Angular Momentum as a Fedosov Quantization on a Sphere

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0. Overview

- Introduction to Moyal *.
- Generalization of Moyal: The Fedosov *.
- An example of the the Fedosov \ast on S^2

1. Introduction

The Moyal *:

 x^a, p_b are classical variables.

$$f(x,p) * g(x,p) \stackrel{def}{=} f \exp\left(\frac{i\hbar}{2}\omega^{AB}\overleftarrow{\partial}_{A}\overrightarrow{\partial}_{B}\right)g$$
$$= fg + \frac{i\hbar}{2}\omega^{AB} \left(\partial_{A}f\right) \left(\partial_{B}g\right) + \cdots$$

where $\partial_A = \frac{\partial}{\partial q^A}$, $q^A = (x^a, p_a)$ and given any function f and g.

 $\omega = \omega_{AB} dq^A dq^B = dp_a \wedge dx^a \quad , \quad A,B = 1,\ldots,2n \quad a,b = 1,\ldots,n$ we note,

$$[x^a, x^b]_* = [p_a, p_b]_* = 0$$
, $[x^a, p_b]_* = i\hbar\delta^a_b$

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Ordinary QM:

An arbitrary operator given by

$$\widehat{A} = \sum_{m,n} A_{a_1 \cdots a_m b_1 \cdots b_n} \widehat{x}^{a_1} \cdots \widehat{x}^{a_m} \widehat{p}^{b_1} \cdots \widehat{p}^{b_n}$$

can be mapped to:

$$A = \sum_{m,n} A_{a_1 \cdots a_m b_1 \cdots b_n} x^{a_1} * \cdots * x^{a_m} * p^{b_1} * \cdots * p^{b_n}$$

via the Weyl Transform (invertible and 1-1)

Moyal * reproduces QM:

$$\hat{H} \leftrightarrow H$$
, $\hat{
ho} \leftrightarrow
ho$ (density matrices), $Tr\left(\hat{A}\right) \leftrightarrow \int \frac{\omega^n}{n!} A$

2. Generalizing Moyal: The Fedosov *

Let M is the phase space associated to N.

Fedosov gives an complicated algorithm to construct * by a perturbative expansion. (I won't go over details –not enough time).

at each point introduce

$$D_A, \omega_{AB} \rightarrow \text{operators } \hat{y}^A \text{ st. } \left[\hat{y}^A, \hat{y}^B \right] = i\hbar\omega^{AB}$$

 \rightarrow some complicated stuff here $\rightarrow \left\lfloor x^{a},x^{b}\right]_{*}$, $[p_{a},p_{b}]_{*}$, $[x^{a},p_{b}]_{*}$

The properties of the Fedosov * are,

- $*: C^{\infty}(M) \times C^{\infty}(M) \to C^{\infty}(M)$ an associative but noncommutative product.
- Invariant under <u>all</u> smooth coordinate transformations.
- In the limit $\hbar \to 0^+ *$ becomes the pointwise multiplication of functions M.
- To first order in \hbar the commutator is the Poisson bracket: $[f,g]_* = i\hbar \{f,g\} + O(\hbar^2)$.

- We could construct the Fedosov * perturbatively given any smooth manifold N.
- When the configuration space $N = E^n$ (\mathbb{R}^n with Euclidean metric δ_{ab}) we get the Moyal *.

*Note: There are some ambiguities in the construction.

3. An Example: **S**²

We calculated the * exactly and got the following commutators,

$$\left[\left[x^a, x^b \right]_* = 0 , \left[x_a, L_b \right]_* = i\hbar\varepsilon_{abc}x^c , \left[L_a, L_b \right]_* = i\hbar\varepsilon_{abc}L^c$$

$$\underline{x} \cdot \underline{x} = 1$$
 , $\underline{x} \cdot \underline{L} = 0$

Most of the ambiguities are related by basis transformations.

$$x^a \to U * x^a * U^{-1}$$

 $L_b \to U * L_b * U^{-1}$

*Note that in general we will have commutators that possibly involve the phase space curvature R and ω .

3. Conclusion

- The Fedosov * seems to be a natural generalization of Moyal.
- We construct an exact solution and write the commutators in the case of S^2 .
- Hopefully we can find more exact solutions for other manifolds.

4. References

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