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Four Dimensional Quantum Hall Effect in Phase Space

(Topological Effects in Spin Orbit Coupling Systems)

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These are preliminary lecture notes, intended only for distribution to participants.

Topological Effects in Spin Orbit Coupling Systems

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Quantum Hall Effect



Quantum Hall Conductance:

$$n = 2 \operatorname{Im} \langle \partial_{\phi} \psi | \partial_{\theta} \psi \rangle$$
$$\sigma_{H} = ne^{2} / h$$

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Outline

- Motivation
- Spin orbital coupling systems
- Topological orbit angular momentum Hall current in Rasha-Dresselhaus S-O coupling system
- 4D Quantum Hall Effect
- Hidden 4D QHE Structure in Luttinger S-O Coupling systems

Motivation

- New topological effects
- Can we realize 4DQHE in solid state systems? --- Emergent physics at low energy
- Connection between 4DQHE and spin orbit coupling systems

Spin Orbit Coupling

spin orbit coupling

$$H = \frac{\hbar}{4m^2c^2} (\vec{\nabla}V \times \vec{P}) \cdot \sigma$$

Solid state physics: band structure

Rasha-Dresselhaus S-O coupling

- Bychkov-Rashba Spin-Orbital Coupling in two dimensional electron gas
 - -- due to confinement in third direction

$$H = \alpha(\vec{P}_y \cdot \sigma_x - \vec{P}_x \cdot \sigma_y)$$



 Dresselhaus Spin Orbital Coupling in two dimensional electron gas

 due to lacking of inversion symmetry

$$H = \beta(\vec{P}_x \cdot \sigma_x - \vec{P}_y \cdot \sigma_y)$$

- 2 dimension
- Spin-1/2

Luttinger Spin-Orbit coupling



3 dimension Spin 3/2

Spin Hall Current



Spin Hall Current in R-D Systems



$$H = \frac{p^2}{2m} + \alpha (p_x \sigma_y - p_y \sigma_x) + \beta (p_x \sigma_x - p_y \sigma_y)$$

$$\begin{aligned}
-\frac{e}{8\pi} & \alpha > \beta \\
\sigma_x^z &= 0 & \alpha = \beta \\
-\frac{e}{8\pi} & \alpha < \beta
\end{aligned}$$

Kubo Formula

$$< J_i^j >= \sum \operatorname{Im} \frac{< P(t), \lambda \mid J_i^j \mid P(t), \lambda' >< P(t), \lambda' \mid \frac{\partial}{\partial t} \mid P(t), \lambda >}{E_{\lambda'}(p(t)) - E_{\lambda}(p(t))}$$

Orbit Angular Momentum Hall Current

 $\beta=0:$ [Lz+Sz, H]=0

Definition:

 $\alpha = 0$: [Lz-Sz, H] = 0

Gauge in Momentum Space

Unitary transformation:

$$H' = U(\vec{k})HU^{+}(\vec{k}) = \frac{p^{2}}{2m} + \lambda(p_{x}, p_{y})\sigma_{z} + U(\vec{k})V(\vec{x})U^{+}(\vec{k})$$
$$U(\vec{k})V(i\partial_{k})U^{+}(\vec{k}) = V(\vec{D}) \qquad D_{i} = i\frac{\partial}{\partial k_{i}} - A_{i} \qquad A_{i} = -iU(\vec{k})\frac{\partial}{\partial k_{i}}U^{+}(\vec{k})$$

• Projection to one band in adiabatic limit
$$H_{z}^{eff} = \frac{p^{2}}{2m_{\lambda}} \pm \lambda(p) + V_{\pm}(\vec{x}) \qquad X_{i} \equiv D_{i} = i\frac{\partial}{\partial p_{i}} - \tilde{A}_{i}(\vec{p})$$
$$Nontrivial gauge after Projection$$

Gauge field: an angular momentum flux

• Gauge field

$$\vec{A} = < L_z > (\frac{p_y}{p^2}, -\frac{p_x}{p^2})$$

$$\oint \vec{A} \cdot d\vec{p} \neq 0$$

$$F \equiv [x, y] = 0$$

$$\int F dp_x dp_y \neq 0$$



Topological Orbit Hall Current

Derivation of orbit hall current:

$$\int \langle Lz \rangle \dot{y}dp_{x}dp_{y} = -ieE \int dp_{x}dp_{y} \langle Lz \rangle [y, H]$$
$$= eE \int dp_{x}dp_{y} \langle Lz \rangle [y, x]$$

Orbit Hall conductance:

$$\theta_x^z = e \int dp_x dp_y < Lz > [y, x] = \frac{e}{8\pi} \frac{\alpha^2 + \beta^2}{|\alpha^2 - \beta^2|}$$

Spin Hall conductance from this derivation is zero! No topological Spin Hall conductance!

Conclusion

- Spin Hall effect in R-D case is not topological effect, However, orbit Hall current is purely topological effect
- Orbit Hall current is larger than spin Hall current in general in R-D case
- Orbit Hall current does not change sign in all coupling parameter region.
- When Rashba coupling is larger than Dresselhaus coupling, spin Hall current and orbit Hall current run in opposite direction. The magnetization direction near edge due to the total angular momentum Hall current is opposite to the prediction from the spin Hall current
- Due to the topological nature, the orbit Hall current is expected to survive in the weak disorder limit.

4DQHE



"Lowest Landau levels" on S⁴

 $l_0^2 = \lim_{R \to \infty} \frac{R^2}{P}$

- Ground degeneracy: $D(p) = \frac{1}{6}(p+1)(p+2)(p+3)$
- Thermodynamic limit

• The problem is effectively six dimensional, given by orbital space
$$S^4$$
 times the isospin space S^2 .

• Filling fraction is given by $D(p) \propto p^3 \propto R^6 \Rightarrow N \propto D(p) \propto R^6$

v = N / D(p)

• Magnetic length:

QH:

$$l_0^2 = \frac{\hbar c}{eB} = \frac{R^2}{S}$$

CP(3)

 S^2

Noncommutative Geometry

- physics at the lowest laudau level

• Noncommutative Geometry in 4D:

$$[X_{\mu}, X_{\nu}] = i l_{0}^{2} \eta_{\mu\nu}^{i} n^{i}$$

Magnetic translations commute up to a SU(2) factor. A noncommutative geometric structure which treats all space coordinates symmetrically

4D Quantum Hall Conductance

• Effective Hamiltonian at LLL: (No kinetic energy)

$$H = V(X)$$

• Quantum Hall Conductance:

$$\dot{X}_{\mu} = l_0^2 \eta_{\mu\nu}^{i} n^i \frac{\partial V(X)}{\partial X_{\nu}}$$

Requirements

- A six dimensional space
- Low energy emergent physics
- Luttinger spin orbit coupling
 - A natural choice: Phase space which has six dimensions
 - Spin Hall effect essentially comes from the physics of 4DQHE

Gauge structure in momentum space

• Unitary transformation:

Dynamic equation

$$\begin{array}{c} \stackrel{\cdot}{X}_{i} = \frac{k_{i}}{m_{\lambda}} - [X_{i}, X_{j}] \frac{\partial V}{\partial X_{j}} \longrightarrow J_{i}^{j}(\lambda) = iTr \{S_{i\lambda} [X_{j\lambda}, X_{k\lambda}]\} eE_{k} \\ \\ \hline \\ Topological term \end{array}$$

Effective Hamiltonian

$$H^{\rm eff} = \frac{k^2}{2m_{\lambda}} + V(\vec{x})$$

$$X_i \equiv D_i = i \frac{\partial}{\partial k_i} - \widetilde{A}_i(\vec{k})$$

•The gauge potential rises from the conventional berry phase in adiabatic processes.



(Dirac monopole)

The key mathematical structure is the noncommutative geometry:

$$[x_i, x_j] = -iF_{ij} \qquad F_{ij} = \epsilon$$

$$F_{ij} = \varepsilon_{ijk} \lambda \frac{k_k}{k^3}$$

Noncommutative Structure on Phase Space

• Define: $X_{\mu} = (x, y, z)$

$$X_{\mu} = (x, y, z, \frac{\lambda}{k}), X_{4} = \frac{\lambda}{k}$$
$$\vec{n} = \frac{\vec{k}}{k}$$

• Noncommutative geometry :

$$[X_{\mu}, X_{\nu}] = i l_{0}^{2} \eta_{\mu\nu}^{i} n^{i}$$



Effective Hamiltonian For Luttinger Coupling

• Effective Hamiltonian is simply a pure potential in 4d space (no ``momentum")

$$H^{\text{eff}} = \frac{1}{8m_{\lambda}X_{4}^{2}} + V(X_{1}, X_{2}, X_{3})$$

With

$$[X_{\mu}, X_{\nu}] = i l_0^2 \eta_{\mu\nu}^{i} n^{i}$$

An exact formulation of 4DQHE at the LLL!!!

Structure of the 4D liquid



What happens at m = infinite

- m= infinite: band is completely flat. There is no potential, i.e. H=0. $\gamma_1 = -2\gamma_2$
- Based on 4DQHE, a new liquid state (non Fermi liquid) could be formed.
- Since electrons are like membranes in the liquid, excitations with factional statistics could emerge.