



the
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ICTP 40th Anniversary

SMR 1564 - 19

SPRING COLLEGE ON SCIENCE AT THE NANOSCALE
(24 May - 11 June 2004)

ELECTRONIC / THERMAL TRANSPORT - Part III

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These are preliminary lecture notes, intended only for distribution to participants.

Electric and Thermal Transport in Nanoscale Materials –Part III

Philip Kim

**Department of Physics
Columbia University**

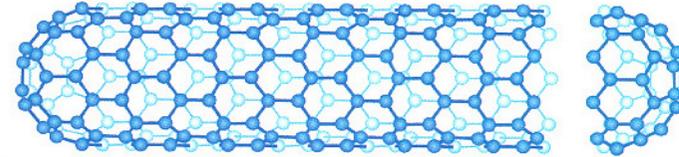


Outline

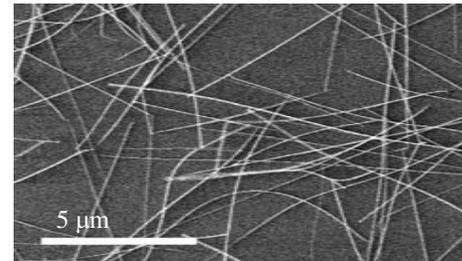
- Charge Transport and Energy Dissipation
- Mesoscopic Heat Transport Measurements
- Mesoscopic Thermoelectric Effects
- Field Effect Transport in 2D Crystallites

Nanoscale Materials

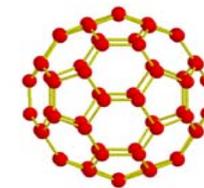
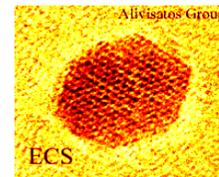
- Carbon Nanotubes



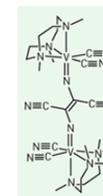
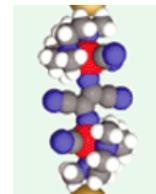
- Nanowires



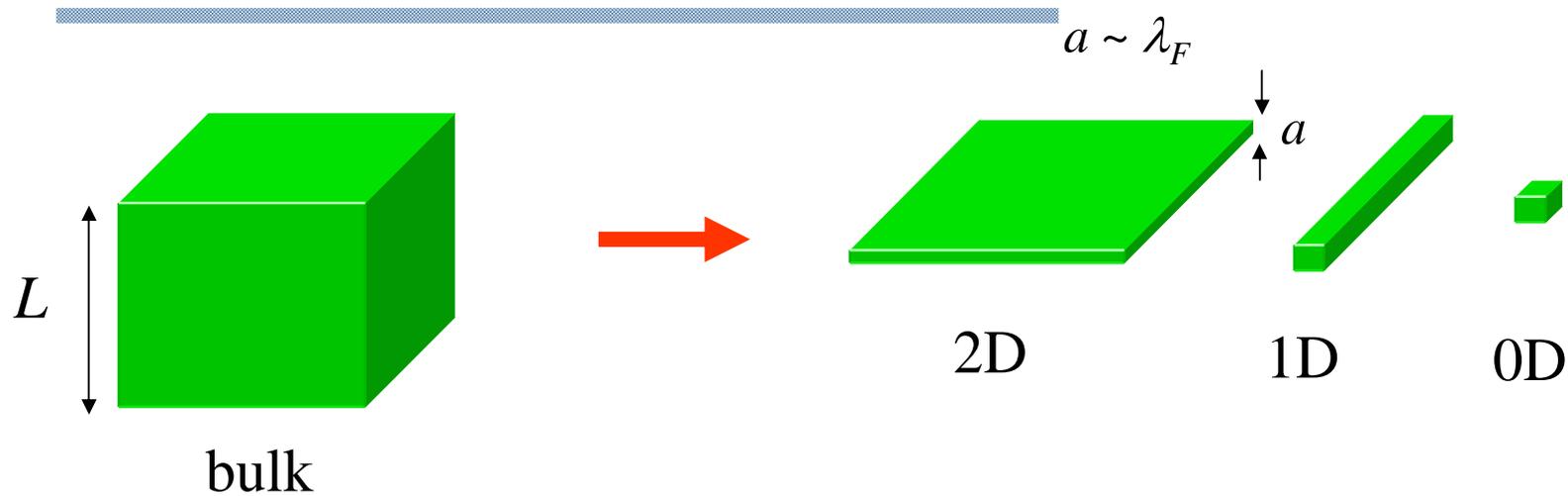
- Nanocrystals, Fullerenes



- Molecules



Electron Transport Low dimensional Nanostructures



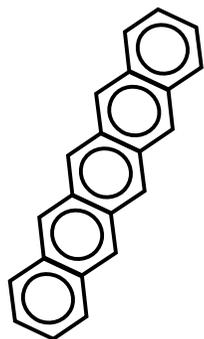
Why low dimensions?

- Fermi surface topology
- Singularities in Density of State
- Strong Interactions
- Strong Fluctuations

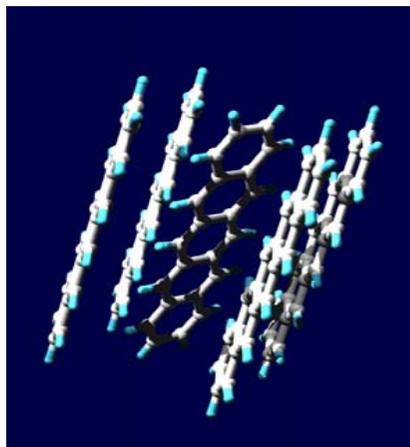
→ New Phenomena

Gate Electric Field Effect: tuning charge density in the samples

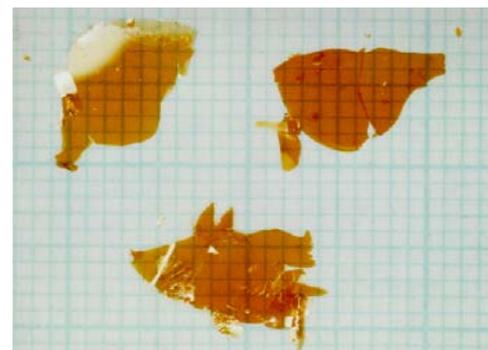
Organic Thin Film : 2D Array of organic molecules



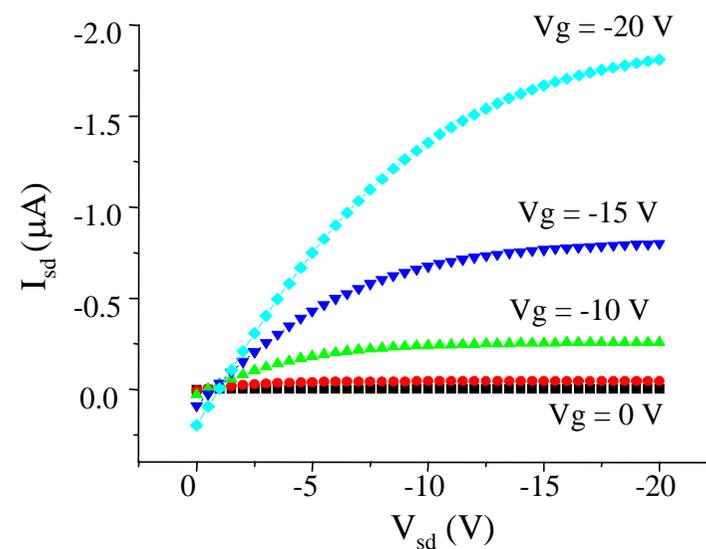
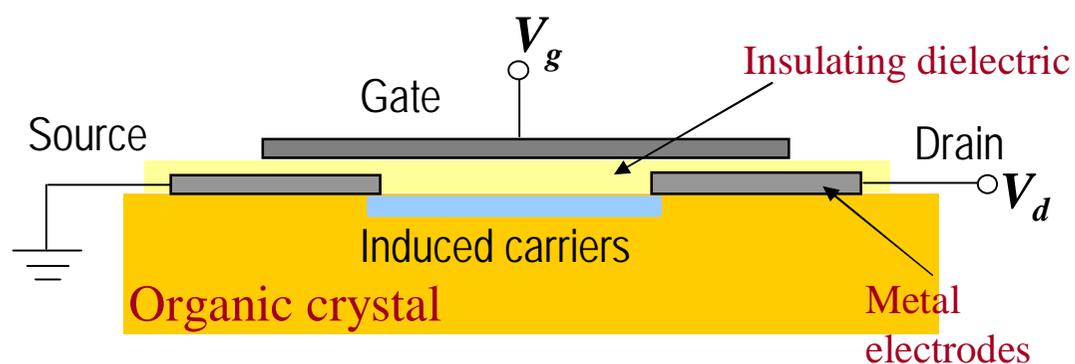
Pentacene



Herringbone packing

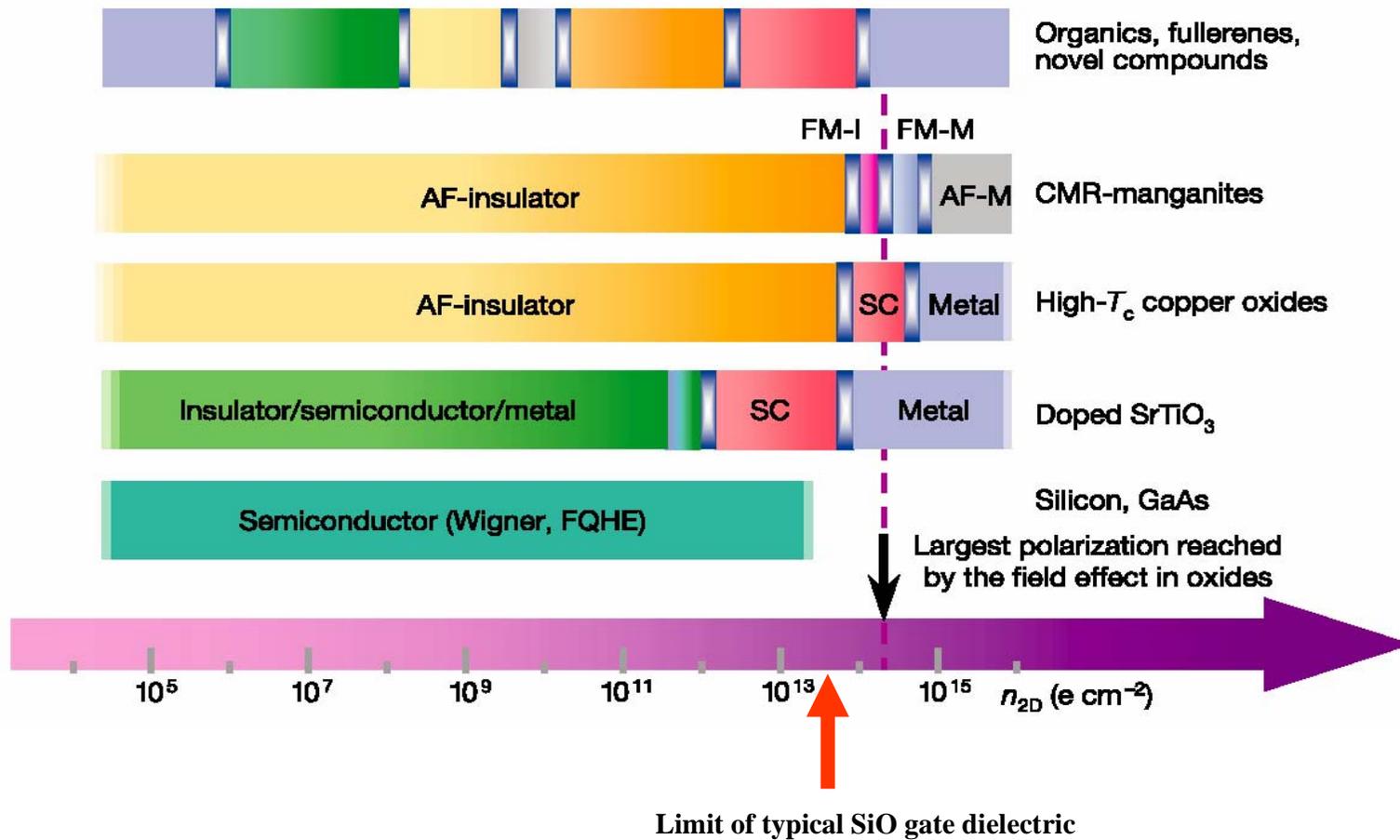


Organic single crystal

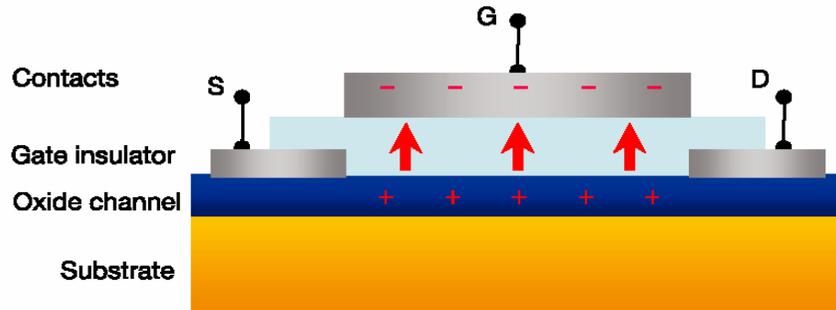


Electric Field Effect in Various Materials

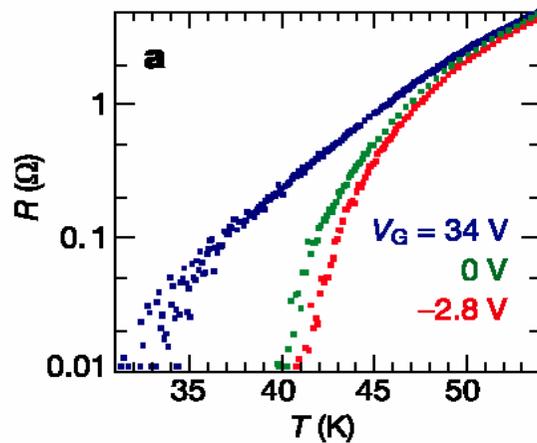
Ahn, Triscone, and Mannhart, Nature 424, 1015 (2003)



Field Effects in Superconducting Films

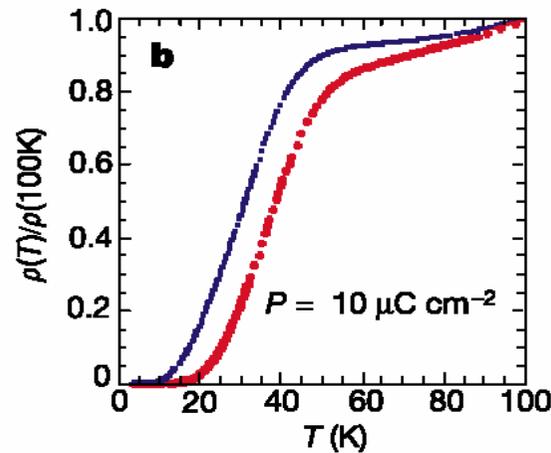


Mannhart, J. High- T_c transistors. *Supercond. Sci. Technol.* **9**, 49–67 (1996).

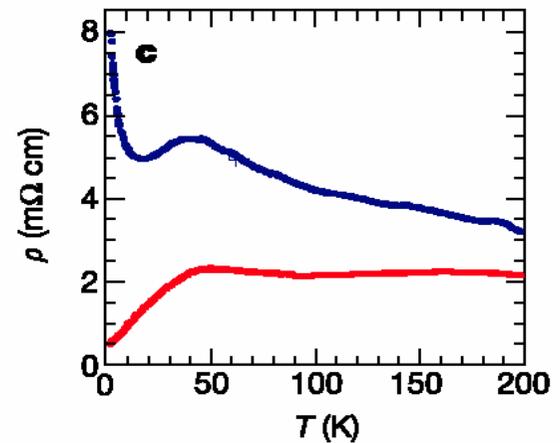


8 nm thick $\text{YBa}_2\text{Cu}_3\text{O}_{7-d}$
/ BaSrTiO_3 gate dielectric

Ahn, C. H. *et al.* Electrostatic modulation of superconductivity in ultra-thin $\text{GdBa}_2\text{Cu}_3\text{O}_7$ films. *Science* **284**, 1152–1155 (1999).

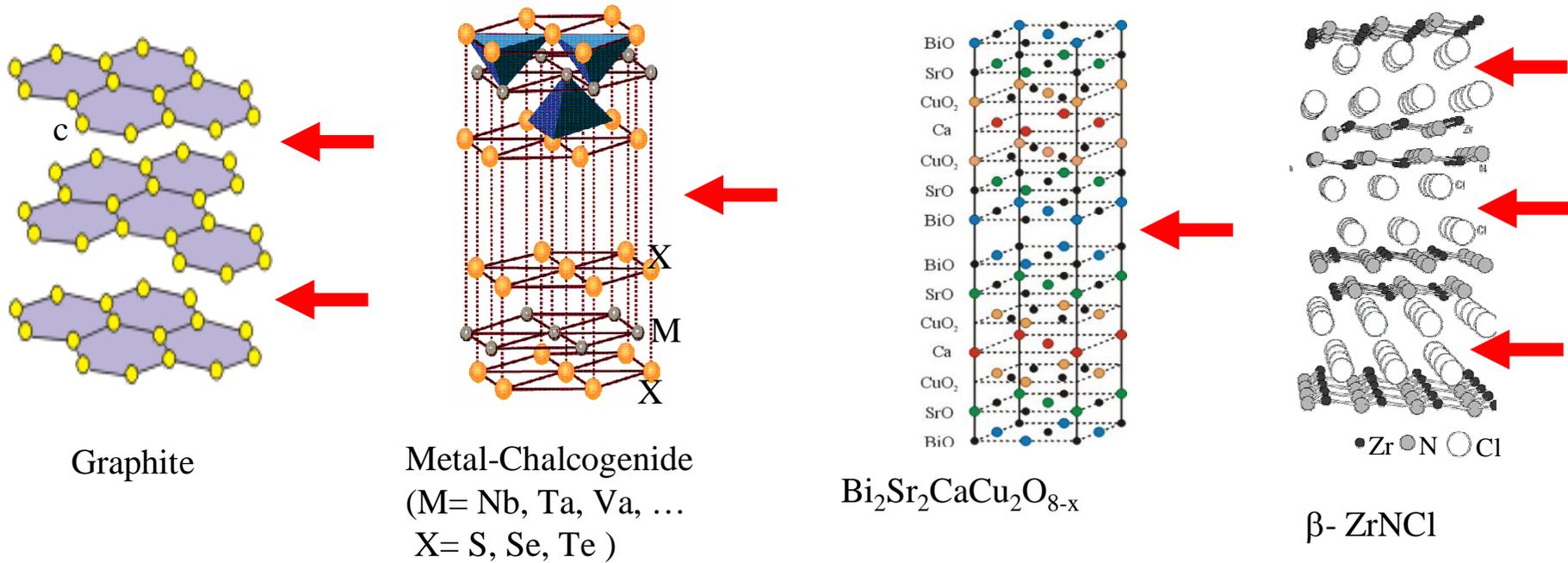


2 nm thick $\text{GdBa}_2\text{Cu}_3\text{O}_{7-d}$
/ PZT layer gate dielectric



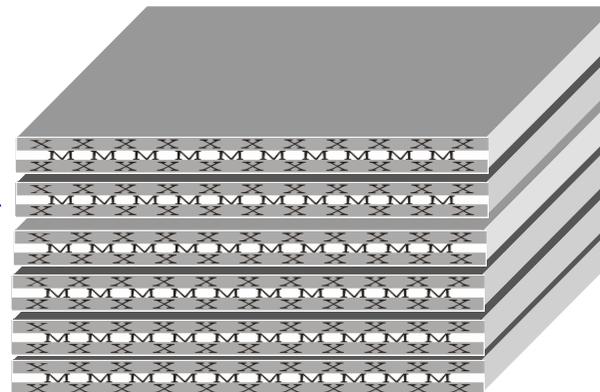
2 nm thick $\text{GdBa}_2\text{Cu}_3\text{O}_{7-d}$, doped
/ PZT layer gate dielectric

Van der Waals Layered Materials

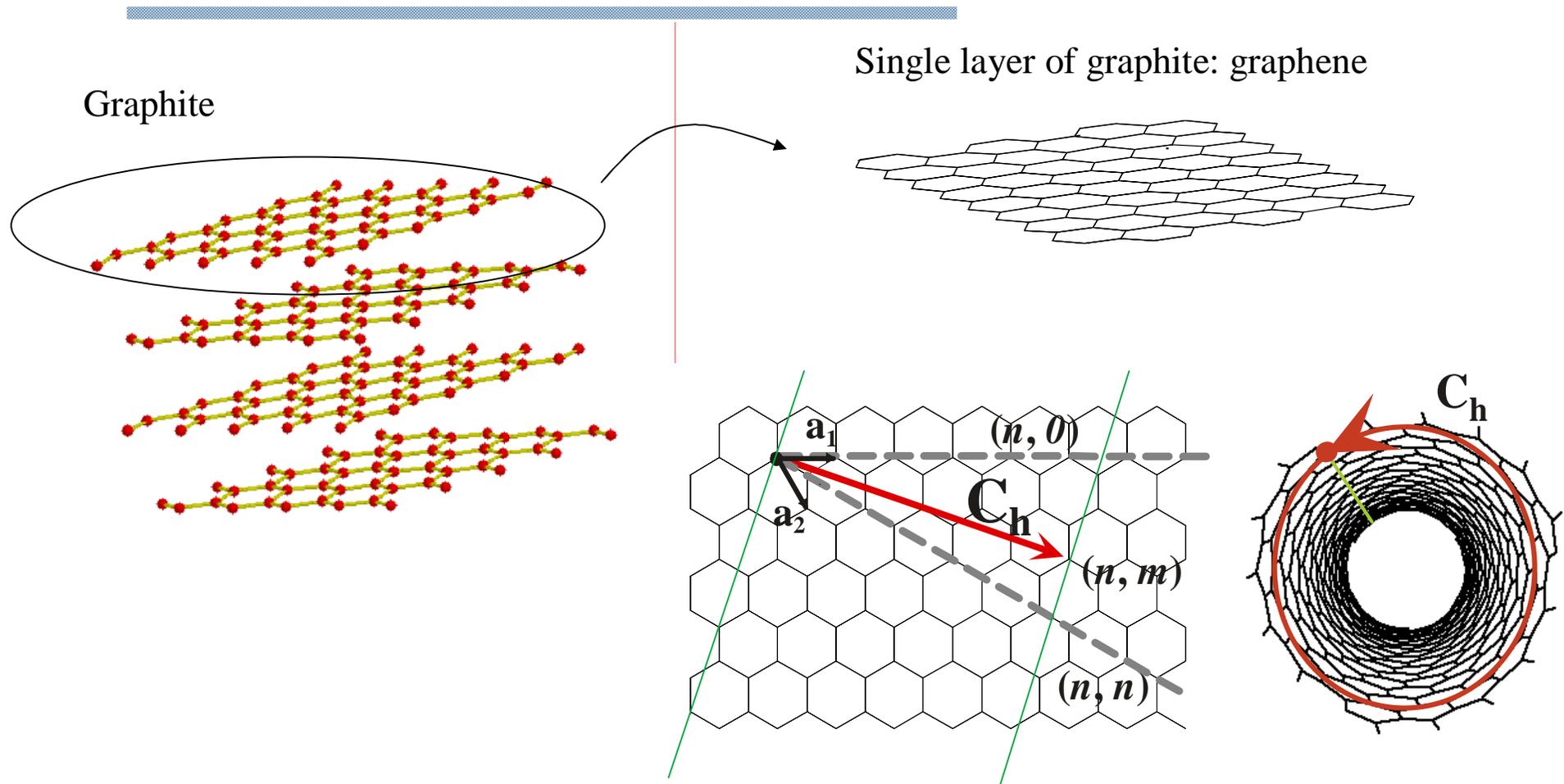


Van der Waals bonding between layers

Covalent bonding within layers



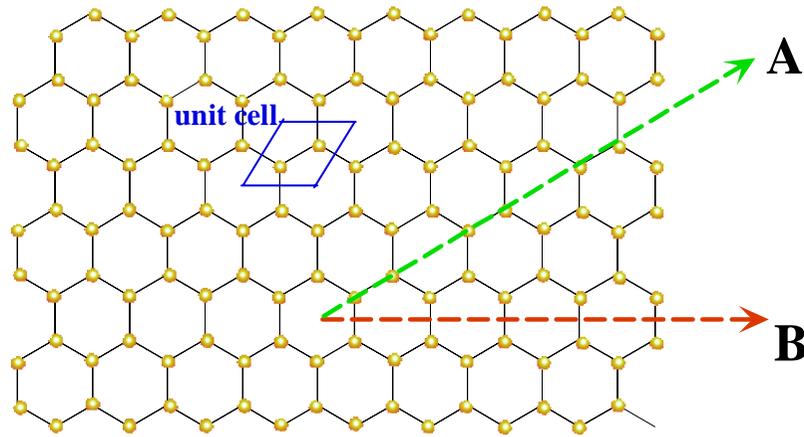
Graphene : Single Sheet of Graphite Layer



Rolling up graphene along C_h imposes a Periodic boundary condition :

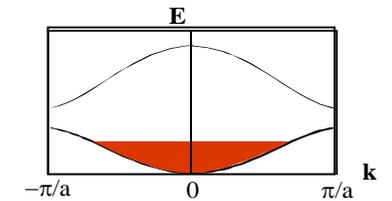
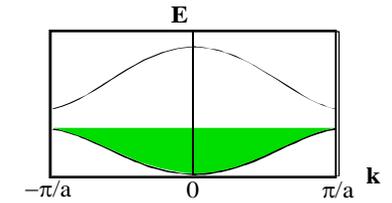
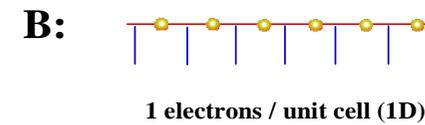
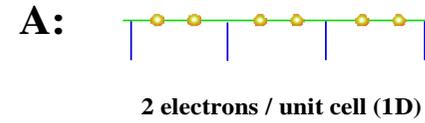
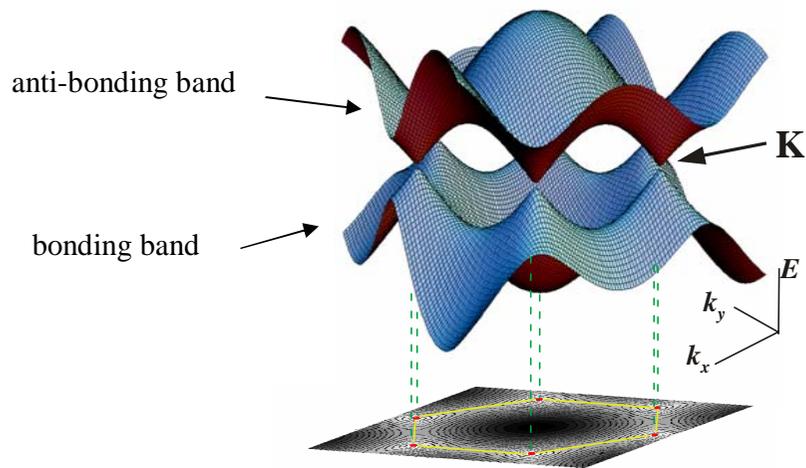
$$C_h \cdot \mathbf{k} = 2\pi q$$

Electronic Band Structure of Graphene

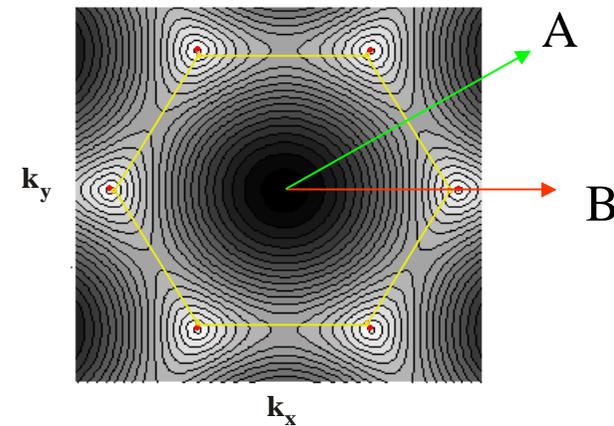


2 electrons / unit cell (2D)

π/π^* band of Graphene

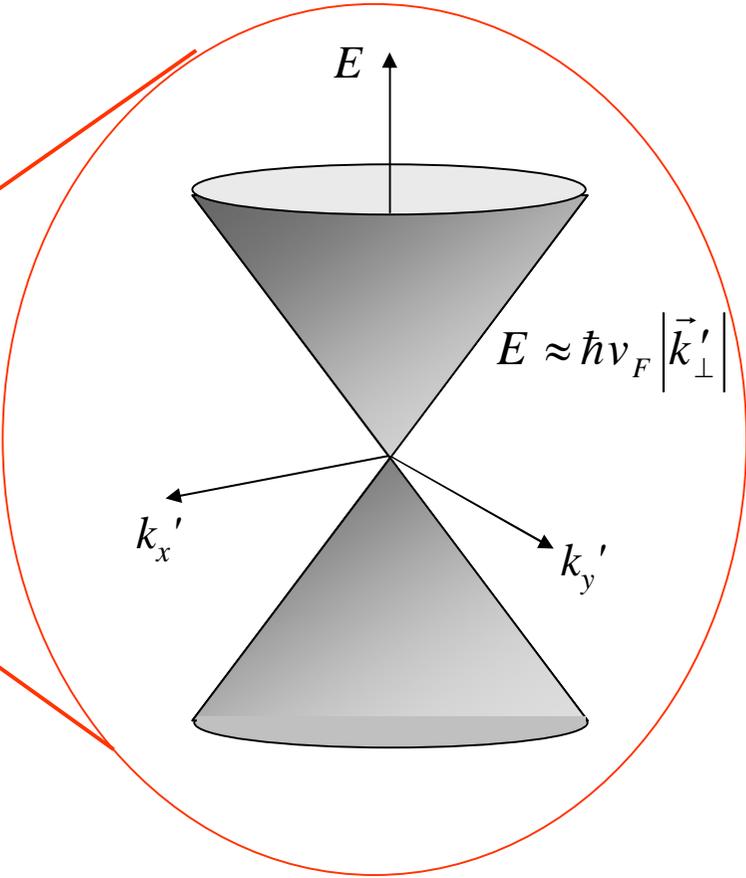
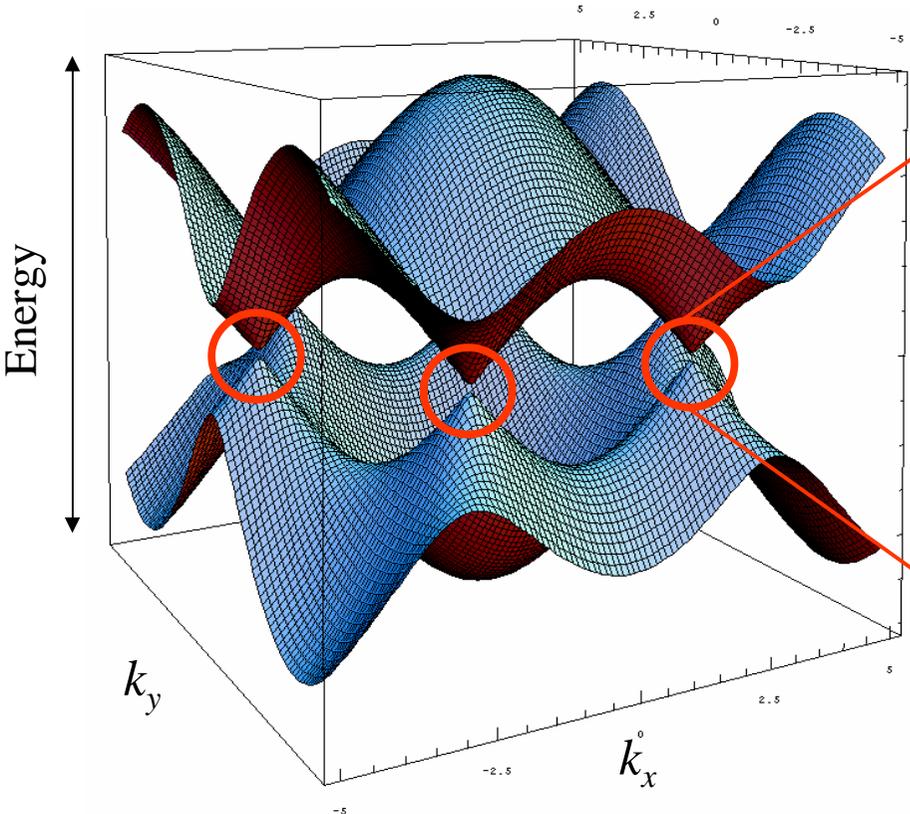


2D Brillouin Zone



Graphene : Dirac Particles in 2D Box

Band structure of graphene



k·p perturbation



Effective Hamiltonian:

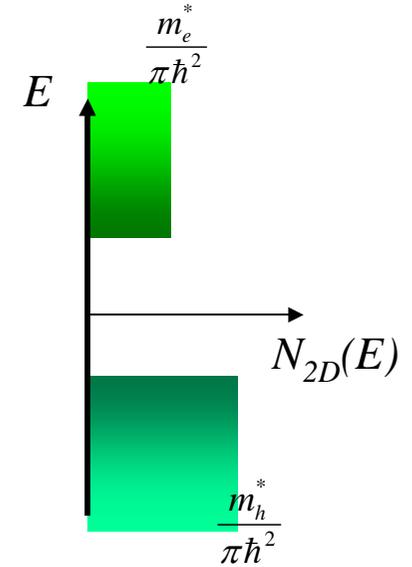
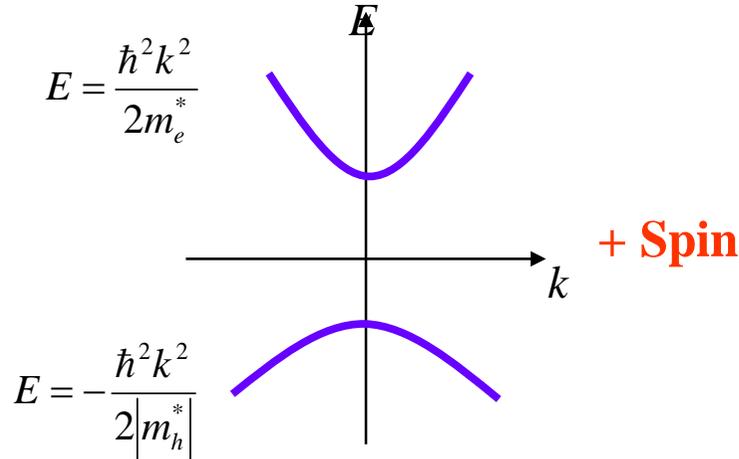
$$H_{eff} = \hbar v_F \vec{\sigma} \cdot \vec{k}'_{\perp}$$

└ "helicity"

Divincenzo and Mele, PRB (1984)

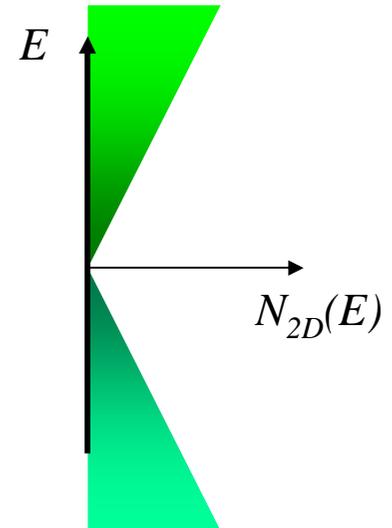
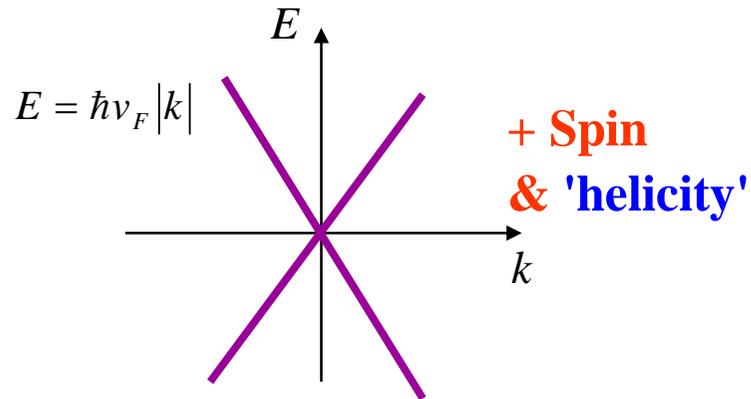
Graphene v.s. Conventional 2D Electron System

Conventional 2D Electron System

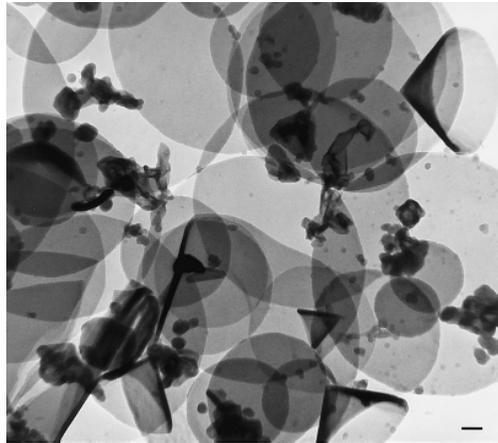


Graphene

- Zero band mass
- Strict electron hole symmetry
- Electron hole degeneracy

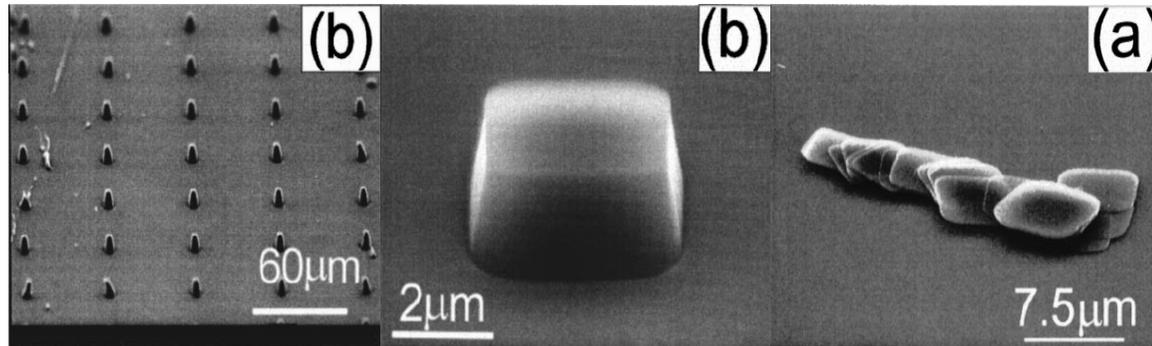


Obtaining thin graphite crystallites – Previous Efforts



A. Krishnan et. al. Nature 1997

By product of carbon nanotubes: disordered



Patterned Graphite Crystallites

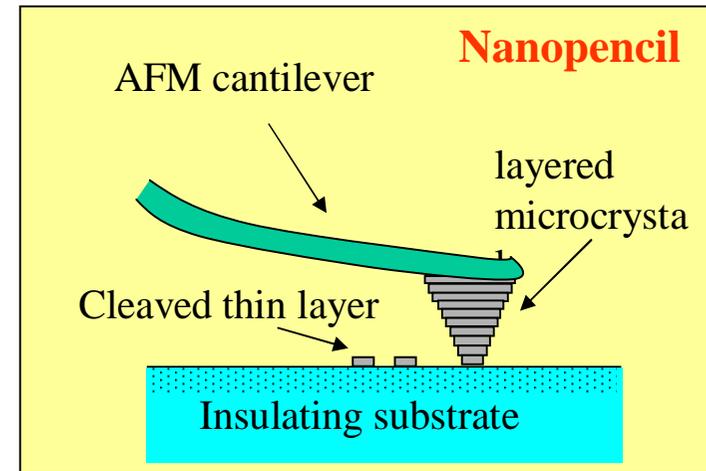
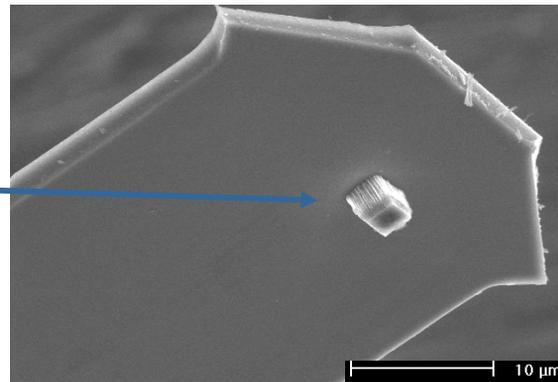
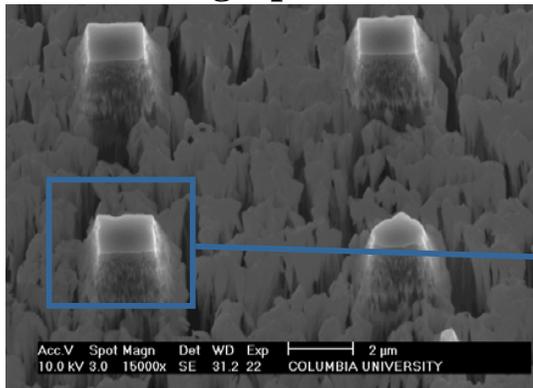
Xuekun Lu et. al. APL 1999

Thinnest crystal reported ~ 30nm

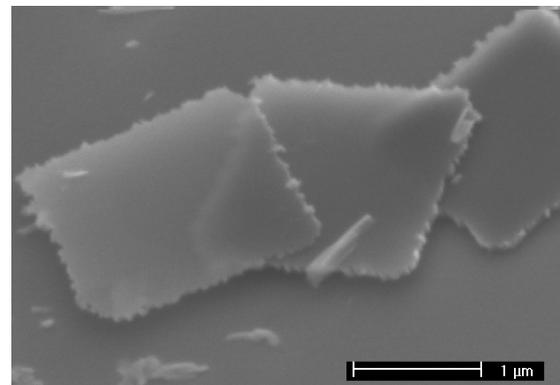
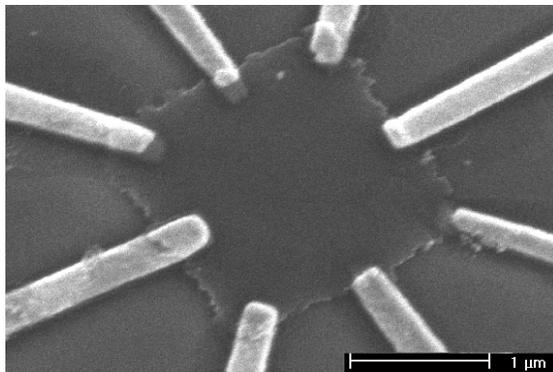
No Systematic transport study

Graphite Nanocrystals Hybrid Device

Patterned graphite blocks



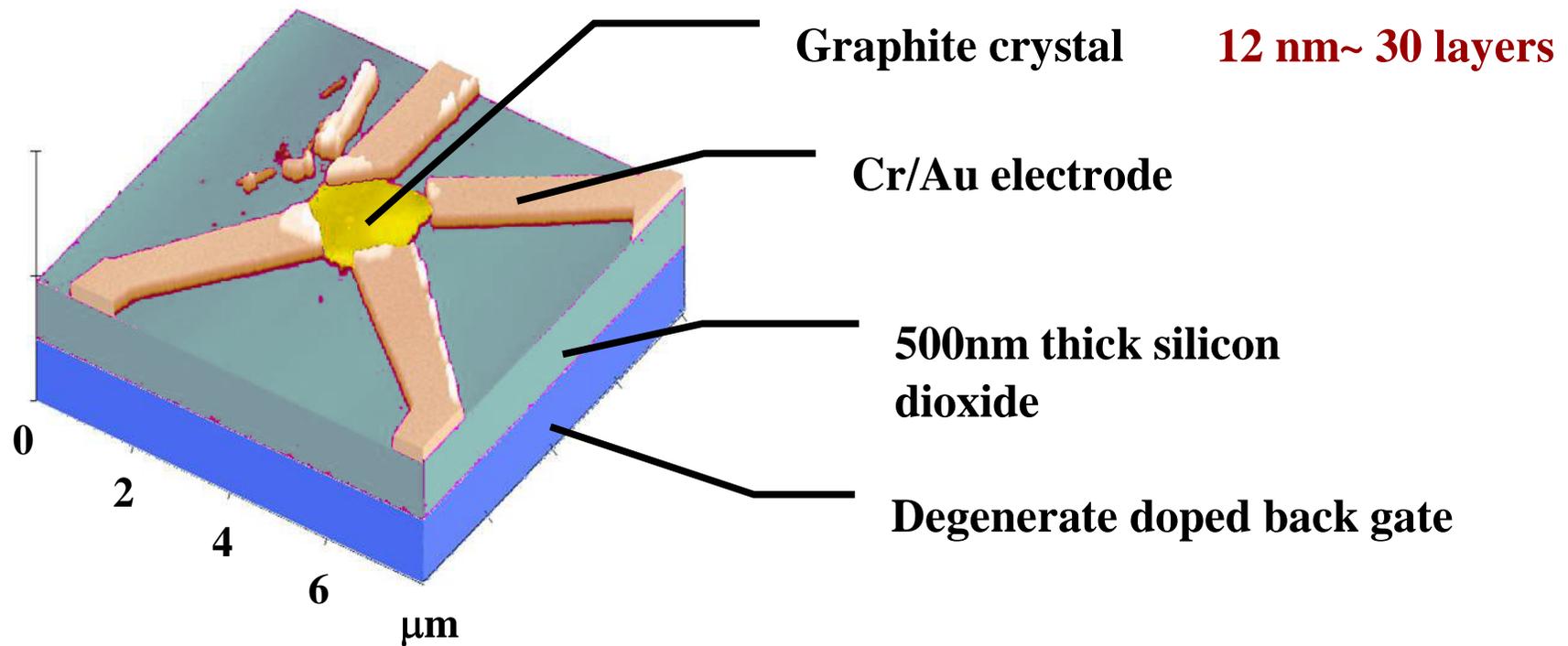
microfabricated device



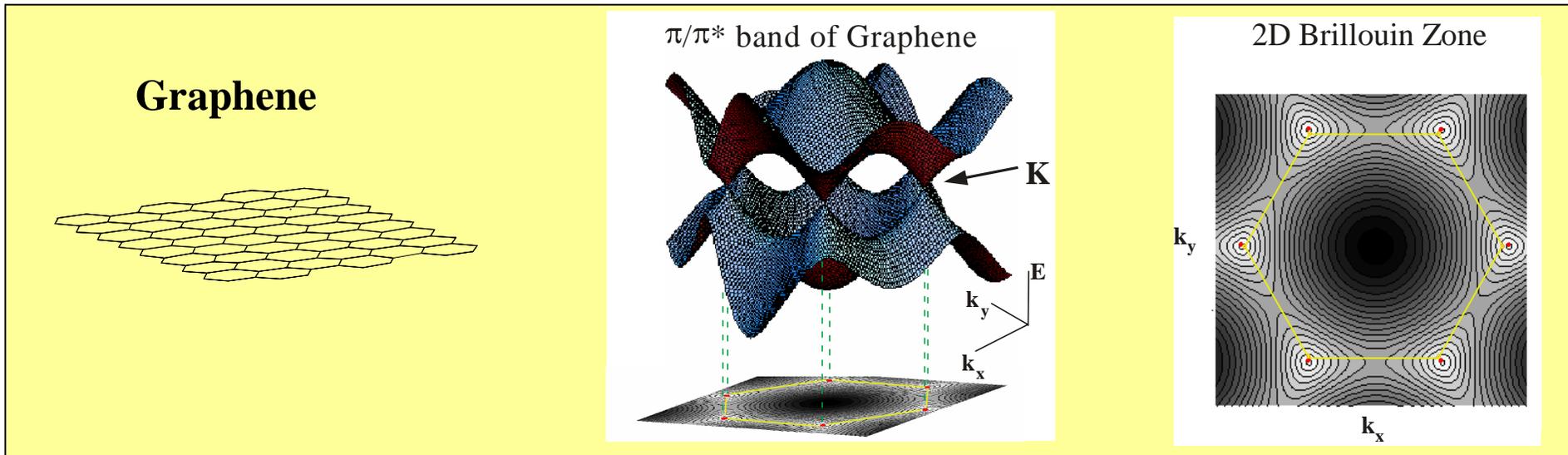
cleaved nanocrystals
5 - 100 nm thick

Device Structure

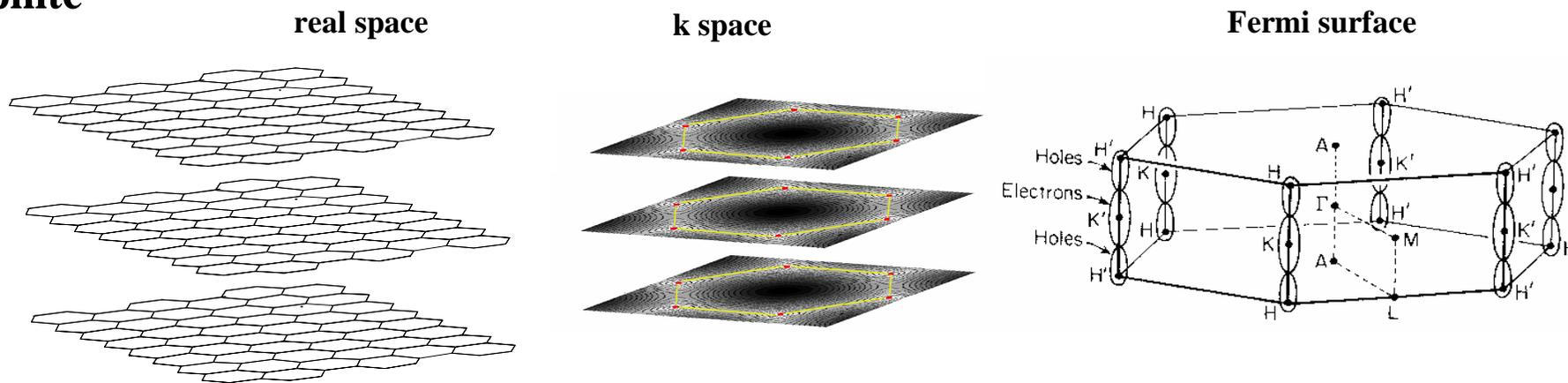
Thin Graphite Crystal Device Fabricated Using Electron Beam Lithography



Band Structure of Graphite

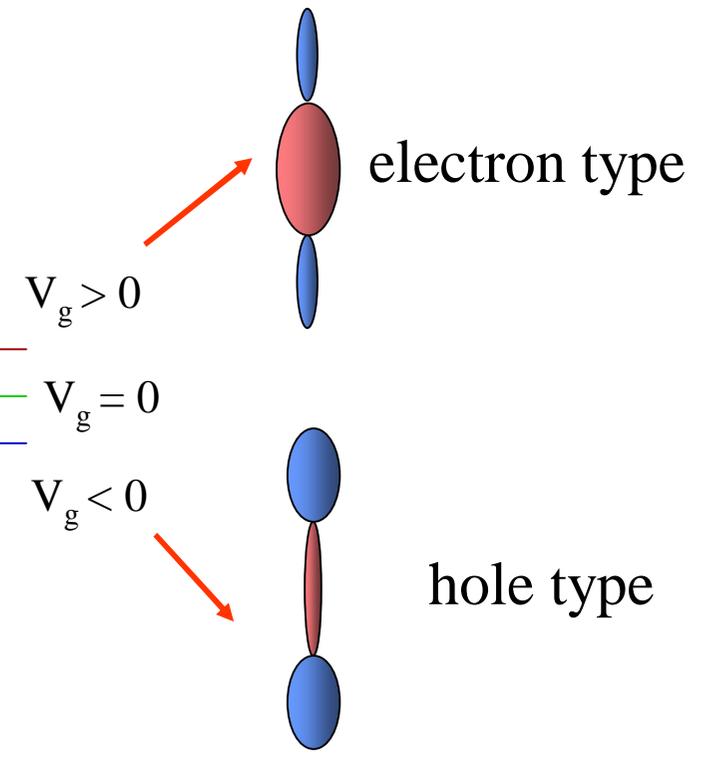
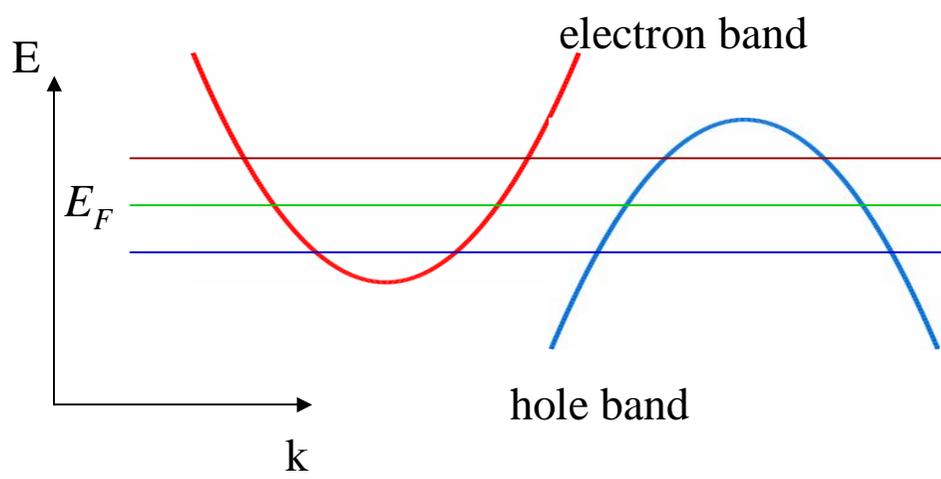
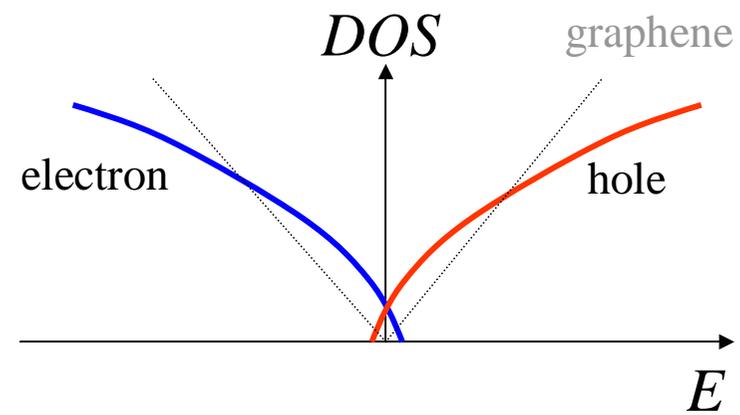
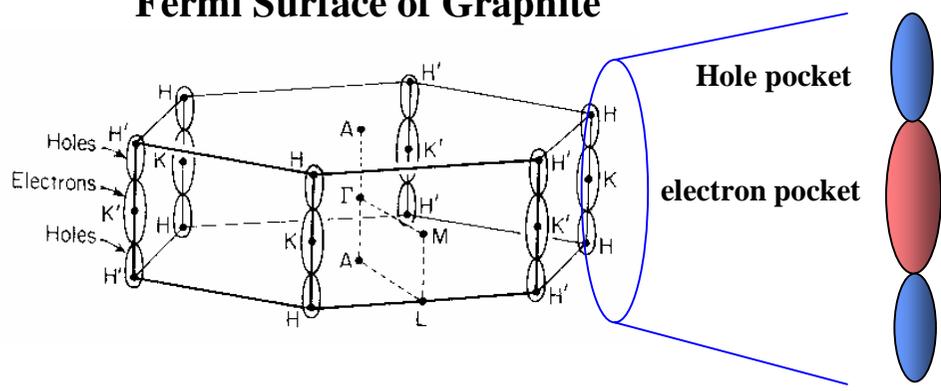


Graphite

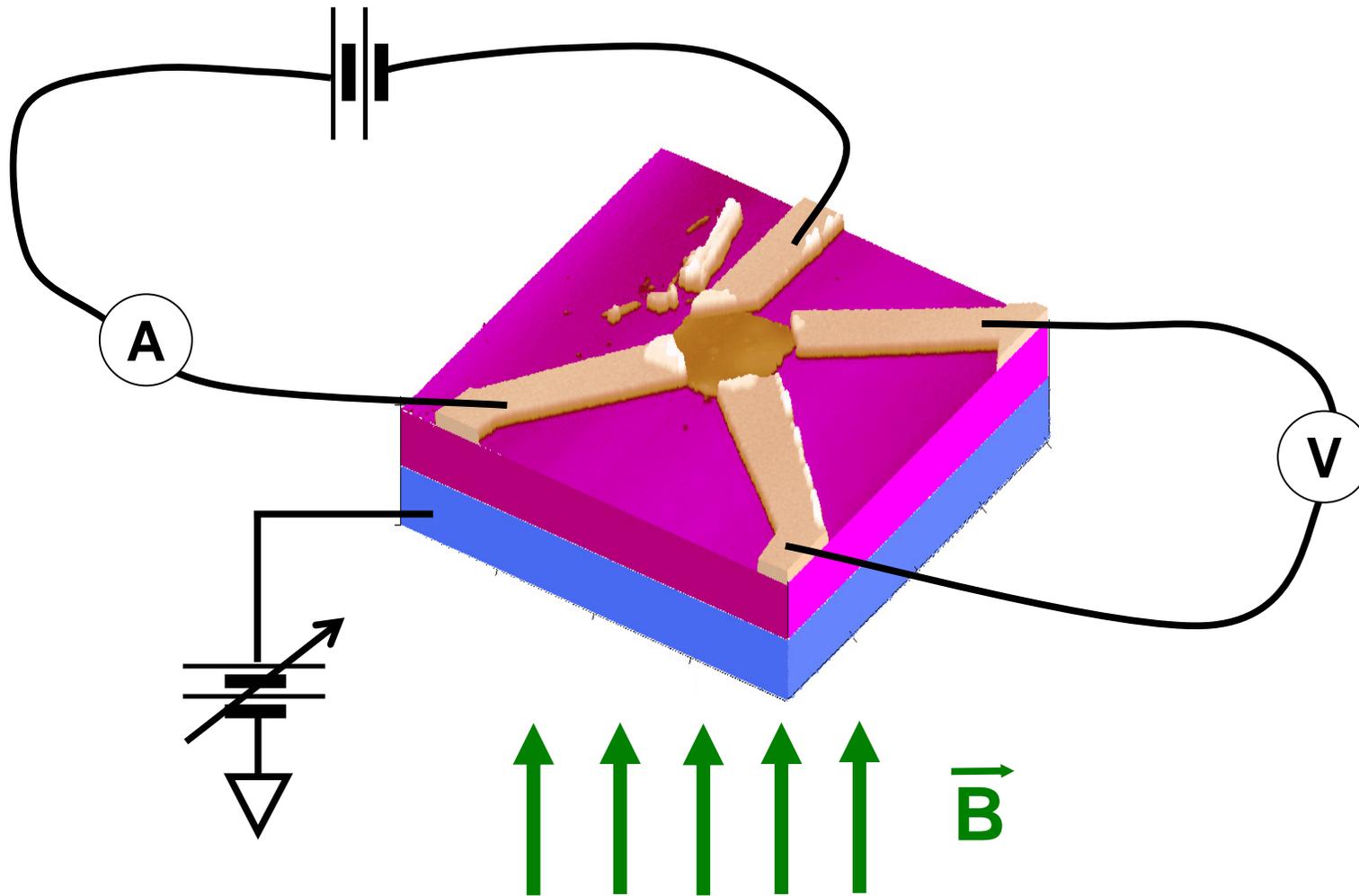


Two Band Model

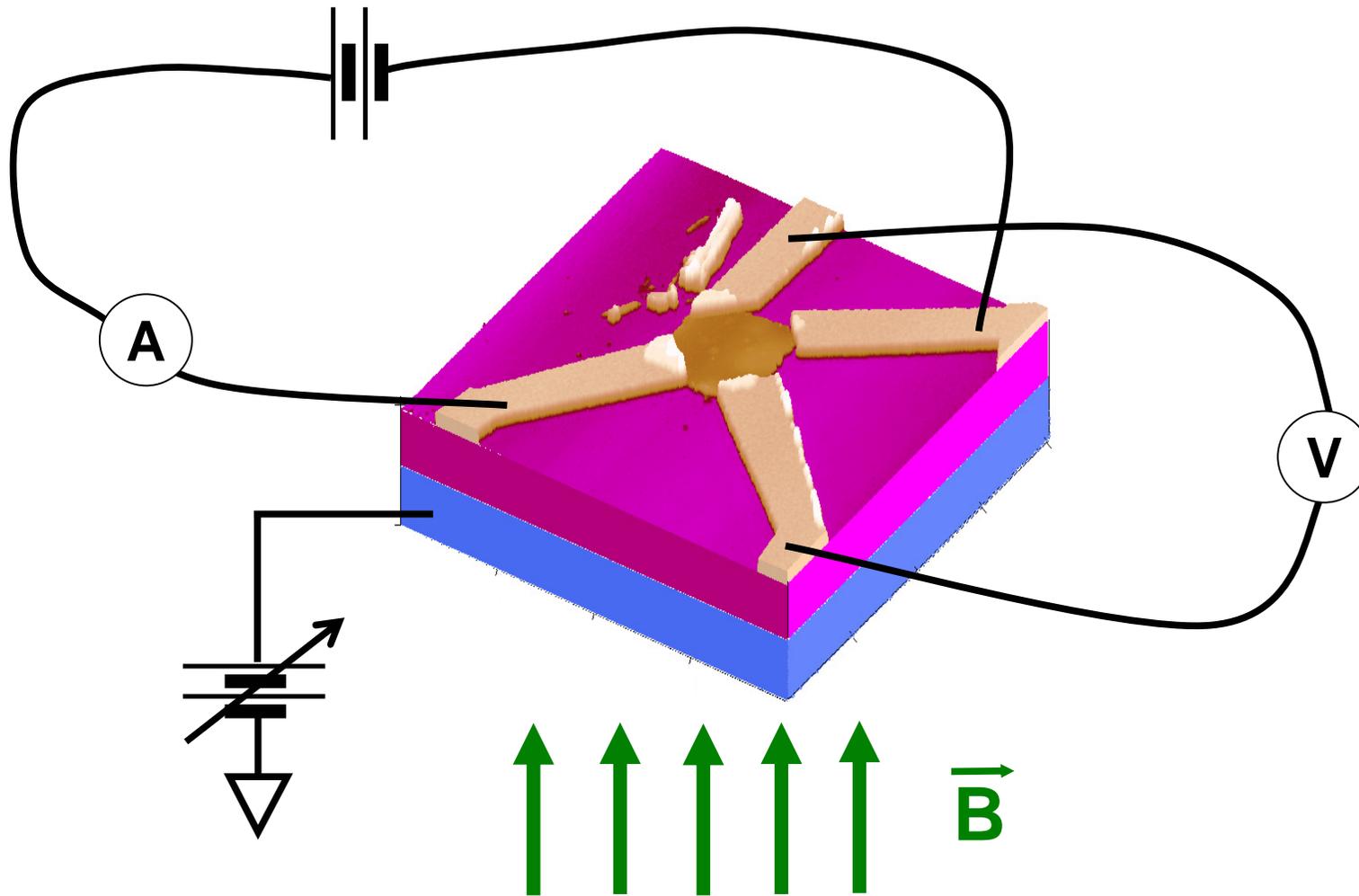
Fermi Surface of Graphite



Magneto-resistance Measurement



Hall Measurement



Magneto and Hall Resistance in Two Band Model

resistivity tensor: $\overleftrightarrow{\rho}_j = \begin{pmatrix} \rho_j & -R_j B \\ R_j B & \rho_j \end{pmatrix}$

j : band index, $R_j = -1/en_j$, $\rho_j = 1/e\mu_j n_j$

$$\rho = \frac{\rho_1 \rho_2 (\rho_1 + \rho_2) + (\rho_1 R_2^2 + \rho_2 R_1^2) B^2}{(\rho_1 + \rho_2)^2 + (R_1 + R_2)^2 B^2}$$

$$R = \frac{R_1 \rho_2^2 + R_2 \rho_1^2 + R_1 R_2 (R_1 + R_2) B^2}{(\rho_1 + \rho_2)^2 + (R_1 + R_2)^2 B^2}$$

Small B ,

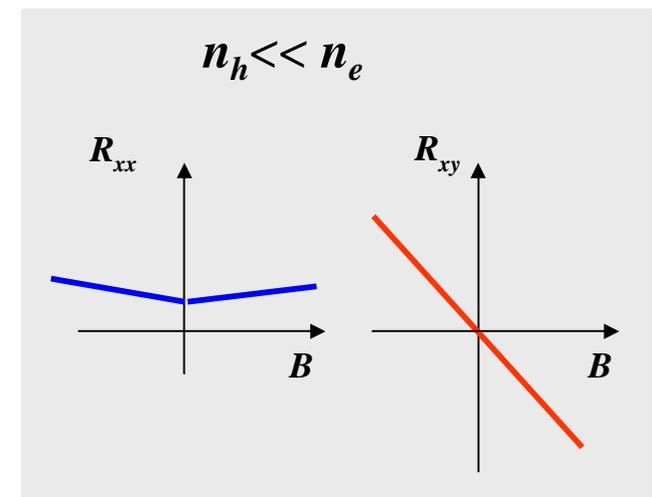
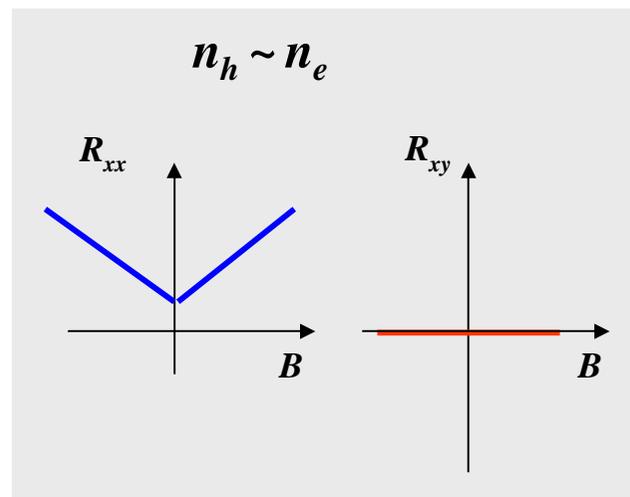
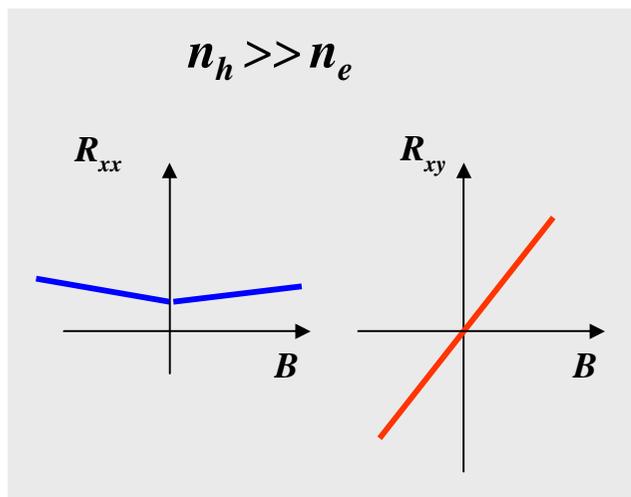
$$\mu_e = \mu_h$$

$$n_e + n_h = \text{const.}$$



Magnetoresistance: $\Delta R_{xx} \sim n_e n_h |B|^\alpha$

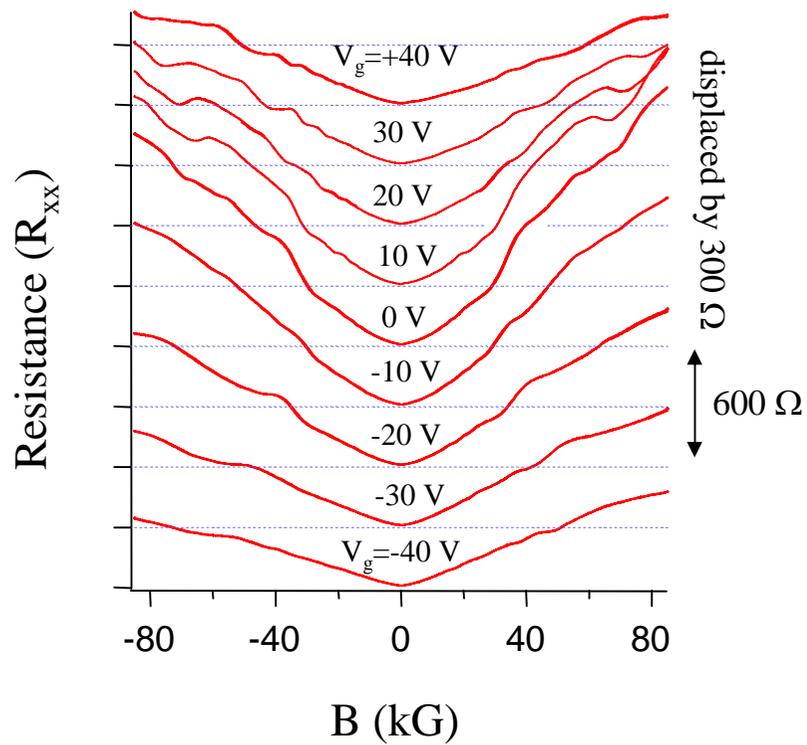
Hall resistance: $R_{xy} \sim (n_h - n_e) B$



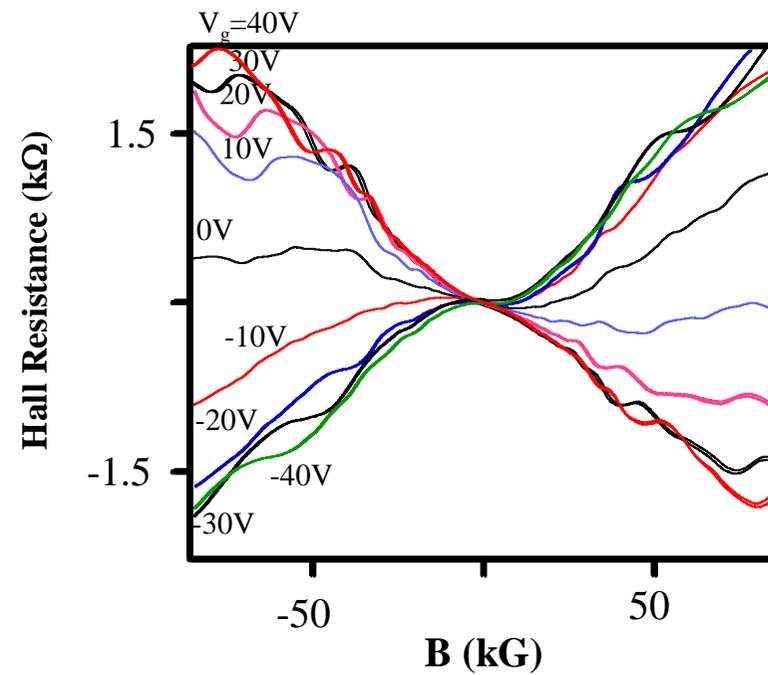
Hall Resistance and Magneto Resistance

Sample: 12 nm thick
T = 1.7 K

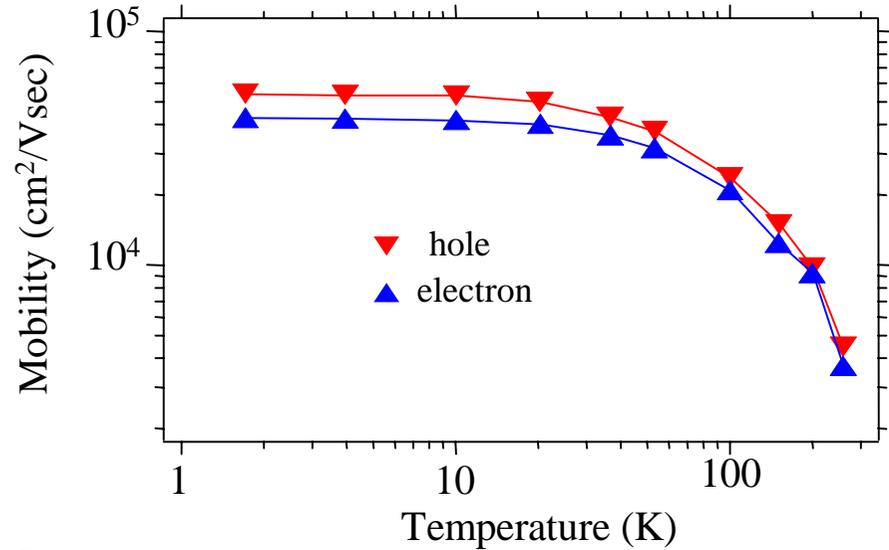
$$\Delta R_{xx} \sim n_e n_h B^\alpha$$



$$R_{xy} \sim (n_h - n_e)B$$



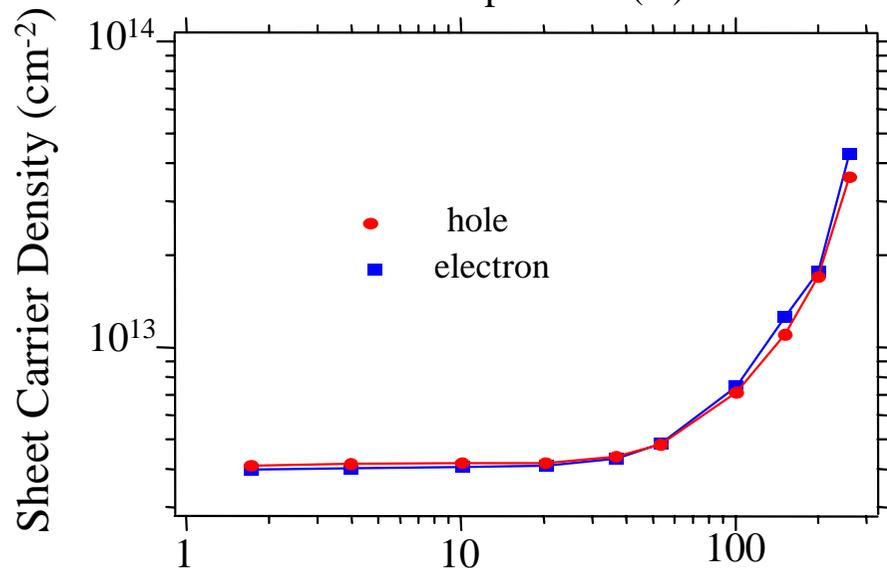
Mobility and Carrier Density



12 nm Graphite Crystal:

Mobility at 1.7K ~ 50000 cm²/Vs

Carrier Density at 1.7K ~ 3.3 × 10¹⁸ cm⁻³
(sheet charge density ~ 4.0 × 10¹² cm⁻²
per a graphene layer ~ 1.2 × 10¹¹ cm⁻²)



Bulk graphite:

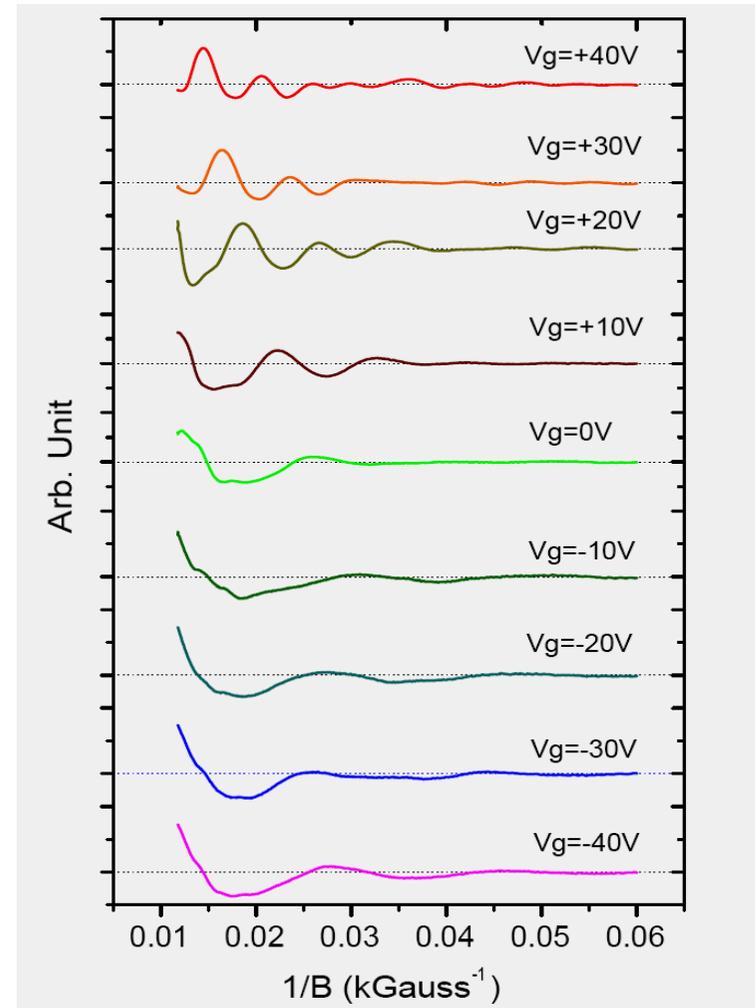
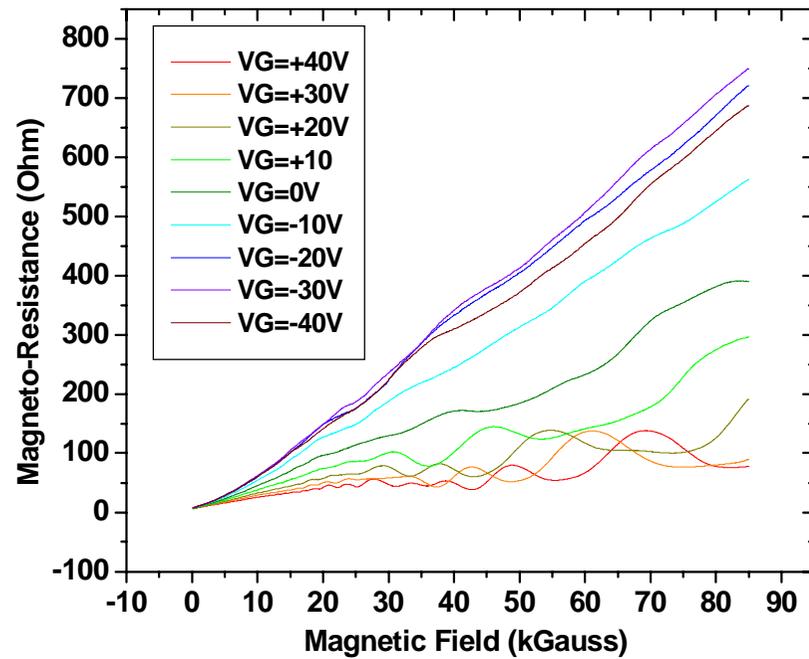
Mobility ~ 70000 cm²/Vs

Carrier Density ~ 3 × 10¹⁸ cm⁻³

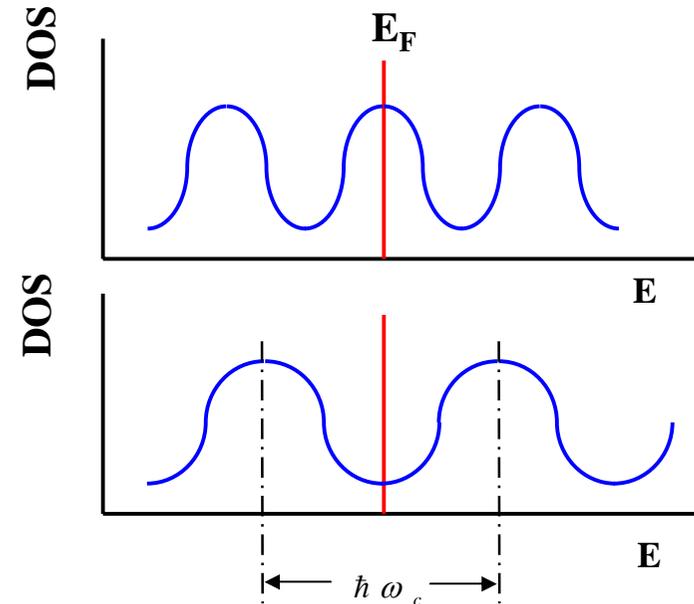
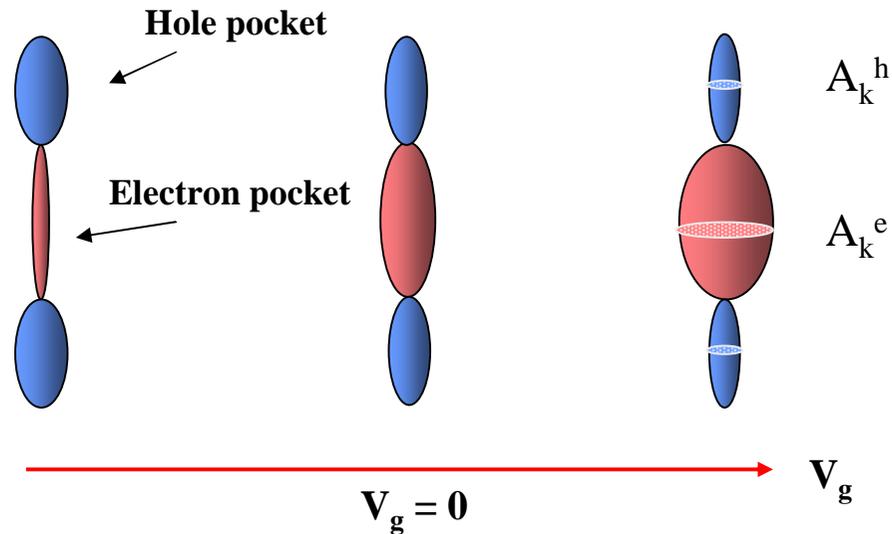
I. Spain et. al., 1967

Shubnikov de Haas Oscillations

Sample: 12 nm thick
 $T = 1.7$ K



Shubnikov de Haas Oscillation



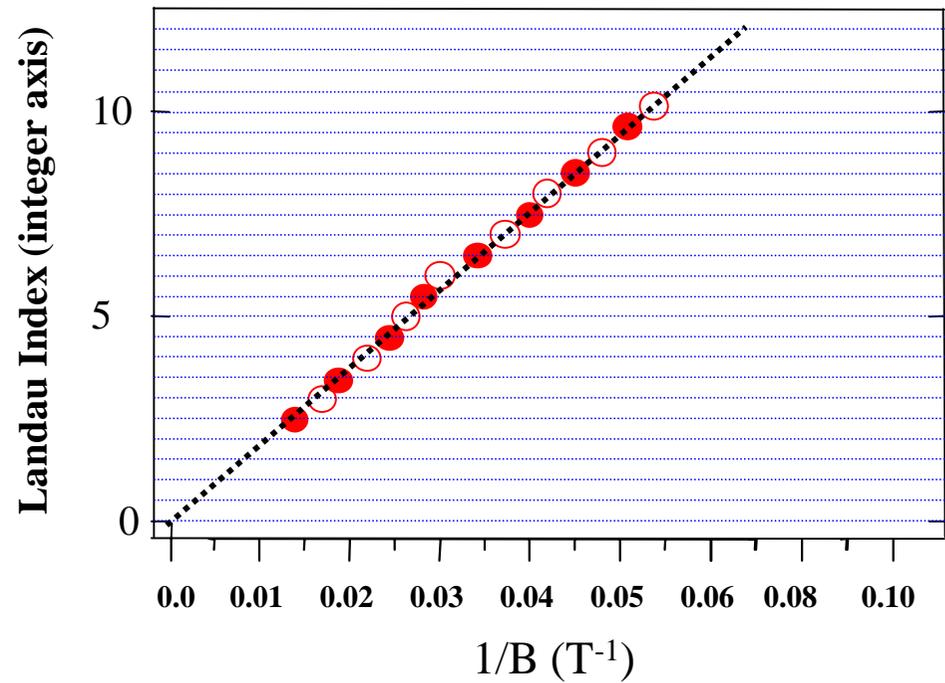
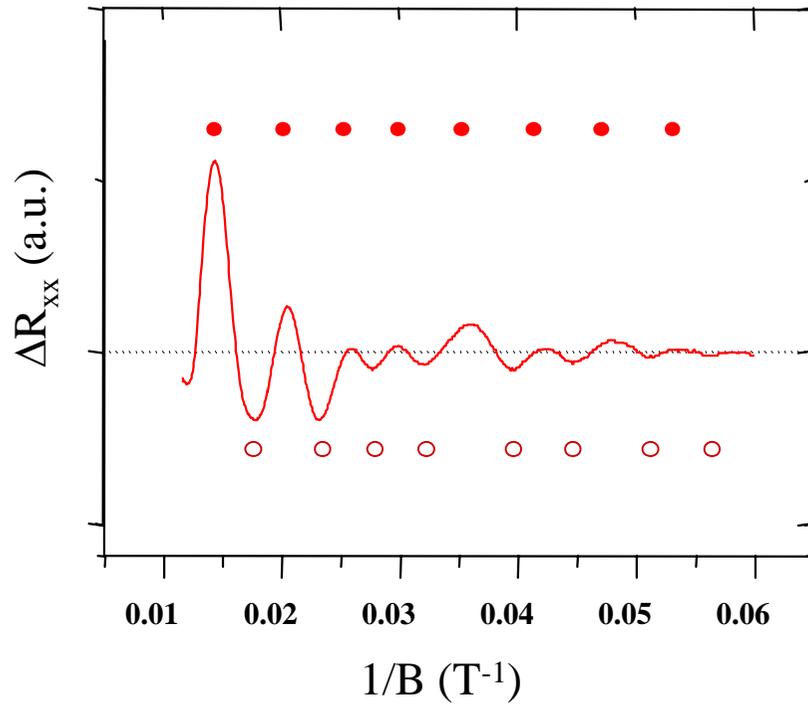
Shubnikov de Haas Oscillation

$$\Delta\left(\frac{1}{H}\right) = \frac{2\pi e}{\hbar c A_k}$$

A_k : maximal cross sectional area of pockets

Magneto-Resistance will show a peak whenever Fermi Level passes a Landau Level

SdH Period and Landau Plot

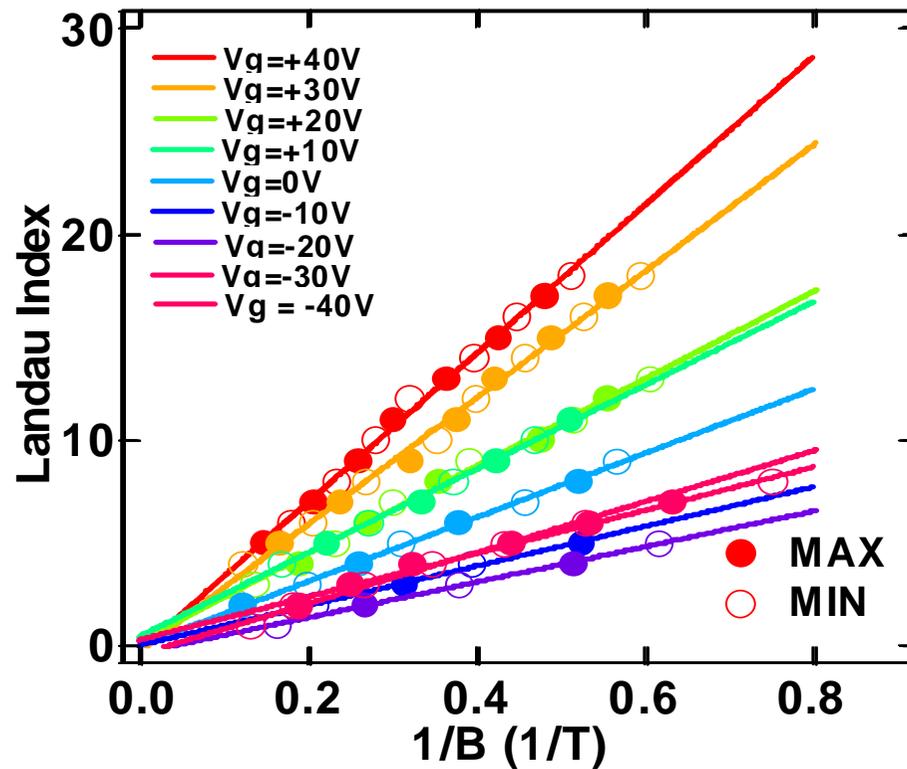


Slope of Landau Plot = Frequency of SdH oscillations

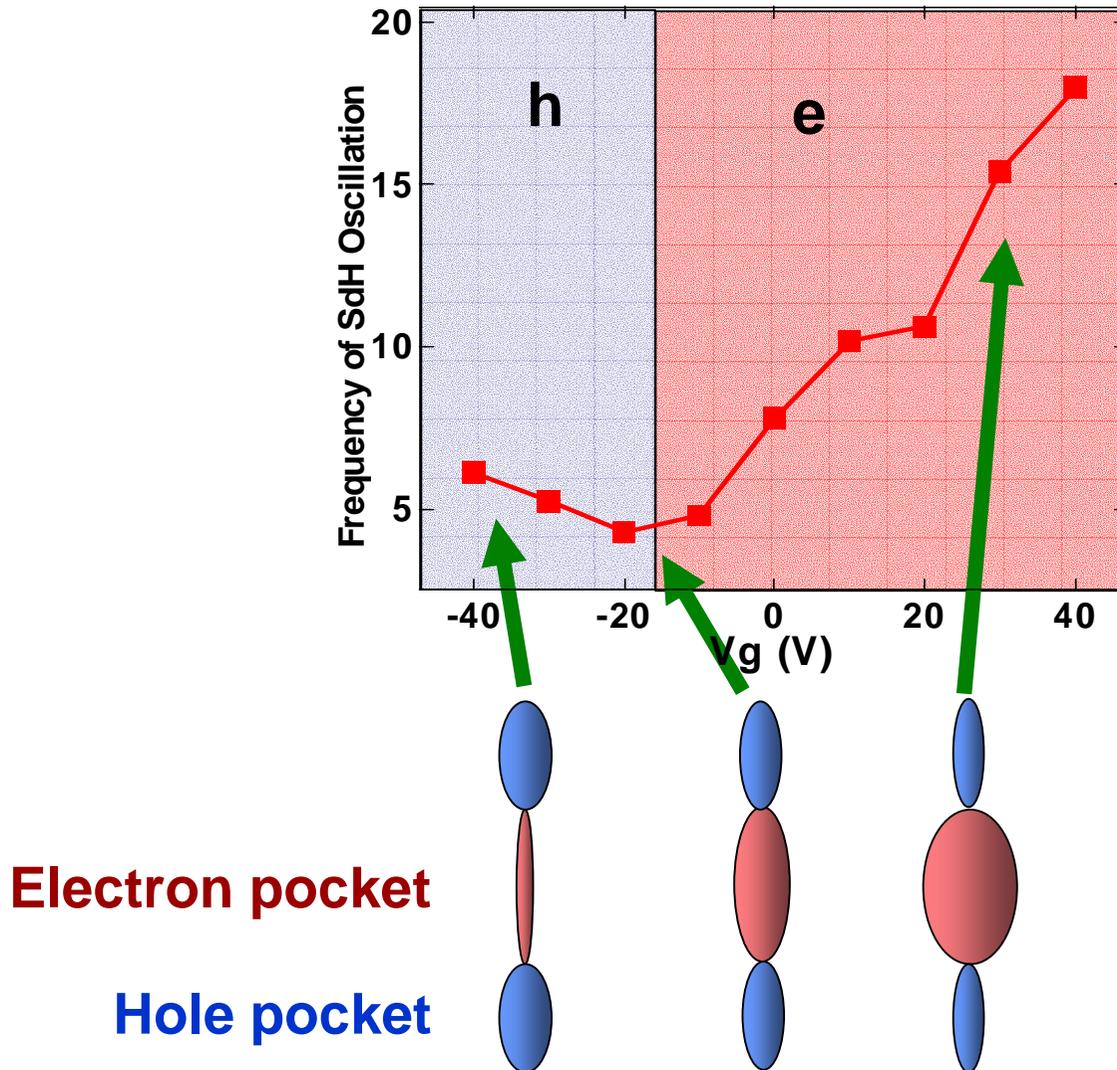
Gate Modulation of Shubnikov de Haas Oscillations

Slope of Landau Plot

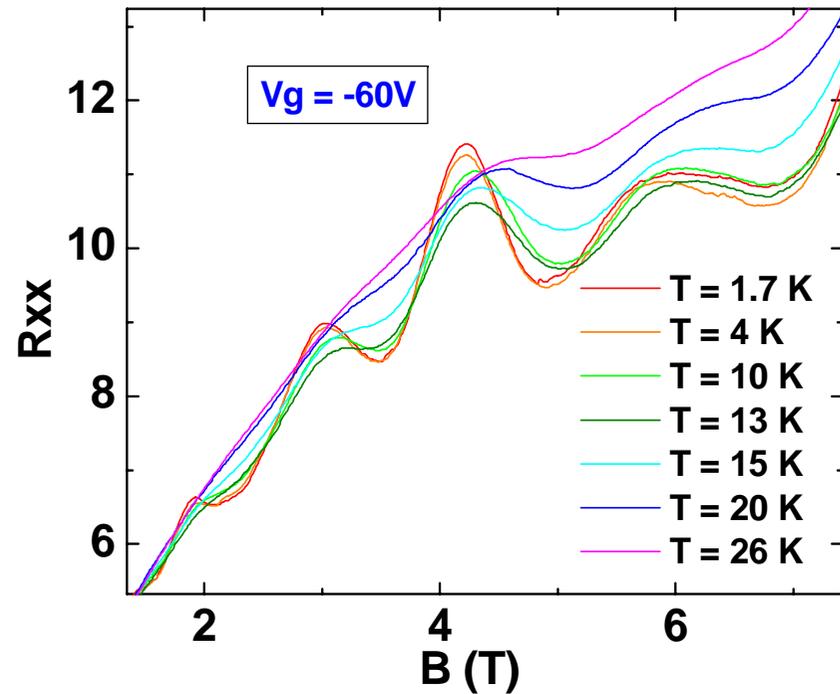
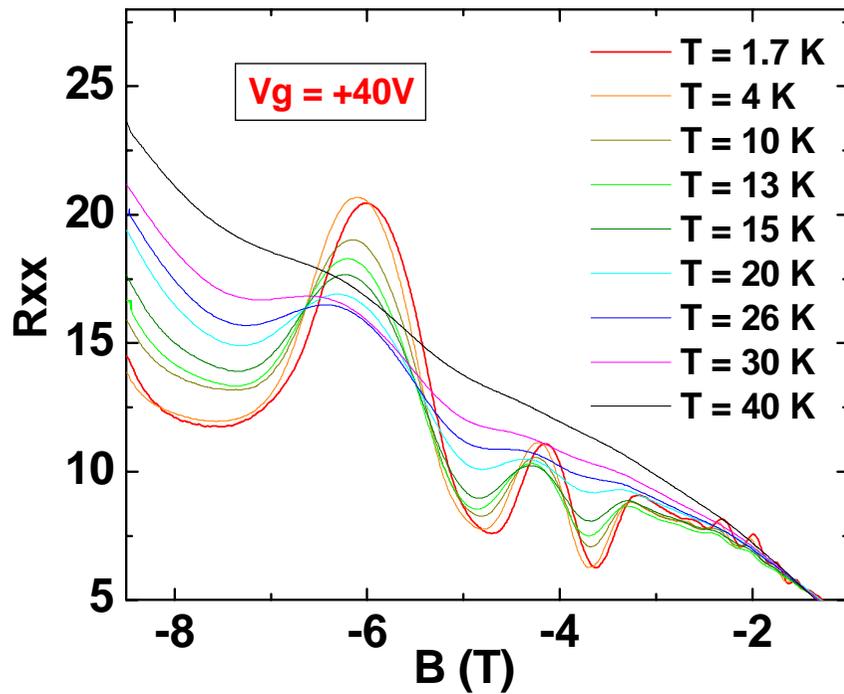
$$\frac{1}{\Delta(1/B)} = \frac{\hbar c}{2\pi e} \cdot A_k$$



Gate Modulation of Shubnikov de Haas Oscillations



Temperature Damping of SdH Oscillations



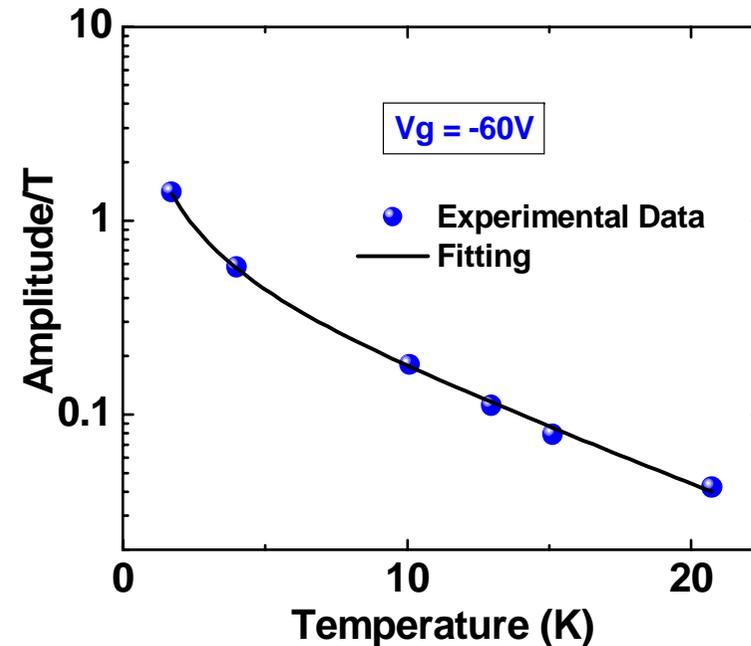
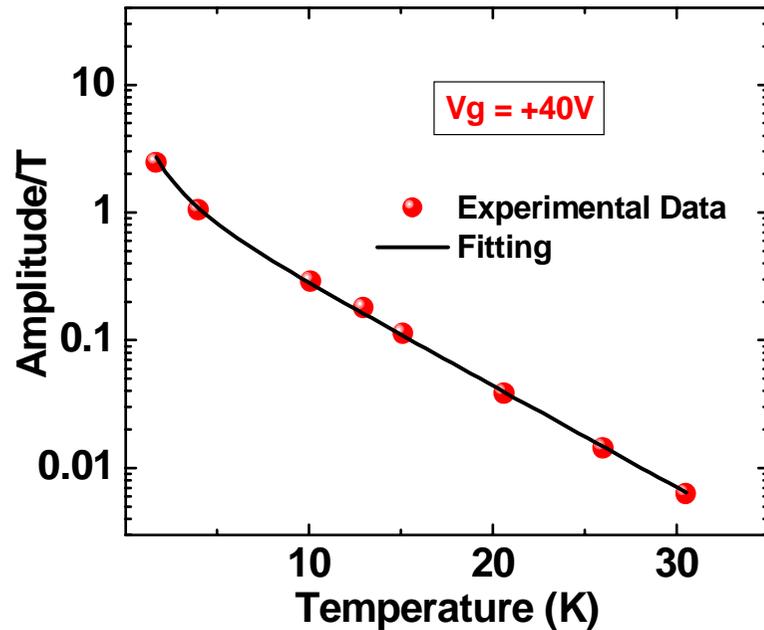
Electrons and Holes can be studied separately

Reduction Factor of SdH
oscillation Amplitude

$$R_T \equiv \frac{2\pi^2 k_B T / \hbar \omega_c}{\sinh(2\pi^2 k_B T / \hbar \omega_c)}$$

$$\omega_c = \frac{eB}{m^*}$$

Effective Mass of Electrons and Holes



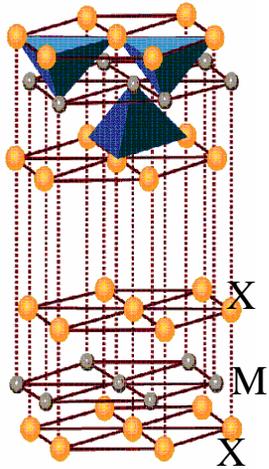
$$\frac{R_T}{T} \equiv \frac{2\pi^2 k_B / \hbar \omega_c}{\sinh(2\pi^2 k_B T / \hbar \omega_c)}$$

$$\omega_c = \frac{eB}{m^*}$$

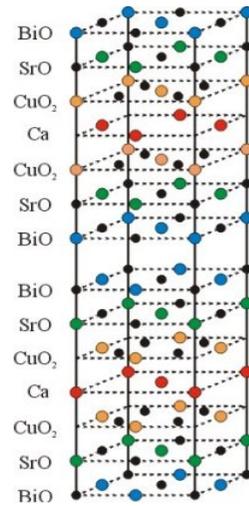
$$m_e^* = 0.05 m_e$$

$$m_h^* = 0.04 m_e$$

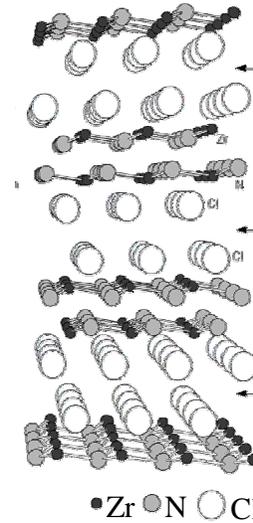
Field Effect in Other Correlated Materials



Metal-Chalcogenide
(M= Nb, Ta, Va, ...
X= S, Se, Te)



$\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8-x}$



$\beta\text{-ZrNCl}$

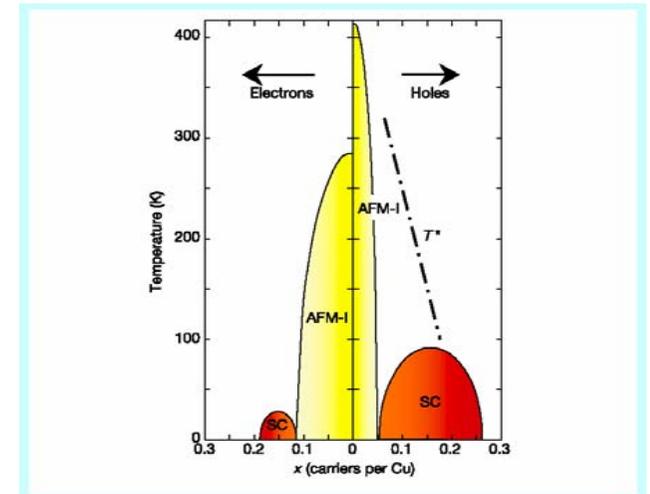
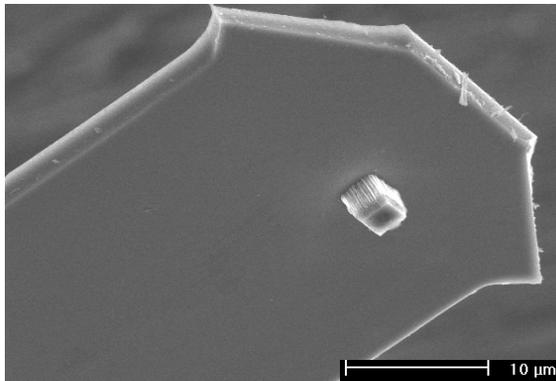
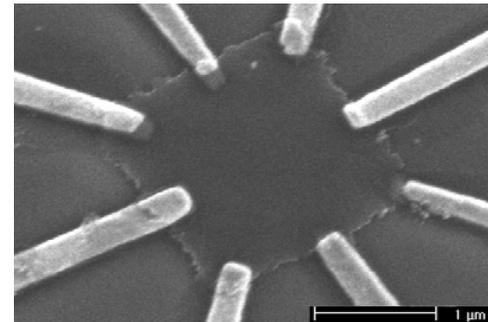
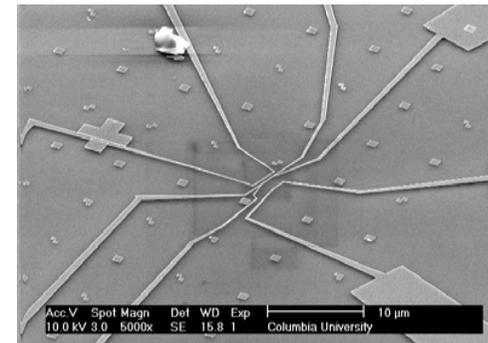
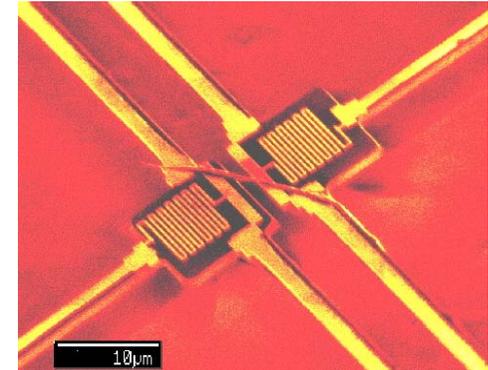
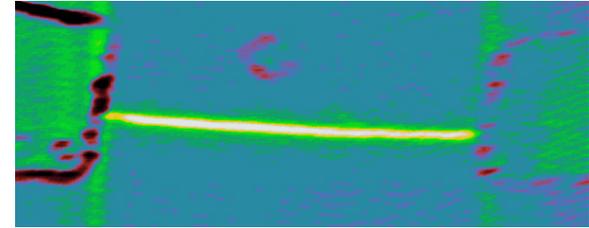


Figure 4 Phase diagram of the high- T_c superconductors. The electron-doped material represents $\text{Nd}_{2-x}\text{Ce}_x\text{CuO}_{4-y}$, and the hole-doped system represents $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$. The insulating, antiferromagnetic phases (AFM-I) are drawn in yellow, and the superconducting phases (SC) are in orange.

Conclusions

- Local temperature measurement energy dissipation in quantum transport limits
- Mesoscopic thermal conductance measurements
- Mesoscopic thermopower measurements
- Magneto-transport measurements in novel 2D-systems



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