SPRING COLLEGE ON SCIENCE AT THE NANOSCALE
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BIOMOLECULES; SIMULATION

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## Pushing back the frontiers of computer simulations

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## Molecular dynamics



Given a potential energy surface:

$$
U\left(R_{1}, R_{2}, \ldots, R_{N}\right)
$$

The dynamics can be determined from Newton's equation:

$$
M_{I} \ddot{\boldsymbol{R}}_{I}=-\nabla U\left(\boldsymbol{R}_{1}, \boldsymbol{R}_{2}, \ldots, \boldsymbol{R}_{N}\right)
$$

## Empirical potentials

- Molecular mechanics: Intramolecular forces: bond stretch, bending, torsion
- Electrostatic interactions: Partial charges, dipoles, polarization
- Van der Waals interactions: Lennard-Jones, Buckingham potential

- Embedded-atom methods: Finnis-Sinclair, Glue model, Daw-Baskes



## Pros and cons

PROS- Efficient- Accurate in specific cases
CONS

- Not transferable- No chemistry


## Dealing with the electrons



$$
H \psi\left(r_{1}, \ldots, r_{M} ; R_{1}, \ldots, R_{N}\right)=E \psi\left(r_{1}, \ldots, r_{M} ; R_{1}, \ldots, R_{N}\right)
$$

## Hartree-Fock

$$
\left.\begin{array}{c}
\Psi\left(x_{1}, \ldots, x_{N}\right)=\frac{1}{\sqrt{N!}} \left\lvert\, \begin{array}{ccc}
\varphi_{1}\left(x_{1}\right) & \cdots & \varphi_{1}\left(x_{N}\right) \\
\vdots & \ddots & \vdots \\
\varphi_{x}\left(x_{1}\right) & \cdots & \varphi_{s}\left(x_{N}\right)
\end{array}\right. \\
{\left[-\frac{1}{2} \nabla+V_{n}(x)+\sum_{j=1} \int \frac{1}{x-x^{\prime} \mid} \varphi_{i}\left(x^{\prime}\right)^{2} d x^{\prime}-J\right]}
\end{array}\right] \varphi_{i}(x)=\varepsilon_{i} \varphi_{i}(x) \quad .
$$

Exchange operator:

$$
J \varphi_{i}(x)=\frac{1}{2} \sum_{j} \int \varphi_{j}(x) \frac{1}{x-x^{\prime}} \varphi_{j}^{*}\left(x^{\prime}\right) \varphi_{i}\left(x^{\prime}\right) d x^{\prime}
$$

## Beyond Hartree-Fock

## Adding correlations

- Perturbation theory MP2, MP4, ...
$\mathrm{N}^{4}, \mathrm{~N}^{5}$
- Configuration interaction
$\exp (\mathrm{N})$
- Coupled clusters
$\mathrm{N}^{6}, \mathrm{~N}^{7}$

Small is beautiful!

## Hohenberg-Kohn

The energy of the ground state of a many-body system is a unique functional of the electron density:

$$
\boldsymbol{E}=\boldsymbol{E}\left[\rho_{e}(r)\right]
$$

The functional is minimum for the ground state density:

$$
\begin{gathered}
\boldsymbol{E}=\boldsymbol{E}\left[\rho_{e}(\boldsymbol{r})\right]-\mu \boldsymbol{N} \\
\frac{\delta \boldsymbol{E}\left[\rho_{e}(\boldsymbol{r})\right]}{\delta \rho_{e}(\boldsymbol{r})}=\mu
\end{gathered}
$$

## Kohn-Sham

$$
\begin{gathered}
\rho_{e}(\boldsymbol{r})=2 \sum_{n} \psi_{n}^{*}(\boldsymbol{r}) \psi_{n}(\boldsymbol{r}) \\
\boldsymbol{E}=-\frac{1}{2} \sum_{n} \int \boldsymbol{d r} \psi_{n}^{*}(\boldsymbol{r}) \nabla \psi_{n}(\boldsymbol{r})+\int d r \rho_{e}(\boldsymbol{r}) V_{e x t}(\boldsymbol{r}) \\
+\frac{1}{2} \int d r d r^{\prime} \rho_{e}(\boldsymbol{r}) \frac{1}{\left|r-\boldsymbol{r}^{\prime}\right|} \rho_{e}\left(\boldsymbol{r}^{\prime}\right)+\boldsymbol{E}_{x c}\left[\rho_{e}(\boldsymbol{r})\right] \\
\left(-\frac{1}{2} \nabla+V_{e x t}(\boldsymbol{r})+V_{H}(\boldsymbol{r})+V_{\mathrm{xc}}(\boldsymbol{r})\right) \psi_{n}(\boldsymbol{r})=\varepsilon_{n} \psi_{n}(\boldsymbol{r}) \\
V_{H}(\boldsymbol{r})=\int d \boldsymbol{r}^{\prime} \frac{1}{\left|r-\boldsymbol{r}^{\prime}\right|} \rho_{e}\left(\boldsymbol{r}^{\prime}\right) \quad V_{x c}(\boldsymbol{r})=\frac{\delta E_{x c}\left[\rho_{e}(\boldsymbol{r})\right]}{\delta \rho_{e}(\boldsymbol{r})}
\end{gathered}
$$

## Born-Oppenheimer

The potential energy surface is defined by the instantaneous ground state electronic energy:

$$
\phi\left(\boldsymbol{R}_{1}, \boldsymbol{R}_{2}, \ldots, \boldsymbol{R}_{N}\right)=\boldsymbol{E}_{0}\left(\boldsymbol{R}_{1}, \boldsymbol{R}_{2}, \ldots, \boldsymbol{R}_{N}\right)
$$

But

$$
E_{0}\left(R_{1}, R_{2}, \ldots, R_{N}\right)
$$

needs to be approximated. We shall choose a theory which has the right balance between accuracy and computational efficiency.

## Ab-initio MD

$$
\begin{gathered}
L=\frac{1}{2} \mu \sum_{n} \int d r \dot{\psi}_{n}(r)^{2}+\sum_{I} \frac{1}{2} M_{I} \dot{R}_{I}^{2}-\boldsymbol{E}_{K S}\left[\psi_{n}, \boldsymbol{R}_{I}\right] \\
+\sum_{n, m} \Lambda_{n, m}\left(\left\langle\psi_{n} \mid \psi_{m}\right\rangle-\delta_{n, m}\right) \\
\omega_{e} \propto \sqrt{\frac{\boldsymbol{E}_{g}}{\mu}} \ll \omega_{I}
\end{gathered}
$$

## Car-Parrinello molecular dynamics

$$
\begin{gathered}
\mu \ddot{\psi}_{i}=-\frac{\delta E^{D F T}}{\delta \psi_{i}^{*}}+\sum_{j} \Lambda_{i j} \psi_{j} \\
M_{I} \ddot{\mathbf{R}}_{I}=-\nabla_{\mathbf{R}_{I}} E^{D F T} \\
\mu_{q} \ddot{\alpha}_{q}=-\frac{\partial E^{D F T}}{\partial \alpha_{q}}
\end{gathered}
$$



# Some problems 

+ Size
+ Time scale
+ Accuracy
* Non Born-Oppenheimer


## Going to larger systems

- New Gaussian-based code QUICKSTEP
n. Field theoretical approach
- $\mathrm{QM} / \mathrm{MM}$
n. DFT-based potentials


## Plane wave basis set

Plane wave expansion: $\psi_{i}(\mathbf{x})=\Sigma_{\mathbf{G}} \mathrm{C}_{i}(\mathbf{G}) \mathrm{e}^{\mathbf{i} \mathbf{G} \mathbf{x}}$


G are the reciprocal space vectors.
The Hilbert space spanned by $\mathbf{P W s}$ is truncated to a cut-off $\mathbf{G}^{2} / 2<E^{\text {cut }}$


## Which basis set?

## Plane Waves

- Orthogonal
- No chemical input
- Convergence easy to check
- Simple algebra
- Memory intensive
- Linear algebra and FFT
- No basis set superposition error


## Gaussians

- Non orthogonal
- Chemical input
- Convergence less easy to check
- More complex
- Reduced memory
- Quantum Chemistry know-how
- Basis set superposition error


## Basis set expansion

Orbitals $\Phi_{i}$ are expanded in a set of $M$ basis functions $\left\{\chi_{\alpha}\right\}$

$$
\Phi_{i}(r)=\sum_{i=1}^{M} c_{\alpha i} \chi \alpha(r)
$$

Basis functions

- Atomic orbital based (Gaussian, Slater, numerical)
- Plane waves
- Grid based, finite elements, wavelets


## Basis set expansion

$$
\begin{array}{rlr}
S_{\alpha \beta} & =\int \chi_{\alpha}^{*}(r) \chi_{\beta}(r) d r & \text { overlap matrix } \\
H_{\alpha \beta} & =\int \chi_{\alpha}^{*}(r) \mathcal{H}(r) \chi_{\beta}(r) d r & \text { Hamiltonian matrix } \\
E_{i j} & =\delta_{i j} \epsilon_{i} & \text { Orbital energies } \\
& H C=S C E &
\end{array}
$$

## Orbitals

$$
\begin{array}{rlr}
C^{\dagger} S C & =1 & \text { Orthogonality } \\
E[C] & =E\left[C^{\prime}\right] & \text { Invariant } \\
C^{\prime} & =C U & \\
U^{\dagger} U & =1 &
\end{array}
$$

## Density matrix

$$
P_{\alpha \beta}=2 \sum_{i=1}^{N_{e} / 2} c_{\alpha i} c_{\beta i}^{*}
$$

Properties of the density matrix

$$
\begin{aligned}
\operatorname{Tr}(\boldsymbol{P S}) & =\sum_{\alpha \beta}^{M} P_{\alpha \beta} S_{\alpha \beta}=N_{e} & & \text { normalisation } \\
\boldsymbol{P} & =\frac{1}{2} \boldsymbol{P S} \boldsymbol{P} & & \text { idempotency }
\end{aligned}
$$

## Density matrix

－Unique
－Electron density

$$
\rho(r)=\sum_{\alpha \beta} P_{\alpha \beta} \chi_{\alpha} \chi_{\beta}^{*}
$$

－Expectation values

$$
\begin{aligned}
\langle\mathcal{O}\rangle & =\operatorname{Tr}(P \mathcal{O}) \\
\sum_{i} \epsilon_{i} & =\operatorname{Tr}(P H)
\end{aligned}
$$

## Gaussian basis

Basis functions

$$
\chi(r)=x^{l} y^{m} z^{n} \exp \left[-\alpha r^{2}\right]
$$

Product of basis functions

$$
\chi(r-A) \chi(r-B)=\tilde{\chi}(r-C)
$$

Localization

$$
\exp \left[-\alpha r^{2}\right] \rightarrow \text { FFT } \rightarrow \exp \left[-\frac{G^{2}}{4 \alpha}\right]
$$

## DFT with gaussians



## The best of both worlds

$$
\begin{array}{cc}
\text { Gaussian } & \begin{array}{c}
\text { PW } \\
\text { (real) }
\end{array} \\
\rho & \begin{array}{c}
\text { PW } \\
\text { (Reciprocal) }
\end{array} \\
\sum_{\mu v} P^{\mu v} \phi_{\mu}(\mathbf{r}) \phi_{v}(\mathbf{r}) \xrightarrow[\text { space }]{\text { real }} \rightarrow n(\mathbf{r}) \xrightarrow[\text { FFT }]{\longrightarrow} n(\mathbf{G})
\end{array}
$$

## Linear scaling



## Standard approach

Solve by diagonalization

$$
\begin{array}{rr}
H_{\mu \nu} C^{v i}=S_{\mu \delta} C^{\delta i} \varepsilon_{i} & i=1 \ldots N \\
\text { Construct } \mathrm{P} & \mu, v, \delta=1 \ldots M \\
P^{\mu \nu}=\sum_{i} C^{\mu i} C^{v i} &
\end{array}
$$

## Orbital rotation

$$
\begin{aligned}
& C(X)=C_{0} \cos \left(\sqrt{X^{T} S X}\right)+X \frac{\sin \left(\sqrt{X^{T} S X}\right)}{\sqrt{X^{T} S X}} \\
& X^{T} S C_{0}=0 \quad C(X)^{T} S C(X)=1 \forall X
\end{aligned}
$$

Minimize the energy with respect to X

## Overall cubic scaling



## Accurate forces



## Checking the eigenstates



## DNA crystal



2388 atoms, 3960 orbitals
DZV(d,p) 22596, TZV(2d,2p) 38688
$675,1100 \mathrm{sec} /$ line search (SP4-32-1.3G)
2.5, $5 \mathrm{~h} /$ total

Not yet fully cubic (45,43,8 \% 3,2,1)
Not yet sparse

## Example: DNA Crystal

2388 atoms, 3960 orbitals, 38688 BSF (TZV(2d,2p)) density matrix, overlap matrix


## One order of magnitude better



## Why should one combine QM and MM ?

Bio-systems are typically very large and catalyse complicated reactions

- Proteins >1000 atoms
- Solvent >10000 atoms
- The active site $\sim 100$ atoms



## Mixed Quantum-Classical QM/MM- Car-Parrinello Simulations



, Fully Hamiltonian

- ע

Q MJPemit CPMI 3.3
( $\rho$ '.ces (2 'boxes), plense weves,
pseudo potentisls, GGASE

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MNJ-PEMt GrajMOSeg j PEMJ, A MEER
1 soode

## Interface Region

Quantum Region
(Car-Parrinello)

Classical Region

# QM/MM- Car-Parrinello Simulations 

- Development of improved QM/MM interfaces:
- pseudo potentials
for boundary atoms
- efficient treatment of long-range electrostatics
- electron spill out problem


## The Bonded Part


-boundary atoms: monovalent pseudopotential

- distances - angles - torsions involving MM and QM atoms come from the force field


## Mixed Quantum-Classical Simulations



- charges, located on the QM atoms, are fitted to the electrostatic fields on the MM atoms due to the electronic charge distribution


Explicit dependence of $q^{\text {ESP }}$ on the positions of all the MM and QM atoms and on the electronic density

## QM/MM

$$
\begin{aligned}
& \text { Quantum mechanics } \\
& L=\frac{1}{2} \sum_{i} \mu \int d \vec{r}\left|\dot{\psi}_{i}\right|^{2}+\frac{1}{2} \sum_{I} M_{I} \dot{\vec{R}}_{I}^{2}-E_{K S}\left[\left\{\psi_{i}\right\}, R_{I}\right]+\sum_{i, j} \Lambda_{i, j}\left(\left\langle\psi_{i} \mid \psi_{j}\right\rangle-\delta_{i, j}\right) \\
& +\frac{1}{2} \sum_{\boldsymbol{I}} m \dot{\vec{R}}_{\boldsymbol{I}^{\prime}}^{2}-\sum_{I^{\prime}<J^{\prime}} V_{L L}\left(\mid \vec{R}_{I^{\prime}}-\vec{R}_{J^{\prime}}\right)-\sum_{I, I^{\prime}} V_{C L}\left(\left|\vec{R}_{I^{\prime}}-\vec{R}_{I^{\prime}}\right|\right)
\end{aligned}
$$

A. Laio, J. VandeVondele and U. Röthlisberger, Chem. Phys. 116, 6941(2002)

## Capturing the complexity



DNA oxidation

## The way of the future?

## $\mathrm{O}(\mathrm{N})$ vs $\mathbf{O}(\mathrm{N} 3)$



Highly parallel and local new field theoretical approach

## The quickstep team

J. Hutter, University of Zurich
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## Why should one combine QM and MM ?

Bio-systems are typically very large and catalyse complicated reactions


- Proteins $>1000$ atoms
- Solvent $>10000$ atoms
- The active site $\sim 100$ atoms



## Mixed Quantum-Classical QM/MM- Car-Parrinello Simulations




- Fully Hamiltonian
- HID drivera (jp Mill

QMAParts CpMID 3.3
( $\rho$ !ocs (2 boxes), plane weves, pэeudo potentials, GGぶs:

n-J nodes
 AMEER
I suode

Interface Region

Quantum Region (Car-Parrinello)

Classical Region

# QM/MM- Car-Parrinello Simulations 

- pseudo potentials
for boundary atoms
- efficient treatment of long-range electrostatics
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## The Bonded Part


-boundary atoms: monovalent pseudopotential

- distances - angles - torsions involving MM and QM atoms come from the force field


## Mixed Quantum-Classical Simulations

$$
\rho(n r 1 i, n r 2 j, n r 3 k)
$$



$$
H_{a l}=\sum_{i \in N M} q_{i} \int d r \frac{\rho(r)}{\left|r-r_{i}\right|}
$$

3D-grid: (NR1,NR2,NR3) NR~100
MM atoms ~10000-100000


NR1*NR2*NR3*MM
DISTANCE CALCULATIONS!!!

- charges, located on the QM atoms, are fitted to the electrostatic fields on the MM atoms due to the electronic charge distribution


Explicit dependence of $q^{\text {ESP }}$ on the positions of all the MM and QM atoms and on the electronic density

## QM/MM

$$
\begin{aligned}
& \text { Quantum mechanics } \\
& L=\left.\frac{1}{2} \sum_{i} \mu f d \vec{r} \psi_{i}\right|^{2}+\frac{1}{2} \sum_{I} M_{I} \dot{\vec{R}}_{I}^{2}-E_{K S}\left[\left\{\psi_{i}\right\rangle, R_{I}\right]+\sum_{i, j} \Lambda_{i, j}\left(\left\langle\psi_{i} \mid \psi_{j}\right\rangle-\delta_{i, j}\right) \\
& +\frac{1}{2} \sum_{\Gamma} m \dot{\vec{R}}_{I^{\prime}}^{2}-\sum_{I^{\prime}<J^{\prime}} V^{\prime} L L\left(\mid \vec{R}_{I^{\prime}}-\vec{R}_{J^{\prime}}\right)-\sum_{I, I^{\prime}} V_{C L}\left(\mid \vec{R}_{I^{\prime}}-\vec{R}_{I^{\prime}}\right)
\end{aligned}
$$

A. Laio, J. VandeVondele and U. Röthlisberger, Chem. Phys. 116, 6941(2002)

## Capturing the complexity



DNA oxidation

## The way of the future?

## $\mathrm{O}(\mathrm{N})$ vs $\mathbf{O}\left(\mathbf{N}^{3}\right)$



Highly parallel and local new field theoretical approach

## The time scale problem

Direct simulation allows only very short runs:
$\sim 10 \mathrm{ps}$ for ab-initio MD, $\sim 10 \mathrm{~ns}$ for classical MD
Many relevant phenomena take place on a larger time scale: chemical reactions, conformational changes, protein folding, etc.

Two-fold strategy:
a) Finding reactive paths
b) Exploring the free energy surface

## Activated events



## The quantum chemical approach

$\times$ Find the saddle point on the PES
$\times$ Use transition state theory
$\times$ Correct for zero point motion

## Many different solutions proposed

Thermodynamic integration
＂Flattening＂the surface
（hyperdynamics，puddle－skimming，umbrella sampling，etc．）
圆 Trajectory－based schemes
（reaction path sampling，Lagrangian action minimization，nudged elastic band，etc．）
圈 Finding the saddle points
（eigenvalue following，dimer method，hessian－based methods，etc．）
圈 Temperature enhanced sampling （histogram reweighting，parallel tempering，etc．）

圆 Etc．etc．

## Driving the reaction

Suppose that the reaction coordinate $\boldsymbol{q}$ is known:

$$
q\left(R_{1}, R_{2}, \ldots, R_{N}\right)=q
$$

We force the reaction by adding a constraint term to the dynamics with a Lagrange multiplier:

$$
\begin{aligned}
& \lambda\left(q\left(R_{1}, R_{2}, \ldots, R_{N}\right)-q\right) \\
& \langle\lambda\rangle \approx \frac{\partial F}{\partial q}
\end{aligned}
$$

The activation energy is:

$$
\Delta F=\int_{q_{A}}^{q^{*}} d q\langle\lambda\rangle_{q}
$$

The right reaction path?


Activated events are often intrinsically multidimensional!!!

## Life is complicated



## Life is complicated

The potential energy surface of a complex system is rough

(C) D. Chandler

## How to explore a multidimensional free energy surface?

Need to be able to escape free energy minima

Our solution:
Non-Markovian coarse-grained dynamics
A. Laio and M. Parrinello, PNAS

## Collective variables

Choose a small set of slow collective variables:

$$
{ }_{i}^{\delta_{i}\left(\vec{R}_{I}\right) \quad i=1, n}
$$

The $\mathbf{s}_{\mathbf{i}}$ :

- Discriminate between reactants and products
- Include all the relevant slow modes
- The reaction coordinate is a linear combination of the $\mathbf{s}_{\mathbf{i}}$


## Examples of collective variables

- Distances
- Angles: bending and torsional
- Coordination numbers: between individual atoms
- or between different species
- Local electric fields
- Number of n-fold rings
- Solvation energy
- Lattice vectors
- Energy
- Etc. etc.


## Probability distribution

$$
P\left(S_{1}, S_{2} \ldots, S_{n}\right)=\frac{\int d \vec{R}_{I} \prod_{i}^{n} \delta\left(S_{i}-s_{i}\left(\vec{R}_{I}\right)\right) e^{-\beta V\left(\vec{R}_{I}\right)}}{\int d \vec{R}_{I} e^{-\beta V\left(\vec{R}_{I}\right)}}
$$

We want to study the free energy as a function of these variables:

$$
\beta F\left(S_{1}, S_{2}, \ldots, S_{n}\right)=-\ln P\left(S_{1}, S_{2}, \ldots, S_{n}\right)
$$

## The algorithm:

-Wherever you go put a "small" Gaussian

- Always move in the direction of the direction that minimizes the sum of $\mathrm{V}(\mathrm{s})$ and all the Gaussians




## NaCl in water

Two minima:
Contact ion pair (metastable)
Dissociated

Collective coordinates: Electric field on $\mathrm{Na}^{+}$
Electric field on $\mathrm{Cl}^{-}$
Distance $\mathrm{Na}^{+} \mathrm{Cl}^{-}$

Classical MD:
Amber force field

## Free energy surface



Transition state


## Dialanine in water

1 dialanine in 287 TIP3P water

AMBER95 force field

Collective coordinates: backbone dihedral angles $\Phi$ and $\Psi$


## Stationary action principle

$$
\begin{gathered}
L=\frac{1}{2} \dot{q}^{2}-V(q) \\
\ddot{a} S=\ddot{a} \int_{t_{A}}^{d_{B}} L(q(t), \dot{q}(t)) d t=0 \\
q\left(t_{A}\right)=q_{A} \\
q\left(t_{B}\right)=q_{B}
\end{gathered}
$$

Can we use this principle to find the trajectories?
Saddle point!

## Folding a small peptide



## Motivation and aims

Oxidatitive damage to DNA is common and has fatal consequences

Guanine, having the lowest oxidation potential among the DNA bases, plays a fundamental role

Does the structure of DNA funnel the reactions towards a unique product?

## A continuous metadynamics

Fictitious kinetic energy

$$
\mathcal{L}=\mathcal{L}_{C P}+\sum_{\alpha} \frac{1}{2} M_{\alpha} \dot{s}_{\alpha}^{2}-\sum_{\alpha} \frac{1}{2} k_{\alpha}\left(S_{\alpha}(\mathbf{R})-s_{\alpha}\right)^{2}-V(\mathbf{s}, t)
$$

$$
\begin{aligned}
& \text { Restrain potential } \\
& \text { If } \sqrt{ }(\mathrm{k} / \mathrm{M}) \ll \\
& \qquad \frac{\partial \mathcal{F}}{\partial s_{\alpha}} \simeq\left\langle k_{\alpha}\left(S_{\alpha}(\mathbf{R})-s_{\alpha}\right)\right\rangle
\end{aligned}
$$

History-dependent potential

From azulene to naphthalene




## Science fiction?



## $\mathrm{G}^{+}$localization how?



## Computational details

## Model

G:C decamer, Z conformation
X-ray structure available
Rich in G and the smallest cell Electronic structure treated by the Kohn-Sham method

BLYP and HCTH functionals, plane waves (70Ry cutoff)

Martin-Troullier pseudopotentials
Car-Parrinello molecular dynamics, CPMD code


## Computational details

The hole localizes on $\mathrm{G}^{+}$and the proton moves to C


The model includes water and counter ions
1194 atoms
Full quantum: 3,960 valence electrons, 408,238 PW!

## The HOMO of DNA



## Fate of $\mathrm{G}^{+} \cdot$ in DNA

In duplex DNA the protonation is not clear
$\mathrm{CH}^{+}$has a pKa of $4.3\left(\mathrm{G}^{+}: \mathrm{pKa} 3.9\right)$
_ Sharing of proton expected

$\mathrm{K}_{\mathrm{exp}}=2.5$ is predicted from experiments in water
Is this relevant for DNA?

## Bridging length scale

Most calculations were done in the QM/MM framework.
Laio A., VandeVondele J. and Roethlisberger U., J.Chem. Phys. 116, 6941(2002)
Due to high polarity of bonds H-capping was used. Added H were decoupled to MM.





Quantum sub-systems of increasing size were used.

The biggest quantum model was the full DNA.

## Protonation state of the bases: gas phase

Collective coordinate used: coordination number ( $\mathrm{C}_{\mathrm{AB}}$ )

$$
C_{A B}=\sum_{i=1}^{N_{B}} \frac{1-\left(\frac{r_{A i}}{\boldsymbol{R}_{A B}}\right)^{n}}{1-\left(\frac{r_{A i}}{\boldsymbol{R}_{A B}}\right)^{m}}
$$



$$
\begin{aligned}
& \mathrm{n}=6, \mathrm{~m}=12 \\
& \text { of } \mathrm{N} 1, \mathrm{~N} 2 \text { and } \mathrm{N} 1, \mathrm{~N} 4 \\
& \text { with respect to } \mathrm{H}
\end{aligned}
$$

Gas Phase $\Delta \mathrm{E}$ :
BLYP $1.9 \mathrm{kcal} / \mathrm{mol}$
HCTH $2.03 \mathrm{kcal} / \mathrm{mol}$
B3LYP/D95* $1.6 \mathrm{kcal} / \mathrm{mol}$
Hutter and Clark, J.Am.Chem.Soc. 118, 7574


## Protonation state of the bases: DNA



## Fate of $\mathrm{G}^{+} \cdot$ in DNA

In duplex DNA the oxidation product is 8 -oxo-guanine ( 8 oxoG)

30,0008 -oxoG in human genome!
Methods Enzymol. 186, 521 (1990)
In water a variety of products are formed



## Oxidation reaction 1



Meta coordinates used:
N1 and water oxygen H-coordination number
Carbon 8 O-coordination number
Rate limited by water autoprotolysis
Catalyzed by $\mathrm{R}_{2} \mathrm{PO}_{4}^{-}$
N1 protonation state matters
$\Delta \mathrm{E}$ N1 deprotonation of 8-OH-G:
$22 \mathrm{kcal} / \mathrm{mol}$ !

## Oxidation reaction



