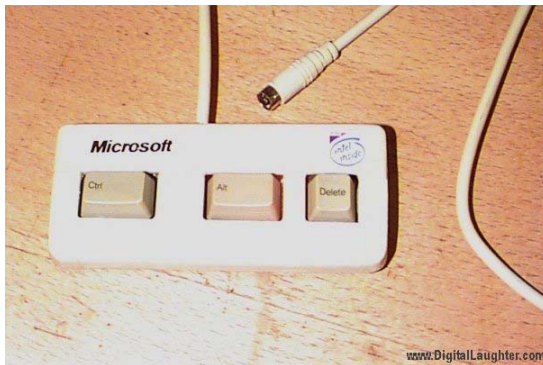


# Dirty Black Hole Evolutions with Creative Engineering

Pablo Laguna

Center for Gravitational Physics and Geometry  
Center for Gravitational Wave Physics  
Penn State University



## Examples:

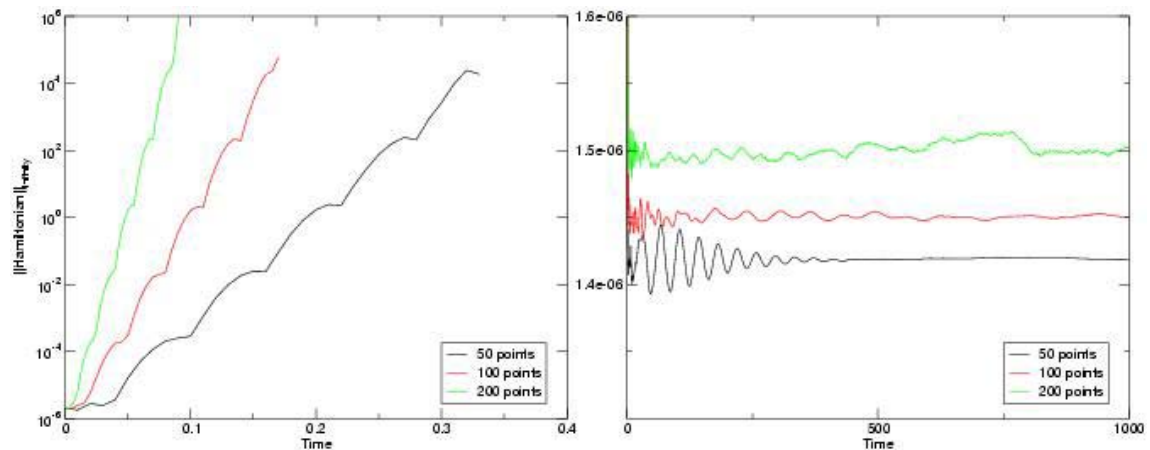
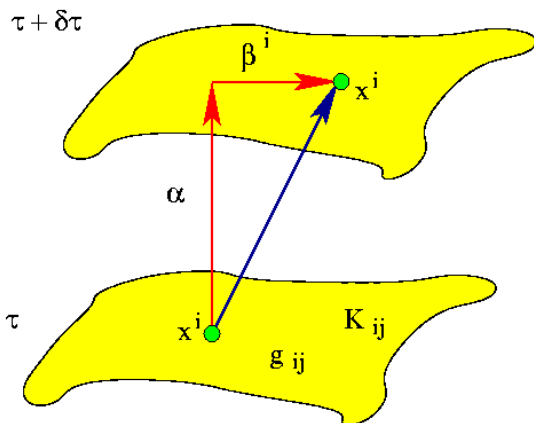
- Conformally flat initial data
- Thin-sandwich data
- Hydro-without-hydro
- Bonazzola's talk

## This Talk:

- 3+1 Conformal Formulations
- Black Hole Excision
- Wobbling Black Hole
- Partial BBH Orbit
- Constraint Violating Modes

$$\begin{aligned} R_{\mu\nu} = 0 & \quad \left\{ \begin{aligned} \partial_o g_{ij} &= -2\alpha K_{ij} \\ \partial_o K_{ij} &= -\nabla_i \nabla_j \alpha + \alpha(R_{ij} - 2K_i^k K_{kj} + K_{ij} K) \end{aligned} \right. \\ G_{\mu\nu} = 0 & \quad \left\{ \begin{aligned} R + K^2 - K_{ij} K^{ij} &= 0 \\ \nabla_j K^{ij} - \nabla^i K &= 0 \end{aligned} \right. \end{aligned}$$

$$\partial_o \equiv \partial_t - \mathcal{L}_\beta$$



Apples with Apples I

$$\begin{aligned}\partial_t \phi &= \eta \\ \partial_t \eta &= \Delta \phi + \nabla_i \phi \nabla^i \phi - \eta^2\end{aligned}$$

- Re-write equations in a manifestly hyperbolic form.
  - Well-posedness will tell you that deviations are bounded by  $\exp(\text{ct})$
  - Theorems and rigorous mathematics are available.
  - Formal excision and outer boundaries.
  - One can talk about initial boundary value problem.
- Re-write equations with **some** hyperbolic flavor.
  - Identify and eliminate troublesome non-linear terms.
  - Sorry, no theorems.
  - Excision and outer boundaries become an art (or cooking).

## BSSN Formulation

$$g_{ij} = g^{1/3} \hat{g}_{ij} = e^{4\Phi} \hat{g}_{ij}$$

$$g = e^{12\Phi}$$

$$K_{ij} = e^{4\Phi} \hat{A}_{ij} + \frac{1}{3} g_{ij} K$$

$$\hat{\Gamma}^i = \hat{g}^{jk} \hat{\Gamma}^i_{jk} = -\partial_j \hat{g}^{ij}$$

$$\partial_o \Phi = -\frac{1}{6} \hat{\alpha} \hat{K} e^{-6(n+k)\Phi}$$

$$\partial_o \hat{g}_{ij} = -2 \hat{\alpha} \hat{A}_{ij} e^{-6(n+a)\Phi}$$

$$\partial_o \hat{K} = -\nabla_i \nabla^i \alpha + \alpha \left( \hat{A}_{ij} \hat{A}^{ij} + \frac{1}{3} \hat{K}^2 \right)$$

$$\partial_o \hat{A}_{ij} = e^{-4\Phi} [-\nabla_i \nabla_j \alpha + \alpha R_{ij}]^{TF} + \alpha (K \hat{A}_{ij} - 2 \hat{A}_{ik} \hat{A}^k_j)$$

$$\partial_o \hat{\Gamma} = -2 \hat{A}^{ij} \nabla_j \alpha + 12 \alpha \hat{A}^{ij} \nabla_j \Phi \dots$$

Laguna & Shoemaker (2002)

$$\begin{aligned}
 \hat{g}_{ij} &= g^{-1/3} g_{ij} \\
 \hat{\Gamma}^i &= \hat{g}^{jk} \hat{\Gamma}^i_{jk} \\
 \hat{A}^i_j &= g^{a/2} A^i_j \\
 \hat{K} &= g^{k/2} K \\
 \hat{\beta}^i &= \beta^i \\
 \hat{\alpha} &= g^{n/2} \alpha \\
 \Phi &= \frac{1}{12} \ln g
 \end{aligned}$$

$$\partial_o \Phi = -\frac{1}{6} \hat{\alpha} \hat{K} e^{-6(n+k)\Phi}$$

Natural choice?  $n = -1, a = 1, k = 1/3$

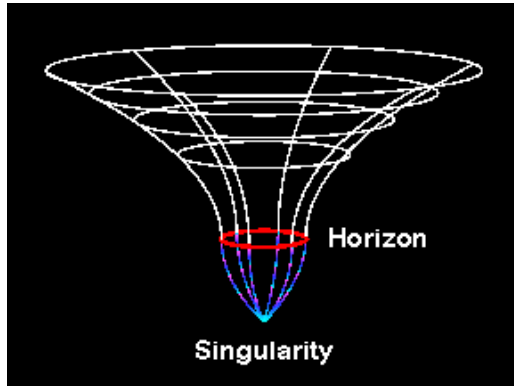
$$\partial_o \hat{g}_{ij} = -2 \hat{\alpha} \hat{A}_{ij} e^{-6(n+a)\Phi}$$

$$\begin{aligned} \partial_o \hat{K} &= -\nabla_i \nabla^i \alpha e^{6k\Phi} + \hat{\alpha} \hat{A}^i_j \hat{A}^j_i e^{-6(n+2a-k)\Phi} \\ &+ \left( \frac{1}{3} - k \right) \hat{\alpha} \hat{K}^2 e^{-6(n+k)\Phi} \end{aligned}$$

$$\begin{aligned} \partial_o \hat{A}^i_j &= -\left( \nabla^i \nabla_j \alpha - \frac{1}{3} \delta^i_j \nabla^k \nabla_k \alpha \right) e^{6a\Phi} \\ &+ \alpha \left[ R^i_j - \frac{1}{3} \delta^i_j \left( \hat{A}^k_l \hat{A}^l_k - \frac{2}{3} \hat{K}^2 \right) \right] e^{6a\Phi} \\ &+ (1 - a) \hat{\alpha} \hat{K} \hat{A}^i_j e^{-6(n+k)\Phi} \end{aligned}$$

$$\begin{aligned} \partial_o \hat{\Gamma}^i &= \left( 8k \hat{\alpha} \hat{K} \hat{g}^{ij} \partial_j \Phi - \frac{4}{3} \hat{\alpha} \hat{g}^{ij} \partial_j \hat{K} \right) e^{-6(n+a)\Phi} \\ &+ \left[ 2 \hat{\alpha} \hat{A}^{jk} \hat{\Gamma}^i_{jk} - 2 \hat{A}^{ij} \partial_j \hat{\alpha} + 12(1 + n) \hat{\alpha} \hat{A}^{ij} \partial_j \Phi \right] e^{-6(n+a)\Phi} \\ &- \left( \frac{2}{3} + \chi \right) \partial_l \beta^l (\hat{\Gamma}^i - \hat{g}^{jk} \hat{\Gamma}^i_{jk}) \end{aligned}$$





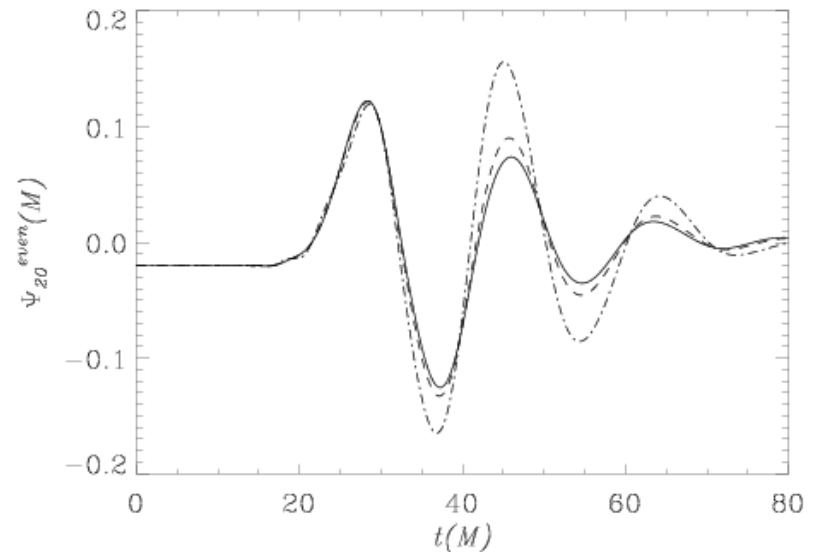
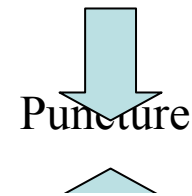
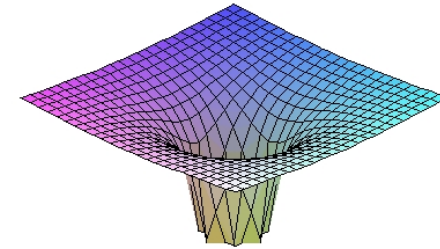
$$ds^2 = -\alpha^2 dt^2 + \phi^4 (dr^2 + r^2 d\Omega^2)$$

$$\alpha = \frac{1 - M/2r}{1 + M/2r}$$

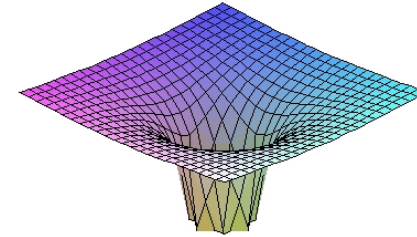
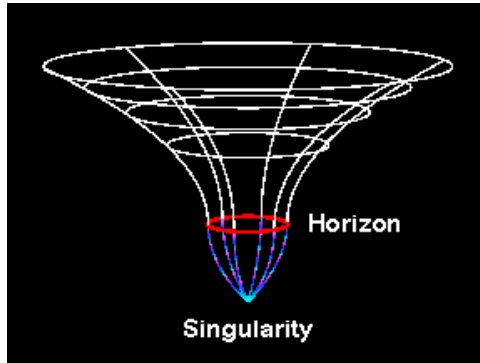
$$\phi = 1 + M/2r$$

BBH Puncture Initial Data  
*Brandt and Brugmann (1997)*

BBH Puncture Evolutions  
*Alcubierre et al (2000,2002)*





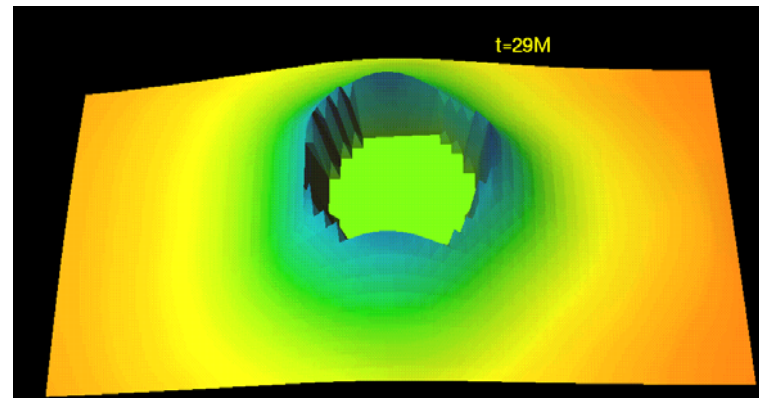
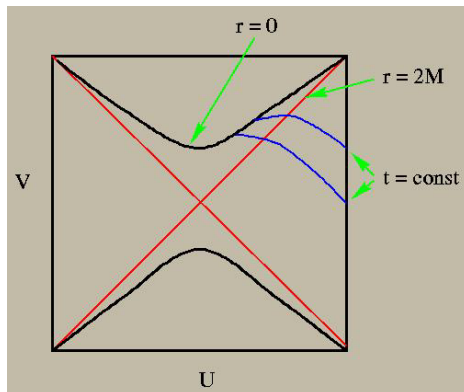


Excision

*Unruh (1984)*

Ingoing-Eddington-Finkelstein

$$ds^2 = - \left(1 - \frac{2M}{r}\right) dt^2 + \frac{4M}{r} dt dr + \left(1 + \frac{2M}{r}\right) dr^2 + r^2 d\Omega^2$$



**Student:** Are boundary conditions needed at the excision boundary?

**Death:** No, if the field variables are purely outgoing.

**S:** How can I be sure that this is indeed the case?

**D:** Work with manifestly hyperbolic formulations and apply rigorous numerical analysis.

**S:** Are the BSSN-based codes in AEI, UIUC and Penn State of this kind?

**D:** I am afraid not my son!

**S:** But their codes seem to work! Is there anything certain in this field ?

**D:** Only death. Bye.

**S:** Wait a moment!

**D:** You all say that. But I grant no reprieves.



## Outgoing Boundary Conditions?

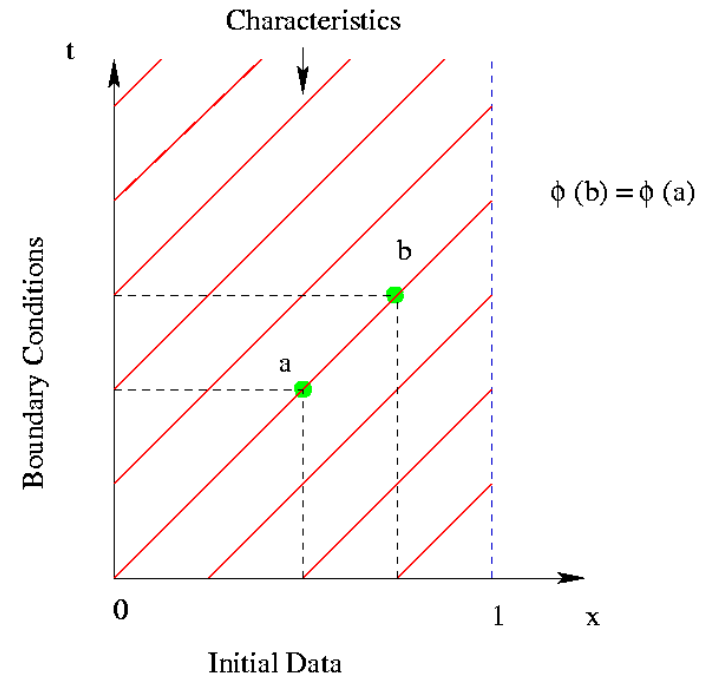
Consider,

$$\partial_t \phi + \partial_x \phi = 0$$

Approximate,

$$\partial_x \phi_i \approx \frac{\phi_{i+1} - \phi_{i-1}}{2 \Delta x}$$

What to do at  $x = 1$ ?



One-sided differencing

$$\partial_x \phi_i \approx \frac{3\phi_i - 4\phi_{i-1} + \phi_{i-2}}{2 \Delta x}$$

Centered differencing & Extrapolation

$$\partial_x \phi_i \approx \frac{\phi_{i+1}^* - \phi_{i-1}}{2 \Delta x}$$

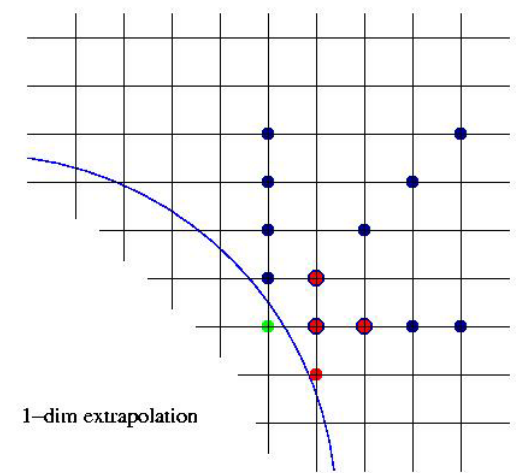
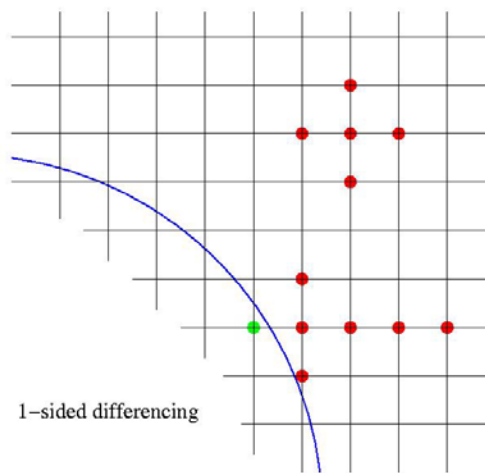
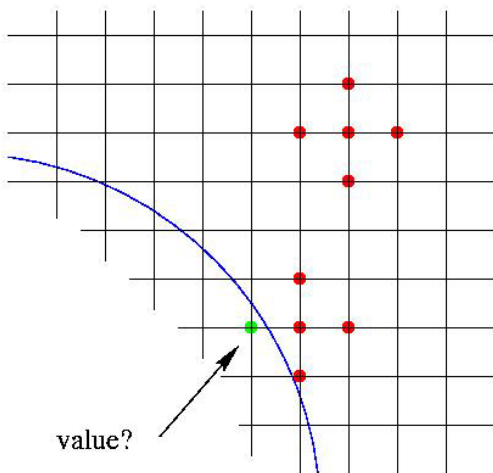
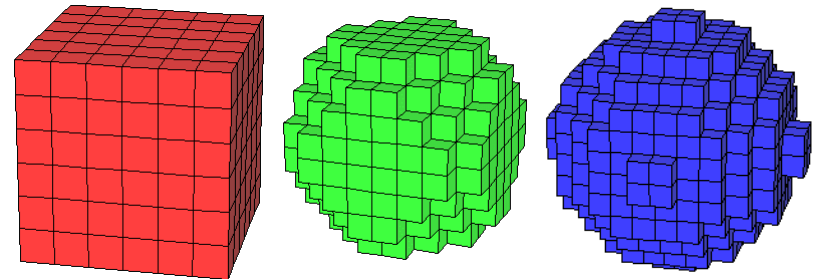
$$\phi_i^* \approx 3\phi_i - 3\phi_{i-1} + \phi_{i-2}$$

## OPTIONS:

- Modify finite-difference stencils at the excision boundary.
- Preserve centered-difference stencils and use extrapolation.

## ISSUES:

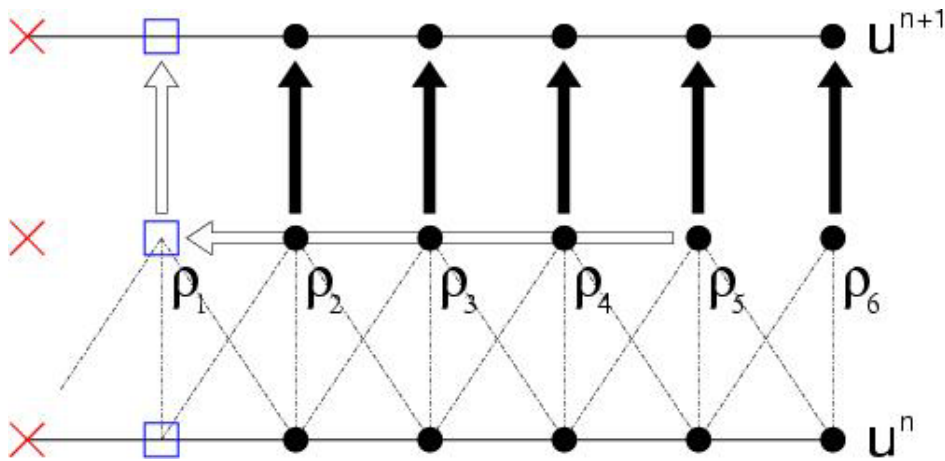
- Accuracy of extrapolation
- Choice of grid-points
- Stability



$$\partial_t u = \rho$$

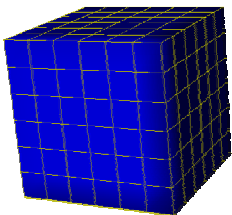


$$u_i^{n+1} = u_i^n + \Delta t \rho_i^{n+1/2}$$



Alcubierre & Brugmann

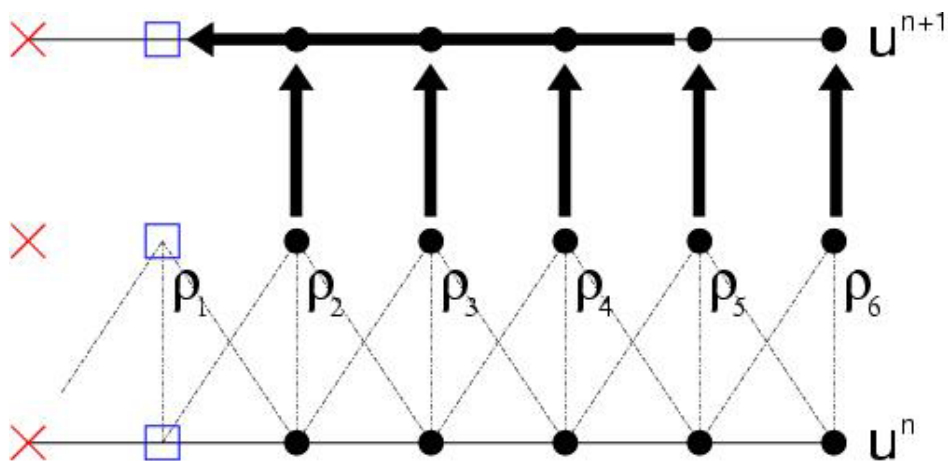
*The Mean-Field (Mean-Field)*  
2011, 2012, 2013, 2014



$$\rho_1^{n+1/2} = \rho_2^{n+1/2}$$

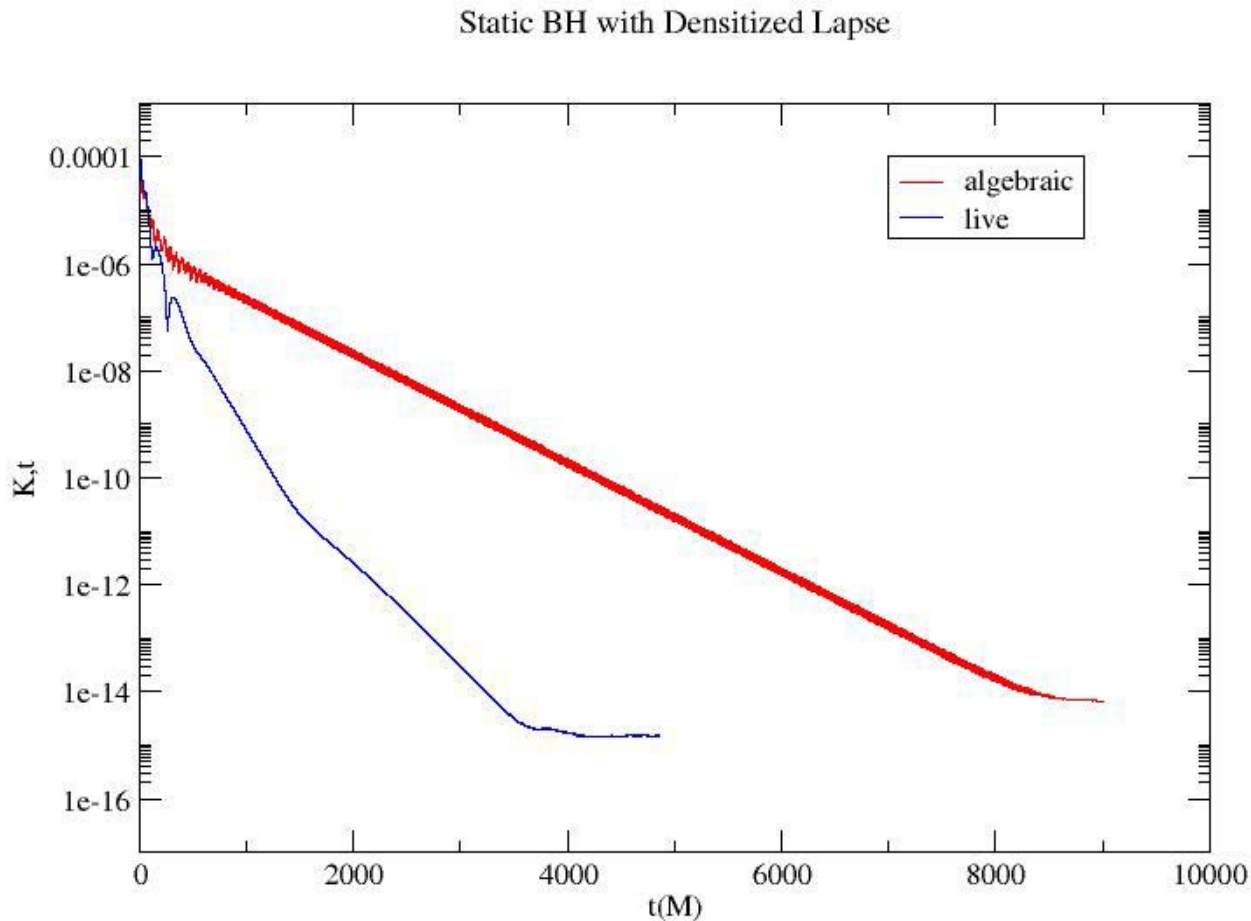


$$0 = \partial_x \rho = \partial_t \partial_x u$$



Solution extrapolation

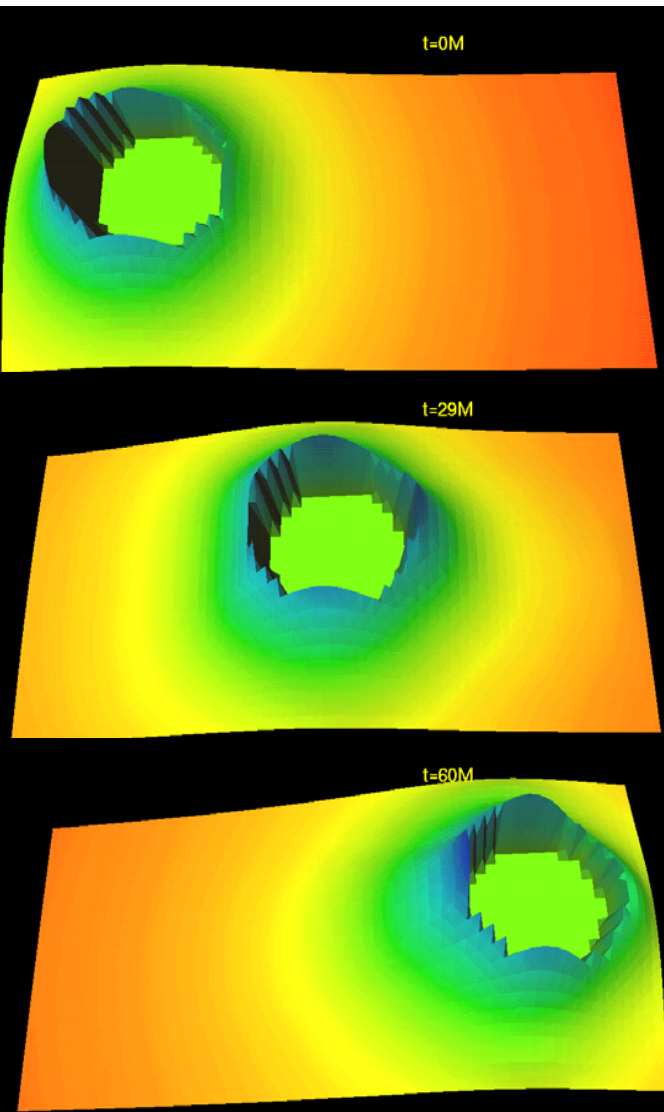
For novel and rigorous excision: O. Reula talk tomorrow.



$$\hat{\alpha} = g^{-1/2} \alpha$$



Boosted BH  
*BBH Alliance*



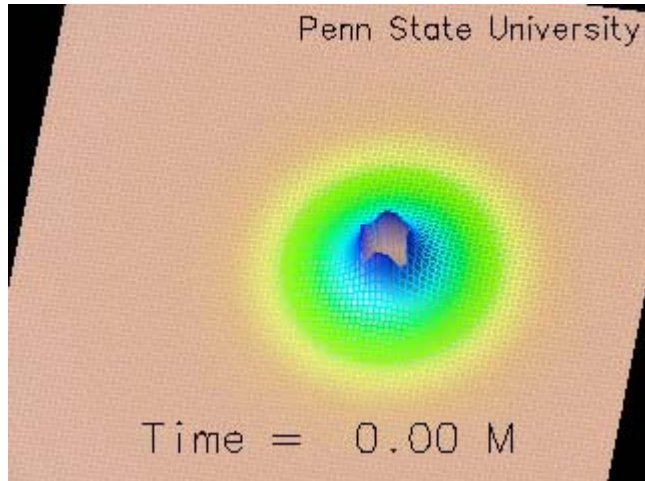
Wobbling BH:

$$\begin{aligned}t &= \bar{t} \\ x^i &= \bar{x}^i + \xi^i(\bar{t}) \\ \beta^i &= \bar{\beta}^i + \dot{\xi}^i\end{aligned}$$

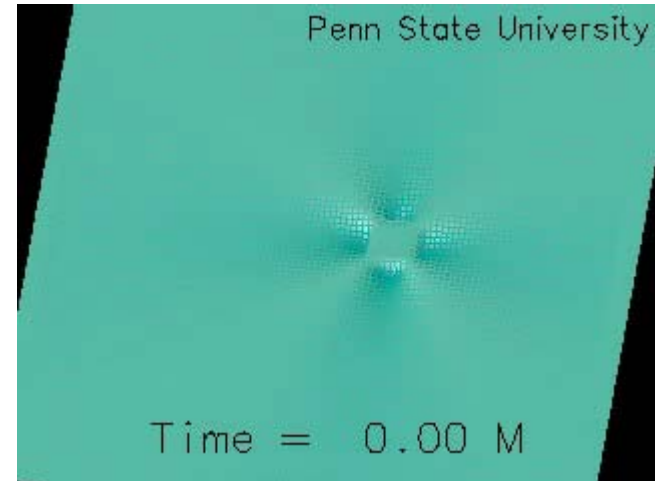
Caution:

- Simple-excision v.1 will not work
- Points appear and disappear
- Better outer boundary is needed

## Circling BH

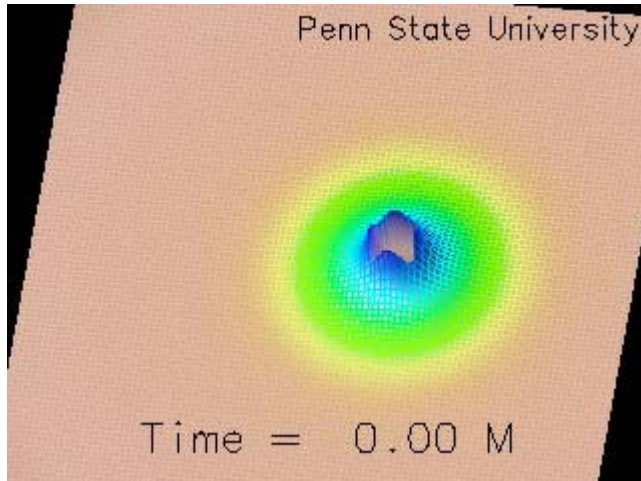


Trace Extrinsic Curvature

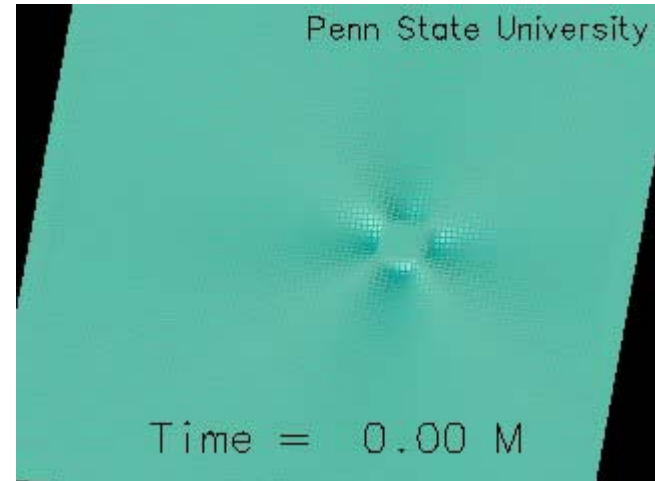


Normalized Hamiltonian Constraint

## Spiraling BH

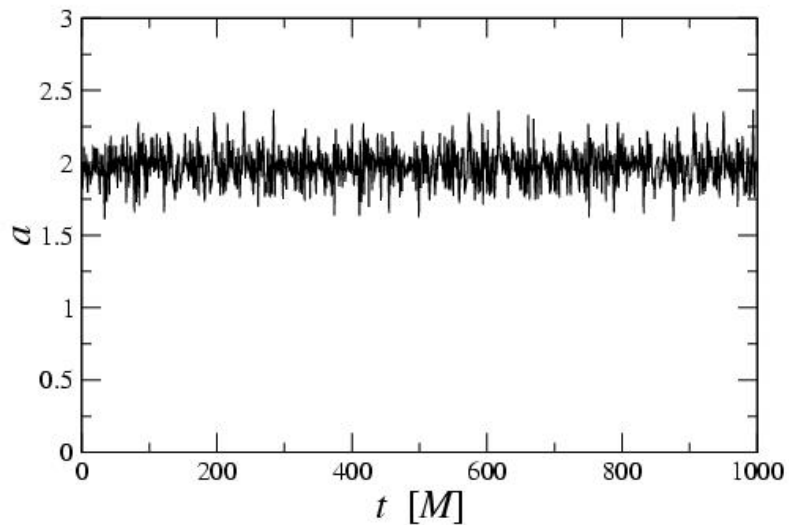


Trace Extrinsic Curvature



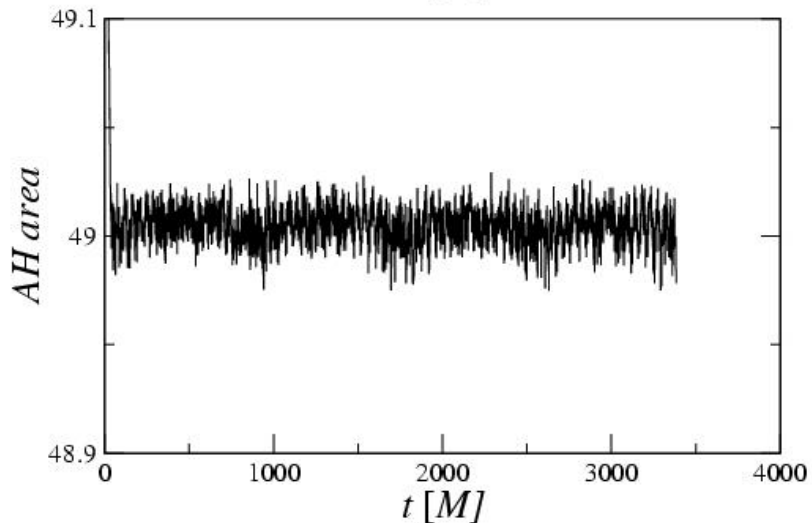
Normalized Hamiltonian Constraint

## Convergence & AH Areas



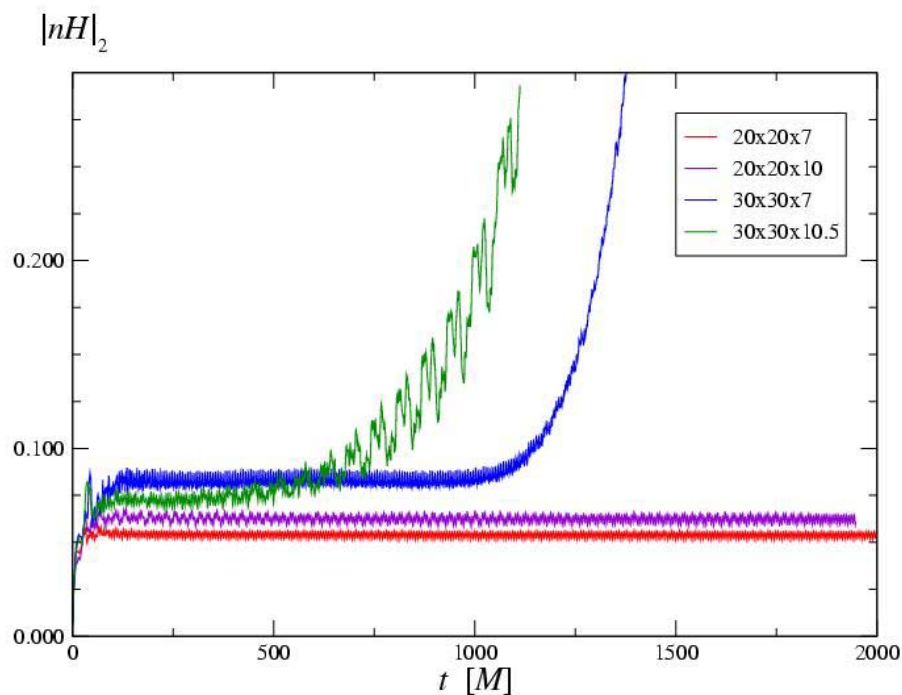
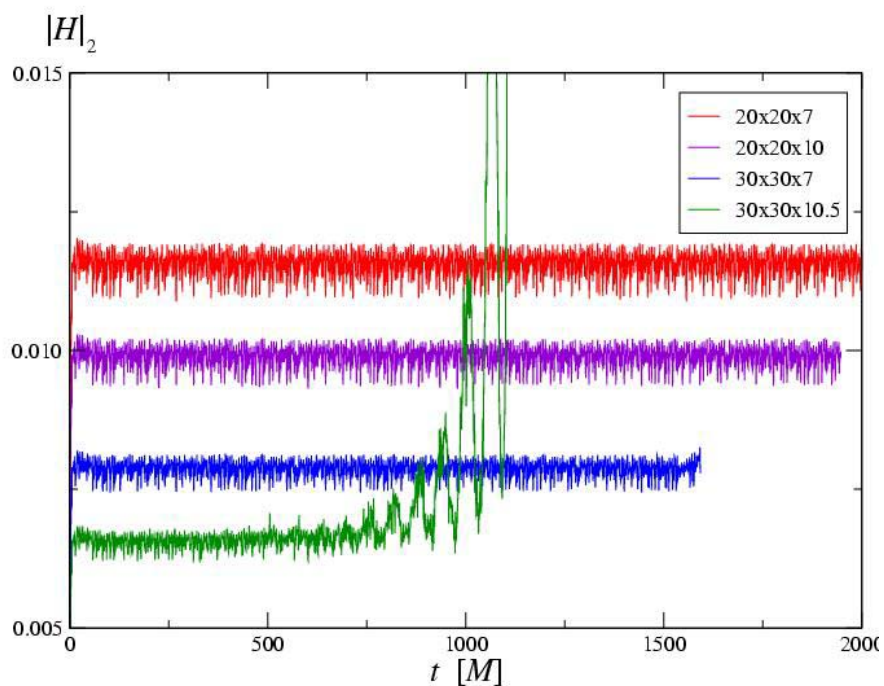
$$K_h = K_{ana} + C h^a$$

$$a = \frac{1}{\ln 2} \ln \left| \frac{K_h - K_{ana}}{K_{h/2} - K_{ana}} \right|$$



$$A_{AH} = 4 \pi (2 M)^2 = 50.265$$

## Boundary Effects





- Superposed boosted IEF holes

$$g_{ij} = \eta_{ij} + 2 H_1 l_i^1 l_j^1 + 2 H_2 l_i^2 l_j^2$$

$$\beta^i = \beta_1^i + \beta_2^i$$

$$\alpha = \alpha_1 \alpha_2$$

- Transform to a rotating frame ( $\Omega$ )

- Extrinsic curvature

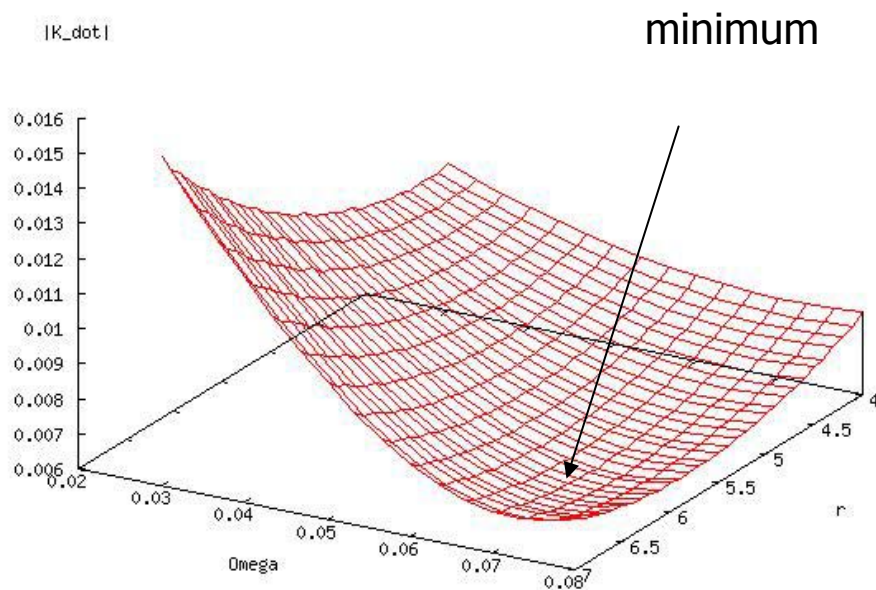
$$K_{ij} = \frac{-1}{2\alpha} (\nabla_i \beta_j + \nabla_j \beta_i)$$

- Three-parameter family ( $v$ ,  $d$ ,  $\Omega$ )

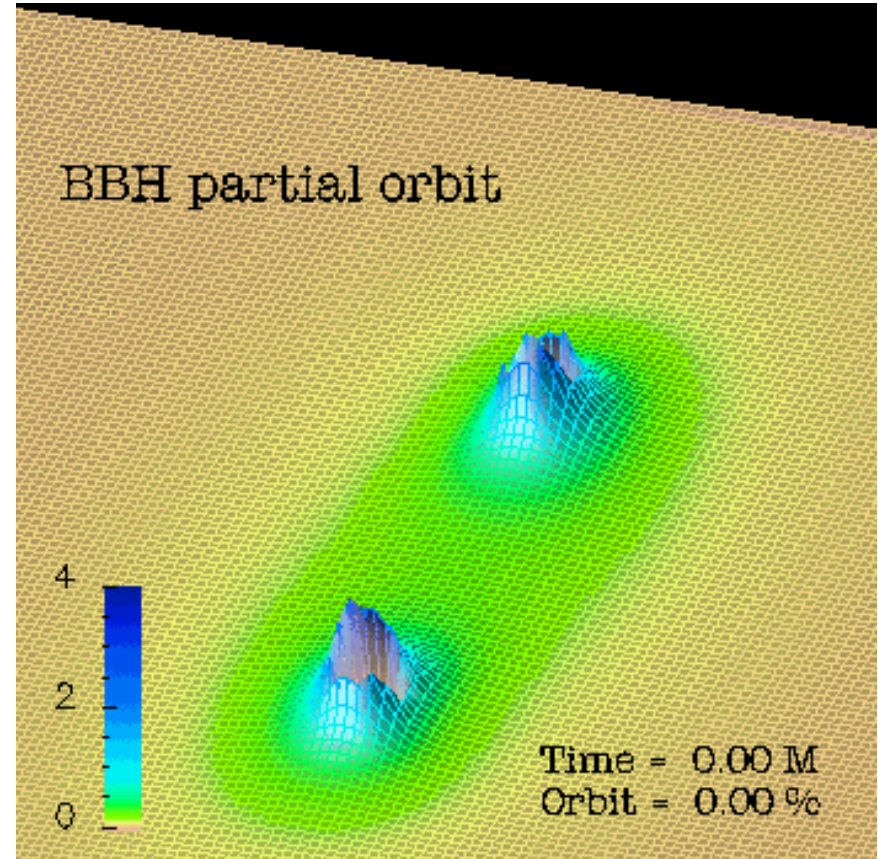
- Find minimum of

$$S = \int a |\partial_t K_{ij}|^2 + b |\rho|^2 + b |J^i|^2 dV$$

- Notice: Data does not satisfy constraints



- Evolution equations:  
Modified BSSN  
[Laguna-Shoemaker CQG 19, 3679 (2002)]
- Initial data:  
Modified Superposed IEF  
[Matzner-Huq-Shoemaker PRD 59, 024015]
- Gauge conditions:  
Live 1+log lapse and Gamma-driver shift  
[Alcubierre, et al gr-qc/0206072]
- Outer boundary:  
 $f = f_0 + u(r-vt)/r$



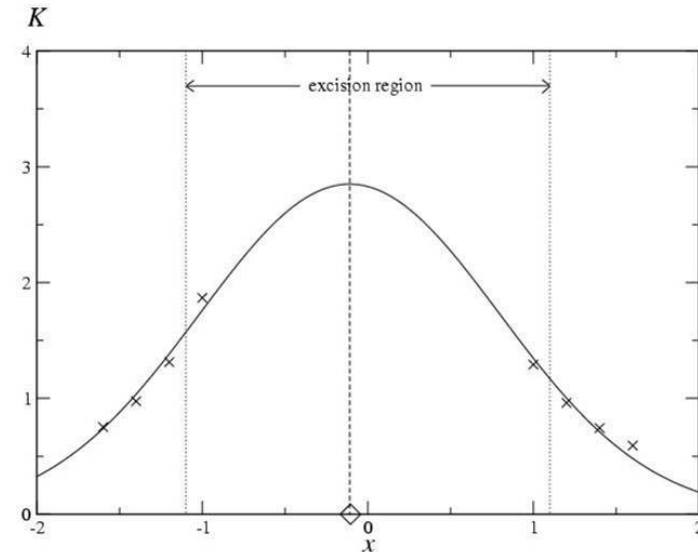
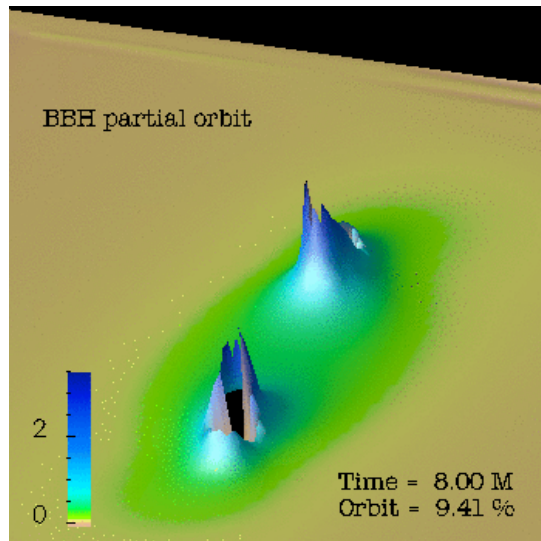
Trace of Extrinsic Curvature  
Separation = 12 M  
Boost Velocity = 0.6

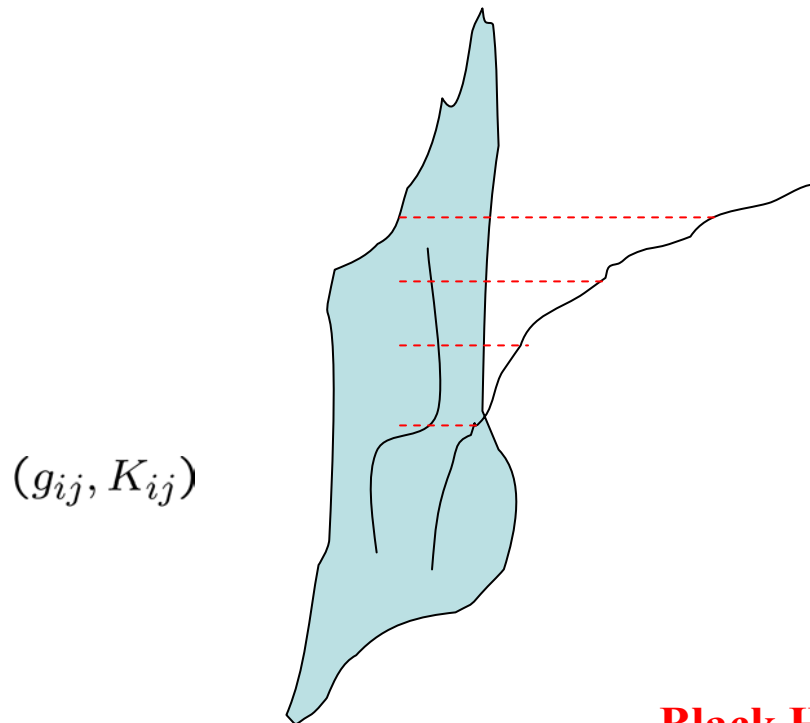


In principle an apparent horizon tracker is needed, but they are too expensive!

Alternative:

Identify the  $1/r$  fall-off near the singularity and follow its changes





Are there solutions to the Einstein evolution equations that violate the constraints and do not exhibit exponential growths?

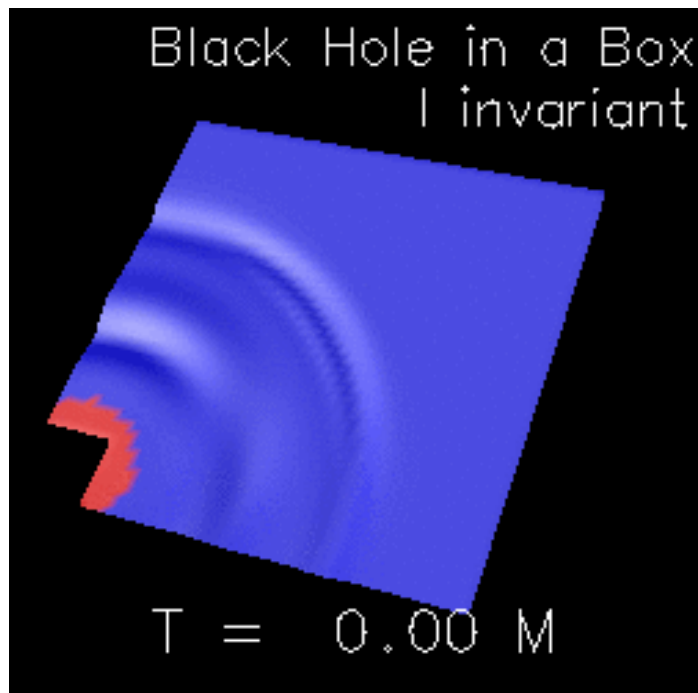
## Black Hole in a Box

Metric Distortion:

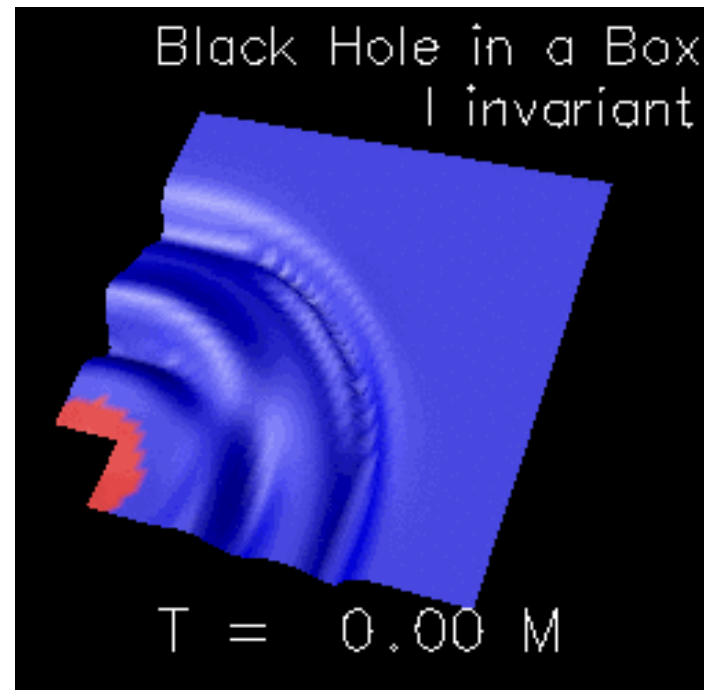
$$\begin{aligned} K &= K_{iEF} [1 + a R(r) Y_{20}] \\ \phi &= \phi_{iEF} [1 + b R(r) Y_{20}] \\ R &= \exp \left[ -\frac{(r - r_o)^2}{\lambda^2} \right] \end{aligned}$$

Mass Blending:

$$\begin{aligned} M &= [1 - c(r)] M_1 + c(r) M_2 \\ c &= \frac{1}{2} + \frac{1}{2} \tanh \left[ \frac{(r - r_o)}{\sigma} \right] \end{aligned}$$

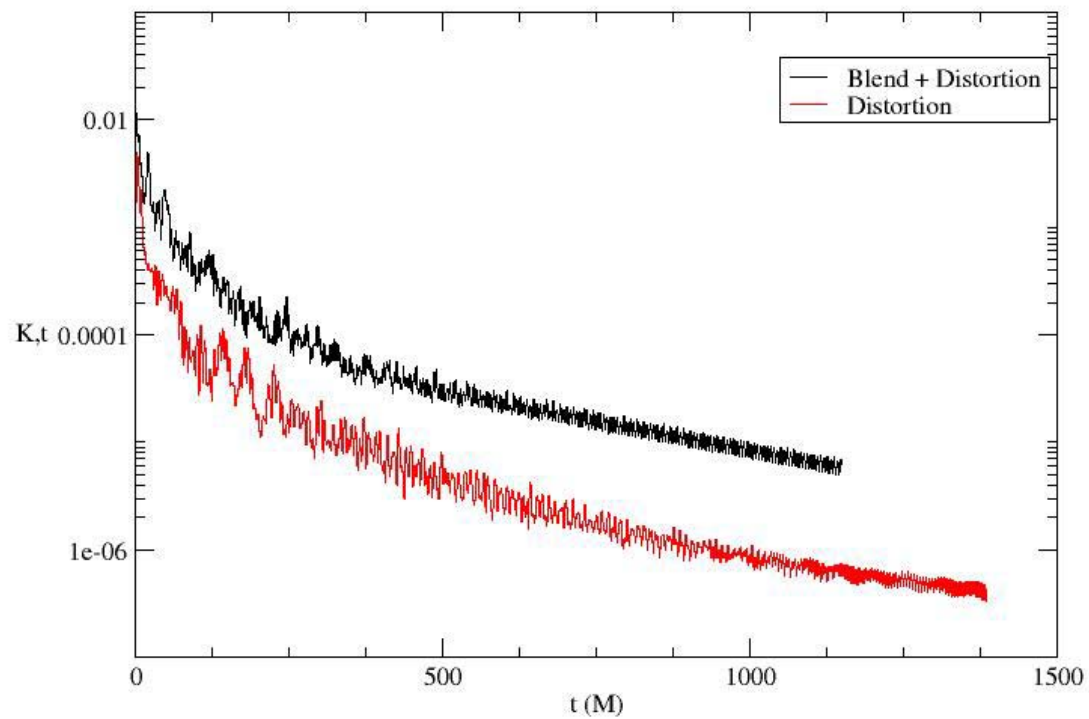


Distortion



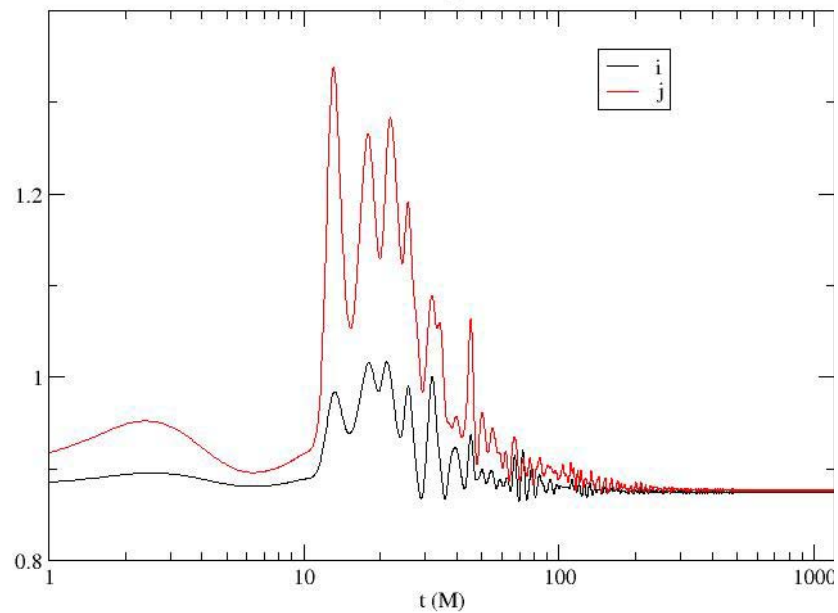
Distortion + Blend

R.H.S. of K,t



I and J invariants

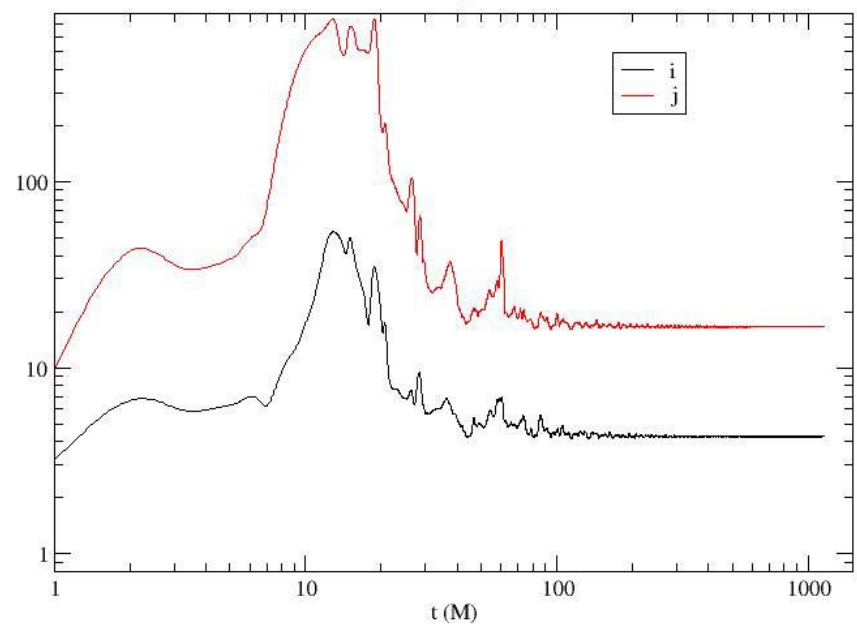
$$i = (I r^3)/3 \quad j = J r^{(9/2)}$$



Distortion

I and J invariants

$$i = (I r^3)/3 \quad j = J r^{(9/2)}$$



Distortion + Blend

$$\text{Recall speciality index: } S = 27 J^2 / I^3$$

- Black hole excision seems to be in a good track.
- Stable constraint-violating evolutions exist. How generic?
- More work on gauges, outer boundary, tracking of singularities, AMR, constrained evolutions, etc.
- More creative engineering and formal mathematics are needed.
- We should not be afraid of using approximations
- Not too far into the future, we will have real simulations of ...

