



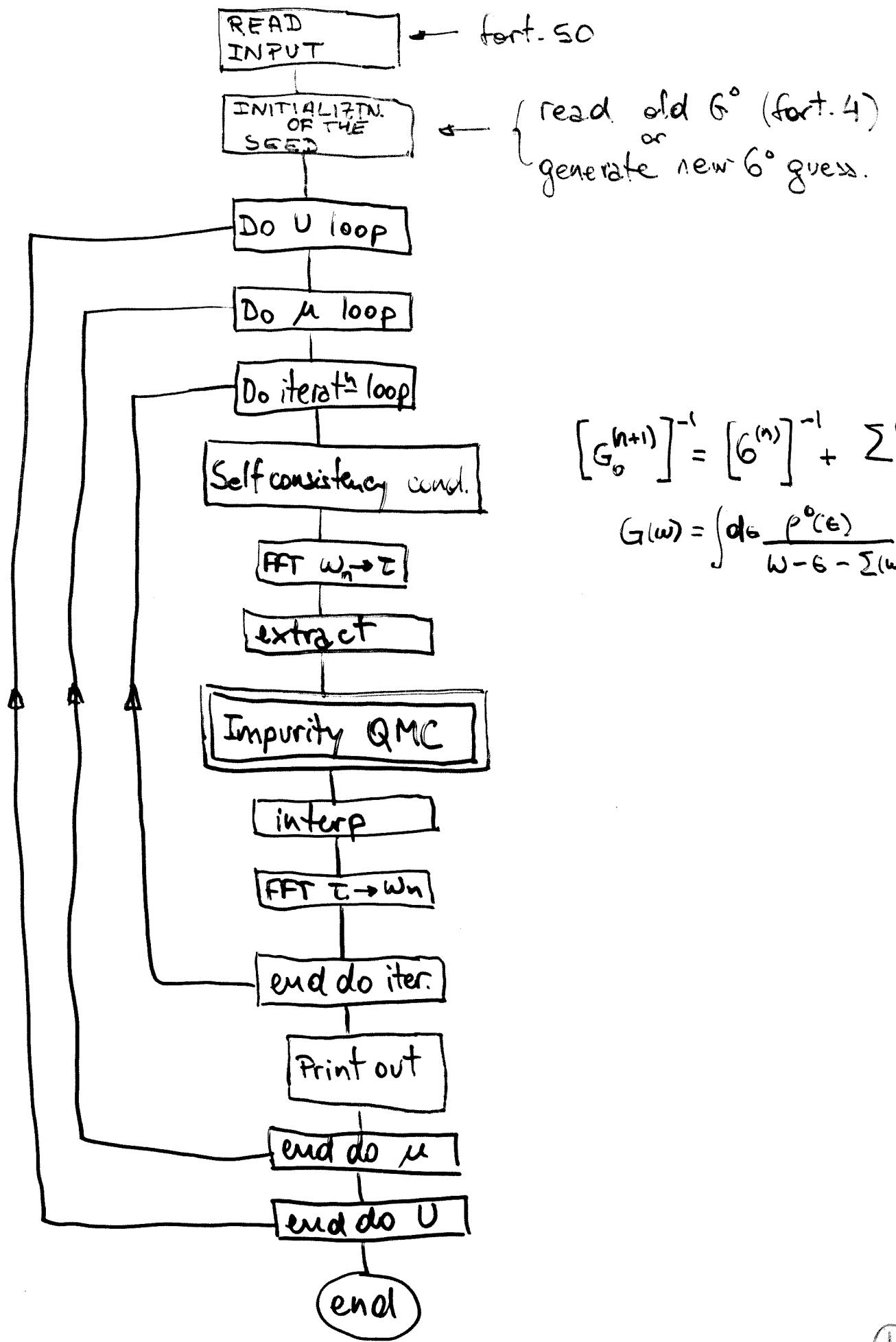
SMR.1667 - 2

Summer School and Miniconference on
**Dynamical Mean-Field Theory for Correlated Electrons:
Applications to Real Materials, Extensions and Perspectives**
25 July - 3 August, 2005

**Metal-insulator transitions
in the Hubbard model**

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QMC code - Flow diagram



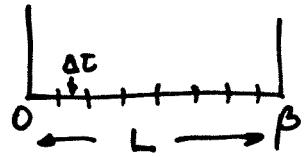
$$[G_0^{(n+1)}]^{-1} = [G^{(n)}]^{-1} + \sum^{(n)}$$

$$G(w) = \frac{d\epsilon}{d\omega} \frac{\rho^0(\epsilon)}{w - \epsilon - \Sigma(w)}$$

QMC

(Blazekanbecker, Scalapino, Sugar '81)
(Hirsch & Fye '87)

$$Z = e^{-C_\sigma^+(\tau) G_0^{-1}(\tau) C_\sigma(\tau') + U M_+(\tau) M_+(\tau)}$$



$$\stackrel{t_0 \rightarrow}{Z} Z = \sum_{S=\pm 1} e^{-C_\sigma^+(\tau) G_0^{-1}(\tau) C_\sigma(\tau') + \lambda S(\tau) (M_+(\tau) - M_-(\tau))}$$

$\lambda = \lambda(U)$

$$Z = \sum_{\{s\}} \det [G_+^{-1}(s)] \det [G_-^{-1}(s)]$$

$$G_{\frac{d}{ds}}^{-1} = G_0^{-1} + \sigma \lambda s \delta_{\tau\tau'}$$

$$G_\sigma = \frac{1}{Z} \sum_{\{s\}} G_\sigma(s) \det [G_+^{-1} G_-^{-1}]$$

$$\langle G_\sigma \rangle = \sum_s G_\sigma(s) P(s) ; R = e^{-\Delta s} = \frac{\det [G_+(s) G_+(s)]}{\det [G_+(s') G_+(s')]}$$

* exact up to the Trotter break-up error $\propto \Delta\tau^2$

rule of thumb $\rightarrow \frac{\Delta\tau U}{2} < 1$

$$\beta = \Delta\tau L \approx \frac{1}{2} 128 = 64$$

$$T_{\text{QMC}} \gtrsim 10^{-2}$$

$\sim 1/2$ hour
~~compute time~~ (93)
Section III (99)

No tricks in the "impurity" routine (QMC
(Fye & Hirsch))

- * autocorrelation ~ 2 so store every other sweep
- * warm up (thermalization of the $S(\tau)$ pseudo spins) ~ 500 sweeps
- * "dirty" vs "clean" inversion

$$G_{\tau\tau'}^{-1} = G_{0\tau\tau'}^{-1} + \lambda S(\tau) \delta_{\tau\tau'}$$

need to store $G_{\tau\tau}$ (invert matrix)

"dirty" \rightarrow Sherman - Morrison formula L^2 operations

"clean" \rightarrow direct inversion (Gauss Jordan) L^3 operations

bottle neck ?

~ 500 dirty sweeps every other clean
is OK in general.

* "wrap around"

$$G_{\xi\tau}^i \quad i = 1, \dots, N_{\text{store}}$$

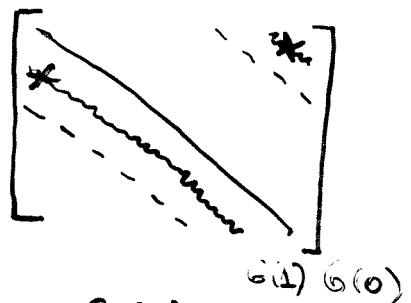
$$G(\tau) = \frac{1}{L \times N_{\text{store}}} \sum_{i=1}^{N_{\text{store}}} \sum_{\tau'}^L G_{\tau+\tau', \tau'}^i$$

$$G(-\tau) = -G(\beta - \tau)$$

$$\tau' + \tau = \text{mod}_L'(\tau' + \tau)$$

$$1, L = 1-L \rightarrow -G(1-\tau+\tau) = -G(1)$$

$$2, 1 = 2-1 \rightarrow G(1)$$



Very nice but... * $T \rightarrow 0$ is numerically expensive
 * imaginary axis results.

\Rightarrow need to analytically continue ^{*noisy*} data to the real axis. It's a mathematically ill-posed problem. Something can be done using Maximum-Entropy methods.

Tricks

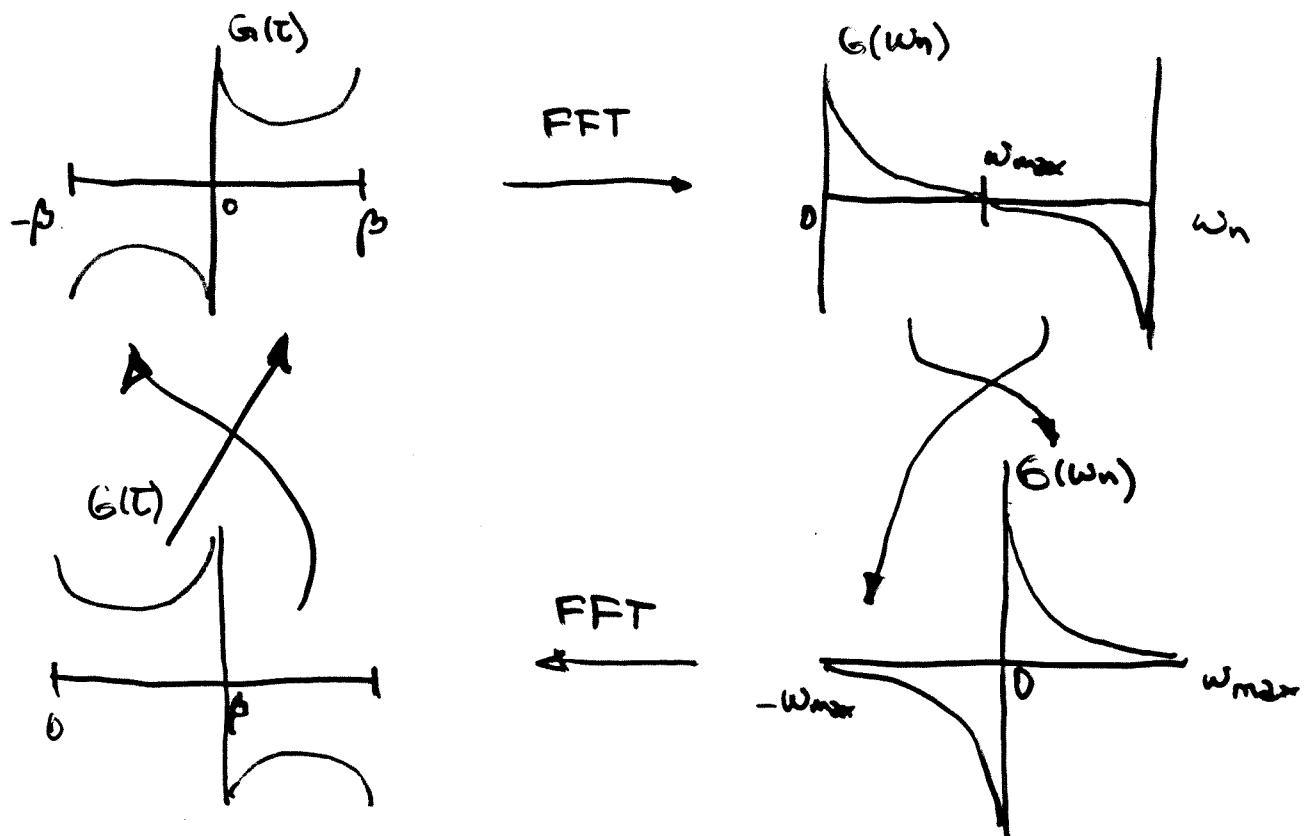
1) Self consistency condition

$$G(\omega_n) = \frac{1}{i\omega_n + \text{sg}(\omega_n) i\Delta}$$

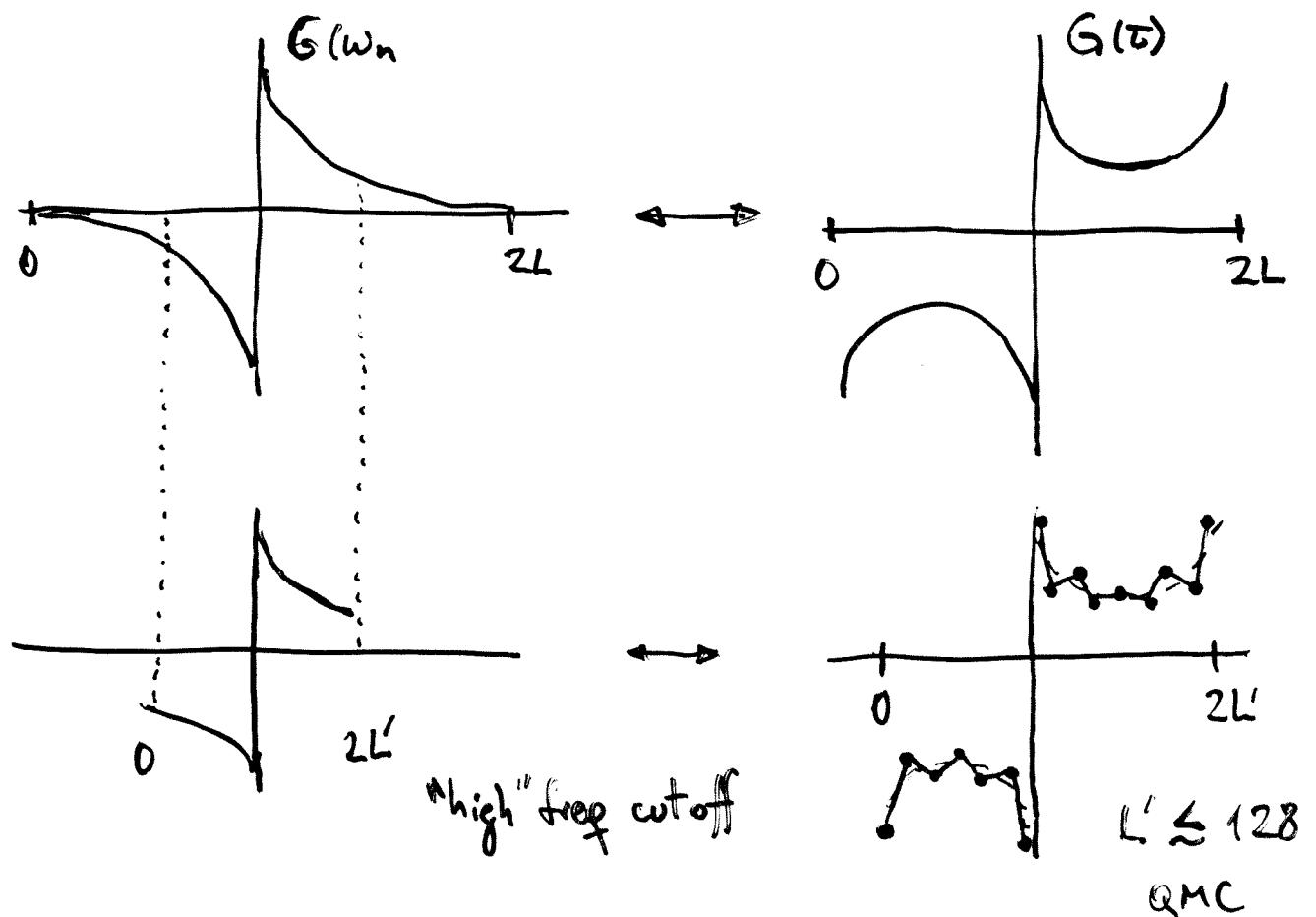
$$\Rightarrow G(\omega_n) = \frac{1}{i\omega_n - \mu - \Sigma(\omega_n) + \text{sg} \sqrt{(i\omega_n - \mu - \Sigma(\omega))^2 - D^2}}$$

$$\text{sg} = \text{sg}(\omega_n - \text{Im}[\Sigma]) \cdot \text{sg}(\text{Im}[\sqrt{\epsilon - \Sigma}])$$

2) FFT \rightarrow cycling



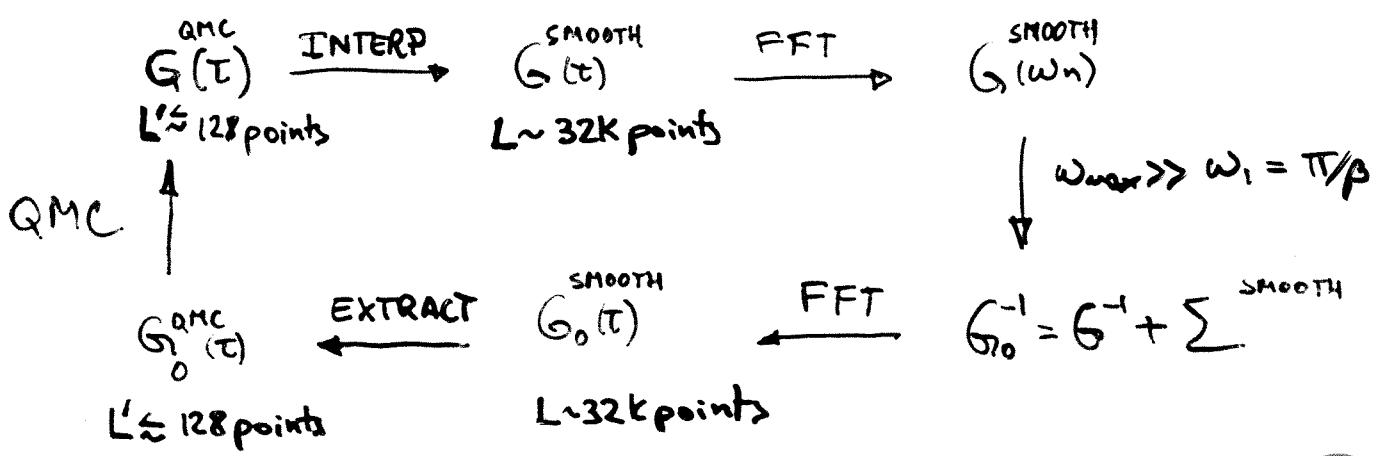
3) FFT → extract - interp



The trick is to interpolate the QMC result $G(t)$

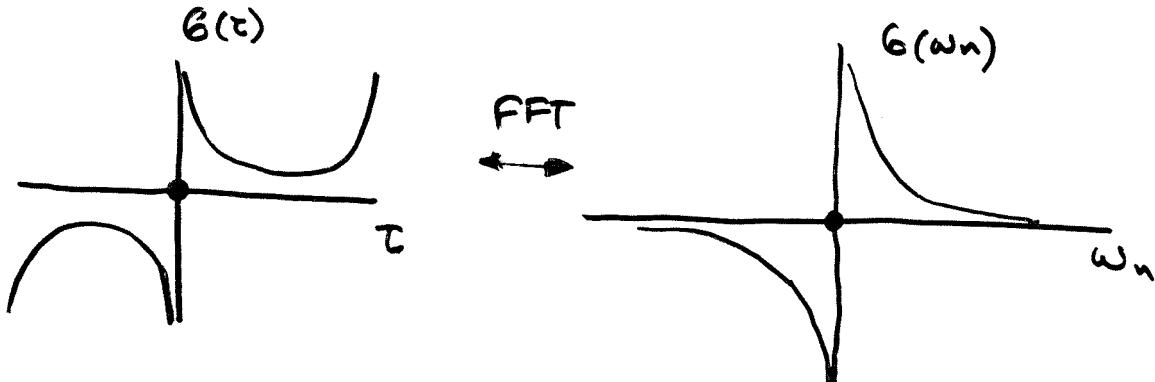
$$G(t) = \frac{1}{\beta} \sum_n e^{i\omega_n t} G(\omega_n) \quad \begin{matrix} \text{is continuous} \\ -\beta \leq t \leq \beta \end{matrix}$$

$$G(\omega_n) = \int_0^\beta e^{-i\omega_n t} G(t) \quad \text{is discrete.}$$



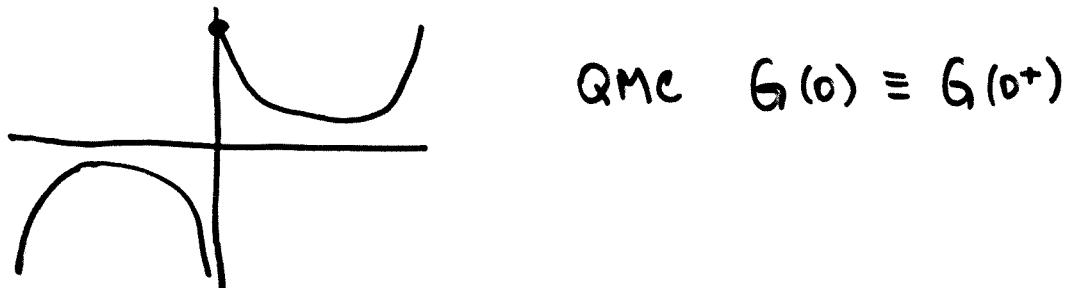
4) FFT \leftrightarrow QMC $G(0^+)$?

p-h symmetry for the sake of the argument.



$$G(0^+) = 1 - n \quad n = \text{part number}$$

$$G(0^+) - G(0^-) = 1$$



Trick $G_0^{\text{QMC}}(0^+) \equiv G_0^{\text{QMC}}(0) = G_0^{\text{smooth}}(0) + \frac{1}{2}$

then

$$G_0^{\text{smooth}}(0) = G_0^{\text{QMC}}(0) - \frac{1}{2}$$

why? FFT for the ω_n or $t=0$ takes the average $0^- \& 0^+$

$$G_0^{\text{smooth}}(0) = \frac{G(0^+) + G(0^-)}{2} = \frac{G(0^+) + G(0^+) - 1}{2} = G(0^+) - \frac{1}{2}$$

* independent of n

Real systems are usually degenerate

d-band $\rightarrow E_g$ - 2 fold
 $\hookrightarrow t_{2g}$ - 3 fold.

\Rightarrow need to implement multi-orbital impurity

Two-band QMC code (2-orbital)

since $t_{vv'} = t_{vv} \delta_{vv'}$

$$G_\sigma(\tau) \rightarrow G_{\sigma\sigma}(\tau)$$

the generalization is easy.

- * Flow diagram is identical
- * Relevant changes in the impurity routine.

1-orbital case

$$e^{-\Delta\tau U n_\uparrow n_\downarrow + \Delta\tau \frac{U}{2} (n_\uparrow + n_\downarrow)} = \frac{1}{2} \sum_{S=\pm} e^{\lambda S (n_\uparrow - n_\downarrow)}$$

2-orbital

$$U n_{1\uparrow} n_{1\downarrow}; U n_{2\uparrow} n_{2\downarrow}; U n_{1\uparrow} n_{2\downarrow}; U n_{2\uparrow} n_{1\downarrow}; U n_{1\downarrow} n_{2\downarrow}; U n_{2\downarrow} n_{1\uparrow}$$

\Rightarrow need to do 6 H-S trans

$$\Rightarrow S_1(\tau); S_2(\tau); S_3(\tau); S_4(\tau); S_5(\tau); S_6(\tau)$$

$$e^{\sum'_{\nu\nu'} - \Delta\tau U (n_{\nu\sigma} - \frac{1}{2}) (n_{\nu'\sigma'} - \frac{1}{2})} = \frac{6U}{4}$$

$$= \left(\frac{1}{2}\right)^6 \sum_{S^1, \dots, S^6} e^{\{\lambda S^1(n_{1\uparrow} - n_{1\downarrow}) + \lambda S^2(n_{2\uparrow} - n_{2\downarrow}) + \dots\}}$$

$$= \left(\frac{1}{2}\right)^6 \sum_{S^1, \dots, S^6} e^{\{\lambda(S^1 + S^3 + S^5)n_{1\uparrow} + \lambda(-S^1 + S^4 + S^6)n_{1\downarrow} + \lambda(S^2 - S^4 - S^5)n_{2\uparrow} + \lambda(-S^2 - S^3 - S^6)n_{2\downarrow}\}}$$

$$G_{1\uparrow}^{-1} = G_{1\uparrow}^{o-1}(\tau, \tau') + \lambda (S^1(\tau) + S^3(\tau) + S^5(\tau)) \delta_{\tau\tau'}$$

$$G_{1\downarrow}^{-1} = G_{1\downarrow}^{o-1}(\tau, \tau') + \lambda (-S^1(\tau) + S^4(\tau) + S^6(\tau)) \delta_{\tau\tau'}$$

$$G_{2\uparrow}^{-1} = \vdots \quad \vdots \quad \vdots$$

$$G_{2\downarrow}^{-1} = \vdots \quad \vdots \quad \vdots$$

$$\Rightarrow xf(4,6) = \begin{pmatrix} 1 & 0 & 1 & 0 & 1 & 0 \\ -1 & 0 & 0 & 1 & 0 & 1 \\ 0 & 1 & 0 & -1 & -1 & 0 \\ 0 & -1 & -1 & 0 & 0 & -1 \end{pmatrix}$$

$$vn(L, 6) = \left(\begin{array}{c|c|c|c|c|c} | & | & | & | & | & | \\ V_1 & V_2 & V_3 & V_4 & V_5 & V_6 \\ | & | & | & | & | & | \end{array} \right)$$

$$V_\tau = \lambda S_\tau$$

$$G = \begin{bmatrix} G_{1\uparrow} & & & \\ & G_{1\downarrow} & & \\ & & G_{2\uparrow} & \\ & & & G_{2\downarrow} \end{bmatrix}$$

$$Z = \sum \det(G_{1\uparrow}^{-1}) \det(G_{1\downarrow}^{-1}) \det(G_{2\uparrow}^{-1}) \det(G_{2\downarrow}^{-1})$$

$S_T^i \rightarrow -S_T^i$ will affect only $2 G^{-1}: v_0 \neq v_0'$

$$R \rightarrow \frac{\det [G_{\sigma}(s) \quad G_{\sigma'}(s)]}{\det [G_{\sigma}(s') \quad G_{\sigma'}(s')]} \quad$$

3 orbitals (t_{2g}) $\binom{6}{2} = 15$ s^i fields

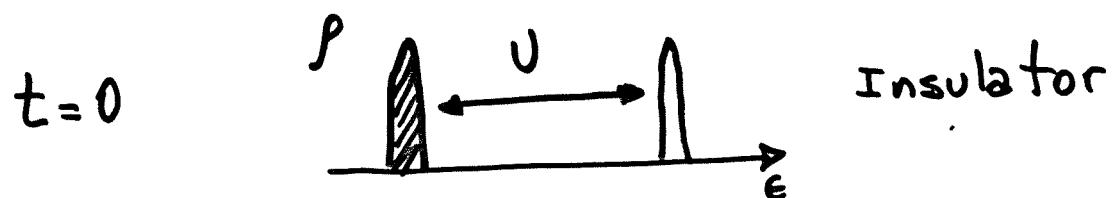
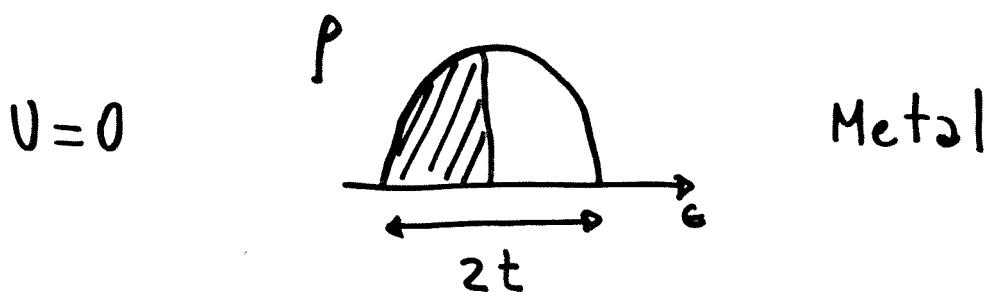
5 orbitals (full d-band) $\binom{10}{2} = 45$ s^i fields

parallelization is easy

Hubbard model has the basic ingredients

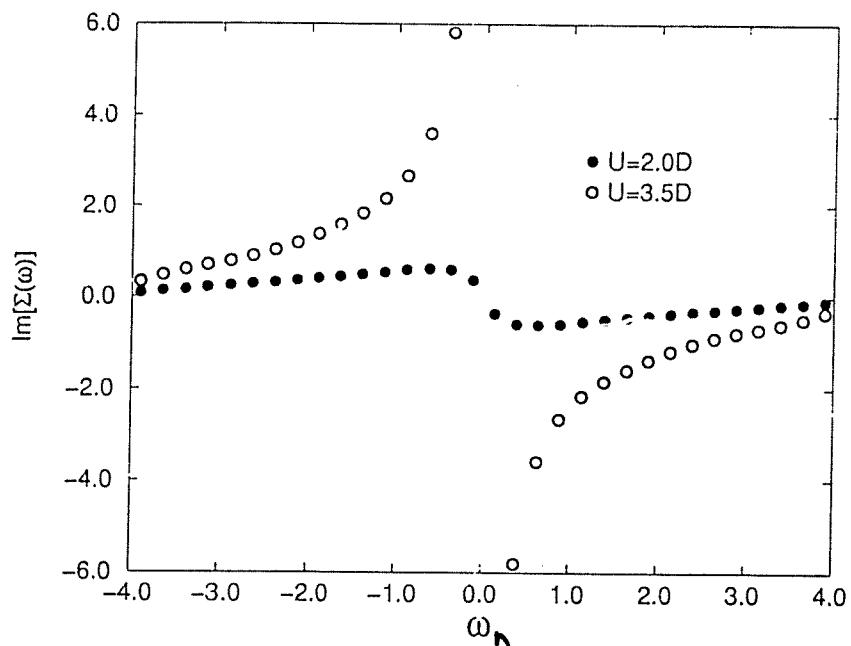
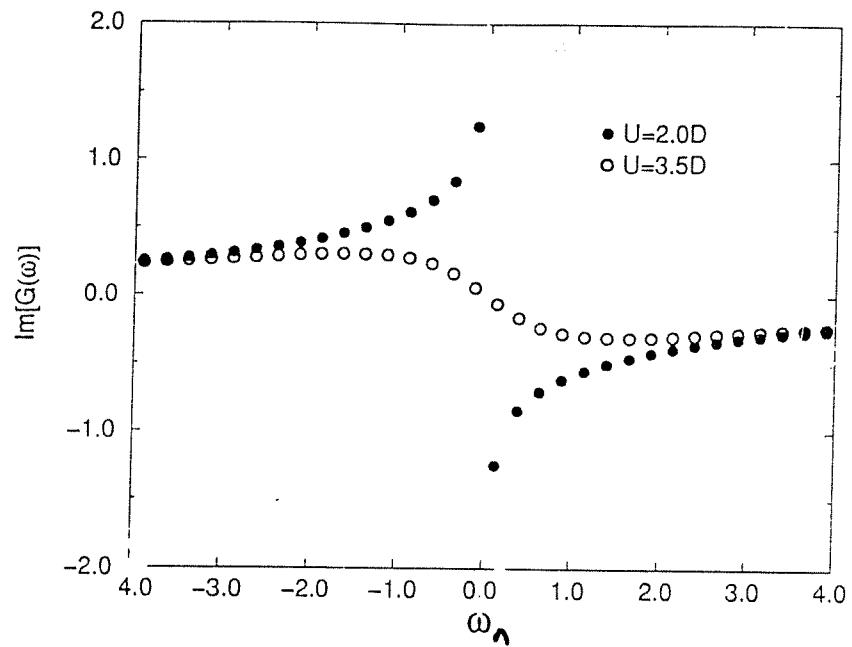
$$H = -t \sum_{\langle i,j \rangle} c_{i\sigma}^+ c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

* PM phase , $1/2$ filling

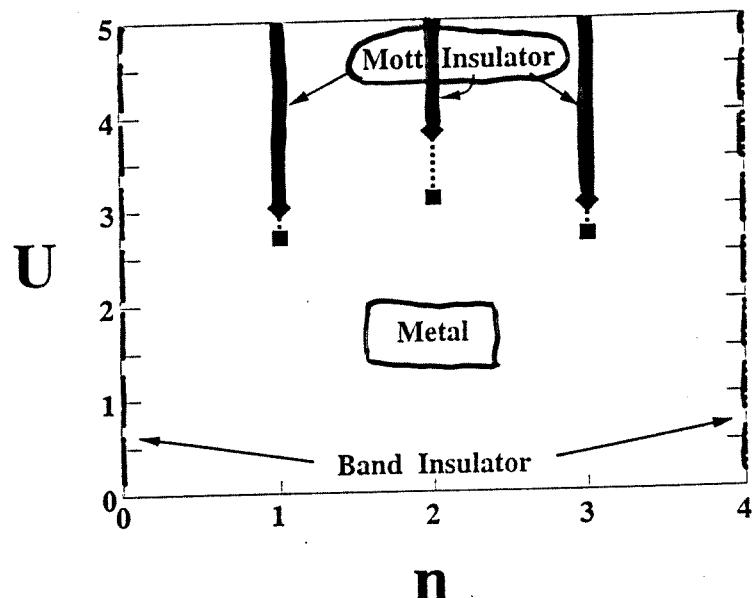


* so far solution exists in $d=1$
(always insulator, $U \neq 0$)

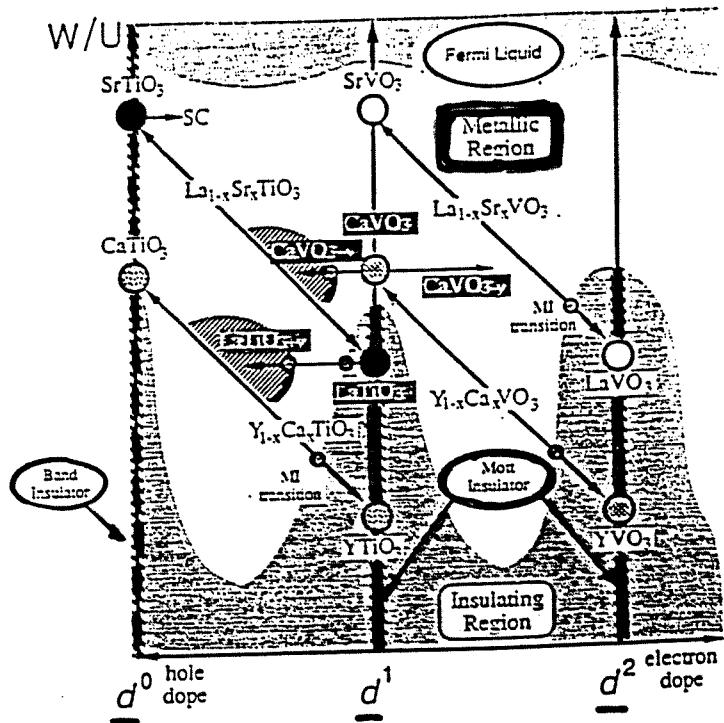
Metal Insulator Transition (QMC)



Phase diagram of 2-band degenerate
Hubbard model (PM)



Schematic Diagram of
Carrier concentration vs. Electron correlation



TI	4+	3+	(2+)
V	5+	4+	3+
Nb	5+	(4+)	3+
Cr	6+		4+

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Nuts & bolts of the QMC - DMFT code

Input file: fort.50

```
0.5d0,0.0d0,1  
1.d0,1,0.1d0,1,0.1d0  
2.0d0  
2.0d0  
2.0d0  
2.0d0  
2.0d0  
2.0d0  
1,1  
1,1  
20000,20000,1000  
  
dtaureal,du,nu  
d,nloop,dmu,nmu,xmu0  
u(1)  
...  
u(6)  
input,imet  
igf,inp  
nsweep,nsweep0,ndirty  
  
input=1 reads seed (fort.20)  
=0 generates new seed  
imet=1 met seed  
=0 ins seed (use only at 1/2 filling)  
igf=1 prints out the gf's  
inp=1 prints out the part number  
nsweep= number of sweeps  
nsweep0= number of sweeps in the last iteration  
ndirty= number of dirty updates
```

Output files

- 'fort.12 G (tau)'
- 'fort.13 Go(tau)'
- 'fort.60 Im G (w)'
- 'fort.61 Re G (w)'
- 'fort.40 Im Go(w)'
- 'fort.41 Re Go(w)'
- 'fort.30 Im Se(w)'
- 'fort.41 Re Se(w)'
- 'fort.80 <docc>'
- 'fort.90 <occ>'
- 'fort.3 seed'