

SMR.1667 - 6

Summer School and Miniconference on  
**Dynamical Mean-Field Theory for Correlated Electrons:  
Applications to Real Materials, Extensions and Perspectives**  
25 July - 3 August, 2005

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## **Approximate Treatments and Cluster Extensions III**

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**Institute for Theoretical Physics**  
**37077 Gottingen**  
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These are preliminary lecture notes, intended only for distribution to participants



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# Cluster Extensions to the DMFT systems



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1. Motivation: Why need a cluster MFT?



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2. Realizations: DCA and CDMFT



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2. Realizations: DCA and CDMFT
3. Selected results for DCA



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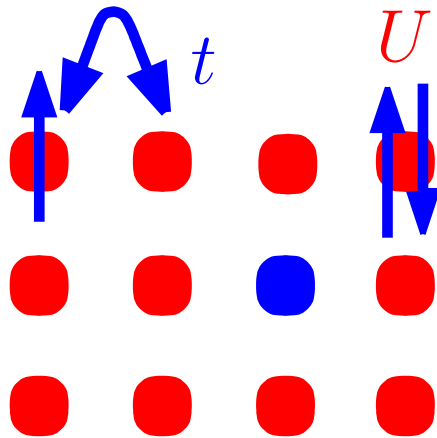


# Cluster Extensions to the DMFT

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## DMFT in a nutshell



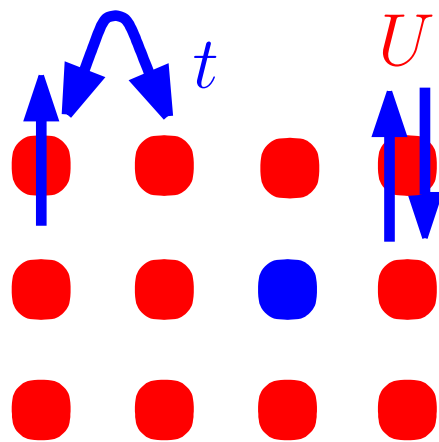
$$H = \sum_{\langle i,j \rangle \sigma} t_{ij} d_{i\sigma}^\dagger d_{j\sigma} + \frac{U}{2} \sum_{i\sigma} n_{i\sigma} n_{i\bar{\sigma}}$$

non-local and dynamical

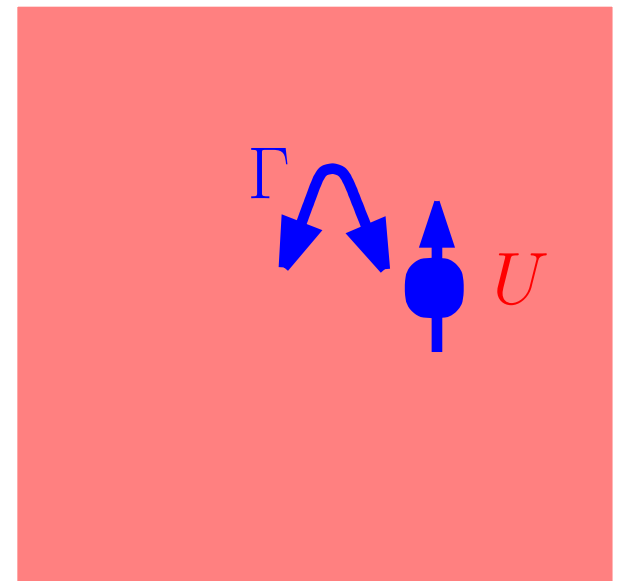




## DMFT in a nutshell



$$\Sigma(\vec{k}, z) \longrightarrow \Sigma(z)$$



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non-local and dynamical

$$H_{\text{DMFT}} = H_{\text{Bath}} + H_U + H_{\text{Hyb}}$$

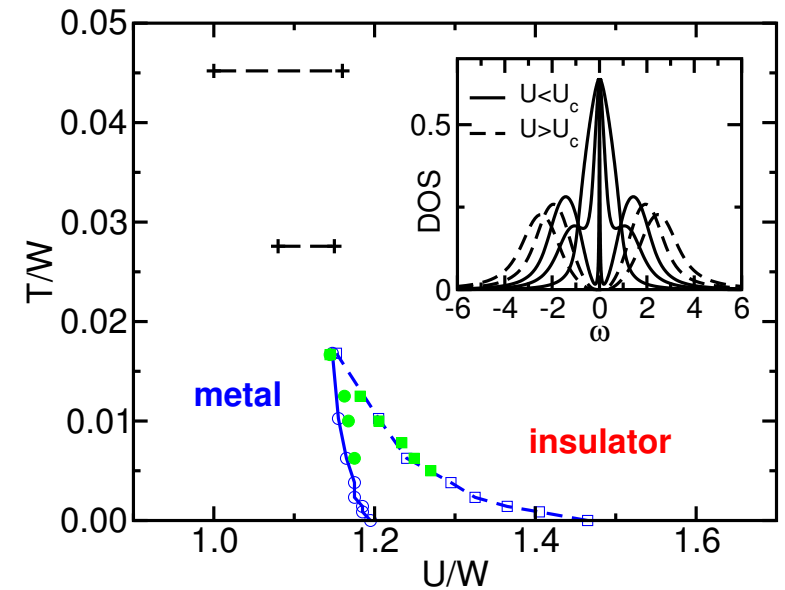
local and dynamical



Key results from DMFT:

- Metal-insulator transition for  $\langle n \rangle = 1$

Georges et al., RMP '96, Bulla et al., PRL '99 & PRB '01

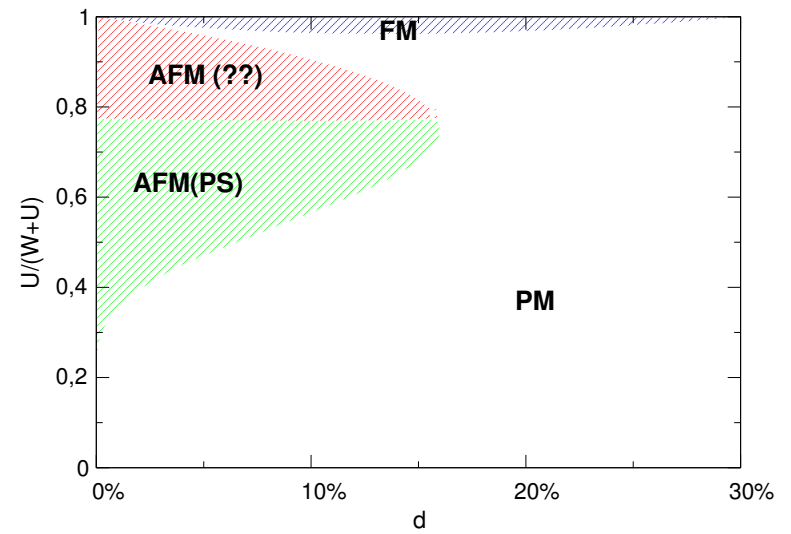




Key results from DMFT:

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Zitzler et al., EPJ '02

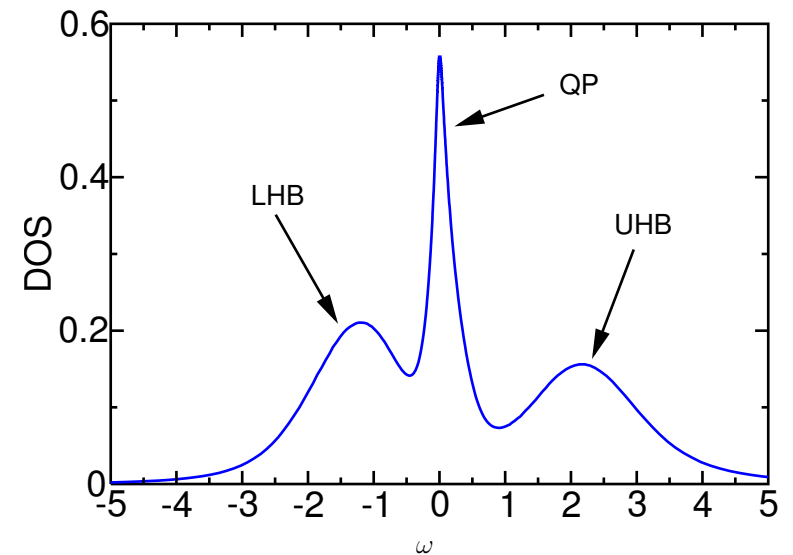




Key results from DMFT:

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TP et al., PRB '93



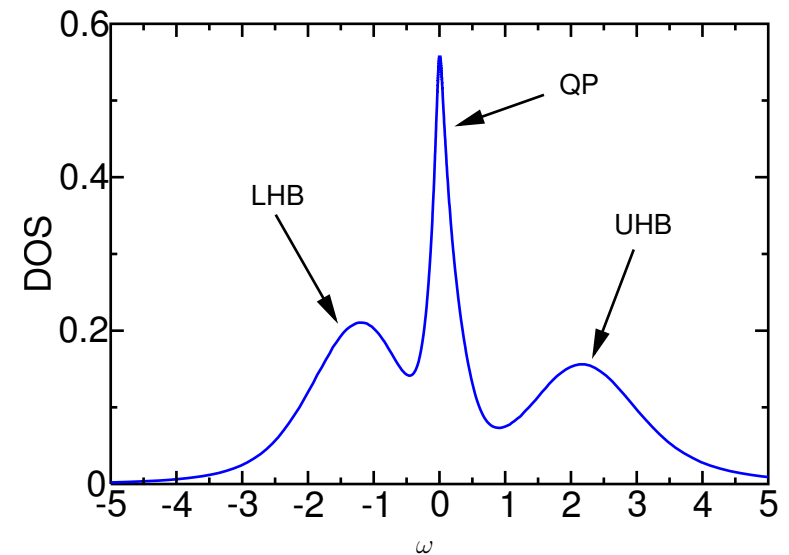


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- ✓ Non-trivial single-particle dynamics
- ✓ Ordered phases, two-particle prop.

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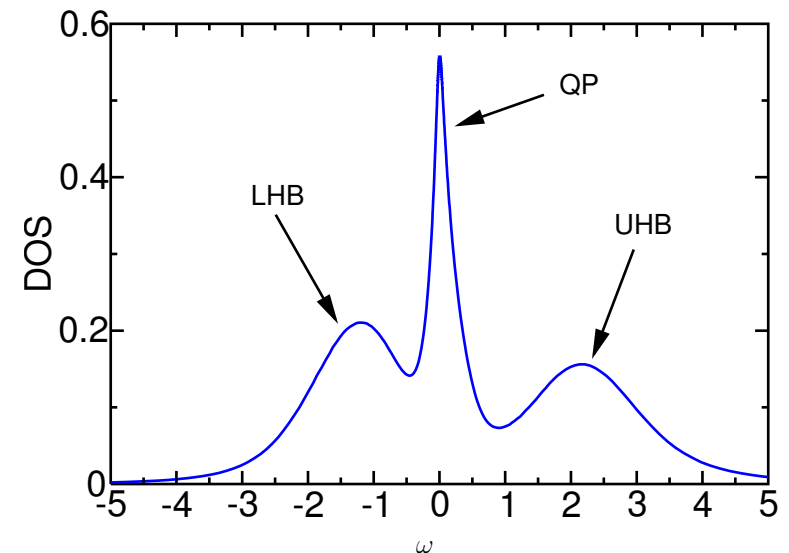




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- ✓ Thermodynamically consistent
- ✓ Non-trivial single-particle dynamics
- ✓ Ordered phases, two-particle prop.
- ✗ Only local dynamics & OPs
- ✗ Insensitiv to dimensionality
- ✗ Violation of Nernst's theorem



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## Where can and will DMFT go wrong?

Simple example: Old acquaintance Hubbard model in  $D = 2$ , at  $\langle n \rangle = 1$

Expected physics:

➡ **Ground state** antiferromagnetic insulator for all  $U > 0$

➡ Mermin-Wagner: **No long-range order** for  $T > 0$

➡ Strong order-parameter fluctuations at low  $T$

(“spin-waves”)



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Answer from DMFT:

- Antiferromagnetic insulator for all  $U > 0$  and  $T < T_N$ ,  $T_N(U) > 0$
- No order-parameter fluctuations as  $T \gtrsim T_N$  (“spin-waves”)

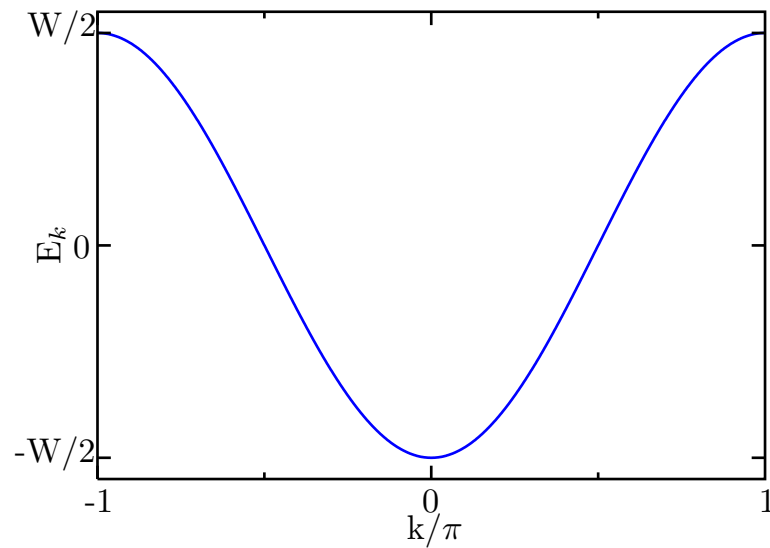




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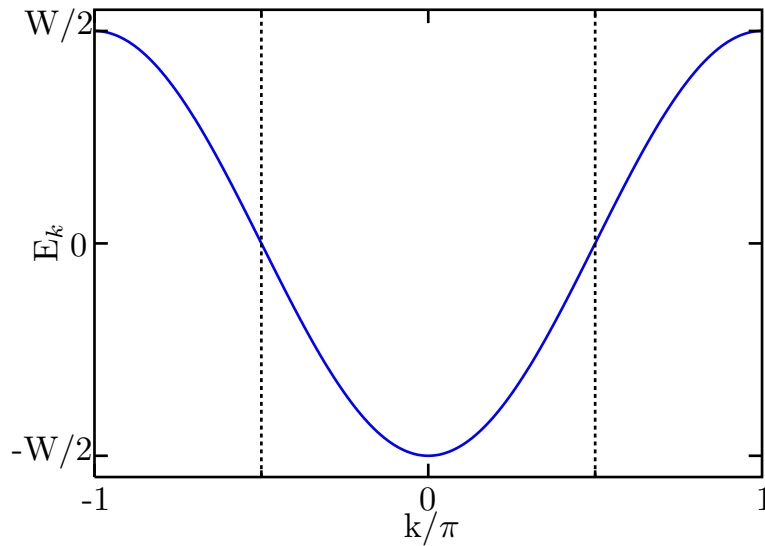


DMFT: Long-range AF order





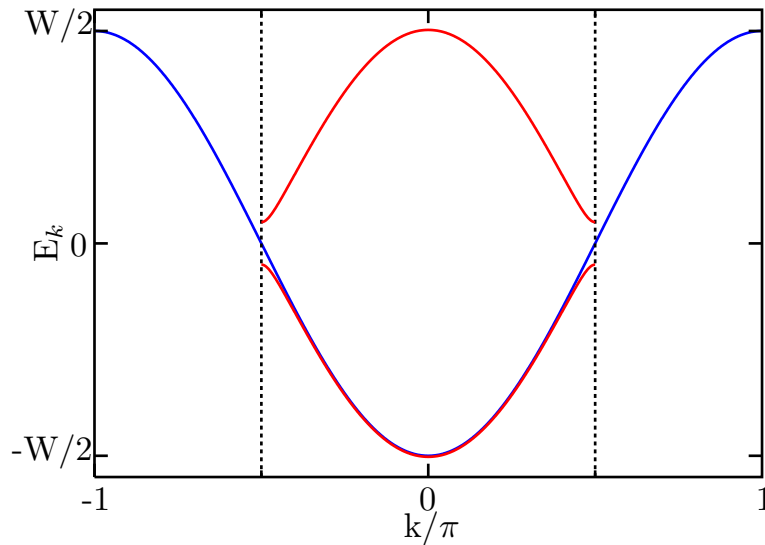
## DMFT: Long-range AF order



broken translational symmetry  
 $\Rightarrow$  reduced Brillouin zone



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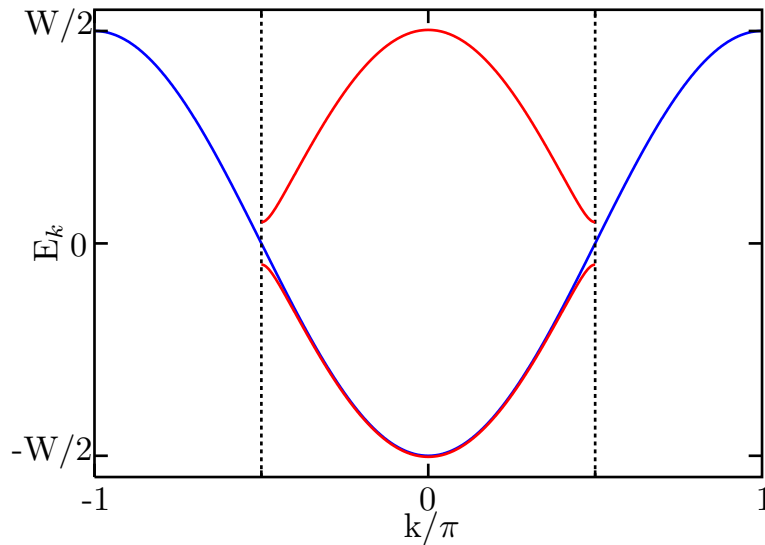


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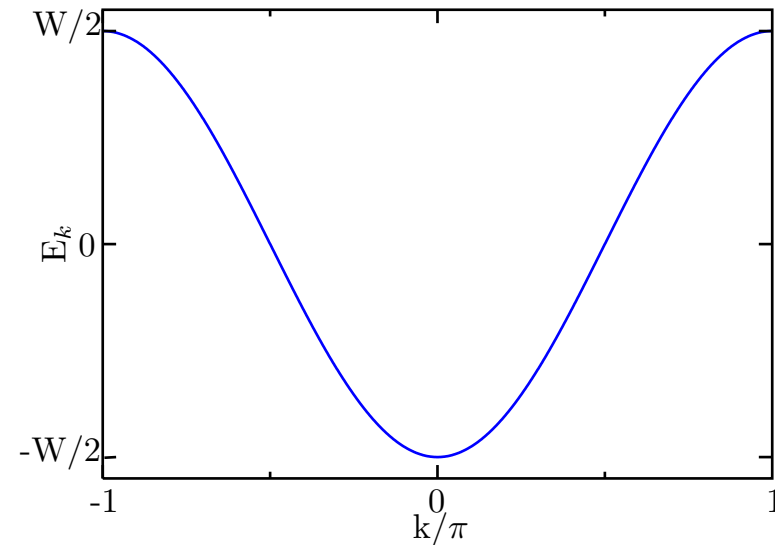
band structure folded back



DMFT: Long-range AF order



True system: OP fluctuations

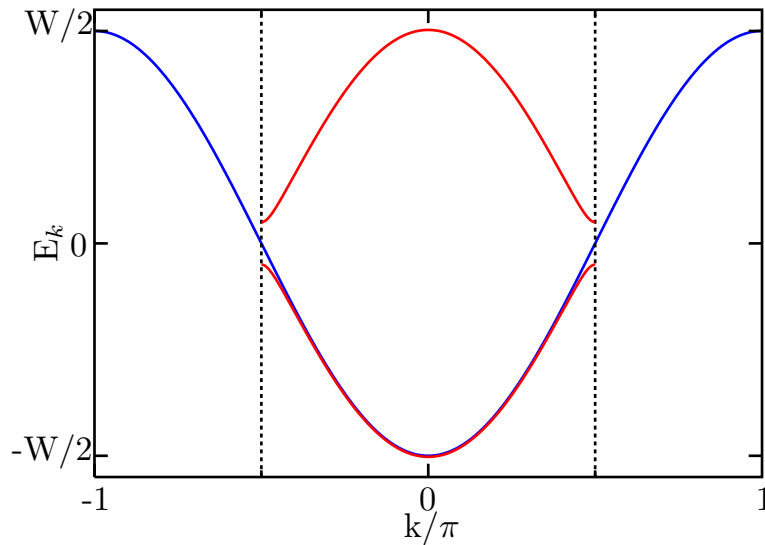


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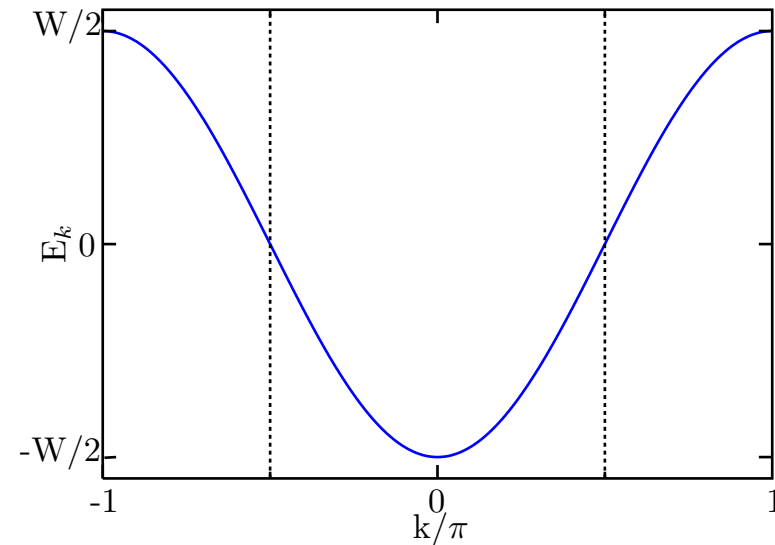


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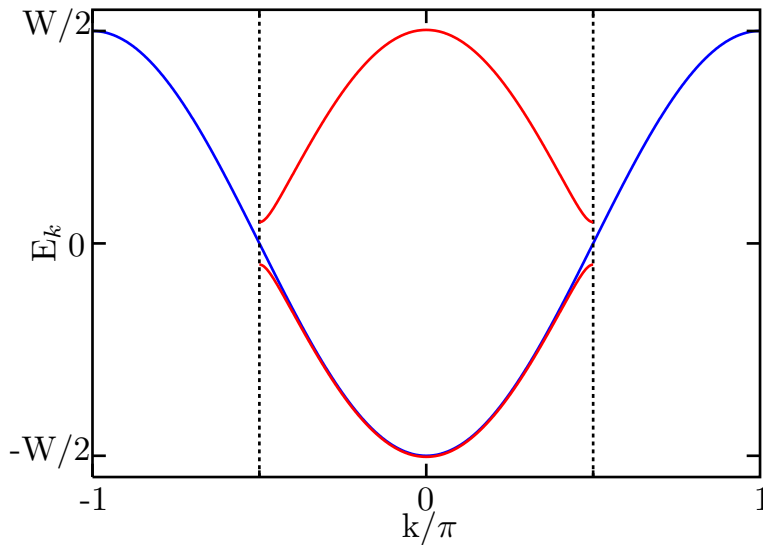
full translational symmetry  
⇒ normal band structure

But dynamical short-ranged order:

band structure folded back



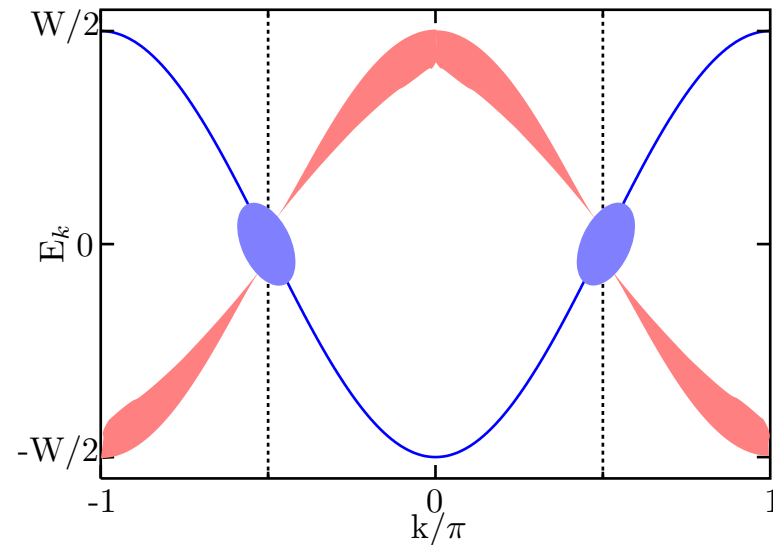
DMFT: Long-range AF order



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band structure folded back

True system: OP fluctuations



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But dynamical short-ranged order:

Shadow bands, pseudo gaps, ...



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How to incorporate non-local effects?





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How to incorporate non-local effects?

- Systematic  $1/d$  corrections to DMFT
  - ✗ Not feasible

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- DMFT + **local** two-particle dynamics

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✓ Pseudo-gap in DOS

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- DMFT + **phenomenological** two-particle dynamics

Nekrassov et al. '05

✓ Pseudo-gap in DOS

✓ True  $\vec{k}$ -dependence

✗ No true microscopic foundation



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More “natural” approach:



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More “natural” approach:

- Not too close to phase transitions
  - ☞ non-local fluctuations typically short-ranged
  - ☞ weak non-local contributions to self-energy



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Relax  $\Sigma(\vec{k}, z) \longrightarrow \Sigma(z)$  to

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with  $N_c > 1$  points  $\vec{K} \in 1. \text{ BZ}$



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Use “molecule” formed by

$N_c > 1$  sites  $\vec{R}_J$  for DMFT

$$\Sigma_{ij}(z) \longrightarrow \Sigma_{IJ}(z)$$





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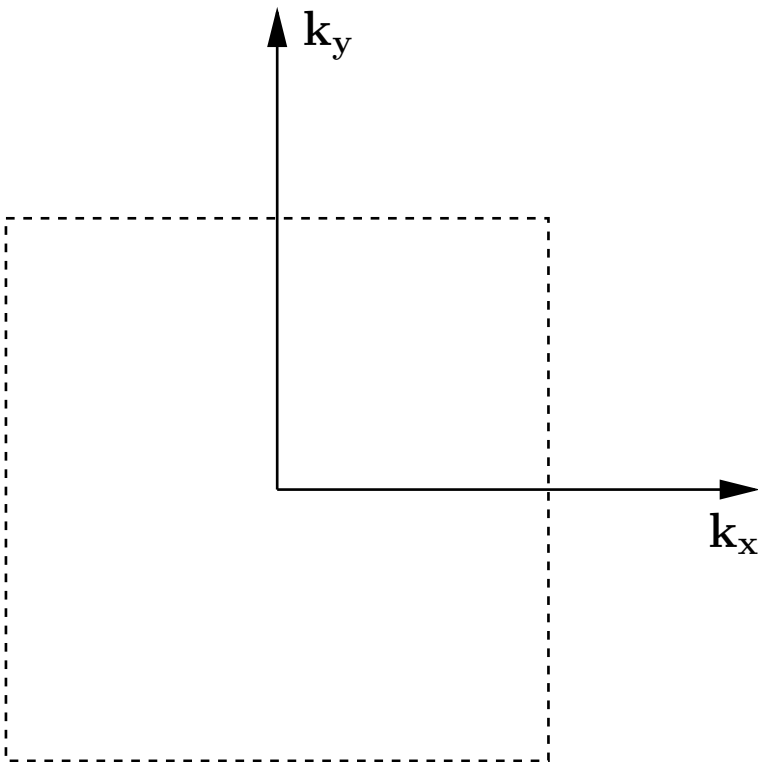
Realization of DCA:

- Reduce resolution in  $\vec{k}$ -space



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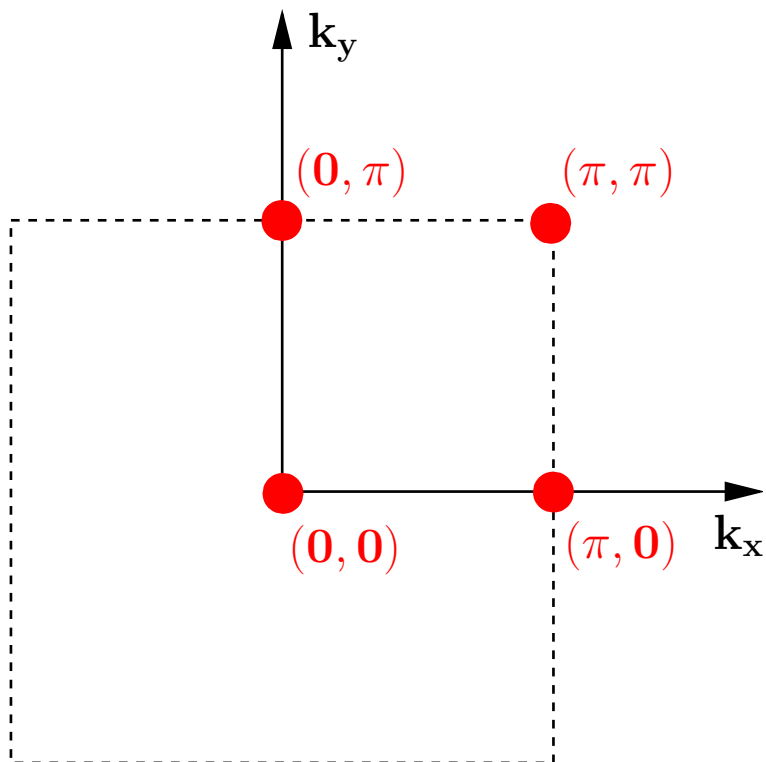
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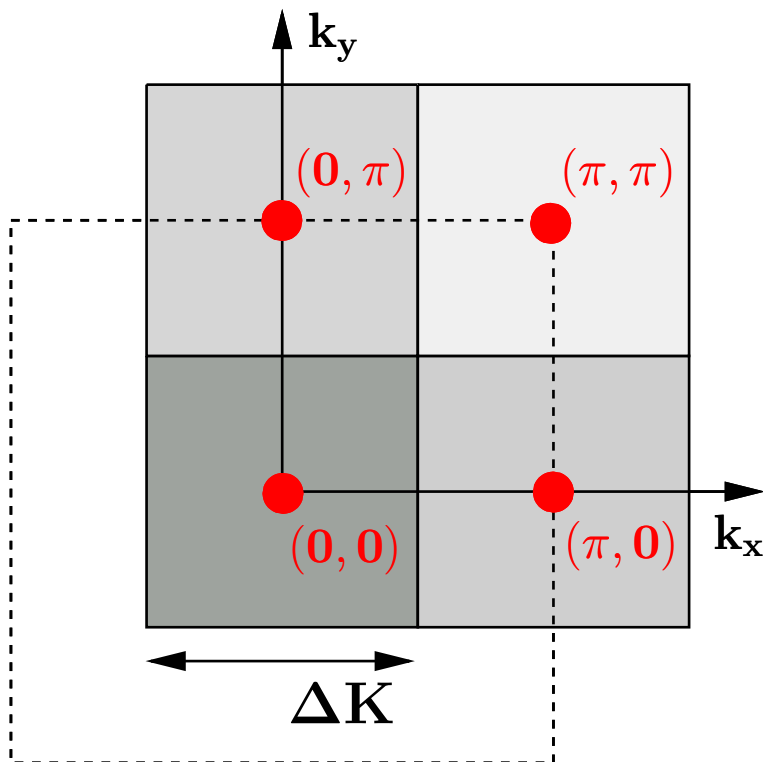


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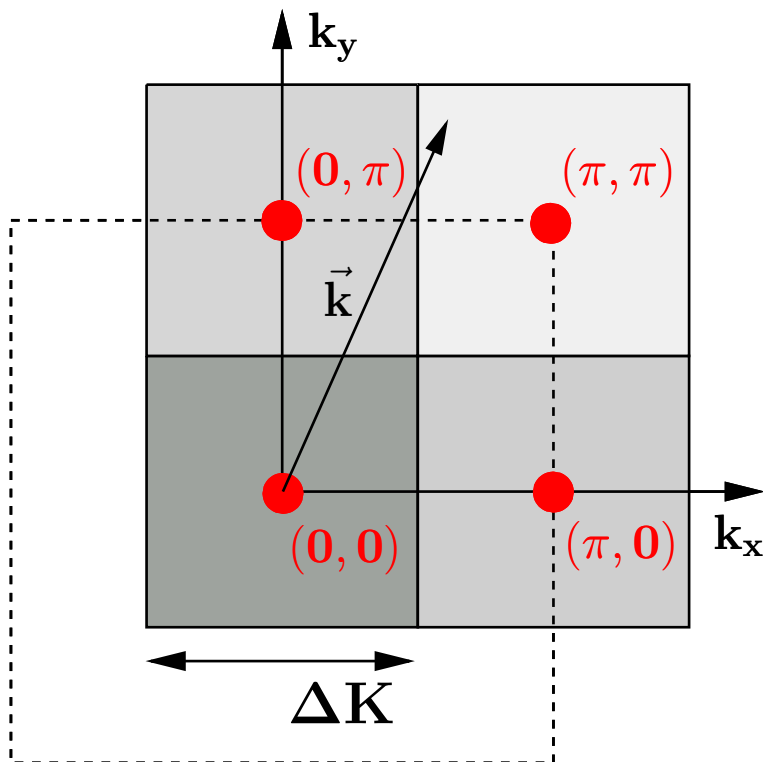
- ❑ Reduce resolution in  $\vec{k}$ -space



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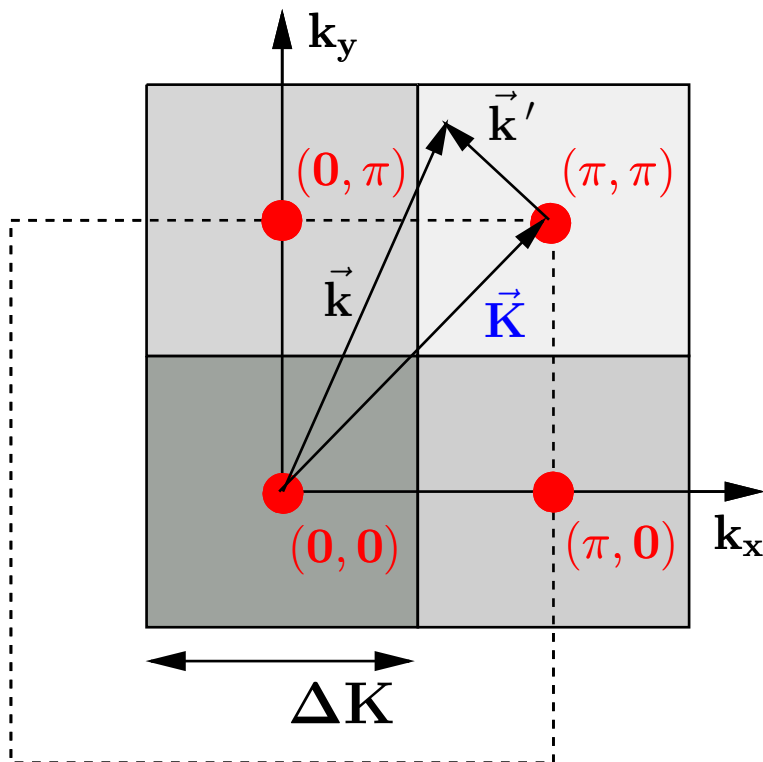
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$$\Sigma(\vec{k}, z)$$

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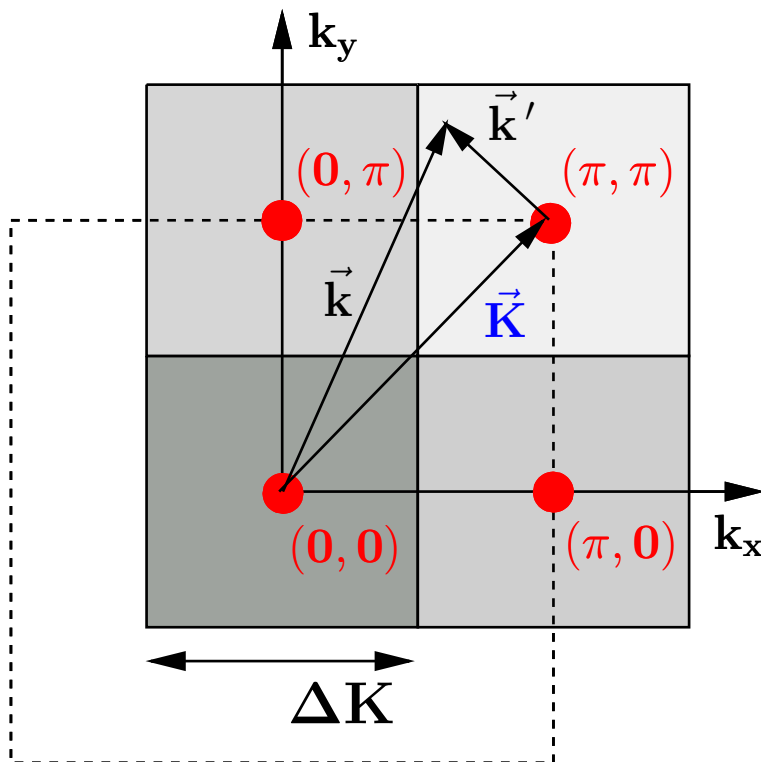
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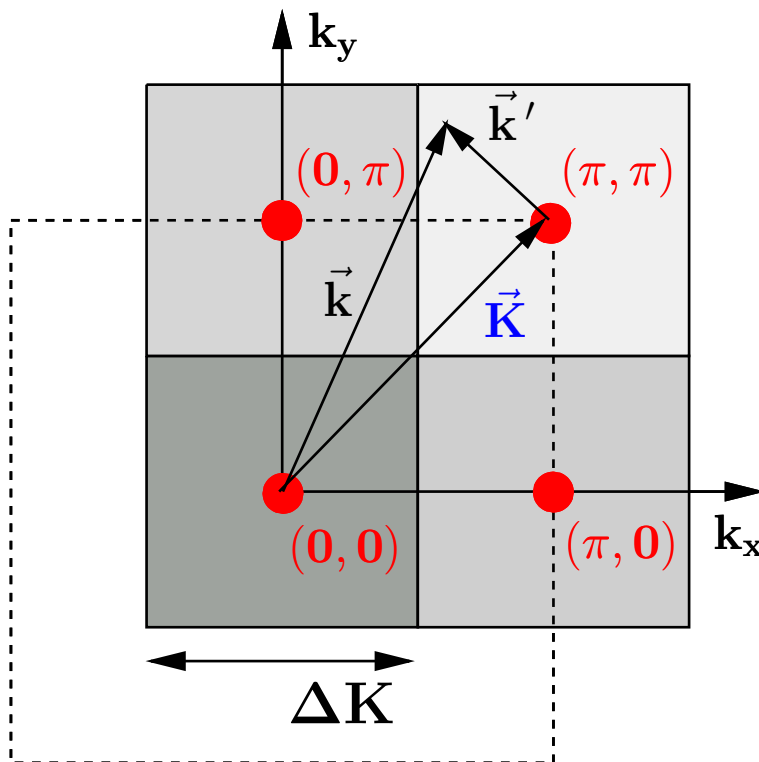
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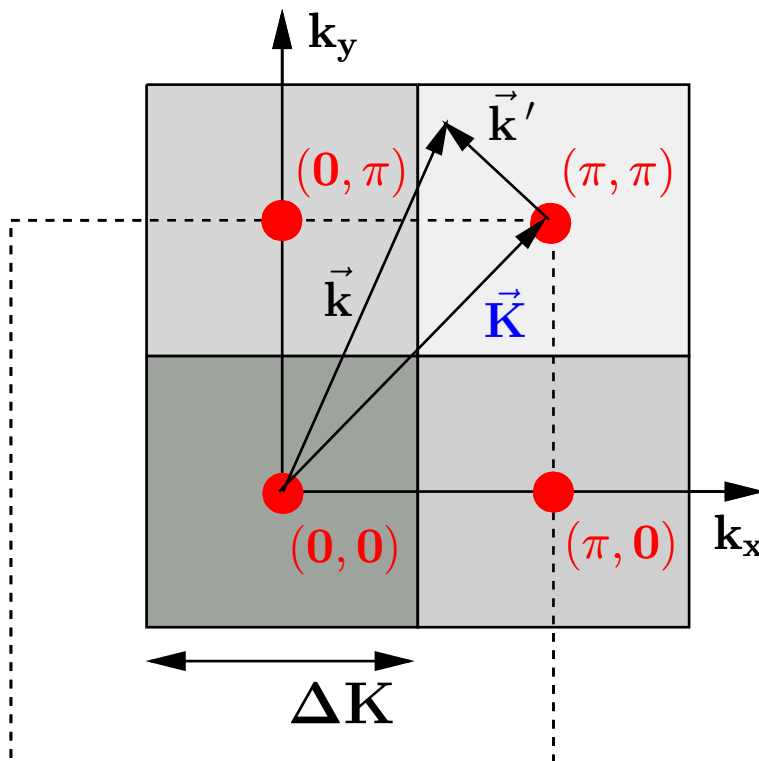
- Perform **coarse graining**:

$$\bar{G}(\vec{K}, z) = \frac{N_c}{N} \sum_{\vec{k}'} \frac{1}{z + \mu - \epsilon_{\vec{K} + \vec{k}'} - \Sigma(\vec{K}, z)}$$



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- Reduce resolution in  $\vec{k}$ -space



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- effective model on periodic cluster**



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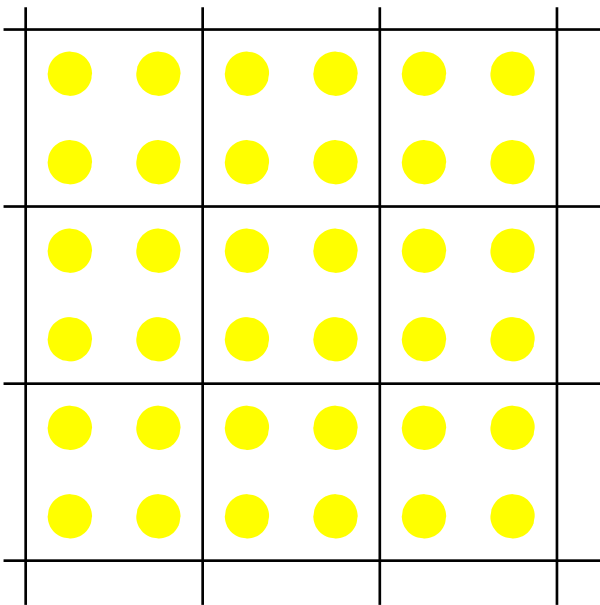
Realization of CDMFT:

- Build clusters in real space



## Realization of CDMFT:

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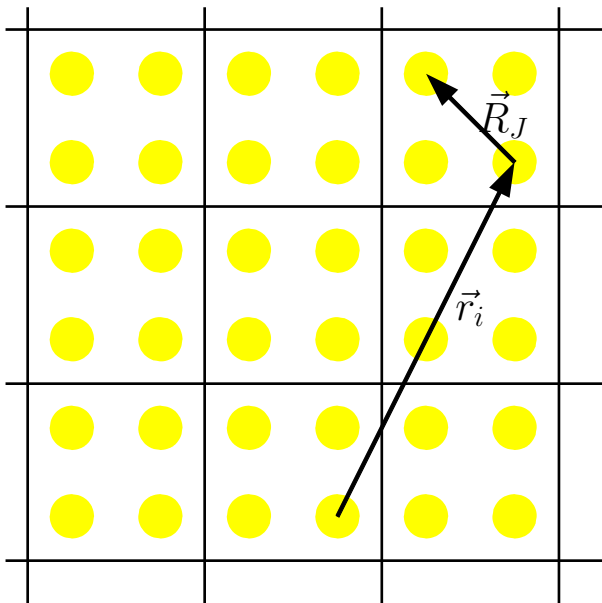


- Choose cluster with  $N_c \times N_c$  points  $\vec{R}_J$



## Realization of CDMFT:

- Build clusters in real space



- Choose **cluster with  $N_c \times N_c$  points  $\vec{R}_J$**

- Break up

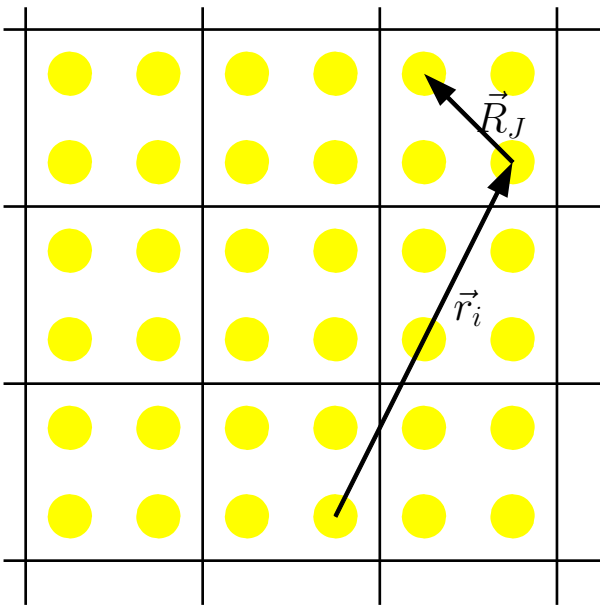
$$[\mathbf{t}]_{ij} = [\mathbf{t}]_{IJ} + [\delta\mathbf{t}](\vec{r}_i - \vec{r}_j)$$

$$[\Sigma(z)]_{ij} \approx [\Sigma(z)]_{IJ}$$



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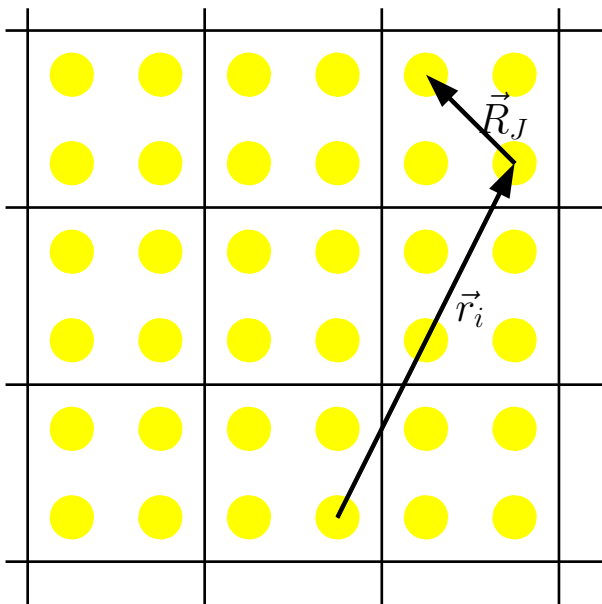
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- **Effective model on open-boundary cluster**



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Implementing the DCA algorithm:

similar for CDMFT with matrices

First guess for  $\Sigma(\mathbf{K})$



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First guess for  $\Sigma(\mathbf{K})$

$$\bar{G}(\mathbf{K}) = \frac{N_c}{N} \sum_{\mathbf{k}'} \frac{1}{\omega - \epsilon_{\mathbf{K}+\mathbf{k}'} + \mu - \Sigma(\mathbf{K})}$$





## Implementing the DCA algorithm:

similar for CDMFT with matrices

$$\mathcal{G}^{-1}(\mathbf{K}) = \bar{G}^{-1}(\mathbf{K}) + \Sigma(\mathbf{K})$$

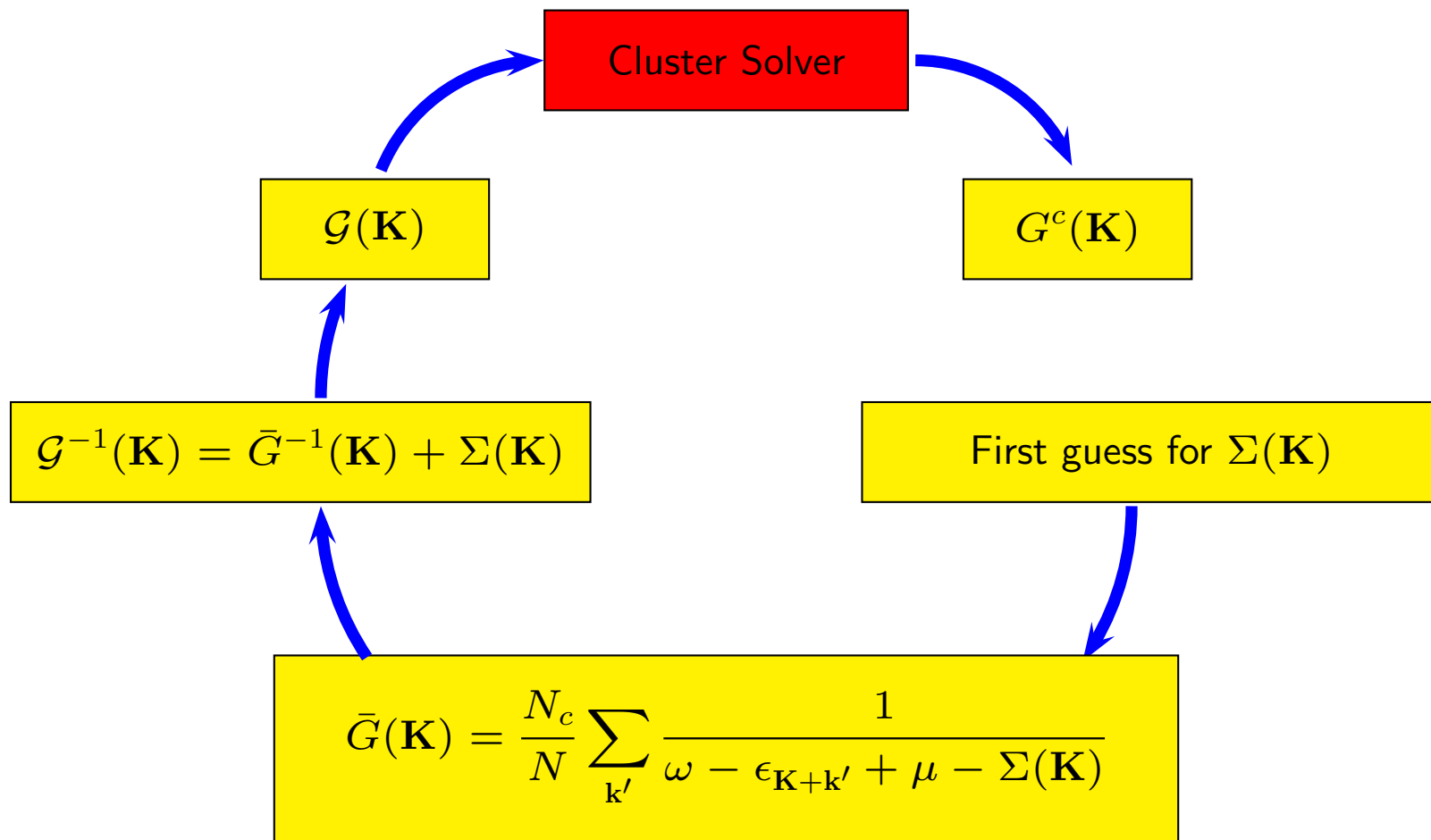
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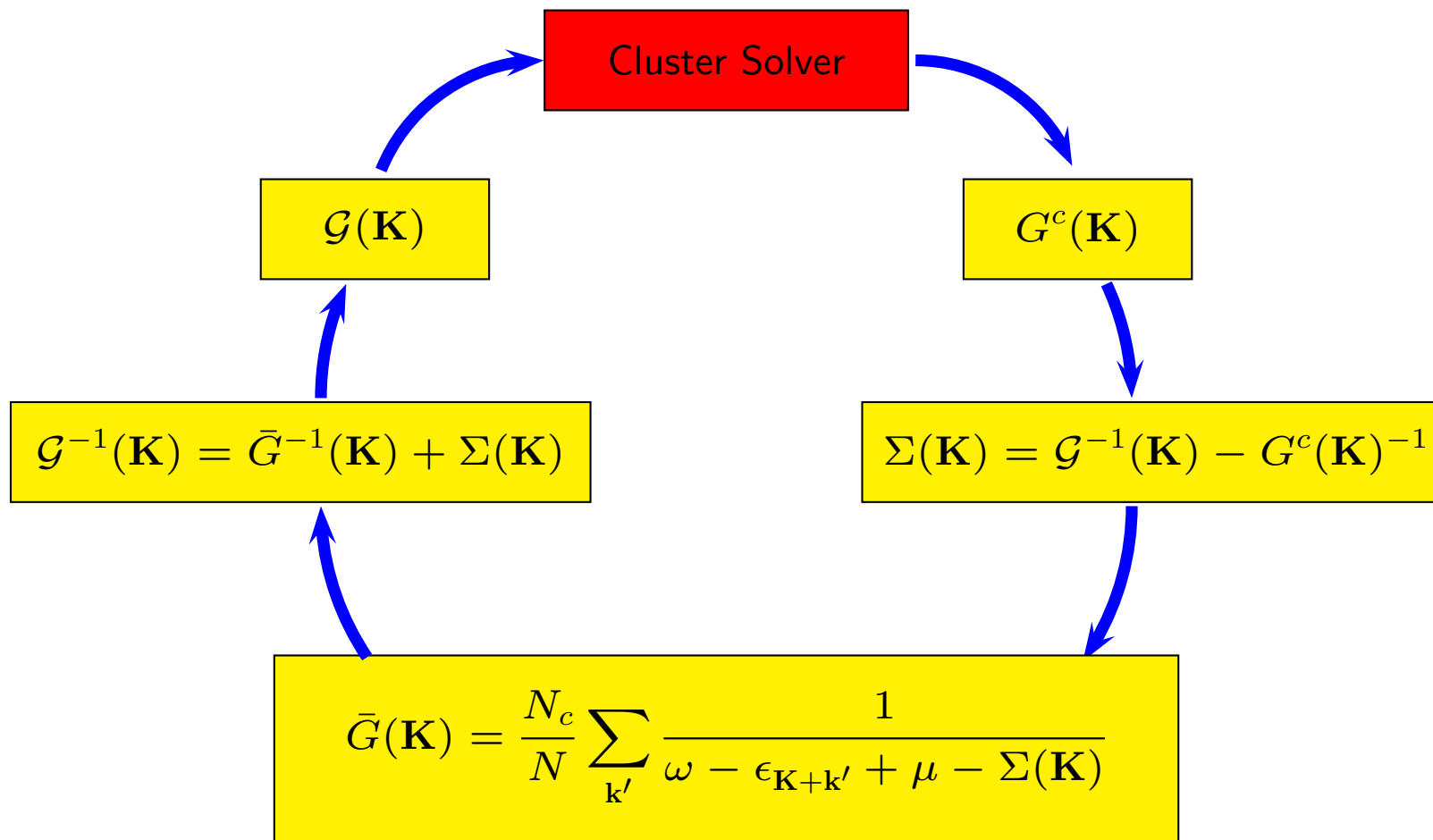
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Implementing the DCA algorithm:

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DCA and CDMFT seem to be nice schemes



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DCA and CDMFT seem to be nice schemes, but:

Structure of effective cluster



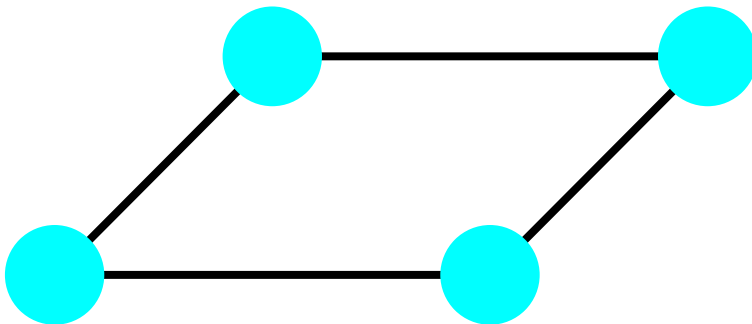
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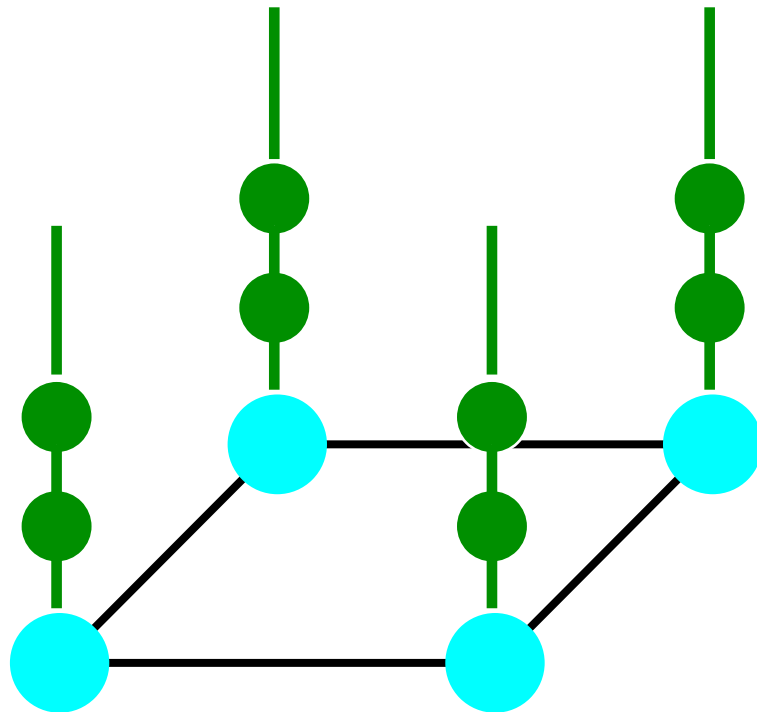
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- $N_c$  sites as in ED/QMC
- Well-developed techniques?





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### Structure of effective cluster

- ➡  $N_c$  sites as in ED/QMC
- ➡ Well-developed techniques?
- ➡ “dynamical mean-field
- ➡  $O(N_c)$  semi-infinite chains
- ➡ Only Hirsch-Fye QMC left



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Another potential problem:

- DCA gives  $\Sigma(\vec{k}, z)$  only as “histogram” on the  $\vec{K}$  values
  - Additional interpolation necessary
  - Reliability close to  $\vec{k}_F$  and  $\mu$ ?





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- CDMFT breaks translational invariance
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- CDMFT breaks translational invariance
  - ☞ Additional “averaging” necessary to obtain  $\Sigma(\vec{k}, z)$
  - ☞ Reliability close to  $\vec{k}_F$  and  $\mu$ ?

General expectation:

“Smoothing” well-defined for smoothly  $\vec{k}$ -dependent  $\Sigma(\vec{k}, z)$



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# Cluster Extensions to the DMFT

1. Motivation: Why need a cluster MFT?
2. Realizations: DCA and CDMFT
- 3. Selected results for DCA**
4. Summary

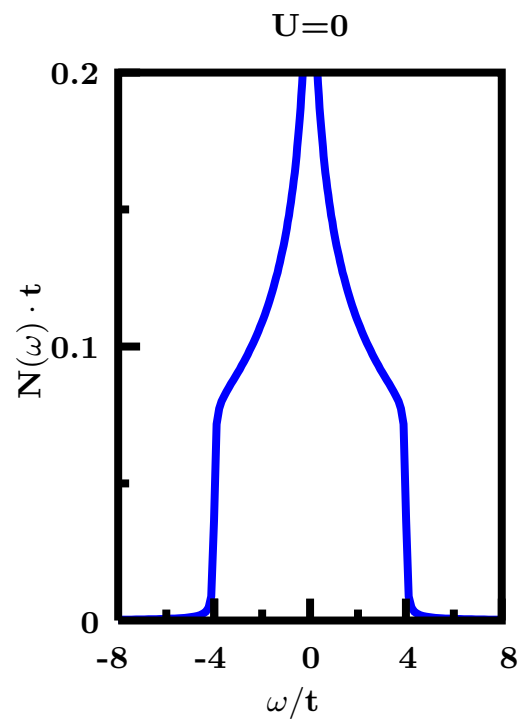


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Maier, TP *et al.* [EPJ B 13](#) ('00)

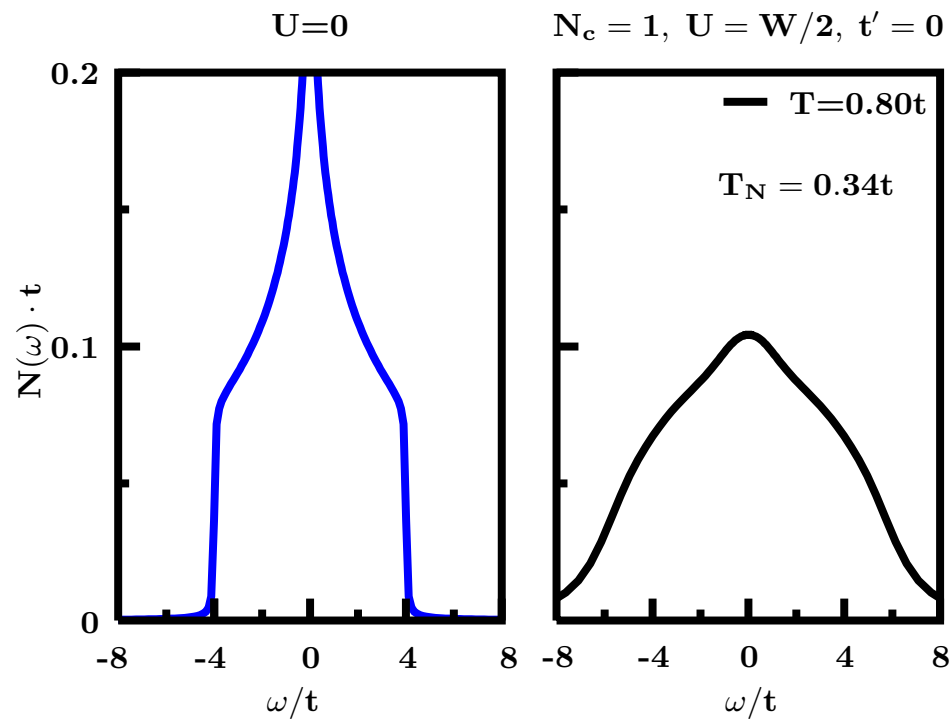


2D Hubbard model:  $N_c = 1$  vs.  $N_c > 1$



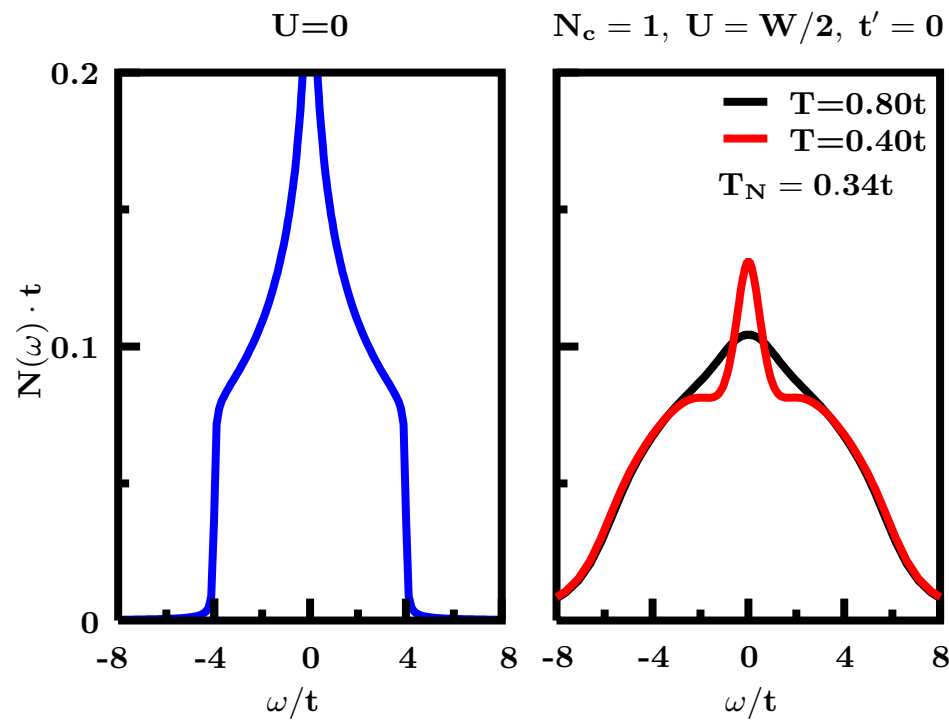


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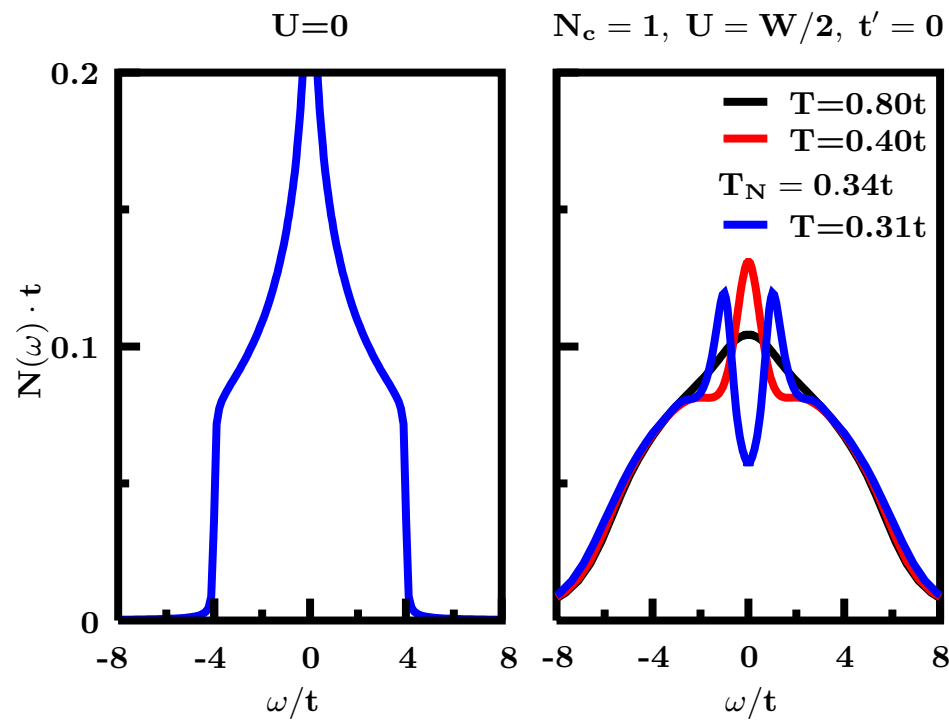


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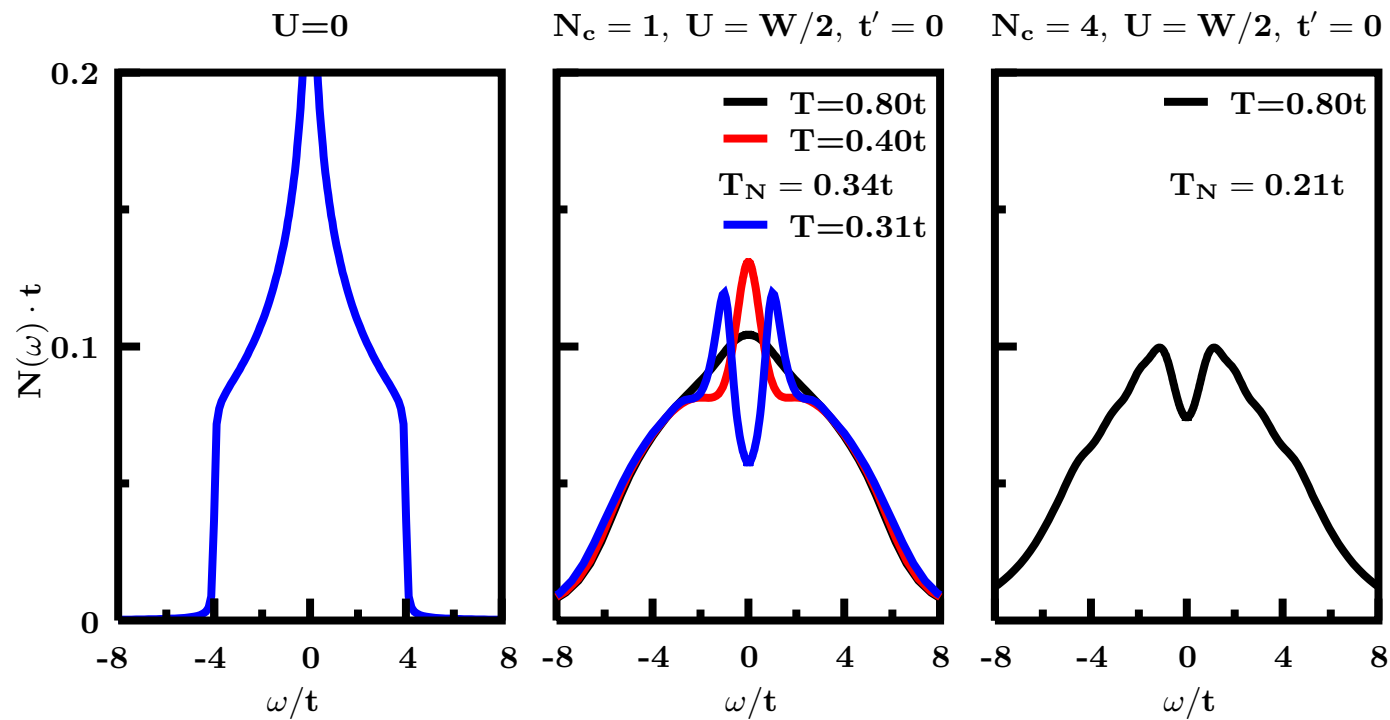


$N_c = 1$ :

➔ No precursor of AF



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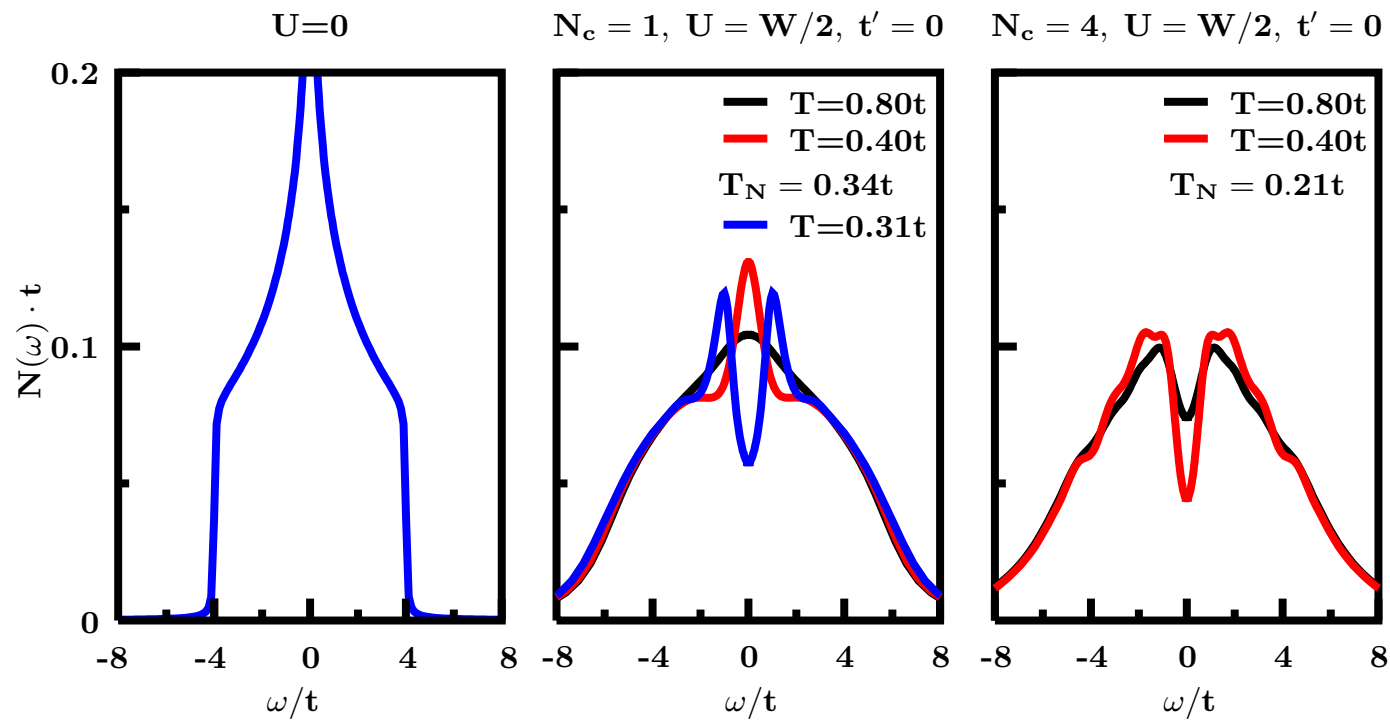
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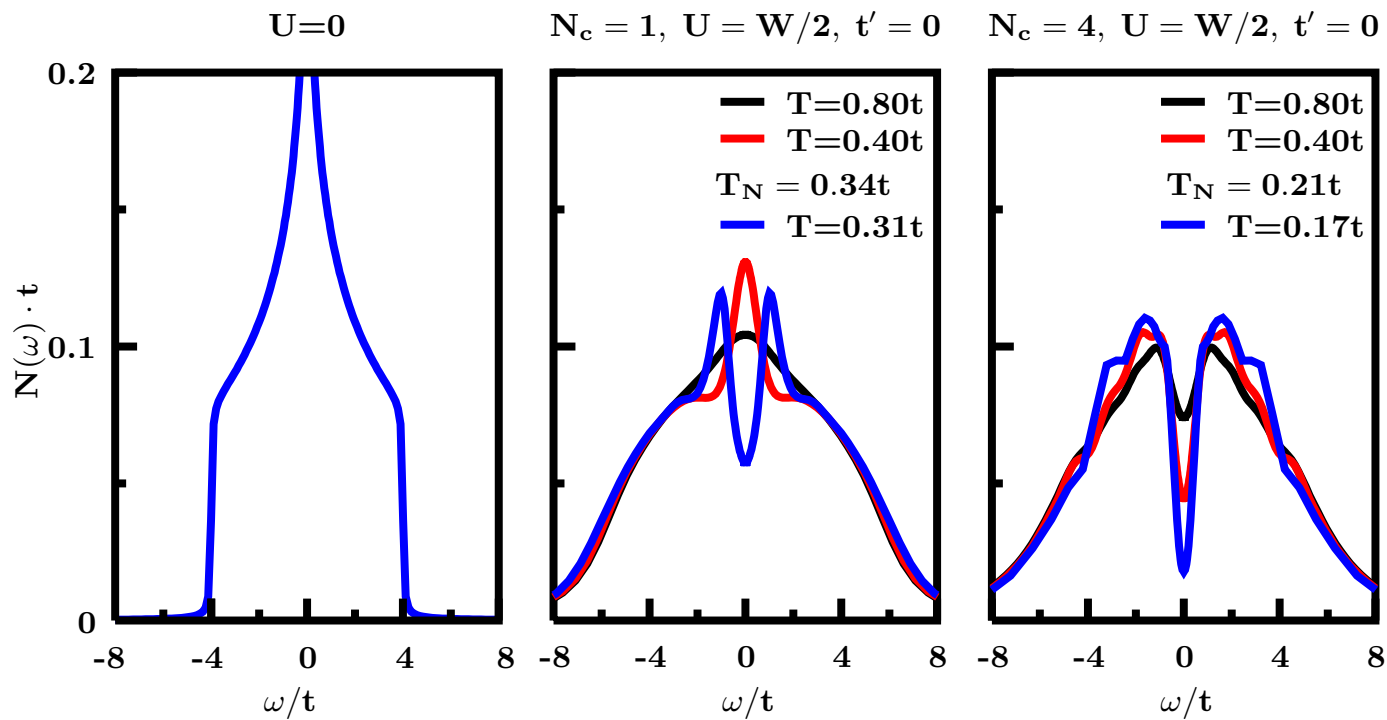
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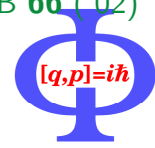
$N_c = 4$ :

➔ Pseudo gap in paramagnet

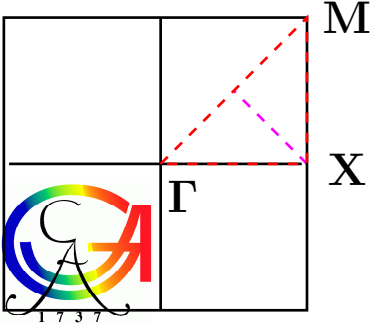


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Maier, TP *et al.*, PRB **66** ('02)

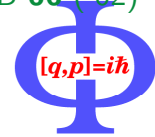


$$N_c = 16, U = W = 8t, t' = -0.2t, t = 0.25\text{eV}$$



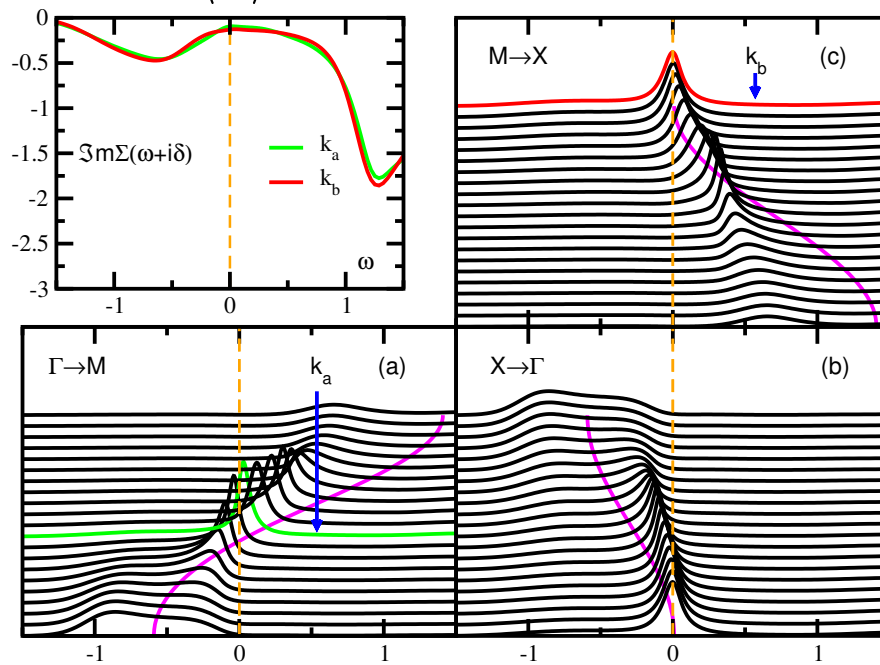
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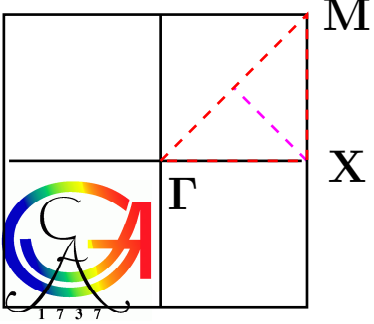


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$$\langle n \rangle = 0.80, T = 370\text{K}$$

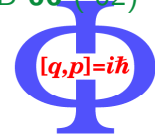


- ❑ Well defined quasi particles
- ❑ Weak  $\vec{k}$ -dependence of  $\Sigma(\vec{k}, \omega)$



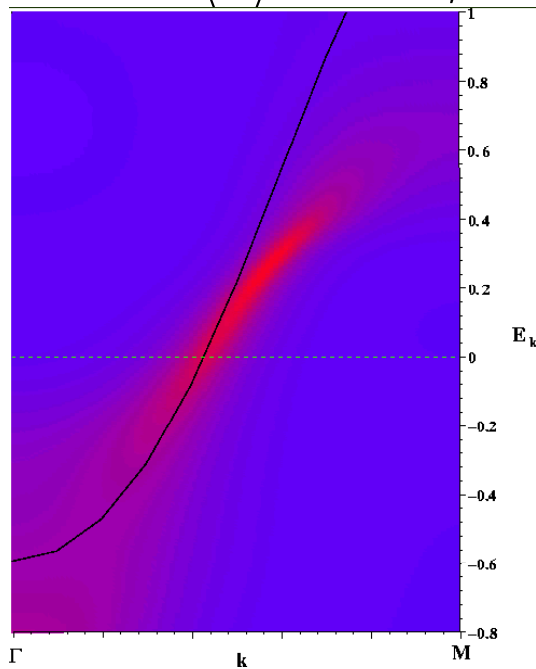
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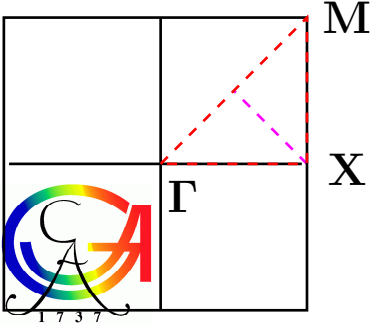


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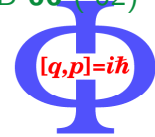


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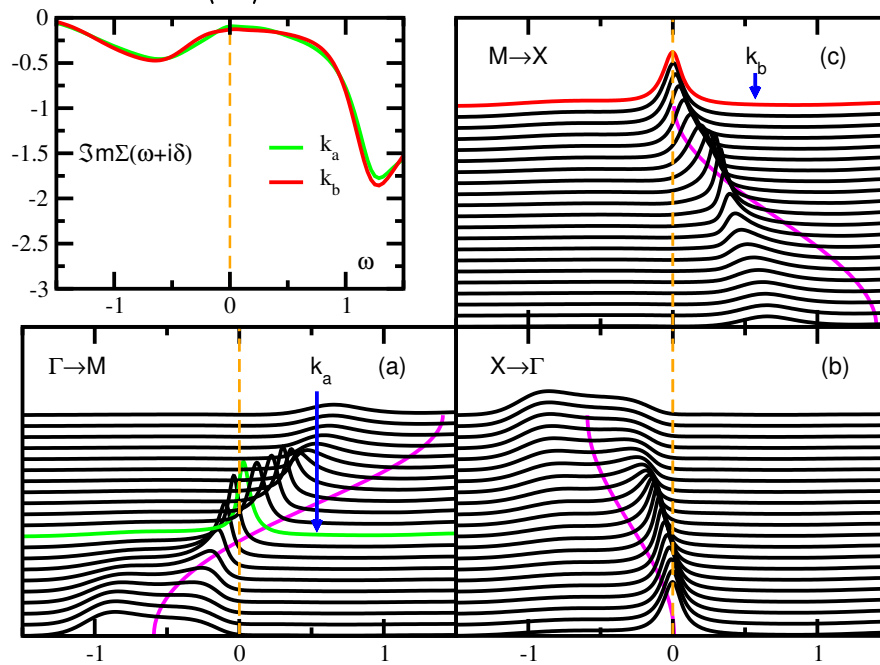
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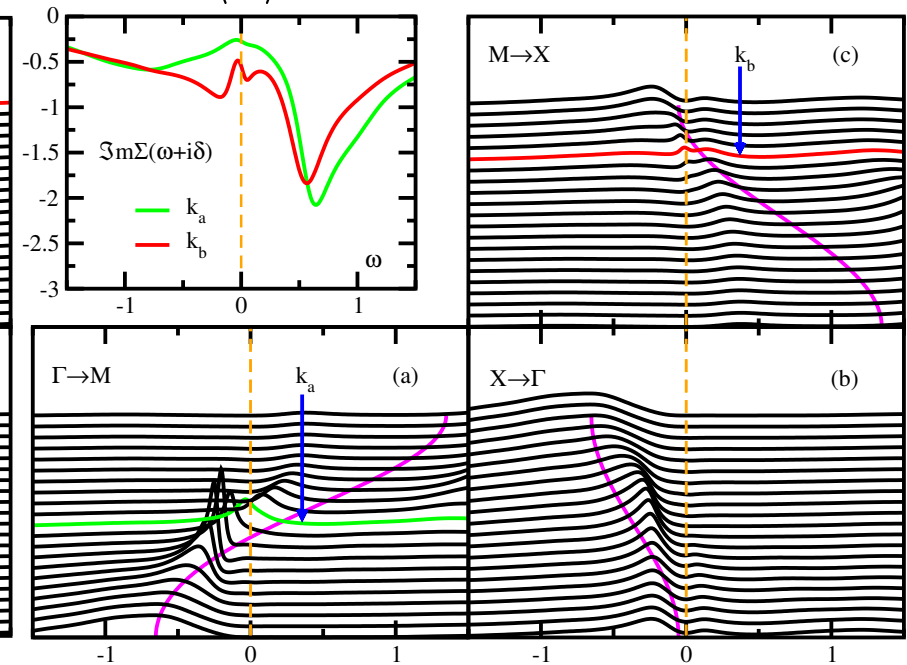


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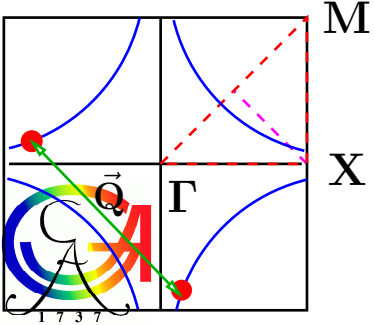


$$\langle n \rangle = 0.95, T = 370K$$



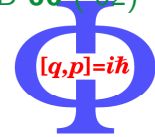
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- Strongly damped structures at X
- Strong  $\vec{k}$ -dependence of  $\Sigma(\vec{k}, \omega)$



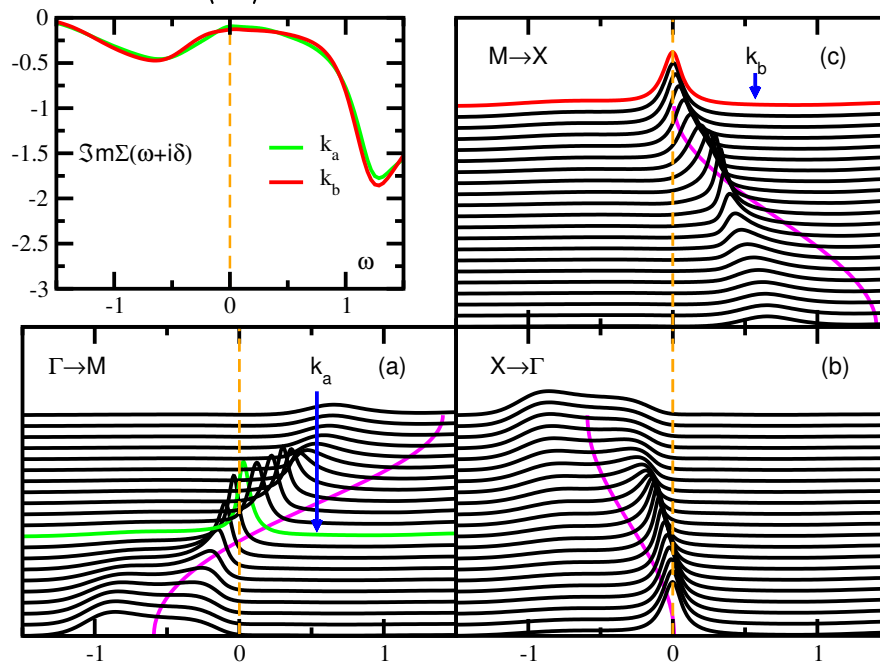
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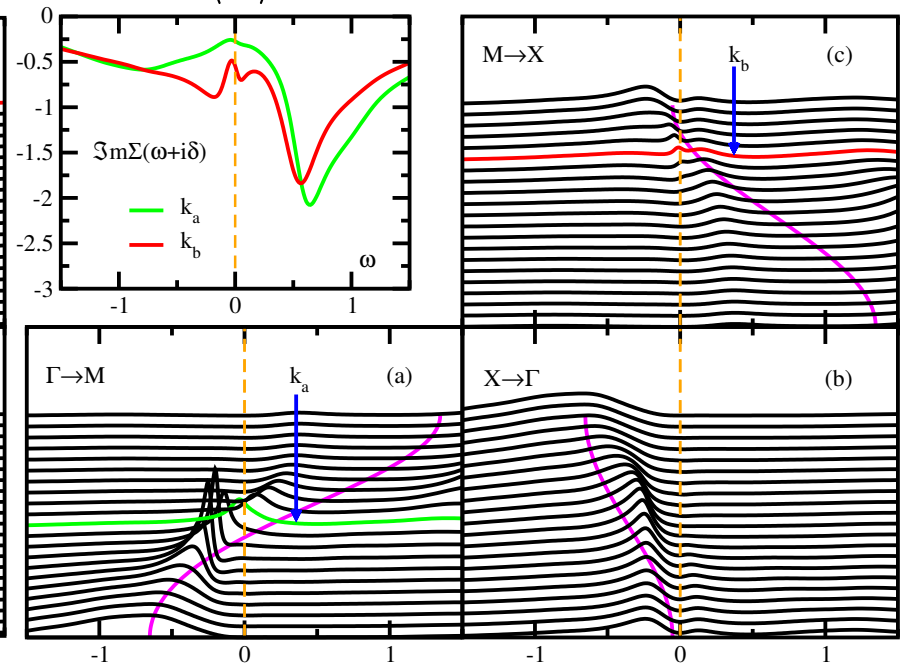


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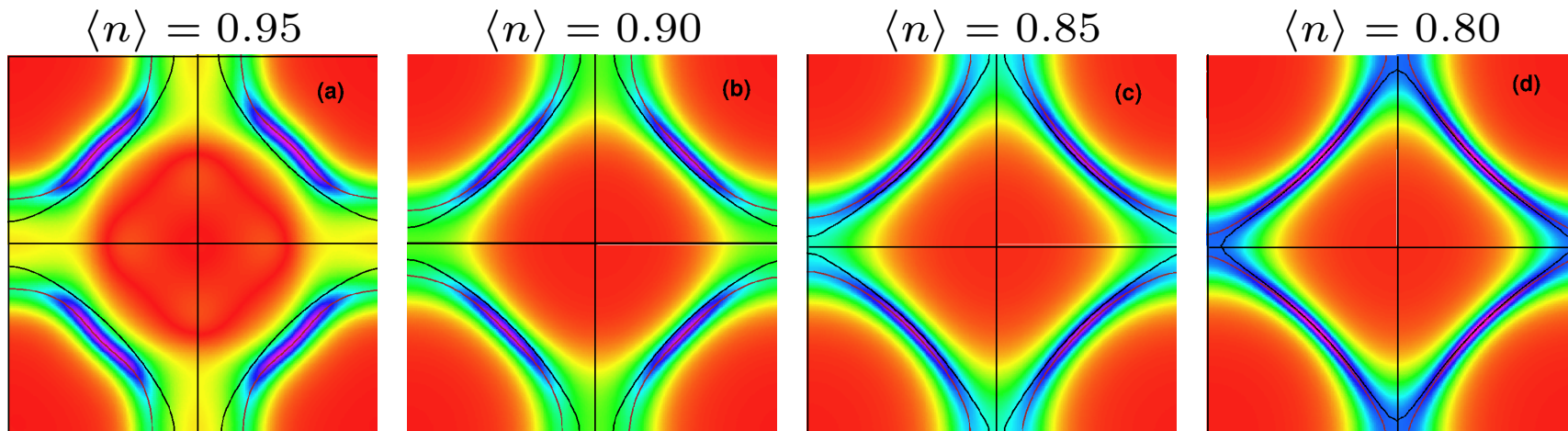
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$$N_c = 16, U = W, t' = -0.2, T = 370\text{K}$$



$$\langle n \rangle \gtrsim 0.9$$

Small FS for  $\langle n \rangle \rightarrow 1$   
 ➔ Luttinger's theorem violated

$$\langle n \rangle \lesssim 0.9$$

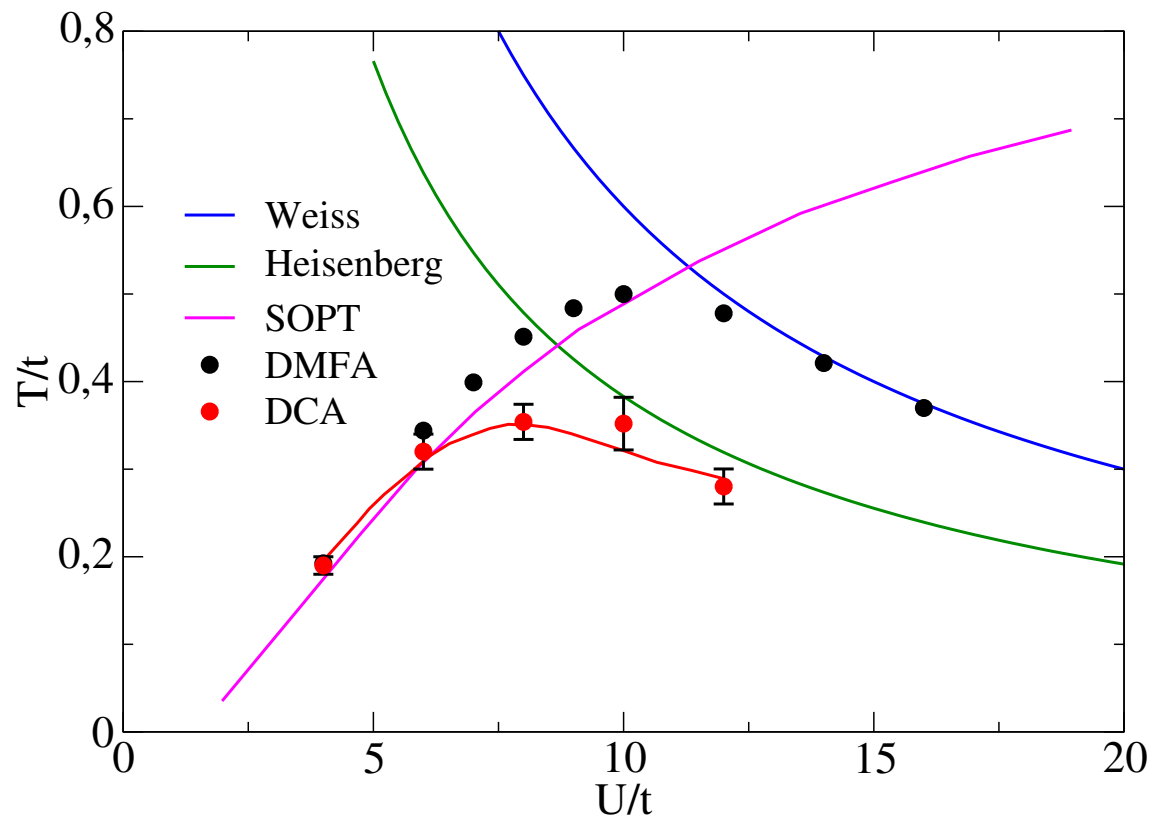
Large FS for  $\langle n \rangle < 0.9$   
 ➔ Luttinger's theorem fulfilled





How does it perform for  $D = 3$ ?

Kent et al. '05





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- Cluster extensions of DMFT:
  - ☞ Systematic inclusion of non-local correlations
  - ☞ Thermodynamic limit
  - ☞ Short-ranged correlations treated exactly
  - ☞ Long-ranged correlations mean-field like
  - ☞ Sensible results with acceptable computational effort

Further reading:

Th. Maier et al., [cond-mat/0404055](#) (RMP in press)