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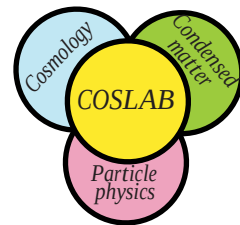
SMR 1646 - 1

**Conference on
Higher Dimensional Quantum Hall Effect, Chern-Simons Theory and
Non-Commutative Geometry in Condensed Matter Physics and Field Theory
1 - 4 March 2005**

***MOMENTUM-SPACE TOPOLOGY OF FERMION ZERO MODES,
CHERN-SIMONS TERM & QUANTUM PHASE TRANSITION***

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These are preliminary lecture notes, intended only for distribution to participants.



Momentum-space topology of fermion zero modes, Chern-Simons term & quantum phase transition

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Trieste, March 2005



* Momentum-space topology in 3+1

Fermi points, chiral fermions, chiral anomaly, Standard Model, Chern-Simons term, spinons & holons



* Momentum-space topology in 2+1

Chern-Simons term, quantum statistics of skyrmions, QHE and spin QHE



* Momentum-space topology of edge states

index theorem



* Momentum-space topology in higher dimensions

Brane fermions



* Quantum phase transitions dictated by momentum-space topology

BEC-BCS and neutrino oscillations

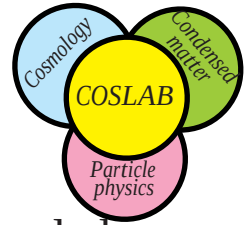


* Conclusion





Example of p -space topology



Compare these two Hamiltonians

Bogoliubov-Nambu
for quasiparticles
in $^3\text{He-A}$ & chiral superconductors

$$H = \begin{pmatrix} \frac{p^2 - p_F^2}{2m} & c_{\perp}(p_x + ip_y) \\ c_{\perp}(p_x - ip_y) & -\frac{p^2 - p_F^2}{2m} \end{pmatrix}$$

Weyl for **right**-handed
neutrino

$$H = + c \boldsymbol{\sigma} \cdot \mathbf{p}$$

$$H = \begin{pmatrix} cp_z & c(p_x + ip_y) \\ c(p_x - ip_y) & -cp_z \end{pmatrix}$$

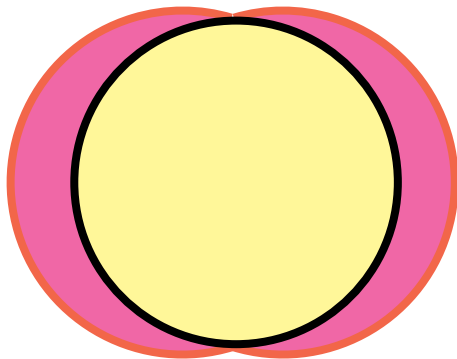
What is common for them?

$$H(\mathbf{p}) = \boldsymbol{\tau} \cdot \mathbf{g}(\mathbf{p})$$

$$E^2(\mathbf{p}) = \mathbf{g}^2(\mathbf{p})$$

1. $E(\mathbf{p}) = 0$ at points (Fermi points)

$$\mathbf{p} = +p_F \mathbf{e}_z \quad N = +1$$



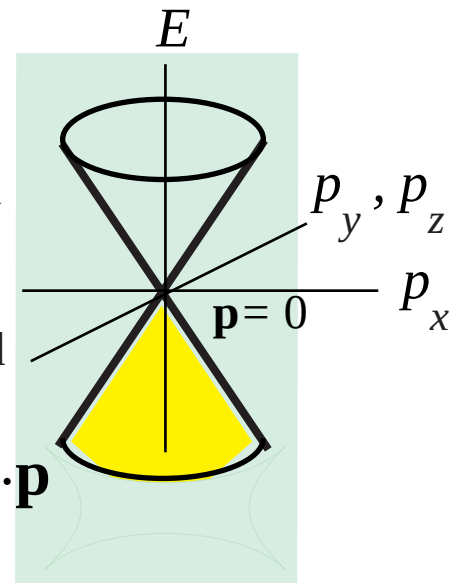
$$\mathbf{p} = -p_F \mathbf{e}_z \quad N = -1$$

$N = +1$
right-handed
neutrino

left-handed
neutrino

$$H = - c \boldsymbol{\sigma} \cdot \mathbf{p}$$

$$N = -1$$



2. Fermi points are topologically stable
& described by topological invariant in momentum space

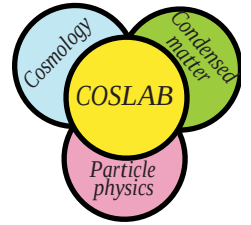
$$N = \frac{1}{8\pi} \epsilon_{ijk} \int_{\text{over 2D surface around Fermi point}} dS^i \mathbf{g} \cdot (\partial^j \mathbf{g} \times \partial^k \mathbf{g})$$

3. Close to Fermi points (quasi)particles are
relativistic **left** or **right**-handed **chiral Weyl fermions**

$$L = e^\mu_a \sigma^a (p_\mu - e A_\mu)$$



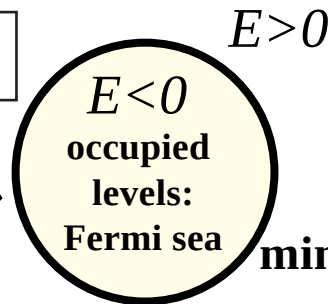
Universality classes of fermionic vacua from p -space topology



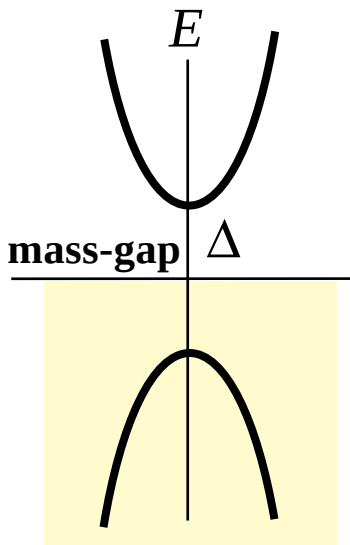
Systems with Fermi surface

$$E(\mathbf{p}) = \frac{p^2 - p_F^2}{2m}$$

Fermi surface $E=0$



$\text{cod} = 1$
co-dimension
 $\text{cod} =$
space dimension D
minus dimension of zeroes;
exist for $D \geq 1$



Fully gapped Fermi systems

no zeroes

Superconductor

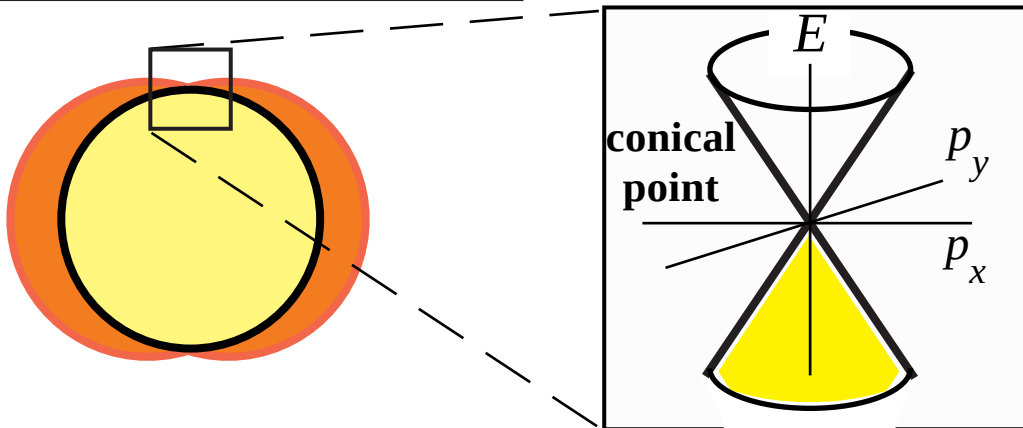
$$E^2(\mathbf{p}) = v_F^2 (\mathbf{p} - \mathbf{p}_F)^2 + \Delta^2$$

Dirac fermions

$$E^2 = p^2 c^2 + M^2$$

p -space topology
is non-trivial
in even
space dimensions
 $D=2,4,6,\dots$

Systems with Fermi points



$$E^2 = p^2 c^2$$

$$H = c \boldsymbol{\sigma} \cdot \mathbf{p}$$

$^3\text{He-A}$
&
Standard Model
fermions

$\text{cod} = 3$
 $D \geq 3$

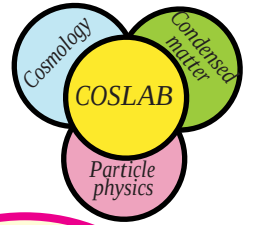
Higher dimensional Fermi points

odd co-dimensions
 $\text{cod} = 5, 7, \dots$

$D \geq 5, 7, \dots$



Fermi surface as vortex in 3+1 p -space



* Fermi gas

$$E = \frac{p^2}{2m} - \mu$$

$$E(\mathbf{p}) = \frac{p^2 - p_F^2}{2m}$$

Fermi surface
 $E=0$

$E>0$

$E<0$

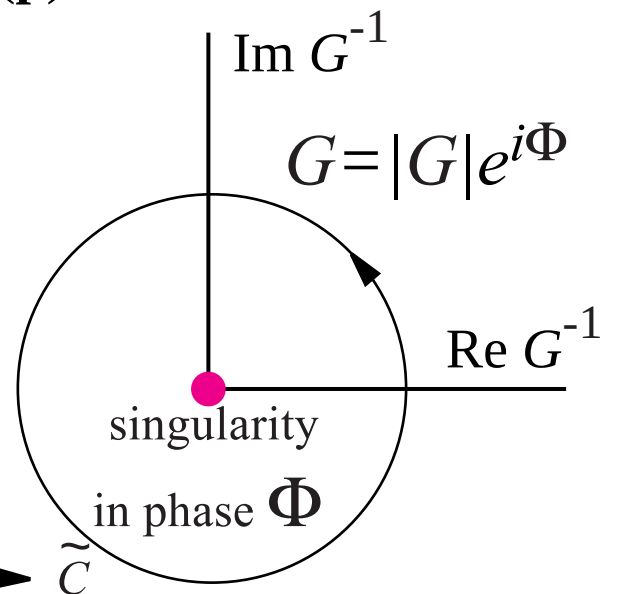
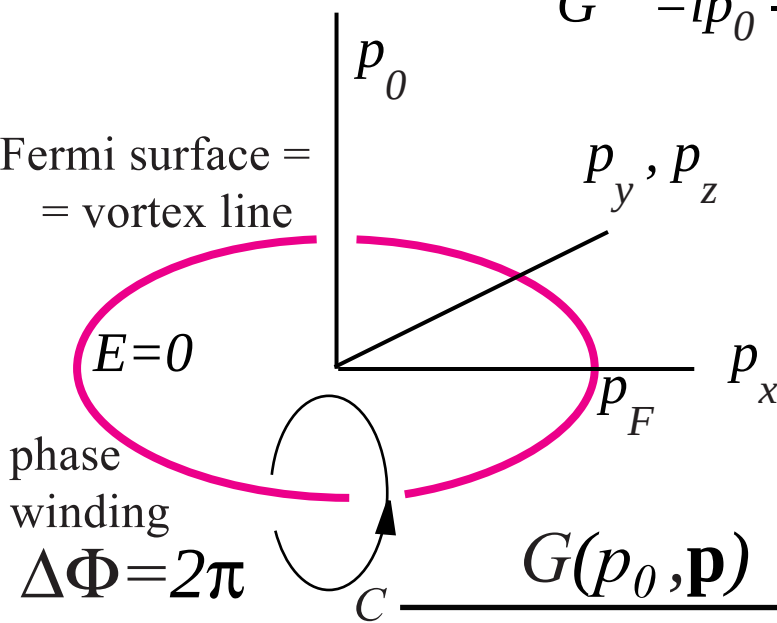
occupied
levels:
Fermi sea

$p=p_F$

Fermi surface is robust to perturbations and interactions
as vortex -- topologically stable singularity of Green function:

$$G^{-1} = ip_0 - E(\mathbf{p})$$

Fermi surface =
= vortex line



* Interacting Fermi system:
general topological invariant

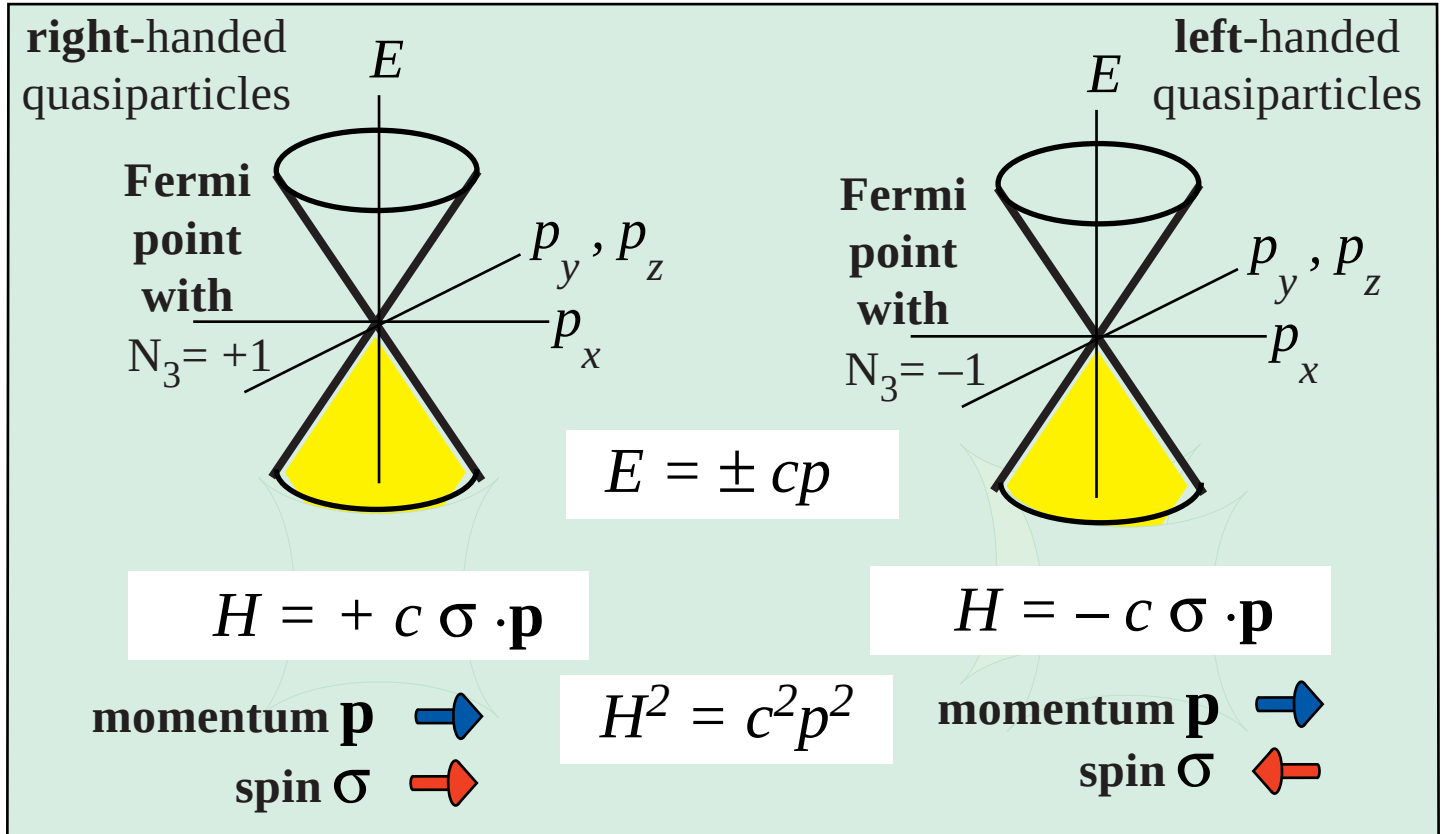
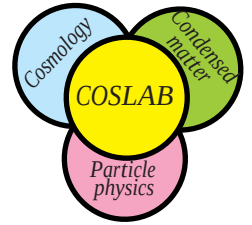
$$N_1 = \frac{1}{2\pi i} \text{tr} \oint_{\text{around Fermi surface}} dp^\mu \mathbf{G} \partial_{p^\mu} \mathbf{G}^{-1}$$



Chiral particles

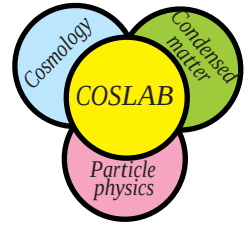
Quasiparticles near **Fermi points** are
relativistic:

left or right-handed chiral Weyl fermions

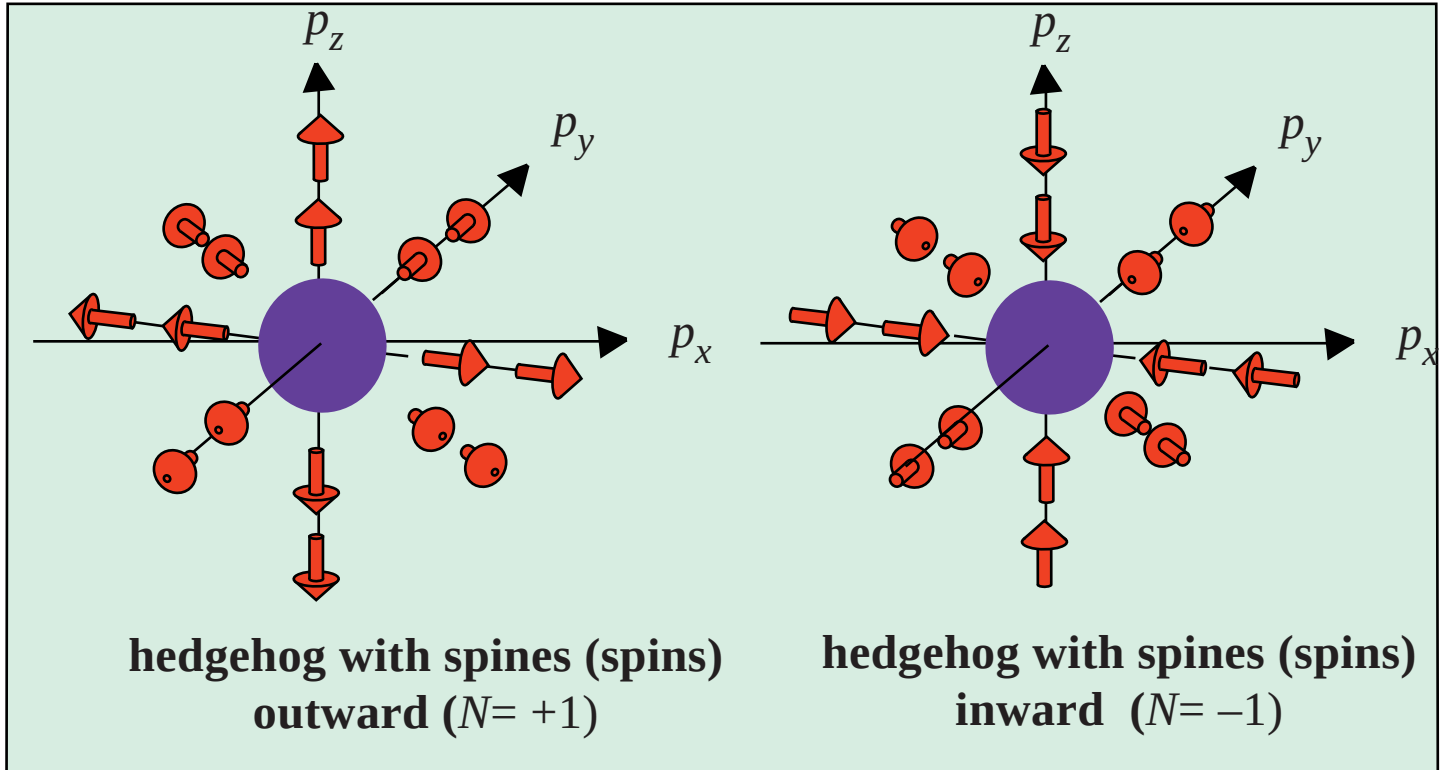




Chiral particles



Topological stability of Fermi point: hedgehog in momentum space



$$H = + c \boldsymbol{\sigma} \cdot \mathbf{p}$$

$$\mathbf{g}(\mathbf{p}) = +c\mathbf{p}$$

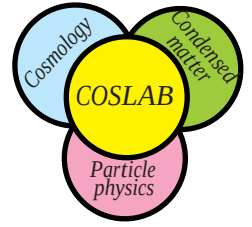
$$H = - c \boldsymbol{\sigma} \cdot \mathbf{p}$$

$$\mathbf{g}(\mathbf{p}) = -c\mathbf{p}$$

$$N = \frac{1}{8\pi} \epsilon_{ijk} \int_{\text{over 2D surface around Fermi point}} dS^i \hat{\mathbf{g}} \cdot (\partial^j \hat{\mathbf{g}} \times \partial^k \hat{\mathbf{g}})$$



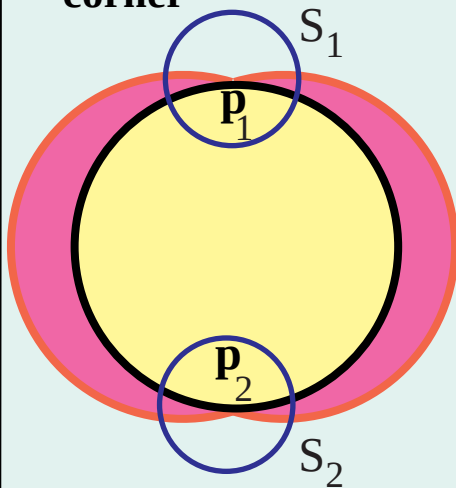
Topological stability of Fermi point (general case)



Topological invariant in 4D momentum space $(\mathbf{p}, p_0)$
in terms of fermionic propagator:
matrix Green's function $\mathbf{G}(\mathbf{p}, p_0)$

$$N = \frac{1}{24\pi^2} \epsilon_{\mu\nu\lambda\gamma} \text{tr} \int_{\text{over 3D surface } S \text{ in 4D momentum space}} dS^\gamma \mathbf{G} \partial^\mu \mathbf{G}^{-1} \mathbf{G} \partial^\nu \mathbf{G}^{-1} \mathbf{G} \partial^\lambda \mathbf{G}^{-1}$$

in
low-energy
corner $\mathbf{G}^{-1} = i p_0 + c \boldsymbol{\sigma} \cdot (\mathbf{p} - \mathbf{p}_1)$



$$N = +1$$

right-handed
particles

$$N = -1$$

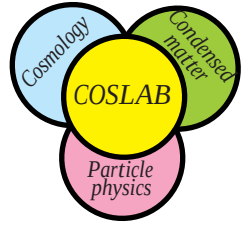
left-handed
particles

$$\mathbf{G}^{-1} = i p_0 - c \boldsymbol{\sigma} \cdot (\mathbf{p} - \mathbf{p}_2)$$

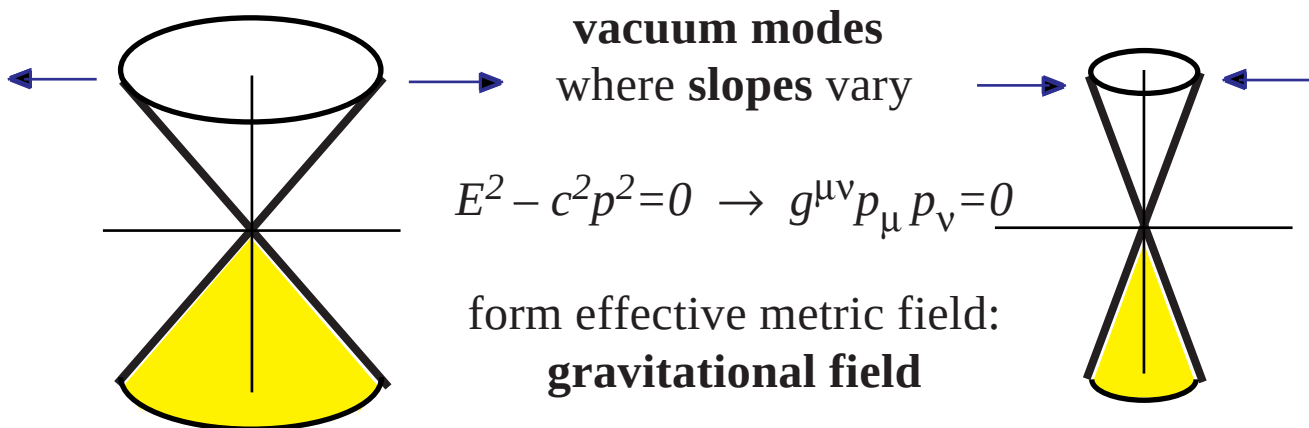
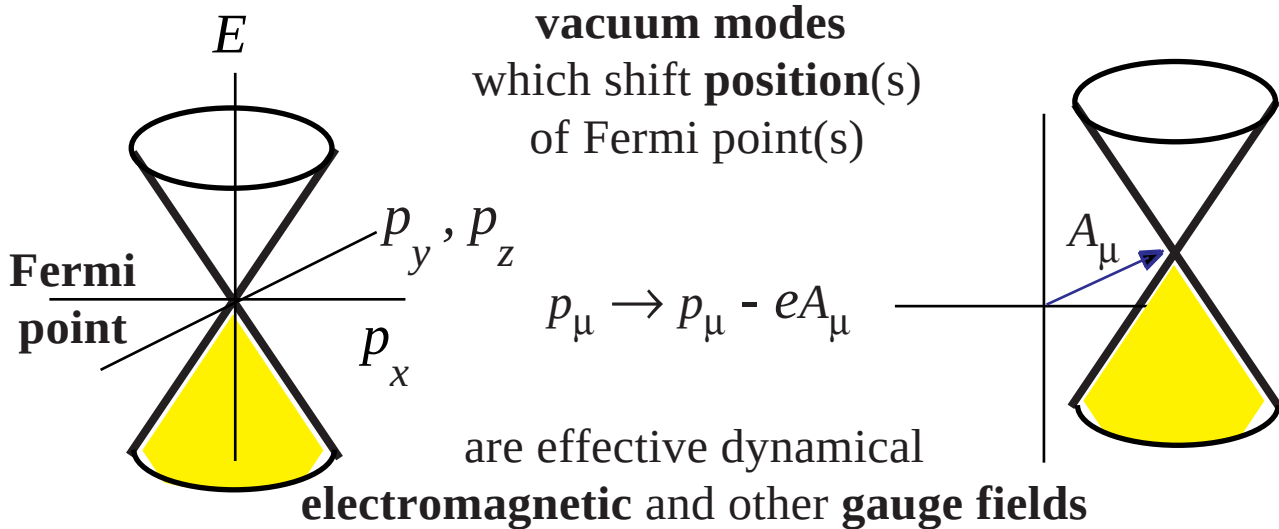
top. invariant
determines
chirality
in low-energy
corner



Collective modes of fermionic vacua of Fermi-point Universality Class: gauge fields & gravity



Vacuum low-energy dynamics cannot destroy the Fermi point.
Shifts $\mathbf{A}$ and slopes g^{ik} are propagating collective modes:



Quasiparticle near **Fermi point** is
left or **right** particle moving in effective
gravitational, electromagnetic, weak fields

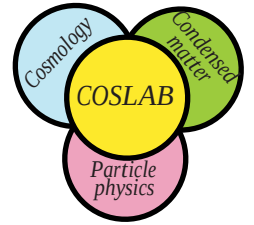
$$g^{\mu\nu} (p_\mu - eA_\mu - e\tau \cdot \mathbf{W}_\mu) (p_\nu - eA_\nu - e\tau \cdot \mathbf{W}_\nu) = 0$$

chiral fermions,
gauge fields and gravity
appear
in low-energy corner
together with spin and
physical laws:
Lorentz and gauge
invariance,
and general covariance

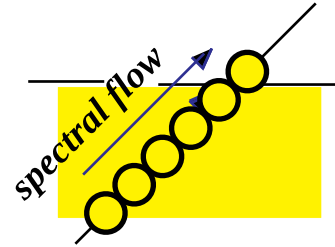


Chiral anomaly

nucleation of fermionic charge from *vacuum*



* **chiral particles: *quarks & leptons in Standard Model***
& *quasiparticles in $^3\text{He-A}$*
are created from vacuum one by one
by spectral flow



creation of momentum
from
 $^3\text{He-A}$ vacuum

creation of baryonic charge
from
Standard Model vacuum

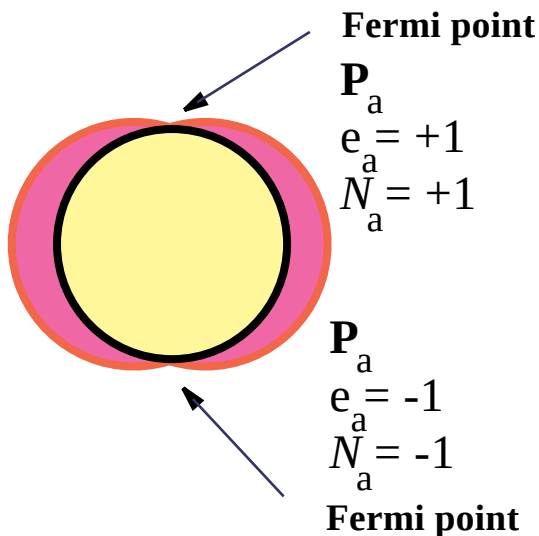
$$\dot{\mathbf{P}} = (1/4\pi^2) \mathbf{B} \cdot \mathbf{E} \sum_a \mathbf{P}_a N_a e_a^2$$

$$\dot{\mathbf{B}} = (1/4\pi^2) \mathbf{B}_Y \cdot \mathbf{E}_Y \sum_a B_a N_a Y_a^2$$

$\mathbf{P}_a$ -- momentum of Fermi point
 e_a -- effective electric charge

B_a -- baryonic charge
 Y_a -- hypercharge

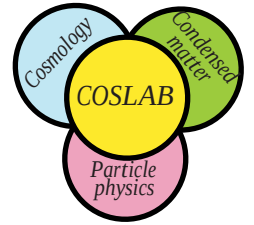
N_a -- topological charge of Fermi point





Chiral anomaly

General anomaly equation
in terms of topological invariant
protected by symmetry :



nucleation of fermionic charge Q
by gauge field A^Y

$$\dot{Q} = (1/4\pi^2) \mathbf{B}_Y \cdot \mathbf{E}_Y N_{YYQ}$$

$$N_{YYQ} = \frac{1}{24\pi^2} \epsilon_{\mu\nu\lambda\gamma} \text{tr } YYQ \int dS^\gamma G \partial^\mu G^{-1} G \partial^\nu G^{-1} G \partial^\lambda G^{-1}$$

Y - charge interacting with gauge field A_μ^Y

$^3\text{He-A}$

Standard Model

$Q = \mathbf{P}_a$ -- momentum of Fermi point

$Q = B_a$ -- baryonic charge

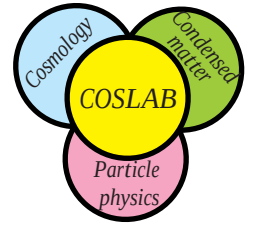
$Y = e_a$ -- effective electric charge

$Y = Y_a$ -- hypercharge



Chiral fermions in Standard Model

Fermi-point Universality Class



Family #1 of quarks and leptons

left particles

right particles

$$SU(3)_C$$

$\begin{matrix} +2/3 \\ \mathbf{u}_L \\ +1/6 \end{matrix}$	$\begin{matrix} -1/3 \\ \mathbf{d}_L \\ +1/6 \end{matrix}$
$\begin{matrix} +2/3 \\ \mathbf{u}_L \\ +1/6 \end{matrix}$	$\begin{matrix} -1/3 \\ \mathbf{d}_L \\ +1/6 \end{matrix}$
$\begin{matrix} +2/3 \\ \mathbf{u}_L \\ +1/6 \end{matrix}$	$\begin{matrix} -1/3 \\ \mathbf{d}_L \\ +1/6 \end{matrix}$

$$SU(3)_C$$

$\begin{matrix} +2/3 \\ \mathbf{u}_R \\ +2/3 \end{matrix}$	$\begin{matrix} -1/3 \\ \mathbf{d}_R \\ -1/3 \end{matrix}$
$\begin{matrix} +2/3 \\ \mathbf{u}_R \\ +2/3 \end{matrix}$	$\begin{matrix} -1/3 \\ \mathbf{d}_R \\ -1/3 \end{matrix}$
$\begin{matrix} +2/3 \\ \mathbf{u}_R \\ +2/3 \end{matrix}$	$\begin{matrix} -1/3 \\ \mathbf{d}_R \\ -1/3 \end{matrix}$

quarks

$SU(2)_L$

$\begin{matrix} 0 \\ \mathbf{v}_L \\ -1/2 \end{matrix}$	$\begin{matrix} -1 \\ \mathbf{e}_L \\ -1/2 \end{matrix}$
---	--

$\begin{matrix} 0 \\ \mathbf{v}_R \\ 0 \end{matrix}$	$\begin{matrix} -1 \\ \mathbf{e}_R \\ -1 \end{matrix}$
--	--

leptons

$$H = -c \boldsymbol{\sigma} \cdot \mathbf{p}$$

$$N = -1$$

$$H = +c \boldsymbol{\sigma} \cdot \mathbf{p}$$

$$N = +1$$

$$N = \frac{1}{24\pi^2} e_{\mu\nu\lambda\gamma} \text{tr} \int_{\text{over 3D surface S in 4D momentum space}} dS^\gamma \mathbf{G} \partial^\mu \mathbf{G}^{-1} \mathbf{G} \partial^\nu \mathbf{G}^{-1} \mathbf{G} \partial^\lambda \mathbf{G}^{-1}$$

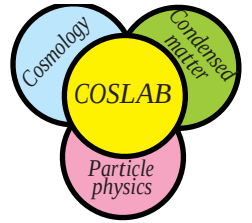
electric charge

Particle_{Left}

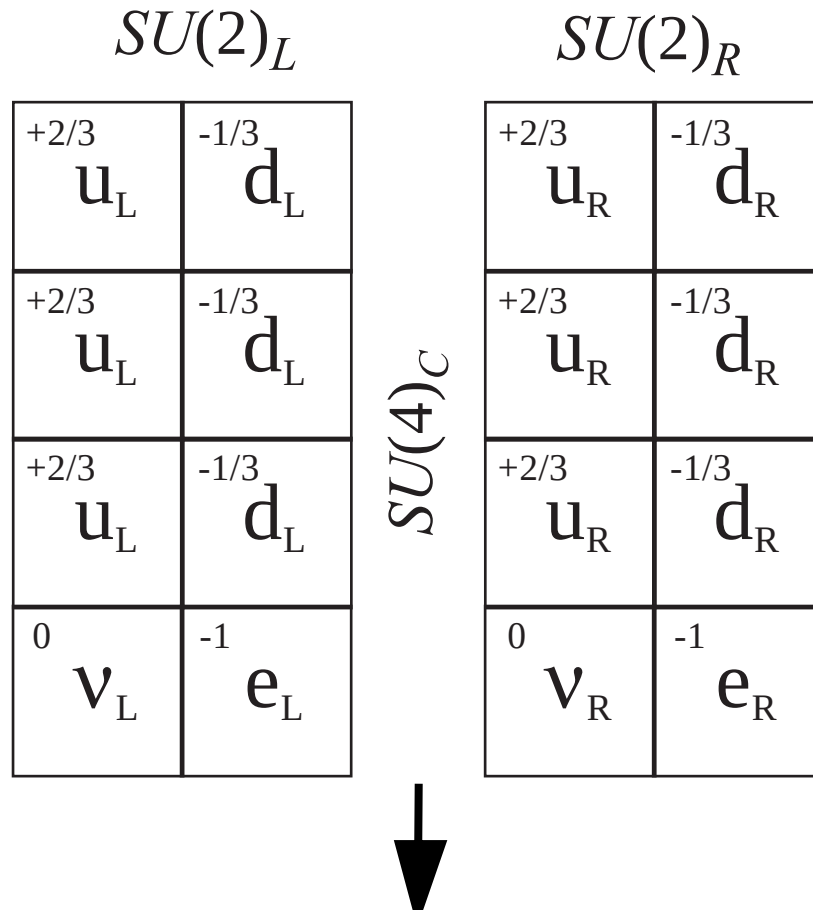
hypercharge



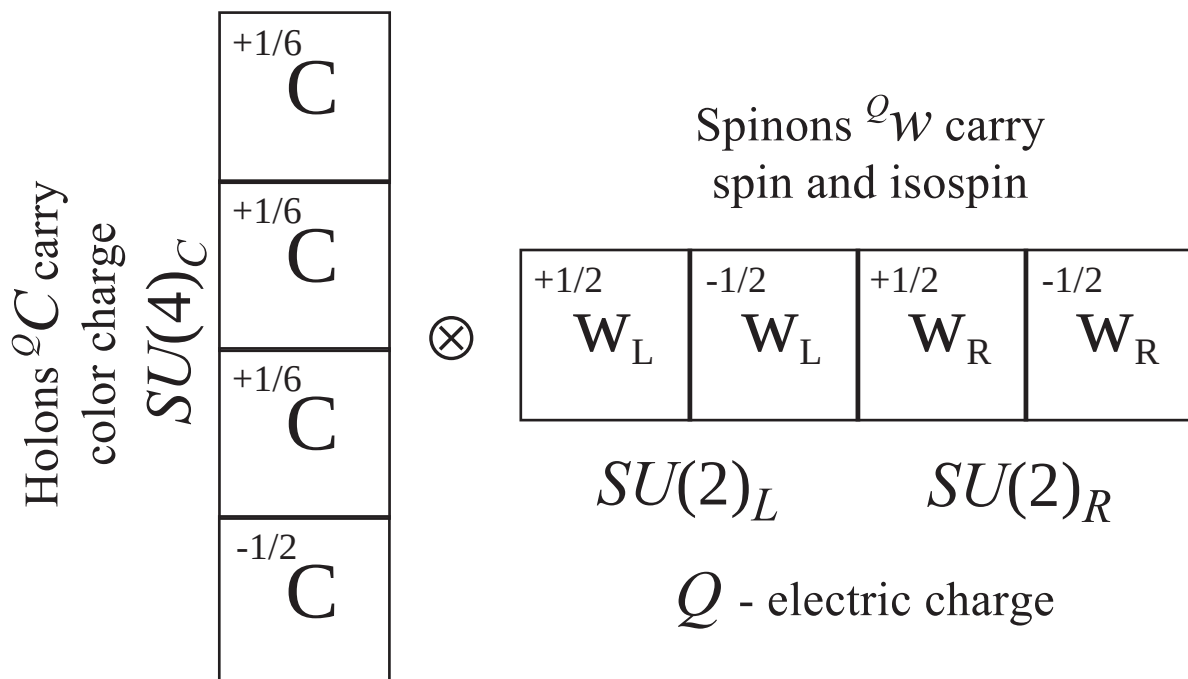
Pati-Salam unification of quarks & leptons & Terazawa spinon-holon model of fermions



lepton as 4-th color



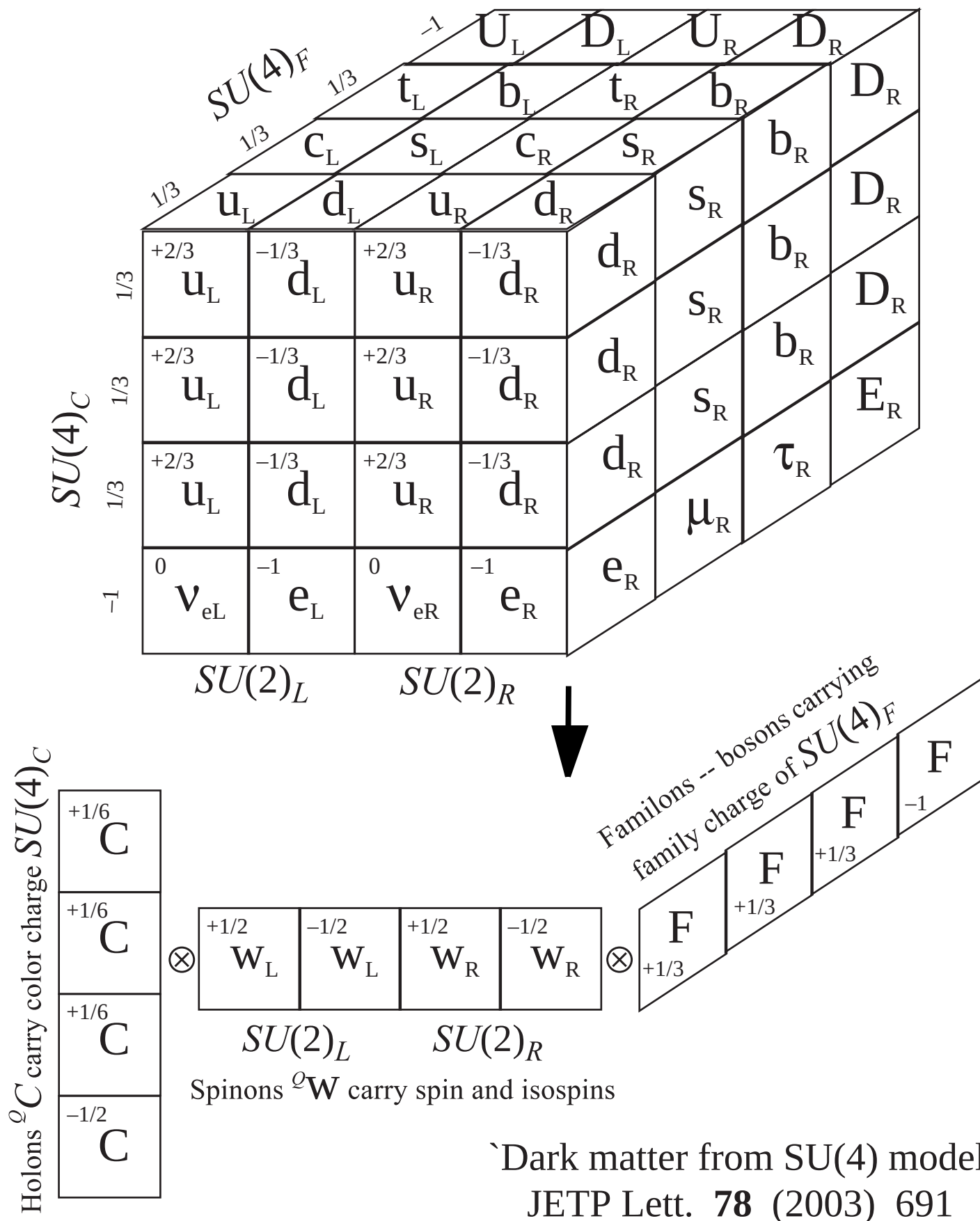
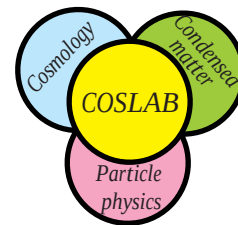
composite model of elementary particles





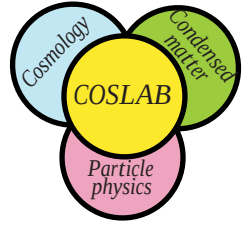
Periodic Table of elementary particles

Extended Pati-Salam model with 4 families





Momentum-space topological invariant in 2+1



general case:

$$N = \frac{1}{24\pi^2} \epsilon_{\mu\nu\lambda} \text{tr} \int d^2p d\omega \mathbf{G} \partial^\mu \mathbf{G}^{-1} \mathbf{G} \partial^\nu \mathbf{G}^{-1} \mathbf{G} \partial^\lambda \mathbf{G}^{-1}$$

simple case:

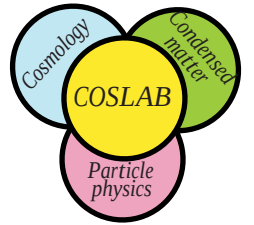
$$\mathbf{G}^{-1} = ip_0 + H(p_x, p_y) \quad , \quad H(\mathbf{p}) = \boldsymbol{\tau} \cdot \mathbf{g}(p_x, p_y)$$

$$H(\mathbf{p}) = \tau_1 g_1(p_x, p_y) + \tau_2 g_2(p_x, p_y) + \tau_3 g_3(p_x, p_y)$$

$$N = \frac{1}{4\pi} \epsilon_{\mu\nu\lambda} \int d^2p \hat{\mathbf{g}} \cdot (\partial_{p_x} \hat{\mathbf{g}} \times \partial_{p_y} \hat{\mathbf{g}})$$



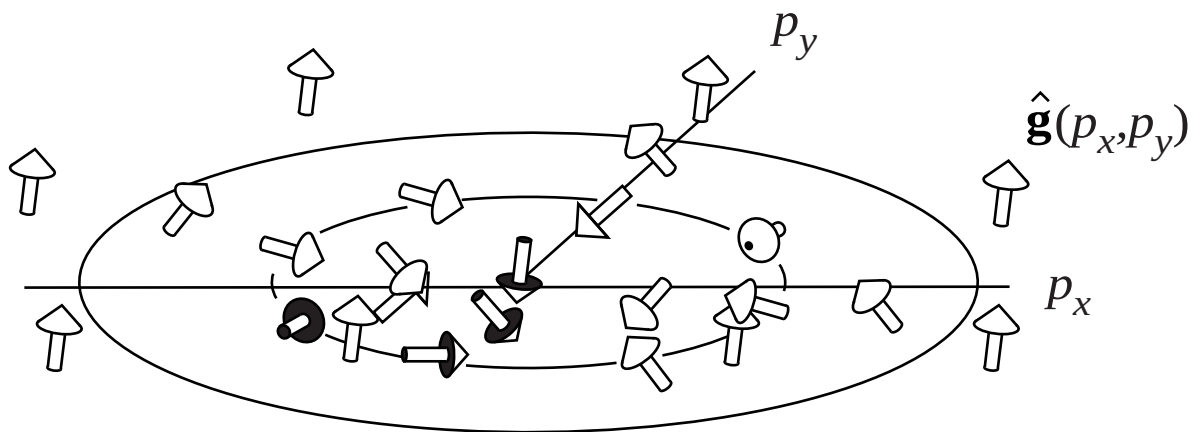
Skyrmion in momentum space



p -wave 2D superconductor & $^3\text{He-A}$ film :

$$H = \begin{pmatrix} \frac{p^2 - p_F^2}{2m} & c(p_x + ip_y) \\ c(p_x - ip_y) & -\frac{p^2 - p_F^2}{2m} \end{pmatrix} \quad \begin{aligned} g_1 &= cp_x & g_2 &= cp_y \\ g_3 &= \frac{p^2 - p_F^2}{2m} \end{aligned}$$

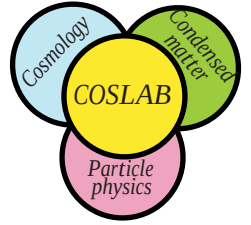
$N = 1$ skyrmion in momentum space



$$N = \frac{1}{4\pi} \epsilon_{\mu\nu\lambda} \int d^2p \, \hat{\mathbf{g}} \cdot (\partial_{p_x} \hat{\mathbf{g}} \times \partial_{p_y} \hat{\mathbf{g}}) = 1$$



N=2 skyrmion in momentum space



$(d_{xx-yy} + i d_{xy})$ -wave 2D superconductor

$$H(\mathbf{p}) = \boldsymbol{\tau} \cdot \mathbf{g}(p_x, p_y)$$

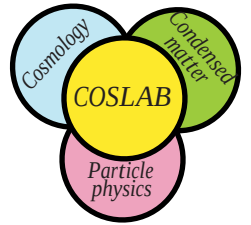
$$g_3 = \frac{p^2 - p_F^2}{2m} \quad g_1 = a(p_x^2 - p_y^2) \quad g_2 = b p_x p_y$$

$$H = \begin{pmatrix} \frac{p^2 - p_F^2}{2m} & a(p_x^2 - p_y^2) + i b p_x p_y \\ a(p_x^2 - p_y^2) - i b p_x p_y & -\frac{p^2 - p_F^2}{2m} \end{pmatrix}$$

$$N = \frac{1}{4\pi} \epsilon_{\mu\nu\lambda} \int d^2p \quad \hat{\mathbf{g}} \cdot (\partial_{p_x} \hat{\mathbf{g}} \times \partial_{p_y} \hat{\mathbf{g}}) = 2$$



Chern-Simons terms & momentum-space invariant (interplay of r -space and p -space topologies)



general case of several gauge fields

$$S_{\text{CS}} = \frac{1}{16\pi} N_{\text{IJ}} e^{\mu\nu\lambda} \int d^2x dt A_{\mu}^{\text{I}} F_{\nu\lambda}^{\text{J}}$$

$\downarrow$
 r -space invariant
 p -space invariant protected by symmetry
 $\downarrow$

$$N_{\text{IJ}} = \frac{1}{24\pi^2} e_{\mu\nu\lambda} \text{tr} Q_{\text{I}} Q_{\text{J}} \int d^2p d\omega \mathbf{G} \partial^{\mu} \mathbf{G}^{-1} \mathbf{G} \partial^{\nu} \mathbf{G}^{-1} \mathbf{G} \partial^{\lambda} \mathbf{G}^{-1}$$

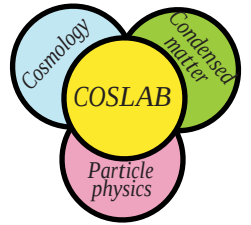
Q_{I} - charge interacting with gauge field A_{μ}^{I}

$Q=e$ for electromagnetic field A_{μ}

$Q=s_z$ for effective spin-rotation field A_{μ}^z ($A_0^z = \gamma H^z$)



Intrinsic spin-current quantum Hall effect & momentum-space invariant



$$\text{spin current } J_x^z = \delta S_{\text{CS}} / \delta A_x^z$$

$$S_{\text{CS}} = \frac{1}{16\pi} N_{\text{IJ}} e^{\mu\nu\lambda} \int d^2x dt A_\mu^I F_{\nu\lambda}^J$$

$$J_x^z = \frac{1}{4\pi} (\gamma N_{\text{ss}} dH^z/dy + N_{\text{se}} E_y)$$

2D singlet superconductor: $N_{\text{ss}} = N/4$

quantized spin Hall conductivity:

$$\sigma_{xy}^s = \frac{\gamma N}{16\pi}$$

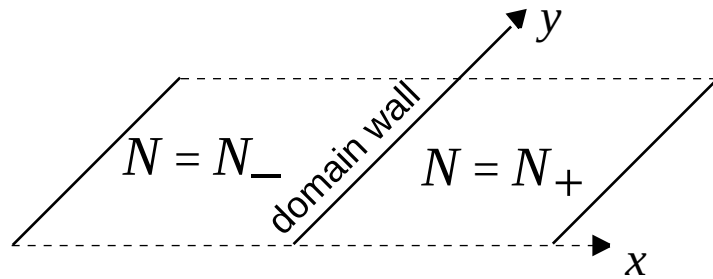
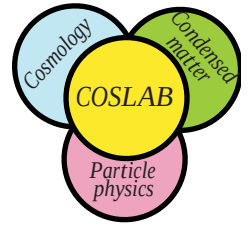
$$s\text{-wave: } N = 0$$

$$p_x + ip_y: N = 2$$

$$d_{xx-yy} + id_{xy}: N = 4$$

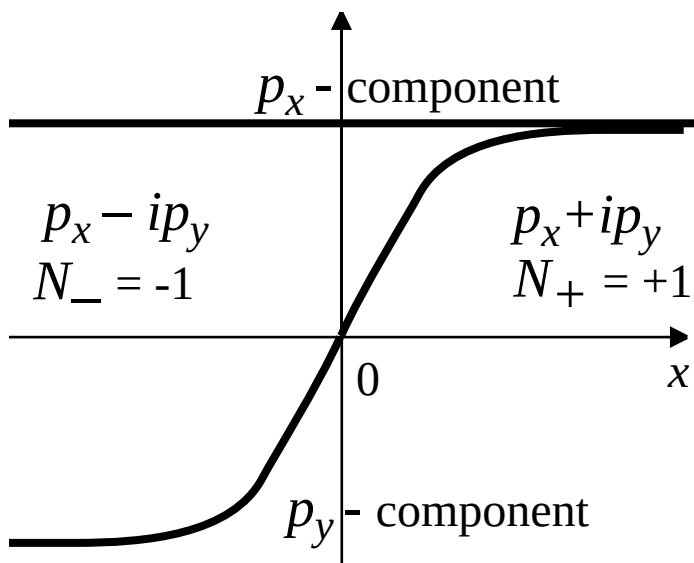


Edge states & momentum-space invariant

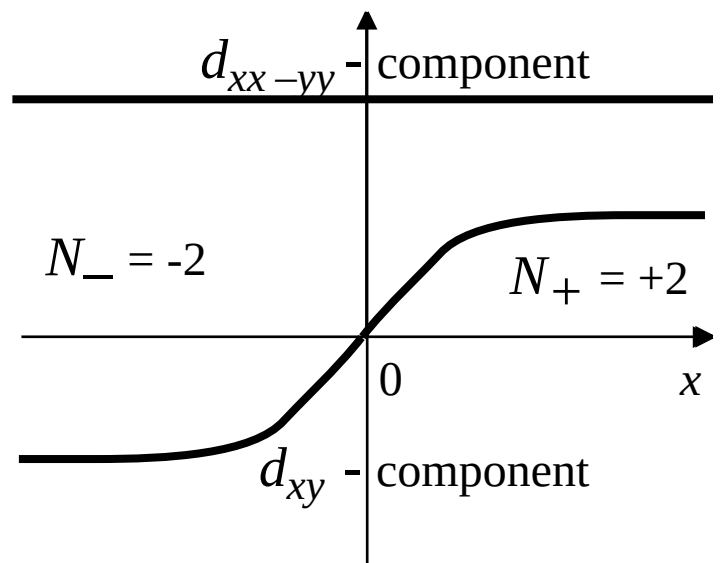


Index theorem: number of edge states $\nu = N_+ - N_-$

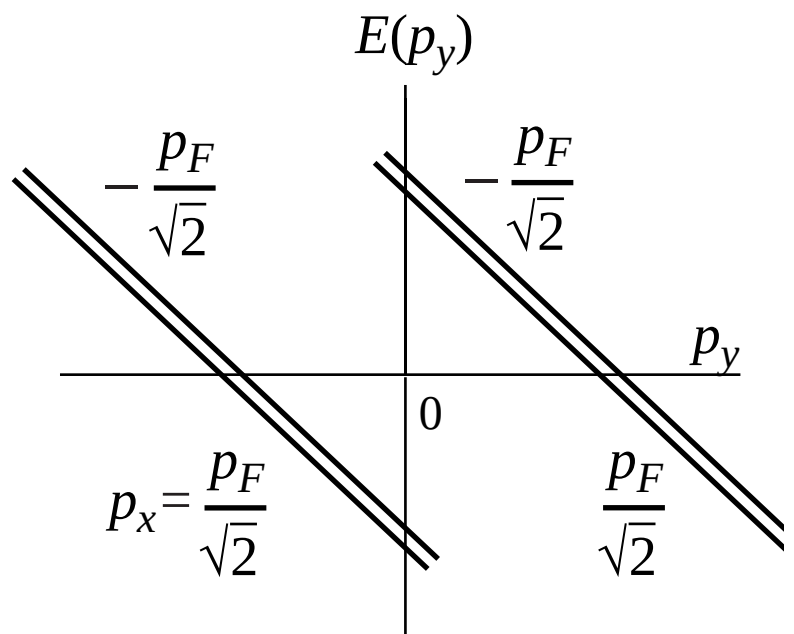
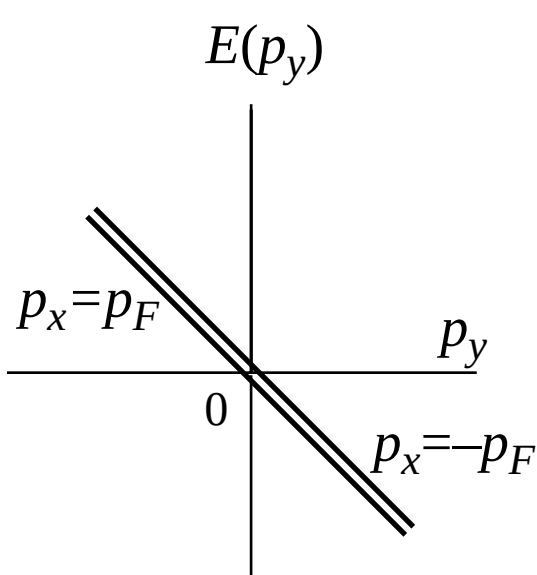
2D superconductors:



Number of edge states $\nu = 2$

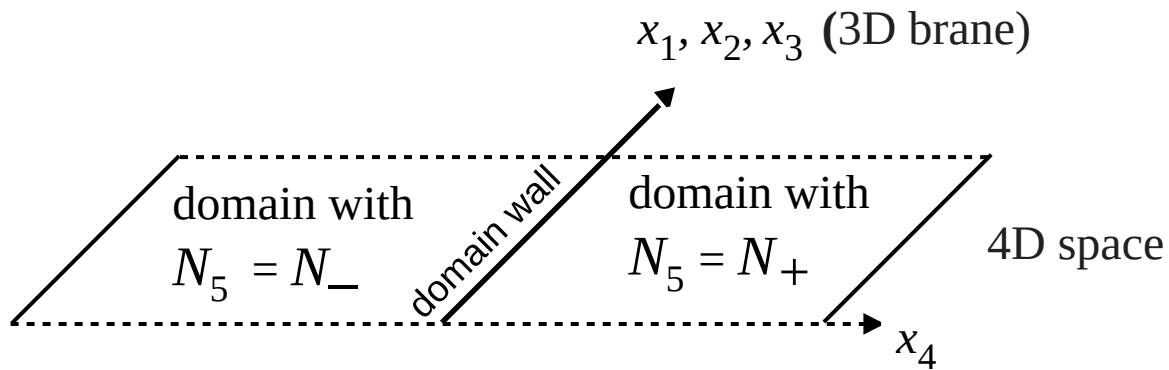
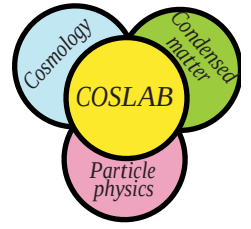


Number of edge states $\nu = 4$





Edge states -- fermion zero modes in higher dimensions & p -space invariant



Index theorem:

number of edge states living at the brane (chiral 3+1 fermions)

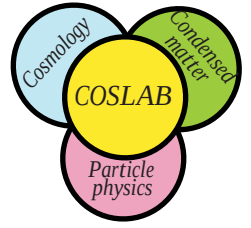
$$\nu = N_+ - N_-$$

$$N_5 = \text{tr} \int d^4 p \, d\omega$$

$$(\mathbf{G} \partial \mathbf{G}^{-1})_{\wedge} (\mathbf{G} \partial \mathbf{G}^{-1})_{\wedge} (\mathbf{G} \partial \mathbf{G}^{-1})_{\wedge} (\mathbf{G} \partial \mathbf{G}^{-1})_{\wedge} (\mathbf{G} \partial \mathbf{G}^{-1})_{\wedge}$$

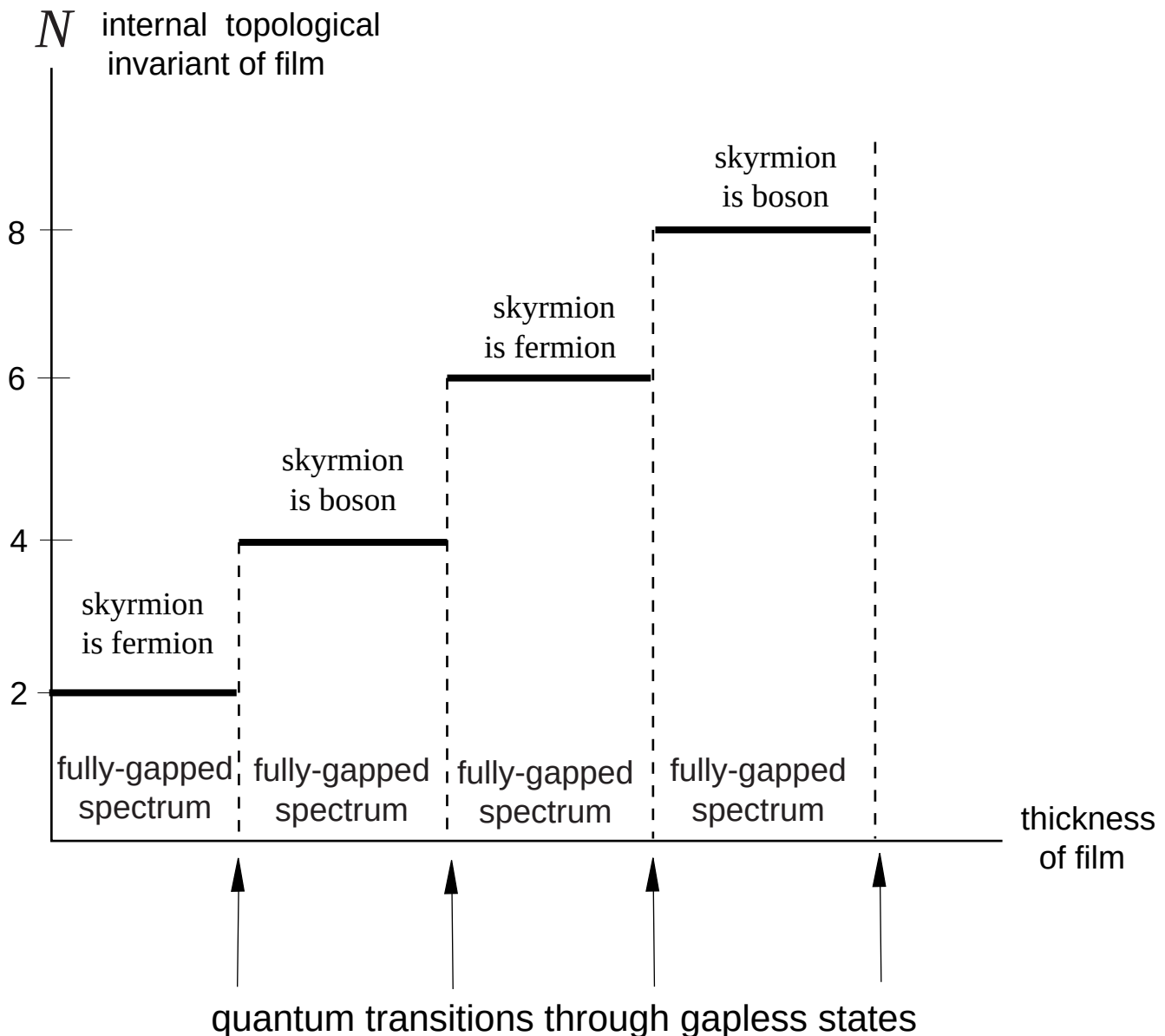


Quantum phase transition at $T=0$ as abrupt change of p -space topology



example of $^3\text{He-A}$ film

$$N = \frac{1}{24\pi^2} \epsilon_{\mu\nu\lambda} \text{tr} \int d^2p d\omega \mathbf{G} \partial^\mu \mathbf{G}^{-1} \mathbf{G} \partial^\nu \mathbf{G}^{-1} \mathbf{G} \partial^\lambda \mathbf{G}^{-1}$$

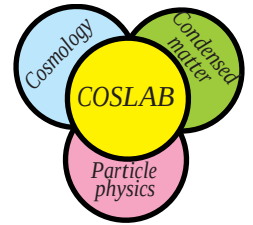


**Quantum statistics of solitons (skyrmions)
abruptly changes at quantum phase transition**

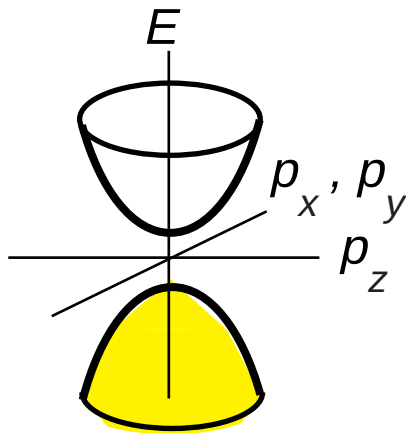
parameter of Chern-Simons action: $\theta = N \pi/2$



BEC-BCS quantum phase transition

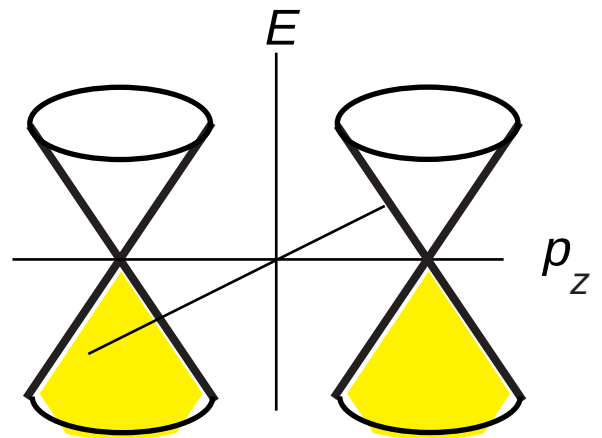


***p*-wave BEC**



**fully gapped spectrum
in strong coupling**

***p*-wave BCS**



**2 Fermi points
in weak coupling regime of BCS**

T (temperature)

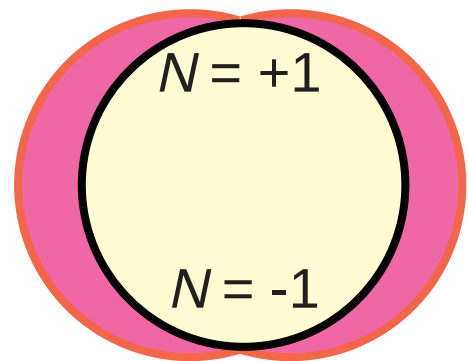
q - chemical potential μ or
inverse coupling $1/g$

fully-gapped
spectrum

spectrum with
two Fermi points

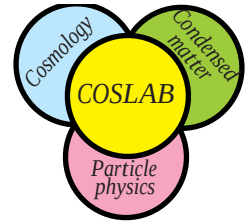
quantum phase transition at $q=q_c$,
marginal Fermi point with $N=0$ appears

for $q>q_c$, marginal Fermi point
has split into two Fermi points
with $N=\pm 1$



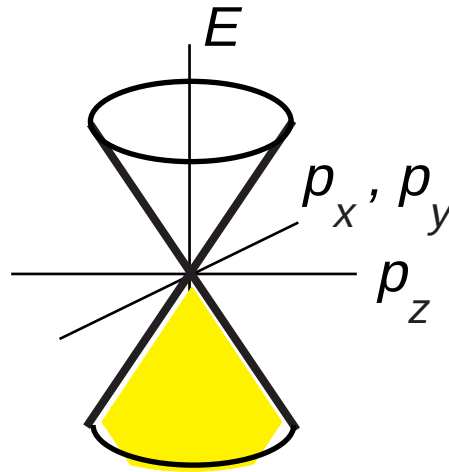


Quantum phase transition in Standard Model and BEC

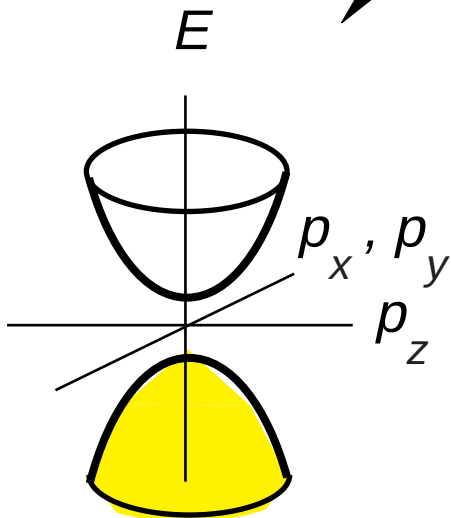


Two topological scenaria in Standard Model

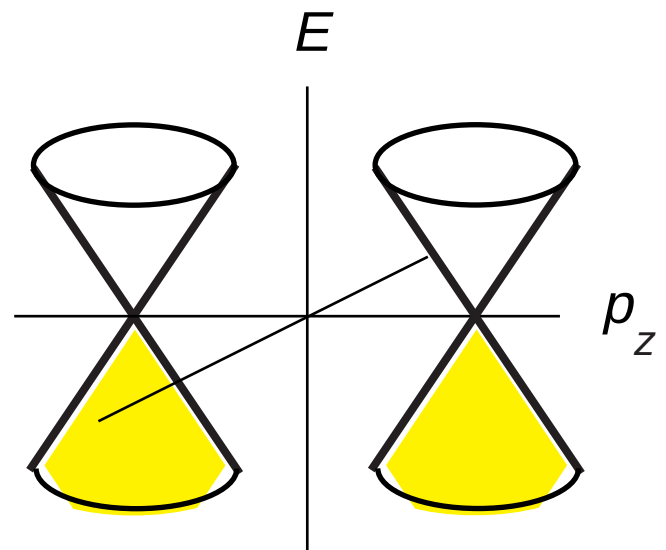
**Spectrum
of chiral (left & right)
Weyl fermions
in Standard Model**



**Marginal
Fermi point
 $N = +1 - 1 = 0$**



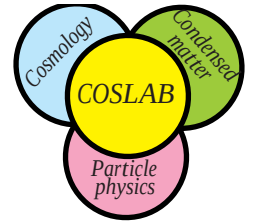
**Marginal Fermi point
disappears,
massive Dirac fermions are formed**



**Marginal Fermi point
splits
into topologically protected
Fermi points with $N = +1$ and $N = -1$**



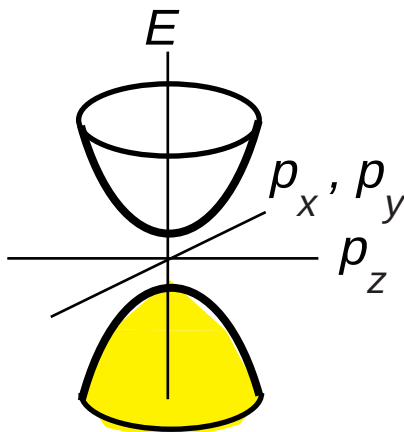
Topological quantum phase transition in Standard Model



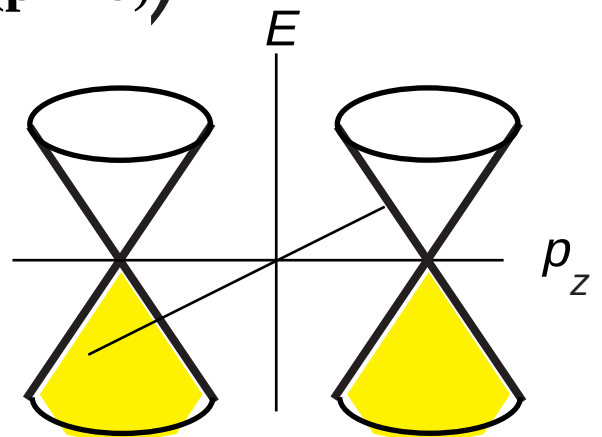
Hamiltonian for neutrino

$$H = \begin{pmatrix} \sigma \cdot (\mathbf{p} - \mathbf{b}) & M \\ M & -\sigma \cdot (\mathbf{p} + \mathbf{b}) \end{pmatrix}$$

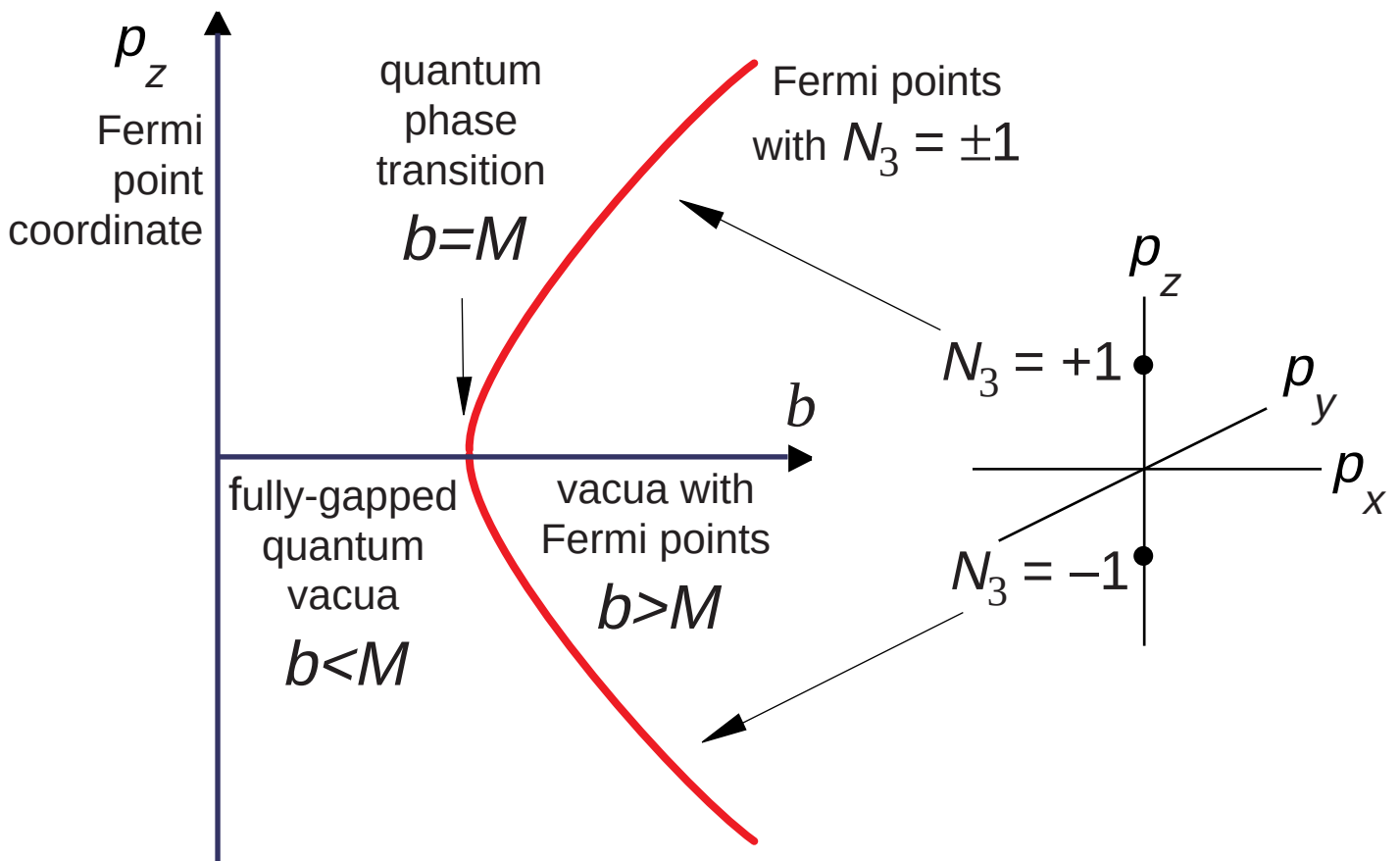
possible scenario
for neutrino
since its mass
 M is small

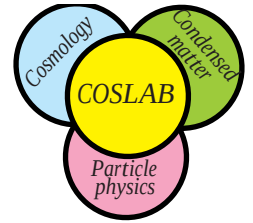


fully gapped at $b < M$



Fermi points $\mathbf{P}_a = \pm \mathbf{b} (1 - M^2/b^2)^{1/2}$
appear at $b = M$ & split at $b > M$





Chern-Simons term in 3+1 action from splitting of Fermi points & *p*-space topology

$$S_{\text{CS}} = e^{\mu\nu\lambda\gamma} k_{\mu} \int d^4x A_{\nu}(x) \partial_{\lambda} A_{\gamma}(x)$$

parameter k is determined
by topology
in momentum space

topological term
in real space

$$\mathbf{k} = \mathbf{N}_{\text{pQQ}} = (1/24\pi^2) \sum_a \mathbf{P}_a N_a Q_a^2$$

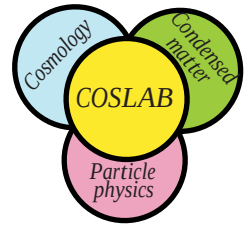
position of
 a -th Fermi point in $\mathbf{p}$ -space

U(1) charge of fermions
living near Fermi point

topological charge of a -th Fermi point



Conclusion



The momentum-space topology determines universality classes of quantum vacua

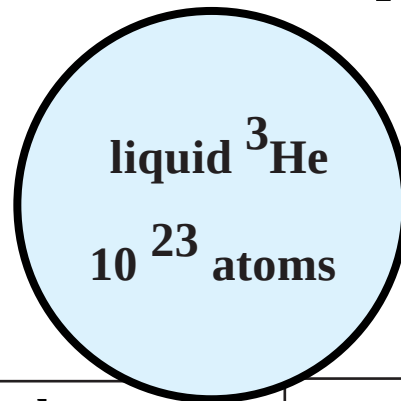
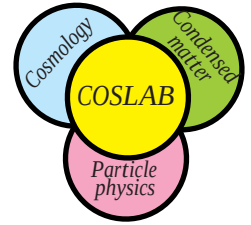
Vacuum of Standard Model belongs to Fermi-point universality class; elementary particles in vacua of this universality class are chiral fermions emerging near Fermi points; gravity and gauge fields are low-energy collective modes, either fundamental or emerging due to Fermi points

The momentum-space topology determines also:

quantization of Hall and spin-Hall conductivity;
prefactor in topological Chern-Simons & Wess-Zumino terms;
quantum statistics of topological objects;
chiral anomaly & vortex dynamics
spectrum of edge states & fermion zero modes on branes & strings;
quantum phase transitions;
existence of mass of Standard-Model fermions;
etc.



Quantum Fermi liquids



complicated many-body system
of strongly interacting
strongly correlated atoms

normal ${}^3\text{He}$, ${}^3\text{He-A}$, ${}^3\text{He-B}$,
normal metals, semiconductors,
superconductors, etc

Effective theory in low temperature limit

Fermionic quasiparticles + Bosonic collective modes = QFT
elementary particles of effective theory

Type of Quantum Field Theory depends on universality class

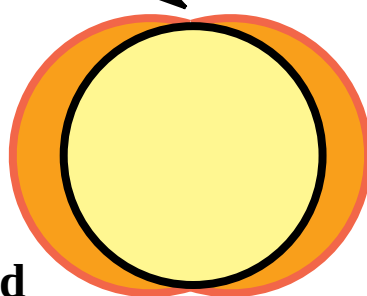
universality class of Fermi points

${}^3\text{He-A}$, Standard Model

Effective theory in low temperature limit

Chiral fermions + Gauge fields & gravity = Relativistic QFT

left-handed
fermions live here



right-handed
fermions live here

emergent phenomena:

Gauge invariance
Lorentz invariance
General covariance (partly)
chiral fermions
gauge fields
gravity
spin