



the

abdus salam
international centre for theoretical physics

SMR.1508 - 34

SUMMER SCHOOL ON PARTICLE PHYSICS

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ASTROPARTICLE PHYSICS

Lecture III

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UHECR: TOP-DOWN SCENARIOS

NOT MUCH SPACE IS LEFT: $E > 1 \cdot 10^{20} \text{ eV}$

GENERAL REQUIREMENTS

- TO EXPLAIN EVENTS AT $E > 1 \cdot 10^{20} \text{ eV}$
- TO EXPLAIN CORRELATIONS WITH BL Lacs

OVERVIEW

• NEW SIGNAL CARRIERS

- STRONGLY INTERACTING NEUTRINOS (ν_{SI})
(EXTRA DIMENSIONS, INSTANTON EFFECT)
- LIGHT STABLE HADRONS (e.g. GLUINO)

MOTIVATION:

NO GZK CUTOFF, SMALL-ANGLE CLUSTERING,
CORRELATION WITH BL LACS

DRAWBACKS:

EXP STATUS AND FLUXES - FOR LIGHT GLUINOS
THEORETICAL STATUS - FOR ν_{SI}

• Z-BURSTS

(RESONANT Z-PRODUCTION: $\nu + \bar{\nu}_{DM} \rightarrow Z^0 \rightarrow \text{hadrons}$)

ELEM.-PARTICLE PHYSICS IS STANDARD (SM)

$$E_\nu^{\text{res}} = \frac{m_Z^2}{2m_\nu} \geq 1.8 \cdot 10^{22} \left(\frac{0.23}{m_\nu} \right) \text{ eV} \text{ AND } E_p > 1.8 \cdot 10^{23} \text{ eV}$$

(needed for ν -production)

• PROBLEM WITH SOURCES AND FLUXES:

E-M CASCADE LIMIT $I_\nu(E) < \frac{c}{4\pi} \frac{\omega_{\text{cas}}}{E^2}$ EXCLUDES

ASTROPHYSICAL SOURCES, TD and SHDR

• TOPOLOGICAL DEFECTS

RELIABLE OBJECTS NATURALLY PRODUCED BY
SYMMETRY BREAKING IN EARLY UNIVERSE.

RELIABLE PRODUCTION OF HE PARTICLES

MANY CANDIDATES: MONOPOLONIUM, MONOPOLES

CONNECTED BY STRINGS, ORDINARY STRINGS,
SUPERCONDUCTING STRINGS, CUSPS, VORTOOLS, NECKLACES

MOST OF THEM ARE EXCLUDED OR DISFAVoured

- SUPERHEAVY DM

QUASISTABLE PARTICLES WITH $M_x > 1 \cdot 10^{13} \text{ GeV}$
PRODUCED AT INFLATION AND ACCUMULATED
IN GALACTIC HALO.

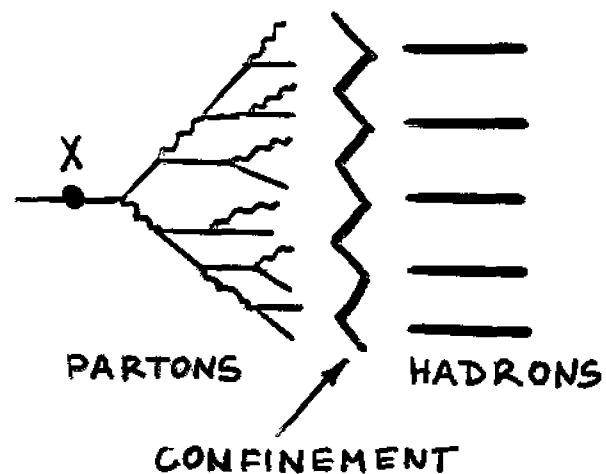
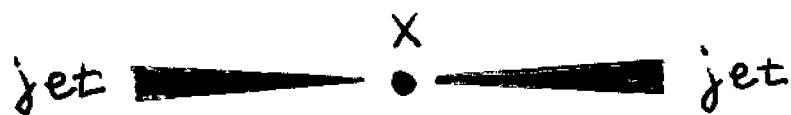
NO GZK CUTOFF

- LORENTZ-INVARIANCE VIOLATION

ENERGY OF CMB PHOTON IN CM-SYSTEM IS ΓE_{CMB}
VIOLATION OF LORENTZ-INVARIANCE CAN RESULT
IN SMALLER E_{cm} AND GZK CUTOFF IS SHIFTED.

- MOST RADICAL POSSIBILITY

X DECAY: ENERGY SPECTRUM



SPECTRUM IS DOMINATED BY ν 'S AND γ 'S
RATIO $\gamma/N \sim 2-3$

CALCULATION OF SPECTRUM

MC SIMULATION

TWO STEPS:

- PERTURBATIVE REGIME



$$x = \frac{2E}{M_X}, Q_0 \sim 1 \text{ GeV} \quad (\text{min virtuality})$$

- HADRONIZATION

$$D_i^h(x, M_x) = \sum_{j=q,g} \int_x^1 \frac{dz}{z} D_i^j\left(\frac{x}{z}, M_x, Q_0\right) f_j^h(z, Q_0)$$

$f_j^h(x, Q_0)$ IS FOUND FROM $D_q^h(x, M_z)$

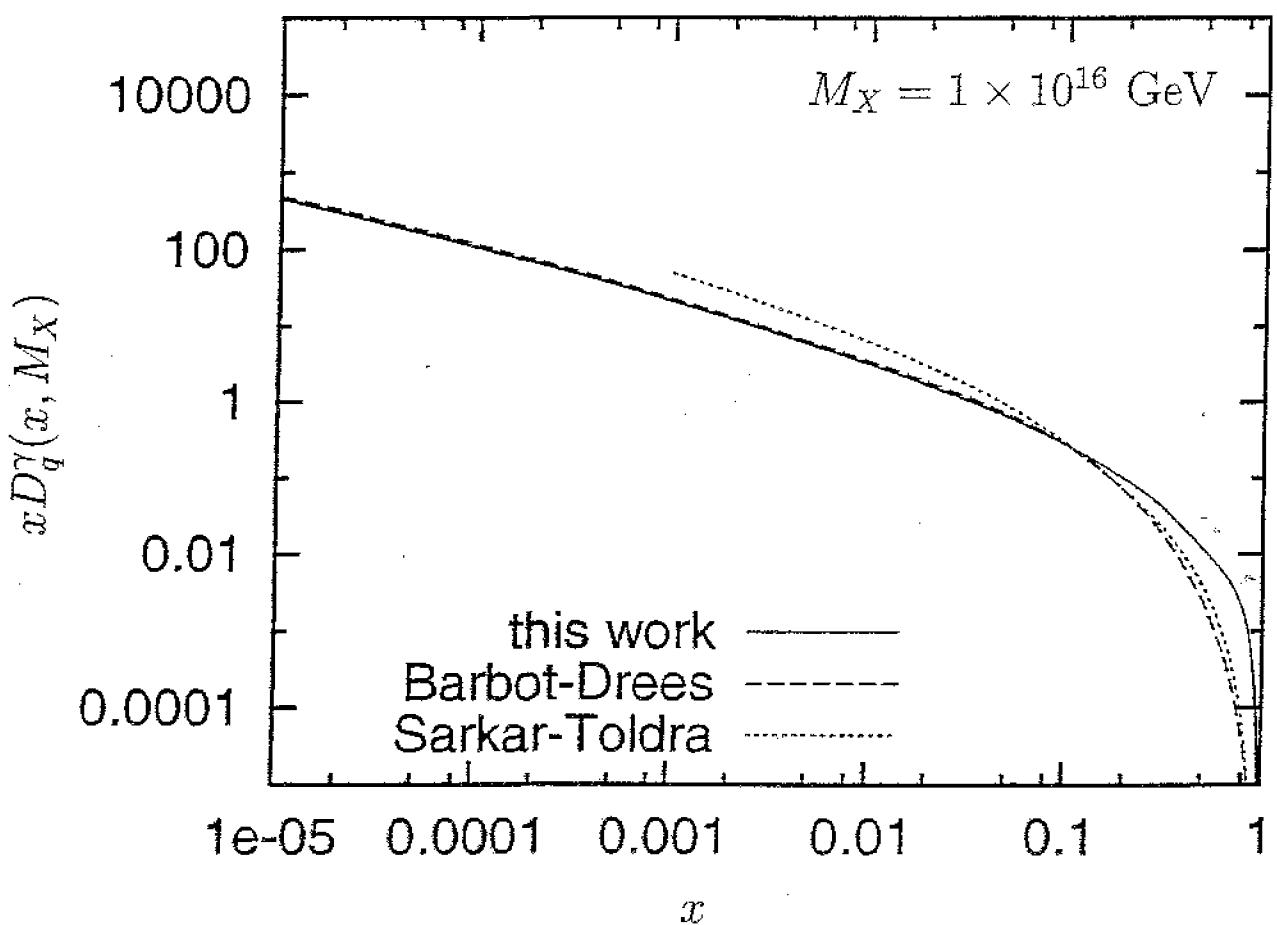
DGLAP EVOLUTION

$$\partial_t D_i^h = \sum_j \frac{\alpha_s(t)}{2\pi} P_{ij} \otimes D_j^h\left(\frac{x}{z}, t\right)$$

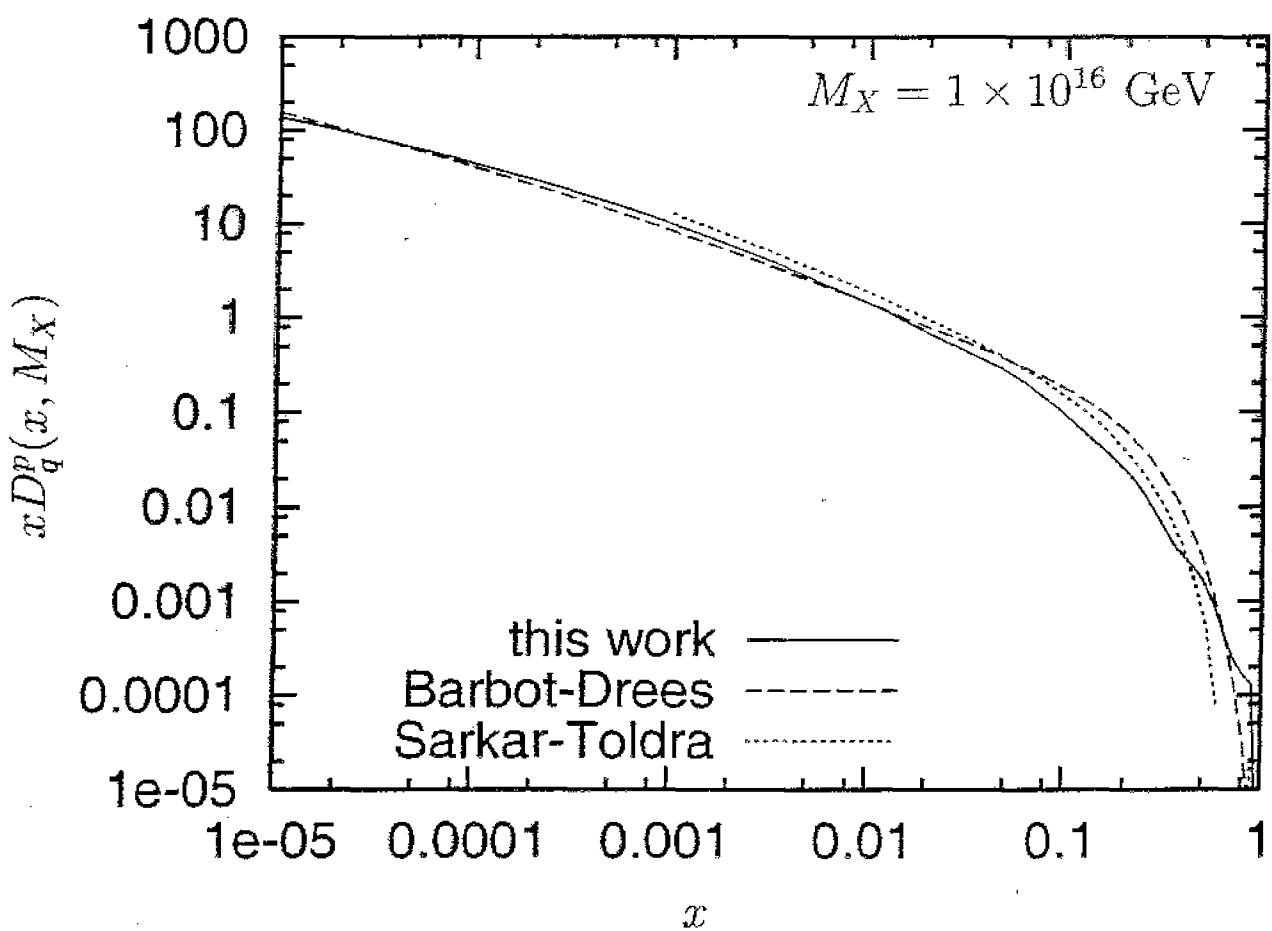
$t = \ln s/s_0$ P_{ij} IS SPLITTING FUNCTION

$$f \otimes g = \int_z^1 \frac{dx}{x} f(x) g\left(\frac{x}{z}, t\right)$$

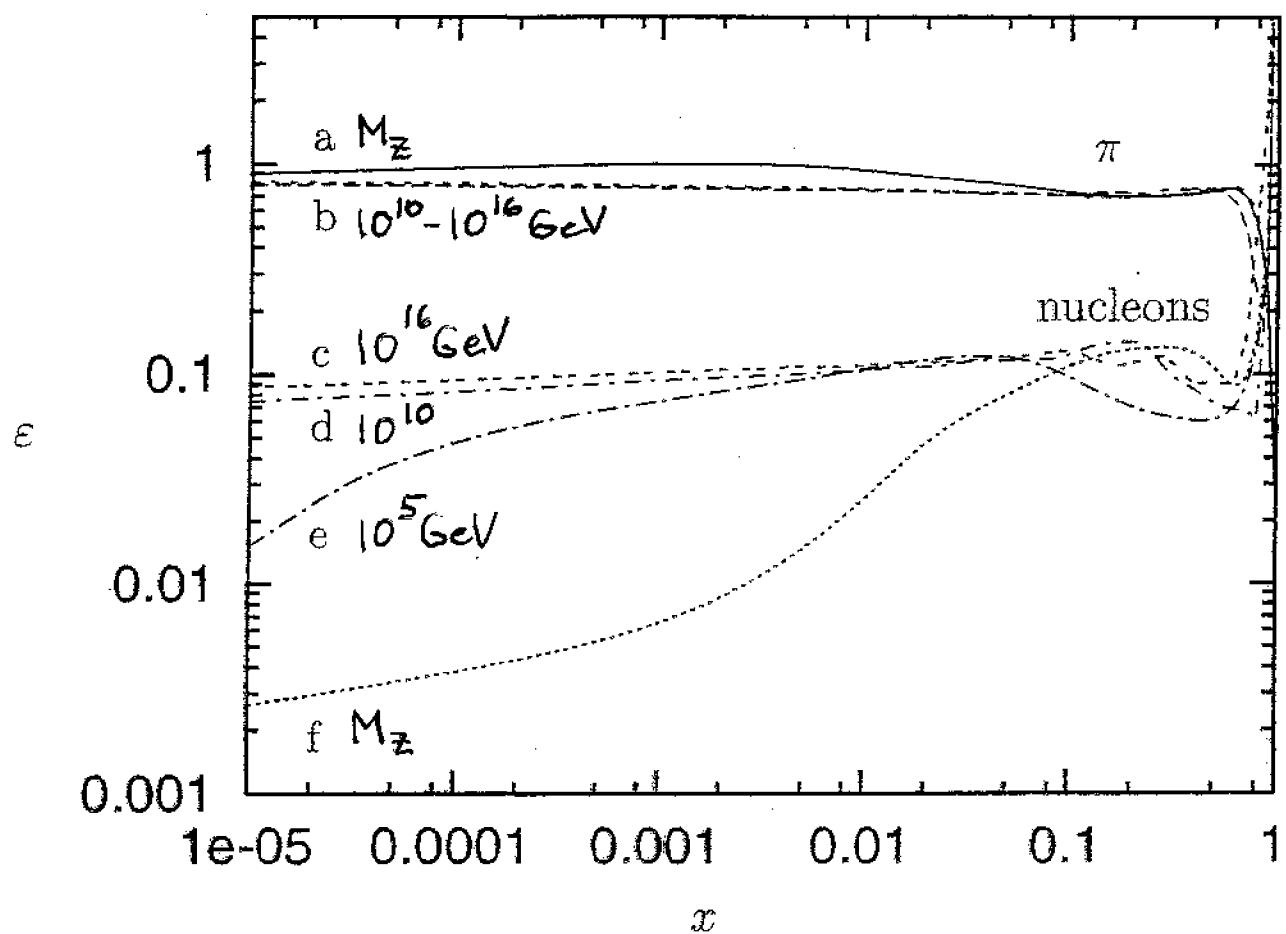
SPECTRA OF PHOTONS
AT THE DECAY OF SH PARTICLE



SPECTRA OF PROTONS
AT THE DECAY OF SH PARTICLE



FRACTION OF PIONS AND NUCLEONS
IN THE DECAY SPECTRA OF SH PARTICLES



Z-BURSTS

Fargion et al 1999
Weiler 1999

$$v + \bar{\nu}_{DM} \rightarrow Z^0 \rightarrow \text{hadrons}$$

$$E_0 = \frac{m_Z^2}{2 m_\nu} = 1.8 \cdot 10^{22} \left(\frac{0.23 \text{ eV}}{m_\nu} \right) \text{ eV}$$

$$\sigma_{\nu f} = \frac{12 \pi}{m_Z^2} \frac{\Gamma_\nu \Gamma_f}{(E_c - m_Z)^2 + \Gamma_Z^2 / 4}$$

$$\gamma_Z = \frac{E_0}{m_Z} = \frac{m_Z}{2 m_\nu} = 2.0 \cdot 10^{11} \left(\frac{0.23 \text{ eV}}{m_\nu} \right)$$

$$\dot{n}_{Z^0} = 4\pi n_{\bar{\nu}_i} \int I_{\nu_i}(E) \sigma(E) dE = 4\pi \sigma_t n_{\bar{\nu}_i} I_{\nu_i}(E_0) E_0$$

$$\sigma_t = 48 \pi f_\nu G_F = 1.29 \cdot 10^{-32} \text{ cm}^2$$

$$f_\nu = \frac{\Gamma_\nu}{G_F m_Z^3} = 0.019$$

PROBLEMS OF Z-BURST MODELS

$$E_\nu = E_0 = 1.8 \cdot 10^{22} \left(\frac{0.23 \text{ eV}}{m_\nu} \right) \text{ eV}$$

$$E_p > 20 E_\nu = 3.6 \cdot 10^{23} \left(\frac{0.23 \text{ eV}}{m_\nu} \right) \text{ eV}$$

IS TOO HIGH FOR ACCELERATION

DECAYS IN THE EXTRAGALACTIC SPACE (TD's, SUPERHEAVY DM PARTICLES) ARE ACCCOMPANIED BY e-m CASCADE

$$I_\nu(E) < \frac{c}{4\pi} \frac{\omega_{cas}}{E^2}; \quad \omega_{cas} \leq 2 \cdot 10^{-6} \frac{\text{eV}}{\text{cm}^3} \text{ (EGRET)}$$

$I_\nu(E_0) < \frac{c}{\pi} \omega_{cas} \frac{m_\nu^2}{m_Z^4}$ IS TOO SMALL TO PRODUCE THE OBSERVED FLUX OF UHECR

$$\dot{n}_Z = 4\pi \sigma_t n_{\nu_i} I_\nu(E_0) \cdot E_0 \leq \sigma_t n_{\nu_i} c \omega_{cas} / E_0$$

$$\dot{n}_\gamma(E_\gamma) < \frac{E_0}{E_\gamma} \dot{n}_Z$$

$$F_\gamma(>E_\gamma) = \frac{1}{4\pi} \dot{n}_\gamma(E_\gamma) \cdot R_\gamma^{\text{abs}} = \frac{1}{4\pi} \sigma_t n_{\nu_i} R_\gamma^{\text{abs}} \frac{c \omega_{cas}}{E_\gamma}$$

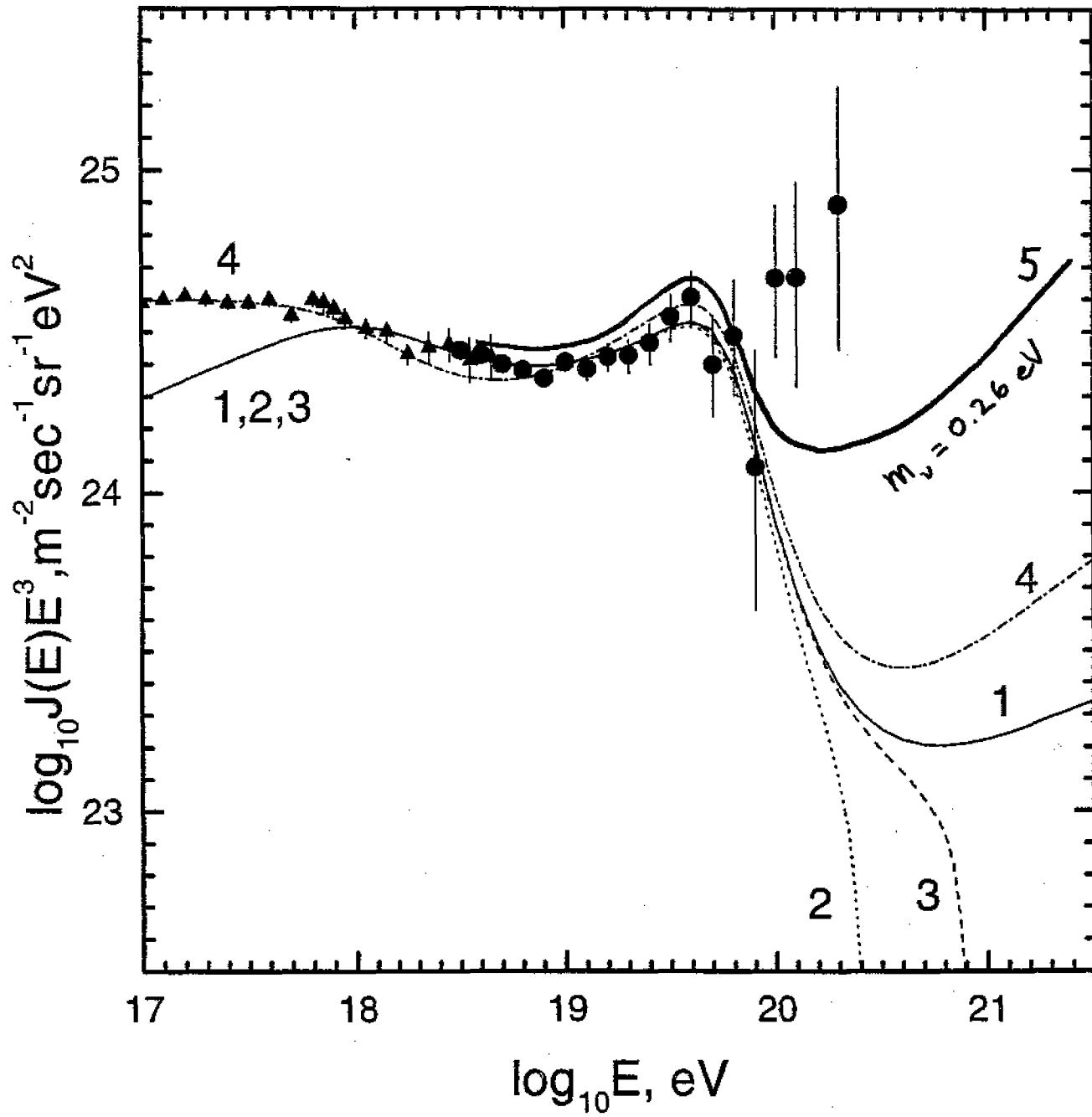
$$n_{\nu_i} = 56 \text{ cm}^{-3} \quad R_\gamma^{\text{abs}} \leq 5 \text{ kpc} \quad E_\gamma = 8 \cdot 10^{19} \text{ eV}$$

$$F_\gamma(>E_\gamma) \leq 5.9 \cdot 10^{-22} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$$

i.e. 2 ORDERS OF MAGN LESS THAN OBSERVED

Z-BURSTS

Fodor, Katz, Ringwald 2001

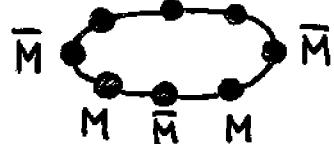


TOPOLOGICAL DEFECTS

- SYMMETRY BREAKING IN EARLY UNIVERSE RESULTS IN PHASE TRANSITIONS (D.A. KIRZHNITZ 1972). THEY ARE ACCCOMPANIED BY TOPOLOGICAL DEFECTS.
- DEPENDING ON SYMMETRY BREAKING, DEFECTS CAN BE IN FORM OF SURFACES (DOMAIN WALLS), LINES (STRINGS), AND POINTS (MONOPOLES).

TD OF INTEREST FOR UHECR

- MONOPOLES: $G \rightarrow H \times U(1)$
- ORDINARY STRINGS: $U(1)$ BREAKING
- MONOPOLES CONNECTED BY STRINGS: $G \rightarrow H \times U(1) \rightarrow H \times \mathbb{Z}_n$
e.g. NECKLACES $\mathbb{Z}_n = \mathbb{Z}_2$
- SUPERCONDUCTING STRINGS



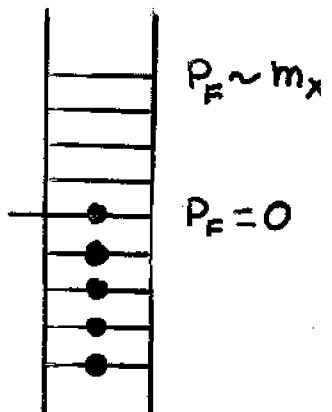
PRODUCTION OF UHECR

- ANNIHILATION OF MONOPOLES
 $M + \bar{M} \rightarrow \text{PARTON CASCADE} \rightarrow \text{PIONS} \rightarrow \text{PHOTONS} + \text{NEUTRINOS}$
- SUPERCONDUCTING STRINGS

$$\frac{dP}{dt} = eE \quad P = eEt \quad P \sim m_x$$

$$J = e^2 E t \quad J \sim em_x$$

$X \rightarrow \text{PARTON CASCADE} \rightarrow \text{PIONS} \rightarrow \gamma + \nu$



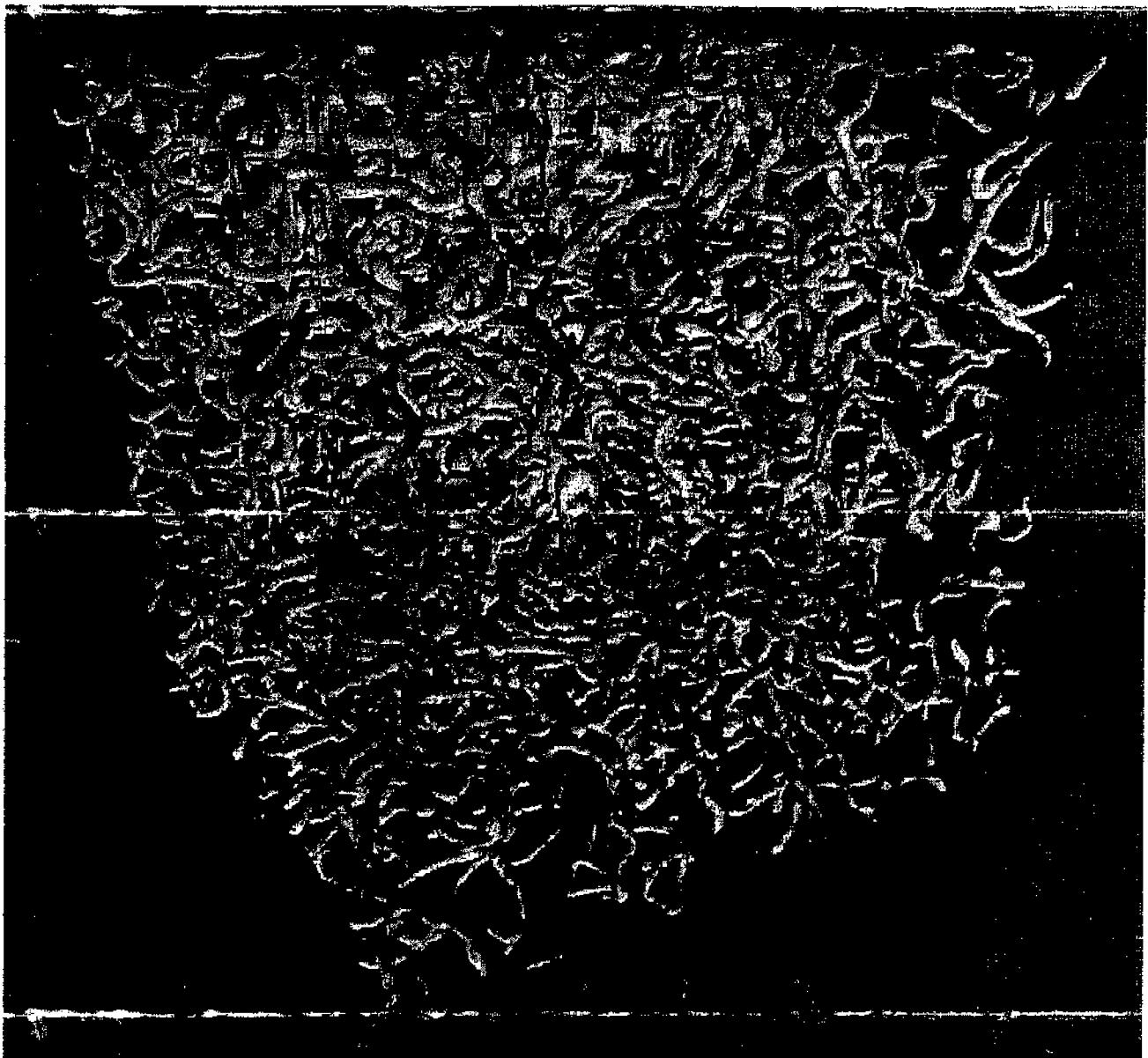


Figure 1. When a body of water freezes over, ice begins forming in many places independently. As the frozen areas grow, they create irregular boundaries between areas of crystalline ice that have different orientations. Phase transitions in the early universe can likewise create topological defects, as regions of space where the transition has proceeded independently run up against each other. The author discusses the possibility that today's universe may, as a result, be riddled with high-energy, gravitationally powerful filaments of primordial material called cosmic strings. Above, a simulation shows a step in the evolution of one region of the universe, in which cosmic strings form from a random initial distribution of phases of a hypothetical field called a Higgs field. (Image above courtesy of Paul Shellard, Cambridge University.)



Figure 2. In a simple model of symmetry breaking, the initial symmetric ground state of the Higgs field (yellow dot) can fall into the left- or right-hand valley of an energy potential (light and dark dots), just as a diner can choose a wine glass from either the left or the right side. In a cosmic phase

FIGURES

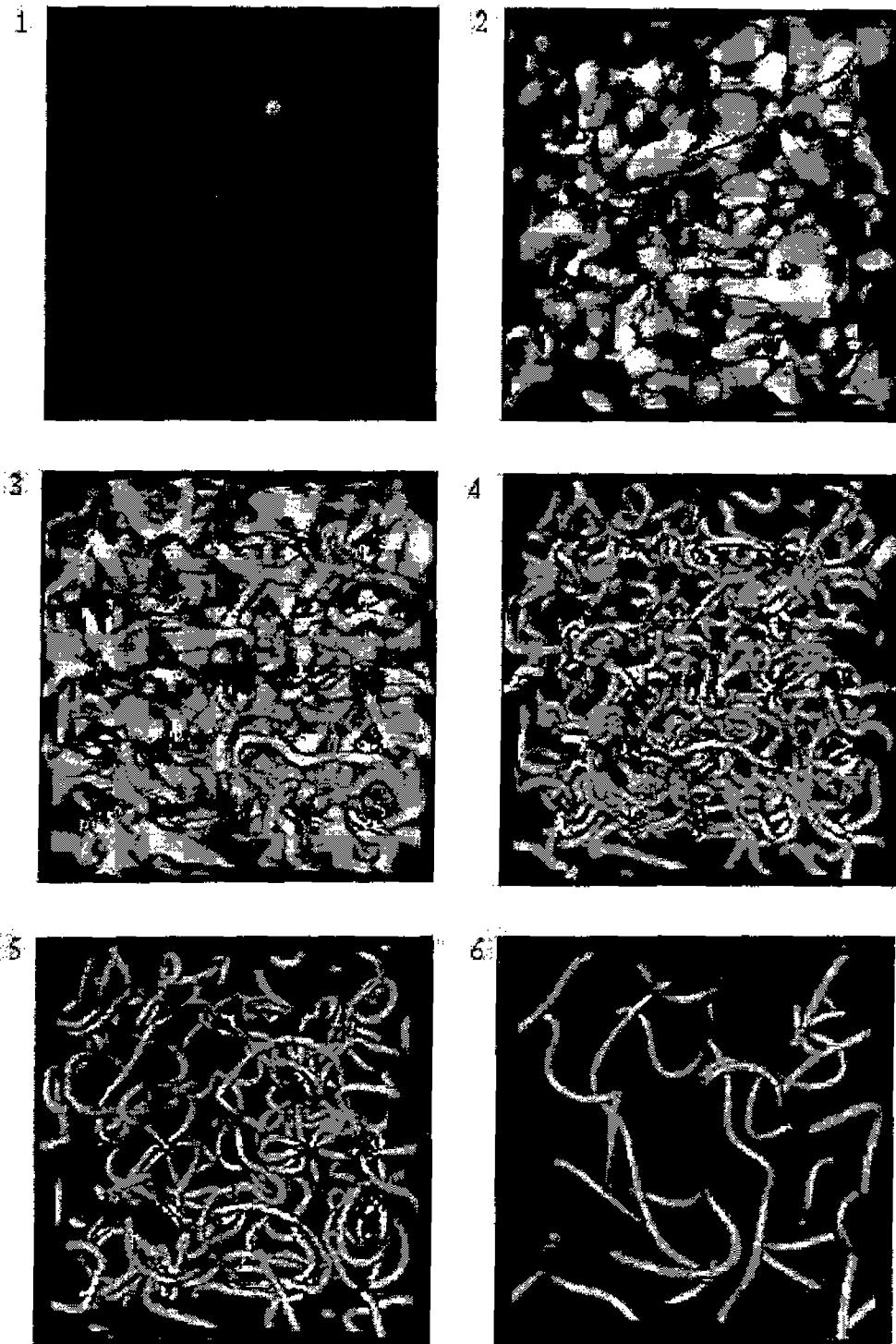


FIG. 1. String network formation in the Abelian-Higgs model using gradient flow (diffusive) evolution. Given random initial phases, a symmetry-breaking phase transition occurs, strings form and then begin to evolve in a scale-invariant manner (the correlation length is $L \propto t^{1/2}$). This dissipative evolution is used to create the initial configuration for a string network with a specified L for subsequent relativistic evolution.

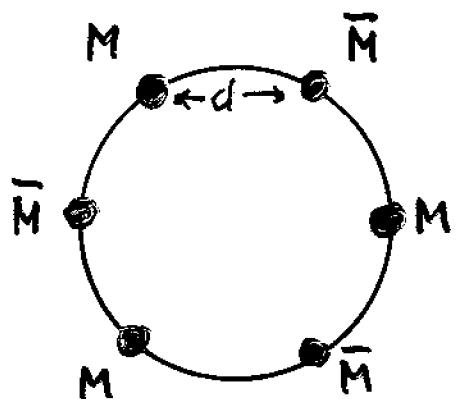
NECKLACES

$$G \rightarrow H \times U(1) \rightarrow H \times \mathbb{Z}_2$$

SYMM BREAKING
SCALE

$$m \sim 4\pi \ell_m/e$$

$$\mu \sim 2\pi \ell_s^2$$



$$r = \frac{m}{\mu d}$$

$$n_x \sim \frac{r^2 M}{t^3 m_x}$$

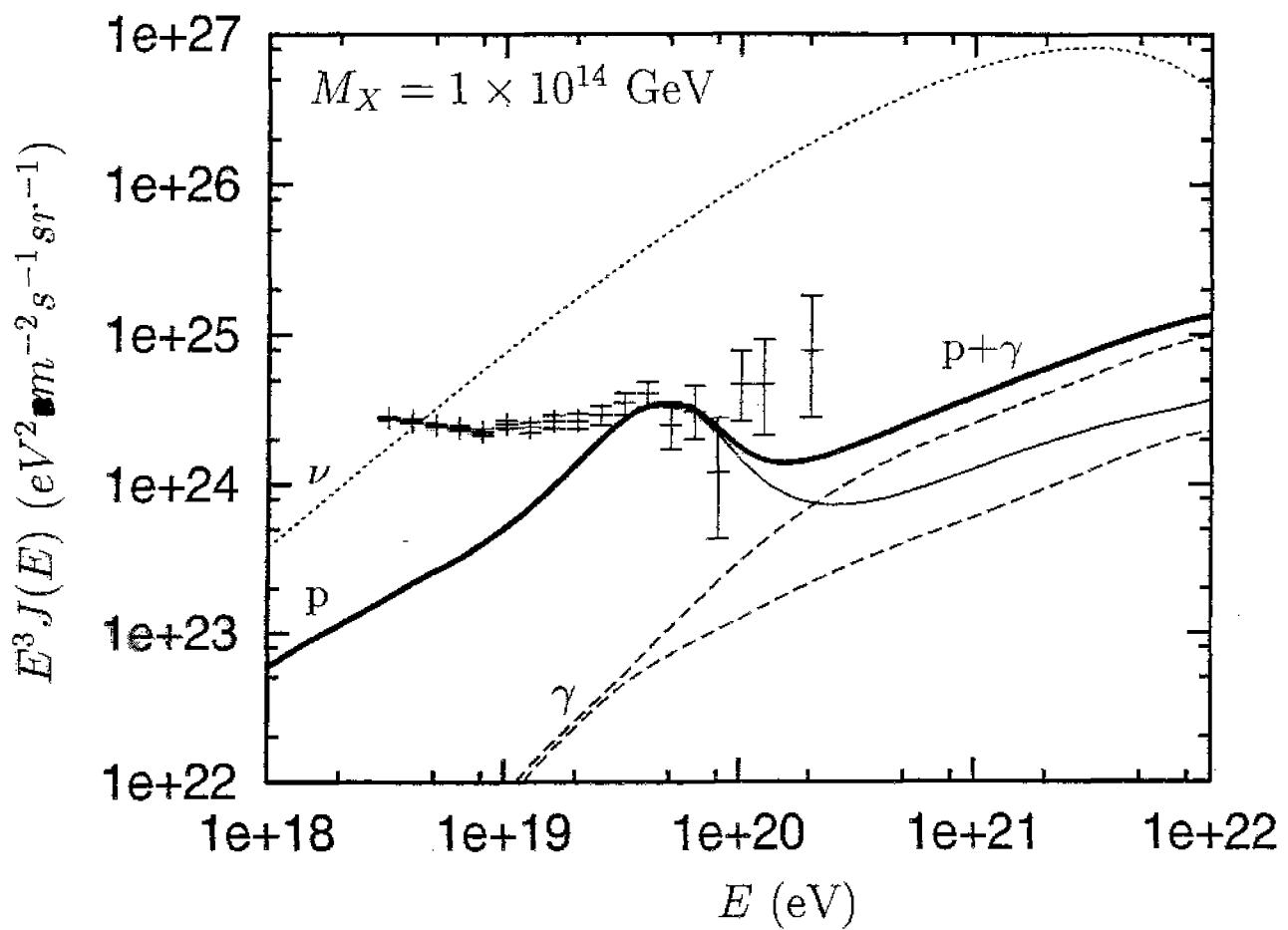
$$\omega \sim m_x t = n_x \sim r^2 \frac{\mu}{t^2}$$

$$\omega_{cas} = \frac{3}{4} f_\pi r^2 \frac{\mu}{t_0} \approx 2 \cdot 10^{-6} \frac{eV}{cm^3} \text{ (EGRET)}$$

$$\text{SEPARATION OF NECKLACES } D \sim \frac{t_0}{\sqrt{F}}$$

$$D \sim \left(\frac{3 f_\pi \mu}{4 t_0^2 \omega_{cas}} \right)^{1/4} t_0 > 10 \left(\frac{1}{10^6 GeV^2} \right)^{1/4} kpc$$

NECKLACES



UHECR FROM SUPERHEAVY DM

- PRODUCTION AT POST-INFLATION EPOCHS
 - (i) VARIETY OF MECHANISMS AT PREHEATING
 - (ii) PRODUCTION BY TIME-VARYING GRAVITATIONAL FIELD
 - NO COUPLING WITH INFLATON, X CAN BE STERILE!
 - PARTICLE MASS m_X COUPLES TO EXPANDING METRIC

$$\mathcal{L} \sim \frac{1}{2} \xi R X^2$$

$$\xi = \frac{1}{6} \text{ CONFORMAL COUPLING}$$

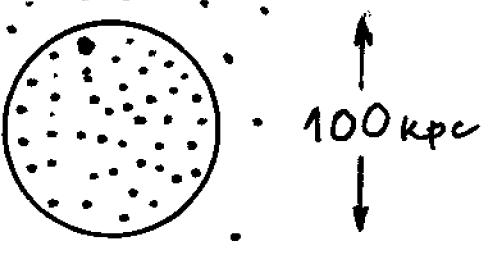
CREATION OCCURS WHEN $H(\epsilon) \sim m_X$,

SINCE $H(\epsilon) \leq m_\phi \sim 10^{13} \text{ GeV}$, $m_X \leq 10^{13} \text{ GeV}$

e.g. $m_X \sim 2-3 \cdot 10^{13} \text{ GeV}$ RESULTS IN $\Omega_X h^2 \sim 0.1$

- LIFETIME $\tau_X > 10^{10} \text{ yr}$
 - e.g. DISCRETE-GAUGE SYMMETRY PROTECTION
 - MANY PARTICLE CANDIDATES

- ACCUMULATION IN THE HALO

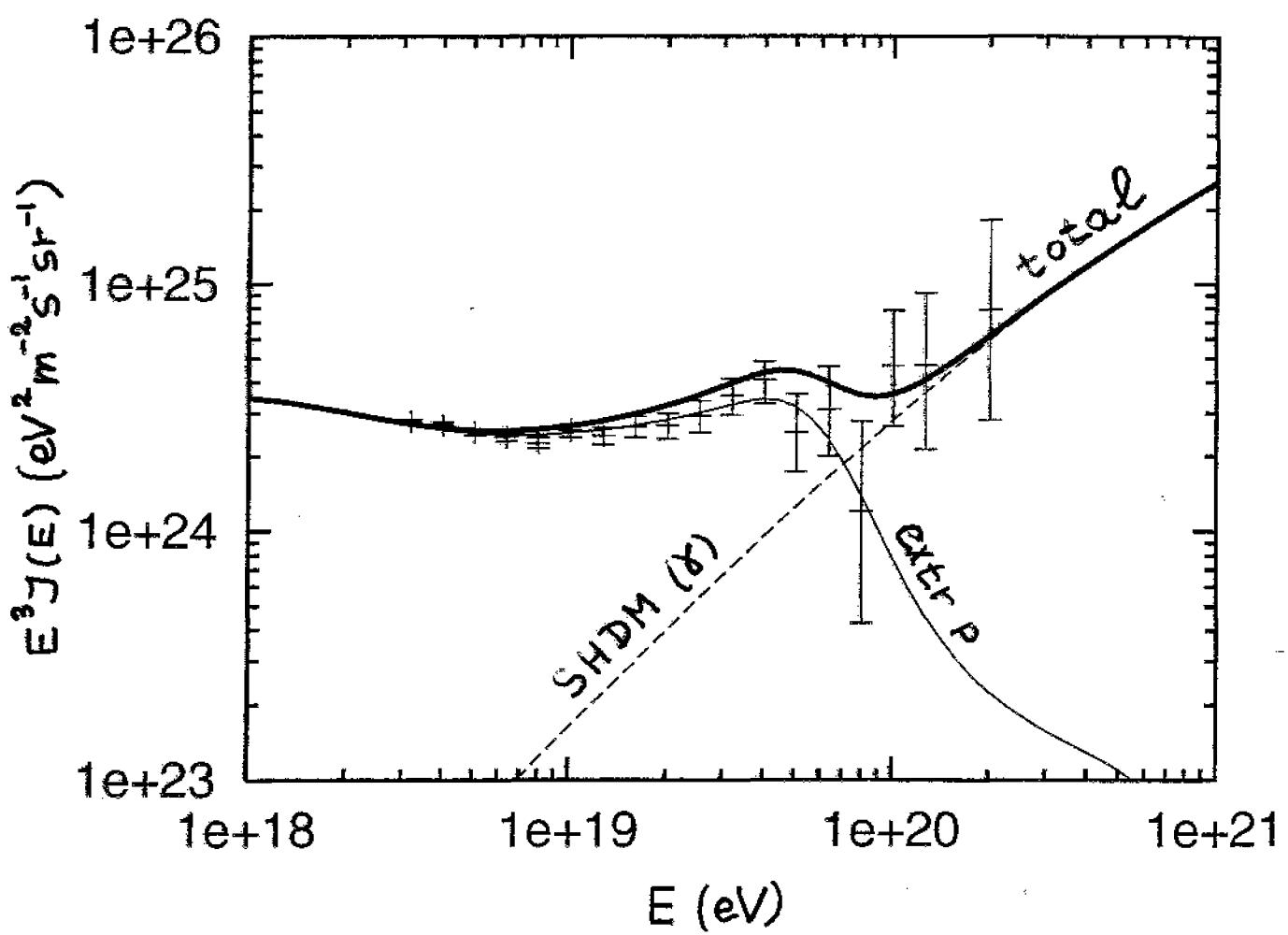


$$\frac{n_X^{\text{halo}}}{n_X^{\text{extr}}} = \frac{s_{\text{CDM}}^{\text{obs}}}{s_{\text{CDM}} s_{\text{cr}}} \approx 2 \cdot 10^5$$

NO GZK CUTOFF!

- CALCULATED ANGULAR CLUSTERING AGREES WITH DATA.
- PREDICTIONS:
 - PRIMARIES ARE γ 's, ANISOTROPY

SHDM



OBSERVATIONAL CONSEQUENCES

- SH RELIC PARTICLES ARE ACCUMULATED IN THE GALACTIC HALO ($n_x^{\text{halo}}/n_x^{\text{ext}} \sim 10^5$) AND THUS UHECR PRODUCED AT THEIR DECAYS HAVE NO GZK CUTOFF.
- UHE PHOTONS (FROM $\pi^0 \rightarrow 2\gamma$) HAVE FLUX ~ 2 TIMES HIGHER THAN NUCLEONS.
- ANISOTROPY DUE TO NON-CENTRAL POSITION OF THE SUN IN THE HALO.

TWO DISTRIBUTIONS OF **DM** IN HALO:

ISO: $\rho(R) = \frac{\rho_0}{1 + (R/R_c)^2}$

NFW: $\rho(R) = \frac{\rho_0}{R/R_s(1+R/R_s)^2}$

(i) PREDICTIONS FOR EXISTING ARRAYS

V.B. and A.MIKHAIEV 1998

G.A.M.TANCO and A.A.WATSON 1998

(ii) RATIO GC/AC

V.B., P.BLASI, A.VILENKO 1998

BENSON, SHIAŁKOWSKI, WOLFENDALE 1998

WITHIN $\theta = 20^\circ$ THIS RATIO IS 3 - 10 FOR DIFFERENT MODELS: PROBLEM FOR THE AUGER DETECTOR.

- ANGULAR CLUSTERING DUE TO CLUMPINESS OF HALO

VIOLATION OF LORENTZ INVARIANCE (LI)

MOTIVATION:

- EXP DATA: PROBABLY NOT (?)
- IS SUPERWEAK LI BREAKING THEORETICALLY FEASIBLE
PROBABLY YES (!)
- ABSOLUTE LORENTZ FRAME AND LI BREAKING.

BREAKING OF LI IMPLIES EXISTENCE OF THE ABSOLUTE LORENTZ FRAME.

DOES EXISTENCE OF THE ABSOLUTE FRAME IMPLY THE VISIBLE LI BREAKING?

e.g. EXISTENCE OF THE FRAME WHERE CMBR IS ISOTROPIC?

NO, IF WE TAKE INTO ACCOUNT THE INTERACTION WITH MICROWAVE PHOTONS.

SPONTANEOUSLY BROKEN LI

- EQUATION OF MOTION REMAIN LORENTZ INVARIANT
- VIOLATION OCCURS SPONTANEOUSLY IN SOLUTIONS, e.g. DUE TO VEV'S.

$$L_{\text{int}} = \frac{e}{M^2} \bar{\Psi} \gamma_\mu \partial_\nu \Psi T_{\mu\nu}$$

Ψ IS FIELD OF VISIBLE SECTOR,

$T_{\mu\nu}$ IS SUPERHEAVY FIELD FROM HIDDEN SECTOR

STRING INTERACTION SET NONZERO VEV FOR T_{00}

$$\langle T_{00} \rangle = v^2,$$

PRODUCING LORENTZ-NONINvariant VACUUM
CONDENSATE ANALOGUE OF CMBR

SPONTANEOUSLY BROKEN L.I. AND MODIFIED DISPERSION RELATIONS

$$\mathcal{L} = \frac{1}{2} i \bar{\Psi} \gamma_\mu \overleftrightarrow{\partial}_\mu \Psi - m \bar{\Psi} \Psi + \mathcal{L}_{\text{int}}$$

e.g. $\mathcal{L}_{\text{int}} = \frac{e}{M^2} T_{\mu\nu} \bar{\Psi} \gamma_\mu \partial_\nu \Psi$

DIRAC EQUATION

$$(i \gamma_\mu \partial_\mu - m - \frac{e}{M^2} T_{\mu\nu} \gamma_\mu \partial_\nu) \Psi = 0$$

AFTER SSB $\langle T_{00} \rangle = v^2$ ASSOCIATE KLEIN-GORDON EQ:

$$(\partial_\mu^2 - m^2 + \frac{e}{M^2} v^2 \partial_0^2) \Psi = 0$$

IN MOMENTUM SPACE $i \partial_\mu \rightarrow p_\mu$

$$p_\mu^2 - m^2 + \frac{e}{M^2} v^2 E^2 = 0$$

THUS, **L.I.** BREAKING RESULTS IN THE MODIFIED DISPERSION RELATION (COLEMAN, GLASHOW 1997)

MODIFIED DISPERSION RELATIONS AND THRESHOLD OF REACTIONS.

$$\gamma + \gamma \rightarrow e^+ + e^-$$

$$E_\gamma \quad \epsilon_\gamma \quad E_e \quad E'_e \quad E_\gamma \gg \epsilon_\gamma$$

$$E_\gamma \xrightarrow{K_\mu} \xleftarrow{K'_\mu} \epsilon_\gamma$$

$$K_\mu + K'_\mu = P_\mu + P'_\mu$$

$$K_\mu^2 + K'_\mu^2 + 2 K_\mu K'_\mu = P_\mu^2 + P'_\mu^2 + 2 P_\mu P'_\mu$$

IN CASE $K_\mu^2 = K'_\mu^2 = 0$ AND $P_\mu^2 = P'_\mu^2 = m_e^2$

$$(E_\gamma \epsilon_\gamma)^{\frac{1}{2}} \gtrsim m_e$$

IN CASE OF MODIFIED DISPERSION RELATIONS

$$P_\mu^2 = E_e^2 - \vec{P}^2 \approx m_e^2 + e \frac{V^2}{M^2} E^2$$

WHEN $E > M m_e / V$ THRESHOLD IS SHIFTED

OTHER ASTROPHYSICAL TESTS OF SPECIAL RELATIVITY

- CONSTANCY OF LIGHT VELOCITY (e.g. IN GRBS).
J. Ellis et al 1999
- HE PHOTON DECAY. S. GLASHOW 1997
- VACUUM CHERENKOV RADIATION. S. GLASHOW 1997

CONCLUSIONS

- OBSERVED UHECR SPECTRA AT $1 \cdot 10^{18} - 8 \cdot 10^{19}$ eV HAVE SIGNATURES OF INTERACTION OF EXTRAGALACTIC PROTONS WITH CMB IN THE FORM OF THE DIP (DUE TO $P + \gamma_{CMB} \rightarrow P + e^+ + e^-$) AND OF THE BEGINNING OF THE GZK CUTOFF. THE CORRELATION WITH BL LACS INDICATE AGN AS UHECR SOURCES.
- NOT MUCH SPACE IS LEFT TO "EXOTIC" PHYSICS. UHECR PROBLEM IS REDUCED TO THE EVENTS AT $E > 1 \cdot 10^{20}$ eV. EXCESS OF THE AGASA EVENTS OVER PREDICTED GZK CUTOFF IS AN INDICATION TO A NEW PHYSICS. IN 2 YEARS AUGER DATA WILL CLARIFY THE SITUATION.
- THREE TOP-DOWN MODELS (TD, SHDM AND Z-BURSTS) PREDICT PHOTON DOMINANCE AT $E \geq 1 \cdot 10^{20}$ eV. THIS PREDICTION WILL BE TESTED BY AUGER.
- SPECTRUM SHAPE ($F(E) \sim E^{-1.9}$) IS NOW A SIGNATURE OF SHDM MODEL. THIS PREDICTION IS IN AGREEMENT WITH THE AGASA EXCESS.
- SHDM MODEL PREDICTS ANISOTROPY CENTERED BY GC. THIS PREDICTION WILL BE TESTED BY AUGER LOCATED IN SOUTHERN HEMISPHERE.
THERE IS A CHANCE OF DISCOVERY OF NEW FUNDAMENTAL PHYSICS IN UHECR!