

Spatial Bose-Einstein Condensation on Inhomogeneous Networks

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Outlook

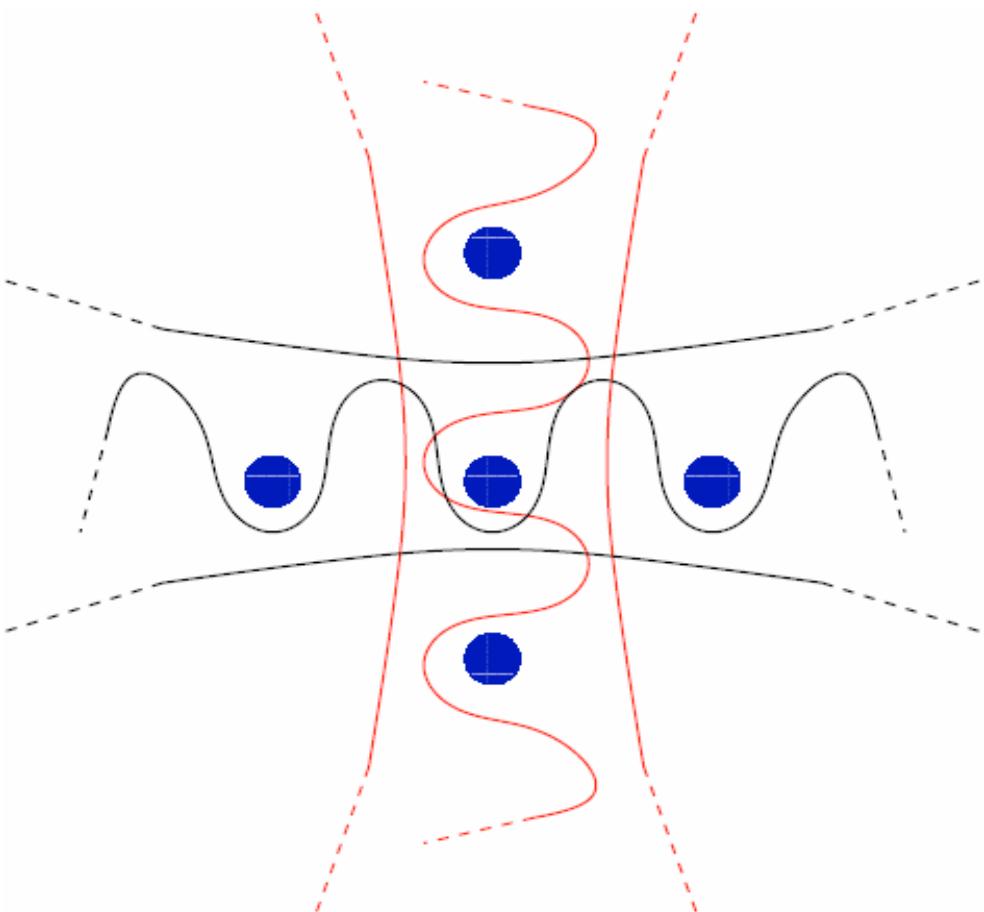
1) Experimental systems:

- Ultracold bosons on **inhomogeneous** optical lattices
 - Superconducting Josephson junctions on **inhomogeneous** insulating supports

2) Thermodynamical properties:

- Macroscopic occupation of the localized ground-state (emergence of spatial Bose-Einstein condensation)

Creating a star network with ultracold bosons



Temperature $\sim 0\text{-}500 \text{ nK}$
Number of particles $\sim 1000\text{-}10000$
Number of wells ~ 100

$$V(x) \approx V_0 \cos^2(kx)$$

$$k = 2\pi/\lambda \quad \lambda \approx 800 \text{ nm}$$

$$E_R = \hbar^2 k^2 / 2m$$

$$V_0 = s \cdot E_R \quad s \approx 10 - 30$$

Creating a comb-shaped network with superconducting Josephson junctions

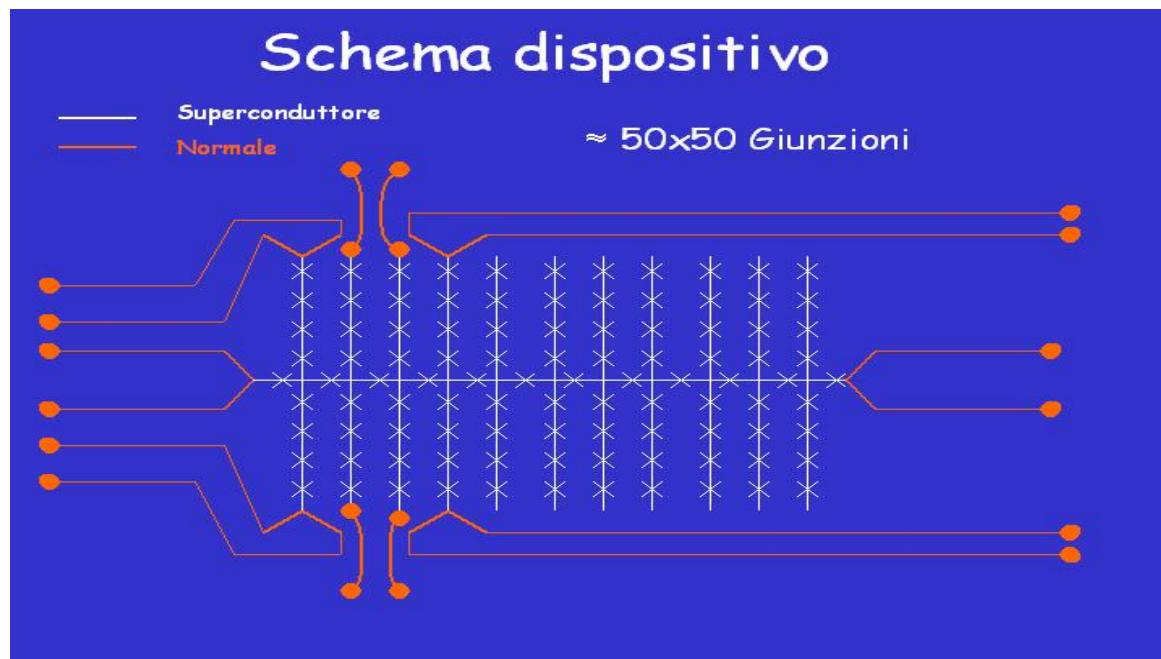


Figure of the experimental setup
(courtesy of M. Cirillo *et al.*)

Critical Josephson energy $E_J = \phi_0 I_C$

Non-interacting bosons on a graph

$$\hat{H} = -t \sum_{i,j} A_{ij} \hat{\mathbf{a}}_i^\dagger \hat{\mathbf{a}}_j$$
$$[\hat{\mathbf{a}}_i, \hat{\mathbf{a}}_j^\dagger] = \delta_{ij} \quad \hat{n}_j = \hat{\mathbf{a}}_j^\dagger \hat{\mathbf{a}}_j$$

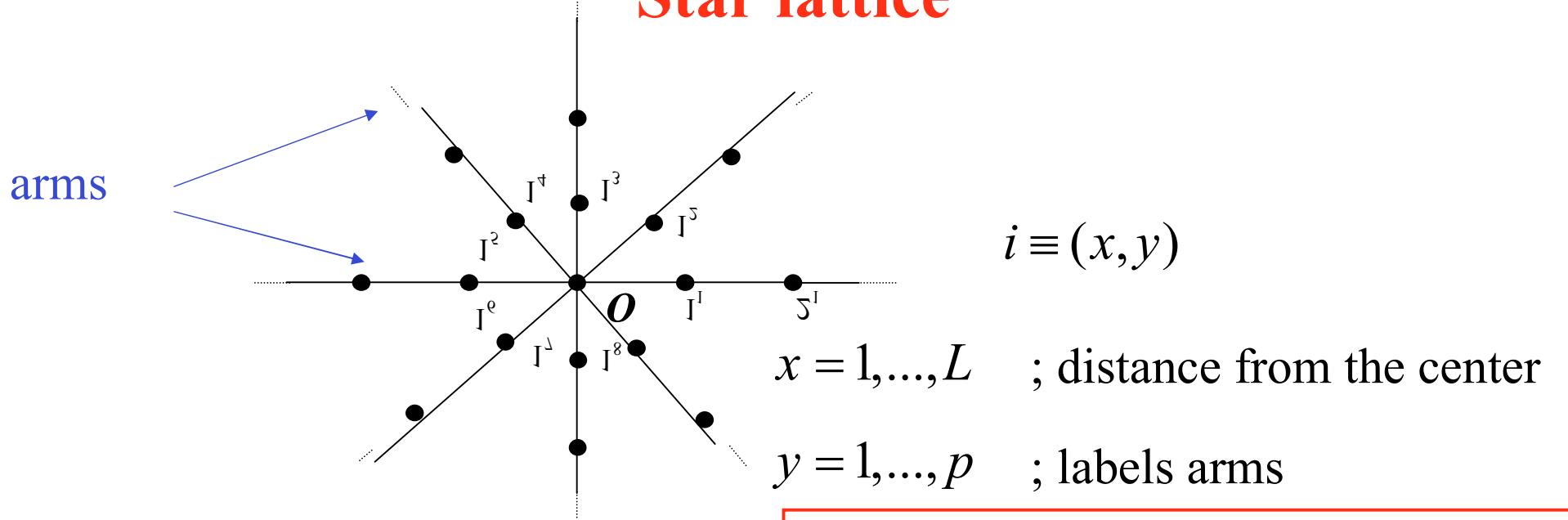
In the following: bosons on discrete structures which are not necessarily regular lattices:

Graphs: Adjacency Matrix $A_{ij} = \begin{cases} 1 & \text{if } i - j \text{ is a link} \\ 0 & \text{otherwise} \end{cases}$

B.E.C. of non interacting bosons on inhomogeneous low-dimensional structures: $d < 2$.

Network geometry naturally induces topological (quantum) order.

Star lattice



Total number of sites: $N_S = pL + 1$

$$-t \sum_j A_{ij} \psi_\nu(j) = E_\nu \psi_\nu(i)$$

z_i coordination number of a given site: **2**

z_O coordination number of the center: **p**

Eigenvalue equation

Spectrum σ formed by N_S states and divided in 3 parts: $\{E_0, \sigma_0, E_+\}$

σ_0 \longrightarrow $pL-1$ delocalized states with $E \in [-2t, 2t]$

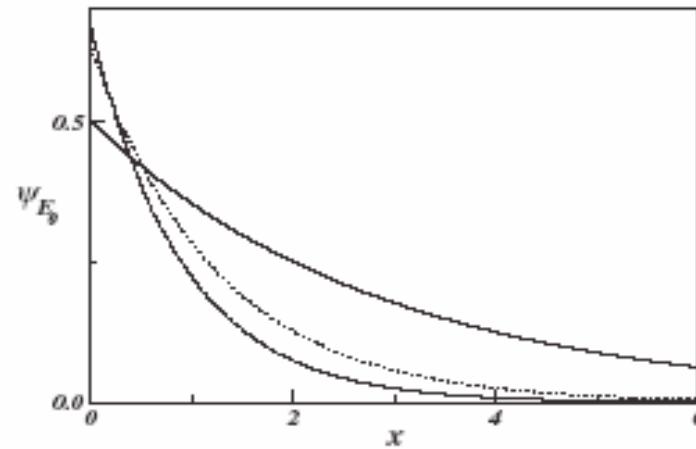
two bound states

$E_0 < -2t$
and
 $E_+ > 2t$

E_0 : localized ground-state

$$\psi_{E_0} = \sqrt{\frac{p-2}{2p-2}} e^{-x/\xi_0}$$

$$\xi_0 = \frac{2}{\log(p-1)}, \quad p > 2$$



$$E_0 = -t \frac{p}{\sqrt{p-1}}$$

**Adding arms
enhances
localization**

$$p=2 \longrightarrow$$

$$E_0 = -2t$$

Linear chain

$$p \neq 2 \longrightarrow$$

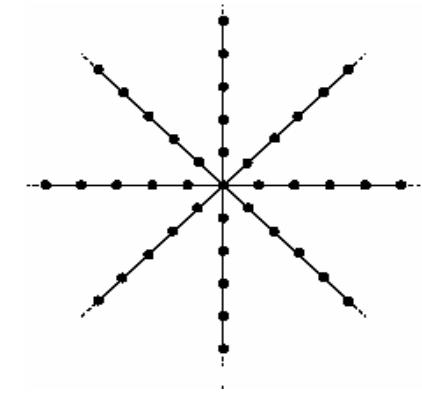
$$\Delta(p) = |E_0| - 2t$$

Gapped spectrum

Thermodynamics for bosons hopping on a star lattice

$$N_T = N_{E_0} + \int_{E \in \sigma_0} \frac{N_S \rho(E)}{z^{-1} e^{\beta(E-E_0)} - 1} + N_{E_+}$$

Ground-state *delocalized states* *excited state*



$$T < T_C \rightarrow \frac{N_{E_0}}{N_S} \neq 0$$

$$T_C \approx \frac{p-2}{2\sqrt{p-1}} \frac{E_J}{k_B}$$

Topology effect

$$f = \frac{N_T}{N_S}$$

$$E_J = 2tf$$

$$\frac{N_{E_0} \left(\frac{T}{T_C} \right)}{N_T} \approx 1 - \frac{T}{T_C}$$

I. Brunelli et al., *J. Phys. B* **37**, S275 (2004)

Conclusions and perspectives

- Topology role in inducing new phases
for bosons on graphs

- The experimental signature of the
**Bose-Einstein condensation is given by the
inhomogeneous distribution of bosons in the
graph below the critical temperature**

- Dream:
**To induce desired macroscopic
coherent behaviors by acting
on the geometry and topology of networks**

Rather new area:

R. Burioni et al., *Europh. Lett.* **52**, 251 (2000)

L. B. Ioffe et al., *Phys. Rev. Lett.* **90**, 107003 (2003)