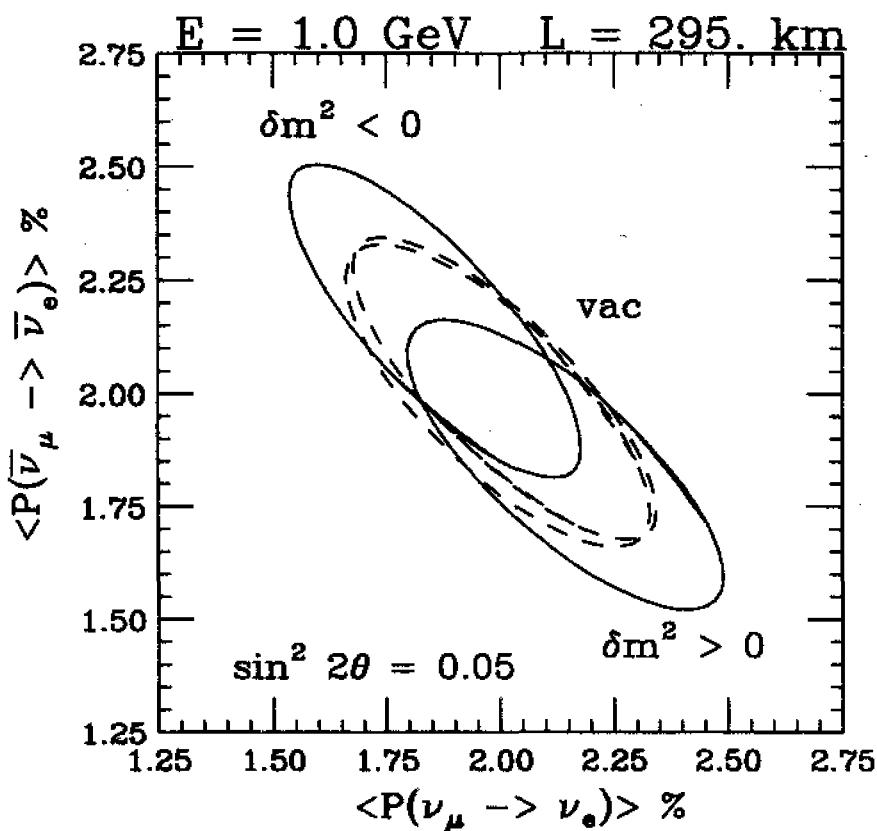
$2^{\text{nd}}$  peak 2 GeV:

New Beamline, New Detector:

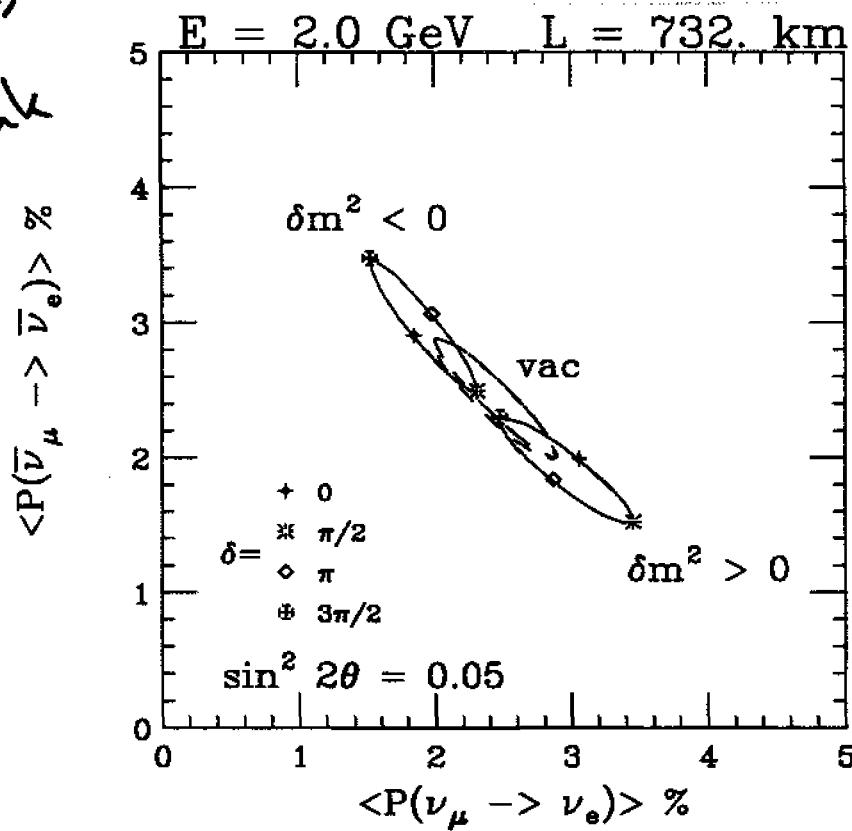
<http://www.neutrino.bnl.gov/>

LOI: hep-ex/0205040

JHF  $\rightarrow$  SK  
1<sup>st</sup> peak

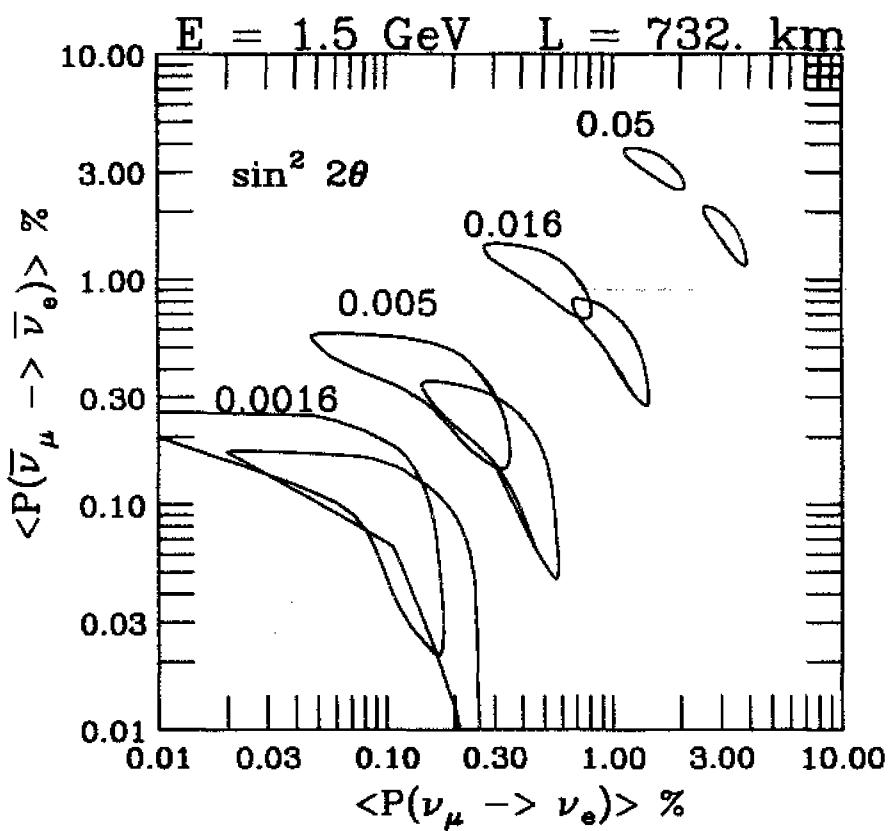


NuMI  $\rightarrow$   
1<sup>st</sup> peak



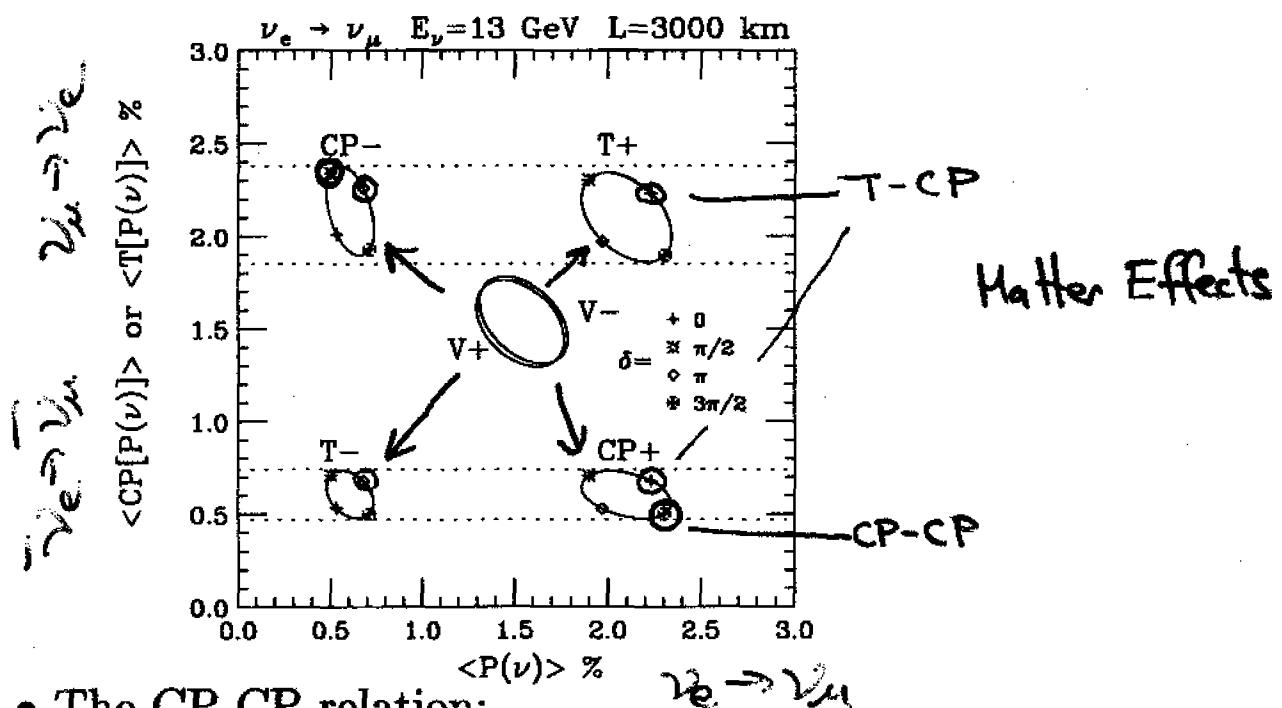
Matter Effects  
separate  
 $\delta m^2 > 0$   
from  
 $\delta m^2 < 0$

As  $\sin^2 2\theta_{13}$  Varies:



NuMI offAxis LOI  
hep-ex/0210005

## Anatomy of the Bi-Probability Plot:



- The CP-CP relation:

$$\begin{aligned}
 & P(\nu_e \rightarrow \nu_\mu; \Delta m_{31}^2, \Delta m_{21}^2, \delta, a) \\
 &= P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu; -\Delta m_{31}^2, -\Delta m_{21}^2, \delta, a) \\
 &\approx P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu; -\Delta m_{31}^2, +\Delta m_{21}^2, \pi + \delta, a)
 \end{aligned}$$

] evolution eqn.

- The T-CP relation:

assuming symmetry matter distribution

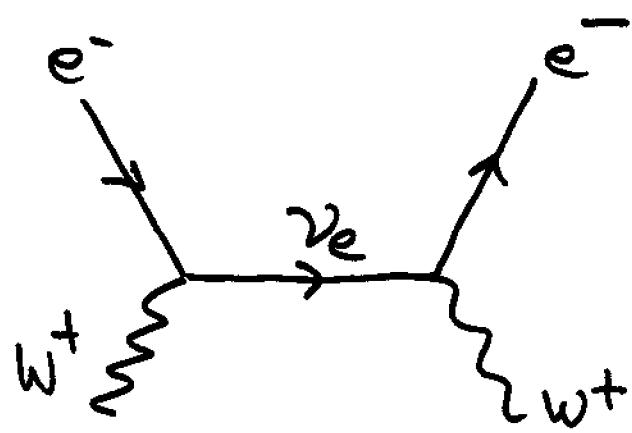
$$\begin{aligned}
 & P(\nu_\mu \rightarrow \nu_e; \Delta m_{31}^2, \Delta m_{21}^2, \delta, a) \\
 &= P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu; -\Delta m_{31}^2, -\Delta m_{21}^2, -\delta, a) \\
 &\approx P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu; -\Delta m_{31}^2, +\Delta m_{21}^2, \pi - \delta, a)
 \end{aligned}$$

] evolution eqn.

- ≈ trade sign of  $\delta m_{12}^2$  for shift by  $\pi$  of  $\delta$ :

$$(\dots) + \delta m_{12}^2 [(\dots) \cos \delta + (\dots) \sin \delta]$$

↑  
Small correction here



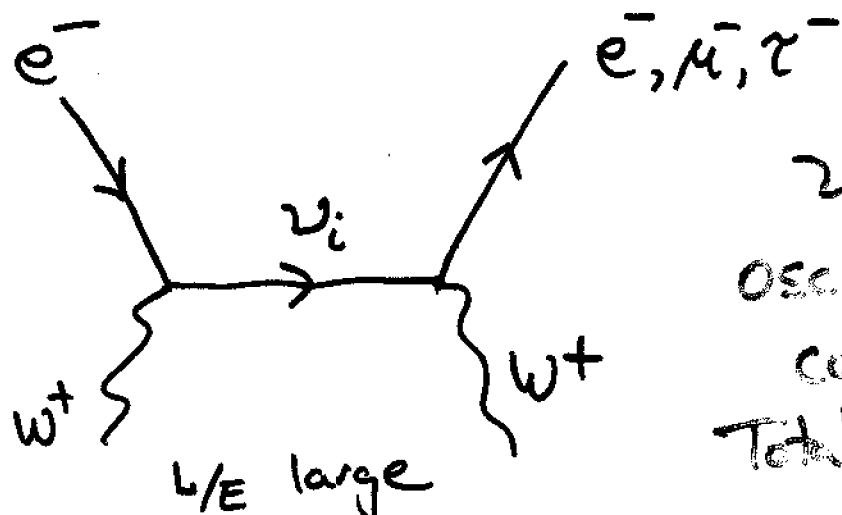
Days of  
Weniger + Salam 21

S.M.

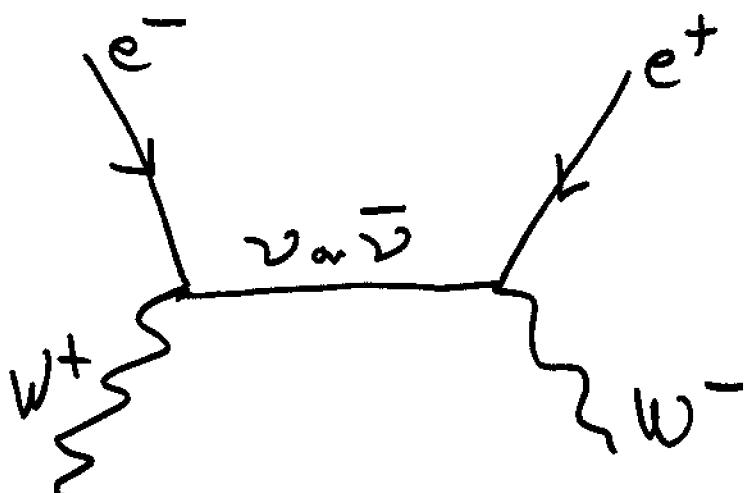
flavor diagonal!

conserves

$l_e, l_\mu, l_\tau$



$\nu$  mass  
oscillations ~~( $\Delta m^2$ )~~  
causes  
Total lepto. #



Majorana Mass  
Lepton # Violating  
(highly suppressed)  
Why ???

# Fermion Mass: (Cartoon:) 27

- Consider a massless electron:

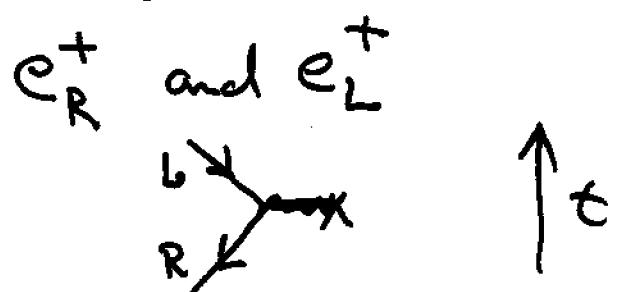
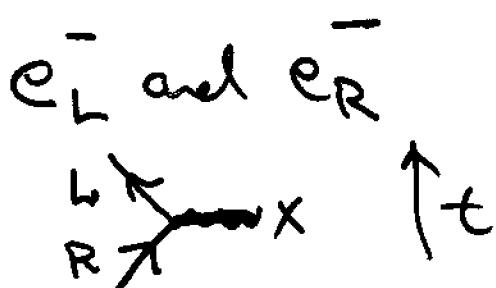
$$\bar{e}_L^- \quad \bar{e}_R^-$$

$$e_R^+ \quad e_L^+$$

- Four States:  $2\bar{e}^- + 2e^+$



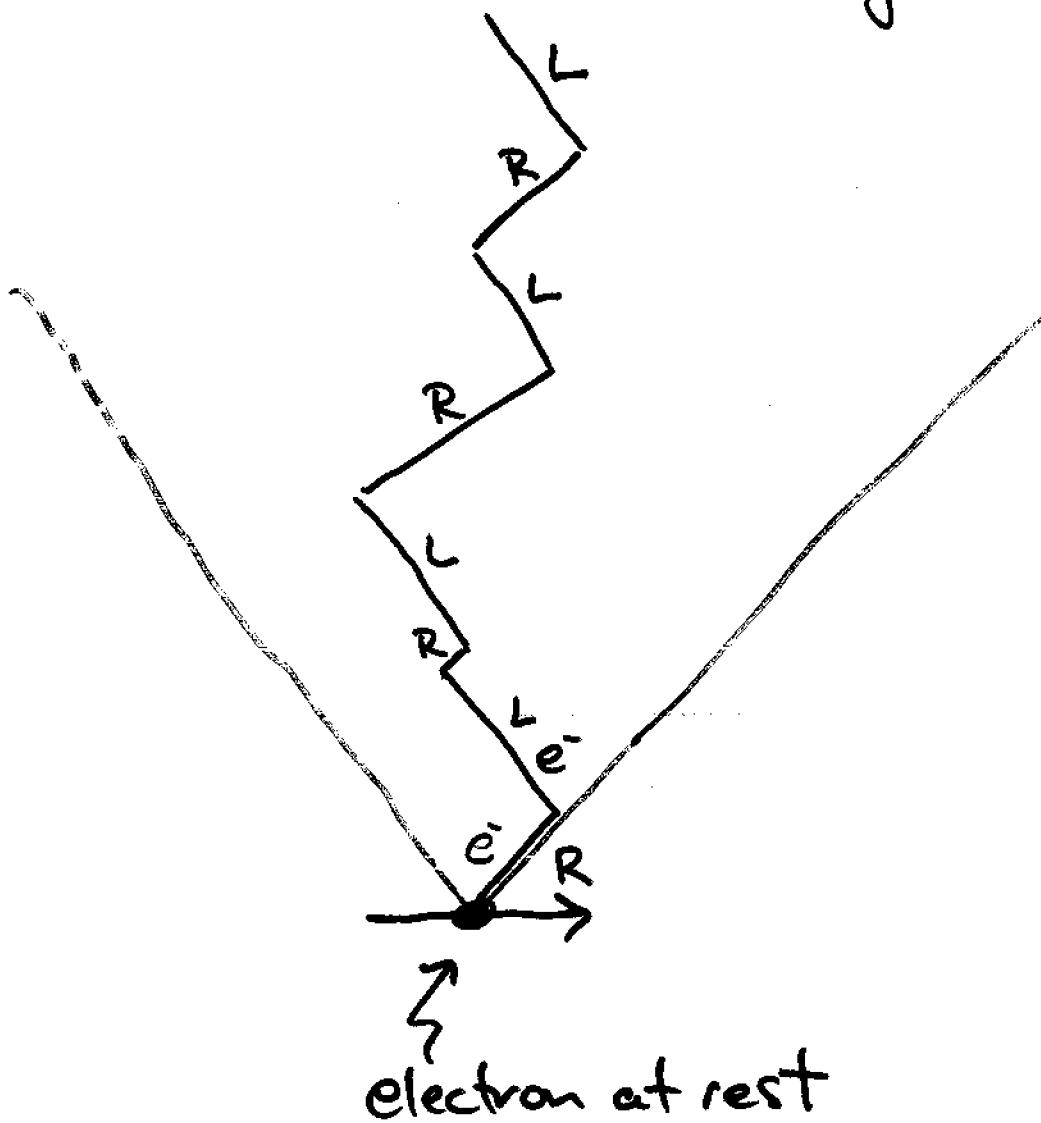
- Mass is a coupling between Left and Right. (chiral states)



CANNOT couple

$e^-_L$  to  $e^+_R$

because of electric charge conservation.



# 29

## Massive Fermion with given Spin:

$$P^2 = M^2 \quad S^2 = -1 \quad P \cdot S = 0$$

(at rest  $P = (M, 0, 0, 0)$   $S = (0, 0, 0, 1)$ )

e.g. { Spin in  $\hat{z}$  direction

$$P = \frac{P+MS}{2} + \frac{P-MS}{2}$$

$$\frac{P+MS}{2} \quad \text{and} \quad \frac{P-MS}{2} \quad \underline{\text{Massless}}$$

$$\left( \frac{P \pm MS}{2} \right)^2 = \frac{P^2 \pm 2P \cdot S + M^2 S^2}{4}$$

$$= \frac{M^2 \pm 2M \cdot 0 + M^2 (-1)}{4}$$

$S^{14}$  Dirac Eqn  
Massive, Spin

$$= 0$$

massless  
spinors

$$U(P, S) = \frac{1+\gamma_5}{2} U\left(\frac{P+MS}{2}\right) + e^{i\phi} \frac{1-\gamma_5}{2} U\left(\frac{P-MS}{2}\right)$$

chiral  
projections

$$U(P, S) = \frac{1 + \gamma_S \phi}{2} U(P)$$

$$= \frac{1 + \gamma_S}{2} \underbrace{\frac{1 + \phi}{2} U(P)}_{Q_+} + \frac{1 - \gamma_S}{2} \underbrace{\frac{1 + \phi}{2} U(P)}_{Q_-}$$

$$Q_{\pm} \bar{Q}_{\pm} = \frac{P \pm MS}{2}$$

and  $\frac{P \pm MS}{2} Q_{\pm} = 0$

Therefore  $Q_{\pm} = U\left(\frac{P \pm MS}{2}\right)$

Relative phase determined by

$$\bar{U}(P, S) U(P, S) = 2M$$

Massive particle at rest,  
Spin in z direction:

31

$$P = \frac{P+MS}{2} + \frac{P-MS}{2}$$

$$(M, 0, 0) = \frac{M}{2} (1, 0, 0, 1) + \frac{M}{2} (1, 0, 0, -1)$$

BIG BOOST IN Z-DIRECTION  $\begin{pmatrix} \gamma & \gamma \\ \gamma & \gamma \end{pmatrix}$

$$\gamma M(1, 0, 0, \beta) = \frac{M}{2} (\gamma + \beta \gamma)(1, 0, 0, 1) + \frac{M}{2} (\gamma - \beta \gamma)(1, 0, 0, -1)$$

$$\text{as } \beta \rightarrow 1 \quad (\gamma + \beta \gamma) \rightarrow 2\gamma \quad (\gamma - \beta \gamma) \rightarrow \frac{1}{2}\gamma$$

$$\text{where } \gamma = E/M$$

$$E(1, 0, 0, \beta) \approx E(1, 0, 0, 1) + \frac{m^2}{E} (1, 0, 0, -1)$$

$$U(P, S) = \frac{1+\gamma_5}{2} U\left(\frac{P+MS}{2}\right) + \frac{1-\gamma_5}{2} U\left(\frac{P-MS}{2}\right)$$

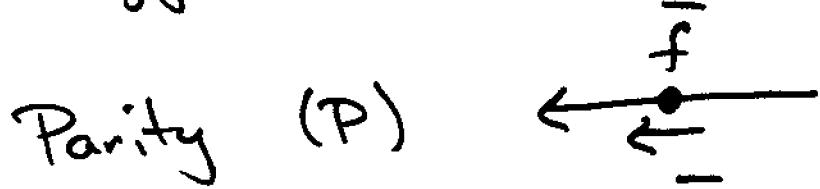
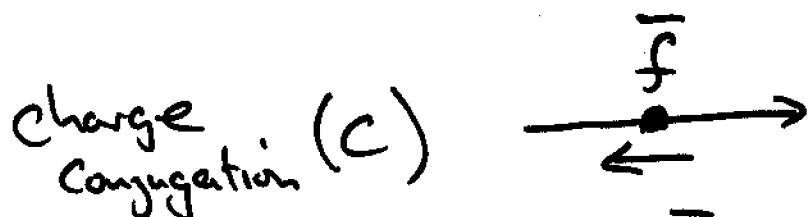
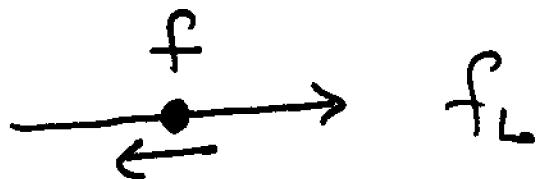
HELICITY  
STATE

RIGHT  
CHIRAL PROJECTIONS

LEFT

HELICITY  $\neq$  CHIRALITY  
for massive particles

## Mass less Fermion:

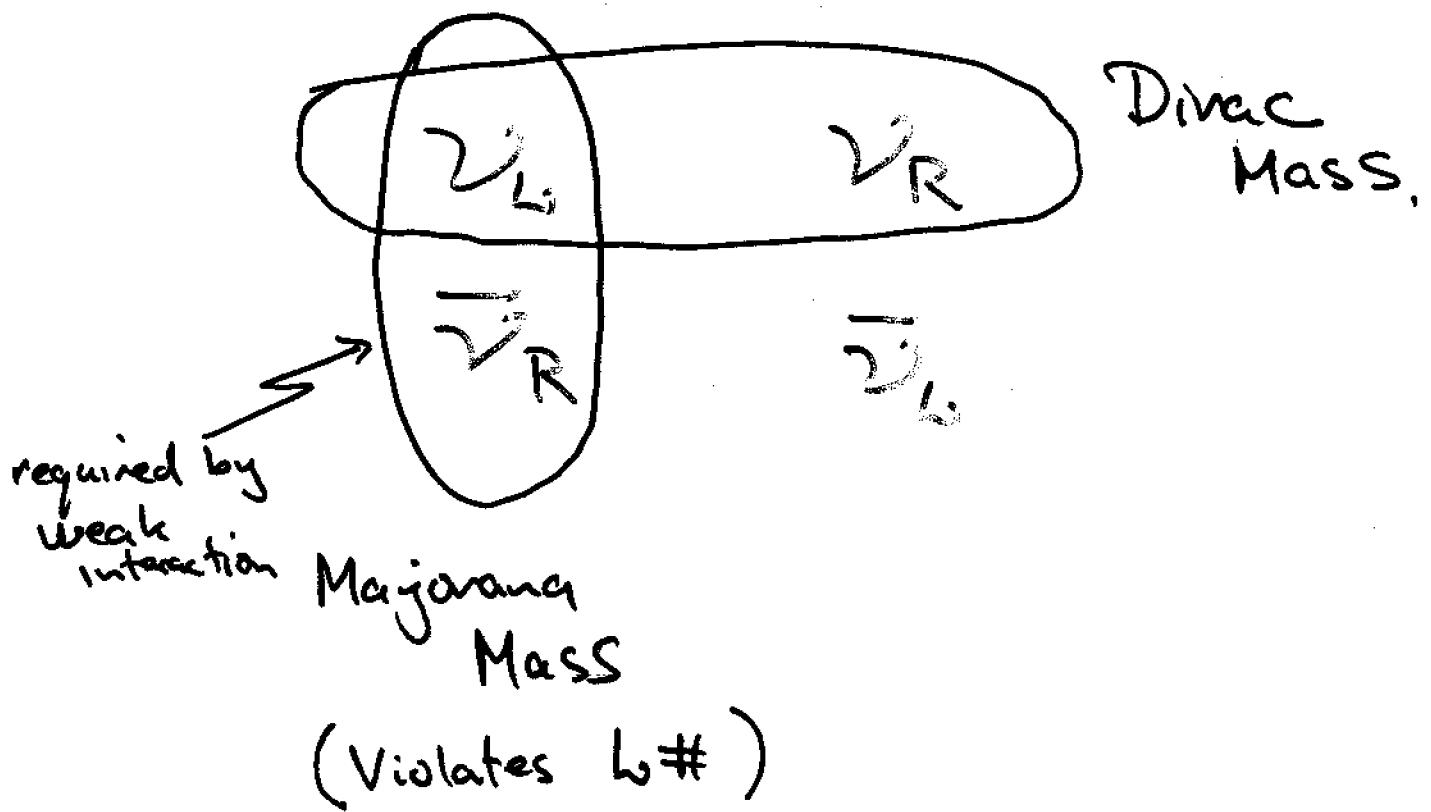


$$CPT \quad f_L \longleftrightarrow \bar{f}_R$$

At a minimum by CPT

$$\bar{f}_L, \bar{f}_R$$

For Neutral Fermions (neutrinos)<sup>3;</sup>



See-Saw Mechanism gives

$$\begin{array}{ccc} \nu_L & \text{and} & N_R \\ (\overleftarrow{\longrightarrow}) & & (\overrightarrow{\longrightarrow}) \\ \nu_R & & N_L \\ \text{light } \frac{m^2}{M} & & M \end{array}$$

Neutrinos being Majorana

is the minimal solution

to giving  $\nu$  mass.

$\nu = \bar{\nu}$  two states  $\nu_L$  and  $\nu_R$

---

If Dirac then there is a global  
symmetry (lepton #) which needs  
to be explained.

$\nu \neq \bar{\nu}$  and four states

weak  $\rightarrow \nu_L$  and  $\bar{\nu}_R$   
 $\nu_R$  and  $\bar{\nu}_L$

Suppression factor:

$$\left( \frac{\bar{M}_2}{E} \right)^2$$

chirality

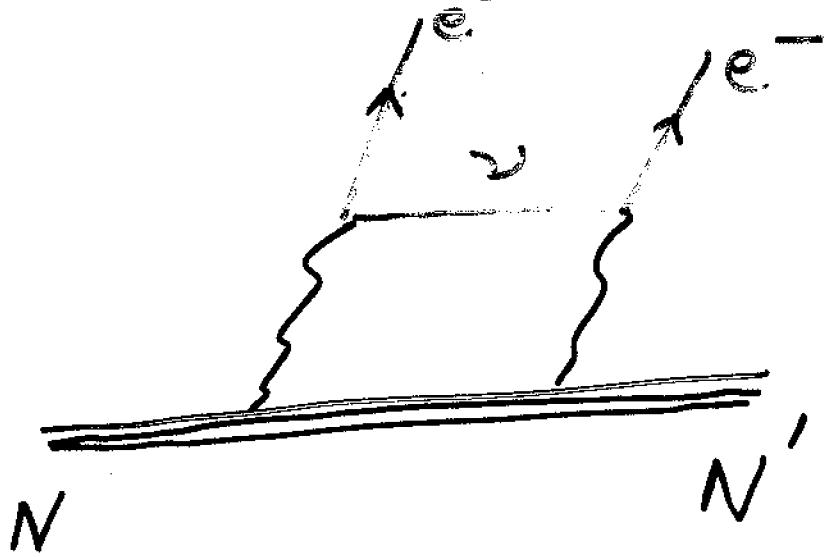
$\neq$

helicity  
for massive  
particles

$$\bar{M}_2 \approx 1 \text{ eV} \quad E = 1 \text{ GeV} \quad \left( \frac{M}{E} \right)^2 \approx 10^{-20}$$

BUT

~~$2\nu\beta\beta$  decay~~



Same Nucleus also has  $2\nu\beta\beta$  decay  
→ end point:

## WHAT WE DON'T KNOW:

- Majorana OR Dirac
- Absolute mass of highest neutrino.  
( except  $< \sim 1\text{eV}$  )
- Size of  $\Theta_{13}$ : ( $\nu_e$  in the "3" state.)  
 $\sin^2 \Theta_{13} < 0.03$
- Is  $\Theta_{23} = \text{or } \leq \frac{\pi}{4}$  the  $\mu \leftrightarrow \tau$  symmetric point.  
( Maximal mixing )  
 $0.35 < \sin^2 \Theta_{23} < 0.65$
- Sign of  $\Sigma_{\text{atm}}$  ( $= + - -$ )  
type of spectrum
- phase  $\delta \leftarrow$  if  $\neq 0$  leads to CP violation
- Number of light Neutrinos: 3 or are there more than 3