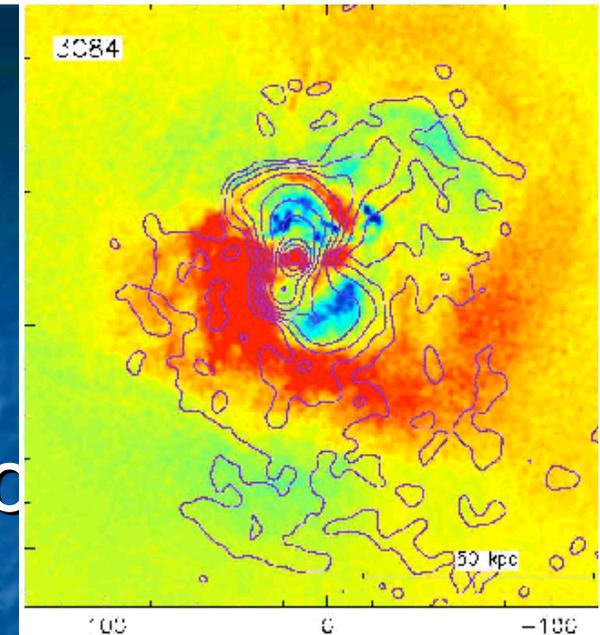


# How AGN Terminate Galaxy Formation

James Binney  
Oxford University

# Introduction

- XMM spectra show `cooling-flow` cooling ( $T_{\min} \sim T_{\text{vir}}/3$ )
- Chandra makes impact of radio galaxies evident
- By what mechanism does an AGN thermostat the ICM?
- What implications does this mechanism have for galaxy formation?  
(astro-ph/0308172; astro-ph/0407238)



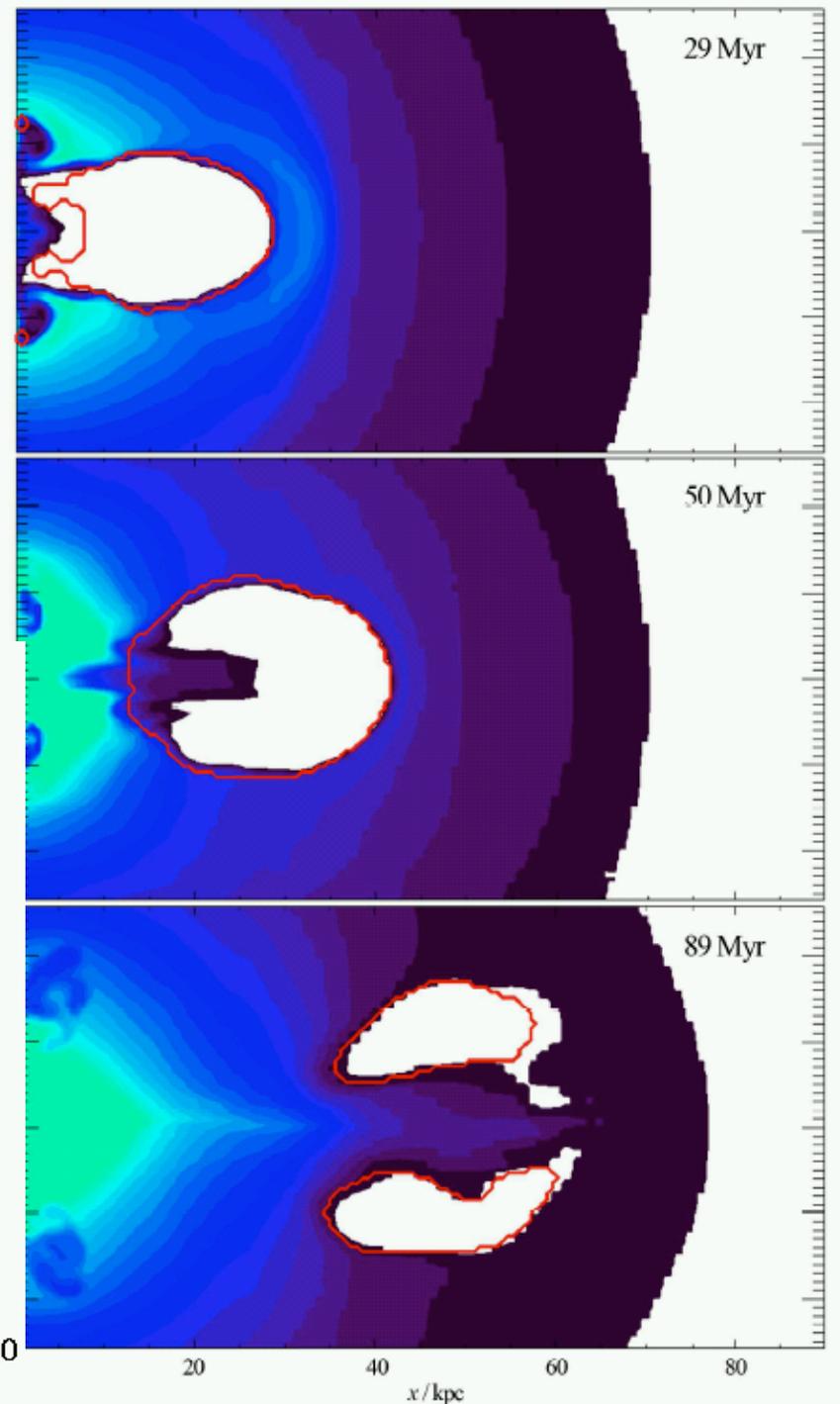
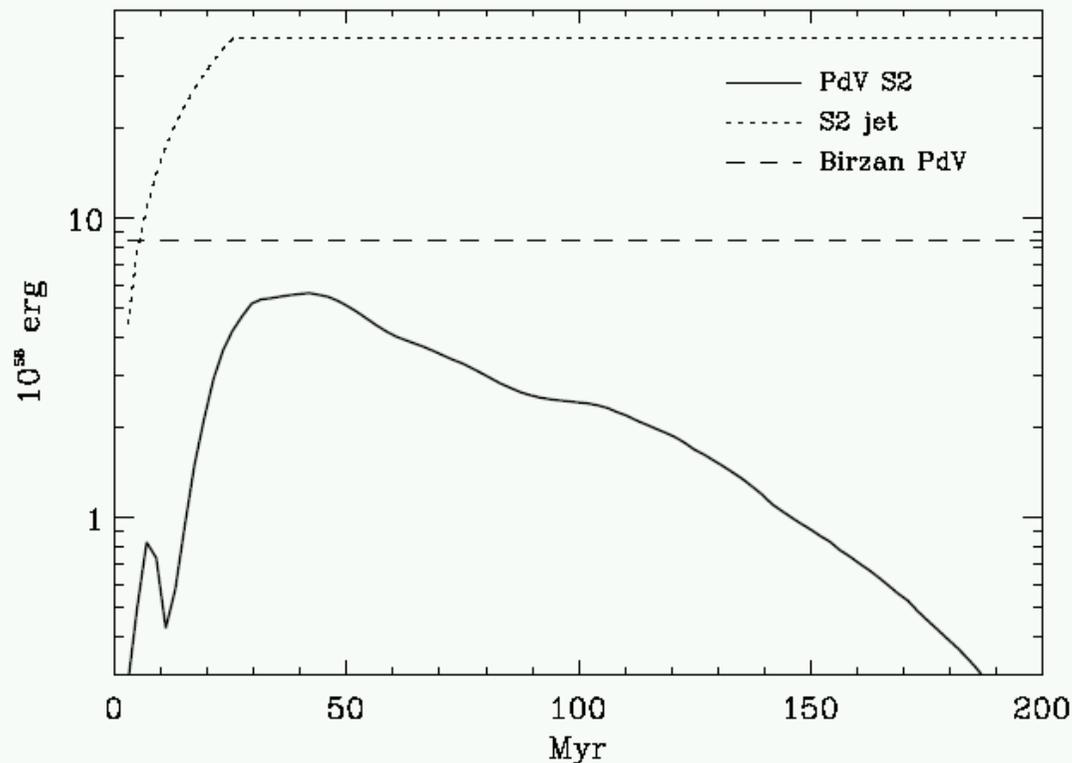
# Anisotropic heating

- Absence of isentropic cores implies energy thermalized at many radii simultaneously
- Expected of jet heating

# Bubble dynamics

- Omma et al (2004), Omma & Binney (2004)  
Omma (2005) PhD thesis  
3D adaptive-grid hydro simulations with  
Bryan-Norman code ENZO (piecewise  
parabolic solver)
- Jets fired into radiatively cooling model of  
Hydra ICM at  $\gg 2\text{kpc}$  (David et al 2000) [movie](#)

- Define cavities by  $\rho < \rho_0/4$
- Evaluate PV
- Peaks at only 10% of actual input



# Statistics of observed cavities

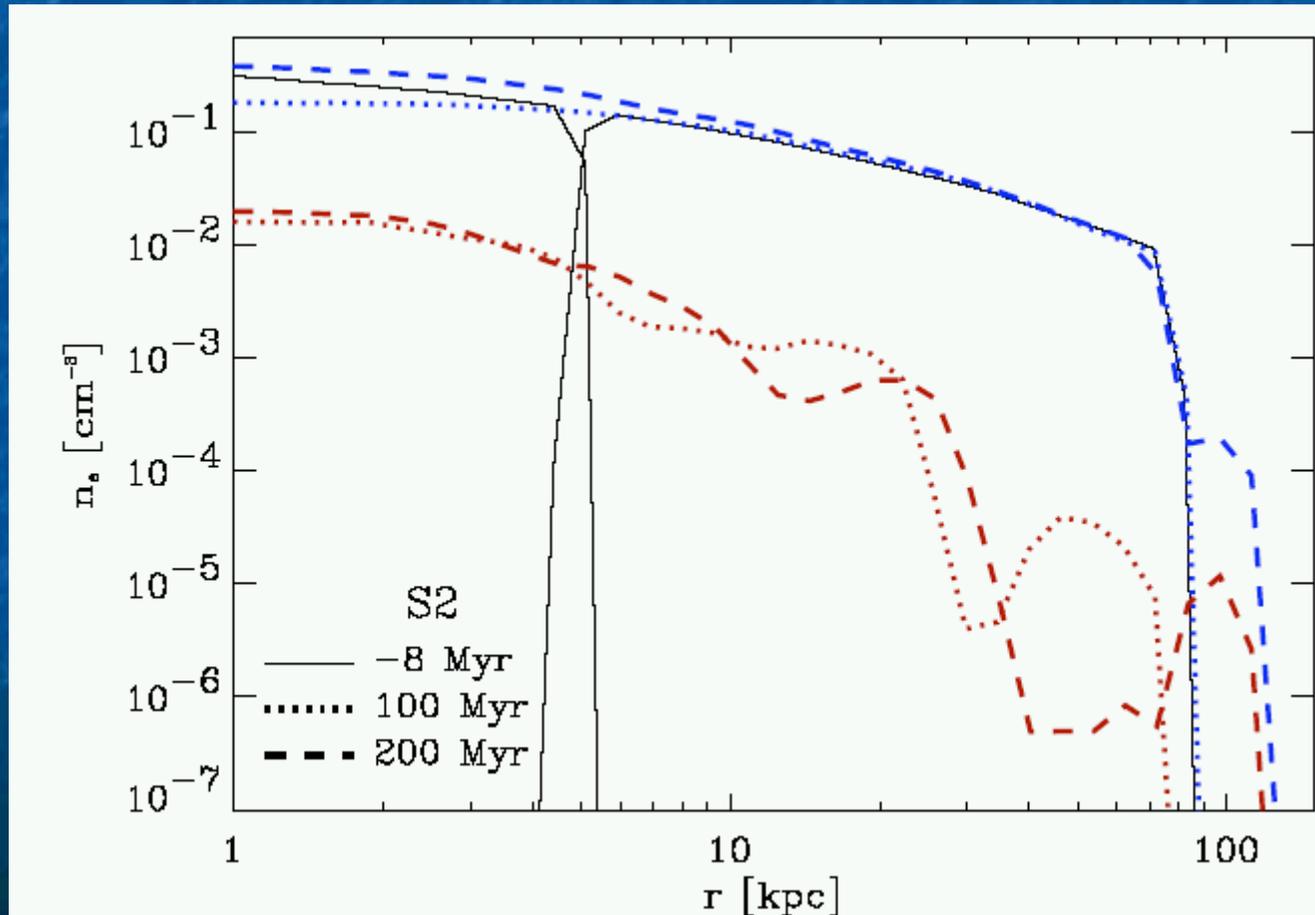
TABLE 1  
PARAMETERS FOR FIVE CLUSTERS WITH CAVITIES.

System	$PV/10^{58}$ erg	$L_X/10^{43}$ erg s $^{-1}$	$\tau$ /Myr	Reference
Hydra A	27	30	88	McNamara et al. (2000); Nulsen et al. (2002)
A2052	4	3.2	122	Blanton et al. (2001)
Perseus	8	27	29	Fabian et al. (2000); Allen et al. (1992)
A2597	3.1	3.8	79	McNamara et al. (2001)
A4059	22	18	119	Huang & Sarazin (1998); Heinz et al. (2002)

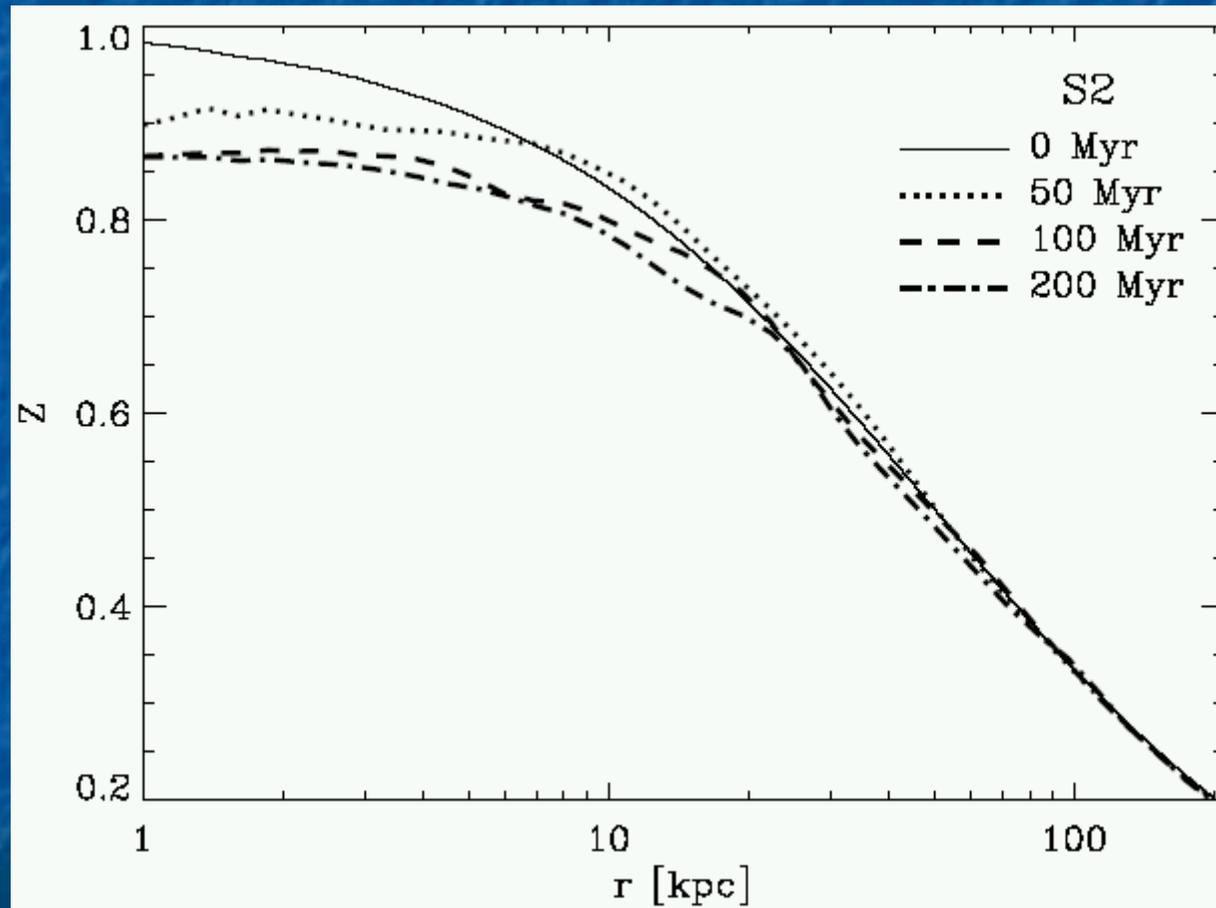
- Rate of heating by cavities  $\sim L_X$  for recurrence time  $> 100$  Myr

# Shifting gas around

- Follow tracer dye from (a)  $r < 5 \text{ kpc}$ , and (b)  $5 < r < 77 \text{ kpc}$



# Effect on Z gradient

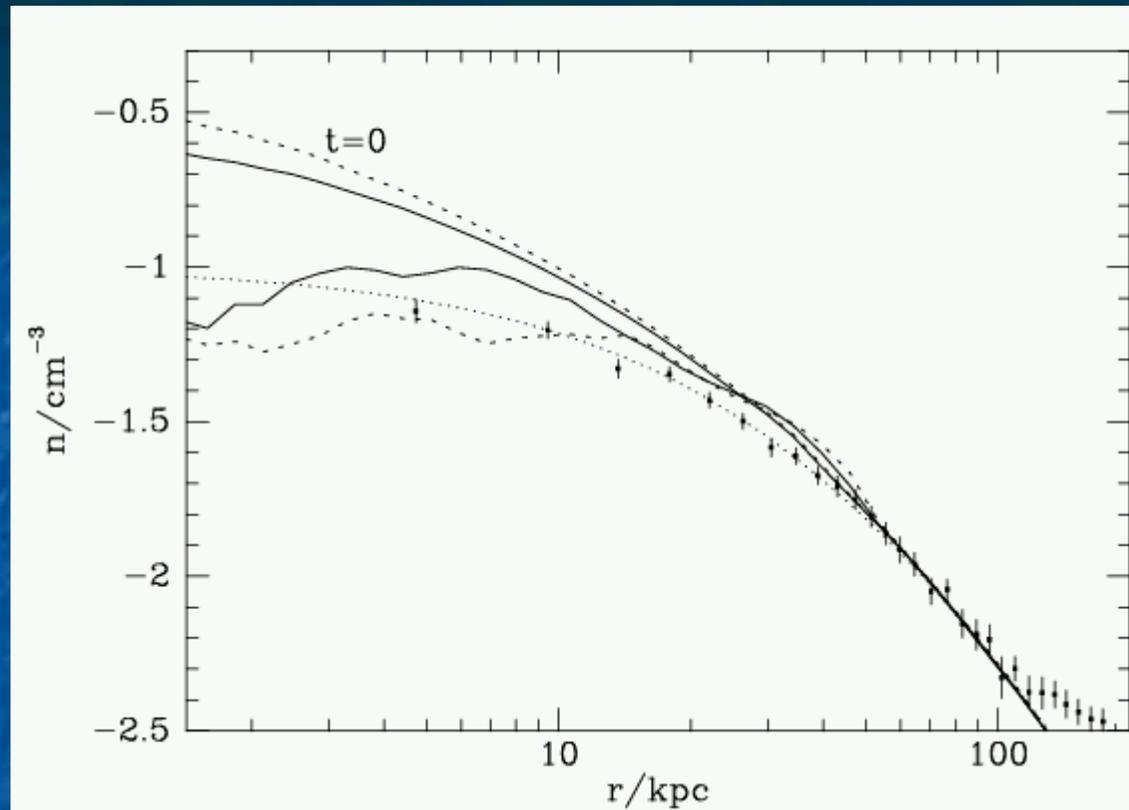


# Fixing the radial density profile

- For steady state,  $E(r)$  must match  $j_x(r)$
- A more powerful jet disrupts further out
- A more concentrated profile disrupts jet further in
- Later jet ignition \_ bigger outburst
- Later ignition \_ more centrally concentrated density profile

# Two Simulations

- Start from present configuration of Hydra
- Wait (i) 262 Myr (ii) 300 Myr
- In (ii) extra  $4 \times 10^{59}$  erg lost to radiation, so add  $8 \times 10^{59}$  erg rather than  $4 \times 10^{59}$  erg as in (i)
- $E_{\text{Bondi}} = 5(M/10^9 M_{\odot})^2 \times 10^{59}$  erg in 262 Myr;  
 $E_{\text{Bondi}} = 7(M/10^9 M_{\odot})^2 \times 10^{59}$  erg in 300 Myr



- Outbursts have undone 300 Myr of cooling
- System with later ignition ends less centrally concentrated
- Implies that systems can oscillate around an attracting profile

# Conclusions so far

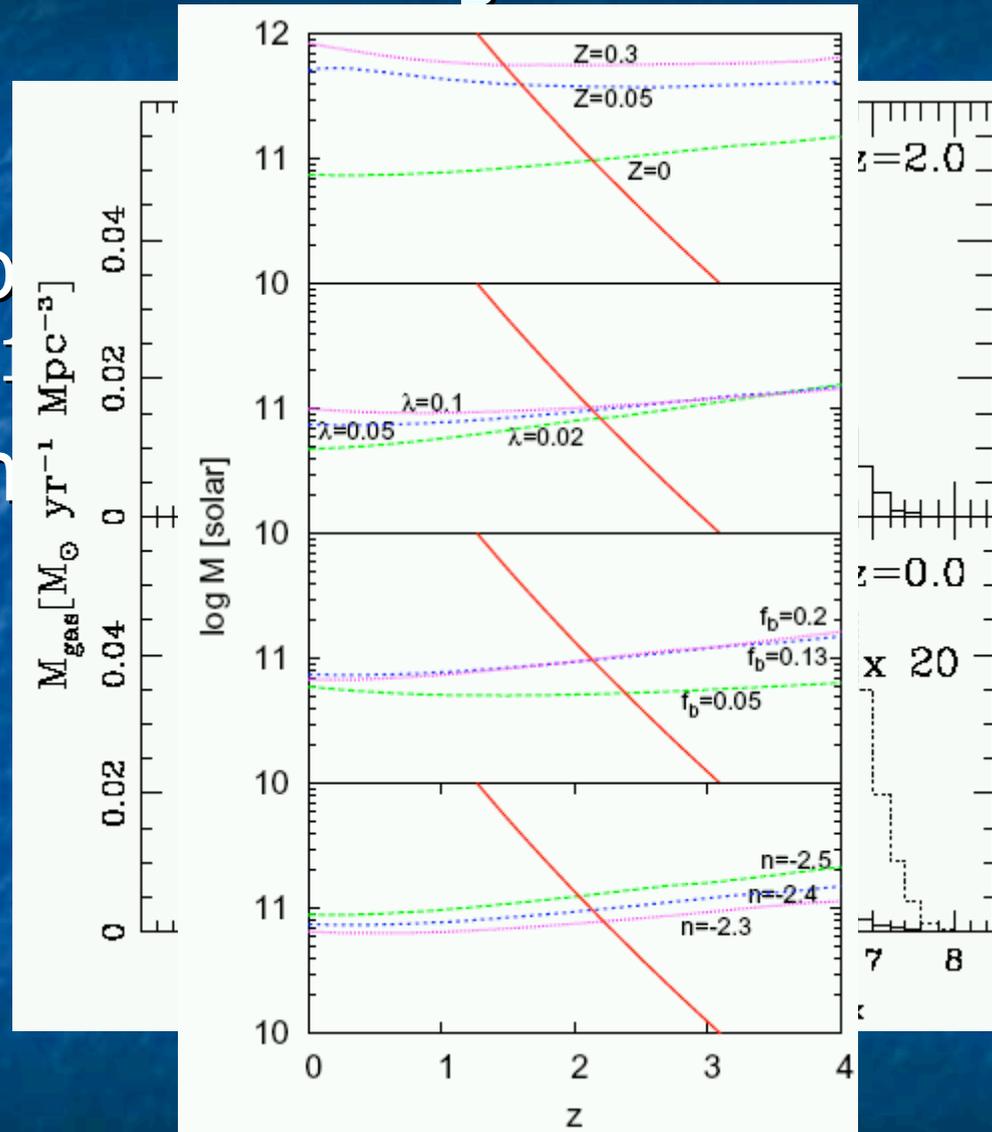
- Bubbles carry only a fraction of the energy injected by radio sources
- Radio sources can readily replace radiated energy
- If energy released in an outburst is proportional to energy radiated since last outburst, post-outburst density profile will oscillate around observed one
- Radio sources thermostat  $T_{\text{vir}}$  atmospheres

# Implications for Galaxy Formation

- In classical theory, galaxies form by cooling of gas trapped at  $T_{\text{vir}}$
- If new picture of cooling-flow dynamics correct, classical theory of GF false

# So how *do* galaxies form?

- Fraction (Binney & Tremaine 1987)
- $f \sim 1$  on

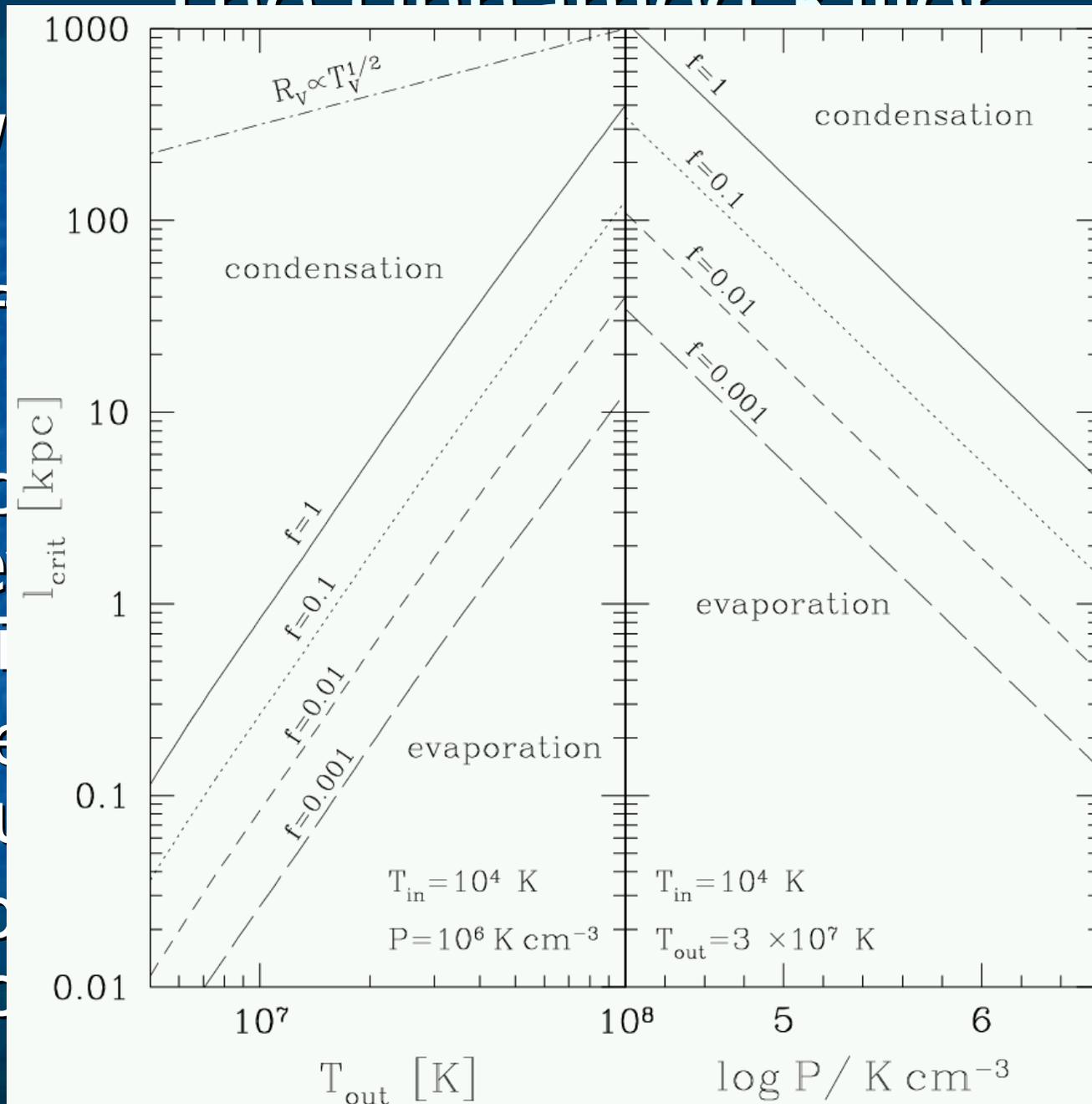


at  $T_i, T_{\text{vir}}$   
(Binney & Tremaine 2003)

OW

# The Unnamed Killer

- W
- &
- Ef
- T
- So
- re
- wi
- He
- clu
- Lo
- ho



plates

y with

to

y

but

use

# Connection with BH growth

- BH growth known to take place in bursts:
- Yu & Tremaine (2002) find
- (i) AGN have radiated in optical/UV as much E as released by all nuclear BHs;
- (ii)  $L \sim L_{\text{Edd}}$  and  $\bar{\alpha} > 0.1$  needed to produce observed quasars from observed BHs

# Lifecycle of supermassive BHs

- @  $L_{\text{Edd}}$   $M \sim \exp(t/t_{\text{Salpeter}})$ ;  $t_{\text{Salpeter}} \sim 25 \text{ Myr}$
- So  $M$  from  $10^3 M_{\odot}$  to  $10^9 M_{\odot}$  with  $14 t_s \sim 0.4 \text{ Gyr}$  and  $10 \text{ Gyr}$  at  $< 0.05 L_{\text{Edd}}$
- Magorrian relation  $M \sim M_{\text{bulge}}$ , & high  $\alpha/\text{Fe}$  of bulges, & high ages of bulges all imply  $L_{\text{Edd}}$  (quasar) phase associated rapid star formation
- Conjecture this is when there is cold gas @ centre
- Episodes end when well deep enough to trap  $10^7 \text{ K}$  gas; then  $\dot{M}$   $0.002$  to  $0.02 M_{\odot}/\text{yr}$  to offset  $10^{43} - 10^{44} \text{ erg/s}$  of  $L_x$

- Stars form from cold gas & then heat all gas
- When  $M < M^*$ , heated gas expelled
- When  $M > M^*$ , it accumulates  
(Larson 79; Dekel & Silk 86)
- Low  $M$  galaxies suppressed by photoionization, evaporation & SN feedback (Efstathiou 92; Dekel & Silk 86; Dekel 04)
- Without AGN heating ejected mass forms very massive objects later (Benson et al 03)
- AGN prevent hot gas cooling

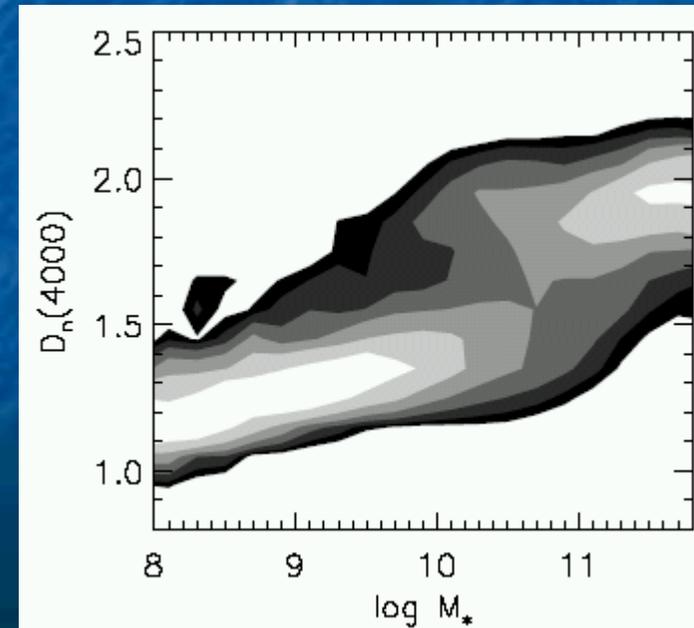
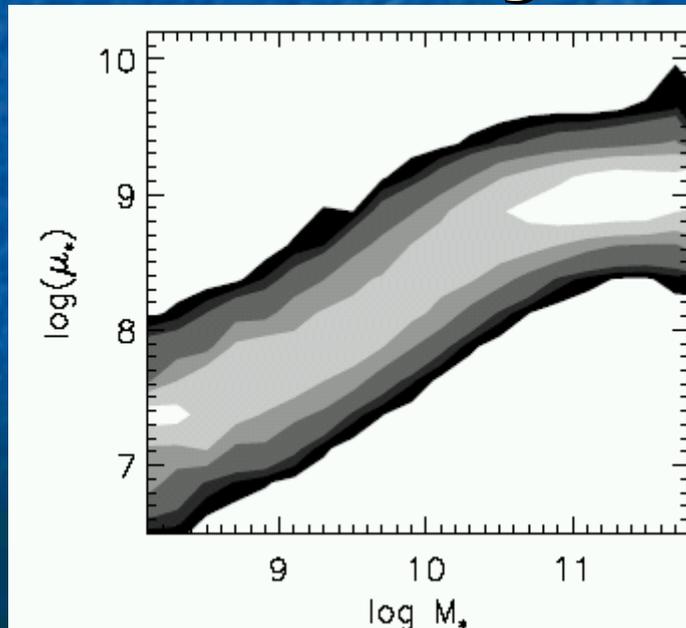
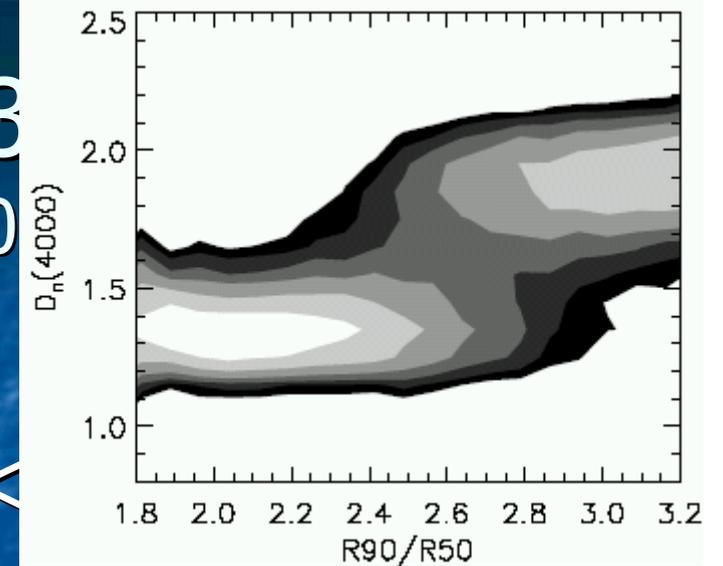
# Conclusions

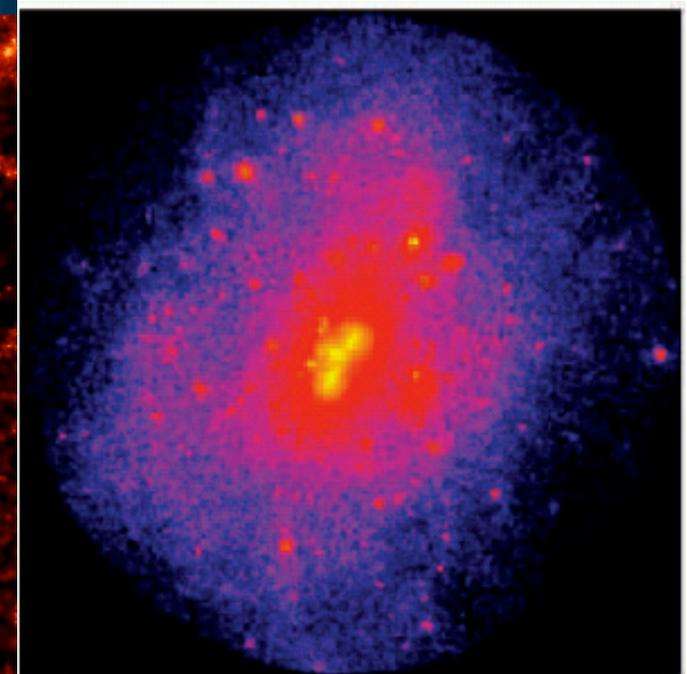
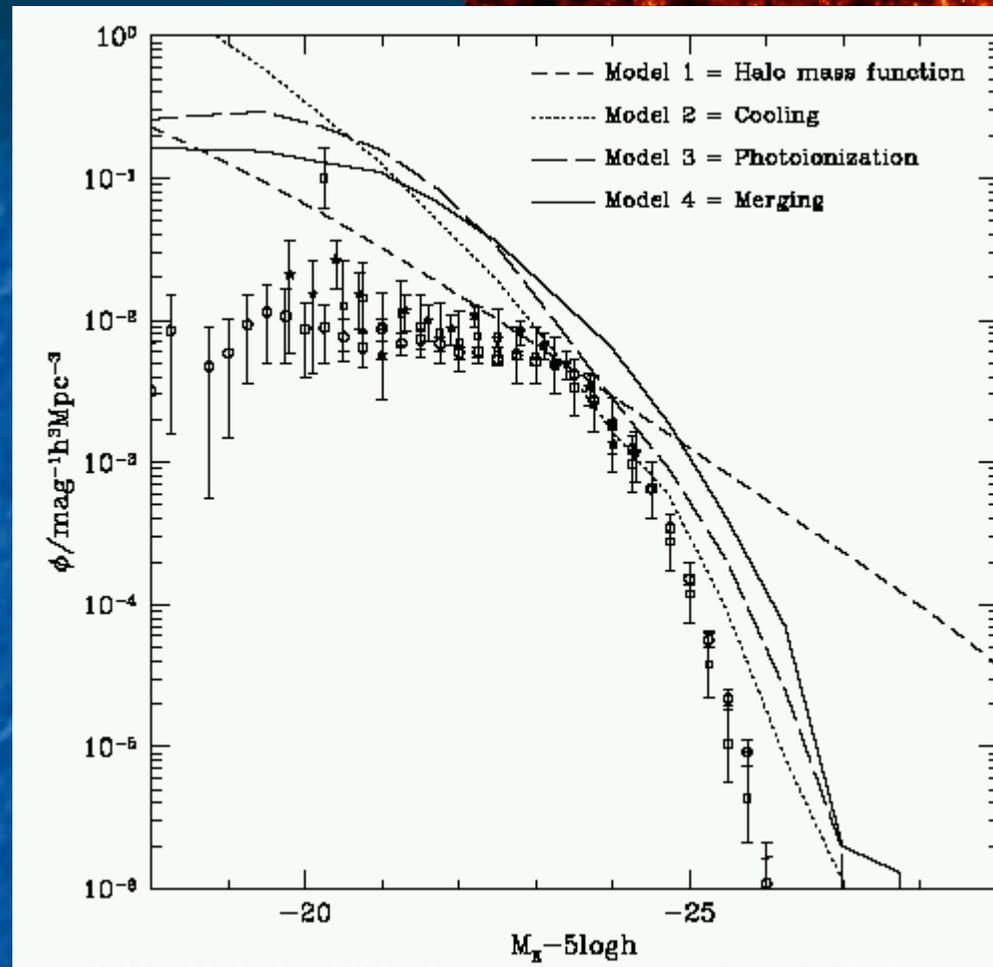
- A CF in most deep potential wells
- Is reheated every  $\sim 100$  Myr
- Almost certainly jet heated – non-adiabatic
- CFs oscillate around a stable radial density profile
- Galaxies do not form from cooling gas
- BHs grow principally from cold gas simultaneously with rapid SF in bulge
- BHs heat CFs when in quiescent state
- Thermostatic control of CFs by AGN prevents formation of galaxies with  $M > M^*$

# Characteristic $M^* = 3$

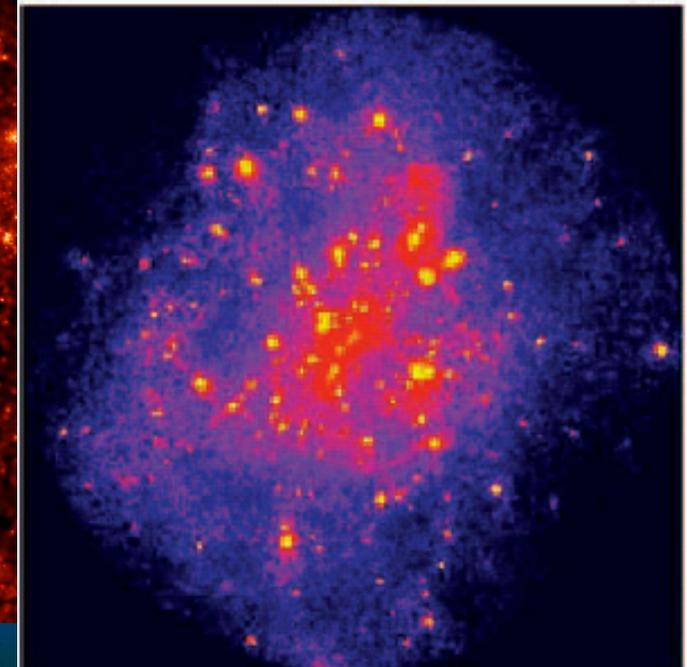
(Kauffmann et al 0

- At  $M > M^*$   $dSB/dM=0$ ; at  $M < M^*$
- At  $M > M^*$  galaxies old; at  $M < M^*$  younger
- At  $M > M^*$  light centrally concentrated





45 kpc



## ■ Spectrum

# Classical Theory of GF

- Gas heated to  $T_{\text{vir}}$  on falling into  $\_$   
(Rees & Ostriker 1977; White & Rees 1978)
- Deep potential wells contain gas @  $T_{\text{vir}}$
- So we should be able to see galaxy formation!

# BH – CF coupling

- Gross mismatch of timescales ( $t_{\text{BH}} \sim 1\text{yr}$ ;  
 $t_{\text{CF}} \sim 100\text{Myr}$ )
- How can BH no bigger than solar system stabilize a system  $>100$  kpc in size?
- Two part answer: (i)  
tight connection BH power to  $L_x$ ; (ii)  
servo system for radial distribution of E

# In M87

- Chandra resolves  $r_{\text{Bondi}}$
- $\dot{M}_{\text{Bondi}} = 0.1 \text{ M.}/\text{yr}$  (Di Matteo et al 03)
- So  $L = 5 \times 10^{44} \text{ erg/s}$  if  $0.1 mc^2$  released
- $L_x(<20\text{kpc}) = 10^{43} \text{ erg/s}$  (Nulsen & Boehringer)
- $L_x(\text{AGN}) < 5 \times 10^{40} \text{ erg/s}$
- $L_{\text{Mech}}(\text{jet}) = 10^{43} - 10^{44} \text{ erg/s}$   
(Reynolds et al 96; Owen et al 00)
- So BH accreting at near  $\dot{M}_{\text{Bondi}}$  & heating on kpc scales with high efficiency (Binney & Tabor 95)

# Bondi accretion

- Area of sonic flow

$$A = 4\pi \left( \frac{GMm_p}{k_B T_0} \right)^2$$

- Particle density

$$n \simeq P/k_B T_0$$

- Accretion rate

$$\dot{M} \simeq nm_p A c_s \simeq 4\pi P (GM)^2 \left( \frac{m_p}{k_B T_0} \right)^{5/2}$$

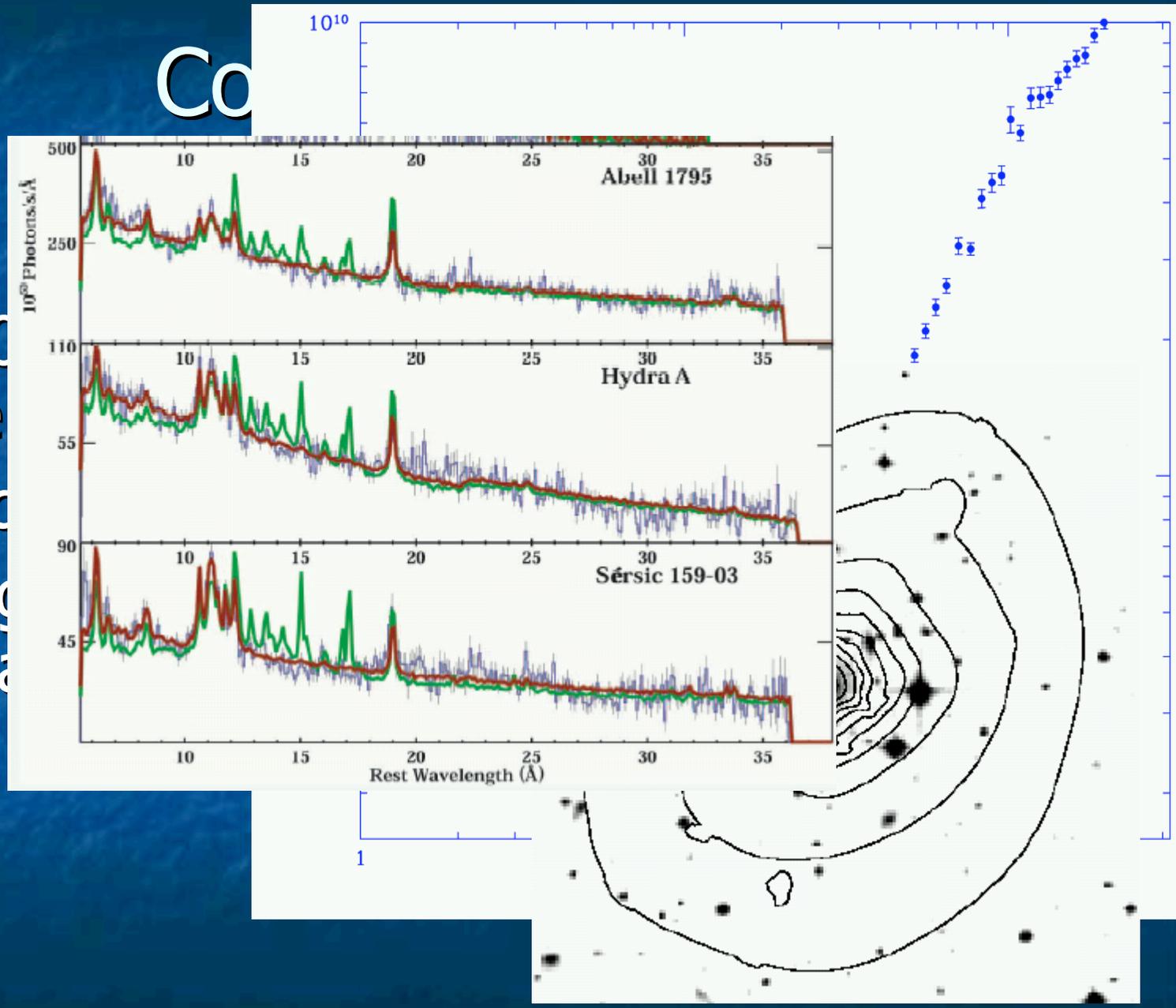
- Luminosity

$$L = \int d^3\mathbf{x} n^2 \Lambda(T) \simeq \int d^3\mathbf{x} \left( \frac{P}{k_B T} \right)^2 \Lambda$$

- So balance possible with  $E \sim \int dt L_x$

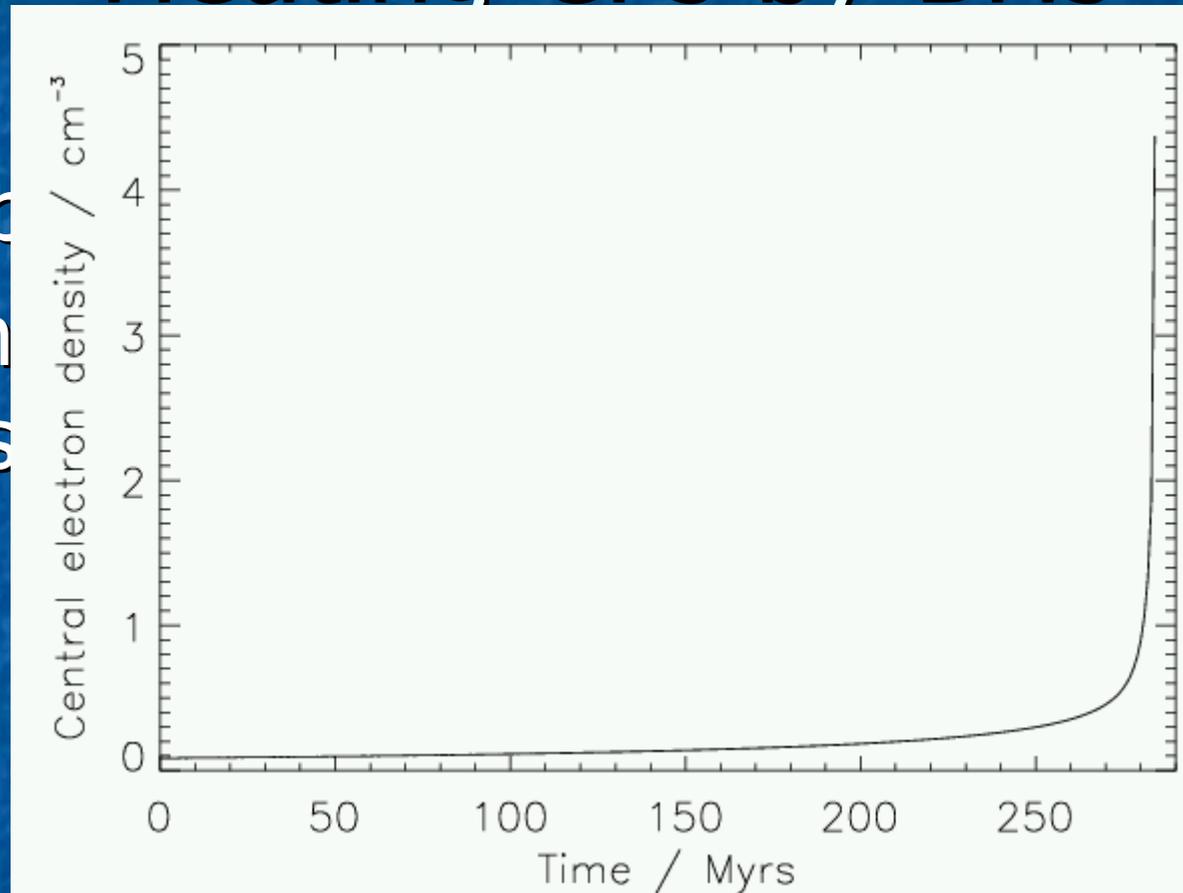
Co

- In  $>70\%$  centre
- $t_{\text{cool}} = 0$
- But if  $g$  spectra



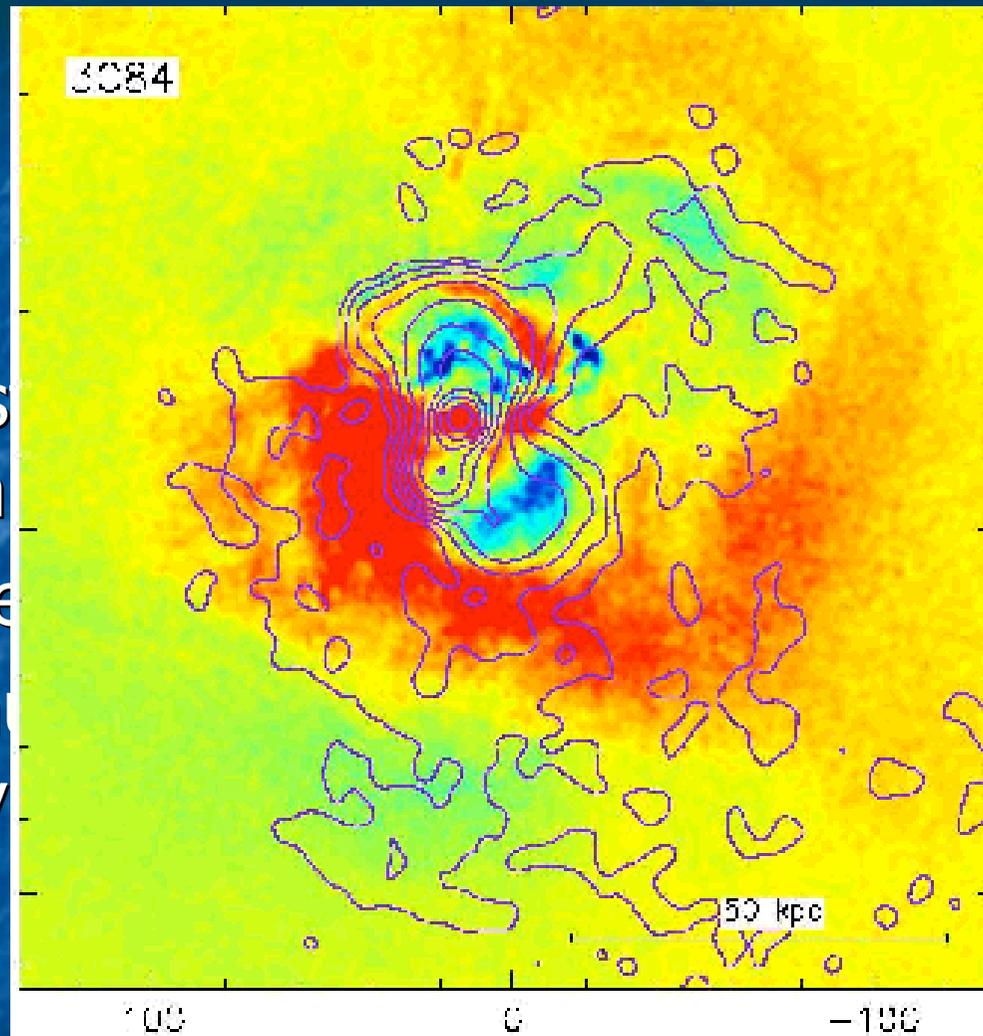
# Heating CFs by BHs

- In ab
- Such
- a res



cool(0)  
voke

- Suggests simultaneous
- Expected
- Now more ray cavity



Y radii

X-

